A framework based on 2-D Taylor expansion for quantifying the impacts of sub-pixel reflectance variance and covariance on cloud optical thickness and effective radius retrievals based on the bi-spectral method

Z. Zhang^{1,2,*}, F. Werner², H.-M. Cho^{2,3}, G. Wind,^{4,5} S. Platnick⁴, A. S. Ackerman⁶, L. Di Girolamo⁷, and A. Marshak⁴, Kerry Meyer^{4,8}

- 1. Physics Department, UMBC, Baltimore, MD, USA
- 2. Joint Center for Earth Systems Technology, UMBC, Baltimore, MD, USA
- 3. Electronics and Telecommunications Research Institute, Korea
- 4. NASA Goddard Space Flight Center, Greenbelt, MD, USA
- 5. Science Systems and Applications, Inc., Maryland 20706, USA
- 6. NASA Goddard Institute for Space Studies, New York City, NY, USA
- 7. Department of Atmospheric Sciences, University of Illinois, Urbana-Champaign, IL, USA
- 8. Universities Space Research Association (USRA), Columbia, MD, USA

For Publication in JGR-Atmosphere

1 Abstract:

2 The bi-spectral method retrieves cloud optical thickness (τ) and cloud droplet effective radius (r_e) simultaneously from a pair of cloud reflectance observations, one in 3 a visible or near infrared (VIS/NIR) band and the other in a shortwave-infrared (SWIR) 4 5 band. A cloudy pixel is usually assumed to be horizontally homogeneous in the retrieval. 6 Ignoring sub-pixel variations of cloud reflectances can lead to a significant bias in the retrieved τ and r_e . In the literature, the retrievals of τ and r_e are often assumed to be 7 8 independent and considered separately when investigating the impact of sub-pixel 9 cloud reflectance variations on the bi-spectral method. As a result, the impact on τ is 10 contributed only by the sub-pixel variation of VIS/NIR band reflectance and the impact on r_e only by the sub-pixel variation of SWIR band reflectance. 11

12 In our new framework, we use the Taylor expansion of a two-variable function to 13 understand and quantify the impacts of sub-pixel variances of VIS/NIR and SWIR cloud reflectances and their covariance on the au and r_{e} retrievals. This framework takes into 14 15 account the fact that the retrievals are determined by both VIS/NIR and SWIR band 16 observations in a mutually dependent way. In comparison with previous studies, it 17 provides a more comprehensive understanding of how sub-pixel cloud reflectance 18 variations impact the τ and r_{e} retrievals based on the bi-spectral method. In particular, 19 our framework provides a mathematical explanation of how the sub-pixel variation in 20 VIS/NIR band influences the r_e retrieval and why it can sometimes outweigh the influence of variations in the SWIR band and dominate the error in r_e retrievals, leading 21 to a potential contribution of positive bias to the r_e retrieval. We test our framework 22 23 using synthetic cloud fields from a large-eddy simulation and real observations from 24 MODIS. The predicted results based on our framework agree very well with the 25 numerical simulations. Our framework can be used to estimate the retrieval uncertainty 26 from sub-pixel reflectance variations in operational satellite cloud products and to help understand the differences in τ and r_e retrievals between two instruments. 27

28

30 1. Introduction

31 Among many satellite-based cloud remote sensing techniques, the bi-spectral 32 solar reflective method ("bi-spectral method" hereafter) is a widely used method to infer cloud optical thickness (τ) and cloud droplet effective radius (r_e) from satellite 33 observation of cloud reflectance [Nakajima and King, 1990]. This method uses cloud 34 35 reflectance measurements from two spectral bands to simultaneously retrieve τ and 36 r_e . One measurement is usually made in the visible or near-infrared (VIS/NIR) spectral 37 region (e.g., 0.64 μ m or 0.86 μ m), where water absorption is negligible and therefore 38 cloud reflection generally increases with τ . The other measurement is usually in the 39 shortwave infrared (SWIR) spectral region (e.g., 2.1 µm or 3.7 µm), where water drops 40 are moderately absorptive and cloud reflectance generally decreases with increasing r_e for optically thick clouds. In practice, the bi-spectral method is often implemented 41 42 utilizing the so-called look-up-table (LUT). A couple of LUT examples are shown in Figure 43 1. Such LUTs contain pre-computed bi-directional cloud reflectances at VIS/NIR and SWIR bands for various combinations of $r_{\!_e}$ and au under different sun-satellite viewing 44 45 geometries and surface reflectances. Given the observed reflectances, the corresponding r_e and τ can be retrieved easily by searching and interpolating the 46 47 proper LUT. The bi-spectral method has been adopted by a number of satellite missions, 48 including Moderate Resolution Imaging Spectroradiometer (MODIS), Visible Infrared Imaging Radiometer Suite (VIIRS), and Spinning Enhanced Visible and Infrared Imager 49 (SEVIRI) for operational retrievals of cloud properties (i.e., τ , r_{e} and derived cloud liquid 50

water path (LWP)) [*Platnick et al.*, 2003; *Roebeling et al.*, 2006; *Minnis et al.*, 2011; *Walther and Heidinger*, 2012]. Given the wide usage of the bi-spectral method, it is
critical to study and understand its limitations and uncertainties.

54 The bi-spectral method makes several important assumptions about the cloud 55 (or cloudy pixels). First, within a cloudy pixel, cloud is assumed to be horizontally 56 homogenous (referred to as the "homogenous pixel assumption"). Second, it is assumed 57 that the pixels are independent from each other, in the sense that there is no net interpixel transport of radiation (often referred to as the "independent pixel assumption, 58 59 IPA"). Under these assumptions, clouds are considered to be "plane-parallel". In 60 addition to plane-parallel cloud assumptions, clouds are often assumed to be vertically 61 homogenous in the operational algorithms. Furthermore, the size spectrum of cloud 62 particles is often assumed to follow certain analytical distributions, such as the single modal gamma or lognormal size distributions [e.g., Nakajima and King, 1990; Dong et 63 al., 1997]. These assumptions may be reasonable for certain types of clouds, such as 64 65 closed-cell, non-precipitating stratocumulus, but become problematic for others, such as broken trade-wind cumuli or precipitating clouds [Di Girolamo et al., 2010; Painemal 66 67 and Zuidema, 2011; Zhang and Platnick, 2011; Liang and Girolamo, 2013; Zhang, 2013]. 68 As elucidated in numerous previous studies, when real clouds deviate from these 69 assumptions, the r_{e} and au retrievals from the bi-spectral method can suffer from large errors and uncertainties [e.g., Várnai and Marshak, 2002; Kato et al., 2006; Marshak et 70 71 al., 2006; Zhang and Platnick, 2011; Zhang et al., 2012; Zhang, 2013; Liang et al., 2015].

The focus of this study is the homogenous pixel assumption. Our objective is to 72 73 develop a unified framework for understanding and quantifying the impacts of sub-pixel level unresolved reflectance variations on r_e and au retrievals based on the bi-spectral 74 75 method. A number of previous studies have already made substantial progress in this 76 direction. It has been known for a long time that at the spatial scale of climate model grids (e.g., $\sim 10^2$ km) approximating inhomogeneous cloud fields with plane-parallel 77 clouds can lead to significant biases in shortwave solar radiation [e.g., Harshvardhan 78 79 and Randall, 1985; Cahalan et al., 1994; Barker, 1996]. Cahalan et al. [1994] described an elegant theoretical framework based on a fractal cloud model to explain the 80 81 influence of small-scale horizontal variability of τ on the averaged cloud reflectance in the visible spectral region $R_{_{V\!I\!S}}$. It is shown that the averaged reflectance $\overline{R_{_{V\!I\!S}}(au_i)}$, 82 where au_i denotes the sub-pixel scale cloud optical thickness, is smaller than the 83 reflectance that corresponds to the averaged cloud optical thickness $\overline{ au_i}$, i.e., 84 $\overline{R_{_{V\!I\!S}}(au_i)} < R_{_{V\!I\!S}}(\overline{ au_i})$. This inequality relation is well known as the "plane-parallel 85 homogenous bias" (referred to as PPHB), which is a result of the non-linear dependence 86 of R_{VIS} on τ i.e., $\frac{\partial^2 R_{vis}}{\partial \tau^2} < 0$. The implication of the PPHB for τ retrievals from R_{VIS} is 87 illustrated using an example shown in Figure 2a. Here, we assume that one half of an 88 inhomogeneous pixel is covered by a thinner cloud with $\tau_{_1}\!=\!5$ and the other half by a 89 thicker cloud with au_2 = 18 (both clouds with r_e = $8\,\mu m$). Because of the PPHB, the 90 retrieved cloud optical thickness $\tau^* = 9.8$ based on the averaged reflectance 91

92 $\overline{R} = [R(\tau_1) + R(\tau_2)]/2$ is significantly smaller than the linear average of the sub-pixel τ , 93 i.e., $\overline{\tau} = 11.5$. The impacts of PPHB on satellite based cloud property retrievals and the 94 implications have been investigated in a number of studies [*Oreopoulos and Davies*, 95 1998; *Pincus et al.*, 1999; *Oreopoulos et al.*, 2007].

96 We note that the variation of cloud reflectance may be a result of varying cloud 97 properties, but may also be caused by 3-D radiative effects. For example, a cloudy pixel 98 can be perfectly homogenous in terms of cloud properties, but the surrounding pixels 99 can cast a shadow on part of this pixel leading to sub-pixel reflectance variation 100 [Marshak et al., 2006]. A variety of such 3-D effects that cannot be explained by the 1-D 101 plane-parallel radiative transfer theory have been identified and their impacts on cloud 102 property retrievals investigated in previous studies [Davis and Marshak, 2010]. In 103 reality, the PPHB is inevitably entangled with the 3-D transfer effects and other 104 uncertainties such as the impact of instrument noise in the retrieval. It is difficult, if not 105 impossible, to separate them. Following the literature, we shall refer to the impact of 106 sub-pixel cloud reflectance variation on cloud property retrievals as the PPHB, while 107 keeping in mind that the sub-pixel cloud reflectance variation can also result from 3-D 108 radiative effects and may not reflect the true variation of sub-pixel cloud properties.

109 Recently, as the interests in aerosol-cloud interactions have grown, there is an 110 increasing attention on the impacts of small-scale cloud variations on the satellite-based 111 r_e retrievals [e.g., *Kato et al.*, 2006; *Marshak et al.*, 2006; *Zhang and Platnick*, 2011; 112 *Zhang et al.*, 2012; *Liang et al.*, 2015]. Marshak et al. [2006] pointed out that similar to 113 the PPHB the non-linear dependence of the SWIR band cloud reflectance R_{SWIR} on r_e

114 can also lead to significant biases on r_e retrievals, which is demonstrated in Figure 2b. Here, one half of an inhomogeneous pixel is covered by a cloud with $r_e = 8 \mu m$ and the 115 other half by a cloud with $r_e = 22 \mu m$. Both parts have the same $\tau = 4.1$. As shown in the 116 figure, the retrieved $r_e^* = 12 \mu m$ based on the averaged reflectance is significantly 117 smaller than the linear average of sub-pixel $\overline{r_e} = 15 \mu m$, similar to the PPHB of τ in 118 Figure 2a. It must be noted that in the framework of Marshak et al. [2006] the retrievals 119 of $r_{\scriptscriptstyle e}$ and au are considered separately and assumed to be independent from one 120 121 another. However, as Marshak et al. [2006] pointed out this assumption is valid only for "large enough" τ and r_e (typically, $r_e > 5 \ \mu m$ and $\tau > 10$). As one can see from the 122 shape of the LUT in Figure 1 the $R_{\scriptscriptstyle SWIR}$ is not completely orthogonal to the $R_{\scriptscriptstyle VIS}$, 123 especially when au is small. As a result, the retrievals of r_e and au are not independent 124 from one another. Marshak et al. [2006] suspected that some cases with large r_e bias in 125 their simulations might be the result of this mutual dependence of $r_{\!_e}$ and au retrievals. 126 127 Zhang and Platnick [2011] showed that the sub-pixel variance of τ can have a significant impact on the r_e retrieval, which is illustrated in the example in Figure 2c. In this 128 hypothetical case, an inhomogeneous pixel is assumed to be covered by a thinner cloud 129 with τ_1 =6 in one half and a thicker cloud with τ_2 =18 in the other. Both clouds have the 130 131 same r_e =14 µm. Note that in this case the sub-pixel reflectance variation is solely caused 132 by the variability in τ . If the r_e retrieval were independent from the τ retrieval, then the retrieved r_e would be 14 μ m. The solid triangle in the figure indicates the location of the 133 $R_{\scriptscriptstyle VIS}$ and $R_{\scriptscriptstyle SWIR}$ averaged over the pixel, i.e., the "observation". The retrieved $au^*=10.8$ 134

is smaller than the averaged $\overline{\tau} = 12$ as a result of the PPHB. However, the retrieved 135 $r_{e}^{*} = 16$ is 2 µm larger than the expected value of 14 µm. This positive bias in the r_{e} 136 retrieval, apparently caused by the sub-pixel variability of τ , cannot be explained by the 137 framework of Marshak et al. [2006] in which the r_e retrieval is assumed to be 138 independent from the τ retrieval. Zhang and Platnick [2011] and Zhang et al. [2012] also 139 found that the magnitude of the positive r_e retrieval bias caused by the sub-pixel 140 variability of τ is dependent on the SWIR band chosen for the r_e retrieval. These studies 141 142 showed that the same sub-pixel τ variability tends to induce larger bias in retrieved r_e using the less absorptive 2.1 μ m band (referred to as $r_{e,2,1}$) than that using the more 143 absorptive 3.7 μ m band (referred to as $r_{e,3,7}$). This spectral dependence provides an 144 important explanation for the fact that the MODIS operational $r_{e2,1}$ retrievals for water 145 clouds are often significantly larger than the $r_{e,3,7}$ retrievals, especially when clouds have 146 large sub-pixel heterogeneity [Zhang and Platnick, 2011; Cho et al., 2015]. 147

The aforementioned studies have undoubtedly shed important light on the 148 impact of sub-pixel cloud variability on $r_{\!_e}$ and au retrievals based on the bi-spectral 149 150 method. However, several questions still remain. For example, an important question is 151 how to reconcile the negative r_e bias discussed in Marshak et al. [Marshak et al., 2006] and the positive r_e bias discussed in Zhang and Platnick [2011] and Zhang et al. [2012]. 152 Indeed, this is the main question we will address in this study. In the light of previous 153 154 studies, here we develop a new mathematical framework to provide a more comprehensive and complete understanding of the impact of sub-pixel cloud variability 155

156 on r_e and τ retrievals based on the bi-spectral method. The paper is organized as 157 follows: We formulate the problem in Section 2. We introduce our mathematical 158 framework in Section 3, test and validate it using two examples in Section 4, and discuss 159 its applications in Section 5.

160 2. Statement of the problem

161 In the bi-spectral method, r_e and τ are retrieved from a pair of cloud reflectance 162 observations, one in VIS/NIR and the other in SWIR. From this point of view, we can 163 define r_e and τ as:

164
$$\tau \equiv \tau \left(R_{VIS}, R_{SWIR} \right)$$
$$r_e \equiv r_e \left(R_{VIS}, R_{SWIR} \right)'$$
(1)

where R_{VIS} and R_{SWIR} are the observed reflectances in the VIS/NIR (denoted by 165 subscript "VIS" for short) and SWIR bands, respectively. Assume that an instrument with 166 167 a relatively coarse spatial resolution observes a horizontally inhomogeneous cloudy pixel in its field of view. The observed cloud reflectances are $\overline{R_{_{VIS}}}$ and $\overline{R_{_{SWIR}}}$, where the 168 169 overbar denotes the spatial average. Now if we use another instrument with a finer spatial resolution to observe the same area covered by the coarser resolution pixel, we 170 can obtain high-resolution observations, $R_{VIS,i}$ and $R_{SWIR,i}$, i = 1, 2, ...N, (the number N 171 172 depends on the relative sizes of the pixels). The high-resolution measurements provide the information on the variance and covariance of R_{VIS} and R_{SWIR} at sub-pixel scale. 173 Each sub-pixel observation $R_{VIS,i}$ and $R_{SWIR,i}$ can be specified as the deviation from the 174 mean value $\overline{R_{VIS}}$ and $\overline{R_{SWIR}}$ as: 175

176

$$R_{VIS,i} = \overline{R_{VIS}} + \Delta R_{VIS,i}$$

$$R_{SWIR,i} = \overline{R_{SWIR}} + \Delta R_{SWIR,i}$$
; $i = 1, 2...N$. (2)

177 It naturally follows that the spatial average $\overline{\Delta R_{VIS,i}} = \overline{\Delta R_{SWIR,i}} = 0$. Based on the coarse-178 resolution reflectance observations $\overline{R_{VIS}}$ and $\overline{R_{SWIR}}$, we can retrieve $\tau(\overline{R_{VIS}}, \overline{R_{SWIR}})$ and 179 $r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})$. From the high-resolution, sub-pixel observations $R_{VIS,i}$ and $R_{SWIR,i}$, we can 180 retrieve $\tau(R_{VIS,i}, R_{SWIR,i})$ and $r_e(R_{VIS,i}, R_{SWIR,i})$. The differences $\Delta \tau$ and Δr_e , defined as:

181
$$\Delta \tau = \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right) - \overline{\tau \left(R_{VIS,i}, R_{SWIR,i} \right)} \Delta r_e = r_e \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right) - \overline{r_e \left(R_{VIS,i}, R_{SWIR,i} \right)},$$
(3)

are considered in this, as well as previous studies, as the biases caused by the homogeneous pixel assumption in r_e and τ retrievals [*Cahalan and Joseph*, 1989; *Marshak et al.*, 2006; *Zhang et al.*, 2012].

185 Consideration of eq. (3) raises a few important questions. What are the sign and magnitude of $\Delta \tau$ and Δr_e ? How do they depend on the sub-pixel $R_{VIS,i}$ and $R_{SWIR,i}$? 186 Addressing these questions could help improve understanding of the biases caused by 187 ignoring the sub-pixel reflectance variation in bi-spectral r_e and τ retrievals. 188 Furthermore, since performing high-resolution retrievals can be computationally 189 expensive, another important question is whether it is possible to estimate 190 $\overline{\tau(R_{_{VIS,i}},R_{_{SWIR,i}})}$ and $\overline{r_{_e}(R_{_{VIS,i}},R_{_{SWIR,i}})}$ from the coarse-resolution retrievals and the 191 192 statistics of sub-pixel reflectance observations, even without doing time-consuming 193 high-resolution retrievals. If this proved possible, then it is a very efficient way to estimate the biases and uncertainty caused by the homogenous pixel assumption. Thesequestions are the focus of this study and will be addressed in the next section.

Before proceeding, we need to clarify two points. First, the $\Delta \tau$ and Δr_e in Eq. (3) 196 are the differences between two sets of retrievals, not the differences between the 197 198 retrievals and "true" cloud properties. As aforementioned, sub-pixel reflectance 199 variations can be due to sub-pixel scale cloud property variation, but may also be caused 200 by 3-D radiative effects. If the former is dominant, then $\Delta \tau$ and Δr_e provide an 201 estimate of the PPHB and can be used to correct the coarse-resolution retrievals to better represent the "true" cloud properties. However, if 3-D effects are the dominant 202 cause of the sub-pixel reflectance variation, then Δau and $\Delta r_{\!_e}$ can be considered a 203 204 quantitative index of the 3-D effects on the retrievals. Second, our scope is to study the connections between retrieval biases $\Delta \tau$ and Δr_e with sub-pixel observations $R_{_{VIS,i}}$ 205 206 and $R_{SWIR,i}$. We simply take $R_{VIS,i}$ and $R_{SWIR,i}$ as given inputs. Here we do not seek to explain the characteristics of $R_{VIS,i}$ and $R_{SWIR,i}$ (e.g., their mean values, variances and 207 covariance), or their dependence on cloud properties. Neither do we try to explain how 208 the 3-D radiative effects and instrument characteristics influence $R_{_{V\!I\!S,i}}$ and $R_{_{S\!W\!I\!R,i}}$. 209

210

3. A unified mathematical framework

In this section, we will introduce a comprehensive framework that is able to reconcile and unify the theoretical understandings provided by Marshak et al. [2006], Zhang and Platnick [2011], and Zhang et al. [2012] To investigate the sign and 215 magnitude of $\Delta \tau$ and Δr_e , we first expand the $\tau(R_{VIS,i}, R_{SWIR,i})$ and $r_e(R_{VIS,i}, R_{SWIR,i})$ into 216 two-dimensional Taylor series of $R_{VIS,i}$ and $R_{SWIR,i}$. Take $r_e(R_{VIS,i}, R_{SWIR,i})$ for example. 217 The expansion is:

$$r_{e}\left(R_{_{VIS,j}},R_{_{SWIR,j}}\right) = r_{e}\left(\overline{R_{_{VIS}}} + \Delta R_{_{VIS,j}},\overline{R_{_{SWIR}}} + \Delta R_{_{SWIR,j}}\right)$$

$$= r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right) + \frac{\partial r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right)}{\partial R_{_{VIS,j}}}\Delta R_{_{VIS,j}} + \frac{\partial r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right)}{\partial R_{_{SWIR,j}}}\Delta R_{_{SWIR,j}} + \frac{1}{2}\frac{\partial^{2} r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right)}{\partial R_{_{SWIR,j}}}\Delta R_{_{SWIR,j}} + \varepsilon$$

$$\frac{1}{2}\frac{\partial^{2} r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right)}{\partial R_{_{VIS},j}}\Delta R_{_{SWIR,j}}}\Delta R_{_{SWIR,j}} + \frac{1}{2}\frac{\partial^{2} r_{e}\left(\overline{R_{_{VIS}}},\overline{R_{_{SWIR}}}\right)}{\partial R_{_{SWIR,j}}}\Delta R_{_{SWIR,j}} + \varepsilon$$

219 where ε is the truncation error if higher order derivative terms are neglected. If we 220 take the spatial average of Eq. (4) and neglect ε , all the linear terms (i.e., $\frac{\partial r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})}{\partial R_e(\overline{R_{VIS}}, \overline{R_{SWIR}})} \neq R_{e_s}$) we risk because $\overline{AR_{e_s}} = \overline{AR_{e_s}} = 0$. Thus

221
$$\frac{\partial r_e(R_{VIS}, R_{SWIR})}{\partial R_{VIS}} \Delta R_{VIS,i} \text{ and } \frac{\partial r_e(R_{VIS}, R_{SWIR})}{\partial R_{SWIR}} \Delta R_{SWIR,i} \text{) vanish because } \overline{\Delta R_{VIS,i}} = \overline{\Delta R_{SWIR,i}} = 0. \text{ Thus}$$

only second order terms in Eq. (4) remain after the spatial average:

223
$$\overline{r_e(R_{_{VIS,i}},R_{_{SWIR,i}})} \approx r_e(\overline{R_{_{VIS}},R_{_{SWIR}}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{VIS}}^2} \sigma_{_{VIS}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{VIS}}\partial R_{_{SWIR}}} \operatorname{cov}(R_{_{VIS}},R_{_{SWIR}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{VIS}}\partial R_{_{SWIR}}} \operatorname{cov}(R_{_{VIS}},R_{_{SWIR}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{VIS}}\partial R_{_{SWIR}}} \operatorname{cov}(R_{_{VIS}},R_{_{SWIR}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{VIS}}\partial R_{_{SWIR}}} \operatorname{cov}(R_{_{VIS}},R_{_{SWIR}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \operatorname{cov}(R_{_{VIS}},R_{_{SWIR}}) + \frac{1}{2} \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2} + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}}^2} + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}^2}} \sigma_{_{SWIR}}^2 + \frac{\partial^2 r_e(\overline{R_{_{VIS}},R_{_{SWIR}}})}{\partial R_{_{SWIR}^2}} + \frac{\partial^2 r_e(\overline{R_{_{SWIR}},R_{_{SWIR}})}{\partial R_{_{SWIR}^2}} + \frac{\partial^2 r_e(\overline{R_{_{SWIR},R_{_{SWIR}},R_{_{SWIR}})}{\partial$$

(5)

224

218

where $\sigma_{VIS}^2 = \overline{\Delta R_{VIS,i}^2}$, $\sigma_{SWIR}^2 = \overline{\Delta R_{SWIR,i}^2}$ are the spatial variances of $R_{VIS,i}$ and $R_{SWIR,i}$, respectively, and $\operatorname{cov}(R_{VIS}, R_{SWIR})$ is the spatial covariance of $R_{VIS,i}$ and $R_{SWIR,i}$. Substituting Eq. (5) into Eq. (3), we obtain the following formula for Δr_e :

$$\Delta r_{e} = r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right) - r_{e} \left(R_{VIS,i}, R_{SWIR,i}\right)$$

$$= -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} \sigma_{VIS}^{2} - \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}} \operatorname{cov}\left(R_{VIS}, R_{SWIR}\right) - \frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \sigma_{SWIR}^{2}$$

(6)

(7)

(8)

229

230 Following the same procedure, we can derive the formula for $\Delta \tau$ as:

231
$$\Delta \tau = \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right) - \overline{\tau \left(R_{VIS,i}, R_{SWIR,i} \right)} = -\frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right)}{\partial R_{VIS}^2} \sigma_{VIS}^2 - \frac{\partial^2 \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right)}{\partial R_{VIS} \partial R_{SWIR}} \operatorname{cov} \left(R_{VIS}, R_{SWIR} \right) - \frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}} \right)}{\partial R_{SWIR}^2} \sigma_{SWIR}^2.$$

232

Eq. (6) and (7) can be combined into a matrix form as follows:

$$234 \qquad \left(\begin{array}{c}\Delta\tau\\\Delta r_{e}\end{array}\right) = \left(\begin{array}{c}-\frac{1}{2}\frac{\partial^{2}\tau\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} & -\frac{\partial^{2}\tau\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{VIS}\partial R_{SWIR}} & -\frac{1}{2}\frac{\partial^{2}\tau\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \\ -\frac{1}{2}\frac{\partial^{2}r_{e}\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} & -\frac{\partial^{2}r_{e}\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{VIS}\partial R_{SWIR}} & -\frac{1}{2}\frac{\partial^{2}r_{e}\left(\overline{R_{VIS}},\overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \\ \end{array}\right) \left(\begin{array}{c}\sigma_{VIS}^{2} \\ cov \\ \sigma_{SWIR}^{2}\end{array}\right).$$

235

236 Eq. (8) is the central equation of our framework for quantifying the impact of sub-pixel reflectance variance on r_e and τ retrievals. Eq. (8) decomposes the impact of 237 sub-pixel cloud reflectance variability on the au and r_e retrievals based on the bi-spectral 238 239 method into two parts: 1) the magnitude of the sub-pixel reflectance variance and covariance specified by the vector $\left(\sigma_{_{V\!I\!S}}^2,\,\mathrm{cov},\,\sigma_{_{SW\!I\!R}}^2
ight)^T$ (referred to as "sub-pixel 240 variance vector") and 2) the matrix of the second-order derivatives of the LUT with 241 respect to $R_{_{V\!I\!S}}$ and $R_{_{S\!W\!I\!R}}$ (referred to as "matrix of 2nd derivatives"). Given the LUT, the 242 matrix of 2nd derivatives can be easily derived from straightforward numerical 243

244 differentiation. An example of such a derived matrix based on the LUT for 0.86 μ m 245 reflectance ($R_{0.86}$) and 2.1 μ m reflectance ($R_{2.1}$) is shown in Figure 3. The values of the 246 2nd derivatives for the grids of LUT are indicated by the color bar. Note that the sign of 247 $\Delta \tau$ or Δr_e is determined both by the 2nd derivatives and the sub-pixel variance vector 248 $(\sigma_{VIS}^2, \text{ cov}, \sigma_{SWIR}^2)^T$. While σ_{VIS}^2 and σ_{SWIR}^2 are positive definite, the covariance term can 249 be negative.

It is clear from Eq. (8) that the τ and r_e retrievals are not only influenced by the sub-pixel variation of the primary band (i.e., R_{VIS} for τ and R_{SWIR} for r_e) but also by the variation of the secondary band (i.e., R_{SWIR} for τ and R_{VIS} for r_e), as well as the covariance of the two bands R_{VIS} and R_{SWIR} . Therefore, it reconciles and unifies the theoretical frameworks in Marshak et al. [2006] and Zhang and Platnick [*Zhang and Platnick*, 2011] and Zhang et al. [2012]. In particular, the impact of the PPHB on τ and r_e , described in Marshak et al. [2006], corresponds to the upper-left term,

257
$$-\frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}}\right)}{\partial R_{_{VIS}}^2}$$
 (Figure 3a), and lower-right term, $-\frac{1}{2} \frac{\partial^2 r_e \left(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}}\right)}{\partial R_{_{SWIR}}^2}$ (Figure 3f), in

258 the 2nd derivatives matrix, respectively. As shown in Figure 3, both terms are generally 259 negative over the most part of LUT, consistent with the finding of Marshak et al. [2006] 260 that ignoring sub-pixel variability tends to result in an underestimation of the pixel 261 average of the retrieved quantity if τ and r_e retrievals are considered separately and 262 independently (i.e., negative $\Delta \tau$ and Δr_e). On the other hand, $\Delta \tau$ and Δr_e are also 263 influenced by other terms in the matrix. Physically, these terms arise from the fact that 264 both $R_{\scriptscriptstyle V\!I\!S}$ and $R_{\scriptscriptstyle S\!W\!I\!R}$ depend not only on au but also $r_{_e}$. For example, the

$$\begin{array}{ll}
265 & -\frac{1}{2} \frac{\partial^2 r_e \left(\overline{R_{VIS}}, \overline{R_{SWR}}\right)}{\partial R_{VIS}^2} \text{ term in Figure 3d is mostly positive in the region of the LUT with } \tau \\
266 between about 1.5 and 20 and r_e between about 10 and 28 µm. This term competes \\
267 with the negative $-\frac{1}{2} \frac{\partial^2 r_e \left(\overline{R_{VIS}}, \overline{R_{SWR}}\right)}{\partial R_{SWR}^2} \text{ term in determining the sign and size of } \Delta r_e$. In
268 some cases, when σ_{VIS}^2 is large as in the example in Figure 2c, the influence of
269 $-\frac{1}{2} \frac{\partial^2 r_e \left(\overline{R_{VIS}}, \overline{R_{SWR}}\right)}{\partial R_{VIS}^2} \text{ may be stronger, leading to a positive } \Delta r_e$, as argued in Zhang and
270 Platnick [2011] and Zhang et al. [2012].
271 Some new terms that have not been explained in previous studies, e.g., the cross
272 terms $-\frac{\partial^2 \tau \left(\overline{R_{VIS}}, \overline{R_{SWR}}\right)}{\partial R_{VIS}} \text{ in Figure 3b and } -\frac{\partial^2 r_e \left(\overline{R_{VIS}}, \overline{R_{SWR}}\right)}{\partial R_{VIS}} \text{ in Figure 3e, also emerge from } \\$$$

273 Eq. (8). These two terms generally have the opposite sign of $-\frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}}\right)}{\partial R_{_{VIS}}^2}$ and

274
$$-\frac{1}{2}\frac{\partial^2 r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})}{\partial R_{SWIR}^2}$$
. Because the covariance cov is generally positive, the cross terms

275 evidently counteract the effects of
$$-\frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}}\right)}{\partial R_{_{VIS}}^2}$$
 and $-\frac{1}{2} \frac{\partial^2 r_e \left(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}}\right)}{\partial R_{_{SWIR}}^2}$ on $\Delta \tau$

276 and Δr_e .

Eq. (8) also provides a quantitative explanation for why sub-pixel inhomogeneity has different impacts on the r_e retrievals based on different SWIR bands (i.e., $r_{e,2.1}$ vs. $r_{e,3.7}$). Figure 4 shows an example of the matrix of 2nd derivatives for the $R_{0.86}$ and $R_{3.7}$ 280 combination. In comparison with the $R_{0.86}$ and $R_{2.1}$ combination in Figure 3, the

281	$-\frac{1}{2}\frac{\partial^2 r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})}{\partial R_{VIS}^2}$ term in Figure 4d is significantly smaller. This suggests that the same
282	sub-pixel inhomogeneity in the 0.86 μm band (i.e., same $\sigma^2_{_{V\!I\!S}}$) has a stronger impact on
283	$r_{e,2.1}$ than it does on $r_{e,3.7}$. Because this term tends to lead to a positive Δr_e bias, it could
284	be an important reason why the MODIS $r_{e,2.1}$ retrievals are often found to be significantly
285	larger than the $r_{e,3.7}$, in particular for inhomogeneous pixels [Painemal and Zuidema,
286	2011; Zhang and Platnick, 2011; Zhang et al., 2012; Cho et al., 2015].
287	As analyzed above, in comparison with previous studies the framework
288	described in Eq. (8) provides us with a more comprehensive explanation of the bias in $ au$
289	and $r_{\!_e}$ retrievals caused by the homogenous pixel assumption. This framework may be
290	useful in a variety of applications. It can be used to quantify Δau and $\Delta r_{_e}$ if the sub-pixel
291	variances and covariance $\left(\sigma_{_{V\!I\!S}}^2,\ ext{cov},\ \sigma_{_{SW\!I\!R}}^2 ight)^{\!T}$ are known, as shown in the example in the
292	next section. The Δau and $\Delta r_{_e}$ can then in turn be used to estimate the uncertainties
293	and potential biases in $ au$ and $r_{_{\! e}}$ retrievals due to ignoring the sub-pixel reflectance
294	variability in the bi-spectral method. Our framework can also be used to understand the
295	differences among retrievals based on instruments with different spatial resolutions.
296	Finally, it is worth mentioning that Eq. (8) can be rewritten in a slightly different
297	form as follows:

$$\begin{pmatrix} \Delta \tau \\ \Delta r_{e} \end{pmatrix} = \begin{pmatrix} -\frac{1}{2} \frac{\partial^{2} \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} \left(\overline{R_{VIS}}\right)^{2} & -\frac{\partial^{2} \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} \tau \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} \\ -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} \left(\overline{R_{VIS}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} \\ +\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} \left(\overline{R_{VIS}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} \\ +\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS}^{2}} \left(\overline{R_{VIS}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{VIS} \partial R_{SWIR}} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{1}{2} \frac{\partial^{2} r_{e} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2}} \left(\overline{R_{SWIR}}\right)^{2} & -\frac{\partial^{2} r_{e} \left(\overline{R_{SWIR}} \cdot \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2} \left(\overline{R_{VIS}} \cdot \overline{R_{SWIR}}\right) & -\frac{\partial^{2} r_{e} \left(\overline{R_{SWIR}} \cdot \overline{R_{SWIR}}\right)}{\partial R_{SWIR}^{2} \left(\overline{R_{SWIR}} \cdot \overline{R_{SWIR}}\right)} \left(\overline{R_{SW$$

(9)

302 where
$$H_{\sigma_{VIS}}^2 = \sigma_{VIS}^2 / \left(\overline{R_{VIS}}\right)^2$$
, $H_{\sigma_{SWIR}}^2 = \sigma_{SWIR}^2 / \left(\overline{R_{SWIR}}\right)^2$, and

 $H_{COV} = \operatorname{cov}(R_{VIS}, R_{SWIR}) / (\overline{R_{VIS}} \cdot \overline{R_{SWIR}})$. Note that $H_{\sigma_{VIS}}$ has been used in previous studies 304 as an index of sub-pixel inhomogeneity, in particular for MODIS cloud retrievals [e.g., *Liang et al.*, 2009; *Di Girolamo et al.*, 2010; *Zhang and Platnick*, 2011; *Zhang et al.*, 2012; *Cho et al.*, 2015]. Therefore, although Eq. (9) and (8) are equivalent, some readers may

307 find
$$\left(H_{\sigma_{VIS}}^2, H_{cov}, H_{\sigma_{SWIR}}^2\right)^T$$
 more familiar than $\left(\sigma_{VIS}^2, \text{ cov}, \sigma_{SWIR}^2\right)^T$.

It is important to point out that Eqs. (8)-(9) hold, no matter whether the subpixel reflectance variations (i.e., non-zero $(\sigma_{VIS}^2, \text{ cov}, \sigma_{SWIR}^2)^T$) are attributable to subpixel scale cloud property variations, 3-D radiative effects, or both. It is the interpretation of the resultant $\Delta \tau$ and Δr_e that is dependent on the circumstances and needs to be made with caution.

Finally, it is important to note that a critical assumption in our derivation is that the truncation error ε in the Taylor expansion is negligible. This term is a summation of all the higher order derivatives. Take r_e for example, the form of the k^{th} order derivative:

316
$$\frac{1}{k!} \frac{d^k r_e}{dR^k} = \sum_{0 \le m \le k} \frac{1}{m!(k-m)!} \frac{\partial^k r_e}{\partial R^m_{vis} \partial R^{k-m}_{vis}} \left(\overline{R_{VIS}}, \overline{R_{SWIR}}\right) \Delta R^m_{VIS} \Delta R^{k-m}_{SWIR} \,. \tag{10}$$

Because there is no analytical solution to the higher order derivatives, we can only assess the validity of this assumption and evaluate the accuracy of our framework numerically, which is done in the next section.

320 4. Numerical tests

In this section, we evaluate the accuracy and limits of our mathematical framework using two examples. The main objective is to assess, through case studies, if the higher order derivative terms are negligible so that our framework in Eq. (8) provides an accurate estimate of the PPHB.

325 4.1. Cloud fields from large-eddy simulation

326 In the first example, we test our framework using a synthetic cloud field simulated 327 from a large-eddy simulations (LES) model (DHARMA) with bin microphysics [Ackerman et al., 2004]. The LES case is based on an idealized case study [Stevens et al., 2010] from 328 329 the Atlantic Trade Wind Experiment (ATEX), with an diagnostic treatment of aerosol, specified to have a uniform number concentration of 40 cm⁻³. The ATEX simulation 330 represents a trade-wind cumulus case under a sharp inversion. The ATEX simulation has 331 a domain size of 9.6 x 9.6 x 3 km, with a uniform horizontal grid of $\Delta x = \Delta y = 100$ m and a 332 333 fixed vertical grid spacing of Δz =40 m. Further details of the model setup for the LES 334 case are provided in Zhang et al. [2012]. The droplet size distributions from the LES are 335 used to drive the radiative transfer simulations. The solar zenith and azimuth angle are 336 set at 20° and 30°, respectively, for the radiative transfer simulations. For simplicity, the surface is assumed to be black. Both 1-D and 3-D radiative transfer simulations were
performed, using the DISORT [*Stamnes et al.*, 1988] and the I3RC model [*Pincus and Evans*, 2009], respectively. We focus on the 3-D results because they are more
representative of real retrievals. The 1-D results are similar and are therefore not shown
here.

342 We first run radiative transfer simulations at the 100-m horizontal resolution of the LES grid. Figure 5a-c show the simulated 100-m cloud bi-directional reflectances at 343 344 nadir viewing angle for the 0.86, 2.1, and 3.7 µm MODIS bands, receptively. Then, the 345 100m reflectances are aggregated to 400 m to simulate the coarse-resolution observations, which are shown in Figure 5d—f. Obviously, for each 400-m pixel we have 346 347 4x4 100m pixels that can be used to derive the variances and covariances of sub-pixel reflectance variances. Figure 6 shows the sub-pixel reflectance variances, $\sigma_{0.86}^2$, $\sigma_{2.1}^2$ and 348 $\sigma_{\scriptscriptstyle 3.7}^2$, and covariances, $\mathrm{cov}(R_{\scriptscriptstyle 0.86},R_{\scriptscriptstyle 2.1})$ and $\mathrm{cov}(R_{\scriptscriptstyle 0.86},R_{\scriptscriptstyle 3.7})$ derived from 100-m 349 reflectances. Because of the large, order-of-magnitude differences between $R_{
m 0.86}$, $R_{
m 2.1}$ 350 and $R_{_{3.7}}$, $\sigma_{_{0.86}}^2$ is substantially larger than $\sigma_{_{2.1}}^2$, which in turn is substantially larger than 351 $\sigma_{_{3.7}}^{_2}$. Both covariances $\mathrm{cov}(R_{_{0.86}},R_{_{2.1}})$ and $\mathrm{cov}(R_{_{0.86}},R_{_{3.7}})$ are generally positive, 352 353 indicating a general positive correlation between SWIR and VIS/NIR band cloud reflectances. This is not surprising because $R_{2.1}$ and $R_{3.7}$ do increase with au when the 354 cloud is optically thin. Only for optically thick clouds do $R_{2.1}$ and $R_{3.7}$ become 355 independent from $R_{0.86}$. Figure 7 shows the reflectance variances and covariances 356 normalized by the mean reflectances squared, i.e., $H^2_{\sigma_{0.86}}$, $H^2_{\sigma_{2,1}}$ and $H^2_{\sigma_{3,7}}$ and 357

358
$$\operatorname{cov}(R_{0.86}, R_{2.1})/(\overline{R_{0.86}}, \overline{R_{2.1}})$$
 and $\operatorname{cov}(R_{0.86}, R_{3.7})/(\overline{R_{0.86}}, \overline{R_{3.7}})$. After the normalization,

359 $H_{\sigma_{0.86}}^2$, $H_{\sigma_{2.1}}^2$ and $H_{\sigma_{3.7}}^2$ are more comparable in terms of magnitude. In addition, cloud 360 edges are seen to have larger sub-pixel inhomogeneity than the center of the cloud, 361 which has also been found in MODIS observations [*Zhang and Platnick*, 2011; *Liang and* 362 *Girolamo*, 2013].

The au retrievals based on the simulated 100-m cloud reflectances ($R_{
m 0.86}$ and $R_{
m 2.1}$ 363 combination) in Figure 5a—b are shown in Figure 8a, which closely follow the $R_{0.86}$ 364 365 observations in Figure 5a. The au retrievals based on the $R_{0.86}$ and $R_{3.7}$ combination are mostly identical and therefore not shown. The $r_{e,2.1}$ and $r_{e,3.7}$ retrievals based on the 366 100-m reflectances are shown in Figure 8b-c. For consistency with the notation in 367 Section 3, we refer to these retrievals as sub-pixel retrievals, i.e., $au(R_{0.86,i},R_{2.1,i})$, 368 $r_e(R_{0.86,i},R_{2.1,i})$ and $r_e(R_{0.86,i},R_{3.7,i})$. The au , $r_{e,2.1}$ and $r_{e,3.7}$ retrievals based on the 369 aggregated 400m reflectances in Figure 5d-f are shown in Figure 8d-f, respectively, 370 which are referred to as pixel-level retrievals $\tau(\overline{R_{0.86,i}}, \overline{R_{2.1,i}})$, $r_e(\overline{R_{0.86,i}}, \overline{R_{2.1,i}})$ and 371 $r_e\left(\overline{R_{0.86,i}},\overline{R_{3.7,i}}\right).$ 372

Having derived both sub-pixel and pixel level retrievals, we first compute the biases caused by the homogenous pixel assumption, $\Delta \tau$ and Δr_e , as expressed in Eq. (3). The results are shown in Figure 9a—c. It can be seen that $\Delta \tau$ is mostly negative over the whole domain, as one would expect based on the PPHB. However, the Δr_e , especially $\Delta r_{e,2.1}$, is predominantly positive, which is the opposite of PPHB but consistent with the findings in Zhang and Platnick [2011] and Zhang et al. [2012]. It should be pointed out that the cloud-free pixels are marked in black in the figure. The pixels in gray are partly cloudy pixels (i.e., one or more 100-m sub-pixels are cloud-free). (Because it is uncertain how cloud-free sub-pixels should be treated in the spatial averages, partly cloudy pixels are excluded from our analysis.)

To assess the accuracy of our framework, we derived the second set of $\Delta \tau$ and Δr_e 383 based on Eq. (8) using the matrix of 2nd derivatives (Figure 3 and Figure 4) and the sub-384 385 pixel reflectance variances and covariances (Figure 6). The results from this method are shown in Figure 9d—f. Evidently, $\Delta \tau$ and Δr_e derived in two different and independent 386 ways agree very well. The correlation coefficients all exceed 0.8 as shown in Figure 9g-387 i. Only those pixels with large sub-pixel inhomogeneity index $H_{\sigma_{\rm 0.86}}$ (>0.5) deviate from 388 the one-to-one line. For these pixels the higher order terms $O(\Delta R^3)$ ignored in Eq. (8), 389 likely impact $\Delta \tau$ and Δr_e . But such cases are relatively rare for this LES scene. The 390 overall excellent agreement clearly demonstrates that our framework is able to provide 391 an accurate quantitative estimation of the biases in au and $r_{_{\!e}}$ retrievals caused by the 392 393 homogenous pixel assumption for overcast pixels.

An advantage of using Eq. (8) is that the bias can be further decomposed into the contributions from each term in the matrix of 2^{nd} derivatives, which help us to better understand the relative importance of various factors in causing the bias. For example, as shown in Figure 10a—c, the τ retrieval bias is dominated by the

398 $-\frac{1}{2} \frac{\partial^2 \tau(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}})}{\partial R_{_{VIS}}^2} \cdot \sigma_{_{VIS}}^2 \text{ term in Eq. (7). As mentioned before, this term corresponds to}$

the PPHB (Figure 2a), which is why the total $\Delta \tau$ in Figure 9 is generally negative. In the

400 case of the
$$r_{e,3.7}$$
 retrieval, both the positive $-\frac{1}{2} \frac{\partial^2 r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})}{\partial R_{VIS}^2} \cdot \sigma_{VIS}^2$ term (Figure 10g)

401 and the negative
$$-\frac{1}{2} \frac{\partial^2 r_e (R_{VIS}, R_{SWIR})}{\partial R_{SWIR}^2} \cdot \sigma_{SWIR}^2$$
 term (Figure 10i) are significant. The former

402 corresponds to the example in Figure 2c, while the latter refers to the example in Figure

403 2b. After summation, the
$$-\frac{1}{2} \frac{\partial^2 r_e(\overline{R_{VIS}}, \overline{R_{SWIR}})}{\partial R_{VIS}^2} \sigma_{VIS}^2$$
 is dominant and leads to the overall

404 positive bias in the $r_{e,3.7}$ retrieval. The bias in the $r_{e,2.1}$ retrieval is even more 405 complicated, as all three terms on the right hand side of Eq. (6) contribute substantially 406 to the bias. Overall, the positive terms in Figure 10d—e dominate the total error budget, 407 leading to a generally positive $\Delta r_{e,2.1}$ in Figure 8.

In the above example, the solar zenith angle is high, with $\theta_0 = 20^{\circ}$. We also tested our framework in a case with low solar zenith angle of $\theta_0 = 60^{\circ}$ and the results are shown in Figure 11. The correlations between the biases from the numerical simulations and those predicted by our framework are substantial, suggesting our framework works equally well for a high sun in this case.

From the above example, one can clearly see that our framework provides a comprehensive explanation of the impact of sub-pixel inhomogeneity on τ and r_e retrievals. As mentioned earlier we have also tested our framework on the retrievals based on reflectance using 1-D radiative transfer, and find the predicted $\Delta \tau$ and Δr_e based on our framework to agree very well with the numerical results (not shown).

418 We'd like to point out here that less sensitivity to sub-pixel heterogeneity in the 419 3.7 μ m channel should not necessarily be equated to less r_e bias in the overall retrieval. For simplicity, our 3.7 µm analysis deals with reflectance only. Thus it assumes that the 420 421 cloud and surface temperatures are known without error, as are the atmospheric 422 emission/correction terms, needed to infer cloud top reflectance from top-of-423 atmosphere measurements of emitted and reflected radiation. Because we are dealing 424 with reflectance only, it is implicitly assumed that the effects of sub-pixel heterogeneity 425 on the cloud temperature retrieval and atmospheric correction are negligible. The validity of this assumption will be assessed in future work. 426

427

428 4.2. MODIS retrieval test

429 In the second example, we test our framework using MODIS observations. The MODIS instrument has 36 spectral bands. The spatial resolution of most bands (bands 430 8-36) is 1 km. Bands 3-7 have a 500-m resolution. Bands 1 and 2 have a 250 m spatial 431 resolution. The current (collection 06) operational MODIS cloud property retrieval 432 products, such as au, r_e and LWP, are made at 1-km resolution. The higher spatial 433 resolution of the 0.86 µm (band 2) and 2.1 µm (band 7) sensors provides us with an 434 435 opportunity to test our framework and investigate the impact of sub-pixel 436 inhomogeneity on the MODIS τ and r_{e} retrievals. For this purpose, we selected a case shown in Figure 12. The granule in Figure 12a was collected by MODIS onboard the 437 Terra satellite on September 9th 2006 over the Gulf of Mexico. We further selected a 438 439 small region off the coast of Louisiana marked in the red box for our test. A zoom-in

view of this small region at the 1km and 500m resolutions is shown in Figure 12b and
Figure 12c, respectively.

442 Similar to the LES example, we first developed two sets of cloud property retrievals, 443 one at a higher spatial resolution of 500 m and the other at a coarser resolution of 1 km. Figure 13a and b show the 500 m resolution τ and r_e retrievals, respectively, based on 444 the combination of 0.86 and 2.1 μ m reflectances for the selected region in Figure 12b. 445 446 The 1 km retrievals are shown in Figure 13c and d. This scene has a cloud fraction of about 72%. In the center of the scene is a cluster of thick clouds with τ around 20 to 30, 447 and r_e ranging mainly between 15µm to 20µm. Note that in our framework the 500 m 448 retrievals are the sub-pixel $\tau(R_{VIS,i}, R_{SWIR,i})$ and $r_e(R_{VIS,i}, R_{SWIR,i})$. The 1 km retrievals are 449 $\tau(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}})$ and $r_e(\overline{R_{_{VIS}}}, \overline{R_{_{SWIR}}})$. To derive the $\Delta \tau$ and Δr_e from our mathematical 450 451 framework in Eq. (8), we compute the sub-pixel reflectance variances and covariances 452 for every 1-km cloudy pixel from the 2x2 500-m sub-pixel reflectance observations. The 453 results are shown in Figure 14. Similar to the LES case, we find that the 0.86 and 2.1 µm cloud reflectances are generally positively correlated over the thin cloud regions. The 454 455 correlation becomes weak (close to zero) over the thick cloud regions. These results 456 indicate that when the cloud is thin, the variability in both 0.86 and 2.1 μ m bands is 457 controlled mainly by τ . The variability of 2.1 μ m cloud reflectances becomes primarily 458 sensitive to r_e only when the cloud becomes optically thick.

459 The difference between the 1 km retrievals and the mean of 500 m retrievals are the 460 biases, $\Delta \tau$ and Δr_e , caused by the homogeneous pixel assumption. Figure 15a and b

show $\Delta \tau$ and Δr_e , respectively, based on Eq. (3). We found that $\Delta \tau$ is mainly negative 461 particularly in the regions with thick clouds, while Δr_e is mainly positive particularly in 462 463 the transition regions from thick to thin clouds. These results are very similar to what we found in the LES scene in Figure 9. Both Δau and Δr_e are shown in Figure 15c and d, 464 respectively. The $\Delta \tau$ and Δr_{e} predicted from Eq. (8) agree reasonably well with the 465 results derived from numerical retrievals in Figure 15a and b. The predicted $\Delta \tau$ based 466 467 on Eq. (9) and the numerical results have a correlation coefficient over 0.85 for all cloudy pixels (over 0.95 for pixels with $\tau > 5$). The correlation coefficient for Δr_{e} is 468 469 significantly lower especially for thin clouds with $\tau < 5$. This is mainly because when the cloud is thin the 2.1 μ m cloud reflectances are not very sensitive to r_e . As a result, the 470 471 retrievals are subject to large uncertainties caused by radiative transfer model 472 uncertainties. If we limited the comparison only to clouds with $\tau > 5$, the correlation coefficient is over 0.70. 473

In summary, our numerical framework work very well for the LES cases, indicating 474 475 that the high-order terms are mostly negligible in these cases. It also works reasonably well for the real MODIS case, especially for the clouds with $\tau > 5$. For thinner clouds, it 476 477 is difficult to tell whether the deviation stems from higher-order terms or retrieval 478 uncertainties. Another factor to consider is that we only have four 500 m sub-pixels for 479 each 1 km pixel, which may be insufficient for deriving meaningful sub-pixel variance and co-variance. As part of ongoing research, we are trying to retrieve τ and r_e from 480 the Advanced Space-borne Thermal Emission and Reflection Radiometer (ASTER) on 481 482 Terra. ASTER has a much greater spatial resolution than MODIS and therefore can

provide much richer information on small scale variability of cloud reflectance [*Zhao and Di Girolamo*, 2006; *Wen et al.*, 2007]. We will further test our framework using ASTER
observations in future work.

486 **5. Summary and Discussion**

487 The impact of unresolved sub-pixel level variation of cloud reflectances is an 488 important source of uncertainty in the bi-spectral solar reflective method. In this study, 489 we develop a mathematical framework for understanding this impact and quantifying the consequent biases, Δau and Δr_e . We show in Eq. (8) that Δau and Δr_e are 490 determined by two factors-the nonlinearity of the LUT and the inhomogeneity of 491 492 reflectances within the pixel. We tested our framework using LES cloud fields and real 493 MODIS observations. The results indicate that, in comparison with previous studies, our 494 framework provides a more comprehensive explanation and also a more accurate 495 estimation of the retrieval biases caused by the sub-pixel level variation of cloud reflectances. Most importantly, it demonstrates that sub-pixel variations in cloud 496 reflectance can lead to both positive and negative values of $\Delta r_{\!_e}.$ In both the LES and 497 MODIS cases that we examined, Δr_e were dominantly positive, hence contributing to the 498 dominantly positive bias in retrieved r_e from resolved cloud variability. 499

500 Our framework could have several applications. For example, it can be used to 501 understand the differences between retrievals made at different spatial resolutions 502 (e.g., MODIS vs. SEVIRI) or based on different spectral reflectances (e.g., MODIS 2.1 μ m 503 vs. 3.7 μ m). It could also useful for estimating retrieval uncertainties. For example, the 504 retrieval uncertainty caused by sub-pixel reflectance variation in the operational 1 km

505 MODIS cloud products can be estimated based on our framework from the 500 m cloud 506 reflectances. It can also be integrated into the operational MODIS retrieval algorithm to 507 determine in real-time whether the high-resolution retrievals (e.g., from 1km to 500m) are necessary for a given pixel. Another useful application is to help the trade-off studies 508 509 for instrument design. For example, the Ocean Color Imager (OCI) is the key instrument planned for NASA's coming Pre-Aerosol, Clouds, and ocean Ecosystem mission 510 (http://decadal.gsfc.nasa.gov/pace.html). An important part of the OCI design trade-off 511 512 study is to determine the optimal spatial resolution for both ocean color and atmospheric observations, including cloud property retrievals. Our framework would be 513 514 highly useful for such study.

Finally, we feel necessary to clarify again that our framework cannot explain or predict 3-D effects, such as the illuminating, shadowing, and photon leaking, which are known to substantially influence cloud reflectances and therefore retrieval results. These effects are beyond the scope of this study. Our framework simply predicts the statistical differences between retrievals with different spatial resolutions, regardless of whether the radiative transfer is 1-D or 3-D.

522 Acknowledgement:

This research is supported by NASA grants NNX14AJ25G and NNX15AC77G. The computations in this study were performed at the UMBC High Performance Computing Facility (HPCF). The facility is supported by the U.S. National Science Foundation through the MRI program (grant nos. CNS-0821258 and CNS-1228778) and the SCREMS program (grant no. DMS-0821311), with additional substantial support from UMBC. The MODIS data are obtained from NASA's Level 1 and Atmosphere Archive and Distribution System (LAADS http://ladsweb.nascom.nasa.gov/).

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642 Figure 1 Examples of the look-up table of cloud bi-directional reflection function as functions of

cloud optical thickness and effective radius, based on the combination of a) 0.86 and 2.1 μm
bands, and b) 0.86 and 3.7 μm bands. Surface is assumed to be Lambertian with a reflectance of

645 0.02. Solar and viewing zenith angle are 45° and 20°, respectively. Relative azimuthal angle is 0°.



Figure 2 a) an example to illustrate the PPHB bias proposed in Cahalan et al. [1994] for au

- 650 retrieval, b) example to illustrate the PPHB bias proposed in Marshak et al. [2006], c) example
- to illustrate the r_e retrieval bias caused by sub-pixel au variability proposed in Zhang and
- Platnick [2011] and Zhang et al. [2012]. See text for details. Solar and view zenith angles are
- assume to be 20° and 0° and relative azimuth angle is assumed to be 30° in these cases.
- 654



660 be 30° in these cases.



 $R_{asc}^{00} = 2^{0} + 10^{0} + 1$



667 ¹/₄ ²/₄ ⁴/₄ ⁴/

669 100-m resolution for the LES cloud field. d)—f) 400-m bi-directional reflectances averaged from

670 100-m resolution simulations.



Figure 6 The sub-pixel reflectance variance a) $\sigma_{0.86}^2$, b) $\sigma_{2.1}^2$, c) $\sigma_{3.7}^2$ and covariances d) $cov(R_{0.86}, R_{2.1})$ and e) $cov(R_{0.86}, R_{3.7})$ for the LES case in Figure 5.







Figure 8 a) τ , b) $r_{e,2.1}$ and c) $r_{e,3.7}$ retrievals based on the 100 m reflectance. d)—f) retrievals

based on the 400 m reflectance.



Figure 9 The a) Δτ, b) Δr_{e,2.1} and c) Δr_{e,3.7} derived based on the Eq. (3). The corresponding
results obtained based on Eq. (8) are shown in d)—f). The pixel-to-pixel comparisons are shown
in g)—i), in which the color indicate the value of the sub-pixel inhomogeneity index H_{σ0.86}.



699
$$-\frac{1}{2} \frac{\partial^2 \tau(\overline{R_{0.86}}, \overline{R_{2.1}})}{\partial R_{0.86} \partial R_{2.1}} \cdot \operatorname{cov}(R_{0.86}, R_{2.1}) \text{ to } \Delta \tau, \mathbf{c} \text{ contribution of } -\frac{1}{2} \frac{\partial^2 \tau(\overline{R_{0.86}}, \overline{R_{2.1}})}{\partial R_{2.1}^2} \cdot \sigma_{2.1}^2 \text{ to } \Delta \tau. \mathbf{d} \text{)}$$

700 contribution of $-\frac{1}{2} \frac{\partial^2 r_e(\overline{R_{0.86}}, \overline{R_{2.1}})}{\partial R_{0.86}^2} \cdot \sigma_{0.86}^2$ to $\Delta r_{e,2.1}$, **e**) contribution of 701 $-\frac{1}{2} \frac{\partial^2 r_e(\overline{R_{0.86}}, \overline{R_{2.1}})}{\partial R_{0.86} \partial R_{2.1}} \cdot \operatorname{cov}(R_{0.86}, R_{2.1})$ to $\Delta r_{e,2.1}$, **f**) contribution of $-\frac{1}{2} \frac{\partial^2 r_e(\overline{R_{0.86}}, \overline{R_{2.1}})}{\partial R_{2.1}^2} \cdot \sigma_{2.1}^2$ to $\Delta r_{e,2.1}$. **g**)

702 contribution of
$$-\frac{1}{2} \frac{\partial^2 \tau \left(\overline{R_{0.86}}, \overline{R_{3.7}}\right)}{\partial R_{3.7}^2} \cdot \sigma_{3.7}^2$$
 to $\Delta r_{e,3.7}$. **h**) contribution of

703
$$-\frac{1}{2} \frac{\partial^2 r_e(R_{0.86}, R_{3.7})}{\partial R_{0.86} \partial R_{3.7}} \cdot \operatorname{cov}(R_{0.86}, R_{3.7}) \text{ to } \Delta r_{e,3.7}, \mathbf{f} \text{ contribution of } -\frac{1}{2} \frac{\partial^2 r_e(R_{0.86}, R_{3.7})}{\partial R_{3.7}^2} \cdot \sigma_{3.7}^2 \text{ to } \Delta r_{e,2.1}.$$





- Figure 12 The a) RGB image of a MODIS granule collected on September 9th 2006 over the Gulf
- of Mexico. A zoom-in view of the region in the red box based on b) 1 km MODIS true color RGB
- image and c) 500 m MODIS true color RGB image..





 $\operatorname{cov}(R_{0.86}, R_{2.1})$ for the MODIS case in Figure 12b.



727True $\Delta \tau_{e,2,1}$ True $\Delta r_{e,2,1}$ (µm)728Figure 15 The a) $\Delta \tau$ and b) Δr_e derived based on the Eq. (3). The corresponding results based729on Eq. (8) are shown in c) and d) and comparisons in e) and f).