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Doplicher/Roberts theorem on the
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A modified approach to the Doplicher/Roberts theorem on the construction of the field algebra and the symmetry group in superselection theory

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Abstract

Let \mathcal{A} be a unital C^* -algebra with trivial center $\mathcal{Z}(\mathcal{A}) = \mathbb{C}1$. Let \mathcal{T} denote a tensorial category of unital endomorphisms of \mathcal{A} equipped with several properties to be explained in the text.

Doplicher and Roberts have shown, among other things, that there is a C^* -algebra $\mathcal{F} \supset \mathcal{A}$ and a compact group \mathcal{G} of automorphisms of \mathcal{F} such that \mathcal{F} is a Hilbert C^* -system over \mathcal{A} w.r.t. \mathcal{G} where \mathcal{A} is the fixed point algebra w.r.t. \mathcal{G} , $\mathcal{G} = \text{stab } \mathcal{A}$ and the objects $\rho \in \text{Ob } \mathcal{T}$ are characterized as the canonical endomorphisms of certain algebraic \mathcal{G} -invariant Hilbert spaces $\mathcal{H}_\rho \subset \mathcal{F}$, see Doplicher/Roberts [1, 2, 3].

The starting point of the approach presented in this paper to point out the mentioned result is an \mathcal{A} -leftmodule $\mathcal{F}_0 := \{\sum_{\rho,j} A_{\rho,j} \Phi_{\rho,j}\}$. ρ runs through a full system of irreducible and mutually disjoint objects of \mathcal{T} , $j = 1, 2, \dots, d(\rho)$, where $d(\rho)$ denotes the statistical dimension of ρ , $\{\Phi_{\rho,j}\}_{j=1}^{d(\rho)}$ is an orthonormal basis of a $d(\rho)$ -dimensional Hilbert space. The system $\{\Phi_{\rho,j}\}_{\rho,j}$ forms a leftmodule basis of \mathcal{F}_0 , the coefficients $A_{\rho,j}$ are members of \mathcal{A} . The strategy is to equip successively \mathcal{F}_0 with a bimodule structure, a product and a $*$ -structure and finally with a C^* -norm $\|\cdot\|_*$. The symmetry group \mathcal{G} appears as the group of all automorphisms of the $*$ -algebra \mathcal{F}_0 leaving the \mathcal{A} -scalar product $\langle F, G \rangle := \sum_{\rho,j} d(\rho)^{-1} A_{\rho,j} B_{\rho,j}^*$ invariant, where $F = \sum_{\rho,j} A_{\rho,j} \Phi_{\rho,j}$, $G = \sum_{\rho,j} B_{\rho,j} \Phi_{\rho,j}$. The field algebra is then given by $\mathcal{F} := \text{clo}_{\|\cdot\|_*} \mathcal{F}_0$.

1 Introduction

A few years ago Doplicher and Roberts established a new duality theory for compact groups which is a deepening of the well-known duality theory of Tannaka/Krein. They developed this theory in connection with the superselection problem of the algebraic quantum field theory (see Doplicher/Roberts [4, 5, 1, 2, 3]). In the algebraic quantum field theory the main part of the superselection problem is to construct an extension $\mathcal{F} \supset \mathcal{A}$ of the quasilocal observable algebra \mathcal{A} , called field algebra, which should describe the sector structure of \mathcal{A} and the transition rules by a symmetry principle

(e.g. a compact group \mathcal{G}) on \mathcal{F} such that, for example, the fusion rules are explained in group-theoretical terms of \mathcal{G} .

In a recent paper [6] (see also Baumgärtel/Wollenberg [7]) the author described a rather simple construction of the field algebra \mathcal{F} as a Hilbert C^* -system over \mathcal{A} (or a crossed product of \mathcal{A} with a group dual of \mathcal{G}), if the group \mathcal{G} is known and if an action of a group dual of \mathcal{G} on \mathcal{A} is given a priori.

The present paper is an attempt to use the method of [6] for a construction of \mathcal{F} and of the group \mathcal{G} (which appears as the invariance group of an \mathcal{A} -scalar product). An essential point of this ansatz is to construct Hilbert spaces \mathcal{H}_α in the \mathcal{A} -bimodule \mathcal{F}_0 to start with, where the Hilbert spaces \mathcal{H}_α correspond bijectively to the objects α of the category \mathcal{T} of endomorphisms of \mathcal{A} .

This approach shows directly the intrinsic structure of \mathcal{F} and \mathcal{G} . On the other hand, it is not (completely) independent of Doplicher/Roberts deep duality theory. At a crucial point in the treatment (introduction of the product structure in \mathcal{F}_0) the author used an important result of Doplicher/Roberts ([2,4.17 Theorem, p.191]) to proceed further.

2 Preliminaries

Let \mathcal{A} be a unital C^* -algebra with center $\mathcal{Z}(\mathcal{A}) = \mathbf{C}1$. By ρ, σ, \dots we denote unital endomorphisms of \mathcal{A} . We use the denotation

$$(\rho, \sigma) := \{T \in \mathcal{A} : T\rho(A) = \sigma(A)T, A \in \mathcal{A}\}$$

for the linear space of all intertwiners between ρ and σ . The identical endomorphism is denoted by ι . If $R \in (\rho, \rho'), S \in (\sigma, \sigma')$ then $R\rho(S) = \rho'(S)R \in (\rho\sigma, \rho'\sigma')$. We use the notation $R \times S := R\rho(S)$. Note that $(\iota, \iota) = \mathbf{C}1$ and $1 \in (\rho, \rho)$ for all ρ , this is indicated by $1 = 1_\rho$.

We consider categories \mathcal{T} of unital endomorphisms equipped with the following properties:

(T1) $\iota \in \text{Ob } \mathcal{T}$.

(T2) If $\rho \in \text{Ob } \mathcal{T}$ and $U \in (\rho, \sigma)$, U unitary, then $\sigma \in \text{Ob } \mathcal{T}$.

(T3) $\rho, \sigma \in \text{Ob } \mathcal{T}$ implies $\rho\sigma \in \text{Ob } \mathcal{T}$.

(T4) The set of arrows between ρ and σ is given by (ρ, σ) .

(T5) If $\rho \in \text{Ob } \mathcal{T}$ and if $0 \neq E \in (\rho, \rho)$ is a selfadjoint projection then there is an isometry $V \in \mathcal{A}$, $VV^* = E$ such that $\sigma \in \text{Ob } \mathcal{T}$ where $\sigma(A) := V^*\rho(A)V$, i.e. $\text{Ob } \mathcal{T}$ is closed w.r.t. subobjects.

(T6) If $\rho, \sigma \in \text{Ob } \mathcal{T}$ then there are isometries $V, W \in \mathcal{A}$ with $VV^* + WW^* = 1$ such that $\tau \in \text{Ob } \mathcal{T}$ where $\tau(A) := V\rho(A)V^* + W\sigma(A)W^*$, i.e. $\text{Ob } \mathcal{T}$ is closed w.r.t. direct sums.

(T7) There is a *permutation structure* on \mathcal{T} , i.e. there is an assignment $\{\rho, \sigma\} \rightarrow \epsilon(\rho, \sigma) \in (\rho\sigma, \sigma\rho)$, $\epsilon(\rho, \sigma)$ unitary, with the following properties:

$$\epsilon(\sigma, \rho)\epsilon(\rho, \sigma) = 1_{\rho\sigma},$$

$$\begin{aligned}
\epsilon(\iota, \rho) &= \epsilon(\rho, \iota) = 1_\rho, \\
\epsilon(\rho\sigma, \tau) &= \epsilon(\rho, \tau) \times 1_\sigma \cdot 1_\rho \times \epsilon(\sigma, \tau), \\
\epsilon(\rho', \sigma') \cdot R \times S &= S \times R \cdot \epsilon(\rho, \sigma), \quad R \in (\rho, \rho'), \quad S \in (\sigma, \sigma').
\end{aligned}$$

(T8) There is a *conjugation structure* on \mathcal{T} , i.e. there is an assignment $\text{Ob } \mathcal{T} \ni \rho \rightarrow \bar{\rho} \in \text{Ob } \mathcal{T}$ such that there are corresponding arrows $R_\rho \in (\iota, \bar{\rho}\rho)$, $S_\rho \in (\iota, \rho\bar{\rho})$ with the following properties:

$$\begin{aligned}
S_\rho &:= \epsilon(\bar{\rho}, \rho)R_\rho, \\
S_\rho^* \times 1_\rho \cdot 1_\rho \times R_\rho &= 1_\rho, \\
R_\rho^* \times 1_{\bar{\rho}} \cdot 1_{\bar{\rho}} \times S_\rho &= 1_{\bar{\rho}}.
\end{aligned}$$

The endomorphism ρ is called *irreducible* if $(\rho, \rho) = \mathbf{C}1$; ρ and σ are called *disjoint* if $(\rho, \sigma) = \{0\}$, they are called *unitarily equivalent* if there is a unitary $U \in (\rho, \sigma)$.

The theory of \mathcal{T} is well-known, it is developed in Doplicher/Roberts [2,p.164 f.]. We recall several properties: $\bar{\rho}$ is unique up to unitary equivalence, the statistical dimension $d(\rho) \in \mathbf{N}$ given by $d(\rho)1 = R_\rho^*R_\rho$ satisfies $d(\rho_1\rho_2) = d(\rho_1)d(\rho_2)$, $d(W_1\rho_1(\cdot)W_1^* + W_2\rho_2(\cdot)W_2^*) = d(\rho_1) + d(\rho_2)$. If ρ_1, ρ_2 are unitarily equivalent then $d(\rho_1) = d(\rho_2)$, furthermore $d(\rho) = d(\bar{\rho})$. If $0 \neq E \in (\rho, \rho)$ is a selfadjoint projection then $d(V^*\rho(\cdot)V) = R_{\bar{\rho}E}^*R_\rho$, where $VV^* = E$.

Each ρ can be decomposed into a finite direct sum of irreducibles, $\rho(\cdot) = \sum_{j=1}^r W_j\rho_j(\cdot)W_j^*$, $W_j^*W_k = \delta_{jk}$, $\sum_j W_jW_j^* = 1$, ρ_j irreducible; (ρ, ρ) is finite-dimensional. If τ is irreducible, ρ arbitrary then (τ, ρ) is a finite-dimensional algebraic Hilbert space in \mathcal{A} , for $A, B \in (\tau, \rho)$ the scalar product (A, B) is given by $(A, B)1 = A^*B$; $\dim(\tau, \rho) =: m(\tau, \rho)$ coincides with the multiplicity of the equivalence class of τ occurring in the decomposition of ρ , i.e.

$$d(\rho) = \sum_{\tau} m(\tau, \rho)d(\tau).$$

For later use recall that $\rho \in \mathcal{T}$ is called *special* if $\det \rho = \iota$ (see [2,p.174] or [7,p.271]).

3 The \mathcal{A} -bimodule \mathcal{F}_0 and the Hilbert spaces \mathcal{H}_ρ

We choose a full system $\hat{\mathcal{T}}$ of irreducible and mutually disjoint endomorphisms $\tau \in \text{Ob } \mathcal{T}$. To each τ we assign a $d(\tau)$ -dimensional Hilbert space \mathcal{H}_τ , $\mathcal{H}_\iota := \mathbf{C}1$. We choose an orthonormal basis $\Phi_{\tau j}$, $j = 1, 2, \dots, d(\tau)$, of \mathcal{H}_τ , $\Phi_\iota := 1$, and we introduce the \mathcal{A} -leftmodule

$$\mathcal{F}_0 := \left\{ \sum_{\tau, j} A_{\tau j} \Phi_{\tau j}, \quad A_{\tau j} \in \mathcal{A}, \text{ finite sum} \right\}$$

where $\{\Phi_{\tau j}\}_{\tau, j}$ is assumed to be a leftmodule basis (then each other system $\{\Psi_{\tau j}\}_{\tau, j}$ of orthonormal bases of \mathcal{H}_τ turns out to be a leftmodule basis). Note that $\mathcal{A} \subset \mathcal{F}_0$. \mathcal{F}_0 is a bimodule w.r.t. the definition

$$\Phi A = \tau(A)\Phi, \quad \Phi \in \mathcal{H}_\tau, \quad A \in \mathcal{A},$$

and extension by linearity to the whole \mathcal{F}_0 . Then the relative commutant $\mathcal{A}' \cap \mathcal{F}_0 := \{F \in \mathcal{F}_0 : AF = FA, A \in \mathcal{A}\}$ equals $\mathbf{C}1$.

Let $\rho \in \mathcal{T}$ and $\tau \in \hat{\mathcal{T}}$. We consider the linear subspace $\text{spa}\{(\tau, \rho)\mathcal{H}_\tau\} \subset \mathcal{F}_0$. Obviously, it is a Hilbert space which can be considered as the tensor product $(\tau, \rho) \otimes \mathcal{H}_\tau$ with the scalar product

$$(A\Phi, B\Psi) := (A, B)_{(\tau, \rho)}(\Phi, \Psi)_{\mathcal{H}_\tau}, \quad A, B \in (\tau, \rho); \Phi, \Psi \in \mathcal{H}_\tau.$$

The sum

$$\sum_{\tau} \text{spa}\{(\tau, \rho)\mathcal{H}_\tau\} =: \mathcal{H}_\rho$$

is always a direct sum. Therefore \mathcal{H}_ρ is a Hilbert space with the scalar product

$$\sum_{\tau} (X_\tau, Y_\tau), \quad X_\tau, Y_\tau \in \text{spa}\{(\tau, \rho)\mathcal{H}_\tau\}.$$

The Hilbert spaces \mathcal{H}_τ satisfy the following properties:

3.1.PROPOSITION. *The assignment*

$$\text{Ob } \mathcal{T} \ni \rho \rightarrow \mathcal{H}_\rho \subset \mathcal{F}_0$$

is one-to-one, $X \in \mathcal{H}_\rho$ iff $XA = \rho(A)X$ for all $A \in \mathcal{A}$ and the dimension of \mathcal{H}_ρ coincides with $d(\rho)$,

$$\dim \mathcal{H}_\rho = d(\rho), \quad \rho \in \mathcal{T}. \quad (1)$$

Proof. Let $X \in \mathcal{H}_\rho$. Then $X = \sum_{\tau} A_\tau \Phi_\tau$ with $A_\tau \in (\tau, \rho)$, $\Phi_\tau \in \mathcal{H}_\tau$. Therefore $XA = \sum_{\tau} A_\tau \Phi_\tau A = \sum_{\tau} A_\tau \tau(A) \Phi_\tau = \rho(A) \sum_{\tau} A_\tau \Phi_\tau = \rho(A)X$. Conversely, let $X \in \mathcal{F}_0$, i.e. $X = \sum_{\tau, j} A_{\tau j} \Phi_{\tau j}$, $A_{\tau j} \in \mathcal{A}$ and $XA = \rho(A)X$. This means

$$\sum_{\tau, j} A_{\tau j} \tau(A) \Phi_{\tau j} = \sum_{\tau, j} \rho(A) A_{\tau j} \Phi_{\tau j}, \quad A \in \mathcal{A},$$

or

$$A_{\tau j} \tau(A) = \rho(A) A_{\tau j}, \quad A \in \mathcal{A},$$

i.e. $A_{\tau j} \in (\tau, \rho)$.

If $\mathcal{H}_\rho = \mathcal{H}_\sigma =: \mathcal{H}$ then $\rho(A)X = \sigma(A)X$ follows for all $A \in \mathcal{A}$, $X \in \mathcal{H}$, in particular for $X \in \mathcal{H}_\tau$ where $(\tau, \rho) \neq \{0\}$, i.e. $\rho(A)\Phi_{\tau j} = \sigma(A)\Phi_{\tau j}$ or $\rho(A) = \sigma(A)$ for all $A \in \mathcal{A}$, i.e. $\rho = \sigma$. Finally, we have

$$\begin{aligned} \dim \mathcal{H}_\rho &= \sum_{\tau} \dim(\text{spa}\{(\tau, \rho)\mathcal{H}_\tau\}) = \sum_{\tau} \dim(\tau, \rho) \dim \mathcal{H}_\tau = \\ &= \sum_{\tau} m(\tau, \rho) d(\tau) = d(\rho). \quad \square \end{aligned}$$

The arrows from (ρ, σ) can be interpreted as linear operators from $\mathcal{L}(\mathcal{H}_\rho \rightarrow \mathcal{H}_\sigma)$ by the assignment

$$(\rho, \sigma) \ni A \rightarrow \phi(A) \in \mathcal{L}(\mathcal{H}_\rho \rightarrow \mathcal{H}_\sigma) : \phi(A)X := AX, \quad X \in \mathcal{H}_\rho. \quad (2)$$

Note that the assignment (2) is injective and that $\phi(A^*)$ is the Hilbert space adjoint of $\phi(A)$:

3.2.PROPOSITION. *Let $A_1, A_2 \in (\rho, \sigma)$. If $\phi(A_1) = \phi(A_2)$ then $A_1 = A_2$. Furthermore,*

$$\phi(A^*) = \phi(A)^*, \quad A \in (\rho, \sigma).$$

Proof. According to the assumption we have $A_1X = A_2X$ for all $X \in \mathcal{H}_\rho$. This implies $A_1Y = A_2Y$ for all $Y \in (\tau, \rho)$, $\tau \in \hat{\mathcal{T}}$, i.e. $(A_1 - A_2)Y = 0$. Choose an orthonormal basis $W_{\tau j} \in (\tau, \rho)$, $j = 1, 2, \dots, m(\tau, \rho)$. Then $\sum_j W_{\tau j} W_{\tau j}^* = E_\tau \in (\rho, \rho)$, where E_τ is the central projection assigned to τ . Recall that $\sum_\tau E_\tau = 1$. So first we obtain $(A_1 - A_2)E_\tau = 0$ for all τ occurring in the decomposition of ρ and second $A_1 - A_2 = 0$. Furthermore, with $A \in (\rho, \sigma)$, $A_1 \in (\tau, \sigma)$, $A_2 \in (\tau, \rho)$ we have $(A_1, AA_2)_{(\tau, \sigma)} = A_1^* AA_2 = (A^* A_1)^* A_2 = (A^* A_1, A_2)_{(\tau, \rho)}$. \square

4 Construction of special isometries (cocycles)

$J(\rho, \sigma) : \mathcal{H}_\rho \otimes \mathcal{H}_\sigma \rightarrow \mathcal{H}_{\rho\sigma}$ and the product structure in \mathcal{F}_0

From (1) and $d(\rho\sigma) = d(\rho)d(\sigma)$ we get immediately $\dim \mathcal{H}_\rho \otimes \mathcal{H}_\sigma = \dim \mathcal{H}_{\rho\sigma}$, i.e. the Hilbert spaces $\mathcal{H}_\rho \otimes \mathcal{H}_\sigma$ and $\mathcal{H}_{\rho\sigma}$ are isomorphic. Within the set of all isometries from $\mathcal{H}_\rho \otimes \mathcal{H}_\sigma$ onto $\mathcal{H}_{\rho\sigma}$ we construct special ones (cocycles) $J(\rho, \sigma)$, equipped with the following properties:

$$(C1) \quad J(\rho, \iota) = J(\iota, \rho) = \text{id}_{\mathcal{H}_\rho},$$

$$(C2) \quad J(\rho\sigma, \tau)(J(\rho, \sigma) \otimes 1_{\mathcal{H}_\tau}) = J(\rho, \sigma\tau)(1_{\mathcal{H}_\rho} \otimes J(\sigma, \tau)),$$

$$(C3) \quad J(\rho', \sigma')\phi(A) \otimes \phi(B) = \phi(A \times B)J(\rho, \sigma), \quad A \in (\rho, \rho'), B \in (\sigma, \sigma').$$

$$(C4) \quad \phi(\epsilon(\rho, \sigma))J(\rho, \sigma) = J(\sigma, \rho)\Theta(\mathcal{H}_\rho, \mathcal{H}_\sigma),$$

where $\Theta(\mathcal{H}_\rho, \mathcal{H}_\sigma)$ denotes the flip isomorphism $\mathcal{H}_\rho \otimes \mathcal{H}_\sigma \rightarrow \mathcal{H}_\sigma \otimes \mathcal{H}_\rho$, and where $\epsilon(\cdot, \cdot)$ are the unitaries of the permutation structure in \mathcal{T} .

Note that one can introduce an equivalence relation within the set of all isometries $J(\rho, \sigma)$ satisfying (C1)-(C3) by

$$J' \sim J \quad \text{if} \quad J'(\rho, \sigma) = V(\rho\sigma)^{-1} J(\rho, \sigma) V(\rho) \otimes V(\sigma)$$

where $\rho \rightarrow V(\rho) \in \mathcal{L}(\mathcal{H}_\rho)$ is a unitary with the property

$$\phi(A)V(\rho) = V(\sigma)\phi(A), \quad A \in (\rho, \sigma).$$

It turns out that $J(\sigma, \rho)\Theta(\mathcal{H}_\rho, \mathcal{H}_\sigma)J(\rho, \sigma)^{-1}$ is an invariant of the equivalence classes and this invariant satisfies the property (T7) of the permutation structure.

We focus on the construction of $J(\rho, \sigma)$ if $\epsilon(\cdot, \cdot)$ is given. We separate three parts of the construction problem:

First we choose a special $\rho \in \mathcal{T}$, we consider the subcategory consisting of all powers ρ^n , $n = 0, 1, 2, \dots$ and we construct $J(\rho^n, \rho^m)$, $n, m = 0, 1, 2, \dots$. At this first stage we use an important result of Doplicher/Roberts [2,4.17 Theorem, p.191].

Second we construct corresponding isometries $J(\alpha, \beta)$ for all $\alpha, \beta \in \mathcal{T}$ which are *dominated* by a special (fixed) ρ , i.e. α, β are assumed to be subobjects of direct sums of powers of ρ .

Third, using the second step, we construct a *unique* function $\mathcal{T} \times \mathcal{T} \ni \{\alpha, \beta\} \rightarrow J(\alpha, \beta) \in \mathcal{L}(\mathcal{H}_\alpha \otimes \mathcal{H}_\beta \rightarrow \mathcal{H}_{\alpha\beta})$ satisfying (C1)-(C4).

The first step: Let $\rho \in \mathcal{T}$ be special. We look for auxiliary isometries Φ_n from \mathcal{H}_{ρ^n} onto \mathcal{H}_ρ^n , equipped with several properties, where \mathcal{H}_ρ^n denotes the n -fold tensor product of \mathcal{H}_ρ .

4.1.PROPOSITION. *Let $\rho \in \mathcal{T}$ be special. Then there is a system $\{\Phi_r\}_{r=1}^\infty$ of isometries $\Phi_r : \mathcal{H}_{\rho^r} \rightarrow \mathcal{H}_\rho^r$ such that the isomorphisms $\beta_r^s : \mathcal{L}(\mathcal{H}_{\rho^r} \rightarrow \mathcal{H}_{\rho^s}) \rightarrow \mathcal{L}(\mathcal{H}_\rho^r \rightarrow \mathcal{H}_\rho^s)$ given by*

$$\beta_r^s(X) := \Phi_s \circ X \circ \Phi_r^{-1}, \quad X \in \mathcal{L}(\mathcal{H}_{\rho^r} \rightarrow \mathcal{H}_{\rho^s}), \quad (3)$$

satisfies

$$(\beta_{r+s}^{r'+s'} \circ \phi)(A \times B) = (\beta_r^{r'} \circ \phi)(A) \otimes (\beta_s^{s'} \circ \phi)(B) \quad (4)$$

for all $A \in (\rho^r, \rho^{r'})$, $B \in (\rho^s, \rho^{s'})$. Moreover,

$$(\beta_n^n \circ \phi)(\epsilon_\rho^{(n)}(\pi)) = \Theta_\rho^{(n)}(\pi), \quad \pi \in \mathbf{P}_n, \quad n = 1, 2, \dots \quad (5)$$

holds, i.e.

$$\Phi_n \circ \phi(\epsilon_\rho^{(n)}(\pi)) = \Theta_\rho^{(n)}(\pi) \circ \Phi_n, \quad \pi \in \mathbf{P}_n.$$

Proof. According to Doplicher/Roberts [2,4.17 Theorem, p.191] there is a closed subgroup \mathcal{G} of $SU(\mathcal{H}_\rho)$ and a unital monomorphism α of a graded C^* -algebra $C^*(\rho)$, constructed by the intertwiner spaces (ρ^r, ρ^s) , which maps $C^*(\rho)$ into the Cuntz algebra $\mathcal{C}_{\mathcal{H}_\rho}$ generated by \mathcal{H}_ρ . One has $\alpha(C^*(\rho)) = \mathcal{C}_\mathcal{G}$ where $\mathcal{C}_\mathcal{G}$ denotes the fixed point algebra of $\mathcal{C}_{\mathcal{H}_\rho}$ under the canonical action of \mathcal{G} and

$$\alpha|_{(\rho^n, \rho^m)} = \mathcal{L}_\mathcal{G}(\mathcal{H}_\rho^n \rightarrow \mathcal{H}_\rho^m).$$

Note that $\mathcal{L}(\mathcal{H}_\rho^n \rightarrow \mathcal{H}_\rho^m)$ is interpreted, as usual, as a subspace of $\mathcal{C}_{\mathcal{H}_\rho}$ and $\mathcal{L}_\mathcal{G}(\mathcal{H}_\rho^n \rightarrow \mathcal{H}_\rho^m)$ denotes the intersection of $\mathcal{C}_\mathcal{G}$ and $\mathcal{L}(\mathcal{H}_\rho^n \rightarrow \mathcal{H}_\rho^m)$, i.e. the set of all intertwiners w.r.t. the (tensorial) action of \mathcal{G} on \mathcal{H}_ρ^n resp. \mathcal{H}_ρ^m . Note further that α satisfies the relations

$$\alpha(A \times B) = \alpha(A) \otimes \alpha(B), \quad A \in (\rho^r, \rho^{r'}), \quad B \in (\rho^s, \rho^{s'}), \quad (6)$$

and

$$\alpha(\epsilon_\rho(\pi)) = \Theta_{\mathcal{H}_\rho}(\pi), \quad \pi \in \mathbf{P}_\infty. \quad (7)$$

Now we form the Fock spaces

$$\mathcal{F} := \bigoplus_{n=0}^\infty \mathcal{H}_{\rho^n}, \quad \mathcal{H}_\rho^0 := \mathbf{C}; \quad \mathcal{F}_1 := \bigoplus_{n=0}^\infty \mathcal{H}_\rho^n, \quad \mathcal{H}_\rho^0 := \mathbf{C}, \quad (\text{Hilbert sums}).$$

By E_n we denote the projection of \mathcal{F} onto its component \mathcal{H}_{ρ^n} ; correspondingly F_n is the projection of \mathcal{F}_1 onto its component \mathcal{H}_ρ^n .

We consider the v. Neumann algebras $\mathcal{M}, \mathcal{M}_1$ on \mathcal{F} resp. \mathcal{F}_1 , generated by $\phi((\rho^n, \rho^m)) \in \mathcal{L}(\mathcal{H}_{\rho^n} \rightarrow \mathcal{H}_{\rho^m})$ resp. by $\mathcal{L}_\mathcal{G}(\mathcal{H}_\rho^n \rightarrow \mathcal{H}_\rho^m)$, i.e.

$$\mathcal{M} := \left(\bigcup_{n,m=0}^\infty E_m \phi((\rho^n, \rho^m)) E_n \right)'' ,$$

$$\mathcal{M}_1 := \left(\bigcup_{n,m=0}^{\infty} F_m \mathcal{L}_{\mathcal{G}}(\mathcal{H}_{\rho^n} \rightarrow \mathcal{H}_{\rho^m}) F_n \right)''.$$

Then the Doplicher/Roberts-isomorphism α induces an isomorphism β between the algebras \mathcal{M} and \mathcal{M}_1 , $\beta : \mathcal{M} \rightarrow \mathcal{M}_1$, where

$$\beta(E_m \phi(A) E_n) = F_m \alpha(A) F_n, \quad A \in (\rho^n, \rho^m),$$

(see for example Kadison/Ringrose [9, p.427]). It has the property

$$\beta(E_m X E_n) = F_m \beta(X) F_n, \quad X \in \mathcal{M}.$$

Obviously, $\mathcal{M}_1 = U(\mathcal{G})'$ where $U(\mathcal{G})$ denotes the canonical representation of \mathcal{G} on \mathcal{F}_1 given by the identical representation on \mathcal{H}_{ρ} : Note that $\mathcal{F}_1 \ni X = \sum_{n,m=0}^{\infty} F_m X F_n$ (strong convergence), hence $X \in U(\mathcal{G})'$ implies $F_m X F_n U(g) = U(g) F_m X F_n$ because $F_n U(g) = U(g) F_n$, for all $g \in \mathcal{G}$. That is, we have $F_m X F_n | \mathcal{H}_{\rho^n} \in \mathcal{L}_{\mathcal{G}}(\mathcal{H}_{\rho^n} \rightarrow \mathcal{H}_{\rho^m})$ or $F_m X F_n = F_m Y F_n$ with $Y \in \mathcal{L}_{\mathcal{G}}(\mathcal{H}_{\rho^n} \rightarrow \mathcal{H}_{\rho^m})$. The other inclusion is trivial.

Thus \mathcal{M}_1 is of type I, therefore, as is well-known (see e.g. [9, p.660ff.]) the isomorphism β is spatial, i.e. there is an isometry $\Phi : \mathcal{F} \rightarrow \mathcal{F}_1$ such that $\beta(X) = \Phi X \Phi^{-1}$, $X \in \mathcal{M}$. Now we have $\beta(E_n) = F_n \beta(E_n) F_n$, i.e. $F_n \beta(E_n) = \beta(E_n) F_n = \beta(E_n)$ or

$$F_n \cdot \Phi E_n \Phi^{-1} = \Phi E_n \Phi^{-1} \cdot F_n = \Phi E_n \Phi^{-1}.$$

Since $\dim \Phi E_n \Phi^{-1} = \dim E_n = \dim F_n$ we obtain $F_n = \Phi E_n \Phi^{-1}$ or $\Phi E_n = F_n \Phi$. This means that Φ is diagonal,

$$\Phi = \text{diag}\{\Phi_n\}_{n=0}^{\infty}$$

where $\Phi_n : \mathcal{H}_{\rho^n} \rightarrow \mathcal{H}_{\rho^n}$, $n = 0, 1, 2, \dots$, is isometric.

Therefore, the isomorphisms

$$\beta_n^m := \Phi_m \phi(A) \Phi_n^{-1}, \quad A \in (\rho^n, \rho^m)$$

satisfy the equation

$$\beta_n^m(\phi(A)) = \alpha(A), \quad A \in (\rho^n, \rho^m),$$

Now from (7) we obtain (5). Furthermore, from (6) we obtain (4): Namely, $A \in (\rho^r, \rho^{r'})$, $B \in (\rho^s, \rho^{s'})$ implies $A \times B \in (\rho^{r+s}, \rho^{r'+s'})$, i.e. we have $\phi(A) \in \mathcal{L}(\mathcal{H}_{\rho^r} \rightarrow \mathcal{H}_{\rho^{r'}})$, $\phi(B) \in \mathcal{L}(\mathcal{H}_{\rho^s} \rightarrow \mathcal{H}_{\rho^{s'}})$, $\phi(A \times B) \in \mathcal{L}(\mathcal{H}_{\rho^{r+s}} \rightarrow \mathcal{H}_{\rho^{r'+s'}})$, i.e.

$$\alpha(A \times B) = \beta_{r+s}^{r'+s'}(\phi(A \times B))$$

and

$$\alpha(A) \otimes \alpha(B) = \beta_r^{r'}(\phi(A)) \otimes \beta_s^{s'}(\phi(B)). \quad \square$$

It would be nice to have a proof for the existence of the β_n^m which is independent of the mentioned theorem of Doplicher/Roberts.

Using the isometries Φ_n , we define isometries from $\mathcal{H}_{\rho^r} \otimes \mathcal{H}_{\rho^s}$ onto $\mathcal{H}_{\rho^{r+s}}$ by

$$J(\rho^r, \rho^s) := \Phi_{r+s}(\Phi_r \otimes \Phi_s)^{-1}.$$

Then we obtain

4.2.PROPOSITION. *The isometries $J(\rho^r, \rho^s)$ satisfy (C1)-(C4).*

Proof. (C1) is obvious. (C2): We have to prove

$$J(\rho^{r+n}, \rho^s)(J(\rho^r, \rho^n) \otimes 1_{\mathcal{H}_{\rho^s}}) = J(\rho^r, \rho^{n+s})(1_{\mathcal{H}_{\rho^r}} \otimes J(\rho^n, \rho^s)).$$

By the definition of $J(\cdot, \cdot)$ we have

$$\Phi_{r+n+s} = J(\rho^{r+n}, \rho^s)\Phi_{r+n} \otimes \Phi_s = J(\rho^r, \rho^{n+s})(\Phi_r \otimes \Phi_{n+s}) \quad (8)$$

and

$$\begin{aligned} J(\rho^{r+n}, \rho^s)\Phi_{r+n} \otimes \Phi_s &= J(\rho^{r+n}, \rho^s)(\{J(\rho^r, \rho^n)(\Phi_r \otimes \Phi_n)\} \otimes \Phi_s) = \\ &= J(\rho^{r+n}, \rho^s)(J(\rho^r, \rho^n) \otimes 1_{\mathcal{H}_{\rho^s}})(\Phi_r \otimes \Phi_n \otimes \Phi_s). \end{aligned}$$

For the right hand side of (8) we get similarly

$$J(\rho^r, \rho^{n+s})(\Phi_r \otimes \{J(\rho^n, \rho^s)(\Phi_n \otimes \Phi_s)\}) = J(\rho^r, \rho^{n+s})(1_{\mathcal{H}_{\rho^r}} \otimes J(\rho^n, \rho^s))(\Phi_r \otimes \Phi_n \otimes \Phi_s).$$

(C3): For brevity we write simply β instead of β_r^r etc. For $A \in (\rho^r, \rho^{r'})$, $B \in (\rho^s, \rho^{s'})$ we have

$$\begin{aligned} J(\rho^{r'}, \rho^{s'})\phi(A) \otimes \phi(B) &= \Phi_{r'+s'}\{(\Phi_{r'}^{-1}\phi(A)) \otimes (\Phi_{s'}^{-1}\phi(B))\} = \\ &= \Phi_{r'+s'}\{(\beta \circ \phi)(A)\Phi_r^{-1} \otimes (\beta \circ \phi)(B)\Phi_s^{-1}\} = \Phi_{r'+s'}((\beta \circ \phi)(A) \otimes (\beta \circ \phi)(B))\Phi_r^{-1} \otimes \Phi_s^{-1} = \\ &= \Phi_{r'+s'}((\beta \circ \phi)(A \times B)\Phi_r^{-1} \otimes \Phi_s^{-1}) = \phi(A \times B)\Phi_{r+s}(\Phi_r^{-1} \otimes \Phi_s^{-1}) = \phi(A \times B)J(\rho^r, \rho^s). \end{aligned}$$

Finally we prove (C4): We calculate

$$\begin{aligned} J(\rho^s, \rho^r)\Theta(\mathcal{H}_{\rho^r}, \mathcal{H}_{\rho^s}) &= \Phi_{r+s}(\Phi_s^{-1} \otimes \Phi_r^{-1})\Theta(\mathcal{H}_{\rho^r}, \mathcal{H}_{\rho^s}) = \Phi_{r+s}\Theta(\mathcal{H}_{\rho^r}^r, \mathcal{H}_{\rho^s}^s)(\Phi_r^{-1} \otimes \Phi_s^{-1}) = \\ &= \Phi_{r+s}\Theta_{\mathcal{H}_{\rho^s}^s}^{(r+s)}((r, s))(\Phi_r^{-1} \otimes \Phi_s^{-1}) = \phi(\epsilon^{(r+s)}((r, s)))\Phi_{r+s}(\Phi_r^{-1} \otimes \Phi_s^{-1}) = \\ &= \phi(\epsilon(\rho^r, \rho^s))J(\rho^r, \rho^s). \quad \square \end{aligned}$$

The second step: We have to construct $J(\alpha, \beta)$ for direct sums α, β of powers of ρ (where ρ is special) and then we have to verify the properties (C1)-(C4). The same procedure is to establish for subobjects of such direct sums. We begin with direct sums: According to their definition we have the following formulas for direct sums α, β, γ , where the isometries W_j, V_k, U_l have the properties mentioned above:

$$\alpha(\cdot) := \sum_j W_j \rho^{m_j}(\cdot) W_j^*, \quad \beta(\cdot) := \sum_k V_k \rho^{n_k}(\cdot) V_k^*, \quad \gamma(\cdot) := \sum_l U_l \rho^{r_l}(\cdot) U_l^*,$$

$$W_j \in (\rho^{m_j}, \alpha), \quad V_k \in (\rho^{n_k}, \beta), \quad U_l \in (\rho^{r_l}, \gamma),$$

$$\mathcal{H}_\alpha = \sum_j W_j \mathcal{H}_{\rho^{m_j}}, \quad \mathcal{H}_\beta = \sum_k V_k \mathcal{H}_{\rho^{n_k}}, \quad \mathcal{H}_\gamma = \sum_l U_l \mathcal{H}_{\rho^{r_l}}.$$

We put

$$J(\alpha, \beta) := \sum_{j,k} \phi(W_j \times V_k) J(\rho^{m_j}, \rho^{n_k}) \phi(W_j^*) \otimes \phi(V_k^*), \quad (9)$$

similarly for the other objects $J(\beta, \gamma)$, $J(\alpha\beta, \gamma)$ etc. Then we assert

4.3.LEMMA. *The isometries (9) satisfy the properties (C1)-(C4).*

Proof. (C1) is obvious. (C3): We have

$$J(\alpha', \beta') = \sum_{j', k'} \phi(\tilde{W}_{j'} \times \tilde{V}_{k'}) J(\rho^{\tilde{m}_{j'}}, \rho^{\tilde{n}_{k'}}) \phi(\tilde{W}_{j'}^*) \otimes \phi(\tilde{V}_{k'}^*)$$

where $\tilde{W}_{j'} \in (\rho^{\tilde{m}_{j'}}, \alpha')$, $\tilde{V}_{k'} \in (\rho^{\tilde{n}_{k'}}, \beta')$. Let $A \in (\alpha, \alpha')$, $B \in (\beta, \beta')$. Then

$$\tilde{W}_{j'}^* A W_j \in (\rho^{m_j}, \rho^{\tilde{m}_{j'}}), \quad \tilde{V}_{k'}^* B V_k \in (\rho^{n_k}, \rho^{\tilde{n}_{k'}}),$$

and we have, according to (C3) for the powers of ρ

$$J(\rho^{\tilde{m}_{j'}}, \rho^{\tilde{n}_{k'}}) \phi(\tilde{W}_{j'}^* A W_j) \otimes \phi(\tilde{V}_{k'}^* B V_k) = \phi(\tilde{W}_{j'}^* A W_j \times \tilde{V}_{k'}^* B V_k) J(\rho^{m_j}, \rho^{n_k}).$$

Now multiplying this equation from the left by $\phi(\tilde{W}_{j'} \times \tilde{V}_{k'})$, from the right by $\phi(W_j^*) \otimes \phi(V_k^*)$ and summing up w.r.t. j', k', j, k , we obtain

$$\begin{aligned} & \sum_{j', k'} \phi(\tilde{W}_{j'} \times \tilde{V}_{k'}) J(\rho^{\tilde{m}_{j'}}, \rho^{\tilde{n}_{k'}}) \phi(\tilde{W}_{j'}^*) \otimes \phi(\tilde{V}_{k'}^*) \cdot \phi(A) \otimes \phi(B) = \\ & = \phi(A \times B) \sum_{j, k} \phi(W_j \times V_k) J(\rho^{m_j}, \rho^{n_k}) \phi(W_j^*) \otimes \phi(V_k^*) \end{aligned}$$

or

$$J(\alpha', \beta') \phi(A) \otimes \phi(B) = \phi(A \times B) J(\alpha, \beta).$$

(C2): Note that

$$\alpha\beta(\cdot) = \sum_{j, k} W_j \times V_k \rho^{m_j+n_k}(\cdot) (W_j \times V_k)^*.$$

Then

$$J(\alpha\beta, \gamma) = \sum_{j, k, l} \phi((W_j \times V_k) \times U_l) J(\rho^{m_j+n_k}, \rho^{r_l}) \phi((W_j \times V_k)^*) \otimes \phi(U_l^*)$$

and we have to calculate

$$\begin{aligned} & J(\alpha\beta, \gamma)(J(\alpha, \beta) \otimes 1_{\mathcal{H}_\gamma}) = \\ & \sum_{j, k, l} \phi(W_j \times V_k \times U_l) J(\rho^{m_j+n_k}, \rho^{r_l}) \phi(W_j \times V_k)^* \otimes \phi(U_l)^* \cdot \\ & \cdot \sum_{j', k'} \phi(W_{j'} \times V_{k'}) J(\rho^{m_{j'}}, \rho^{n_{k'}}) (\phi(W_{j'}^*) \otimes \phi(V_{k'}^*)) \otimes 1_{\mathcal{H}_\gamma} \\ & = \sum_{j, k, l} \sum_{j', k'} \phi(W_j \times V_k \times U_l) J(\rho^{m_j+n_k}, \rho^{r_l}) \cdot \\ & \cdot (\phi((W_j \times V_k)^* W_{j'} \times V_{k'}) J(\rho^{m_{j'}}, \rho^{n_{k'}}) (\phi(W_{j'}^*) \otimes \phi(V_{k'}^*)) \otimes \phi(U_l)^*) \\ & = \sum_{j, k, l} \phi(W_j \times V_k \times U_l) J(\rho^{m_j+n_k}, \rho^{r_l}) (J(\rho^{m_j}, \rho^{n_k}) \otimes 1_{\mathcal{H}_{\rho^{r_l}}}) (\phi(W_j)^* \otimes \phi(V_k)^* \otimes \phi(U_l)^*). \end{aligned}$$

On the other hand,

$$J(\alpha, \beta\gamma)(1_{\mathcal{H}_\alpha} \otimes J(\beta, \gamma)) = \sum_{j,k,l} \phi(W_j \times V_k \times U_l) J(\rho^{m_j}, \rho^{n_k+r_l})(1_{\mathcal{H}_{\rho^{m_j}}} \otimes J(\rho^{n_k}, \rho^{r_l}))(\phi(W_j)^* \otimes \phi(V_k)^* \otimes \phi(U_l)^*).$$

Hence the validity of (C2) for powers of ρ implies its validity for direct sums. (C4): We assert that

$$J(\beta, \alpha)\Theta(\mathcal{H}_\alpha, \mathcal{H}_\beta) = \phi(\epsilon(\alpha, \beta))J(\alpha, \beta). \quad (10)$$

We calculate the left hand side: it equals

$$\begin{aligned} & \sum_{k,j} \phi(V_k \times W_j) J(\rho^{n_k}, \rho^{m_j}) \phi(V_k)^* \otimes \phi(W_j)^* \Theta(\mathcal{H}_\alpha, \mathcal{H}_\beta) = \\ & = \sum_{k,j} \phi(V_k \times W_j) J(\rho^{n_k}, \rho^{m_j}) \Theta(\mathcal{H}_{\rho^{m_j}}, \mathcal{H}_{\rho^{n_k}}) \phi(W_j)^* \otimes \phi(V_k)^* \\ & = \sum_{k,j} \phi(V_k \times W_j) \phi(\epsilon(\rho^{m_j}, \rho^{n_k})) J(\rho^{m_j}, \rho^{n_k}) \phi(W_j)^* \otimes \phi(V_k)^* \\ & = \sum_{k,j} \phi(\epsilon(\alpha, \beta)) \phi(W_j \times V_k) J(\rho^{m_j}, \rho^{n_k}) \phi(W_j)^* \otimes \phi(V_k)^* = \phi(\epsilon(\alpha, \beta)) J(\alpha, \beta), \end{aligned}$$

i.e. (10) is proved. \square

Now we turn to the case of subobjects of direct sums. Let α, β be direct sums and let α', β' denote subobjects of α, β , correspondingly. Then we have

$$\alpha'(\cdot) = W^* \alpha(\cdot) W, \quad \beta'(\cdot) = V^* \beta(\cdot) V, \quad W \in (\alpha', \alpha), \quad V \in (\beta', \beta), \quad W^* W = V^* V = 1.$$

Now we put

$$J(\alpha', \beta') := \phi(W \times V)^* J(\alpha, \beta) \phi(W) \otimes \phi(V). \quad (11)$$

Note that

$$\alpha' \beta'(\cdot) = W^* \times V^* (\alpha \beta)(\cdot) W \times V.$$

Then we have

4.4.LEMMA. *The isometries (11) satisfy the properties (C1)-(C4).*

Proof. (C1) is obvious. (C3): Let α'_1, β'_1 be subobjects of α_1, β_1 , respectively:

$$\alpha'_1(\cdot) = W_1^* \alpha_1(\cdot) W_1, \quad \beta'_1(\cdot) = V_1^* \beta_1(\cdot) V_1, \quad W_1 \in (\alpha'_1, \alpha_1), \quad V_1 \in (\beta'_1, \beta_1),$$

$$W_1^* W_1 = V_1^* V_1 = 1.$$

Choose $A \in (\alpha', \alpha'_1), B \in (\beta', \beta'_1)$. Then

$$W_1 A W^* \in (\alpha, \alpha_1), \quad V_1 B V^* \in (\beta, \beta_1)$$

and we have, according to Lemma 4.3.,

$$J(\alpha_1, \beta_1) \phi(W_1 A W^*) \otimes \phi(V_1 B V^*) = \phi(W_1 A W^* \times V_1 B V^*) J(\alpha, \beta).$$

Multiplying this equation from the left by $\phi(W_1^* \times V_1^*)$, and from the right by $\phi(W \otimes V)$ we obtain

$$\begin{aligned} \phi(W_1^* \times V_1^*)J(\alpha_1, \beta_1)\phi(W_1) \otimes \phi(V_1) \cdot \phi(A) \otimes \phi(B) &= \\ &= \phi(A \times B)\phi(W^* \times V^*)J(\alpha, \beta)\phi(W) \otimes \phi(V) \end{aligned}$$

or

$$J(\alpha'_1, \beta'_1)\phi(A) \otimes \phi(B) = \phi(A \times B)J(\alpha, \beta).$$

(C2): Let α', β' be as before and let γ' be a subobject of γ , i.e. $\gamma'(\cdot) = U^*\gamma(\cdot)U$, $U \in (\gamma', \gamma)$, $U^*U = 1$. We calculate

$$\begin{aligned} J(\alpha'\beta', \gamma')(J(\alpha', \beta') \otimes 1_{\mathcal{H}_{\gamma'}}) &= \\ \phi((W \times V) \times U)^*J(\alpha\beta, \gamma)(\phi(W \times V) \otimes \phi(U))(\phi(W \times V)^*J(\alpha, \beta)(\phi(W) \otimes \phi(V)) \otimes 1_{\mathcal{H}_{\gamma'}}) &= \\ = \phi(W \times V \times U)^*J(\alpha\beta, \gamma)(\phi((W \times V)(W \times V)^*)J(\alpha, \beta)(\phi(W) \otimes \phi(V)) \otimes \phi(U)) &= \\ = \phi(W \times V \times U)^*J(\alpha\beta, \gamma)(\phi(WW^* \times VV^*)J(\alpha, \beta)(\phi(W) \otimes \phi(V)) \otimes \phi(U)) &= \\ = \phi(W \times V \times U)^*J(\alpha\beta, \gamma)(J(\alpha, \beta) \otimes 1_{\mathcal{H}_{\gamma}})\phi(W) \otimes \phi(V) \otimes \phi(U) &= \\ = \phi(W \times V \times U)^*J(\alpha, \beta\gamma)(1_{\mathcal{H}_\alpha} \otimes J(\beta, \gamma))\phi(W) \otimes \phi(V) \otimes \phi(U) & \end{aligned}$$

because of Lemma 4.3. and finally this equals $J(\alpha', \beta'\gamma')(1_{\mathcal{H}_{\alpha'}} \otimes J(\beta', \gamma'))$.

(C4): We have to prove

$$J(\beta', \alpha')\Theta(\mathcal{H}_{\alpha'}, \mathcal{H}_{\beta'}) = \phi(\epsilon(\alpha', \beta'))J(\alpha', \beta').$$

So we calculate the left hand side, it equals

$$\begin{aligned} \phi(V \times W)^*J(\beta, \alpha)\phi(V) \otimes \phi(W)\Theta(\mathcal{H}_{\alpha'}, \mathcal{H}_{\beta'}) &= \phi(V \times W)^*J(\beta, \alpha)\Theta(\mathcal{H}_\alpha, \mathcal{H}_\beta)\phi(W) \otimes \phi(V) \\ &= \phi(V \times W)^*\phi(\epsilon(\alpha, \beta))J(\alpha, \beta)\phi(W) \otimes \phi(V) = \\ &= \phi(\epsilon(\alpha', \beta'))\phi(W^* \times V^*)J(\alpha, \beta)\phi(W) \otimes \phi(V) = \phi(\epsilon(\alpha', \beta'))J(\alpha', \beta'). \quad \square \end{aligned}$$

The third step consists in the construction of a unique function

$$\{\alpha, \beta\} \rightarrow J(\alpha, \beta), \quad \alpha, \beta \in \mathcal{T},$$

such that $J(\alpha, \beta)$ satisfies all the properties (C1)-(C4). First let α, β be fixed. We consider the system of *all* isometries Φ from $\mathcal{H}_\alpha \otimes \mathcal{H}_\beta$ onto $\mathcal{H}_{\alpha\beta}$. It is denoted by $M(\alpha, \beta)$. Obviously, $M(\alpha, \beta)$ is a compact space (essentially the unitary group of all $d(\alpha)d(\beta) \times d(\alpha)d(\beta)$ -matrices). Then the cartesian product

$$M := \prod_{\alpha, \beta \in \mathcal{T}} M(\alpha, \beta)$$

is a compact space w.r.t the Tychonov topology.

Let $\mathcal{T}_{fin} := \{\alpha_1, \alpha_2, \dots, \alpha_r\}$ be a finite set of objects. Let $M(\mathcal{T}_{fin}) \subset M$ denote the subset consisting of all functions $J(\rho, \sigma)$ such that (C1)-(C4) is satisfied if $\rho, \sigma \in \mathcal{T}_{fin}$. The foregoing considerations (the first and the second step) show that $M(\mathcal{T}_{fin}) \neq \emptyset$ for each \mathcal{T}_{fin} , because the elements of \mathcal{T}_{fin} are dominated by some special object ρ .

Note that the intersection $\bigcap_{\mathcal{T}_{fin} \in \mathcal{C}} M(\mathcal{T}_{fin})$ for a finite collection $\mathcal{C} = \{\mathcal{T}_{fin}\}$ is also nonempty because

$$\bigcap_{\mathcal{T}_{fin} \in \mathcal{C}} M(\mathcal{T}_{fin}) \supseteq M\left(\bigcup_{\mathcal{T}_{fin} \in \mathcal{C}} \mathcal{T}_{fin}\right)$$

and $\bigcup_{\mathcal{T}_{fin} \in \mathcal{C}} \mathcal{T}_{fin}$ is again a finite subset of objects. Therefore we obtain

4.5.LEMMA. *There is a function $J(\alpha, \beta)$ satisfying all the properties (C1)-(C4).*

Proof. It is obvious because, according to Cantor's theorem,

$$\bigcap_{\mathcal{T}_{fin}} M(\mathcal{T}_{fin}) \neq \emptyset$$

where now \mathcal{T}_{fin} runs through *all* finite subsets of objects. A member $J(\cdot, \cdot) \in \bigcap M(\mathcal{T}_{fin})$ satisfies the properties (C1)-(C4) for all $\alpha, \beta \in \mathcal{T}$. \square

We note that the argument using Cantor's theorem is already used by Doplicher/Roberts in the course of their treatment.

Now the goal of the second part of this section is the introduction of a product structure in \mathcal{F}_0 , using the special isometries $J(\alpha, \beta) : \mathcal{H}_\alpha \otimes \mathcal{H}_\beta \rightarrow \mathcal{H}_{\alpha\beta}$ constructed before.

Let $\Phi \in \mathcal{H}_\rho, \Psi \in \mathcal{H}_\sigma$, where ρ, σ are irreducible. Then we define

$$\Phi \cdot \Psi := J(\rho, \sigma)(\Phi \otimes \Psi). \quad (12)$$

The extension of this definition to arbitrary elements of \mathcal{F}_0 is given by the distributive law. A straightforward calculation shows then that (12) is also true for elements of arbitrary Hilbert spaces $\mathcal{H}_\alpha, \mathcal{H}_\beta$ (see e.g. [6,p.365]).

4.6.PROPOSITION. *The product in \mathcal{F}_0 , defined by (12) and by its extension by the distributive law, is associative.*

Proof. Obvious by the definition, because the tensor product is associative and (C2) holds. \square

We remark that corresponding formulas already appear in the treatment of Doplicher/Roberts: the formulas (SM1-SM3) and (6.2) in [2,p.303].

5 The $*$ -structure in the algebra \mathcal{F}_0

The introduction of a $*$ -structure in \mathcal{F}_0 consists of two steps. In the first one we introduce the concept of a *conjugated basis* in the Hilbert space $\mathcal{H}_{\bar{\rho}}$ where $\bar{\rho}$ denotes the representer of the *conjugated class* w.r.t. ρ .

Let $R_\rho \in (\iota, \bar{\rho}\rho)$ be a conjugate belonging to the pair $\rho, \bar{\rho}$. Then $R_\rho \in \mathcal{H}_{\bar{\rho}\rho}$ holds and therefore

$$R_\rho = J(\bar{\rho}, \rho)u, \quad u \in \mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_\rho.$$

We choose an orthonormal basis $\{\Phi_k\}_{k=1}^d$ in \mathcal{H}_ρ and an orthonormal basis $\{\Psi_j\}_{j=1}^d$ in $\mathcal{H}_{\bar{\rho}}$, where $d = d(\rho) = d(\bar{\rho})$. Then

$$u = \sum_{j,k} u_{jk} \Psi_j \otimes \Phi_k$$

where $\{u_{jk}\}$ is a $d \times d$ -matrix with complex entries. Now we have

5.1.LEMMA. *The $d \times d$ -matrix $\{u_{jk}\}$ is unitary.*

Proof. We use the relation

$$R_\rho^* \times 1_{\bar{\rho}} \cdot 1_{\bar{\rho}} \times S_\rho = 1_{\bar{\rho}}$$

where $S_\rho = \epsilon(\bar{\rho}, \rho)R_\rho$. We have $S_\rho = J(\rho, \bar{\rho})v$ with $v \in \mathcal{H}_\rho \otimes \mathcal{H}_{\bar{\rho}}$ and

$$J(\rho, \bar{\rho})\Theta(\mathcal{H}_{\bar{\rho}}, \mathcal{H}_\rho)u = \phi(\epsilon(\bar{\rho}, \rho))J(\bar{\rho}, \rho)u = \epsilon(\bar{\rho}, \rho)R_\rho = S_\rho,$$

therefore $v = \Theta(\mathcal{H}_{\bar{\rho}}, \mathcal{H}_\rho)u$. We have $\phi(1_{\bar{\rho}} \times S_\rho) \in \mathcal{L}(\mathcal{H}_{\bar{\rho}} \rightarrow \mathcal{H}_{\bar{\rho}\rho\bar{\rho}})$ hence $\phi(1_{\bar{\rho}} \times S_\rho)x_{\bar{\rho}} = 1_{\bar{\rho}} \times S_\rho \cdot x_{\bar{\rho}}$, $x_{\bar{\rho}} \in \mathcal{H}_{\bar{\rho}}$. Furthermore,

$$\phi(1_{\bar{\rho}} \times S_\rho)1_{\bar{\rho}} = \phi(1_{\bar{\rho}} \times S_\rho)J(\bar{\rho}, \iota) = J(\bar{\rho}, \rho\bar{\rho})\phi(1_{\bar{\rho}} \otimes \phi(S_\rho))$$

hence

$$\begin{aligned} \phi(1_{\bar{\rho}} \times S_\rho)x_{\bar{\rho}} &= J(\bar{\rho}, \rho\bar{\rho})x_{\bar{\rho}} \otimes S_\rho = J(\bar{\rho}, \rho\bar{\rho})(x_{\bar{\rho}} \otimes J(\rho, \bar{\rho})v) \\ &= J(\bar{\rho}, \rho\bar{\rho})(1_{\bar{\rho}} \otimes J(\rho, \bar{\rho}))x_{\bar{\rho}} \otimes v = J(\bar{\rho}\rho, \bar{\rho})(J(\bar{\rho}, \rho) \otimes 1_{\bar{\rho}})x_{\bar{\rho}} \otimes v. \end{aligned}$$

Now there is a representation $v = \sum_{j,k} v_{jk}\Phi_j \otimes \Psi_k$ and replacing v by this expression we obtain

$$\begin{aligned} \phi(1_{\bar{\rho}} \times S_\rho)x_{\bar{\rho}} &= J(\bar{\rho}\rho, \bar{\rho})(J(\bar{\rho}, \rho) \otimes 1_{\bar{\rho}})x_{\bar{\rho}} \otimes \sum_{j,k} v_{jk}\Phi_j \otimes \Psi_k \\ &= J(\bar{\rho}\rho, \bar{\rho})(J(\bar{\rho}, \rho) \otimes 1_{\bar{\rho}}) \sum_{j,k} v_{jk}x_{\bar{\rho}} \otimes \Phi_j \otimes \Psi_k \\ &= J(\bar{\rho}\rho, \bar{\rho}) \sum_{j,k} v_{jk}(J(\bar{\rho}, \rho)x_{\bar{\rho}} \otimes \Phi_j) \otimes \Psi_k \end{aligned}$$

and

$$\begin{aligned} \phi(R_\rho^* \times 1_{\bar{\rho}})\phi(1_{\bar{\rho}} \times S_\rho)x_{\bar{\rho}} &= \\ \phi(R_\rho^* \times 1_{\bar{\rho}})J(\bar{\rho}\rho, \bar{\rho}) \sum_{j,k} v_{jk}(J(\bar{\rho}, \rho)x_{\bar{\rho}} \otimes \Phi_j) \otimes \Psi_k &= \\ = J(\iota, \bar{\rho})\phi(R_\rho^*) \otimes \phi(1_{\bar{\rho}}) \sum_{j,k} v_{jk}(J(\bar{\rho}, \rho)x_{\bar{\rho}} \otimes \Phi_j) \otimes \Psi_k &= \\ = \phi(R_\rho^*) \otimes \phi(1_{\bar{\rho}}) \sum_{j,k} v_{jk}(J(\bar{\rho}, \rho)x_{\bar{\rho}} \otimes \Phi_j) \otimes \Psi_k. & \end{aligned}$$

Now for $x \in \mathcal{H}_{\bar{\rho}\rho}$, $y \in \mathcal{H}_{\bar{\rho}}$ we have

$$\phi(R_\rho^*) \otimes \phi(1_{\bar{\rho}})(x \otimes y) = \phi(R_\rho^*)x \otimes \phi(1_{\bar{\rho}})y = (R_\rho, x)_{\mathcal{H}_{\bar{\rho}\rho}}1 \otimes y = (R_\rho, x)_{\mathcal{H}_{\bar{\rho}\rho}}y$$

because $\phi(R_\rho^*)x = R_\rho^* \cdot x = (R_\rho, x)_{\mathcal{H}_{\bar{\rho}\rho}}1$. In particular, we put $x := J(\bar{\rho}, \rho)(x_{\bar{\rho}} \otimes \Phi_j)$, $y := \Psi_k$. Then we obtain

$$\phi(R_\rho^*) \otimes \phi(1_{\bar{\rho}})(J(\bar{\rho}, \rho)(x_{\bar{\rho}} \otimes \Phi_j)) \otimes \Psi_k = (R_\rho, J(\bar{\rho}, \rho)(x_{\bar{\rho}} \otimes \Phi_j))_{\mathcal{H}_{\bar{\rho}\rho}}\Psi_k =$$

$$\begin{aligned}
&= (J(\bar{\rho}, \rho)u, J(\bar{\rho}, \rho)(x_{\bar{\rho}} \otimes \phi_j))_{\mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_{\rho}} \Psi_k = (u, x_{\bar{\rho}} \otimes \Phi_j)_{\mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_{\rho}} \Psi_k = \\
&= \left(\sum_{j', k'} v_{j'k'} \Psi_{k'} \otimes \Phi_{j'}, x_{\bar{\rho}} \otimes \Phi_j \right)_{\mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_{\rho}} \Psi_k = \sum_{j', k'} \bar{v}_{j'k'}(\Psi_{k'}, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} (\Phi_{j'}, \Phi_j)_{\mathcal{H}_{\rho}} \Psi_k = \\
&= \sum_{k'} \bar{v}_{jk'}(\Psi_{k'}, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} \Psi_k
\end{aligned}$$

hence

$$\begin{aligned}
&\phi(R_{\rho}^* \times 1_{\bar{\rho}}) \phi(1_{\bar{\rho}} \times S_{\rho}) x_{\bar{\rho}} = \\
&\sum_{j, k} v_{jk} \phi(R_{\rho}^*) \otimes \phi(1_{\bar{\rho}}) (J(\bar{\rho}, \rho) x_{\bar{\rho}} \otimes \Phi_j) \otimes \Psi_k = \sum_{j, k} v_{jk} \left(\sum_{k'} \bar{v}_{jk'}(\Psi_{k'}, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} \right) \Psi_k = \\
&\sum_k \left(\sum_{j, k'} v_{jk} \bar{v}_{jk'}(\Psi_{k'}, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} \right) \Psi_k = x_{\bar{\rho}} = \sum_k (\Psi_k, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} \Psi_k.
\end{aligned}$$

This implies

$$\sum_{j, k'} v_{jk} \bar{v}_{jk'}(\Psi_{k'}, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}} = (\Psi_k, x_{\bar{\rho}})_{\mathcal{H}_{\bar{\rho}}}, \quad x_{\bar{\rho}} \in \mathcal{H}_{\bar{\rho}},$$

i.e.

$$\sum_{j, k'} v_{jk} \bar{v}_{jk'} \Psi_{k'} = \Psi_k, \quad k = 1, 2, \dots, d,$$

or

$$\sum_j v_{jk} \bar{v}_{jk'} = \delta_{kk'},$$

i.e. the matrix $\{v_{jk}\}$ is unitary, hence $\{u_{jk}\}$ is also unitary. \square

Using the second conjugation relation

$$S_{\rho}^* \times 1_{\rho} \cdot 1_{\rho} \times R_{\rho} = 1_{\rho},$$

one obtains the same result, i.e. the unitarity of the matrix $\{u_{jk}\}$ of the vector $u \in \mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_{\rho}$ w.r.t. arbitrary chosen orthonormal bases. However, in this case the unitarity relation reads

$$\sum_k v_{jk} \bar{v}_{j'k} = \delta_{jj'},$$

i.e. the difference is only that between the orthogonality of rows and columns.

Using the unitary matrix $\{u_{jk}\}$ one can define a transformed orthonormal basis in $\mathcal{H}_{\bar{\rho}}$ by the formula

$$\Psi'_j := \sum_{j'} u_{j'j} \Psi_{j'}.$$

Obviously, $\{\Psi'_j\}_{j=1}^d$ is an orthonormal basis in $\mathcal{H}_{\bar{\rho}}$. However, w.r.t. this basis the formula for the vector $u \in \mathcal{H}_{\bar{\rho}} \otimes \mathcal{H}_{\rho}$ reads now

$$u = \sum_j \Psi'_j \otimes \Phi_j.$$

We define the concept *conjugated basis* by this relation.

5.2.DEFINITION. Let $\{\Phi_j\}_{j=1}^d$ be an orthonormal basis of the Hilbert space \mathcal{H}_ρ . An orthonormal basis $\{\Psi_k\}_{k=1}^d$ of $\mathcal{H}_{\bar{\rho}}$ is said to be *conjugated* w.r.t. $\{\Phi_j\}_{j=1}^d$ if

$$u = \sum_{j=1}^d \Psi_j \otimes \Phi_j, \quad R_\rho = J(\bar{\rho}, \rho)u.$$

Then the foregoing considerations based on Lemma 5.1. can be formulated as follows:

5.3.LEMMA. Let $\rho \in \mathcal{T}$ and let $\bar{\rho} \in \mathcal{T}$ be conjugated w.r.t. ρ . Let $\{\Phi_{\rho j}\}_{j=1}^{d(\rho)}$ be an orthonormal basis in \mathcal{H}_ρ . Then there exists a conjugated basis $\{\Psi_{\bar{\rho} j}\}_{j=1}^{d(\rho)}$ in $\mathcal{H}_{\bar{\rho}}$, i.e.

$$R_\rho = J(\bar{\rho}, \rho) \left(\sum_{j=1}^{d(\rho)} \Psi_{\bar{\rho} j} \otimes \Phi_{\rho j} \right).$$

In general a conjugated basis is not unique. However, note that if ρ is irreducible then R_ρ hence u , where $R_\rho = J(\bar{\rho}, \rho)u$, are unique up to a phase factor λ , $|\lambda| = 1$. Therefore we obtain

5.4.LEMMA. If ρ is irreducible then the conjugated basis is unique up to a phase factor λ , $|\lambda| = 1$.

Proof. Let ρ be irreducible. Then R_ρ is unique up to a factor λ , $|\lambda| = 1$ (see Doplicher/Roberts [2, Lemma 2.2., p.165]). Then u is also unique up to this factor. Let Φ_j be an orthonormal basis of \mathcal{H}_ρ and let Ψ_k, Ψ'_k be conjugated bases. Then we have $\sum_j \Psi'_j \otimes \Phi_j = \lambda \sum_j \Psi_j \otimes \Phi_j$, $|\lambda| = 1$ hence $\sum_j (\Psi'_j - \lambda \Psi_j) \otimes \Phi_j = 0$ and it follows that $\Psi'_j = \lambda \Psi_j$. \square

Now we proceed to the second step, the definition of the $*$ -structure in \mathcal{F}_0 .

5.5.DEFINITION. Let $\rho \in \hat{\mathcal{T}}$ be irreducible and let $\Phi_{\rho j}$, $j = 1, \dots, d$ be an orthonormal basis in \mathcal{H}_ρ . The conjugated basis in $\mathcal{H}_{\bar{\rho}}$ is denoted by $\Phi_{\bar{\rho} j}$, $j = 1, \dots, d$. Then $\Phi_{\rho j}^*$ is defined by

$$\Phi_{\rho j}^* := R_\rho^* \Phi_{\bar{\rho} j}, \quad j = 1, \dots, d.$$

Note that this definition makes sense. Namely, if $\Psi_{\bar{\rho} j}$ is a second conjugated basis then $\Psi_{\bar{\rho} j} = \lambda \Phi_{\bar{\rho} j}$, $|\lambda| = 1$, and $\sum_j \Psi_{\bar{\rho} j} \otimes \Phi_{\rho j} = \lambda \sum_j \Phi_{\bar{\rho} j} \otimes \Phi_{\rho j}$ hence $J(\bar{\rho}, \rho) \sum_j \Psi_{\bar{\rho} j} \otimes \Phi_{\rho j} = \lambda J(\bar{\rho}, \rho) \sum_j \Phi_{\bar{\rho} j} \otimes \Phi_{\rho j}$ or $\hat{R}_\rho = \lambda R_\rho$ and $\hat{R}_\rho^* \Psi_{\bar{\rho} j} = \bar{\lambda} R_\rho^* \lambda \Phi_{\bar{\rho} j} = |\lambda|^2 R_\rho^* \Phi_{\bar{\rho} j} = R_\rho^* \Phi_{\bar{\rho} j}$.

Additionally we define

$$(A\Phi_{\rho j})^* := \Phi_{\rho j}^* A^*$$

and extend the definition by linearity to the whole algebra \mathcal{F}_0 . We have

5.6.LEMMA. The relations

$$(\Phi_{\rho j}^*)^* = \Phi_{\rho j}, \tag{13}$$

$$(\Phi_{\rho j} \cdot \Phi_{\rho' j'})^* = \Phi_{\rho' j'}^* \cdot \Phi_{\rho j}^*, \tag{14}$$

hold.

Proof. First we prove (13): We have

$$\begin{aligned} (\Phi_{\rho j}^*)^* &= (R_\rho^* \Phi_{\bar{\rho} j})^* = \Phi_{\bar{\rho} j}^* R_\rho = R_{\bar{\rho}}^* \Phi_{\rho j} R_\rho = \\ R_{\bar{\rho}}^* \rho(R_\rho) \Phi_{\rho j} &= R_{\bar{\rho}}^* \times 1_\rho \cdot 1_\rho \times R_\rho \cdot \Phi_{\rho j} = \Phi_{\rho j} \end{aligned}$$

because of the conjugation relation (T8) for R_ρ , where $R_{\bar{\rho}} := S_\rho$.

Second we prove (14): Here we have

$$\Phi_{\rho j} \cdot \Phi_{\rho' j'} = \sum_k W_k \psi_k, \quad \psi_k := \sum_j \lambda_j \Phi_{\tau_k j} \in \mathcal{H}_{\tau_k}, \quad W_k \in (\tau_k, \rho \rho')$$

where the $\tau_k \in \hat{\mathcal{T}}$ are irreducible. Therefore

$$(\Phi_{\rho j} \cdot \Phi_{\rho' j'})^* = \sum_k \psi_k^* W_k^* = \sum_k R_{\tau_k}^* \bar{\psi}_k W_k^* = \sum_k R_{\tau_k}^* \bar{\tau}_k (W_k^*) \bar{\psi}_k, \quad (15)$$

where $\bar{\psi}_k := \sum_j \bar{\lambda}_j \Phi_{\bar{\tau}_k j}$. On the other hand we get

$$\Phi_{\rho' j'}^* \Phi_{\rho j}^* = R_{\rho'}^* \Phi_{\bar{\rho}' j'}^* R_\rho^* \Phi_{\bar{\rho} j}^* = R_{\rho'}^* \bar{\rho}'(R_\rho^*) \Phi_{\bar{\rho}' j'}^* \Phi_{\bar{\rho} j}^*$$

and we have

$$\Phi_{\bar{\rho}' j'}^* \Phi_{\bar{\rho} j}^* = \sum_k \hat{W}_k \bar{\psi}_k$$

where \hat{W}_k is the conjugated arrow w.r.t. W_k , $\hat{W}_k \in (\bar{\tau}_k, \bar{\rho}' \bar{\rho})$, see [2,p.167] or [6,p.257]. So we obtain

$$\Phi_{\rho' j'}^* \Phi_{\rho j}^* = \sum_k R_{\rho'}^* \bar{\rho}'(R_\rho^*) \hat{W}_k \bar{\psi}_k. \quad (16)$$

Comparing (15) and (16) we see that it is sufficient to prove

$$R_{\tau}^* \bar{\tau}(W^*) = R_{\rho'}^* \bar{\rho}'(R_\rho^*) \hat{W},$$

where τ is irreducible, $W \in (\tau, \rho \rho')$, and $\hat{W} \in (\bar{\tau}, \bar{\rho}' \bar{\rho})$ is the conjugated arrow. Note that

$$R_{\rho'}^* \bar{\rho}'(R_\rho^*) = R_{\rho \rho'}^*.$$

This means that it is sufficient to prove

$$R_{\tau}^* \bar{\tau}(W^*) = R_{\sigma}^* \hat{W}, \quad (17)$$

where τ is irreducible, σ arbitrary, $W \in (\tau, \sigma)$, $\hat{W} \in (\bar{\tau}, \bar{\sigma})$ the conjugated arrow. But (17) is the characteristic relation for the conjugated arrow (see Doplicher/Roberts [2, Lemma 2.5., p.167] or [7, (12), p.257]). \square

The relations (13) and (14) ensure that the $*$ -structure given by Definition 5.5 satisfies the usual properties, i.e. \mathcal{F}_0 is a $*$ -algebra.

The orthonormal bases of the Hilbert spaces \mathcal{H}_α , $\alpha \in \mathcal{T}$, satisfy important relations.

5.7.PROPOSITION. Let $\alpha \in \mathcal{T}$ be arbitrary and let $\{\Phi_{\alpha j}\}_{j=1}^{d(\alpha)}$ be an orthonormal basis of \mathcal{H}_α . Then the relations

$$\Phi_{\alpha j}^* \Phi_{\alpha k} = \delta_{jk} 1, \quad (18)$$

$$\sum_j \Phi_{\alpha j} \Phi_{\alpha j}^* = 1 \quad (19)$$

hold.

Proof. First we prove (18) for $\rho \in \hat{\mathcal{T}}$. We have $\Phi_{\rho j}^* \Phi_{\rho k} = R_\rho^* \Phi_{\bar{\rho} j} \Phi_{\rho k}$. Note that

$$\phi(R_\rho^*)X = R_\rho^* X = (R_\rho, X)_{\mathcal{H}_{\bar{\rho}}} 1, \quad X \in \mathcal{H}_{\bar{\rho}}.$$

Thus

$$\begin{aligned} \Phi_{\rho j}^* \Phi_{\rho k} &= (R_\rho, \Phi_{\bar{\rho} j} \Phi_{\rho k})_{\mathcal{H}_{\bar{\rho}}} 1 = \\ &= \left(\sum_{l=1}^{d(\rho)} \Phi_{\bar{\rho} l} \Phi_{\rho l}, \Phi_{\bar{\rho} j} \Phi_{\rho k} \right)_{\mathcal{H}_{\bar{\rho}}} 1 = \\ &= \sum_{l=1}^{d(\rho)} (\Phi_{\bar{\rho} l}, \Phi_{\bar{\rho} j})_{\mathcal{H}_{\bar{\rho}}} (\Phi_{\rho l}, \Phi_{\rho k})_{\mathcal{H}_\rho} 1 = \sum_{l=1}^{d(\rho)} \delta_{lj} \delta_{lk} 1 = \delta_{jk} 1. \end{aligned}$$

Now, using the definition of the Hilbert spaces \mathcal{H}_α , $\alpha \in \mathcal{T}$, it is easy to extend the formula to arbitrary $\alpha \in \mathcal{T}$.

Second we prove (19), again first for $\rho \in \hat{\mathcal{T}}$. We calculate

$$\begin{aligned} \sum_{j=1}^{d(\rho)} \Phi_{\rho j} \Phi_{\rho j}^* &= \sum_{j=1}^{d(\rho)} \Phi_{\rho j} R_\rho^* \Phi_{\bar{\rho} j} = \sum_{j=1}^{d(\rho)} \rho(R_\rho^*) \Phi_{\rho j} \Phi_{\bar{\rho} j} \\ &= \rho(R_\rho^*) \sum_{j=1}^{d(\rho)} \Phi_{\rho j} \Phi_{\bar{\rho} j} = \rho(R_\rho^*) R_{\bar{\rho}} = 1_\rho \times R_\rho^* \cdot R_{\bar{\rho}} \times 1_\rho = 1_\rho, \end{aligned}$$

according to the conjugation relation. Again, using the definition of the Hilbert spaces \mathcal{H}_α , the formula can be easily extended to arbitrary $\alpha \in \mathcal{T}$. \square

Proposition 5.7. means that the Hilbert spaces $\mathcal{H}_\alpha \subset \mathcal{F}_0$, $\alpha \in \mathcal{T}$, are algebraic Hilbert spaces of the $*$ -algebra \mathcal{F}_0 . In particular, (19) means that the so-called support of \mathcal{H}_α equals 1. Note that the Hilbert spaces \mathcal{H}_ρ for irreducibles $\rho \in \hat{\mathcal{T}}$ together with \mathcal{A} are generating for \mathcal{F}_0 w.r.t. \mathcal{F}_0 as an \mathcal{A} -leftmodule.

Recall that according to Propositions 3.1. and 3.2. we have bijections

$$\text{Ob } \mathcal{T} \in \alpha \longleftrightarrow \mathcal{H}_\alpha \subset \mathcal{F}_0,$$

$$\mathcal{A} \supset (\alpha, \beta) \longleftrightarrow \phi((\alpha, \beta)) \subset \mathcal{L}(\mathcal{H}_\alpha \rightarrow \mathcal{H}_\beta).$$

Taking into account the results in sections 4 and 5 we can conclude that in fact by these bijections there is established a categorial isomorphism: First we have

$$\mathcal{H}_{\alpha\beta} = \mathcal{H}_\alpha \cdot \mathcal{H}_\beta$$

where $\mathcal{H}_\alpha \cdot \mathcal{H}_\beta$ denotes the algebraic Hilbert space $\text{spa}\{\mathcal{H}_\alpha \cdot \mathcal{H}_\beta\}$. Second, according to the original definition of the algebraic Hilbert spaces \mathcal{H}_α , we get: If

$$\gamma(\cdot) = V\alpha(\cdot)V^* + W\beta(\cdot)W^*,$$

where $V \in (\alpha, \gamma)$, $W \in (\beta, \gamma)$, $VV^* + WW^* = 1$, $V^*V = W^*W = 1$, then

$$\mathcal{H}_\gamma = V\mathcal{H}_\alpha + W\mathcal{H}_\beta.$$

Third we obtain: the category $\{\mathcal{H}_\alpha\}_\alpha$ is equipped with a conjugation, given by the conjugation $\alpha \rightarrow \bar{\alpha}$.

Therefore, at this stage we can apply the Tannaka-Krein result. This means there is a compact group \mathcal{K} such that $\hat{\mathcal{K}}$ is in one-to-one correspondence to $\{\mathcal{H}_\rho\}_\rho$ for a full collection of representers $\rho \in \hat{\mathcal{T}}$. So to each $\alpha \in \mathcal{T}$ we get a unitary representation $U_\alpha(\mathcal{K})$ of \mathcal{K} on \mathcal{H}_α which is irreducible for $\alpha \in \hat{\mathcal{T}}$. Note that \mathcal{K} acts on $\mathcal{H}_{\bar{\alpha}}$, $\alpha \in \mathcal{T}$, by the *conjugated* representation w.r.t. \mathcal{H}_α , i.e.

$$U_{\bar{\alpha}}(k) = \overline{U_\alpha(k)}, \quad k \in \mathcal{K},$$

where $\bar{}$ means complex conjugation of the representing matrices.

The special structure of \mathcal{F}_0 allows to define an action of \mathcal{K} on \mathcal{F}_0 by the definition

$$U(k)\left(\sum_{\rho,j} A_{\rho j} \Phi_{\rho j}\right) := \sum_{\rho,j} A_{\rho j} U_\rho(k) \Phi_{\rho j}.$$

Obviously we have

$$U(k)(X) = U_\alpha(k)X, \quad X \in \mathcal{H}_\alpha,$$

because from

$$X = \sum_{j=1}^r B_j \Phi_j, \quad B_j \in (\rho_j, \alpha), \quad \Phi_j \in \mathcal{H}_{\rho_j}, \quad \rho_j \text{ irreducible},$$

we obtain

$$U(k)(X) = \sum_{j=1}^r B_j U_{\rho_j} \Phi_j = U_\alpha(k) \sum_{j=1}^r B_j \Phi_j = U_\alpha(k)X.$$

Therefore it turns out that $U(k)$ is even an automorphism of \mathcal{F}_0 , because for $\Phi_1 \in \mathcal{H}_\alpha$, $\Phi_2 \in \mathcal{H}_\beta$ we obtain

$$U(k)(\Phi_1 \cdot \Phi_2) = U_{\alpha\beta}(k)(\Phi_1 \Phi_2) = (U_\alpha(k)\Phi_1)(U_\beta(k)\Phi_2) = U(k)\Phi_1 \cdot U(k)\Phi_2.$$

Moreover we have

$$(U(k)(F))^* = U(k)(F^*), \quad F \in \mathcal{F}_0.$$

Namely, let $\{\Phi_{\rho j}\}_{j=1}^{d(\rho)}$ be an orthonormal basis of \mathcal{H}_ρ and

$$U_\rho(k)\Phi_{\rho j} = \sum_{j'=1}^{d(\rho)} u_{j'j}^{(\rho)}(k)\Phi_{\rho j'}.$$

Then we get

$$\begin{aligned}
(U(k)(\Phi_{\rho j}))^* &= (U_\rho(k)\Phi_{\rho j})^* = \left(\sum_{j'=1}^{d(\rho)} u_{j'j}^{(\rho)}(k)\Phi_{\rho j'} \right)^* \\
&= \sum_{j'=1}^{d(\rho)} \overline{u_{j'j}^{(\rho)}(k)}\Phi_{\rho j'}^* = \sum_{j'=1}^{d(\rho)} \overline{u_{j'j}^{(\rho)}}R_\rho^*\Phi_{\bar{\rho}j'} = \\
&R_\rho^* \sum_{j'=1}^{d(\rho)} \overline{u_{j'j}^{(\rho)}}\Phi_{\bar{\rho}j'} = R_\rho^*(U_{\bar{\rho}}(k)\Phi_{\bar{\rho}j}) = \\
R_\rho^*U(k)(\Phi_{\bar{\rho}j}) &= U(k)(R_\rho^*\Phi_{\bar{\rho}j}) = U(k)(\Phi_{\rho j}^*).
\end{aligned}$$

This relation can be easily extended to be valid for arbitrary $F = \sum_{\rho,j} A_{\rho j}\Phi_{\rho j}$.

Furthermore, the fixed point algebra of $U(\mathcal{K})$ within \mathcal{F}_0 equals \mathcal{A} , because if $\mathcal{F}_0 \ni F = \sum_{\rho,j} A_{\rho j}\Phi_{\rho j}$ and $U(k)(F) = F$ for all $k \in \mathcal{K}$ then we obtain

$$\sum_{\rho,j} A_{\rho j}U_\rho(k)\Phi_{\rho j} = \sum_{\rho,j,j'} A_{\rho j}u_{j'j}^{(\rho)}(k)\Phi_{\rho j'} = \sum_{\rho,j} A_{\rho j}\Phi_{\rho j},$$

hence

$$\sum_{j'=1}^{d(\rho)} u_{jj'}^{(\rho)}(k)A_{\rho j'} = A_{\rho j}$$

follows for all $k \in \mathcal{K}$, $\rho \in \hat{T}$, $j = 1, 2, \dots, d(\rho)$. If $\rho \neq \iota$, by integration over \mathcal{K} we get $A_{\rho j} = 0$.

In the next section we characterize the group \mathcal{K} as the group of all automorphisms of \mathcal{F}_0 leaving a characteristic \mathcal{A} -scalar product invariant.

6 The \mathcal{A} -scalar product and the symmetry group

Let $F, G \in \mathcal{F}_0$. Then $F = \sum_{\rho,j} A_{\rho j}\Phi_{\rho j}$, $G = \sum_{\rho,j} B_{\rho j}\Phi_{\rho j}$ where $\{\Phi_{\rho j}\}_{j=1}^{d(\rho)}$ denotes, as before, an orthonormal basis in \mathcal{H}_ρ ; $A_{\rho j}, B_{\rho j} \in \mathcal{A}$. Recall $\Phi_\iota = 1$. The coefficient A_ι of F is called the \mathcal{A} -component of F , resp. B_ι is the \mathcal{A} -component of G . We define an \mathcal{A} -scalar product in the \star -algebra \mathcal{F}_0 by

$$\langle F, G \rangle_{\mathcal{A}} := \sum_{\rho,j} \frac{1}{d(\rho)} A_{\rho j} B_{\rho j}^*.$$

Obviously, $\langle F, G \rangle_{\mathcal{A}} \in \mathcal{A}$ and $\langle F, G \rangle_{\mathcal{A}}$ is independent of the special choice of the orthonormal bases in the Hilbert spaces \mathcal{H}_ρ . The \mathcal{A} -scalar product $\langle F, G \rangle_{\mathcal{A}}$ satisfies the properties

$$\begin{aligned}
\langle AF, BG \rangle_{\mathcal{A}} &= A \langle F, G \rangle_{\mathcal{A}} B^*, \\
\langle F, F \rangle_{\mathcal{A}} &\geq 0, \\
\langle F, F \rangle_{\mathcal{A}} = 0 &\text{ iff } F = 0.
\end{aligned}$$

We have

6.1.LEMMA. Let $F, G \in \mathcal{F}_0$. Then $\langle F, G \rangle_{\mathcal{A}}$ coincides with the \mathcal{A} -component of FG^* .

Proof. It is sufficient to prove this for $F := \Phi_{\alpha j}$, $G := \Phi_{\beta k}$, $\alpha, \beta \in \hat{T}$ irreducible. Recall that there is a decomposition

$$\Phi_{\alpha j} \cdot \Phi_{\beta k} = \sum_{\tau, l} K_{\alpha j \beta k}^{\tau l} \Phi_{\tau l}$$

of the product $\Phi_{\alpha j} \cdot \Phi_{\beta k}$ where $K_{\alpha j \beta k}^{\tau l} \in (\tau, \alpha\beta)$. We have

$$\Phi_{\alpha j} \Phi_{\beta k}^* = \Phi_{\alpha j} R_{\beta}^* \Phi_{\beta k} = \alpha(R_{\beta}^*) \Phi_{\alpha j} \Phi_{\beta k}^*.$$

Therefore, the \mathcal{A} -component of $\Phi_{\alpha j} \Phi_{\beta k}^*$ coincides with $\alpha(R_{\beta}^*) K_{\alpha j \beta k}^{\iota}$. On the other hand, by definition we have

$$\langle \Phi_{\alpha j}, \Phi_{\beta k} \rangle_{\mathcal{A}} = \frac{1}{d(\alpha)} \delta_{\alpha\beta} \delta_{jk} 1.$$

So it is sufficient to prove

$$\alpha(R_{\beta}^*) K_{\alpha j \beta k}^{\iota} = \frac{1}{d(\alpha)} \delta_{\alpha\beta} \delta_{jk} 1. \quad (20)$$

We check easily that

$$\alpha(R_{\beta}^*) K_{\alpha j \beta k}^{\iota} \in (\beta, \alpha)$$

because $K_{\alpha j \beta k}^{\iota} \in (\beta, \alpha\bar{\beta}\beta)$ and $\alpha(R_{\beta}^*) \in (\alpha\bar{\beta}\beta, \alpha)$. Therefore we have: if $\alpha \neq \beta$ then $\alpha(R_{\beta}^*) K_{\alpha j \beta k}^{\iota} = 0$ and (20) is true. Now let $\alpha = \beta$. Then we have to prove

$$\alpha(R_{\alpha}^*) K_{\alpha j \bar{\alpha} k} = \frac{1}{d(\alpha)} \delta_{jk} 1.$$

Since the left hand side is a member of (α, α) , we have

$$\alpha(R_{\alpha}^*) K_{\alpha j \bar{\alpha} k} = \lambda_{jk} 1, \quad \lambda_{jk} \in \mathbf{C}.$$

Note that

$$K_{\alpha j \bar{\alpha} k}^{\iota} = J(\alpha, \bar{\alpha}) u_{jk}, \quad u_{jk} \in \mathcal{H}_{\alpha} \otimes \mathcal{H}_{\bar{\alpha}}, \quad u_{jk} = \sum_{p, q} c_{jk}^{pq} \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} q}.$$

Let $X \in \mathcal{H}_{\alpha}$. We calculate

$$\begin{aligned} \alpha(R_{\alpha}^*) K_{\alpha j \bar{\alpha} k}^{\iota} X &= 1_{\alpha} \times R_{\alpha}^* \cdot J(\alpha \bar{\alpha}, \alpha) (K_{\alpha j \bar{\alpha} k}^{\iota} \otimes X) \\ &= 1_{\alpha} \times R_{\alpha}^* \cdot J(\alpha \bar{\alpha}, \alpha) (J(\alpha, \bar{\alpha}) u_{jk} \otimes 1_{\mathcal{H}_{\alpha}} X) \\ &= 1_{\alpha} \times R_{\alpha}^* \cdot J(\alpha \bar{\alpha}, \alpha) (J(\alpha, \bar{\alpha}) \otimes 1_{\mathcal{H}_{\alpha}}) (u_{jk} \otimes X) = \\ &= 1_{\alpha} \times R_{\alpha}^* \cdot J(\alpha, \bar{\alpha} \alpha) (1_{\mathcal{H}_{\alpha}} \otimes J(\bar{\alpha}, \alpha)) \left(\sum_{p, q} c_{jk}^{pq} \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} q} \otimes X \right) = \end{aligned}$$

$$\begin{aligned}
& J(\alpha, \iota)(1_{\mathcal{H}_\alpha} \otimes R_\alpha^*)(1_{\mathcal{H}_\alpha} \otimes J(\bar{\alpha}, \alpha)) \left(\sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} q} \otimes X \right) = \\
& (1_{\mathcal{H}_\alpha} \otimes R_\alpha^* J(\bar{\alpha}, \alpha)) \left(\sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} q} \otimes X \right) = \sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} (R_\alpha, J(\bar{\alpha}, \alpha) \Phi_{\bar{\alpha} q} \otimes X) \mathcal{H}_{\bar{\alpha}\alpha} = \\
& \sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} \left(J(\bar{\alpha}, \alpha) \sum_{l=1}^{d(\alpha)} \Phi_{\bar{\alpha} l} \otimes \Phi_{\alpha l}, J(\bar{\alpha}, \alpha) \Phi_{\bar{\alpha} q} \otimes X \right)_{\mathcal{H}_{\bar{\alpha}\alpha}} = \\
& \sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} \sum_{l=1}^{d(\alpha)} (\Phi_{\bar{\alpha} l} \otimes \Phi_{\alpha l}, \Phi_{\bar{\alpha} q} \otimes X) \mathcal{H}_{\bar{\alpha}\alpha} \mathcal{H}_\alpha = \sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} \sum_{l=1}^{d(\alpha)} (\Phi_{\bar{\alpha} l}, \Phi_{\bar{\alpha} q}) \mathcal{H}_{\bar{\alpha}} (\Phi_{\alpha l}, X) \mathcal{H}_\alpha \\
& = \sum_{p,q} c_{jk}^{pq} \Phi_{\alpha p} (\Phi_{\alpha q}, X) \mathcal{H}_\alpha = \lambda_{jk} X.
\end{aligned}$$

For $X := \Phi_{\alpha s}$, we obtain the equation

$$\sum_p c_{jk}^{ps} \Phi_{\alpha p} = \lambda_{jk} \Phi_{\alpha s}$$

hence

$$c_{jk}^{ps} = \delta^{ps} \lambda_{jk}$$

follows. That is, we obtain

$$u_{jk} = \lambda_{jk} \sum_p \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} p}$$

and

$$K_{\alpha_j \bar{\alpha} k}^\iota = \lambda_{jk} J(\alpha, \bar{\alpha}) \left(\sum_p \Phi_{\alpha p} \otimes \Phi_{\bar{\alpha} p} \right) = \lambda_{jk} R_{\bar{\alpha}}.$$

Therefore we can write

$$\Phi_{\alpha j} \cdot \Phi_{\bar{\alpha} k} = \lambda_{jk} R_{\bar{\alpha}} + \sum_{\gamma \neq \iota} \sum_l K_{\alpha_j \bar{\alpha} k}^{\gamma l} \Phi_{\gamma l}. \quad (21)$$

In other words, we have: The \mathcal{A} -component of $\Phi_{\alpha j} \cdot \Phi_{\bar{\alpha} k}$ equals $\lambda_{jk} R_{\bar{\alpha}}$. From (21) we obtain

$$R_{\bar{\alpha}}^* \cdot \Phi_{\alpha j} \cdot \Phi_{\bar{\alpha} k} = \lambda_{jk} R_{\bar{\alpha}}^* R_{\bar{\alpha}} + \sum_{\gamma \neq \iota} \sum_l R_{\bar{\alpha}}^* K_{\alpha_j \bar{\alpha} k}^{\gamma l} \Phi_{\gamma l}.$$

Recall $R_{\bar{\alpha}} \in \mathcal{H}_{\alpha \bar{\alpha}}$ and $\Phi_{\alpha j} \Phi_{\bar{\alpha} k} \in \mathcal{H}_{\alpha \bar{\alpha}}$. Therefore we get

$$R_{\bar{\alpha}}^* \Phi_{\alpha j} \Phi_{\bar{\alpha} k} = (R_{\bar{\alpha}}, \Phi_{\alpha j} \Phi_{\bar{\alpha} k}) \mathcal{H}_{\alpha \bar{\alpha}} 1 = \lambda_{jk} d(\alpha) 1 + \sum_{\gamma \neq \iota} \sum_l R_{\bar{\alpha}}^* K_{\alpha_j \bar{\alpha} k}^{\gamma l} \Phi_{\gamma l}.$$

Now we calculate

$$(R_{\bar{\alpha}}, \Phi_{\alpha j} \Phi_{\bar{\alpha} k}) \mathcal{H}_{\alpha \bar{\alpha}} = \left(J(\alpha, \bar{\alpha}) \sum_{l=1}^{d(\alpha)} \Phi_{\alpha l} \otimes \Phi_{\bar{\alpha} l}, J(\alpha, \bar{\alpha}) (\Phi_{\alpha j} \otimes \Phi_{\bar{\alpha} k}) \right)_{\mathcal{H}_{\alpha \bar{\alpha}}}$$

$$= \sum_{l=1}^{d(\alpha)} (\Phi_{\alpha l}, \Phi_{\alpha j})_{\mathcal{H}_\alpha} (\Phi_{\bar{\alpha} l}, \Phi_{\bar{\alpha} k})_{\mathcal{H}_{\bar{\alpha}}} = \sum_{l=1}^{d(\alpha)} \delta_{lj} \delta_{lk} = \delta_{jk}$$

and we obtain

$$\delta_{jk} 1 = \lambda_{jk} d(\alpha) 1 + \sum_{\gamma \neq l} \sum_l R_{\bar{\alpha}}^* K_{\alpha j \bar{\alpha} k}^{\gamma l} \Phi_{\gamma l}.$$

Since the system $\{\Phi_{\rho j}\}_{\rho j}$ forms an \mathcal{A} -leftmodule basis we finally get

$$\lambda_{jk} = \frac{1}{d(\alpha)} \delta_{jk}$$

and (20) is proved for $\alpha = \beta$, too. \square

6.2.COROLLARY. *The relation*

$$\langle X, YF^* \rangle_{\mathcal{A}} = \langle XF, Y \rangle_{\mathcal{A}}; \quad X, Y, F \in \mathcal{F}_0,$$

holds.

Proof. We already know that $\langle F, G \rangle_{\mathcal{A}}$ equals the \mathcal{A} -component of FG^* . Now

$$X(YF^*)^* = XFY^*$$

holds. Therefore their \mathcal{A} -components also coincide. \square

Next we introduce a C^* -norm on \mathcal{F}_0 . The corresponding arguments are standard and follow the lines of Doplicher/Roberts, e.g. in [1, p.88 f.]; see also [7, p.217 f.] and [7, p.286 f.].

The first step is to introduce a norm in \mathcal{F}_0 by the definition

$$|F|_{\mathcal{A}} := \|\langle F, F \rangle_{\mathcal{A}}\|^{1/2}. \quad (22)$$

Obviously $|A|_{\mathcal{A}} = \|A\|$ for $A \in \mathcal{A}$ because $\|AA^*\|^{1/2} = \|A\|$. Note that $|\cdot|_{\mathcal{A}}$ fails to be a C^* -norm. Therefore in a second step we define the following norm on \mathcal{F}_0 :

$$\|F\|_* := |\pi(F)|, \quad F \in \mathcal{F}_0, \quad (23)$$

where $(\pi(F))X := XF^*$, $X \in \mathcal{F}_0$ and $|\pi(F)|$ means the operator norm w.r.t. the norm $|\cdot|_{\mathcal{A}}$ on \mathcal{F}_0 , i.e.,

$$|\pi(F)| = \sup_{|X|_{\mathcal{A}} \leq 1} |XF^*|_{\mathcal{A}}.$$

$\|\cdot\|_*$ turns out to be a C^* -norm on \mathcal{F}_0 . Note that the C^* -property of $\|\cdot\|_*$ follows from Corollary 6.2. which says that

$$\langle X, \pi(F)Y \rangle_{\mathcal{A}} = \langle \pi(F^*)X, Y \rangle_{\mathcal{A}}, \quad X, Y, F \in \mathcal{F}_0.$$

Note further that $\|\cdot\|_*$ on \mathcal{F}_0 is an extension of the C^* -norm $\|\cdot\|$ on $\mathcal{A} \subset \mathcal{F}_0$, i.e.

$$\|A\|_* = \|A\|, \quad A \in \mathcal{A}.$$

Moreover, we have

$$|F|_{\mathcal{A}} \leq \|F\|_*, \quad F \in \mathcal{F}_0.$$

Finally we form the C^* -algebra

$$\mathcal{F} := \text{clo}_{\|\cdot\|_*}(\mathcal{F}_0).$$

Now we turn to a special automorphism group \mathcal{G} of \mathcal{F} . First we consider again the $*$ -subalgebra $\mathcal{F}_0 \subseteq \mathcal{F}$ which is dense in \mathcal{F} and consider automorphisms g of \mathcal{F}_0 requiring invariance of the \mathcal{A} -scalar product $\langle \cdot, \cdot \rangle_{\mathcal{A}}$:

6.3.DEFINITION. The automorphism $g \in \text{aut } \mathcal{F}_0$ is said to be a *symmetry* if

$$\langle gF_1, gF_2 \rangle_{\mathcal{A}} = \langle F_1, F_2 \rangle_{\mathcal{A}}, \quad F_1, F_2 \in \mathcal{F}_0. \quad (24)$$

The collection of all automorphisms $g \in \text{aut } \mathcal{F}_0$ of this type forms a group $\mathcal{G} \neq \emptyset$, because the identical automorphism $\text{id}: \text{id } F := F$ belongs to \mathcal{G} . The group \mathcal{G} is called the *symmetry* group.

6.4.LEMMA. *Each automorphism $g \in \mathcal{G}$ can be uniquely extended to an automorphism of \mathcal{F} so that \mathcal{G} can be considered as a subgroup of $\text{aut } \mathcal{F}$, $\mathcal{G} \subseteq \text{aut } \mathcal{F}$.*

Proof. From (22) and (24) it follows immediately that $|gF|_{\mathcal{A}} = |F|_{\mathcal{A}}$, $F \in \mathcal{F}_0$. Furthermore, we have $\|gF\|_* = \|F\|_*$ for all $F \in \mathcal{F}_0$ because from (23) we obtain

$$\begin{aligned} \|gF\|_* &= |\pi(gF)| = \sup_{|X|_{\mathcal{A}} \leq 1} |X(gF)^*|_{\mathcal{A}} = \sup_{|X|_{\mathcal{A}} \leq 1} |g\{(g^{-1}X)F^*\}|_{\mathcal{A}} \\ &= \sup_{|X|_{\mathcal{A}} \leq 1} |(g^{-1}X)F^*|_{\mathcal{A}} = \sup_{|Y|_{\mathcal{A}} \leq 1} |YF^*|_{\mathcal{A}} = |\pi(F)| = \|F\|_*. \end{aligned}$$

Therefore, by continuous (isometric) extension each g can be extended to the whole C^* -algebra \mathcal{F} , yielding an automorphism of \mathcal{F} . \square

We emphasize that the symmetry group \mathcal{G} coincides with the *stability group* $\text{stab } \mathcal{A} \subset \text{aut } \mathcal{F}$ of \mathcal{F} . Recall that $\text{stab } \mathcal{A} := \{g \in \text{aut } \mathcal{F} : gA = A \text{ for all } A \in \mathcal{A}\}$.

First we have

6.5.LEMMA. *Each Hilbert space \mathcal{H}_α , $\alpha \in \mathcal{T}$, is invariant w.r.t. $g \in \text{stab } \mathcal{A}$.*

Proof. Recall that $X \in \mathcal{H}_\alpha$ iff $\alpha(A)X = XA$ for all $A \in \mathcal{A}$. Therefore we have

$$\alpha(A) \cdot gX = gX \cdot A, \quad X \in \mathcal{H}_\alpha, A \in \mathcal{A}, g \in \text{stab } \mathcal{A},$$

i.e. $gX \in \mathcal{H}_\alpha$. \square

Now we prove

6.6.LEMMA. *The relation*

$$\mathcal{G} = \text{stab } \mathcal{A}$$

holds.

Proof. First we prove $\mathcal{G} \subseteq \text{stab } \mathcal{A}$. Let $A \in \mathcal{A}$. Then we have to prove that $gA = A$ for all $g \in \mathcal{G}$. Note that a priori $gA \in \mathcal{F}_0$. We put $\mathcal{F}_0 \ni F = \sum_{\alpha, j} (F)_{\alpha j} \Phi_{\alpha j}$ with fixed orthonormal bases $\{\Phi_{\alpha j}\}_{j=1}^{d(\alpha)}$. For example $(F)_i$ is the \mathcal{A} -component of F . Recall that $\langle F, 1 \rangle_{\mathcal{A}} = (F)_i$. Then

$$\langle gA, \Phi_{\alpha j} \rangle_{\mathcal{A}} = \langle A, g^{-1} \Phi_{\alpha j} \rangle_{\mathcal{A}} = A(g^{-1} \Phi_{\alpha j})_i^* = 0$$

for $\alpha \neq \iota$, $j = 1, 2, \dots, d(\alpha)$ because $g\mathcal{H}_\alpha \subseteq \mathcal{H}_\alpha$, i.e. $gA \in \mathcal{A}$. This implies

$$gA = (gA)_\iota = \langle gA, 1 \rangle_{\mathcal{A}} = \langle gA, g1 \rangle_{\mathcal{A}} = \langle A, 1 \rangle_{\mathcal{A}} = A$$

hence $gA = A$ follows.

Second we prove $\text{stab } \mathcal{A} \subseteq \mathcal{G}$. Let $g \in \text{stab } \mathcal{A}$. From Lemma 6.5. we get

$$g\Phi_{\alpha j} = \sum_{j'} u_{j'j}^{(\alpha)}(g)\Phi_{\alpha j'},$$

where $\{u_{j'j}^{(\alpha)}(g)\}_{j',j=1}^{d(\alpha)}$ is a unitary $d(\alpha) \times d(\alpha)$ -matrix because

$$(g\Phi_{\alpha j})^* g\Phi_{\alpha j'} = g(\Phi_{\alpha j}^* \Phi_{\alpha j'}) = g(\delta_{jj'}1) = \delta_{jj'}1.$$

Now we have

$$\langle gF_1, gF_2 \rangle_{\mathcal{A}} = \sum_{\rho, j, \sigma, k} A_{\rho j} \langle g\Phi_{\rho j}, g\Phi_{\sigma k} \rangle_{\mathcal{A}} B_{\sigma k}^*.$$

In the case $\rho \neq \sigma$ we have $\langle g\Phi_{\rho j}, g\Phi_{\sigma k} \rangle_{\mathcal{A}} = 0 = \langle \Phi_{\rho j}, \Phi_{\sigma k} \rangle_{\mathcal{A}}$. In the other case $\rho = \sigma$ we get

$$\begin{aligned} \langle g\Phi_{\rho j}, g\Phi_{\rho k} \rangle_{\mathcal{A}} &= \left\langle \sum_{j'} u_{j'j} \Phi_{\rho j'}, \sum_{k'} u_{k'k} \Phi_{\rho k'} \right\rangle_{\mathcal{A}} = \sum_{j'k'} u_{j'j} \bar{u}_{k'k} \langle \Phi_{\rho j'}, \Phi_{\rho k'} \rangle_{\mathcal{A}} \\ &= \sum_{j'k'} u_{j'j} \bar{u}_{k'k} \frac{1}{d(\rho)} \delta_{j'k'} = \frac{1}{d(\rho)} \sum_{k'} u_{k'j} \bar{u}_{k'k} = \frac{1}{d(\rho)} \delta_{jk} = \langle \Phi_{\rho j}, \Phi_{\rho k} \rangle_{\mathcal{A}}, \end{aligned}$$

hence $\langle gF_1, gF_2 \rangle_{\mathcal{A}} = \langle F_1, F_2 \rangle_{\mathcal{A}}$ follows. \square

Further it turns out that \mathcal{G} is compact.

6.7.PROPOSITION. *The symmetry group \mathcal{G} is compact w.r.t. the pointwise norm topology.*

Proof. The Hilbert spaces \mathcal{H}_α , $\dim \mathcal{H}_\alpha = d(\alpha)$, are invariant w.r.t. \mathcal{G} and they generate \mathcal{F} , together with \mathcal{A} . Therefore, according to Doplicher/Roberts [8, Lemma 3.2., p.292] \mathcal{G} is compact w.r.t. the pointwise norm topology. \square

6.8.PROPOSITION. *Let $\mathcal{F}_0 \ni F = \sum_{\rho, j} F_{\rho j} \Phi_{\rho j}$, where the system of orthonormal bases $\{\Phi_{\rho j}\}_{j=1}^{d(\rho)}$ is fixed. Then*

$$\|F_{\rho j}\| \leq d(\rho) \|F\|_*, \quad j = 1, 2, \dots, d(\rho), \rho \in \hat{T}, \quad (25)$$

i.e. $F_{\rho j}$ depends continuously on F .

Proof. We have

$$\frac{1}{d(\rho)} \|F_{\rho j}\|^2 = \frac{1}{d(\rho)} \|F_{\rho j} F_{\rho j}^*\| \leq \left\| \sum_{\alpha, j} d(\alpha)^{-1} F_{\alpha j} F_{\alpha j}^* \right\| = \|\langle F, F \rangle_{\mathcal{A}}\|^2 = \|F|_{\mathcal{A}}\|^2 \leq \|F\|_*^2.$$

This proves (25). \square

6.9.COROLLARY. *The \mathcal{A} -scalar product $\langle F_1, F_2 \rangle_{\mathcal{A}}$ on \mathcal{F}_0 is continuous w.r.t. the C^* -norm $\|\cdot\|_*$.*

Proof. According to Lemma 6.1., we have $\langle F_1, F_2 \rangle_{\mathcal{A}} = (F_1 F_2^*)_{\iota}$. So it is sufficient to prove that F_{ι} depends continuously on F . But this assertion is a special case of Proposition 6.8. \square

Corollary 6.9. implies that the \mathcal{A} -scalar product has a continuous extension to \mathcal{F} w.r.t. $\|\cdot\|_{\star}$.

Obviously, on \mathcal{F}_0 we can define projections Π_{ρ} , $\rho \in \hat{T}$, by the definition

$$\mathcal{F}_0 \ni F = \sum_{\sigma, j} A_{\sigma j} \Phi_{\sigma j} \rightarrow \Pi_{\rho} F := \sum_{j=1}^{d(\rho)} A_{\rho j} \Phi_{\rho j} \in \text{spa}\{\mathcal{A}\mathcal{H}_{\rho}\}.$$

Note that this definition does not depend on the special orthonormal basis in \mathcal{H}_{ρ} .

$\Pi_{\rho} F$ depends continuously on F w.r.t. $\|\cdot\|_{\star}$, in particular we get

$$\|\Pi_{\rho} F\| \leq d(\rho)^2 \|F\|_{\star}, \quad F \in \mathcal{F}_0.$$

Therefore, Π_{ρ} can be continuously extended on \mathcal{F} as a projection and we obtain the relations

$$\begin{aligned} 1 &\leq \|\Pi_{\rho}\|_{\star} \leq d(\rho)^2, \\ \Pi_{\rho}(A_1 F A_2) &= A_1 \Pi_{\rho}(F) A_2, \quad f \in \mathcal{F}, A_1, A_2 \in \mathcal{A}, \\ \Pi_{\rho} \Pi_{\sigma} &= \Pi_{\sigma} \Pi_{\rho} = \delta_{\rho\sigma} \Pi_{\rho}, \end{aligned}$$

for a fixed system of representers $\rho, \sigma \in \hat{T}$.

6.10.COROLLARY. *The relative commutant $\mathcal{A}' \cap \mathcal{F}$ is trivial, $\mathcal{A}' \cap \mathcal{F} = \mathbb{C}1$.*

Proof. Let $F \in \mathcal{F}$ and $AF = FA$ for all $A \in \mathcal{A}$. Then

$$A \cdot \Pi_{\rho} F = \Pi_{\rho}(AF) = \Pi_{\rho}(FA) = \Pi_{\rho} F \cdot A, \quad \rho \in \hat{T}.$$

This means $\Pi_{\rho} F \in \mathcal{A}' \cap \mathcal{F}_0$, i.e. $\Pi_{\rho} F = \lambda_{\rho} 1$, hence for $\rho \neq \iota$ one has $\lambda_{\rho} = 0$ and this implies $F = \lambda_{\iota} 1$. \square

6.11.LEMMA. *The relation*

$$\Pi_{\rho} \mathcal{F} = \text{spa}\{\mathcal{A}\mathcal{H}_{\rho}\}, \quad \rho \in \hat{T},$$

is valid.

Proof. Obviously, $\Pi_{\rho} \mathcal{F}_0 = \text{spa}\mathcal{A}\mathcal{H}_{\rho}$. Since $\Pi_{\rho} \mathcal{F} = \text{clo}_{\|\cdot\|_{\star}} \Pi_{\rho} \mathcal{F}_0$ it is sufficient to prove that $\text{spa}\mathcal{A}\mathcal{H}_{\rho}$ is closed w.r.t. $\|\cdot\|_{\star}$. Let $F_n = \sum_{j=1}^{d(\rho)} A_{\rho j}^{(n)} \Phi_{\rho j} \in \text{spa}\mathcal{A}\mathcal{H}_{\rho}$ and $\|F_n - F_m\|_{\star} \rightarrow 0$, $n, m \rightarrow \infty$. The limit is denoted by F_{∞} . Then from $|F_n - F_m|_{\mathcal{A}}^2 \leq \|F_n - F_m\|_{\star}^2$ we get

$$\begin{aligned} \|\langle F_n - F_m, F_n - F_m \rangle_{\mathcal{A}}\| &= \left\| \sum_{j=1}^{d(\rho)} (A_{\rho j}^{(n)} - A_{\rho j}^{(m)})(A_{\rho j}^{(n)} - A_{\rho j}^{(m)})^{\star} \right\| \\ &\geq \|(A_{\rho j}^{(n)} - A_{\rho j}^{(m)})(A_{\rho j}^{(n)} - A_{\rho j}^{(m)})^{\star}\| = \|A_{\rho j}^{(n)} - A_{\rho j}^{(m)}\|^2 \rightarrow 0 \end{aligned}$$

hence $F_{\infty} = \sum_{j=1}^{d(\rho)} A_{\rho j}^{\infty} \Phi_{\rho j}$ and $F_{\infty} \in \text{spa}\mathcal{A}\mathcal{H}_{\rho}$. \square

6.12.LEMMA. *The projection Π_ρ commutes with the action of \mathcal{G} , i.e.*

$$\Pi_\rho(gF) = g\Pi_\rho F, \quad F \in \mathcal{F}, g \in \mathcal{G}. \quad (26)$$

Proof. It is sufficient to prove (26) for all elements $F \in \mathcal{F}_0$ because (26) then follows using extension by continuity. But since \mathcal{H}_ρ is invariant w.r.t. \mathcal{G} and $gA = A$ for all $A \in \mathcal{A}$, we have $g \sum_{\rho,j} A_{\rho j} \Phi_{\rho j} = \sum_{\rho,j} A_{\rho j} g \Phi_{\rho j}$, hence $g(\Pi_\rho F) = \sum_{j=1}^{d(\rho)} A_{\rho j} g \Phi_{\rho j} = \Pi_\rho(gF)$ follows. \square

6.13.PROPOSITION. *The fixed point algebra of \mathcal{G} coincides with \mathcal{A} .*

Proof. First we note that $U(\mathcal{K})$ is a subgroup of \mathcal{G} , $U(\mathcal{K}) \subseteq \mathcal{G}$. This is obvious because

$$\langle U(k)F_1, U(k)F_2 \rangle_{\mathcal{A}} = \langle F_1, F_2 \rangle_{\mathcal{A}}, \quad k \in \mathcal{K}, F_1, F_2 \in \mathcal{F}.$$

Since \mathcal{K} acts irreducibly on \mathcal{H}_ρ , $\rho \in \hat{\mathcal{T}}$, and \mathcal{H}_ρ is invariant w.r.t. \mathcal{G} it follows that \mathcal{H}_ρ is irreducible w.r.t. \mathcal{G} , too. Now, according to Lemma 6.12, from $gF = F$ for all $g \in \mathcal{G}$ we obtain $g\Pi_\rho F = \Pi_\rho F$, i.e. the Π_ρ -components of F are also invariant. So we have to prove $\Pi_\rho F = 0$ for all $\rho \neq \iota$. Now $\Pi_\rho F = \sum_{j=1}^{d(\rho)} A_{\rho j} \Phi_{\rho j}$ and $g(\Pi_\rho F) = \sum_{j,j'=1}^{d(\rho)} A_{\rho j} u_{j'j}^{(\rho)}(g) \Phi_{\rho j'}$, hence $A_{\rho j} = \sum_{j'=1}^{d(\rho)} u_{j'j}^{(\rho)}(g) A_{\rho j'}$ follows. Integrating this over \mathcal{G} we obtain $A_{\rho j} = 0$. \square

6.14.COROLLARY. *The groups $U(\mathcal{K})$ and \mathcal{G} coincide.*

Proof. First one has to prove that $\text{spec } \mathcal{G} = \hat{\mathcal{G}}$, i.e. all irreducible representations of \mathcal{G} occur. This is obtained by considering the scalar functions $H_1^* g(H_2)$, $H_1, H_2 \in \mathcal{H}_\alpha$, $\alpha \in \mathcal{T}$. Using the Stone-Weierstra approximation theorem one gets that these functions span a dense set in $C(\mathcal{G})$ hence in $L^2(\mathcal{G})$. Then one can conclude that $U(\mathcal{K}) = \mathcal{G}$ (cf. Doplicher/Roberts [8, 3.3.Theorem, p.293]). \square

We note that also a direct calculation shows that the representations induced by \mathcal{G} on \mathcal{H}_α and $\mathcal{H}_{\bar{\alpha}}$, $\alpha \in \mathcal{T}$, are mutually conjugated. Namely we have $R_\alpha = \sum_{j=1}^{d(\alpha)} \Phi_{\bar{\alpha}j} \Phi_{\alpha j}$ hence $R_\alpha = \sum_{j=1}^{d(\alpha)} (g \Phi_{\bar{\alpha}j})(g \Phi_{\alpha j})$ and one obtains immediately

$$u_{j'j}^{(\bar{\alpha})}(g) = \overline{u_{j'j}^{(\alpha)}(g)}, \quad g \in \mathcal{G}.$$

Finally we note that the objects $\alpha \in \mathcal{T}$ turn out to be the restrictions to \mathcal{A} of the so-called *canonical endomorphisms* w.r.t. \mathcal{H}_α ,

$$\alpha(X) = \sum_{j=1}^{d(\alpha)} \Phi_{\alpha j} X \Phi_{\alpha j}^*, \quad X \in \mathcal{A}, \alpha \in \mathcal{T}.$$

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