# A Comparative Study of Carrier-Frequency Modulation Techniques for Conducted EMI Suppression in PWM Converters

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Abstract—A rigorous mathematical analysis and a comparative study of carrier-frequency modulation (CFM) techniques for the conducted electromagnetic interference (EMI) suppression in pulsewidth-modulated converters is presented. CFM techniques dither the switching period with a small amplitude variation around the nominal value, so that the harmonic power is redistributed over the spectrum of concern. Two types of dithering signals, including the periodic and random signals, are investigated in this paper. The operational characteristics as well as the input and output power spectra of the converters with the two modulating signals are compared. In particular, their characteristics in the low- and high-frequency harmonic power redistribution will be depicted. It is shown that random CFM (RCFM) gives a more effective way to disperse the harmonics around the switching frequency than the periodic CFM (PCFM) with the same frequency deviation. However, RCFM introduces higher low-frequency harmonics than the PCFM at the converter output. Furthermore, effects of the resolution filter bandwidth in the electromagnetic compatibility analyzer on conducted EMI measurement is discussed. The validity of the analyses is confirmed experimentally by using a dc/dc buck converter operating in continuous conduction mode.

Index Terms—DC-DC power conversion, power electronics, pulsewidth modulation, random switching techniques, switching circuits.

# I. INTRODUCTION

A S THE electromagnetic spectrum becomes crowded, many countries have imposed electromagnetic compatibility (EMC) regulations that must be met before electronic products can be sold legally. The most direct way of solving electromagnetic interference (EMI) problems is to reduce the emission from the noise source [1].

EMI is a serious problem in power electronics circuits because of their fast switching characteristics. Extensive research has been conducted on developing various soft-switching techniques for different applications to reduce both switching

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loss and radiated EMI [2]. By smoothing the switching transients and reducing the ringing problems, radiated EMI emission can be lessened. Although those techniques are particularly effective for reducing high-frequency emissions, they are unsuitable for suppressing switching-frequency-related emissions (i.e., conducted EMI). Apart from using passive noise filters and shields in the circuit, another approach of EMI suppression is based on modulating a parameter, such as the switching frequency, pulsewidth, or pulse position, of the pulsewidth-modulated (PWM) signals that apply to the switching devices either in a periodic [3] or random manner [4]–[8]. The parameter is dithered with small amplitudes around the nominal value, so that the harmonic power is redistributed over the spectrum of concern. Among all possible parameters, modulating the frequency, namely, the carrier-frequency modulation (CFM), is a very effective solution because the switching frequency is not fixed at a constant value and, thus, the spectral component at the original switching frequency will be diminished in the input and output spectra.

As reported in [3], periodic CFM (PCFM) can reduce the conducted EMI with a suitable selection of the modulating frequency. This method introduces sidebands at the original discrete frequency components in the standard PWM scheme with much smaller magnitude. So far, effects of the resolution filter bandwidth in the EMC analyzer on the discrepancies between theoretical and measured spectra have not been addressed. Moreover, as discussed in [3], Carson's rule is inapplicable for studying high-frequency behaviors, due to the overlap of the sideband harmonics of different switching frequency harmonics. In this paper, rigorous mathematical analysis of the spectra will be given.

For the random switching schemes, random CFM (RCFM) with fixed duty cycle is the best option because: 1) all discrete harmonics can be significantly reduced; 2) the harmonic power can be spread over as a continuous noise spectrum of small magnitude; and 3) the low-frequency harmonic injection is minimal [6].

This paper presents a comparative study of using PCFM and RCFM on the conducted EMI suppression in PWM converters. Issues addressed include the operational properties and the input and output power spectral densities of the converters with the two modulation techniques. In particular, their characteristics in low- and high-frequency harmonic power redistribution are depicted. The effects of the resolution filter bandwidth in the EMC analyzer on the conducted EMI measurement will be emphasized. The validity of the analyses is verified experimentally with a dc/dc buck converter.

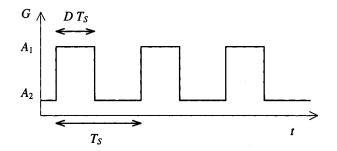


Fig. 1. Typical PWM function.

#### II. PCFM AND RCFM

Fig. 1 shows a typical PWM function G(t) with the duty cycle D and the switching period  $T_s$ . G(t) has two discrete levels (namely,  $A_1$  and  $A_2$ ), which are applicable to describe the behaviors of switching converters. Mathematically, G(t) can be expressed by a Fourier series [9]

$$G(t) = \sum_{n = -\infty}^{\infty} C_n e^{j\theta_n} \tag{1}$$

where  $C_n$  and  $\theta_n$  are the magnitude and phase of the nth harmonic, respectively.

 $C_n$  can be calculated by

$$C_{n} = \frac{1}{T_{s}} \int_{-T_{s}/2}^{T_{s}/2} G(t) e^{-j2\pi n f t} dt$$

$$= \frac{j}{2\pi n f T_{s}} \left[ A_{1} \left( e^{-j2\pi n f D T} - 1 \right) + A_{2} \left( e^{-j2\pi n f T} - e^{-j2\pi n f D T} \right) \right]. \quad (2)$$

### A. Theoretical Power Spectrum With PCFM

With the PCFM, the switching frequency of G(t) is modulated by a cosinusoidal function of amplitude  $\Delta f_s$  and frequency  $f_m$ . In practice,  $\Delta f_s$  represents the maximum frequency deviation in the switching frequency and  $f_m$  is the modulating frequency. The instantaneous frequency  $f_{\text{inst},n}$  of the nth harmonic is equal to the time derivative of  $\theta_n$ . That is,

$$f_{\text{inst},n} = \frac{1}{2\pi} \cdot \frac{d\theta_n}{dt}$$

$$= n[f_s + \Delta f_s \cos(2\pi f_m t)]$$
(3)

where  $f_s = 1/T_s$  is the nominal switching frequency. By using (3),

$$\theta_n = 2\pi n \int_0^t [f_s + \Delta f_s \cos(2\pi f_m \lambda)] d\lambda$$
$$= 2\pi n f_s t + n\beta \sin(2\pi f_m t) \tag{4}$$

where  $\beta$  is the ratio of the peak instantaneous frequency deviation to the frequency of the modulating signal, namely, the modulation index in [3] and [11]. It is defined as

$$\beta = \frac{\Delta f_s}{f_m}. (5)$$

Substituting (4) into (1) gives

$$G(t) = \sum_{n=-\infty}^{\infty} C_n e^{jn2\pi f_s t} \left\{ \cos[n\beta \sin(2\pi f_m t)] + j \sin[n\beta \sin(2\pi f_m t)] \right\}.$$
 (6)

By using the Jacobi equations in [10], the cosine and sine terms in (6) can be expressed as

$$\cos[n\beta \sin(2\pi f_m t)] = J_0(n\beta) + \sum_{k=\text{even}}^{\infty} 2J_k(n\beta)\cos(2\pi k f_m t)$$

$$\sin[n\beta \sin(2\pi f_m t)] = \sum_{k=\text{odd}}^{\infty} 2J_k(n\beta)\sin(2\pi k f_m t)$$
 (7)

where  $J_k(\cdot)$  is the kth order Bessel function. It can be expanded by the following series:

$$J_{k}(n\beta) = \left(\frac{n\beta}{2}\right)^{k} \left[\frac{1}{k!} - \frac{\left(\frac{n\beta}{2}\right)^{2}}{1!(k+1)!} + \frac{\left(\frac{n\beta}{2}\right)^{4}}{2!(k+2)!} - \frac{\left(\frac{n\beta}{2}\right)^{6}}{3!(k+1)!} + \cdots\right]. \quad (8)$$

Numerical values of the Bessel functions with respect to  $n\beta$  up to the tenth order are tabulated in Table I. The following two properties are also noted:

$$J_k(n\beta) \approx 0, \qquad k > n\beta + 1,$$
 (9a)

and

$$\sum_{k=-\infty}^{\infty} J_k^2(n\beta) = 1 \qquad \forall \, n\beta. \tag{9b}$$

Substituting (7) into (6) gives

$$G(t) = \sum_{n=-\infty}^{\infty} C_n e^{j2\pi n f_s t}$$

$$\times \left[ J_0(n\beta) + \sum_{k=1}^{\infty} J_k(n\beta) e^{j2\pi k f_m t} \right.$$

$$+ \sum_{k=\text{even}}^{\infty} J_k(n\beta) e^{-j2\pi k f_m t}$$

$$- \sum_{k=\text{odd}}^{\infty} J_k(n\beta) e^{-j2\pi k f_m t} \right]$$

$$= \sum_{n=-\infty}^{\infty} C_n \left\{ J_0(n\beta) e^{j2\pi n f_s t} \right.$$

$$+ \sum_{k=1}^{\infty} J_k(n\beta) \left[ e^{j2\pi (n f_s + k f_m) t} \right.$$

$$+ (-1)^k e^{j2\pi (n f_s - k f_m) t} \right] \right\}. \tag{10}$$

nβ	$J_0$	$J_1$	$J_2$	$J_3$	$J_4$	$\overline{J_5}$	$J_6$	$J_7$	$J_8$	$J_9$
0	1.00	-	_	_	-	-	-	-	-	-
0.5	0.94	0.24	0.03	-	-	-	-	-	-	-
1	0.77	0.44	0.11	0.02	~	<del>-</del>	-	-	-	-
2	0.22	0.58	0.35	0.12	0.03	-	-	-	-	-
3	-0.26	0.34	0.49	0.31	0.13	0.04	0.01	-	-	-
4	-0.40	-0.07	0.36	0.43	0.28	0.13	0.05	0.02	-	-
5	-0.18	-0.33	0.05	0.36	0.39	0.26	0.13	0.05	0.02	-
6	0.15	-0.28	-0.24	0.11	0.36	0.36	0.25	0.13	0.06	0.02

TABLE I Numerical Values of the First Ten Order  $J_k$  With Respect to  $n\beta$ 

The magnitude of G(t) at frequency f,  $G(f,\beta)$ , can be shown to be

$$G(f,\beta) = \sum_{n=-\infty}^{\infty} C_n \left\{ J_0(n\beta)\delta(f - nf_s) + \sum_{k=1}^{\infty} J_k(n\beta) \left[ \delta(f - nf_s - kf_m) + (-1)^k \delta(f - nf_s + kf_m) \right] \right\}.$$
(11)

Its power spectrum  $S_G(f,\beta)$  in the positive frequency range is equal to

$$S_{G}(f,\beta) = 2 \left| \sum_{n=1}^{\infty} C_{n} \left\{ J_{0}(n\beta)\delta(f - nf_{s}) + \sum_{k=1}^{\infty} J_{k}(n\beta) \left[ \delta(f - nf_{s} - kf_{m}) + (-1)^{k} \delta(f - nf_{s} + kf_{m}) \right] \right\} \right|^{2},$$

$$f > 0. \tag{12}$$

Thus,  $S_G(f,\beta)$  contains infinite discrete harmonics and is dependent on  $\beta$  and  $f_m$ .

For the standard PWM scheme (i.e.,  $\beta = 0$ ),  $J_0(0)$  exists only. Referring to Table I, (12) becomes

$$S_G(f,\beta) = 2\sum_{n=1}^{\infty} |C_n|^2 \delta(f - nf_s).$$
 (13)

As expected, the power spectrum consists of the discrete harmonics at the multiples of  $f_s$ .

For the PCFM scheme (i.e.,  $\beta \neq 0$ ),  $J_k(n\beta)$  leads to upper and lower sidebands at the multiples of  $f_s$ . Although the number of side frequencies is infinite, the bandwidth  $B_n$  of the sideband of the nth harmonic can be approximated by (9a) as

$$B_n = 2(n\beta + 1)f_m \tag{14}$$

where  $B_n \in [nf_s - (n\beta + 1)f_m, nf_s + (n\beta + 1)f_m].$ 

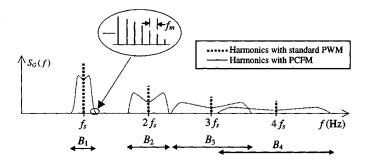


Fig. 2. Spectral structure of  $S_G$  with PCFM.

If  $n\beta$  is sufficiently large,  $B_n$  can further be approximated by

$$B_n \approx 2n\beta f_m = 2n\Delta f_s, \qquad \beta \gg 1.$$
 (15)

Compared with Carson's rule [3], (12) gives a more general formula that describes the spectrum of the PWM waveform. It is possible to study the spectrum at high-order switching harmonics and large  $\Delta f_s$ , in which case the side frequencies of different switching harmonics may overlap. As will be discussed in Section II-C, it is possible to predict accurately the measured power spectrum using (12).

Fig. 2 depicts the spectral structure of  $S_G$ . PCFM generates the upper and lower sidebands with the bandwidth  $B_n$  at the nth multiples of  $f_s$ . Each sideband consists of side frequencies, which are  $f_m$  apart and have magnitudes much smaller than the value at the center frequency in the standard PWM scheme. The number of side frequencies in  $B_n$  is equal to the order of  $J_k$ . Mathematically, using a larger value of  $\beta$  results in more side frequencies. As boldfaced in Table I, the maximum absolute values of  $J_k$ ,  $|J_k(n\beta)|_{\text{max}}$ , under each  $n\beta$  determine the peak magnitude in the considered sideband. They give a relative measure on the harmonic reduction as compared to the standard PWM scheme. Fig. 3 shows the relationships of  $|J_k(n\beta)|_{\text{max}}$ against  $n\beta$ . Basically,  $|J_k(n\beta)|_{\text{max}}$  decreases as  $n\beta$  increases. However, its rate of decrement decreases with the increase in  $n\beta$ . With the above phenomena, the following properties can be concluded.

- P1) Harmonic suppression effect increases with the increase in  $\beta$
- P2) Suppression of high-order harmonics is more effective than the low-order ones.
- P3) Insignificant enhancements in harmonic suppression are observed when  $n\beta$  is sufficiently large.

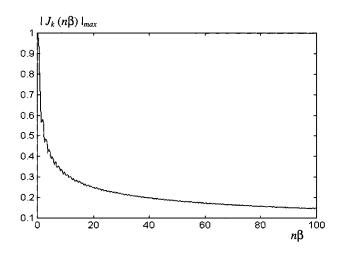


Fig. 3. Relationships of  $|J_k(n\beta)|_{\text{max}}$  against  $n\beta$ .

TABLE II
COMPONENT VALUES AND SPECIFICATIONS OF THE BUCK CONVERTER

Inductor	940μΗ	Output voltage	5V
Capacitor	4.7μF	Nominal switching frequency $f_s$	100kHz
Output Load	10Ω	Duty-cycle D	0.4

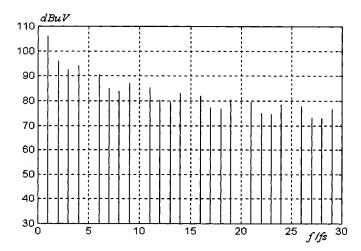


Fig. 4. Theoretical input current power spectrum with  $\beta = 0$ .

The above properties are illustrated with an example of a buck converter. Its specifications are tabulated in Table II. The values of  $A_1$  and  $A_2$  in Fig. 1 are 0.48 and 0 A, respectively, in the following analysis. Fig. 4 shows the theoretical input current power spectrum with  $\beta$  equals zero (i.e., standard PWM scheme). It contains discrete harmonics at the multiples of the switching frequency only. Figs. 5 and 6 show the power spectra with  $\beta=2$ ,  $\beta=8$ , and  $\beta=16$ , respectively, under two testing conditions.

The first testing condition is that  $f_m$  is fixed at 3 kHz with  $\Delta f_s = 6$  kHz,  $\Delta f_s = 24$  kHz, and  $\Delta f_s = 48$  kHz. The results in Fig. 5 reveal that harmonic suppression is more effective with a larger  $\beta$  (i.e,  $\beta = 8$  and  $\beta = 16$ ). This confirms P1). Comparing Figs. 4 and 5(b) shows a reduction of 10 dB in

the low-frequency range and 20 dB in the high-frequency range. This confirms P2).

Fig. 6 shows the results in the second test when  $\Delta f_s$  equals 12 kHz with  $f_m=6$  kHz,  $f_m=1.5$  kHz, and  $f_m=750$  Hz, respectively. The results are similar to those in Fig. 5. They indicate that harmonic suppression is highly dependent on the value of  $\beta$  and is relatively independent on the combinations of  $f_m$  and  $\Delta f_s$ . Conversely, it will be shown later that the measured spectra are dependent on  $f_m$  and  $\Delta f_s$ , because of the interaction between the resolution filter in the EMC analyzer and  $S_G$ .

It can be observed that the results for  $\beta=16$  in Figs. 5(c) and 6(c) are similar to those of Figs. 5(b) and 6(b), respectively, especially at the high-frequency range. This demonstrates P3), that the spectra will be similar when  $n\beta$  is sufficiently large.

# B. Theoretical Power Spectral Density (PSD) With RCFM

For the RCFM,  $T_s$  in Fig. 1 is dithered in random manner and D is kept constant [6]. Instead of using power spectrum, PSD is considered in the analysis [5]. The PSD  $S_G(f,\Re)$  under RCFM has been shown to be

$$S_G(f, \Re) = \frac{1}{E[T_k]} \left\{ E[|G(f)|]^2 + 2 \operatorname{Re} \left\{ \frac{E[G(f)e^{j2\pi f T_k}]E[G^*(f)]}{1 - E[e^{j2\pi f T_k}]} \right\} \right\}$$
(16)

where  $G^*(f)$  is the complex conjugate of G(f) and  $T_k$  is the instantaneous switching period. The probability density function of  $T_k$ ,  $P(T_k)$ , is of uniform distribution with the expected value of  $T_s$ . It is defined as

$$P(T_k) = \frac{1}{T_2 - T_1} = \frac{1}{\Re T_s} \tag{17}$$

where  $T_1$  and  $T_2$  are the minimum and maximum possible switching periods, respectively, and  $\Re$  is the randomness level.

This scheme spreads the discrete harmonics over the frequency spectrum. As  $\Re$  increases, the harmonics are gradually spread over. However, no significant improvements in spreading high-frequency harmonics are observed when  $\Re$  is larger than a certain value, such as 0.15 in [6]. Moreover, low-frequency harmonic components increase as  $\Re$  increases.

## C. Discrepancies Between Theoretical and Measured Spectra

The bandwidth  $BW_{\rm res}$  of the resolution filter in the EMC analyzer significantly affects the measured power spectrum.  $BW_{\rm res}$  might introduce discrepancies between the theoretical and measured spectra. In the following discussions, the HP 8590EM series EMC analyzer is used for illustration.

As shown in [12], attenuation of the resolution filter is in Gaussian characteristics. The approximate ratio between the bandwidth at -60 dB (denoted by  $BW_{-60}$ ) and  $BW_{\rm res}$  is 15. That is,

$$\frac{BW_{-60}}{BW_{\rm res}} \approx 15. \tag{18}$$

In order to predict the displayed spectrum on the EMC analyzer, a triangle-like filter response shown in Fig. 7 is proposed to

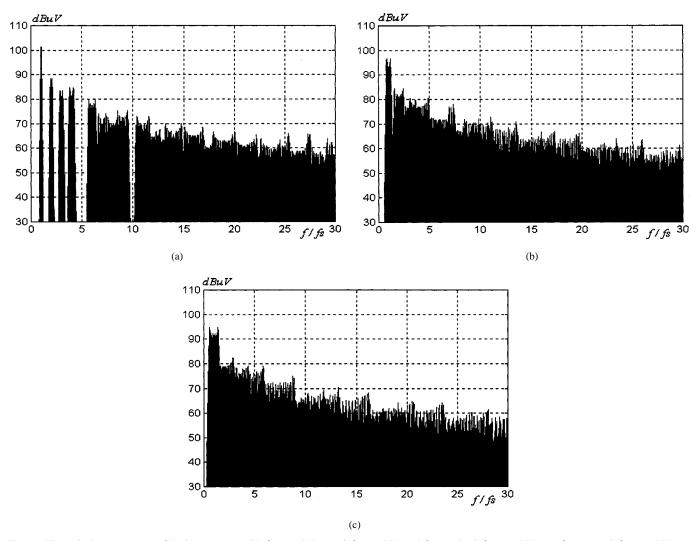


Fig. 5. Theoretical power spectra of the input current with  $f_m=3$  kHz. (a)  $\Delta f_s=6$  kHz and  $\beta=2$ . (b)  $\Delta f_s=24$  kHz and  $\beta=8$ . (c)  $\Delta f_s=48$  kHz and  $\beta=16$ .

model the filter attenuation. The filter response function W(f) is defined as follows:

$$W(f) = \begin{cases} W_R(f), & \text{for } f \ge f_0 \\ W_L(f), & \text{for } f \le f_0 \end{cases}$$
 (19)

where  $f_0$  is the center frequency over  $BW_{\rm res},\,W_R(f)=10^{-(0.4/BW_{\rm res})(f-f_0)}=-(4/BW_{\rm res})(f-f_0)$  dB, and  $W_L(f)=10^{(0.4/BW_{\rm res})(f-f_0)}=-(4/BW_{\rm res})(f-f_0)$  dB. Hence, the measured power at  $f_0$  by the analyzer is equal to

$$S_{G,\text{meas}}(f_0, \beta) = \int_{f_0 - 7.5BW_{\text{res}}}^{f_0 + 7.5BW_{\text{res}}} S_G(f, \beta) |W(f)|^2 df$$

$$= S_G(f_0, \beta)$$

$$+ \int_{f_0 + \Delta}^{f_0 + 7.5BW_{\text{res}}} S_G(f, \beta) W_R(f)^2 df$$

$$+ \int_{f_0 - \Delta}^{f_0 - 7.5BW_{\text{res}}} S_G(f, \beta) W_L(f)^2 df$$
(20)

By using the property in (9b), the total power of the nth order harmonic in the fixed-frequency PWM scheme is equal to the sum of the discrete powers over  $B_n$  in PCFM. The number of side frequencies  $N_{sf}$  that falls within  $BW_{\rm res}$  can be approximated by

$$N_{sf} = \frac{15BW_{\text{res}}}{f_m} = \frac{15BW_{\text{res}}}{\Delta f_s} \beta. \tag{21}$$

The measured spectrum will be similar to the theoretical calculation if  $N_{sf}$  is small. This can be achieved by either reducing  $BW_{\mathrm{res}}$  or increasing  $f_m$  and  $\Delta f_s$ . According to some international EMC regulations, such as the Comite International Special des Perturbations Radioelectriques (CISPR),  $BW_{\mathrm{res}}$  is fixed in the measurement. Thus, the most viable approach is to adjust  $f_m$  and  $\Delta f_s$ . However, if  $f_m$  is too large,  $\beta$  will be small and the magnitude of the side frequencies will increase. If  $\Delta f_s$  is too large, it will be discussed later that low-frequency harmonics will be introduced. Table III shows the interactions between  $S_G$  and the resolution filter under two testing conditions. First  $f_m$  is fixed.  $BW_{\mathrm{res}}$  will enclose more harmonic power in  $S_G$  with

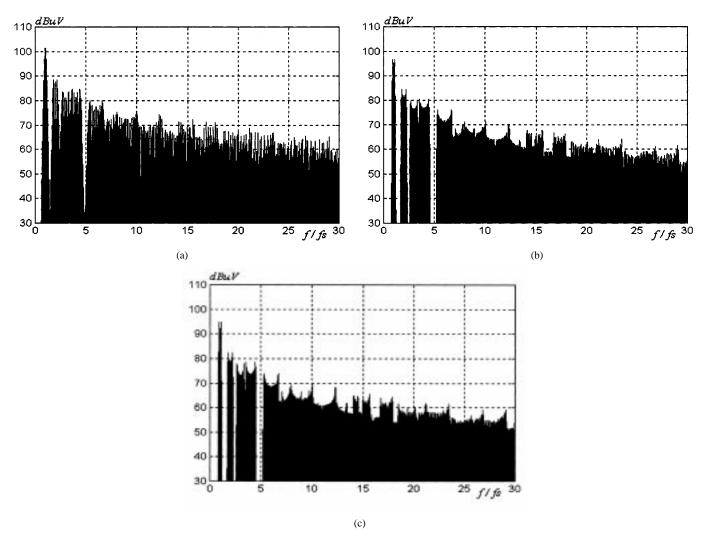


Fig. 6. Theoretical power spectra of the input current with  $\Delta f_s=12$  kHz. (a)  $f_m=6$  kHz and  $\beta=2$ . (b)  $f_m=1.5$  kHz and  $\beta=8$ . (c)  $f_m=750$  Hz and  $\beta=16$ .

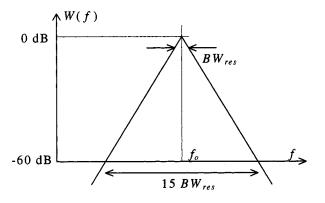


Fig. 7. Model of the triangle-like filter response in the EMC analyzer.

smaller  $\beta$  (i.e., small  $\Delta f_s$ ). Second  $\Delta f_s$  is fixed. The measured  $S_G$  is insensitive to  $\beta$ . The measured spectrum is unable to differentiate the effectiveness of  $\beta$  in the actual EMI suppression.

In order to demonstrate the above observations, the theoretical power spectra shown in Figs. 5 and 6 are used to compare the measured power spectra, which are calculated by (20). With  $BW_{\rm res}$  equals 9 kHz, the measured power spectra for Figs. 5

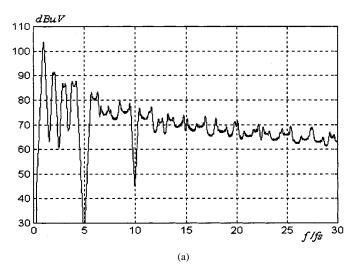
and 6 are shown in Figs. 8 and 9, respectively. As expected, the measured results are higher than the corresponding envelopes in Figs. 5 and 6. Comparing Fig. 8(a) with Fig. 8(b), the results vary with  $\Delta f_s$  (and, hence,  $\beta$ ). Increasing  $\Delta f_s$  can effectively disperse the low-order harmonics. However, with the same  $\Delta f_s$  and different  $\beta$  in Fig. 9, the predicted results in the two cases are similar, confirming the argument of the discrepancies.

Finally,  $\Delta f_s$  is not a prime factor in high-frequency harmonic suppression in all cases. Since  $B_n$  increases with the harmonic order, the side frequencies in  $B_n$  will overlap inevitably with the lower band  $B_{n-1}$ . The overlapped regions increase the PSD over  $BW_{\rm res}$  and become relatively insensitive to  $\Delta f_s$ . Conversely, since the overlapping effect is small at the low-order harmonics, the measured spectra vary with  $\Delta f_s$  (Fig. 8). Carson's rule can be applied to explain the EMI suppression with the increase in  $\Delta f_s$  [3].

In the RCFM, as frequency components randomly appear, the measured spectrum is dependent on the spectral density inside  $BW_{\rm res}$ . If  $\Re$  is sufficiently large, the PSD of the measured PWM waveform is uniformly distributed. The measured spectrum will increase as  $BW_{\rm res}$  increases.

PCFM at Low  $\beta$  High  $\beta$ Fixed  $f_m$  And Different  $\Delta f_s$ Different  $f_m$  and Fixed  $\Delta f_s$ 

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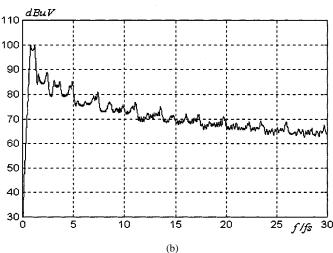


Fig. 8. Comparison of the predicted power spectra with  $BW_{\rm rcs}=9$  kHz and  $f_m=3$  kHz. (a)  $\Delta f_s=6$  kHz and  $\beta=2$ . (b)  $\Delta f_s=24$  kHz and  $\beta=8$ .

## D. Low-Frequency Noise

A dc/dc converter can be considered as a low-pass filter fed by different types of input sources [6]. Let H(f) be the filter

transfer characteristic for buck and transimpedance for boost and buck-boost converters. The power spectrum of the noise output  $S_{n_o}(f,\beta)$  is

$$S_{n_0}(f,\beta) = S_G(f,\beta)|H(f)|^2, \qquad f \neq 0$$
 (22)

where  $S_G$  is the diode voltage for buck converters and inductor current for boost and buck–boost converters.

Since  $|H(f)|^2$  rolls off at the high-frequency range, the rms value of noise ripple  $v_{n_o}$  can be approximated by

$$v_{n_o}(\beta) = \left[ \int_0^{2f_s} S_G(f,\beta) |H(f)|^2 df \right]^{1/2}$$

$$\cong \left[ \sum_{n=1}^2 \sum_{k=0}^{n\beta+1} |C_n J_k(n\beta) H(nf_s \pm kf_m)|^2 \right]^{1/2}$$
(23)

for the PCFM.

As discussed above, the lower sideband of the first harmonic is subject to  $\Delta f_s$ . Using large  $\Delta f_s$  will introduce low-frequency noise at the converter output. In order to keep a tight output regulation, the lowest harmonic spectra in PCFM has to be far beyond the cutoff frequency  $f_{\rm cutoff}$  of H(f). Hence,

$$\Delta f_s \ll f_s - f_{\text{cutoff}}.$$
 (24)

With the RCFM, it has been shown in [6] that the PSD of the output noise is

$$v_{n_o}(\Re) = \left[ \frac{4f_s}{N} \sum_{k=1}^{N} S_G\left(\frac{2kf_s}{N}, \Re\right) \left| H\left(\frac{2kf_s}{N}\right) \right|^2 \right]^{1/2}$$
 (25)

where N is the number of frequency points within  $[0, 2f_s]$ .

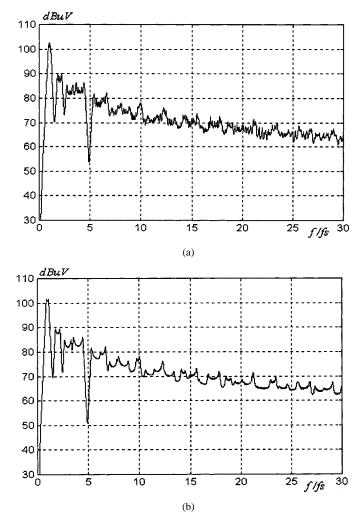


Fig. 9. Comparison of the predicted power spectra with  $BW_{\rm res}=9$  kHz and  $\Delta f_s=12$  kHz. (a)  $f_m=6$  kHz and  $\beta=2$ . (b)  $f_m=1.5$  kHz and  $\beta=8$ .

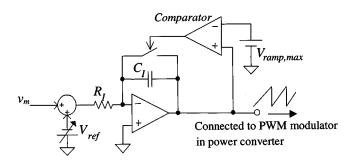
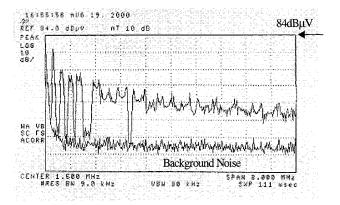


Fig. 10. Implementation of the ramp VCO.

Low-frequency noise will inherently introduce when  $\Re$  is nonzero. Selection of  $\Re$  is based on the harmonic spreading effects and the low-frequency noise injection.

### III. EXPERIMENTAL VERIFICATIONS AND COMPARISONS

The same converter, as tabulated in Table II, is investigated. CFM is achieved by using the ramp voltage-controlled oscillator (VCO) shown in Fig. 10. By integrating the signal that combines a modulating signal  $v_m$  (which is either periodic or random) and



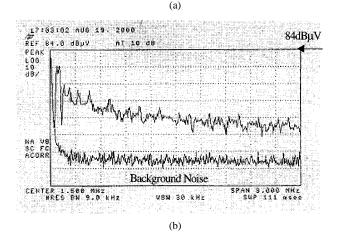


Fig. 11. Measured power spectrum of the input current with PCFM (10 dB/div). (a)  $f_m=3$  kHz,  $\Delta f_s=6$  kHz, and  $\beta=2$ . (b)  $f_m=3$  kHz,  $\Delta f_s=24$  kHz, and  $\beta=8$ .

a reference dc voltage  $V_{\rm ref}$ , a triangular waveform is generated. The peak value  $V_{\rm ramp,max}$  of the output is defined as

$$V_{\text{ramp,max}} = \frac{1}{R_I C_I} \int_0^{T_s} (-V_{\text{ref}} + v_m) dt$$
 (26)

where  $T_s$  is the generated switching period. The modulated switching frequency equals

$$f_s = G_{VCO}(-V_{ref} + v_m) = f_s + \Delta f_s \tag{27}$$

where  $G_{\rm VCO}=1/(R_IC_IV_{\rm ramp,max})$  is the gain of the VCO,  $f_s=-G_{\rm VCO}V_{\rm ref}$ , and  $\Delta f_s=G_{\rm VCO}\,v_m$ .

Fig. 11 shows the measured power spectrum of the input current under PCFM with the same condition as in Fig. 8. The EMC analyzer is an HP 8591EM [13] and  $BW_{\rm res}$  equals 9 kHz. The conversion factor of the current probe is 0.05 V/A (which introduces a 26 dB offset in the overall spectrum). As the cutoff frequency of the output filter is 332 Hz, the value of  $\Delta f_s$  satisfies (24). It can be seen that the measured results agree closely with the predictions. In the RCFM, the measured spectrum under the same frequency deviation is examined. The equivalent random level  $\Re$  equals (1/76-1/124)/(1/100)=0.5. The theoretical and measured results are shown in Fig. 12.

Comparing Fig. 11 with Fig. 12, both PCFM and RCFM give similar profiles in the measured spectra, especially in the high-frequency range. However, for the low-frequency harmonics,

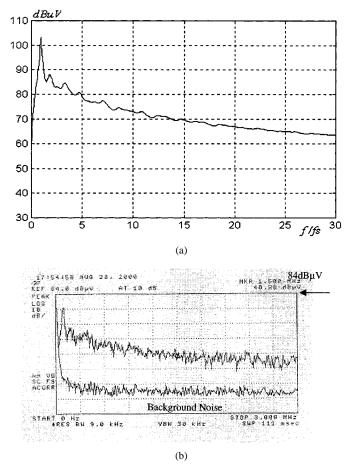


Fig. 12. Power spectra of the input current with RCFM and  $BW_{\rm rcs}=9$  kHz. (a) Theoretical prediction. (b) Measured result.

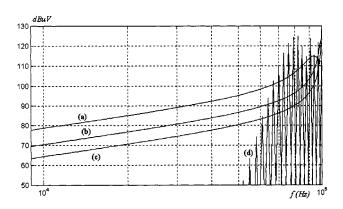
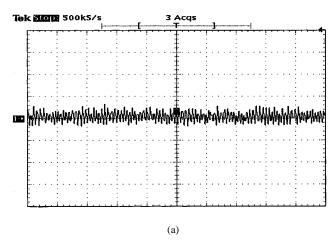


Fig. 13. Comparison of the low-frequency characteristics of the diode voltage with  $BW_{\rm res}=200$  Hz. (a): RCFM,  $\Re=0.5$ . (b): RCFM,  $\Re=0.2$ . (c): RCFM,  $\Re=0.1$ . (d): PCFM.

the performance of the two schemes are different. The theoretical spectra of the diode voltage with the two schemes are shown in Fig. 13. Basically, RCFM gives higher low-frequency harmonics than the PCFM. No harmonics will be generated beyond the sideband of the first-order-switching harmonic in PCFM [Fig. 13(d)]. Fig. 14 shows the output voltage of the converter under the two schemes. It can be observed that the RCFM introduces higher low-frequency variation than the PCFM, thus confirming the theoretical prediction. Although low-frequency harmonics in the RCFM can be reduced with a smaller  $\Re$  [Fig. 13(b)



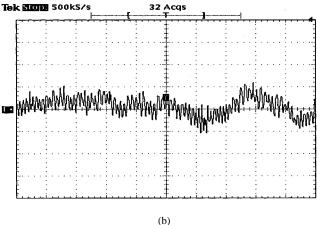


Fig. 14. Measured output ripple (50 mV/div) (timebase: 100  $\mu$ s/div). (a) PCFM. (b) RCFM.

and (c)], the effectiveness of the frequency spreading will be lowered [6] and the harmonic components around the switching frequency will be increased. Nevertheless, with the same frequency deviation, RCFM gives a more effective way to disperse the harmonics around the switching frequency than the PCFM.

Due to different spectral behaviors of the two schemes, they can be applied to applications of different requirements. PCFM is more suitable for applications that require tight output regulation, like dc/dc converters, since it introduces lower undesirable low-frequency harmonics at the output. On the other hand, RCFM is more suitable for applications that do not require tight output regulation, like power-factor correctors, since output regulation is usually performed in the second stage and high-frequency harmonic suppression is more effective.

# IV. CONCLUSION

A comparative study on using PCFM and RCFM schemes for EMI suppression of PWM converters has been presented. Both schemes can provide a considerable reduction in the conducted EMI, as compared to the standard PWM scheme. With the same frequency deviation in both schemes, RCFM gives a higher high-frequency suppression than the PCFM. However, RCFM introduces higher low-frequency harmonics than the PCFM at the converter output. Rigorous mathematical derivations on the spectral characteristics of the PCFM have

been developed. In addition, due to the finite bandwidth of the resolution filter in EMC analyzers, discrepancies between the theoretical and experimental results have been explained. Selections of the modulation frequency and the frequency deviation in PCFM have been discussed. Finally, theoretical predictions are confirmed by experimental results of a buck converter operating in continuous conduction mode.

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