

A Highly Elastic, Capacitive Strain Gauge Based on Percolating Nanotube Networks

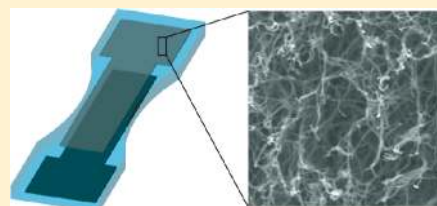
Daniel J. Cohen,^{*,†} Debkishore Mitra,[†] Kevin Peterson,[‡] and Michel M. Maharbiz[‡]

[†]Departments of Bioengineering and [‡]Electrical Engineering and Computer Science, University of California, Berkeley, California 94720, United States

S Supporting Information

ABSTRACT: We present a highly elastic strain gauge based on capacitive sensing of parallel, carbon nanotube-based percolation electrodes separated by a dielectric elastomer. The fabrication, relying on vacuum filtration of single-walled carbon nanotubes and hydrophobic patterning of silicone, is both rapid and inexpensive. We demonstrate reliable, linear performance over thousands of cycles at up to 100% strain with less than 3% variability and the highest reported gauge factor for a device of this class (0.99). We further demonstrate use of this sensor in a robotics context to transduce joint angles.

KEYWORDS: Nanotubes, percolation, sensor, strain, elastomer



Conventional, metal-foil strain gauges are limited to applications involving both relatively stiff substrates and strains under 5% (above which the gauge mechanically fails).^{1,2} As such, they are particularly ill-suited for use in the low-stiffness, high-strain environments that characterize many medical applications and, increasingly, soft robotics systems.^{3–8} Recently, many efforts to address this problem have focused on carbon nanotube-based percolation networks that are variously coupled to elastomeric matrices.^{9–14} Percolating networks consist of a collection of conductive particles in two- or three-dimensions that overlap with each other to create a conductive network. Such a network is able to maintain conductivity even when deformed due to continual contact among the conductive particles. The critical concentration, or percolation threshold, required to achieve stable percolation is a direct function of the geometry of these particles and scales with aspect ratio.¹⁵ Carbon nanotubes are particularly well suited for use in percolation networks as they have both a high intrinsic conductivity and aspect ratios as high as 1:1000.

The most common approach to percolation sensors relies on piezoresistance (strain induces a change in the resistance of the network). A variety of techniques have been developed for producing these networks using carbon nanotubes. Some of the earliest approaches relied on carefully dispersing nanotubes in an unpolymerized polymer before curing the ensemble to produce a conductive composite.^{7,16} While quite effective for flexible electronics, this approach results in significant reinforcement of the elastic modulus, thereby making the material too stiff for stretchable applications.¹⁶ More recent techniques involve infiltration of an elastomer into a vertically aligned array of nanotubes grown on a wafer, direct contact transfer of wafer-grown or vacuum filtered nanotubes microcontact stamping, airbrushing of dispersed nanotubes onto a substrate, and manual assembly of sheets of aligned nanotubes onto the surface of an elastomeric support.^{2,17–20} While effective,

piezoresistive designs can be difficult to tune as they inherently rely on the mechanical stability of the percolation network and are susceptible to hysteresis and variable gauge factors (the normalized change in resistance divided by the applied strain) as the network adjusts over time.

An alternative approach relies on the piezocapacitance of dielectric elastomers sandwiched by percolation electrodes. Here, any deformation that brings the electrodes closer together results in an increase in capacitance, and vice versa. While these devices are most often used as low-strain pressure sensors that transduce forces normal to the sensor surface, they can also be used to transduce planar, tensile strains (Figure 1a).⁸ To date, such planar strain piezocapacitors have been used to transduce strains up to 30% by using either nanotubes or conductive polymers.^{20–22} While we also opted to use nanotube percolation electrodes, our design was optimized for cyclic, planar strains up to 100% (Figure 1b–c).

The general mechanism by which such a device couples planar strain to a change in capacitance relies on Poisson contraction. Uniaxial, planar strain results in a Poisson's ratio-mediated contraction of the orthogonal axes that brings the two electrodes closer together resulting in a corresponding increase in capacitance. The simplified, linear mechanics of this process (Figure 1a) are described by eqs 1–3

$$\frac{\Delta L}{L} = \epsilon_z = \frac{\sigma_z}{E_{\text{silicone}}} \quad (1)$$

$$\epsilon_x = \epsilon_y = \nu \epsilon_z \quad (2)$$

Received: November 16, 2011

Revised: March 8, 2012

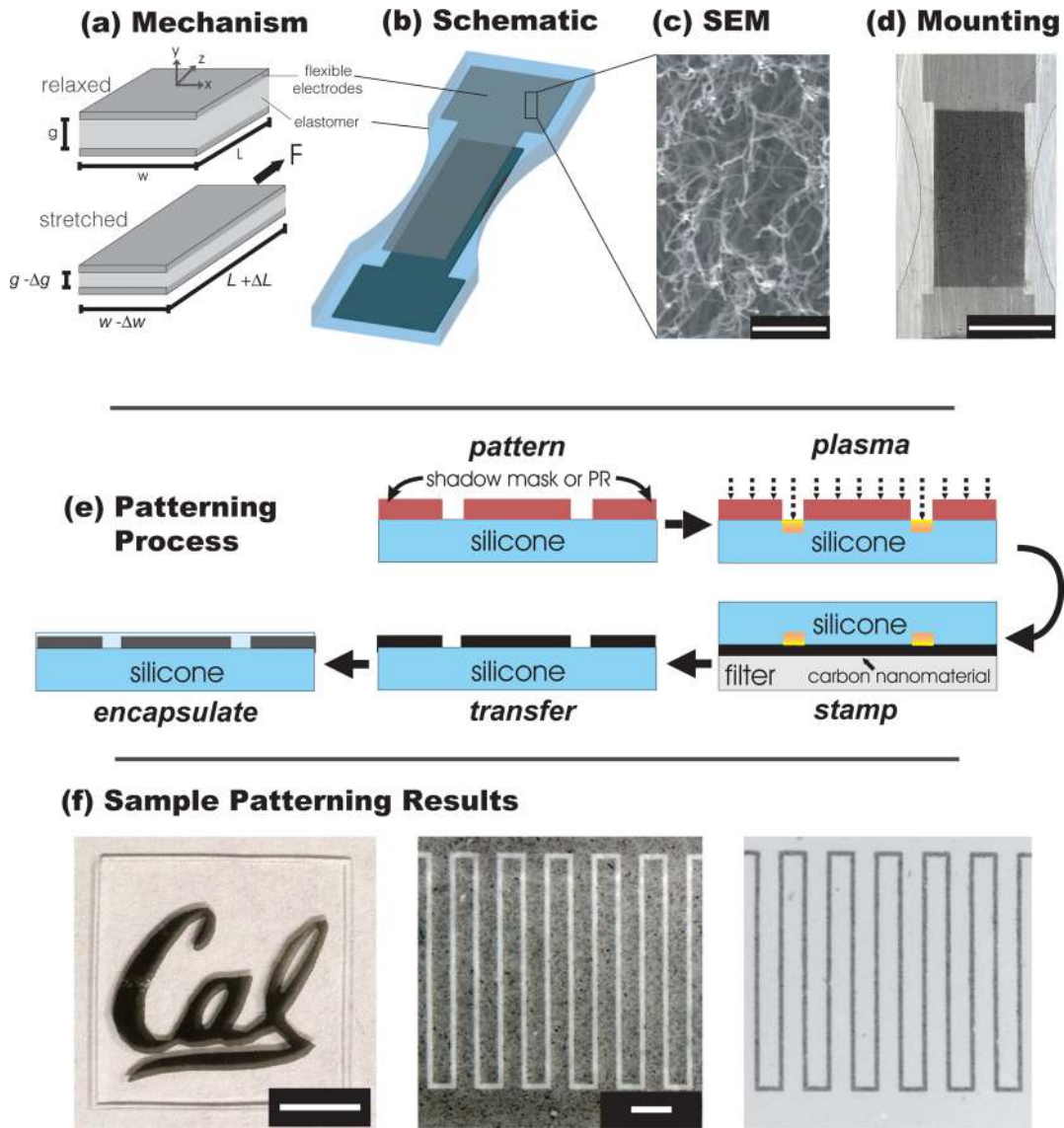


Figure 1. Mechanism, design, and fabrication of a Poisson Capacitor. (a) Poisson contraction converts planar strain to a decrease in the gap distance between the percolating electrodes and an increase in capacitance. (b) Schematic of our device geometry; all sensing is carried out in the middle region of the sensor. (c) SEM data demonstrating percolation of the nanotubes within the electrode; scale bar is 500 nm. (d) Close-up image of the sensing region of the device. Darker region is due to overlap of electrodes, and striations are due to texture of the background; scale bar is 0.75 cm. (e) Process flow for preparing a sensor. (f) Results from patterning. (left) Produced using atmospheric plasma and paper shadow-mask, fringing due to shadows; scale bar is 5 mm. (center) RIE oxygen plasma and photoresist was used to produce 300 μm wide digits; scale bar is 750 μm. (right) Nanotubes left on the filter after previous transfer, line width is 100 μm; same scale as left.

$$\begin{aligned}
 L_{\text{stretch}} &= L + \Delta L = L + \epsilon_z L \\
 w_{\text{stretch}} &= w - \Delta w = w - \nu \epsilon_z w \\
 g_{\text{stretch}} &= g - \Delta g = g - \nu \epsilon_z g
 \end{aligned}
 \tag{3}$$

75 Here, ν is Poisson's ratio (~ 0.5 for silicone elastomers); ϵ is
 76 strain; σ is stress; E is the elastic modulus of silicones; and L , w ,
 77 and g are the initial dimensions of the capacitor. Equation 1 is
 78 Hooke's law relating modulus, stress, and strain. While models
 79 taken from nonlinear elasticity theory are generally used for
 80 large deformations, the simple linear model described here has
 81 an equivalent goodness-of-fit ($R^2 = 0.997$) to nonlinear models
 82 when fitted to mechanical testing data taken from our sample
 83 devices (Figure S1 and Supporting Information).²³

84 Equations 2 and 3 describe how Poisson's ratio couples strain
 85 applied in one axis to opposite strains induced in the other two

axes. If we relate this to the simplest model for parallel plate
 capacitance, we arrive at eq 4 86
87

$$\begin{aligned}
 \Delta C_{\text{simple}} &= \epsilon_0 \epsilon_{\text{silicone}} \left(\frac{w_{\text{stretch}} L_{\text{stretch}}}{g_{\text{stretch}}} \right) \\
 &= \epsilon_0 \epsilon_{\text{silicone}} \left(\frac{w(1 - \nu \epsilon_z) L(1 + \nu \epsilon_z)}{g(1 - \nu \epsilon_z)} \right) \\
 &= \epsilon_0 \epsilon_{\text{silicone}} (1 + \epsilon_z) \frac{wL}{g}
 \end{aligned}
 \tag{4}$$

Here, ϵ_0 is the permittivity of free space, and $\epsilon_{\text{silicone}}$ is the
 relative permittivity of silicone. This simple electromechanical
 equation directly relates the change in capacitance of the 90
 parallel plate geometry to the applied strain and the initial 91

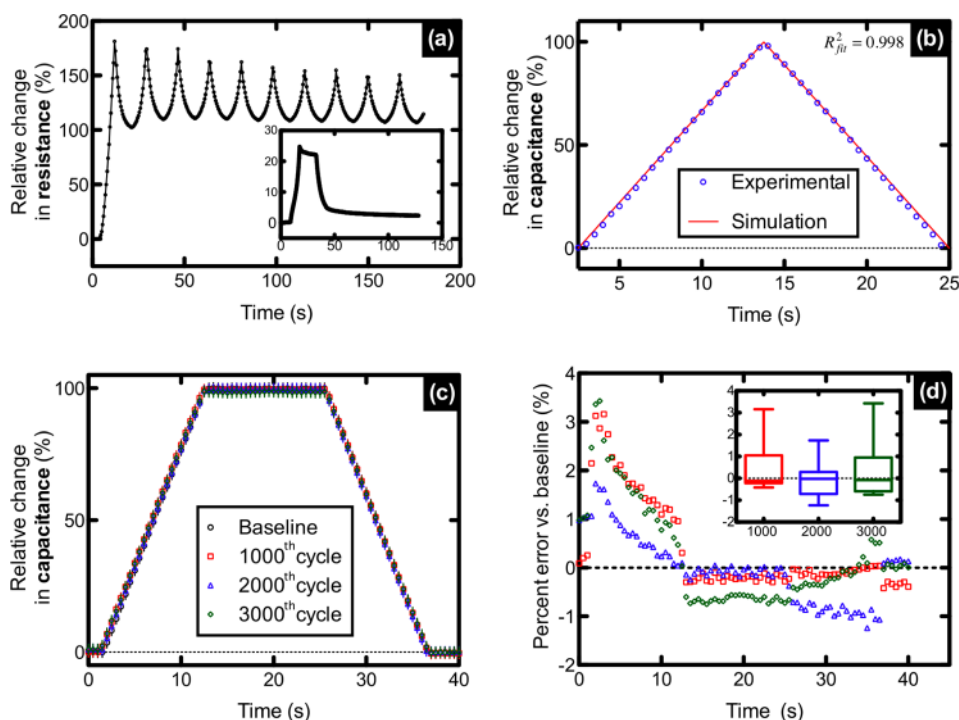


Figure 2. Performance characterization of Poisson capacitor strain gauge. (a) Resistive performance of a single percolation electrode when undergoing 100% cyclic strain. Main plot indicates significant hysteresis, while inset shows both relaxation during a 100% step strain and subsequent hysteresis during recovery. (b) Single, 100% strain cycle of the capacitive sensor (blue) overlaid (red) by the fit from parallel plate model with Palmer correction (red). (c) Demonstration of repeatability over 3000 cycles of 100% strain. After every 1000 cycles, the sample was subjected to a 100% step strain before being relaxed back to basal strain. (d) Stability plot showing how much sensor performance deviated from baseline performance over 3000 cycles. Legend is the same as for (c).

92 geometry of the sensor. It is interesting to note that the final
 93 equation, despite resulting from Poisson contraction, does not
 94 depend on the value of Poisson's ratio as long as the material is
 95 assumed to be isotropic. The equation further predicts a linear
 96 output, assuming that the overlap geometry remains stable.
 97 This implies that, regardless of resistive hysteresis (i.e.,
 98 piezoresistive effects, changes in the percolation network
 99 nanostructure), the sensor performance remains stable as
 100 long as the effective overlap area between percolation
 101 electrodes remains stable. While this model describes the
 102 general mechanics and agrees with that used by Loh et al. to
 103 describe a strain gauge based on a stiff polymer layer with
 104 nanotube/gold particle electrodes strained at 1%, it does not
 105 take into account the fringing fields normally obtained with
 106 finite-size parallel plate capacitors. To take these into account,
 107 we used the Palmer correction factor, which is a multiplier that
 108 acts on all terms in the simple model, for all of simulations (eq
 109 S1 and Supporting Information).²⁴

110 Our design, shown in mounted form in Figure 1d, relies on
 111 producing carbon nanotube percolation electrodes that are
 112 stable over large deformations and can be precisely patterned
 113 onto the elastomeric substrate such their shape and orientation
 114 can be specified independently of those of the elastomer. We
 115 chose silicone as the substrate both due its excellent elastomeric
 116 properties and its ability to facilitate patterning on both sides of
 117 a sheet via the hydrophobic to hydrophilic transition that
 118 results from atmospheric (or oxygen) plasma treatment.²⁵ As
 119 raw single-walled carbon nanotubes (SWNT) are highly
 120 hydrophobic, they will preferentially adhere to the hydrophobic
 121 (untreated) regions of the patterned silicone (Figure 1f).

Our fabrication process is summarized in Figure 1e. First, 122
 275 μm thick silicone substrates were laser-cut into the test 123
 structures (Figure 1b). The pattern resolution necessary for the 124
 sensor electrodes was achievable using laser-cut sticker-paper 125
 masks that were applied to both sides of the silicone substrates. 126
 The masked substrate was positioned on its side in an 127
 atmospheric plasma cleaner to allow the plasma access to both 128
 sides of the silicone to render the unmasked regions 129
 hydrophilic. Upon removal from the chamber, the masks 130
 were removed and the substrate was ready to receive 131
 nanotubes. The percolation networks were produced using 132
 SWNT that were dispersed in a surfactant solution via 133
 sonication and subsequently vacuum filtered and collected on a 134
 20 nm pore-size filtration membrane.^{26,27} The result of this 135
 process is a percolation network resting on top of a filtration 136
 membrane. By bringing both sides of the plasma-treated 137
 silicone into contact with these percolation networks, we 138
 achieved direct nanotube transfer to the untreated regions of 139
 the silicone and produced the overlapping electrodes necessary 140
 for a parallel plate capacitor. Should submillimeter features be 141
 required, the shadow mask can be replaced with a photoresist 142
 layer that can be removed via acetone immersion after plasma 143
 treatment. Between these two masking techniques, we have 144
 successfully demonstrated patterns ranging from centimeters in 145
 size to 100 μm wide serpentes (Figure 1f). 146

As a final step to protect and stabilize the newly transferred 147
 nanotube layers, we used an airbrush to coat the substrates with 148
 thin layers of silicone while avoiding spraying the ends of the 149
 device (the contact pads).²⁸ All of our tested sensors were 150
 designed with a 1.7 cm \times 0.75 cm electrode overlap area and an 151
 initial silicone spacer thickness of 275 μm . During testing, the 152

153 contact pads were clamped against copper foil leads connected
154 to an electrical impedance spectroscopy system (EIS). While
155 we used EIS to precisely characterize the device performance, a
156 capacitive Wheatstone bridge can be used as a lightweight,
157 portable alternative. Once connected to the EIS, the sensor was
158 placed in a customized mechanical stretcher capable of
159 cyclically stretching the sensor to 100% strain at 2 mm/s.
160 The fabrication process with special attention to the patterning
161 method is discussed in greater detail in Supporting Information.
162 Sensors made in this fashion were cycled between 0 and
163 100% strain (Figure 2a). For comparison, we first tested the
164 piezoresistive performance of our electrodes. As anticipated,
165 there was significant resistive hysteresis that resulted in a
166 permanent increase in baseline resistance of slightly over 100%
167 relative to the initial resistance (~ 15 k Ω by two-point
168 measurement). In addition to hysteresis, there was pronounced
169 viscoelastic behavior (inset, Figure 2a), where the sample was
170 strained to 100% and held at that strain for 10 s before being
171 relaxed. During this process, it is clear that the percolation
172 network undergoes stress relaxation. Lastly, there was a
173 continual decrease in the piezoresistive gauge factor with
174 cycling (~ 0.5 over the first 10 cycles to ~ 0.25).

175 By contrast, when used as capacitive strain sensors the
176 performance was markedly stable and reliable. The Poisson
177 capacitors averaged a baseline capacitance of 16 pF ($n = 6$;
178 measured by EIS). All measurements used an AC excitation
179 frequency of 10 kHz (selected by examining the frequency
180 response curve, Supporting Information Figure S2). Devices
181 subjected to 100% cyclic strain (Figure 2b) exhibited a relative
182 change in capacitance (blue circles) that closely matched ($R^2 =$
183 0.998) that predicted by the linear elastic model with the
184 Palmer correction factor (see Supporting Information). This
185 stable performance occurs despite the wide variability in the
186 resistive properties of the two percolation electrodes. Moreover,
187 the performance under increasing tension versus that under
188 decreasing tension is essentially symmetric. When comparing
189 strain versus normalized change in capacitance for these two
190 cases, the average slopes differed by $<2\%$.

191 The actual stability of the device was assessed over the course
192 of 3000 cycles at 100% cyclic strain (Figure 2c). During testing,
193 a new sensor was mounted and immediately subjected to a
194 “step-and-hold” test where it was stretched at 2 mm/s up to
195 100% strain and held for ~ 12.5 s before reversing the direction
196 of the stretcher. This data set was taken to be the baseline
197 response curve. Following this, the sensor was cycled from 0 to
198 100% strain with a sawtooth wave profile. Every 1000 cycles
199 (up to 3000 total cycles), the sensor was subjected to an
200 additional step-and-hold test to assess its stability. When used
201 as a Poisson capacitor, the device exhibited stable performance
202 with little hysteresis (Figure 2d). The offset error did not
203 exceed 3% of the baseline, and the trend was not monotonic,
204 implying a relatively stable network absent of significant
205 hysteresis. To our knowledge, variability of within 3% at over
206 3000 cycles represents some of the lowest values reported for
207 stretchable strain gauges (a typical value for a recent nanotube
208 resistive strain gauge was over 10% hysteresis after the first
209 cycle).¹⁵ Moreover, if the data are normalized with respect to
210 the data taken at 1000 cycles as opposed to the first cycle this
211 error drops to less than 1.5% (data not shown), implying that
212 there is slight settling that occurs over the first 1000 cycles.

213 The sensitivity of the Poisson capacitor lies very near to the
214 theoretical limit for an elastomeric parallel plate capacitor,
215 which predicts a gauge factor of 1. Here, the gauge factor is

defined as $(\Delta C/C_0)/\epsilon$. Calculating the mean gauge factor from 216
the data in Figure 2c, we obtain a gauge factor of 0.99 that is 217
uniform throughout the entire 100% strain range. A similar 218
piezocapacitor recently reported was tested to 30% strain with a 219
gauge factor of 0.4.²⁰ The only sensor to our knowledge that 220
explored similar strain (100+%) and cyclic testing relies on a 221
nanotube piezoresistor and has a gauge factor of 0.86 for strains 222
less than 40%, and 0.06 for strains greater than 60%.² 223
Additionally, strain gauge sensitivity can vary with temperature, 224
and we controlled for thermally induced expansion and 225
dielectric constant variation by conducting all experiments at 226
25 °C. However, we also directly calculated the sensor’s 227
sensitivity to temperature by using published values for silicone 228
material properties and adding both thermal expansion and 229
thermally induced dielectric constant variation into the Palmer 230
model.²⁹ These calculations suggest that the sensor’s capacitive 231
output will vary on the order of 0.01 pF/°C. Given that the 232
sensor operates on the order of 10 pF, this represents a 233
sensitivity of 0.1%/°C, meaning that it should be fairly resistant 234
to temperature variations. 235

As a demonstration of an alternative to traditional rigid 236
transducers and encoders for robotics, we built a proof-of- 237
concept for robotics applications where size, weight, and power 238
strongly constrain design options. Two such examples where 239
these constraints are crucial are the MEDIC and RoACH 240
(platforms of centimeter-scale walking and running robots). 241
Here, the exoskeleton of the robot consists entirely of origami- 242
style composite laminates that have been cut and folded to 243
produce from four to six “legs” in the form of four-bar 244
linkages.^{30–32} This process results in a strong and light robot 245
(RoACH is only 2.4 g) but also places a premium on sensing 246
and actuation components. At present, the legs are controlled 247
via open-loop, contractile shape memory alloy actuators and 248
passive return springs. Using a Poisson capacitor instead of a 249
return spring would additionally allow feedback control, in turn 250
allowing much more consistent locomotion and more complex 251
behaviors. 252

Using the same smart composite manufacturing process as 253
used with MEDIC and DASH, we built a scaled up version off 254
of leg linkage and attached the Poisson capacitor in the place of 255
the return spring. Supporting Information video S1 presents the 256
full range of motion of the linkage when coupled to the sensor 257
and stretching apparatus. Figure 3 presents stills of this motion 258
sequence above a plot mapping “limb angle” to relative change 259
in capacitance. To do this, the system was cycled through 80% 260
strain of the sensor (limited by the geometry of the linkage), 261
during which time the 4-bar linkage rotates the “limb” through 262
a wide arc while the sensor stretches. In essence, this allows us 263
to transduce changes in joint angles without relying on an angle 264
encoder or any other traditional, rigid sensing component such 265
as an linear variable differential transformer. While a simple 266
demonstration, this type of sensing can clearly be scaled to a 267
variety of different joints and linkages in a diverse array of 268
robotics platforms. 269

We have demonstrated a high-strain elastomeric, parallel- 270
plate capacitive strain gauge that relies on the Poisson effect to 271
translate uniaxial strain into a scaled deformation that brings 272
the two percolating electrodes closer together. By operating in 273
this fashion, we bypass a number of the typical problems 274
underlying percolating nanotube electrodes and piezoresistive 275
designs such as hysteresis and a variable gauge factor. The 276
device was fabricated using a novel hydrophobicity patterning 277
technique that can be performed rapidly (<20 min for full 278

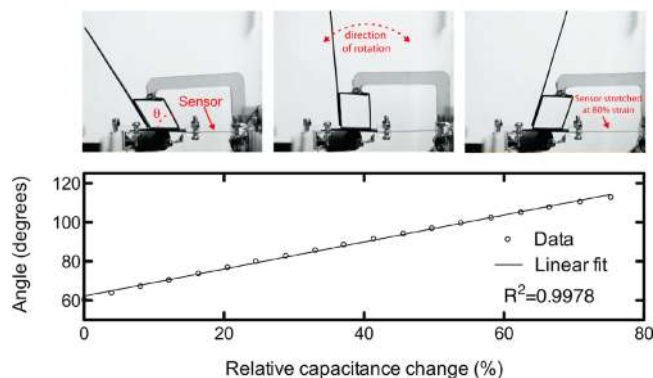


Figure 3. Demonstration of sensor as angle transducer for four-bar linkage. The image sequence is a time-lapse of the motion of a four-bar linkage limb. The sensor, attached to the base of the linkage, stretches in proportion to the amount by which the limb rotates. The plot shows the mapping between sensor output and angle of the limb. The “L”-shaped structure physically pulls on the mobile part of the four-bar linkage. See Supporting Information video S1 for the full range of motion.

assembly) with minimal equipment and at low cost (<\$1/sensor) making it an attractive option for anyone needing to rapidly prototype this type of sensor (or similar nanotube devices). Despite the anticipated, unstable resistive properties of our stretchable electrodes, the capacitive output of the sensor remains stable to within 3% over 3000 cycles. Further, the sensor exhibits the highest reported gauge factor ($GF = 0.99$) of any strain gauge capable of reversibly undergoing 100% strains or higher. Coupling this to the linear performance throughout the entire 100% strain cycle means that is readily calibrated and adaptable to practical sensing applications, as we have demonstrated by incorporating it into a typical robotic linkage as an alternative to an angle encoder.

ASSOCIATED CONTENT

Supporting Information

Characterization of the silicone mechanical properties, detailed description of the fabrication processes, discussion of the patterning technique, a discussion of the Palmer model for parallel plate capacitors, capacitance versus frequency response of the sensor, and a video of the four-bar linkage leg going through its full range of motion. This material is available free of charge via the Internet at <http://pubs.acs.org>.

AUTHOR INFORMATION

Corresponding Author

*E-mail: daniel.cohen@berkeley.edu.

Notes

The authors declare no competing financial interest.

ACKNOWLEDGMENTS

We thank Cary Pint, Nikta Fakhri, Nick Kohut, and Gabriel Lavella for their support and advice. Thanks also to Patrick Goodwill and Paul Lum for assistance with equipment. Funding for this project came from the National Defense Science and Engineering graduate fellowship (NDSEG) and the National Science Foundation.

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