A CONTINUUM OF H- TO He-RICH TIDAL DISRUPTION CANDIDATES WITH A PREFERENCE FOR E+A GALAXIES

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ABSTRACT

We present the results of a Palomar Transient Factory (PTF) archival search for blue transients that lie in the magnitude range between "normal" core-collapse and superluminous supernovae (i.e., with $-21 \leq M_{R \text{ (peak)}} \leq -19$). Of the six events found after excluding all interacting Type IIn and Ia-CSM supernovae, three (PTF09ge, 09axc, and 09djl) are coincident with the centers of their hosts, one (10iam) is offset from the center, and a precise offset cannot be determined for two (10nuj and 11glr). All the central events have similar rise times to the He-rich tidal disruption candidate PS1-10jh, and the event with the best-sampled light curve also has similar colors and power-law decay. Spectroscopically, PTF09ge is He-rich, while PTF09axc and 09djl display broad hydrogen features around peak magnitude. All three central events are in low star formation hosts, two of which are E+A galaxies. Our spectrum of the host of PS1-10jh displays similar properties. PTF10iam, the one offset event, is different photometrically and spectroscopically from the central events, and its host displays a higher star formation rate. Finding no obvious evidence for ongoing galactic nuclei activity or recent star formation, we conclude that the three central transients likely arise from the tidal disruption of a star by a supermassive black hole. We compare the spectra of these events to tidal disruption candidates from the literature and find that all of these objects can be unified on a continuous scale of spectral properties. The accumulated evidence of this expanded sample strongly supports a tidal disruption origin for this class of nuclear transients.

Key words: accretion, accretion disks - galaxies: nuclei - quasars: supermassive black holes

Online-only material: color figures, machine-readable table

1. INTRODUCTION

The peak luminosities of novae $(-10 \leq M_R \leq -4)$, supernovae (SNe; $-19 \leq M_R \leq -14$), and superluminous SNe (SLSNe; $-21 \leq M_R$; see Gal-Yam 2012 for a review) span a wide but discontinuous range. Discoveries of new types of transients (e.g., van Dyk et al. 2000; Valenti et al. 2009; Foley et al. 2009; Perets et al. 2010; Kasliwal et al. 2011) have been filling

the luminosity gap between novae and SNe. The gap between SNe and SLSNe, however, is less explored (Figure 1). Here we present the results of a search for such transients in the Palomar Transient Factory (PTF; Law et al. 2009; Rau et al. 2009) archive, focusing on events originally classified as corecollapse SNe.

PTF is an untargeted survey that discovered and classified more than 500 core-collapse SNe between the years 2009 and 2012 (iPTF²⁴ continues as the successor of this survey).

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	Table 1
A	All PTF Events Originally Classified as Type II SNe due to the Presence of Broad H or a Blue Continuum in Their Spectra,
	Having a Peak Absolute Magnitude between -19 and -21

Name	Туре	α (J2000)	δ (J2000)	Redshift	Discovery	Discovery	Peak M_R	Host Offs	et [mas]	ND
					Date	m_R		P48	SDSS	
09ge	He	14:57:03.10	+49:36:40.8	0.064	2009 May 7	19.22 ± 0.05	-19.64 ± 0.02	84 ± 42	82 ± 81	1.0
09axc	Н	14:53:13.06	+22:14:32.2	0.1146	2009 Jun 20	20.83 ± 0.12	-19.53 ± 0.04	119 ± 60	78 ± 84	0.9
09djl	Н	16:33:55.94	+30:14:16.3	0.184	2009 Jul 24	20.76 ± 0.10	-20.20 ± 0.10	175 ± 155	84 ± 109	0.8
10iam	Н	15:45:30.85	+54:02:33.0	0.109	2010 May 22	20.73 ± 0.26	-20.13 ± 0.02	929 ± 67	934 ± 81	11.4
10nuj	Н	16:26:24.70	+54:42:21.6	0.132	2010 Jun 13	21.29 ± 0.17	-19.31 ± 0.05	304 ± 68	217 ± 82	2.6
11glr	Н	16:54:06.13	+41:20:14.8	0.207	2011 May 28	21.44 ± 0.29	-19.85 ± 0.08	1199 ± 1390	941 ± 758	1.2

Notes. Here we exclude interacting Type IIn and Ia-CSM events. Coordinates are based on P48 images astrometrically calibrated to SDSS (with typical errors of 0''.1). Peak magnitudes refer to the brightest measured photometric point and its error. The offset of each target with respect to its host center was determined using relative image registration and the normalized distance (ND) is noted (see Section 3.1 for details). We separate the sample into three classes: central events (ND \leq 1), offset events (ND \geq 3), and intermediate events (1 < ND < 3) for which we cannot robustly conclude a coincidence or separation from their host center.

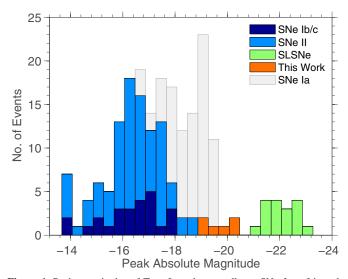


Figure 1. Peak magnitudes of Type Ia and core-collapse SNe from Li et al. (2011; see Filippenko 1997 for a review of SN types), SLSNe from Gal-Yam (2012), and the events presented in this work.

(A color version of this figure is available in the online journal.)

Of those, we search for events originally classified as Type II SNe due to the presence of broad (a few thousand km s⁻¹) hydrogen lines or a blue continuum in their spectra and displaying a peak absolute magnitude $-21 \leq M_R \leq -19$. We exclude from this sample all Type IIn events (SNe showing narrow emission lines; e.g., Schlegel 1990; Kiewe et al. 2012) and Ia-CSM events (e.g., Hamuy et al. 2003; Silverman et al. 2013) known to extend into this magnitude range (e.g., Stoll et al. 2011) likely due to emission powered by interaction with a dense circumstellar medium. These events were typed using both visual and automatic classifications carried out with the SN spectral fitting codes Superfit (Howell et al. 2005) and SNID (Blondin & Tonry 2007). We are left with the six events whose spectra do not match those of known types of interacting SNe. These events are presented in Table 1 and Figure 2.

Several types of transients not necessarily related to massive stars could still fulfill all of our search criteria. One such example is the flare resulting from a tidal disruption event (TDE; Rees 1988). A TDE can occur when a star passes close enough to a supermassive black hole (SMBH) and is

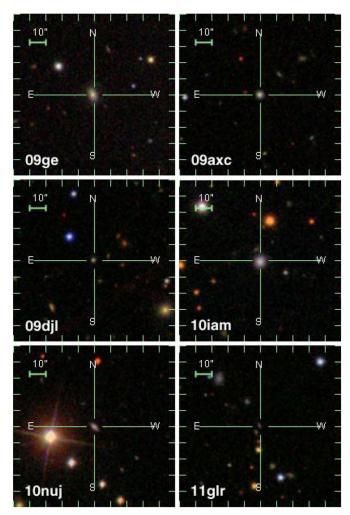


Figure 2. SDSS DR10 images centered at the locations of the events listed in Table 1.

(A color version of this figure is available in the online journal.)

torn apart by gravitational tidal forces. For a star of mass M_* and radius R_* , and a SMBH with mass M_{BH} , this will occur at $R_T \sim (M_{BH}/M_*)^{1/3} R_*$. Part of the disrupted star is accreted onto the SMBH and part is unbound. If the mass of the

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SMBH satisfies²⁵

$$M_{\rm BH} \lesssim 10^8 \ M_{\odot} \cdot \left(\frac{R_*}{R_{\odot}}\right)^{3/2} \left(\frac{M_*}{M_{\odot}}\right)^{-1/2},\tag{1}$$

then the tidal disruption radius is greater than the SMBH event horizon and a flare of radiation is expected to be observed. The observational signature of this flare depends on the structure of the accreting debris disk and on the morphology of the ejected material (Ulmer 1999; Bogdanović et al. 2004; Strubbe & Quataert 2009; Guillochon et al. 2014).

The expected rate of TDEs lies in the range of 10^{-5} to 10^{-4} events per galaxy per year (e.g., Donley et al. 2002; Wang & Merritt 2004; Kesden 2012; but see also Alexander 2012 for a discussion on sources of uncertainty for this rate). A few TDE candidates were found in X-ray data from *ROSAT* (Komossa & Bade 1999; Donley et al. 2002; Halpern et al. 2004), *Chandra* (Komossa et al. 2004), and *XMM-Newton* (Esquej et al. 2007); in γ -ray and X-ray data from *Swift* (Bloom et al. 2011; Burrows et al. 2011; Levan et al. 2011; Zauderer et al. 2011; Cenko et al. 2012); in UV data from *GALEX* (Gezari et al. 2006; Gezari et al. 2009); and in optical data from the Sloan Digital Sky Survey (SDSS²⁶; van Velzen et al. 2011). The high-energy emission is associated with a jet pointing in our line of sight, while the soft X-ray/UV-optical flare is usually identified with thermal emission coming from an accretion flow.

One expected observational signature of such accretionpowered emission is a light curve decline rate of $t^{-5/3}$ which follows from a $t^{-5/3}$ expected decay in the accretion rate (Rees 1988; Evans & Kochanek 1989; Phinney 1989).²⁷ The $t^{-5/3}$ mass accretion rate assumes a uniform energy distribution for the post-disrupted material. Lodato et al. (2009) show that this assumption fails for certain stellar density profiles and that TDE light curves may deviate from $t^{-5/3}$ at early times. An expanding disk may alter this rate also at late times (Shen & Matzner 2014) and it may be entirely different for partial disruptions (Guillochon & Ramirez-Ruiz, 2013). Finally, Strubbe & Quataert (2009) warn that even if the mass accretion rate does decay as $t^{-5/3}$, the flux in a given band will not necessarily follow the same rule.

Gezari et al. (2012; hereafter G12) presented the joint discovery of PS1-10jh in the optical and UV by Pan-STARRS1 (PS1) and the *GALEX* Time Domain Survey (TDS), respectively, and identified it as a likely TDE. The event is a luminous (bolometric $L_{\text{peak}} \sim 10^{44}$ erg s⁻¹) and blue ($T_{\text{blackbody}} \sim 30,000$ K) transient declining on a timescale of a few months. It is the first TDE candidate to have its rise to peak well sampled and the light curve is roughly consistent with the models of Lodato et al. (2009).

Chornock et al. (2014) report on the PS1 and *GALEX* TDS discovery of PS1-11af, which is similar to PS1-10jh but with a lower effective temperature. They find good fits to models with an accreted mass as small as $0.002 M_{\odot}$, leading them to suggest that PS1-11af was a partial TDE (i.e., the star was not fully disrupted). Their spectra show no clear features except for a broad UV absorption component (possibly associated with Mg II), but their wavelength coverage cannot be used to rule

out any H α emission. PS1-10jh, in contrast, displayed broad emission lines at rest-frame wavelengths of 3203 Å and 4686 Å, interpreted as He II, with no obvious signs of hydrogen.

G12 explain the lack of observed hydrogen in PS1-10jh with stellar winds or stripping of the star during previous passages near the SMBH (Bogdanovic et al. 2014 discuss the plausibility of the latter scenario). Guillochon et al. (2014) claim that hydrogen would not be visible at early times even if it were present in the disrupted star. In their model, the optical emission is dominated by the bound material in a regime where the hydrogen is fully ionized and therefore not observable until long after peak.

Wang et al. (2011) searched for TDE candidates in SDSS by looking for narrow high-ionization coronal lines. The spectrum of one of their objects, SDSS J074820.66+471214.6 (hereafter SDSS J0748), showed also broad emission features around He II 4686 Å *and* around H α . Yang et al. (2013) re-observed SDSS J0748 several years after the initial SDSS spectrum and found that the broad features had disappeared, confirming their transient nature. However, with no light curve for this event, it is hard to compare it to PS1-10jh directly.

PS1-10jh was also detected by PTF as PTF10onn and marked as a possible TDE by the autonomous software framework Oarical (Bloom et al. 2012). With a peak magnitude of $M_r =$ -19.5, it would have come up in the archival search discussed here, but a spectrum was never obtained as part of PTF followup. We do, however, find very similar events. One (PTF09ge; Kasliwal et al. 2009) displays a nearly identical He II emission feature as SDSS J0748 but with no hydrogen and with a light curve very similar to PS1-10jh. Two more of the events found in our search (PTF09axc and PTF09djl) show broad hydrogen features starting from the earliest spectra (taken a few days after peak magnitude).

These two H-rich events happen to be located in the centers of rare E+A galaxies (Dressler & Gunn 1983). Such galaxies show no emission lines that are indicative of ongoing or recent star formation. They are thus not likely to host core-collapse SNe that originate in massive, short-lived stars. Our spectrum of the host galaxy of PS1-10jh shows similar features. E+A galaxies (sometimes referred to as K+A) are so called because the Balmer absorption features in their spectra (characteristic of A stars) appear superimposed on an old K star or (E)arlytype galaxy population. The A stars would have been formed in an episode of star formation that ceased abruptly \sim 1 Gyr ago (Dressler & Gunn 1983), possibly following a merger.

Recently, Prieto et al. (2014) and Holoien et al. (2014) reported the discovery of ASASSN-14ae as a likely TDE. We find that it displays spectral features similar to those seen in SDSS J0748, with broad emission in both He II and H.

We note that another type of transient expected to occur exclusively in galaxy centers was recently suggested by Balberg et al. (2013). They propose that hypervelocity stellar collisions could result in bright SN-like transients. However, their prediction for the peak magnitudes of such flares are at the lower end of the known SN luminosity scale, much fainter than the events discussed here.

We present photometric and spectroscopic data of all six events found in our archival search in Section 2, and analyze these observations in Section 3. In Section 4, we briefly discuss the events that could not be robustly associated with the center of their host. We then focus on the three central events in Section 5, comparing them to the TDE candidates mentioned above. We conclude in Section 6.

²⁵ Using a Newtonian calculation for a non-spinning black hole (see, e.g., Kesden 2012 for a relativistic derivation yielding higher limits for rotating black holes).

²⁶ http://www.sdss3.org/

²⁷ This light curve decline rate is similar to that of Type IIL SNe, but those events typically peak at lower luminosities (Arcavi et al. 2012) compared to what is expected for a TDE.

 Table 2

 Photometric Observations (Upper Limits Mark 3σ Non-detections)

Object	Telescope	Filter	JD	Mag	Error
09ge	P48	g	2454910.954	>22.001	
09ge	P48	g	2454917.759	>21.349	
09ge	P48	g	2454918.904	>21.597	
09ge	P48	g	2454921.74	>21.384	
09ge	P48	g	2454921.817	>21.532	
09ge	P48	g	2454975.742	17.833	0.008
09ge	P48	g	2454975.803	17.852	0.011

(This table is available in its entirety in a machine-readable form in the online journal. A portion is shown here for guidance regarding its form and content.)

2. OBSERVATIONS

All the events from our archival search were discovered by the Palomar 48 inch Oschin Schmidt Telescope (P48) as part of the PTF survey using the Mould *R*-band filter. The discovery details are presented in Table 1.

2.1. Photometry

We obtained photometric observations in the R and g bands using P48, and in g, r, and i bands with the Palomar 60 inch telescope (P60; Cenko et al. 2006). Initial processing of the P48 images was conducted by the Infrared Processing and Analysis Center (IPAC; Laher et al. 2014). Photometry was extracted using a custom PSF fitting routine (e.g., Sullivan et al. 2006), which measures the transient flux after image subtraction (using template images taken before the outburst or long after it faded). We calibrate our light curves to the SDSS system using SDSS observations (Ahn et al. 2012) of the same field when possible and to USNO-B (Monet et al. 2003) reference magnitudes otherwise. We average magnitudes obtained from the same filter and instrument taken on the same night. We correct for Galactic extinction using the Schlafly & Finkbeiner (2011) maps, extracted via the NASA Extragalctic Database (NED²⁸). Distance moduli are calculated from spectroscopic redshifts, measured using narrow host features. A cosmological model with $H_0 = 70 \,\mathrm{km \, s^{-1} \, Mpc^{-1}}$, $\Omega_m = 0.3$ and $\Omega_{\Lambda} = 0.7$ is assumed throughout. Our photometry is presented in the AB system in Table 2 and Figure 3.

Due to gaps in the photometry, we note that all phases stated hereafter (regarding the PTF sample) should be considered with few-day uncertainties.

2.2. Adaptive Optics Imaging

We imaged the host galaxy of PTF09djl on 2013 July 5 (several years after the transient faded) using the Laser Guide Star Adaptive Optics (AO) system (Wizinowich et al. 2006) and the NIRC2 camera on the Keck II 10 m telescope in the K_p band with the 40" square field of view "wide" camera.

The image frames were dark-subtracted and flat-fielded in the standard manner using custom Python scripts. The effects of fringing were removed by subtracting a sky fringe scaled to the sky brightness of each frame. Since 2007, a variable glow has been present at the lower right corner of the "wide" camera that cannot be adequately removed by calibration. A triangular section of the frames was masked before registering and coadding the images. The resulting image is shown in Figure 6.

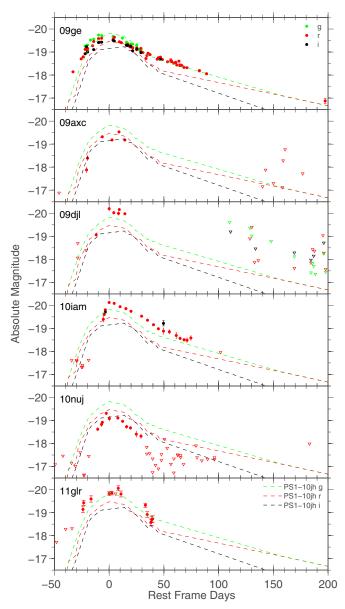


Figure 3. Light curves of PTF09ge, PTF09axc, PTF09djl, PTF10iam, PTF10nuj, and PTF11glr from P48 and P60. Most display light curve shapes consistent with those of PS1-10jh (G12; dashed lines). PTF10iam shows a uniquely fast rise to peak and PTF10nuj shows a sudden luminosity drop \sim 30 days after peak.

(A color version of this figure is available in the online journal.)

2.3. Spectroscopy

Spectra of the events were obtained with the Double Beam Spectrograph (DBSP; Oke & Gunn 1982) mounted on the Palomar 200 inch telescope (P200) and the Low Resolution Imaging Spectrometer (LRIS; Oke et al. 1995) mounted on the Keck I 10 meter telescope (Table 3). The data were reduced using standard IRAF²⁹ and IDL routines. Host galaxy spectra were obtained in 2013, after all transient emission had faded (except for PTF10iam, for which a host spectrum was obtained by the SDSS in 2002 and downloaded via DR10; Ahn et al. 2014). Spectra of the transients are presented in Figures 4 and 12, and

²⁸ http://ned.ipac.caltech.edu/

²⁹ IRAF, the Image Reduction and Analysis Facility, is a general purpose software system for the reduction and analysis of astronomical data. IRAF is written and supported by the National Optical Astronomy Observatories (NOAO) in Tucson, Arizona.

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 Table 3

 Spectroscopic Observations

	1			
Object	UT Date	Phase	Telescope	Instrument
09ge	2009 May 20	-24	P200	DBSP
09axc	2009 Jul 22	7	Keck I	LRIS
09djl	2009 Aug 25	2	Keck I	LRIS
09djl	2009 Sep 23	31	Keck I	LRIS
09djl	2009 Oct 24	62	Keck I	LRIS
10iam	2010 Jun 8	5	Keck I	LRIS
10iam	2010 Jul 7	34	Keck I	LRIS
10iam	2010 Jul 18	45	P200	DBSP
10nuj	2010 Jul 14	11	P200	DBSP
11glr	2011 Jun 29	-16	Keck I	LRIS
11glr	2011 Aug 28	45	Keck I	LRIS
09ge Host	2013 Sep 9		Keck I	LRIS
09axc Host	2013 May 9		Keck I	LRIS
09djl Host	2013 May 9		Keck I	LRIS
10nuj Host	2013 May 9		Keck I	LRIS
11glr Host	2013 Oct 4		Keck I	LRIS

Note. The phase is denoted in days relative to *R*-band maximum.

host galaxy spectra are shown in Figure 8. Digital versions of our spectra are available online through the Weizmann Interactive Supernova data REPository (WISeREP³⁰; Yaron & Gal-Yam 2012).

2.4. Radio

We observed the field of PTF09axc with the Jansky Very Large Array (VLA) on 2014 June 28. The observations were performed at both 3.5 GHz (*S* band) and 6.1 GHz (*C* band) in the D configuration. We used J1513+2388 as a phase calibrator and 3C286 as a flux calibrator. The data were analyzed using standard calibration scripts within the CASA software (McMullin et al. 2007).

3. THE SAMPLE

3.1. Offsets from Host Centers

After registering the images of each event, we measure its position from a host-subtracted image taken near peak magnitude and compare it to the host centroid measured in a co-added reference image using the IRAF task IMEXAMINE. We estimate the error in the measured offset of the transient using the scatter in centroid determinations (when varying the fit type and radii parameters of IMEXAMINE between 2 and 7 pixels) and the registration error (typically ≤ 0.03 pixels, corresponding to $\leq 0''.03$ in the P48 images).

We also measure the host galaxy centroids in the SDSS *r*band images and register them to the P48 images. Here the registration errors are closer to $\sim 0''.06$, but the finer pixelscale of the SDSS images contributes to a more accurate host centroid determination for the brighter hosts. We repeat these measurements for the P48 photometric data of PS1-10jh, for comparison.

The offsets of our events (as well as that of PS1-10jh), using the comparison to both P48 host centers and SDSS host centers, are presented in Figure 5.

We normalize each apparent offset to its galaxy *expRad_r* parameter (an estimate of the half-light radius of the galaxy) from SDSS DR10. We then divide each normalized offset by its

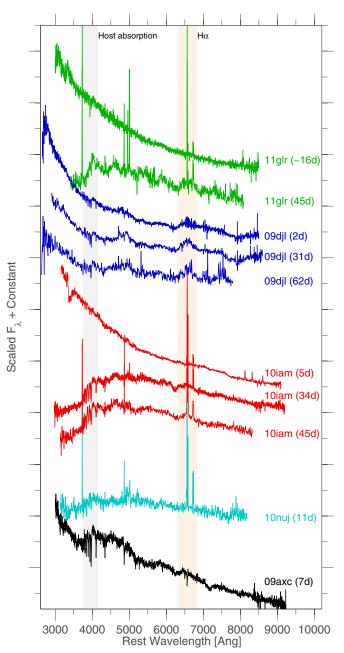


Figure 4. Spectra of PTF09axc, PTF09djl, PTF10iam, PTF10nuj, and PTF11glr. Phases are shown relative to peak. Host absorption features and the broad H α region are marked with thick bands.

(A color version of this figure is available in the online journal.)

error and call this parameter ND (for "Normalized Distance"). ND measures how many sigma each event is from its host center, in terms of its half-light normalized distance. The results are shown in Table 1.

We find that one event (PTF10iam), with ND = 11.4, is obviously offset from the center of its host. Three events (PTF09ge, PTF09axc, and PTF09djl) are coincident with the centers of their hosts within the errors (ND \leq 1). For the two remaining events (PTF10nuj and PTF11glr), with 1 < ND < 3, we do not conclude whether they are coincident with their host centers or not.

For PTF09djl, we use the AO imaging of its host to further constrain its position relative to the host center. We find a WCS solution for the PTF09djl AO image using Aladin

³⁰ http://www.weizmann.ac.il/astrophysics/wiserep

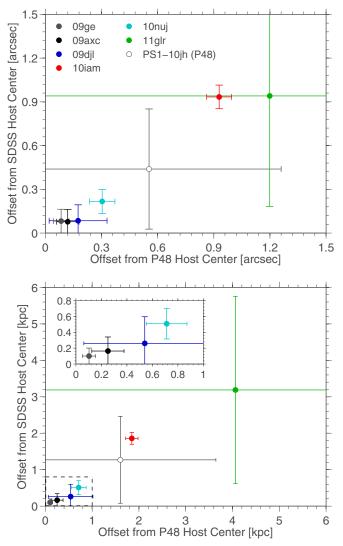


Figure 5. Offsets in arcseconds (top) and kiloparsecs (bottom) between the locations of our events and the centers of their hosts, determined both from P48 and from SDSS images of the hosts registered to P48 images of the transients. Error bars denote 1σ uncertainties. The region enclosed in the dashed box in the bottom panel is shown in greater detail in the inset. Three events (PTF09ge, PTF09axc, and PTF09djl) are found to be coincident with the centers of their hosts with relatively small errors, one event (PTF10iam) is found to be offset from its host center, and it is not possible to determine either an offset nor a robust coincidence with the center for the remaining two events (PTF10nuj and PTF11glr). The position of PS1-10jh (from its P48 detection as PTF10onn) is shown for comparison.

(Bonnarel et al. 2000) by comparing to 2MASS stars in the field and register this image to the P48 resampled SDSS image based on its WCS solution. We then adjust the registration using a mean offset of three objects in the field (one of which is the host galaxy of PTF09djl itself; see Figure 6). We superimpose the SN position from P48 onto this registered AO image, taking into account the centroid errors and offsetting errors. The results are shown in the inset of Figure 6. We conclude that the position of PTF09djl remains consistent with the center of its host under this analysis.

3.2. Blackbody Fits

We fit a combination of a scaled host galaxy spectrum and a blackbody function to the spectra of each event. The best-

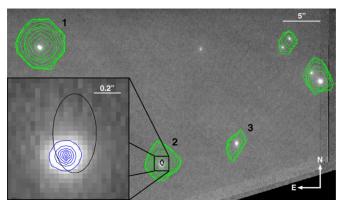


Figure 6. NIRC2 AO image of the position of PTF09djl taken on 2013 July 5 (after the transient was no longer visible). The contours from the P48 resampled SDSS image are shown in green after WCS registration and shifting to match the centroids of the numbered sources. The P48 position of PTF09djl, taking into account the centroid and shifting errors (1σ), is shown in the black ellipse. Inset: Enlarged area around the host of PTF09djl, with the AO image contours shown in blue and the position of PTF09djl marked in black.

(A color version of this figure is available in the online journal.)

fit effective temperatures and radii (for blackbody fits yielding T > 6000 K, so as not to be dominated by lines over the continuum) are displayed in Figure 7. We are not able to correct for unknown host extinction and therefore consider the measured temperatures to be lower limits.

Two of the central events, PTF09ge and PTF09djl, display high effective temperatures (consistent with PS1-10jh) and smaller radii compared the rest of the sample. The non-central events are cooler initially and continue to cool with time. The third central event, PTF09axc, displays an intermediate temperature and radius.

3.3. Host Galaxies

We obtain *ugriz* magnitudes of all host galaxies from SDSS DR10 and analyze them using Z-PEG (Le Borgne & Rocca-Volmerange 2002), which is based on the spectral synthesis code PEGASE.2 (Fioc & Rocca-Volmerange 1997). Z-PEG fits the observed galaxy magnitudes with SED templates (of SB, Im, Sd, Sc, Sbc, Sb, Sa, S0, and E spectral types) to obtain the stellar mass and star formation rate. We assume a Salpeter (1955) initial mass function and fit also a foreground dust screen varying in color excess from E(B - V) = 0 to 0.2 magnitudes. The results are presented in Table 4.

Spectra of the host galaxies are presented in Figure 8. All of them display Balmer absorption features. While the host galaxies of PTF10iam, PTF10nuj, and PTF11glr show emission lines, those of PTF09axc and PTF09djl (which are two of the three events coincident with the centers of their hosts) display spectra similar to E+A galaxies (Dressler & Gunn 1983), with no strong indications of ongoing or recent star formation. The host of the third central event (PTF09ge) is similar to E+A galaxies but shows some H α emission. We do not see narrow coronal emission lines in any of our host spectra taken after the transient had faded (see Wang et al. 2012; Yang et al. 2013).

We fit the host spectra to stellar templates provided by the MILES empirical stellar library (Sánchez-Blázquez et al. 2006; Vazdekis et al. 2010), using PPXF (Cappellari & Emsellem 2004) and GANDALF (Sarzi et al. 2006). We measure the star formation rate (SFR) from the H α luminosity using the conversion of Kennicutt (1998). We calibrate the gas-phase metallicity in our sample using the PP04 "N2" method (Pettini & Pagel 2004)

	Properties of the Host Galaxies of Our Archival Sample								
Host		Photometric Analysis	Spectroscopic Analysis						
	M ($10^{10} M_{\odot}$)	SFR $(M_{\odot} \text{ yr}^{-1})$	$sSFR (10^{-10} yr^{-1})$	SFR $(M_{\odot} \text{ yr}^{-1})$	$12 + \log(O/H)$	[M/H]	Age (Gyr)		
09ge	1.05 (1.03, 1.35)	n/a	n/a	0.10 (0.05)	8.8732 (0.0638)	-0.196	7.035		
09axc	1.23 (1.16, 1.28)	<16.11	<13.10	0.04 (0.02)	n/a	-0.356	4.469		
09djl	1.86 (1.07, 3.73)	3.42 (1.22, 4.19)	1.84 (0.34, 3.02)	n/a	n/a	-0.218	4.461		
10iam	2.94 (2.83, 4.39)	8.51 (6.53, 9.40)	2.90 (1.59, 3.02)	2.86 (1.41)	8.672 (0.233)	-1.114	5.717		
10nuj	1.22 (1.18, 2.89)	3.86 (2.59, 4.03)	3.16 (1.05, 3.16)	8.11 (3.74)	8.698 (0.070)	-0.284	5.743		
11glr	0.30 (0.28, 0.65)	1.09 (0.73, 1.15)	3.58 (1.44, 3.58)	1.22 (0.56)	8.305 (0.080)	n/a	n/a		

Table 4

Notes. Obtained using Z-PEG SED template fits to SDSS ugriz photometry of the hosts and from the host spectral features using the same methods discussed by Pan et al. (2014). Values in parenthesis (when noted) describe the lower and upper limits deemed acceptable by Z-PEG or 1σ errors from the spectral analysis. Cases where no good fits were found are denoted by "n/a."

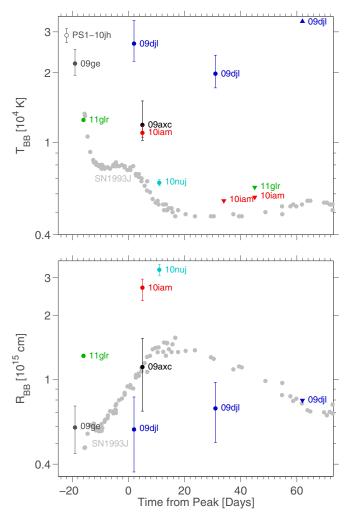


Figure 7. Top: best-fit blackbody temperatures (when a good fit was possible) obtained from the optical host-subtracted spectra (triangles denote upper or lower limits). No host-extinction correction is performed. The blackbody temperatures of PS1-10jh (G12) and SN1993J (Richmond et al. 1994), assuming low extinction, are shown for comparison. Bottom: blackbody radii estimated from the fits

(A color version of this figure is available in the online journal.)

following the procedure described in Kewley & Ellison (2008). We further determine the mass-weighted stellar metallicity and age by fitting the full stellar continuum with PPXF and weight averaging all the stellar templates. The SFR determined for galaxies showing insignificant or no H α lines in their spectrum

should be interpreted with caution here as it depends on the stellar absorption fitted by the model templates. For more details on these procedures see Pan et al. (2014). All best-fit parameters are presented in Table 4.

4. THE NON-CENTRAL EVENTS

PTF10iam is the only event in our sample for which a clear offset from the center of its host can be determined. It shows a faster rise to peak luminosity compared to the other events (Figure 3) and it displays a broad absorption feature near rest wavelength 6200 Å (Figure 4), not seen in the spectra of the central events. If interpreted as high velocity H α , this feature may be an indication of interaction (Chugai et al. 2007), which does not manifest itself in narrow emission lines. In any case, it is clear that this off-center event is different photometrically and spectroscopically from the rest of the sample. We discuss PTF10iam in detail in a companion paper (I. Arcavi et al., in preparation).

For PTF10nuj and PTF11glr, we are not able to measure or rule out an offset of their position relative to their host center. We leave the detailed analysis of the photometry and spectroscopy of these two events for a future publication.

5. THE CENTRAL EVENTS AS TDE CANDIDATES

We now focus on the three events coincident with the centers of their hosts: PTF09ge, PTF09axc, and PTF09djl.

5.1. Could They Be AGNs?

A transient associated with the center of a galaxy may be related to galactic nuclei activity. However, the spectra taken during the outbursts (and those taken several years later) do not show obvious emission lines typical of active galactic nuclei (AGNs).

To evaluate the central emission more carefully, we subtract stellar spectral templates from the host galaxy integrated spectra using the Bruzual & Charlot (2003) stellar population synthesis models via the STARLIGHT code (Cid Fernandes et al. 2005). We find no signs of AGN emission lines in the host spectra of PTF09ge and PTF09djl, but for the host of PTF09axc, we detect an [O III] 5007 Å luminosity of $2.4 \pm 0.3 \times 10^{39}$ erg s⁻¹ (Figure 9). We find an $[O_{III}]/Hbeta$ ratio of >3.4, which indicates that the host galaxy of PTF09axc may contain a very weak AGN.

We find no X-ray sources in the ROSAT All-Sky Survey at the positions of our targets between 1990 July and 1991 February. Comparing to other ROSAT sources in the field, and using the Voges et al. (1999) conversions assuming a hardness ratio

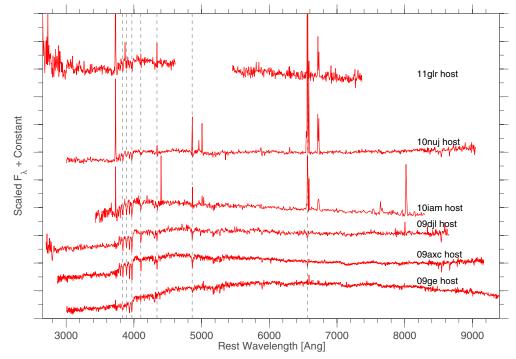


Figure 8. Host galaxy spectra for our archival sample (that of PTF10iam is from SDSS DR10). The hydrogen Balmer series, evident in absorption in E+A galaxies, is marked (dashed lines). The hosts of PTF09ge, 09axc and 09djl (the same events coincident with the centers of their hosts) do not show any strong emission lines, indicating no (or very low) ongoing star formation, and display spectra similar to those of E+A galaxies. (A color version of this figure is available in the online journal.)

Table 5
ROSAT and Swift X-Ray Observations at the Locations
of Our Three Nuclear Transients

Object	ROSAT 0.1-2.5 keV	Swift XRT 0.3–10 keV
	(7/1990 to 2/1991)	(3/2014 to 4/2014)
	(erg s^{-1})	(erg s^{-1})
09ge	$< 4.33 \times 10^{42}$	$<2.27 \times 10^{42}$
09axc	$< 1.86 \times 10^{43}$	$7.13^{+12.22}_{-3.06} \times 10^{42}$
09djl	$< 6.28 \times 10^{43}$	$<1.91 \times 10^{43}$

Notes. Two of the three nuclear transients are not detected in either telescope (3σ upper limits shown). One (PTF09axc) is detected in the XRT data. This luminosity is marginally consistent with a weak AGN (Heckman et al. 2005) and too bright for an X-ray binary (Hornschemeier et al. 2005).

(HR1) <0.5, we derive 3σ upper limits for the X-ray flux of at the positions of PTF09ge, PTF09axc, and PTF09djl as listed in Table 5.

We obtained target-of-opportunity X-ray observations of the fields of PTF09ge, PTF09axc, and PTF09djl with the *Swift* satellite (Gehrels et al. 2004) during 2014 March–April (i.e., five years after the outbursts). Data obtained by the onboard X-Ray Telescope (XRT; Burrows et al. 2005) was analyzed with the automated gamma-ray burst pipeline outlined in Evans et al. (2009). We find no X-ray flux at the positions of PTF09ge and PTF09djl (the derived limits listed Table 5 assume a power-law spectrum with a photon index of 2), but we do detect X-ray emission at the position of PTF09axc corresponding to a luminosity of $7.13^{+12.22}_{-3.06} \times 10^{42}$ erg s⁻¹. This luminosity, together with that in the host [O III] 5007 Å line discussed above, is roughly consistent with AGNs on the low-luminosity end of the Heckman et al. (2005) local sample. The X-ray luminosity

is likely too high (given the stellar mass of this galaxy) for an accreting binary origin (Hornschemeier et al. 2005). We currently cannot constrain any time variability in this X-ray signal and therefore are not able to determine if it is related to PTF09axc directly.

Van Velzen et al. (2011) rejected a few of their TDE candidates as AGNs based on photometric variability beyond the season containing the flare. Here we find no evidence for additional eruptions of our events during the years of available PTF coverage (Figure 10).

We conclude that these three outbursts are not likely due to AGNs, though the host of PTF09axc may also contain an extremely weak AGN.

5.2. A Sequence of H- to He-rich Events

PTF09ge shows very similar photometric and spectroscopic behavior to PS1-10jh (Figures 11 and 12), identified by G12 as a likely He-rich TDE. Both objects display broad He II emission superimposed on a blue continuum. We find a ~1000 km s⁻¹ blueshift in the He II 4686 Å emission peak of PTF09ge. A similar blue "wing" was observed in PS1-10jh (see inset of Figure 12), suggesting that the PS1-10jh line profile is made of two components: an extended blueshifted component (seen also in PTF09ge) and an intermediate-width component (not seen in PTF09ge, but apparent also in the Hα profile of ASASSN-14ae; Holoien et al. 2014). For PTF09ge, we further identify a possible broad absorption feature redshifted by ~3000 km s⁻¹, but this could be due to remaining narrow Fe II 5018 Å and 5169 Å contamination from the host.

The single-band, sparsely sampled light curves of PTF09axc and PTF09djl show similar rise times to PTF09ge (Figure 3) but their spectra exhibit broad H α emission (Figure 4). Van Velzen et al. (2011) noted possible hydrogen emission in

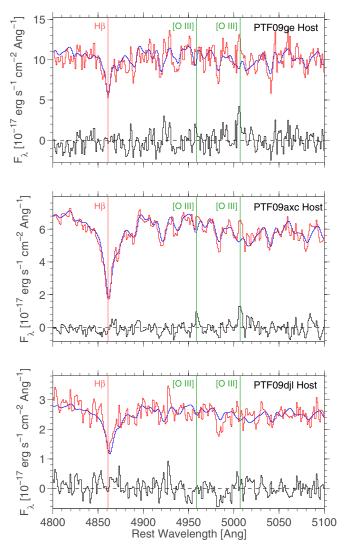


Figure 9. Host spectra (red), fitted stellar spectral templates (blue), and residuals (black). The spectra and templates are shifted in flux for clarity. PTF09axc shows an [O III] 5007 Å emission line in the residual spectrum, marginally consistent with a weak AGN.

their TDE2 as well, albeit narrower and less prominent than in our events.

Following the announcement of ASASSN-14ae by Prieto et al. (2014), and later discovery by iPTF, we obtained spectra of it with DBSP on P200 on 2014 February 1 and 2014 April 4 (Figure 13). The later spectrum shows emission lines of both H and HeII, similar to those noted for SDSS J0748 by Wang et al. (2011).

In total, these seven TDE candidates span a continuous sequence of spectral types, from He-dominated (PS1-10jh, PTF09ge) to H-dominated (PTF09djl, PTF09axc, and possibly TDE2) through intermediate H+He events (SDSS J0748, ASASN-14ae). The spectra of these events have different continuum shapes (either due to intrinsic differences between the events or because of extinction differences). To isolate the differences in the emission features, we remove a second-order polynomial fit from each spectrum and present the continuum-subtracted spectra of this sample in Figure 14.

The He-dominated spectra in Figure 14 are taken at earlier phases compared to the rest, suggesting the features observed

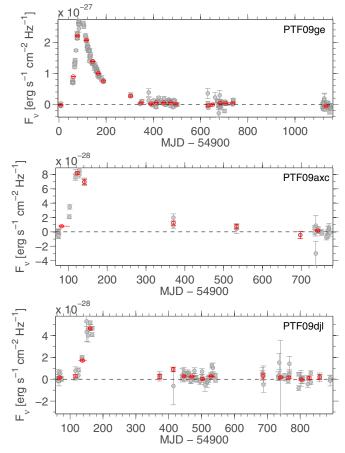


Figure 10. Long-term host-subtracted *R*-band P48 light curves showing no subsequent activity at the locations of PTF09ge, PTF09axc, and PTF09djl for a few hundred days after the original outburst (gray: raw data; red: 30 day binned data).

(A color version of this figure is available in the online journal.)

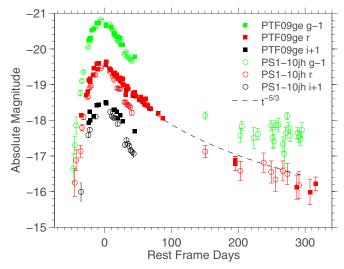


Figure 11. *g*, *r*, and *i* band light curves of PTF09ge from P48 and P60, compared to those of PS1-10jh (G12). Both events show very similar photometric behavior. A $t^{-5/3}$ decline rate is also shown.

(A color version of this figure is available in the online journal.)

could be time-dependent. However, PS1-10jh remained Hedominated out to 254 days post-peak (G12), while ASASSN-14ae displays H from its first spectrum (Holoien et al. 2014; Figure 13).

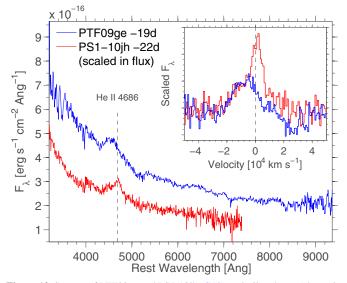


Figure 12. Spectra of PTF09ge and PS1-10jh (G12) at similar phases (shown in days relative to peak). The PS1-10jh spectrum is scaled in flux to the distance of PTF09ge. Both events show very similar spectral properties, namely, broad He II 4686 Å on top of a blue continuum. Inset: the He II 4686 Å line profile of PTF09ge and PS1-10jh after host and continuum subtraction.

5.3. Host Galaxies

We present host galaxy spectra for these seven TDE candidates in Figure 15. The properties of the hosts of the PTF events were discussed in Section 3.3. Here we add our own spectrum of the host galaxy of PS1-10jh (obtained with LRIS on Keck I on 2014 April 29), the TDE2 post-flare spectrum from van Velzen et al. (2011), the SDSS J0742 post-flare spectrum from Yang et al. (2013), and the ASASN-14ae pre-flare spectrum from SDSS DR10. All host galaxies show blue Balmer sequence absorption features typical of E+A galaxies, though some display $H\alpha$, [O II], and [O III] in emission.

The SDSS host galaxy spectrum of ASASSN-14ae, in particular, shows prominent [O III] 5007 Å emission. We perform the same analysis described in Section 5.1 to extract the non-stellar component of the host spectrum (Figure 16). We detect [O III] 5007 Å and H α at 5.2 × 10⁻¹⁶ and 1.7 × 10⁻¹⁶ erg s⁻¹ cm⁻², respectively, in the stellar-subtracted spectrum and find no emission from H β and [O II] 3727 Å to 3σ upper limits of 7×10^{-17} and 3×10^{-17} erg s⁻¹ cm⁻², respectively. For [N II] 6583 Å, contamination from sky emission make it difficult to determine the existence of the line and we set a conservative 3σ upper limit of 10^{-16} erg s⁻¹ cm⁻² for it. The measured [O III]/H β and $[N_{II}]/H\alpha$ ratios are >7.4 and <0.59, respectively. Such a high $[O_{III}]/H\beta$ ratio suggests the presence of central accreting black hole (Baldwin et al. 1981). The low luminosity of [O III] 5007 Å (at 2.3×10^{39} erg s⁻¹) is consistent with the host containing a weak AGN.

We repeat the analysis of Section 3.3 for the hosts of these TDE candidates and present the deduced parameters in Table 6 together with those of the PTF TDE candidates from Table 4 for easy comparison. We translate our measured stellar masses to bulge masses through the Gadotti (2009) conversion using the SDSS R90/R50 ratios³¹ in the *r* band. We then use the relation of Häring & Rix (2004) to translate the bulge mass to

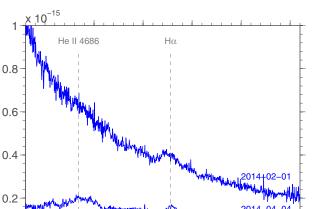


Figure 13. Spectra of ASASSN-14ae. The early spectrum displays broad H α on top of a blue continuum. An additional broad component at He II 4686 Å appears in the later (cooler) spectrum.

6000

Rest Wavelength [Ang]

7000

8000

9000

(A color version of this figure is available in the online journal.)

5000

 F_{λ} [erg s⁻¹ cm⁻² Ang⁻¹]

0

4000

the SMBH mass. The R90 and R50 values for the host galaxy of PS1-10jh have very large errors due to the low luminosity of this galaxy. We therefore adopt the SMBH mass obtained by G12 in their scaling from the host galaxy mass. All SMBH masses are presented in Table 6.

5.4. Line Widths

The velocity widths of the broad emission lines seen during a TDE could be indicative of the region from which the lines are being emitted.

Bound material in a Keplerian orbit at the tidal radius R_T will have a (circular) velocity

$$v_T \approx 43700 \left(\frac{M_{\rm BH}}{10^6 M_{\odot}}\right)^{1/3} \left(\frac{\rho_*}{\rho_{\odot}}\right)^{1/6} \rm km \ s^{-1},$$
 (2)

where ρ_* is the average density of the disrupted star. Assuming $M_{\rm BH} \propto \sigma_G^{\alpha}$ (where σ_G is the velocity dispersion of the galaxy), and $L \propto \sigma_G^4$ (from the Faber–Jackson relation; Faber & Jackson, 1976), it follows that

$$v_T \propto (L_{\text{host}})^{\alpha/12} \rho_*^{1/6}.$$
 (3)

For typical values of α (e.g., $\alpha = 4.42$; Kormendy & Ho 2013), the stellar density will have a small influence and a correlation between the measured TDE line widths and the host galaxy luminosity is expected, assuming the measured velocity indeed represents v_T . Guillochon et al. (2014), however, show that bound material could extend to larger radii than the standard truncated disk model assumes. Therefore bound material could have much lower velocities than v_T .

Another option for the origin of the emission lines is the outflowing unbound material. Strubbe & Quataert (2009) find that the most energetic unbound material will move at velocities of approximately

$$7500 \left(\frac{R_T}{R_P}\right) \left(\frac{M_{\rm BH}}{10^6 \ M_{\odot}}\right)^{\frac{1}{6}} \left(\frac{M_*}{M_{\odot}}\right)^{\frac{1}{3}} \left(\frac{R_*}{R_{\odot}}\right)^{-\frac{1}{2}} \,\rm km \, s^{-1}.$$
 (4)

 $^{^{31}}$ R90 and R50 are the radii enclosing 90% and 50% of the galaxy light, respectively.

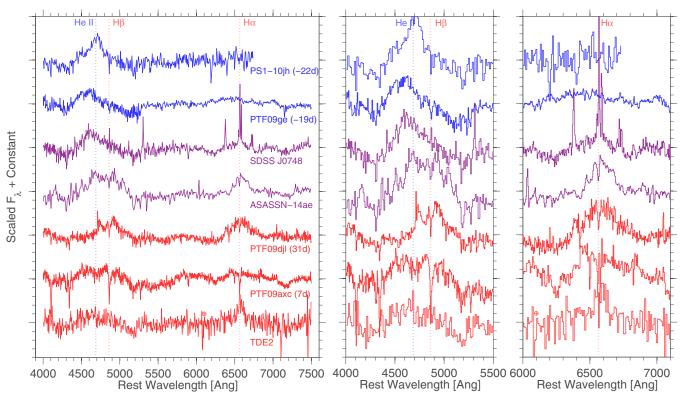


Figure 14. Continuum-subtracted spectra of our three PTF TDE candidates together with PS1-10jh (G12), ASASSN-14ae, SDSS J0748 (Wang et al. 2011), and TDE2 (van Velzen et al. 2011). Phases are shown relative to peak. A progression from He-rich to H-rich events is apparent. The middle and right panels present more detailed views of the regions around the marked lines.

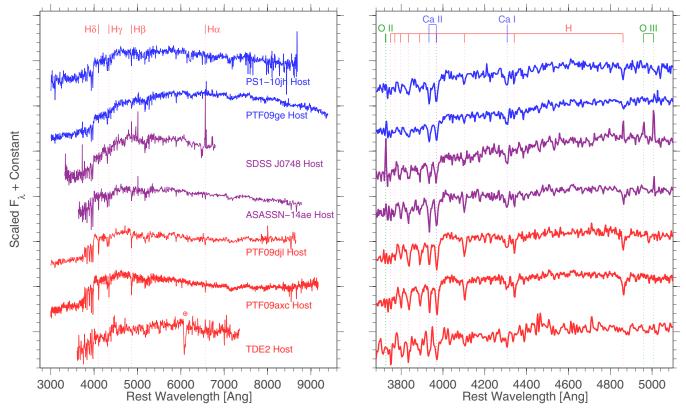


Figure 15. Host galaxy spectra of the TDE candidates from Figure 14 (the spectrum of SDSS J0748 is from Yang et al. 2013; the spectrum of the ASASSN-14ae host is from SDSS DR10). Most show similar signs of low ongoing star formation. Balmer series absorption lines identify some of these hosts as rare E+A galaxies. The right panel presents a more detailed view of the marked lines.

(A color version of this figure is available in the online journal.)

 Table 6

 Same as Table 4 for the Host Galaxies of the PTF TDE Candidates, Together with the Additional TDE Candidates Presented in Figure 14

Host		Spectroscopic Analysis						
	$\mathop{\rm M}\limits_{(10^{10}\ M_\odot)}$	$M_{ m BH} \ (10^6 \ M_{\odot})$	$\frac{\text{SFR}}{(M_{\odot} \text{ yr}^{-1})}$	sSFR (10 ⁻¹⁰ yr ⁻¹)	$\frac{\text{SFR}}{(M_{\odot} \text{ yr}^{-1})}$	12 + log(O/H)	[M/H]	Age (Gyr)
PTF09ge	1.05 (1.03, 1.35)	$5.65^{+3.02}_{-0.98}$	n/a	n/a	0.10 (0.05)	8.873 (0.064)	-0.196	7.035
PTF09axc	1.23 (1.16, 1.28)	$2.69^{+0.66}_{-0.64}$	<16.11	<1.31	0.04 (0.02)	n/a	-0.356	4.469
PTF09djl	1.86 (1.07, 3.73)	$3.57^{+9.97}_{-2.96}$	3.42 (1.22, 4.19)	1.84 (0.34, 3.02)	n/a	n/a	-0.218	4.461
PS1-10jh	0.67 (0.48, 1.12)	4^{+4}_{-2}	1.21 (0.29, 1.28)	1.81 (0.30, 2.02)	n/a	n/a	-0.215	5.599
SDSS J0748	3.40 (2.79, 3.57)	$11.78^{+2.29}_{-3.56}$	2.76 (2.58, 3.58)	0.81 (0.76, 1.22)	0.35 (0.40)	8.760 (1.723)	-0.199	6.265
ASASSN-14ae	0.60 (0.52, 0.8)	$2.45^{+1.55}_{-0.74}$	n/a	n/a	0.02 (0.01)	n/a	-0.407	5.544
TDE2	9.33 (8.43, 13.06)	$35.52^{+55.31}_{-25.80}$	1.01 (0.53, 1.22)	0.11 (0.04, 0.13)	n/a	n/a	-0.319	4.627

Note. We add the SMBH mass calculated using the Gadotti (2009) and Häring & Rix (2004) relations (for PS1-10jh, we adopt the value from G12).

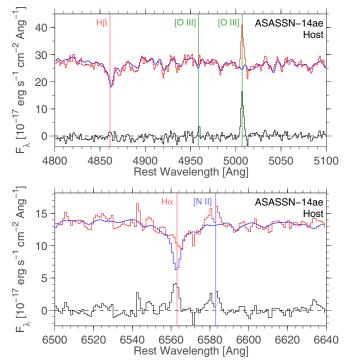


Figure 16. Same as Figure 9, but for the host galaxy of ASASSN-14ae. The $[O_{III}]/H\beta$ ratio and the $[O_{III}]$ luminosity are consistent with a weak AGN. (A color version of this figure is available in the online journal.)

For main-sequence stars, the stellar mass and radius nearly cancel, but the R_T/R_P factor here may smear any remaining correlation.

Strubbe & Quataert (2009) note that lines from the outflowing material alone will be bulk blueshifted or redshifted. It is therefore possible that a broadening of the lines would be caused by a combination of emission from bound and the outflowing material.

We fit a Gaussian function to each broad emission feature in continuum-subtracted spectra of the TDE candidates presented in Figure 14 and use it to estimate the velocity width of these lines. We change the continuum subtraction parameters and take the scatter in the fitted width as its error. Our measured 1σ line widths are presented in Table 7 (our values for the line widths of TDE2 and PS1-10jh are consisted with those reported by van Velzen et al. 2011 and G12, respectively, though they quote FWHM values while we use 1σ).

Table 7Measured 1σ Line Widths of the H or He Emission Feature

Object	Line	Phase (days)	Line Width (km s ⁻¹)	
PS1-10jh	Не II 4686 Å	-22	5430 ± 1460	
PTF09ge	Не II 4686 Å	-19	10070 ± 670	
SDSS J0748	Не п 4686 Å	n/a	9950 ± 510	
ASASSN-14ae	Ηα	n/a	3600 ± 175	
PTF09axc	Ηα	7	11890 ± 220	
PTF09djl	Ηα	2-62	6530 ± 350	
TDE2	Ηα	n/a	3440 ± 1100	

Notes. Phases are shown relative to peak brightness. For PTF09djl, the average velocity from the three spectra is shown, as we see no evolution in the width of the main component within the errors.

The velocities are lower than the expected v_T (at the tidal radius) and correspond to bound (circular) orbits at distances of ~20–80 $(M_{BH}/10^6 M_{\odot})^{2/3} (\rho_*/\rho_{\odot})^{1/3}$ tidal radii from the SMBH, or to outflowing unbound trajectories.

We plot the observed line velocities versus host galaxy magnitudes and versus derived SMBH masses (Figure 17, top and bottom panels, respectively). For the top panel, we use K-corrected (Chilingarian et al. 2010)³² host galaxy gband magnitudes. We translate these magnitudes to SMBH mass using the Bernardi et al. (2003) coefficients for the Faber–Jackson relation and assuming the $M - \sigma$ relation found by Kormendy & Ho (2013). For the bottom panel, we calculate the SMBH masses from our derived galaxy stellar masses using the Gadotti (2009) and Häring & Rix (2004) relations (see Section 5.3). In both panels, we plot the expected correlations for bound material at several times the tidal radius for a Sun-like disrupted star and for unbound material with $R_T = R_P$ (Strubbe & Quataert 2009). The overall scale of the velocities can be seen to be consistent with large bound radii or unbound velocities (as stated by G12 for PS1-10jh).

There does not seem to be a robust correlation of the line widths with either host galaxy magnitude or SMBH mass, encompassing all events. This suggests that the simple association of velocities to bound circular Keplerian orbits is likely incorrect. Additional outflowing components, non-circular orbits, and variability in the stellar properties could all smear out the correlation.

³² Obtained using g - r colors through the "*K*-corrections calculator" at http://kcor.sai.msu.ru/.

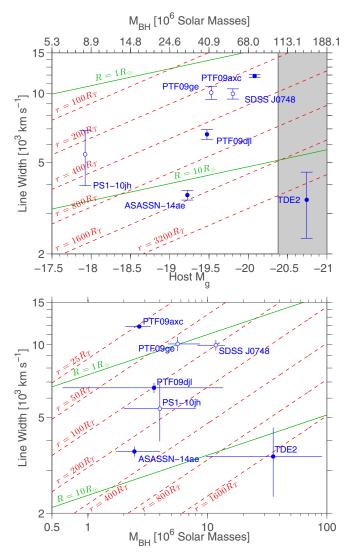


Figure 17. TDE-candidate line widths of H (filled circles) and He (empty circles) emission lines vs. host galaxy K-corrected *g*-band absolute magnitudes (top) and SED-derived SMBH masses (bottom). The dashed red lines represent the expected velocity correlations for bound material at different radii (assuming a Sun-like star). The solid green lines are the Strubbe & Quataett (2009) velocities for the most energetic outflowing material assuming $R_P = R_T$ for a Sun-like star or red giant as noted. For the top panel, we use the Bernardi et al. (2003) coefficients for the Faber–Jackson relation to derive these lines and to derive the top x-axis values from the host galaxy *g*-band absolute magnitudes. For the bottom panel, the SMBH masses are taken from Table 6 (i.e., calculated using the Gadotti (2009) and Häring & Rix (2004) relations from our derived host stellar masses). The shaded region in the top panel denotes SMBH masses above $10^8 M_{\odot}$, for which a TDE is not expected to be observed (assuming a Sun-like star and a non-rotating SMBH).

5.5. The Double-peak Ha Profile of PTF09djl

For PTF09djl, we find an additional redshifted component to the H α emission feature extending out to high velocities (Figure 18). The observed structure is reminiscent of the doublepeaked line profiles usually explained by Keplerian disk models (e.g., Chen et al. 1989). We construct a circular disk model following Strateva et al. (2003), which reproduces the shape but not the location of the profile (Figure 18). While the model emission peaks are symmetric around the rest wavelength, in PTF09djl, one peak is at the rest wavelength while the second is redshifted. We therefore have to shift the model profile to fit

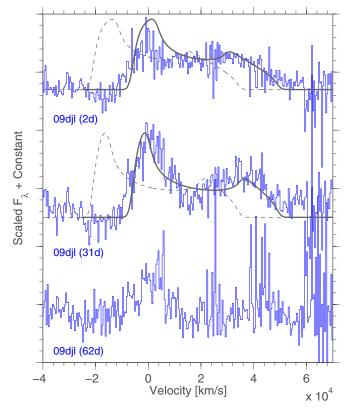


Figure 18. H α profile in the spectra of PTF09djl after host and continuum subtraction, corrected to the host redshift determined by its narrow absorption features. The data are binned to improve signal to noise, and phases are shown relative to peak. A red tail (first spectrum) and distinct red component (second spectrum) can be seen. A circular Keplerian disk model, following Strateva et al. (2003), is shown for inner and outer radii of 70 and 300 Schwarzschild radii (first spectrum) and 70 and 170 Schwarzschild radii (second spectrum), respectively. Both models assume an emissivity index of -3, a local turbulent broadening of 1200 km s⁻¹, and an inclination angle of 50 degrees. The models are shifted to the red by 15, 000 km s⁻¹ (solid lines) from their original position (dashed lines).

(A color version of this figure is available in the online journal.)

it to the observed spectrum. The disk model would thus have to include a bulk motion component to explain these observations.

Another possibility is that the geometry is more complex than a disk, with the emission related to the unbound debris, such as in the models considered by Bogdanović et al. (2004). Their models reproduce the general profile shape but with much lower velocities (though the velocities can be increased using smaller inner radii; T. Bogdanovic 2014, private communication).

The double-peaked structure is not observed in the H β line nor is it seen in the other events (though this could be due to a low inclination angle for the disk model there).

5.6. Radio Non-detections of PTF09axc

Our VLA observations of PTF09axc resulted in a null detection at both the 3.5 GHz and 6.1 GHz with an rms of 110 μ Jy and 50 μ Jy, respectively. At the redshift of PTF09axc, these limits imply 3 σ luminosity upper limits of 1.2 × 10²⁹ erg Hz⁻¹ s⁻¹ and 5.3 × 10²⁸ erg Hz⁻¹ s⁻¹ at these bands.

So far, the only two TDE candidates ever to be detected in the radio are those found in γ -rays by *Swift*: *Swift* J1644+57 (Bloom et al. 2011; Burrows et al. 2011; Levan et al. 2011; Zauderer et al. 2011) and *Swift* J2058+05 (Cenko et al. 2012). The high energy emission of these events was suggested to originate in

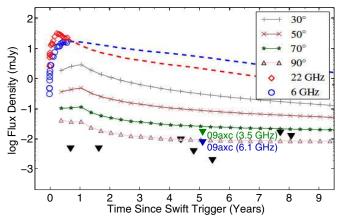


Figure 19. Radio non-detection limits of PTF09axc, five years after the eruption, overplotted on Figure 1 from van Velzen et al. (2013). The open symbols show the observed radio light curve of *Swift* J1644+57. The dashed lines are predicted light curves from Berger et al. (2012) assuming a jet energy of 10^{52} erg. The van Velzen et al. (2013) 5 GHz off-axis light curve predictions are also shown. The triangles denote 2σ upper limits on the radio flux of the van Velzen et al. (2013) sample (black) and our data (green and blue). All data and models are scaled to the redshift of *Swift* J1644+57 (z = 0.354). An off-axis jet is very unlikely for PTF09axc.

a relativistic jet pointed in our direction. The radio emission of such a jet should be observable also in off-axis cases.

Recently, van Velzen et al. (2013) observed seven TDE candidates (including PS1-10jh and TDE2) in the radio and did not detect emission for any of them. The non-detection limits were deep enough to rule out most off-axis jet angles for these events and led van Velzen et al. (2013) to conclude that it is not likely that all TDEs launch jets. Our limits for PTF09axc strengthen this conclusion (Figure 19).

6. SUMMARY AND CONCLUSIONS

We presented the results of a search for blue transients with peak magnitudes between those of SNe and SLSNe in the PTF non-interacting core-collapse SN sample. Of the six events found, we focus on PTF09ge, PTF09axc, and PTF09djl, the three that are coincident with the centers of their host galaxies. These events, selected amongst core-collapse SN candidates only by their peak magnitude, color, centrality in their hosts, and lack of narrow emission lines indicative of interaction-powered emission, turn out to be similar photometrically to the TDE candidate PS1-10jh (G12) and are all found to reside in E+A (or E+A-like) galaxies.

The lack of obvious AGN features in the spectra of the outbursts and the host galaxies as well as the magnitude of the outbursts and their non-recurrence all disfavor an AGN interpretation for these events. The lack of recent or ongoing star formation in all the hosts rules out core-collapse SNe, which originate in massive, short-lived stars. In addition, the blue spectra imply higher temperatures than those typically observed in SNe. We conclude that these events are likely TDEs.

One of the events is similar spectroscopically to the Herich TDE candidate PS1-10jh, while the other two show broad hydrogen emission features, as expected by some models for TDE's involving a H-rich star. A spectrum of the recent TDE candidate ASASSN-14ae displays *both* H and He emission features. Comparing these events to literature data of two additional TDE candidates, we find a sequence of H- to Hedominated spectral features. This indicates that either a viewingangle effect or some continuous parameter(s) of the disrupted star, its orbit, or the SMBH may generate a continuum of TDE spectral types.

Van Velzen et al. (2011) calculate that PTF could potentially detect \sim 13 TDE flares per year. However, there exists a strong bias against spectroscopically following transients in centers of galaxies by PTF or any transient survey that is wary of contamination by AGNs and subtraction artifacts around galactic cores. Additional biases inherent to the survey (such as alternating observing strategies and follow up prioritizations) make it impossible at this time to quantify the true expected rate of spectroscopically observed TDEs in PTF.

The preference of these events for E+A hosts is intriguing. If indeed E+As are post-merger galaxies (Zabludoff et al. 1996), then their TDE rates could be enhanced by the interaction of two central SMBHs. Chen et al. (2009) find that a binary SMBH can increase the TDE rate by several orders of magnitude, even up to \sim one event per year (though only for a relatively short period). Some E+A galaxies may be the result of interaction rather than mergers (Goto 2005; Yamauchi et al. 2008). In this case, the TDE rate may be enhanced by dynamical interactions perturbing the orbits of stars near the SMBH. Alternatively, a unique stellar population in the cores of E+A galaxies may influence the TDE rate. A large fraction of evolved extended stars, for example, would increase the rate of TDEs since lower density stars are more easily disrupted (such stars may also experience multiple partial disruptions; MacLeod et al. 2013). We note, however, that our light curves are not as slowly evolving as those expected for the tidal disruptions of evolved stars (MacLeod et al. 2012).

We note that when excluding SNe classified as Type IIn, we may have introduced a bias against TDEs in star-forming hosts (where the narrow emission lines from the host would have lead to the misclassification of the transient as a IIn SN). However, the line profiles of IIn SNe are different than those of galaxies (e.g., Kiewe et al. 2012), and given that we use both visual and well-established SN spectral fitting tools to identify them and that half the events that made the cut into our initial sample are in star-forming hosts, we conclude that any bias against such hosts is small, if at all present. Even if such a bias were strong, it would introduce an over-representation of passive galaxies in general and not exclusively rare E+As as we see here.

We detect X-ray flux at the position of PTF09axc, which is marginally consistent with an extremely weak AGN given our measured [O III] luminosity (Heckman et al. 2005) and highly inconsistent with an accreting binary origin (Hornschemeier et al. 2005). The high X-ray luminosity and low mass of the host galaxy are very similar to that of X171206.83+640830.7 (Hornschemeier et al. 2005). Curiously, this source may be variable and its host galaxy is also a post-merger (Davis et al. 2003). If some TDEs can emit X-rays at late time, it may be that X171206.83+640830.7 was also a TDE in an E+A galaxy.

One of the events in our PTF sample (PTF10iam) is clearly offset from the center of its host and its light curve and spectral behavior are different than those of the central events. Its host galaxy is also different, showing signs of star formation. We discuss this transient as a possible interacting SN (displaying broad high velocity H α absorption rather than the typical narrow emission lines) in a companion paper (I. Arcavi et al., in preparation).

The combined sample of nuclear transients presented here, tied together by the PTF events, now strongly supports a TDE origin for all of these objects, spanning a continuum of spectral classes and preferring low star-forming E+A-like host galaxies.

LSST will discover thousands of TDEs per year (van Velzen et al. 2011), but UV observations are required to accurately constrain the temperature and energetics of these transients. The proposed *ULTRASAT* mission could discover hundreds of TDEs per year in the UV (Sagiv et al. 2014). Understanding how to interpret TDE observations will enable them to be used to study accretion physics, stellar populations, and otherwise-quiescent black holes.

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