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[Yunhui Ge](#), [David F. Hahn](#), [David L. Mobley](#)





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A Benchmark of Electrostatic Method Performance in Relative Binding Free Energy Calculations

Yunhui Ge,[†] David F. Hahn,[‡] and David L. Mobley^{*,¶}

[†]*Department of Pharmaceutical Sciences, University of California, Irvine, CA 92697, USA*

[‡]*Computational Chemistry, Janssen Research & Development, Turnhoutseweg 30, Beerse
B-2340, Belgium*

[¶]*Departments of Pharmaceutical Sciences and Chemistry, University of California, Irvine,
CA 92697, USA*

E-mail: dmobley@moblelab.org

ABSTRACT

Relative free energy calculations are fast becoming a critical part of early stage pharmaceutical design, making it important to know how to obtain the best performance with these calculations in applications which could span hundreds of calculations and molecules. In this work, we compared two different treatments of long-range electrostatics, Particle Mesh Ewald (PME) and Reaction Field (RF), in relative binding free energy calculations using a non-equilibrium switching protocol. We found simulations using RF achieve comparable results as those using PME but gain more efficiency when using CPU and similar performance using GPU. The results from this work encourage more use of RF in molecular simulations.

INTRODUCTION

The lead optimization stage is an important part of pharmaceutical drug discovery, involving optimization of several key chemical and biophysical properties in order to ensure candidate compounds have adequate selective binding to their target while also having appropriate other properties to potentially become a new pharmaceutical. Lead optimization efforts always require adequate ligand binding affinity for the target, making this a critical design criterion, and one which is a target for predictive methods. Alchemical relative binding free energy (RBFE) calculations based on molecular dynamics (MD) simulations and statistical mechanics have shown promise in providing reliable predictions to guide experimental work in the context of real drug discovery projects.¹⁻³

RBFE calculations compare the potency of two structurally related ligands by transforming one ligand into another via an unphysical or alchemical pathway. This transformation is performed in both the protein-ligand complex and in the solution state to form a closed thermodynamic cycle. The RBFE of the ligands simulated can be calculated through two opposite legs in this cycle.^{4,5}

To calculate free energies, these alchemical transformations can be performed via either an equilibrium or non-equilibrium protocol. In general, a number of intermediate simulations are conducted along an alchemical path between the two physical end states, with these simulations being either equilibrium or nonequilibrium depending on the approach chosen. While the free energy difference of interest depends only on the physical end states in the limit of adequate sampling, these intermediate states serve to help obtain converged results and provide sampling which is hopefully adequate. Equilibrium free energy calculations run an equilibrium simulation at each intermediate state as well as the end states which are physically meaningful. In contrast, the non-equilibrium protocol simulates the end states at equilibrium, potentially spending considerably more time there, and only runs short simulations to switch between end states. However, running a large number of switching trials is critical in this case. It is still under debate which of the two protocols is more efficient,

with different studies drawing different conclusions,⁶⁻⁹ and the choice of protocol is beyond the scope of this paper. This work follows the protocol deployed in a previous work¹⁰ in which a non-equilibrium approach was used, though our results may generalize to equilibrium approaches as well.

Generally, modern MD engines (e.g., AMBER, GROMACS, CHARMM, etc) support different approaches for simulations, including RBEF calculations. Among these software packages, GROMACS is an open source package which is widely used for molecular simulations and reproduces calculated free energies within about 0.2 kcal/mol.¹¹ Moreover, tools (e.g., pmx^{12,13}) allow easy high-throughput applications of GROMACS RBEF simulations, providing a workflow spanning from initial coordinate files to final free energy estimates. This leads to relatively user-friendly RBEF calculations using GROMACS.

Long-range electrostatics interactions are critical in modeling molecular motions in simulations. Due to the computational complexity of such long-range interactions, it becomes challenging to design accurate and efficient methods to describe such interactions. An excellent review of methods for computing the long-range electrostatics interactions in biomolecular simulations can be found here.¹⁴

Among a number of existing methods, Particle Mesh Ewald (PME)^{15,16} is perhaps the most broadly used. PME and other Particle-Mesh methods (e.g., particle-particle particle-mesh method (P3M)^{17,18}) are based on the Ewald approach which is a classic method to exactly calculate the electrostatic potential¹⁹ and is chosen as a starting point for further adjustments for better efficiency. In particle-mesh (PM) methods (e.g., PME, P3M) the interaction potential is split into a short-range term (direct sum) and a long-range term (reciprocal sum). The direct part converges fast using a fixed cutoff, and the efficiency of the method is determined by the reciprocal part. The original PME and the smoothed version, smooth PME (which is the implementation of PME method in GROMACS), assign the charges to a grid using Lagrangian interpolation and cardinal B-spline interpolation, respectively. The reciprocal energy and forces can be obtained efficiently using fast Fourier

transforms (FFT). Historically, the PME methods were inspired by the P3M method and later it is shown that they can be mathematically transformed into each other.²⁰ The algorithms details along with the similarities and differences of the two methods are reviewed by previous work.^{17,21–25} It is also somewhat controversial in the literature as to whether P3M and PME should be considered completely equivalent (and share the same name, e.g. P3M); here, however, we do not enter this debate and simply refer to the approach as PME, as this is how the implementation we use is described in the GROMACS documentation.

The PME method finds a balance between the accuracy and efficiency and is widely used in molecular simulations. Additionally, a number of other approaches to long-range electrostatics use similar adaptations of Ewald-based techniques.^{17,18,26,27}

While widely used in explicit solvent simulations, Ewald based approaches including PME may also introduce periodicity-induced artifacts which overstabilize the native structure of inherently non-periodic systems. Such artifacts have been extensively explored in previous work and the conclusions are sometimes contradictory suggesting such artifacts are dependent on the system and simulation conditions.^{28–37}

Another popular option to compute electrostatics interactions is Reaction Field (RF).^{38–40} The RF method only computes the interactions up to a cutoff distance and implicitly treats any interactions beyond the cutoff in a mean-field manner using an appropriate dielectric constant. In this way, electrostatics interactions can be calculated efficiently while hopefully maintaining enough accuracy. However for relatively inhomogeneous systems (e.g., membranes), using RF may introduce artifacts and lead to incorrect modeling of the system. For example, lattice sum electrostatics can be used to represent effectively infinite membranes when simulating membrane proteins, but RF requires a continuum treatment of long-range interactions. Still, RF sees considerable use in the field and can be appealing for suitable system geometries.^{41–47}

In this work, we compare the performance of PME and RF methods in RBF calculations. There is extensive literature comparing RF and Particle-Mesh methods (e.g., PME, P3M) in

plain MD simulations on different systems (e.g., liquids, proteins, etc)^{37,40,48-74} which is well summarized in the work of Reif and Oostenbrink.⁶⁵ Our focus in this work is specifically on whether these methods yield equivalent relative binding free energies, given the tremendous industrial interest in such calculations. While both RF and PME methods are widely used in free energy calculations,^{41-47,75-92} they have not yet been compared head-to-head to see if they yield equivalent results in RBFE calculations, to our knowledge.

METHODS

Selected targets. We selected the targets TYK2⁹³ and CDK2⁹⁴ which are part of several RBFE benchmark studies^{2,95,96} including the set commonly referred to as Schrödinger’s ”JACS set”⁹⁵ from a key paper in JACS on large-scale free energy calculations. Among all 8 target systems in the JACS set, TYK2 has a moderate system size (\sim 60000 atoms including water and ions) which yields representative results in performance. CDK2 is the largest system (\sim 110000 atoms including water and ions) among these targets and we selected it to verify the trends observed in TYK2 simulations. For TYK2, we used all 24 edges in our calculations and randomly selected 6 edges of CDK2. The ligands simulated in this work are neutral so no net charge change is involved in each perturbation. A table of successfully simulated perturbations can be found in Table S1, S2.

Molecular Dynamics Simulations. The simulations were performed using GROMACS (2021-dev-20200320-89f1227-unknown) with a patch optimizing PME performance on GPU (<https://gerrit.gromacs.org/c/gromacs/+/13382>).

For each perturbation, two sets of simulations were prepared: solvated ligands and ligand-protein complexes. The initial ligand and protein structures were obtained from a previous published work.¹⁰

The ligand was parameterized using Open Force Field version 1.0.0 (codenamed ”Parsley”).⁹⁷ Hybrid structures and topologies for the ligand pairs were generated using pmx^{12,13}

following a single topology approach. The workflow established a mapping between atoms of two ligands based on the maximum common substructure and conformational alignment while minimizing perturbation and stabilizing the system.

The AMBER ff99sb*ILDN force field⁹⁸⁻¹⁰⁰ was used for protein parameterization. Dodecahedral boxes were filled with TIP3P explicit solvent model solvated ligand pairs/ligand-protein complexes and counterions (150 mM NaCl).

For each perturbation, two states were prepared for both in-solution/bound state ligands: state A and state B, representing ligand 1 and ligand 2, respectively. An energy minimization was first performed, followed by a 10 ps NVT equilibration at 298K. Then the production equilibrium simulation (in the NPT ensemble) was performed for 6 ns at 298 K and a pressure of 1 bar. 80 snapshots were extracted from the production simulation. For each snapshot, a non-equilibrium transformation from state A to B (and vice versa) was performed during 50 ps. For each perturbation, 3 replicas of the series of simulations described above were performed leading to a total of 60 ns simulation data to calculate the free energy differences for the ligands in their in-solution/bound states.

The stochastic dynamics thermostat was used to control the temperature in the simulations. The Parrinello-Rahman barostat¹⁰¹ was applied to keep the pressure constant. All bond lengths were constrained using the LINCS algorithm.¹⁰² The van der Waals interactions were smoothly switched off between 1.0 and 1.1 nm. A dispersion correction for energy and pressure was used. The non-bonded interactions for the alchemical transitions were treated with a modified soft-core potential.¹⁰³

Two different methods implemented in GROMACS were used to treat long-range electrostatic interactions: Particle Mesh Ewald (PME) and Reaction Field (RF). For simulations using PME, a direct space cutoff of 1.1 nm and a Fourier grid spacing of 0.12 nm were used. For simulations using RF, a dielectric constant (ϵ_{rf}) of 78.3, the dielectric constant of water at 298 K was used.¹⁰⁴

Both CPU and GPU simulations were performed using PME or RF electrostatics. In the

rest of the paper, we denote CPU-PME, CPU-RF, GPU-PME and GPU-RF to represent the hardware and methods for long-range electrostatic interactions treatment.

RESULTS

Predicted relative binding free energies from simulations using RF show good agreement with those from PME. The relative binding free energies ($\Delta\Delta G$) were calculated for a set of modifications of TYK2 (24 $\Delta\Delta G$ values in total) which is also included in a standard test set for relative free energy calculations (the set commonly referred to as Schrödinger’s ”JACS set”⁹⁵ from a key paper in JACS on large-scale free energy calculations). Figure 1 summarizes the computed values using RF/PME on CPU/GPU. The uncertainty estimates were performed by 1000 bootstrapping trials and are reported in Figure 1 as $x_{x_{low}}^{x_{high}}$ where x is the mean value, x_{high} and x_{low} indicate 95% confidence intervals. The averaged root-mean-square error (RMSE) and mean unsigned error (MUE) of $\Delta\Delta G$ for CPU-PME versus CPU-RF are 0.45 and 0.34 kcal/mol and for GPU-PME versus GPU-RF are 0.52, 0.40 kcal/mol, respectively (Figure 1a-b). A cross-platform (CPU vs GPU) comparison also shows essentially the same level of agreement between results using RF and PME (Figure 1c-d) suggesting the same trend holds true on both CPU and GPU.

In this work, we are focused on the comparison between RF and PME methods on relative binding free energy calculations. While comparing calculated $\Delta\Delta G$ with experimental measured values is more important to benchmarking force fields/methodologies, it is not the top interest in this work. However, we summarize our calculated values and the corresponding experimental results in Table S3. In conclusion, the calculated values of both methods are similarly accurate when compared to the experimental values.

Simulations using RF are more efficient than using PME. Given that PME and RF achieve a good level of agreement of calculated $\Delta\Delta G$ values, our focus shifts to computational efficiency, where RF is generally less computationally demanding than using PME^{79,105} when

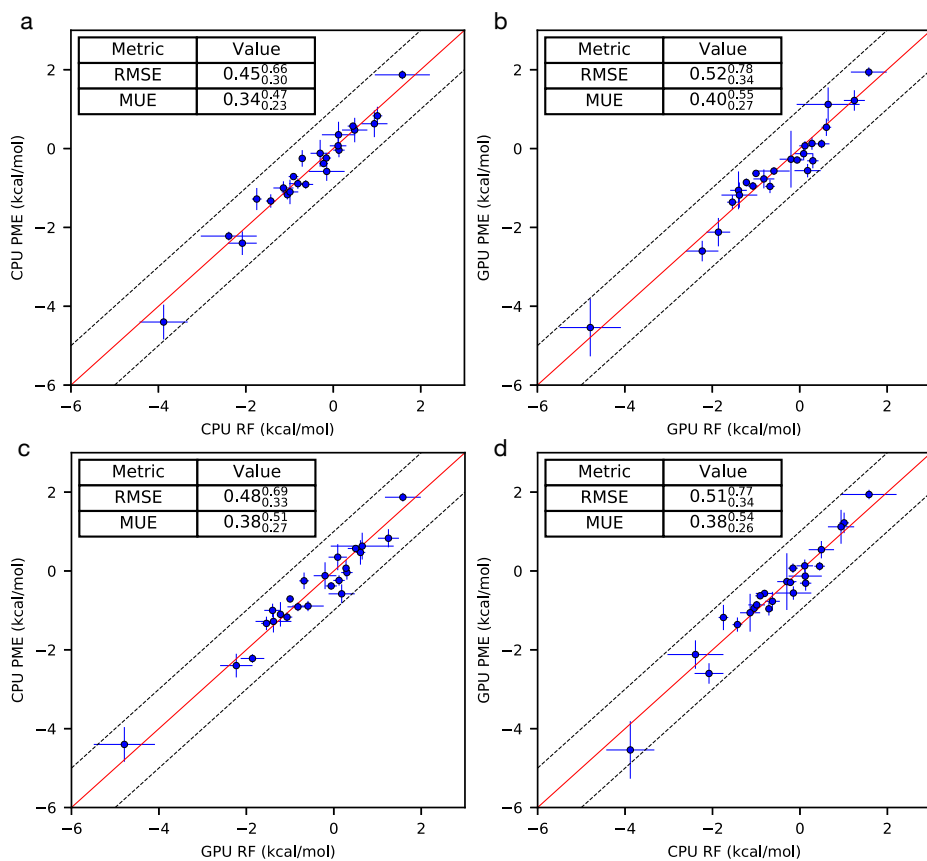


Figure 1: Calculated $\Delta\Delta G$ values of TYK2 ligands using RF/PME on CPU/GPU. Overall, a good agreement is achieved between using PME and RF on CPU/GPU. The uncertainty estimates are calculated by bootstrapping, using 100 bootstrap samples.

the same cut-off is used since more computational time is needed for PME in reciprocal space and tabulated interactions. The simulation performance (in the unit of ns/day) was analyzed from the 6-ns equilibrium NPT simulation and is summarized in Figure 2 where different colors represent using RF/PME on CPU/GPU. The uncertainties were estimated using the standard deviations across different edges. As mentioned in Methods, the simulations were performed in-solution and ligand-bound state in which 5908 ± 221 atoms and 62290 ± 6 atoms were simulated, respectively. For the (less costly) in-solution ligand simulations, using RF on CPU is $\sim 30\%$ faster than using PME on average (Figure 2a) and is similar to PME when simulated on GPU considering the uncertainties (Figure 2b). For the more costly bound state simulation, a similar trend is observed in CPU simulations (Figure 2c).

However, using RF on GPU is $\sim 10\%$ slower than using PME on average (Figure 2d).

The results from Figure 2 show that using RF is faster than PME in most cases for the different system sizes and hardware tested here. It is notable that the GROMACS version used in these simulations was specifically optimized for PME performance on GPU (<https://gerrit.gromacs.org/c/gromacs/+/13382>). Thus, it is not surprising that PME outperformed RF on GPUs (slightly) in our tests here. However, the difference between PME and RF on GPU performance is only minor ($\sim 10\%$). Possibly a similar optimization of GROMACS for RF on GPUs could yield substantial performance gains.

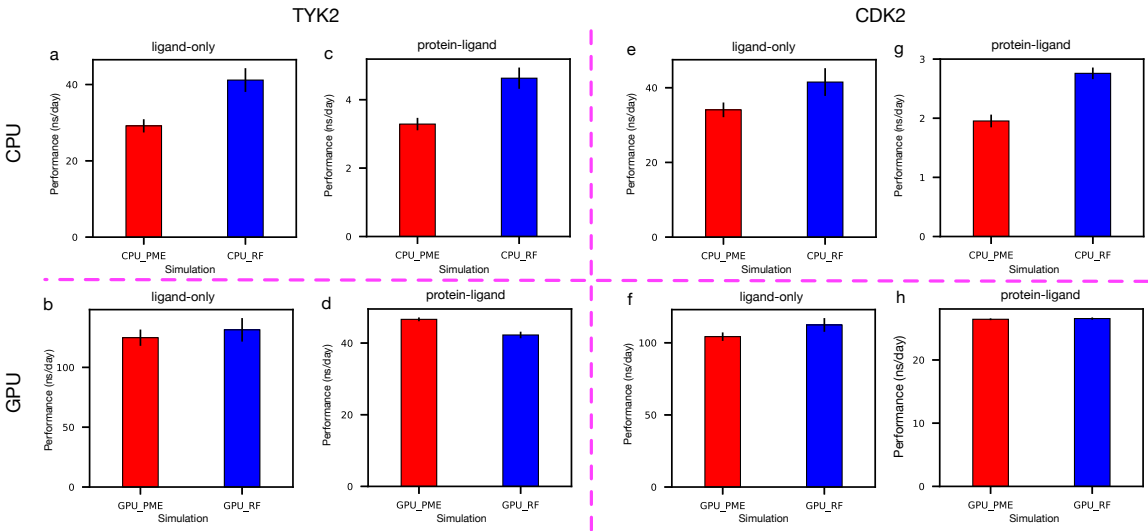


Figure 2: Performance difference between using RF/PME on CPU/GPU in TYK2 simulations of (a-b) ligand-only in-solution, (c-d) protein-ligand complex and in CDK2 simulations of (e-f) ligand-only in-solution, (g-h) protein-ligand complex.

To verify the observed trends in TYK2 simulations, 6 selected edges from CDK2 were simulated using the same protocol described in METHOD. CDK2 is a larger system than TYK2 and has 5602 ± 373 and 106908 ± 5 atoms in in-solution and ligand-bound state simulation, respectively. Similar to the results of TYK2, simulations using RF are also faster ($\sim 20\text{-}40\%$) than using PME on CPU (Figure 2e,g). Notably, GPU simulations using RF achieves a similar efficiency to PME (Figure 2f,h) which means even with optimizations, PME still cannot surpass RF in efficiency (Figure 2d).

Figure 2 shows that using RF is more efficient than PME on CPUs and is comparable to PME in GPU simulations. This is remarkable given the fact the GROMACS used in this work is optimized for PME performance on GPU. RF could perhaps become even more promising with further optimizations.

On average, the protein-ligand complex simulation of each TYK2 perturbation took 31.7 hours using RF and 47.3 hours using PME on 8 CPUs ($\sim 30\%$ performance gain using RF), with the protocol used here. With same setup, the ligand-only simulation of TYK2 took 3.5 and 4.8 hours using RF and PME, respectively (also $\sim 30\%$ performance gain using RF). When using 1 GPU, a calculation time of 3.4 hours using RF and 3.1 hours using PME was needed for protein-ligand complex simulations ($\sim 10\%$ performance gain using PME). In ligand-only simulation on 1 GPU, an average time of 1.1 and 1.2 hours was needed using RF and PME, respectively ($\sim 10\%$ performance gain using RF). A summary of simulation wallclock time can be found in Table S4 and S5.

Here, while we find that PME and RF achieve comparable results within uncertainty, this is with present-day force fields, and may not always be the case in future studies. It is not too difficult to imagine a scenario where force fields might be optimized for the best performance with a particular model of long range electrostatics, such as a force field designed for use with PME, or one for use with RF. Once force fields are tuned for specific electrostatics treatments, it is reasonable to expect that results might differ in quality if the electrostatics treatment is changed. To some extent, we already see such tuning taking place for water models (e.g. the SPC/E, TIP4P-Ew, and TIP5P-E models designed for use with Ewald-based electrostatics¹⁰⁶⁻¹¹⁰), though not yet for other components of the force field to our knowledge. Still, results from these water studies indicate that electrostatics treatment, or at least the switch from cutoff based to lattice-sum electrostatics, impacts bulk properties enough that the choice of electrostatics treatment should perhaps be considered part of the force field. So, perhaps future force fields should require a particular choice of long range electrostatics treatment, and separate Ewald-based and RF-based versions should be fitted.

CONCLUSION

The treatment of long-range electrostatic interactions is critical for a correct modeling of (bio)molecular systems in molecular dynamics simulations. Due to the long-range nature and N^2 -scaling of electrostatic interactions, they are computationally the most demanding terms in the force field evaluation. Both PME and RF are popular methods of treating long range electrostatic interactions. PME is widely used and attempts to find a good balance between efficiency and accuracy, but we find that RF achieves higher efficiency, although it may have some limitations in relatively inhomogeneous systems such as for membrane simulations.

We expect the results observed in this work will be generalizable to other simulation packages and force fields, in that results will not be highly dependent on PME vs RF, but (a) that is outside of scope for this work, and (b) details will depend on the exact implementation in those software packages. Previous work has shown that different simulation engines are able to take in the same force field and same systems and give identical energies (with PME) within a reasonable tolerance, which would strongly suggest the implementations are equivalent or nearly so,^{111,112} at least for PME.

The conclusions presented in this work are true only for neutral perturbations and are not guaranteed in other perturbations (e.g., charge-changing mutations). When charged species are involved, different correction schemes must be applied to correct the bias induced by electrostatics finite-size effects depending on electrostatic treatment methods used (e.g., RF, PME).¹¹³⁻¹¹⁸

This work is focused on comparing the results and performance of PME and RF for relative binding free energy calculations. We found the calculated $\Delta\Delta G$ values are in good agreement using PME and RF whereas simulations using RF are comparable or faster than PME on both CPU and GPU. Taken the results presented here, we suggest that RF may be a promising option for relative free energy simulations because, at least in GROMACS, it is less computationally demanding while retaining comparable accuracy to PME. This advantage

may be particularly helpful in cases where a large number of simulations are needed (e.g., in the lead optimization stage of the drug discovery process). Thus we recommend free energy calculations with RF be considered as a viable option, at least for homogeneous systems.

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Supporting Information Available

Supporting information (simulated system size, calculated/experimental relative binding free energies, simulation wallclock time) is available free of charge via the Internet at <http://pubs.acs.org>.

Data and Software Availability

Datasets. All input files for simulations, data and scripts for analysis are freely available at <https://github.com/MobleyLab/PME-RF-benchmark>.

Software. Simulations were performed using the open-source molecular dynamics package GROMACS with a specific patch optimizing PME performance on GPUs. Analysis was performed using pmx and PLBenchmarks. The versions of both packages used in this work are freely distributed at <https://github.com/MobleyLab/PME-RF-benchmark/tree/main/SI/analysis> with installation instructions.

Notes

D.L.M. is a member of the Scientific Advisory Board of OpenEye Scientific Software and an Open Science Fellow with Silicon Therapeutics.

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Supplementary Information

Supporting Tables

Table S1: Simulated perturbations of TYK2 and the number of particles.

Ligand1	Ligand 2	No. atoms (ligand-only)	No. atoms (protein-ligand)
ejm-44	ejm-55	5954	62290
ejm-49	ejm-31	6521	62288
ejm-31	ejm-46	5888	62281
jmc-28	jmc-27	5886	62282
ejm-42	ejm-48	6014	62290
ejm-31	ejm-43	5950	62286
ejm-50	ejm-42	5630	62272
ejm-42	ejm-55	5633	62275
jmc-23	ejm-46	5884	62277
ejm-31	ejm-45	5638	62283
ejm-55	ejm-54	5729	62278
ejm-45	ejm-42	5639	62284
ejm-31	ejm-48	5960	62284
ejm-47	ejm-31	5906	62290
ejm-47	ejm-55	5915	62290
ejm-44	ejm-42	5956	62295
jmc-23	jmc-27	5886	62276
ejm-43	ejm-55	5951	62281
jmc-23	jmc-30	5884	62277
jmc-28	jmc-30	5887	62286
ejm-42	ejm-54	5640	62282
ejm-49	ejm-50	6517	62290
jmc-23	ejm-55	5887	62289
ejm-31	jmc-28	5889	62291
jmc-23	jmc-27	5886	62276

Table S2: Simulated perturbations of CDK2 and the number of particles.

Ligand1	Ligand 2	No. atoms (ligand-only)	No. atoms (protein-ligand)
22	1h1r	4978	106892
loi u	26	5885	106905
26	1h1q	5762	106902
17	1h1q	4987	106904
loi y	loi9	5675	106908
17	21	4984	106895

Table S3: Calculated $\Delta\Delta G$ values (kcal/mol) and the experimental measured values of simulated perturbations of TYK2. The uncertainty estimates were performed by 1000 bootstrapping trials and are reported as $x_{x_{low}^{high}}$ where x is the mean value, x_{high} and x_{low} indicate 95% confidence intervals.

Ligand1	Ligand 2	$\Delta\Delta G$ (CPU- PME)	$\Delta\Delta G$ (GPU- PME)	$\Delta\Delta G$ (CPU-RF)	$\Delta\Delta G$ (GPU-RF)	$\Delta\Delta G$ (EXP)
ejm-44	ejm-55	-4.40±0.44	-4.54±0.73	-3.88±0.55	-4.79±0.70	-1.79
ejm-49	ejm-31	-0.89±0.13	-0.57±0.09	-0.81±0.20	-0.59±0.36	-1.79
ejm-31	ejm-46	-0.25±0.21	-0.96±0.17	-0.71±0.07	-0.68±0.11	-1.77
jmc-28	jmc-27	-0.71±0.06	-0.63±0.06	-0.91±0.11	-1.00±0.08	-0.30
ejm-42	ejm-48	0.83±0.23	1.22±0.26	1.01±0.11	1.25±0.24	0.78
ejm-31	ejm-43	1.87±0.11	1.94±0.12	1.58±0.63	1.58±0.41	1.28
ejm-50	ejm-42	-0.24±0.13	0.07±0.12	-0.16±0.11	0.12±0.13	-0.80
ejm-42	ejm-55	-1.17±0.12	-0.95±0.10	-1.05±0.07	-1.07±0.12	0.57
jmc-23	ejm-46	-0.04±0.18	-0.31±0.19	0.13±0.13	0.30±0.13	0.39
ejm-55	ejm-54	-1.10±0.31	-0.86±0.10	-0.99±0.19	-1.22±0.07	-1.32
ejm-45	ejm-42	0.57±0.07	0.12±0.11	0.45±0.12	0.50±0.18	-0.22
ejm-31	ejm-48	0.63±0.34	1.12±0.43	0.94±0.30	0.65±0.72	0.54
ejm-47	ejm-31	-0.12±0.34	-0.27±0.72	-0.30±0.22	-0.20±0.26	0.16
ejm-47	ejm-55	-0.91±0.11	-0.77±0.24	-0.63±0.17	-0.82±0.24	0.49
ejm-44	ejm-42	-2.22±0.11	-2.12±0.36	-2.39±0.64	-1.86±0.27	-2.36
jmc-23	jmc-27	0.07±0.21	0.13±0.16	0.11±0.19	0.28±0.06	0.42
ejm-43	ejm-55	-2.40±0.30	-2.60±0.26	-2.08±0.33	-2.23±0.37	-0.95
jmc-23	jmc-30	-0.38±0.09	-0.29±0.06	-0.22±0.10	-0.06±0.13	0.76
jmc-28	jmc-30	-1.00±0.17	-1.06±0.48	-1.14±0.23	-1.40±0.19	0.04
ejm-42	ejm-54	-1.33±0.17	-1.36±0.18	-1.43±0.11	-1.54±0.12	-0.75
ejm-49	ejm-50	-1.28±0.28	-1.18±0.32	-1.75±0.11	-1.38±0.41	-1.23
jmc-23	ejm-55	0.35±0.33	-0.13±0.17	0.12±0.38	0.09±0.21	2.49
ejm-31	ejm-45	0.47±0.31	0.54±0.22	0.49±0.29	0.61±0.11	-0.02
ejm-31	jmc-28	-0.58±0.24	-0.56±0.17	-0.15±0.41	0.18±0.30	-1.44
RMSE		1.09 ^{1.43} _{0.77}	1.16 ^{1.56} _{0.82}	1.04 ^{1.35} _{0.76}	1.20 ^{1.61} _{0.83}	
MUE		0.86 ^{1.14} _{0.61}	0.93 ^{1.27} _{0.66}	0.86 ^{1.10} _{0.64}	0.94 ^{1.27} _{0.66}	

Table S4: Simulation wallclock time (min) of TYK2.

Simulation	complex	ligand-only
CPU-PME	2838.2±328.1	288.1±27.6
GPU-PME	185.4±4.12	69.0±3.8
CPU-RF	1903.7±334.2	211.6±20.5
GPU-RF	204.7±5.9	66.0±4.8

Table S5: Simulation wallclock time (min) of CDK2.

Simulation	complex	ligand-only
CPU-PME	4447.3±423.7	252.8±24.7
GPU-PME	326.8±6.5	83.1±3.8
CPU-RF	3143.0±204.5	210.2±21.9
GPU-RF	325.9±11.9	77.2±4.3