



# A planetary-scale land–sea breeze circulation in East Asia and the western North Pacific

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The diurnal wind variation over the East Asian continent is commonly considered to be a combination of a land–sea breeze near the coast and a mountain–valley breeze along the slopes of the Tibetan Plateau. The local land–sea breeze along the coastline typically spans <100 km into the ocean. However, a detailed examination of the global reanalysis data suggests that this local land–sea breeze circulation apparently couples with the global-scale diurnal atmospheric pressure tide to produce a planetary-scale land–sea breeze with a spatial scale of ~1000 km over the western North Pacific. Computations of the momentum budget and equivalent potential temperatures indicate that the atmospheric diurnal tidal wave contributes the most to this circulation feature. A diagnosis of the water vapour budget further suggests that the convergence of water vapour flux, which is related to the convergence of low-level wind induced by the seasonal change of diurnal tidal wave, leads to different times of occurrence of maximum diurnal rainfall over East Asia between summer and winter. Copyright © 2010 Royal Meteorological Society

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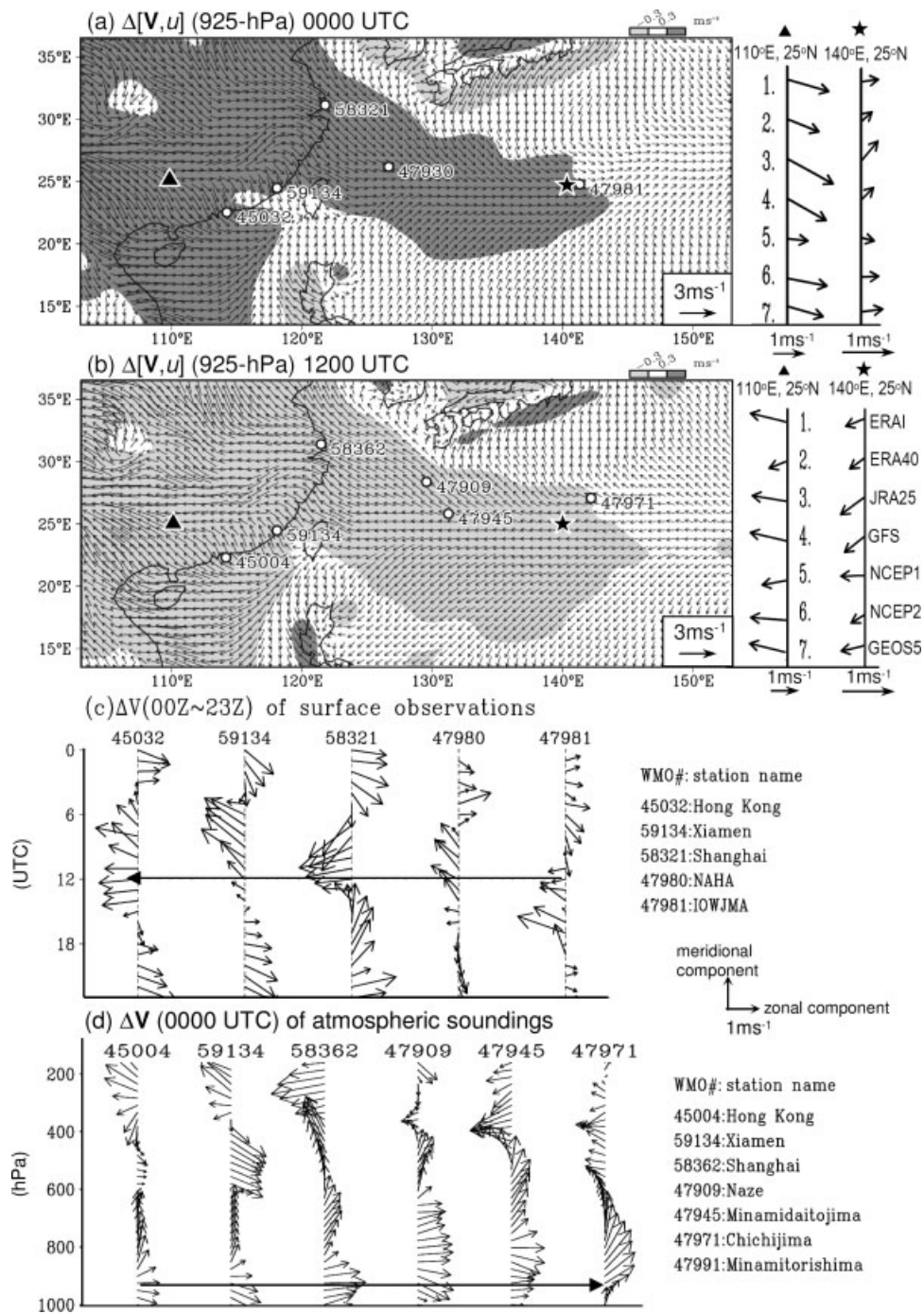
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## 1. Introduction

The periodic incoming solar heating in the atmosphere generates global tidal waves which can be revealed from variables such as pressure, wind, temperature, radiation fluxes, and even precipitation. The tidal waves evolve with periods of 24 and 12 h, commonly referred to as the diurnal and semi-diurnal oscillations respectively (Wallace and Hartranft, 1969; Haurwitz and Cowley, 1973; Hamilton, 1980; Hsu and Hoskins, 1989; Deser and Smith, 1998; Dai and Wang, 1999). On a regional scale, the pressure variations are geographically dependent and subject to the modulation of mountain–valley and land–sea differential heating, as well as latent heating in moist convection (Hamilton, 1981; Krishnamurti and Kishtawal, 2000; Ciesielski and Johnson, 2008). The East Asia–Western North Pacific

(hereafter EAWNP) region with complex mountain–valley and land–sea distributions exhibits large diurnal variations in surface wind (Krishnamurti and Kishtawal, 2000; Yu *et al.*, 2008) and precipitation (Ramage, 1952; Zhao *et al.*, 2005; Dai *et al.*, 2007; Yu *et al.*, 2007a, 2007b; Li *et al.*, 2008; Zhou *et al.*, 2008). The diurnal surface wind variation is important for determining the timing of maximum diurnal precipitation (Yeh and Chen, 1998; Dai and Deser, 1999; Wu *et al.*, 2009). Compared to the extensive literature on the diurnal variation of precipitation, however, the diurnal variation of surface winds over the EAWNP region has attracted less attention.

During the summer months (June, July and August, or JJA), a typical diurnal variation of lower-tropospheric winds over the EAWNP region between morning (Figure 1(a)) and evening (Figure 1(b)), in terms of the 925 hPa



**Figure 1.** (a) The JJA mean of 925 hPa wind anomalies at 0000 UTC from GEOS5 (left panel) and two single-point wind anomalies ((110°E, 25°N) marked by triangle and (140°E, 25°N) marked by star) at 0000 UTC from seven reanalyses (right panel). (b) is similar to (a), but for 1200 UTC. The scale of stippled zonal wind speed is given atop (a)–(b). (c) The hourly surface wind anomalies of five selected surface stations in summer. (d) The multiple-level wind anomalies observed by six selected atmospheric soundings at 0000 UTC in summer. The WMO ID and name of selected stations (marked in (a)) and atmospheric soundings (marked in (b)) are provided in the right panel of (c) and (d). The scale of wind vector for (a)–(b) is given at bottom right and for (c)–(d) is shown in the right middle panel between them.

wind anomalies from the daily mean, reveals distinct features including the local land–sea breeze (hereafter LSB) circulation over the South China Sea and coastal East Asia, mountain–valley winds between the Tibetan Plateau and eastern China, and pronounced wind anomalies over the ocean between 120°E and 145°E. These features altogether form a planetary-scale LSB-like circulation which has not been examined in detail hitherto, although a concept of large-scale diurnal variation in the global surface winds was proposed in Dai and Deser (1999) based on the observations that surface wind divergence tends to peak

in the morning over the major continents but in the late afternoon over the oceans. Previous studies suggested that a typical LSB circulation (e.g. LSB along the coast of the Philippines: Figure 1(a)–(b)), as a result of the difference in sensible heat flux over land and sea, only varies on a horizontal scale of <100 km (Schmidt, 1947; Neumann and Mahrer, 1971; Rotunno, 1983). In other words, to produce a diurnal variation of circulations of such a large scale as that revealed in Figure 1, an extra forcing mechanism in addition to the land–sea differential heating must exist. We therefore hypothesize that the pressure gradient force

from the large-scale global atmospheric pressure tidal wave is a major driving forcing of this planetary-scale LSB-like circulation.

The objective of the present study is to examine this hypothesis, which is important because a proper conceptual model for such a large-scale LSB-like circulation has significant implications for the improvement of weather and climate simulations. As the diurnal variation of low-level winds controls that of precipitation, this study also investigates how such a large-scale LSB-like circulation affects the diurnal precipitation change. Analyses are performed on both station data and global reanalyses, which are introduced in section 2. Results of this study are presented in sections 3 and 4, followed by a summary and discussion in section 5.

## 2. Data and Methodology

The analysis utilizes wind observations from surface stations (WMO ID and locations are given in Figure 1(a)), atmospheric soundings (located in Figure 1(b)), and meteorological data from seven global reanalyses including NCEP1, NCEP2, ERA40, ERAI, JRA25, GEOS5 and GFS (abbreviations and references described in Table I). The analysis covers the time period from 1989 to 2009 during summer (JJA) and winter (December, January and February, or DJF). The precipitation analysis uses the TRMM dataset (Tropical Rainfall Measuring Mission: Simpson *et al.*, 1996) which has been shown to resemble rain-gauge observations (Zhou *et al.*, 2008) with a good representation of the diurnal rainfall variability (Hong *et al.*, 2005). In addition, precipitation records of surface stations are extracted from the National Climatic Data Center<sup>†</sup>. Other information, including the period of record, temporal resolution and spatial resolution for all the datasets, is also listed in Table I.

In this study, anomalies of a given variable from the GEOS5 reanalyses, surface stations, and atmospheric soundings at a specific synoptic time step (e.g. 0000 UTC; 1200 UTC) are obtained by subtracting the daily means from their 3-hourly maps, 24-hourly observations, and twice-daily observations, respectively. The same procedure is carried out for the other reanalyses except that the map interval is 6 hours. The diurnal and semi-diurnal harmonic components of anomaly variables are obtained from Fourier analysis. Hereafter, the anomalies, diurnal (i.e. first harmonic) and semi-diurnal (i.e. second harmonic) cycles of a given variable  $X$  are denoted by  $\Delta X$ ,  $S1(X)$  and  $S2(X)$ , respectively.

## 3. The Planetary-Scale LSB-like Circulation over the EAWNP Region

During the summer months, the planetary-scale LSB-like circulation in the morning (0000 UTC; 0700–0900 local standard time LST) (Figure 1(a)) exhibits westerly wind anomalies over the eastern Tibetan Plateau around 110°E and offshore winds up to 145°E. The reverse occurs in the evening (1200 UTC; 1900–2100 LST) (Figure 1(b)), with easterly wind anomalies spanning from 145°E westward to

100°E. Such a morning land-breeze-like and evening sea-breeze-like circulation covers a range of a few thousand kilometres between the coastline of East Asia (hereafter EA) and ~145°E. This planetary-scale circulation is discernible not only in GEOS5 (left panel of Figure 1(a)–(b)) but in all other reanalyses (right panel of Figure 1(a)–(b)), surface stations (Figure 1(c)), and atmospheric soundings (Figure 1(d)). Because the results from all reanalyses are mutually consistent, to avoid redundancy we only show the GEOS5 reanalyses in the following discussion. The selection of GEOS5 is for two reasons. First, because the spatial resolution of GEOS5 (~66.7 km in longitude  $\times$  50 km in latitude: Table I) is smaller than the horizontal scale of a typical LSB circulation (<100 km), GEOS5 is suitable for capturing features of local LSB which need to be resolved when we examine its possible coupling with the global atmospheric pressure tide later. Second, the 3-hourly GEOS5 data give a higher sampling rate for both the diurnal and semi-diurnal harmonics to be analysed than other 6-hourly reanalyses (Dai and Deser, 1999; Chen, 2005).

As revealed from the soundings over the EAWNP region (Figure 1(d)), the phase reversal of wind direction in the vertical implies ascending motion over the WNP ocean and downward motion over the EA continent which, subsequently, introduces a zonal mass flux circulation between the continent and the ocean. Because the differential heating between mountains and valleys only causes the diurnal wind change over land (Krishnamurti and Kishtawal, 2000) and the land–sea differential heating only induces LSB along the coast (Neumann and Mahrer, 1971), other forcing mechanisms for causing such a large-scale diurnal circulation must exist.

Dai and Deser (1999) proposed a mechanism responsible for a large-scale diurnal surface wind divergence change between day and night, which is similar to that for the large-scale monsoon circulation change between summer and winter. However, it takes weeks or even months to develop a giant thermal contrast between land and ocean that forms a monsoon system (e.g. Ramage, 1971). It is unlikely that the diurnally-driven land–sea thermal contrast can induce an LSB-like circulation on such a large scale within 12 hours. Deser and Smith (1998) examined the tropical oceanic zonal wind variations with the diurnal time-scale and suggested that the pressure gradient force (hereafter PGF) produced by the semi-diurnal component of atmospheric tides is likely the driving mechanism of such wind anomalies. It is therefore possible that the aforesaid extra forcing to induce a planetary-scale LSB-like circulation over the EAWNP region, as that shown in Figure 1, is linked to the diurnally varying global atmospheric pressure tide.

As seen from the distribution of mean sea-level pressure (hereafter SLP) anomalies, denoted as  $\Delta SLP$  (Figure 2(a)), the global pressure tidal wave is a planetary-scale phenomenon dominated by a spatial wave-number-2 structure. It was noted that the pressure tidal wave experiences regular oscillations at the diurnal and semi-diurnal periods (Chapman and Lindzen, 1970). More than 96% of the total variability of global  $\Delta SLP$  (Figure 2(a)) is explained by the combination of its diurnal  $S1(SLP)$  (Figure 2(b)) and semi-diurnal  $S2(SLP)$  (Figure 2(c)) harmonics. In general,  $\Delta SLP$  over the tropical open ocean is mostly contributed by its related  $S2(SLP)$  and less by  $S1(SLP)$  (Dai and Wang, 1999). In the middle latitudes, the magnitude of  $\Delta SLP$  is apparently larger over the EA

<sup>†</sup>Information is available at <http://www.ncdc.noaa.gov/oa/climate/climatedata.html#hourly>.



Table I. List of observations and reanalysis datasets.

Dataset	Data period	Temporal resolution	Spatial resolution	Source information
WMO surface stations	1950 to present	Hourly		<ul style="list-style-type: none"> <li>• <a href="http://dss.ucar.edu/datasets/ds564.0/">http://dss.ucar.edu/datasets/ds564.0/</a></li> </ul>
WMO sounding stations	1973 to present	Twice daily		<ul style="list-style-type: none"> <li>• <a href="http://weather.uwyo.edu/upperair/sounding.html">http://weather.uwyo.edu/upperair/sounding.html</a></li> <li>• Simpson <i>et al.</i> (1996)</li> </ul>
Tropical Rainfall Measuring Mission (hereafter, TRMM)	1998 to present	3 hours	0.25° (long) × 0.25° (lat)	<ul style="list-style-type: none"> <li>• Kalnay <i>et al.</i> (1996)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds090.0/">http://dss.ucar.edu/datasets/ds090.0/</a></li> </ul>
National Centers for Environmental Prediction (NCEP)-NCAR Reanalysis (hereafter, NCEP1)	1948 to present	6 hours	2.5° (long) × 2.5° (lat)	<ul style="list-style-type: none"> <li>• Kalnay <i>et al.</i> (1996)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds090.0/">http://dss.ucar.edu/datasets/ds090.0/</a></li> </ul>
NCEP-Department of Energy Atmospheric Model Intercomparison Project (AMIP-II) Reanalysis (hereafter, NCEP2)	1979 to present	6 hours	2.5° (long) × 2.5° (lat)	<ul style="list-style-type: none"> <li>• Kanamitsu <i>et al.</i> (2002)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds091.0/">http://dss.ucar.edu/datasets/ds091.0/</a></li> </ul>
The European Centre for Medium-Range Weather Forecasts (ECMWF) Re-Analysis (hereafter, ERA40)	Sep 1957 to Aug 2002	6 hours	2.5° (long) × 2.5° (lat)	<ul style="list-style-type: none"> <li>• Uppala <i>et al.</i> (2005)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds120.0/">http://dss.ucar.edu/datasets/ds120.0/</a></li> </ul>
The ECMWF Interim Re-Analysis (hereafter, ERAI)	1989 to present	6 hours	0.702° (long) × 0.703° (lat)	<ul style="list-style-type: none"> <li>• Simmons <i>et al.</i> (2007)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds627.0/">http://dss.ucar.edu/datasets/ds627.0/</a></li> </ul>
Japanese 25-year Reanalysis Project (hereafter, JRA25)	1979 to present	6 hours	2.5° (long) × 2.5° (lat)	<ul style="list-style-type: none"> <li>• Onogi <i>et al.</i> (2007)</li> <li>• <a href="http://dss.ucar.edu/datasets/ds625.1/">http://dss.ucar.edu/datasets/ds625.1/</a></li> </ul>
Goddard Earth Observing System Data Assimilation System Version 5 (hereafter, GEOS5)	1979 to Aug 2006	3 hours	0.667° (long) × 0.5° (lat)	<ul style="list-style-type: none"> <li>• <a href="http://gmao.gsfc.nasa.gov/pubs/docs/GEOS5_104606-Vol27.pdf">http://gmao.gsfc.nasa.gov/pubs/docs/GEOS5_104606-Vol27.pdf</a></li> </ul>
NCEP Global Forecast System (hereafter, GFS)	June 2004 to present	6 hours	0.5° (long) × 0.5° (lat)	<ul style="list-style-type: none"> <li>• Kanamitsu <i>et al.</i> (1991)</li> </ul>

continent than the WNP region. Such distribution of  $\Delta SLP$  can result in a positive longitudinal pressure gradient over the EAWNP region which has the same sign as the positive longitudinal pressure gradient associated with S1(SLP) (Figure 2(b)) but opposite sign to the negative longitudinal pressure gradient caused by S2(SLP) (Figure 2(c)). In other words, over the EAWNP region the land–sea difference of  $\Delta SLP$  is more likely caused by S1(SLP) rather than S2(SLP). More information about the temporal–spatial evolutions of  $\Delta SLP$  and its S1 and S2 harmonic modes were documented by Dai and Wang (1999) and Hoinka (2007).

The role of pressure tidal wave in affecting a large-scale LSB-like circulation over the EAWNP regions can be examined through the following momentum budget equation:

$$\frac{\partial V}{\partial t} = \underbrace{-f \bullet \mathbf{k} \times \mathbf{V}}_{(CF)} - \underbrace{\frac{1}{\rho} \nabla p}_{(PGF)} + \underbrace{\frac{\partial}{\partial z} (-V'w')}_{(VDIF)} + \underbrace{\text{residual}}_{(RES)}, \quad (1)$$

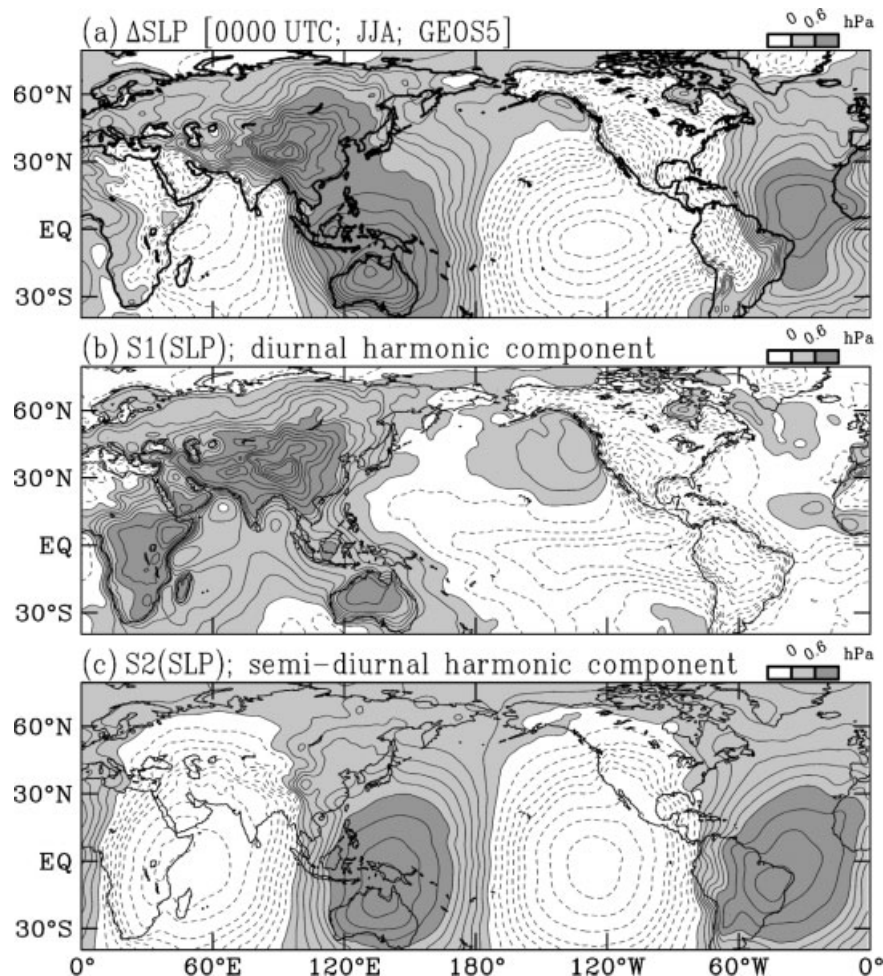
where  $\mathbf{V}$ ,  $f$ ,  $\rho$ ,  $p$  and  $w$  are the atmospheric horizontal wind vector, Coriolis parameter, air density, pressure and the vertical velocity.  $V'w'$  is the vertical turbulent flux of

momentum, as defined in Holton (1992). Equation (1) consists of the Coriolis force term (hereafter, CF), the PGF term, the vertical diffusion term (VDIF), and the residual term (RES). Using these terms, wind variations can be reconstructed by integrating Eq. (1):

$$u = \int_{t-3}^t (fv - \frac{1}{\rho} \frac{\partial p}{\partial x} - \frac{\partial u'w'}{\partial z} + \text{residual}) dt, \quad (2)$$

$$v = \int_{t-3}^t (-fu - \frac{1}{\rho} \frac{\partial p}{\partial y} - \frac{\partial v'w'}{\partial z} + \text{residual}) dt. \quad (3)$$

In response to the temporal resolution of GEOS5, a 3 h time interval is used here for the time integration for the reconstructed winds. According to Eq. (2), the east–west land–sea distribution of S1(SLP) in Figure 2(b), which forms substantial longitudinal PGF between a high-pressure centre over the EA continent and a low-pressure centre over the WNP ocean (i.e.  $S1(-\partial p/\partial x) > 0$ ), likely produces anomalous flows that resemble a giant land-breeze-like circulation (i.e.  $S1(u) > 0$ ). In contrast, the distribution of S2(SLP) in Figure 2(c) likely generates easterly wind anomalies over the EAWNP region, which is opposite to the offshore wind shown in Figure 1(a).



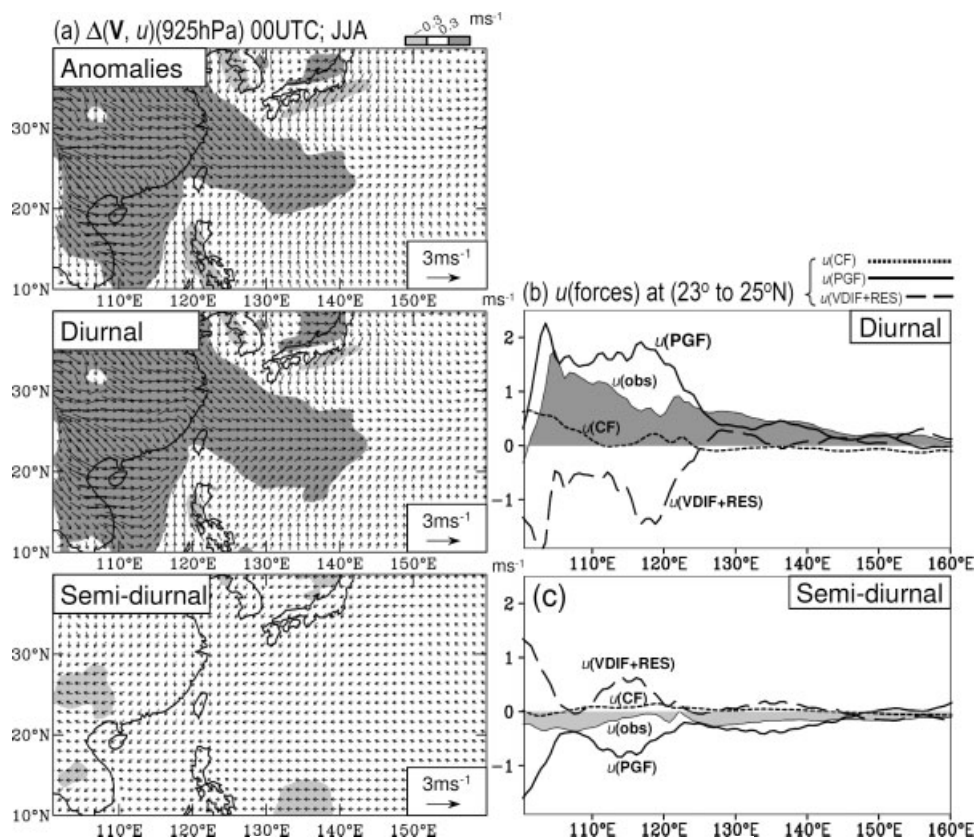
**Figure 2.** (a) The mean sea-level pressure anomalies ( $\Delta$ SLP) for 0000 UTC of 1989–2009 summers (JJA) from GEOS5. (b) and (c) are the diurnal harmonic and semi-diurnal harmonic component of (a), respectively. The contour interval in all three panels is 0.1 hPa.

To examine which kinematic forcing term contributes most to the formation of large-scale LSB-like circulation, the 925 hPa wind anomalies induced by different forcing terms based on Eqs (2)–(3) are reconstructed. Because  $23^{\circ}$ – $25^{\circ}$ N is a latitudinal zone that reveals the furthest eastward extent of the east–west LSB-like circulation (Figure 3(a)), the reconstructed  $u$ -wind (denoted as  $u(\text{forcing terms})$ ) averaged over  $23^{\circ}$ – $25^{\circ}$ N is plotted in Figure 3(b) for a detailed discussion. Moreover, because  $u(\text{RES})$  is two orders smaller than the other terms (not shown), we combine  $u(\text{RES})$  together with  $u(\text{VDIF})$  in Figure 3(b) for evaluation.

As revealed in Figure 3(a), the observed planetary-scale LSB-like circulation over the EAWNP region is mainly contributed by its diurnal component which explains more than 80% of the total variability of the zonal wind speed at 0000 UTC, based on the spatial variance of  $S1[u(925 \text{ hPa})]$  divided by that of  $\Delta[u(925 \text{ hPa})]$  over the entire domain of Figure 3. Compared to the  $u(\text{forces})$  at  $23^{\circ}$ – $25^{\circ}$ N (Figure 3(b)),  $u(\text{PGF})$  has the same order as  $u(\text{VDIF}+\text{RES})$  and is one order larger than  $u(\text{CF})$  in both S1 and S2. This result suggests that CF is not important in forming the east–west LSB-like wind at 0000 UTC or 1200 UTC. However, Yu *et al.* (2008) found CF to be important for the meridional wind variations (not shown) that are related to the formation of the East Asia low-level jet over the EA continent. A further comparison between  $S1[u(\text{PGF})]$  and  $S1[u(\text{VDIF}+\text{RES})]$  shows that the former has an opposite direction but larger magnitude than the latter, indicating

that  $S1(\text{PGF})$  is the most important forcing source for  $S1[u(925 \text{ hPa})]$ .  $S2[u(\text{PGF})]$ , however, is almost cancelled out by  $S2[u(\text{VDIF}+\text{RES})]$ , resulting in a small variation of  $S2[u(925 \text{ hPa})]$  which contributes less than  $S1[u(925 \text{ hPa})]$  to the observed  $\Delta[u(925 \text{ hPa})]$ . These results suggest that the kinematic forcing of the planetary-scale east–west LSB-like circulation over the EAWNP region is mainly caused by  $S1(\text{SLP})$ . While for areas where the meridional wind components are more evident than zonal wind components (e.g. the area ( $115^{\circ}$ – $120^{\circ}$ E,  $10^{\circ}$ – $20^{\circ}$ N): Figure 3(a)), it is noted from the horizontal distribution of reconstructed ( $u, v$ ) at 925 hPa (not shown) that CF and (VDIF+RES) are more important kinematic forcing terms as compared with PGF.

In addition to these dynamic forcing mechanisms, thermal forcing is also likely to have a contribution to the planetary-scale LSB-like circulation. Such a thermal forcing can be inferred from past studies regarding the formation of  $S1(\text{SLP})$ . It was noted that the global  $S1(\text{SLP})$ , which consists of migrating and non-migrating components, can be excited by orographic features and longitudinal variations in water vapour and ozone (Forbes and Garrett, 1979). Lindzen (1967) suggested that the major forcing of migrating tides is due to solar radiation absorption by water vapour in the troposphere and stratospheric ozone. Groves and Wilson (1982) further pointed out that water vapour heating provides a much larger contribution than ozone heating



**Figure 3.** (a) The JJA mean of 925 hPa wind anomalies and its diurnal and semi-diurnal harmonic components at 0000 UTC from GEOS5. The scale of stippled zonal wind speed is given atop (a). Based on the momentum budget equation, zonal winds induced by the Coriolis force (i.e.  $u(\text{CF})$ ), pressure gradient force (i.e.  $u(\text{PGF})$ ) and vertical diffusion with residual terms (i.e.  $u(\text{VDIF+RES})$ ) are reconstructed and a 3 h time interval is used for the time integration. (b) The diurnal component of reconstructed and observed  $u$ -winds (i.e.  $u(\text{obs})$ ) averaged between 23°N and 25°N. (c) is similar to (b), but for the semi-diurnal component.

to migrating tides. In contrast to migrating tides, the non-migrating tides can be excited by latent heat released by meteorological events in the troposphere or longitudinal variations of sensible heat due to land–sea differences and topography (Hamilton, 1981; Kato *et al.*, 1982; Forbes and Groves, 1987; Tsuda and Kato, 1989). Lieberman and Leovy (1995) further indicated that sensible heat contributed the most to the non-migrating surface pressure tide. Because both the mountain–valley and land–sea differential heating are thermal forcing mechanisms that can excite the non-migrating component of  $S1(\text{SLP})$ , it is likely that the thermal forcing exciting the migrating component of  $S1(\text{SLP})$  found by past studies (e.g. Lindzen, 1967; Groves and Wilson, 1982) is also generating the large-scale LSB-like circulation. Numerical experiments and analytical investigation will be needed to support this argument further.

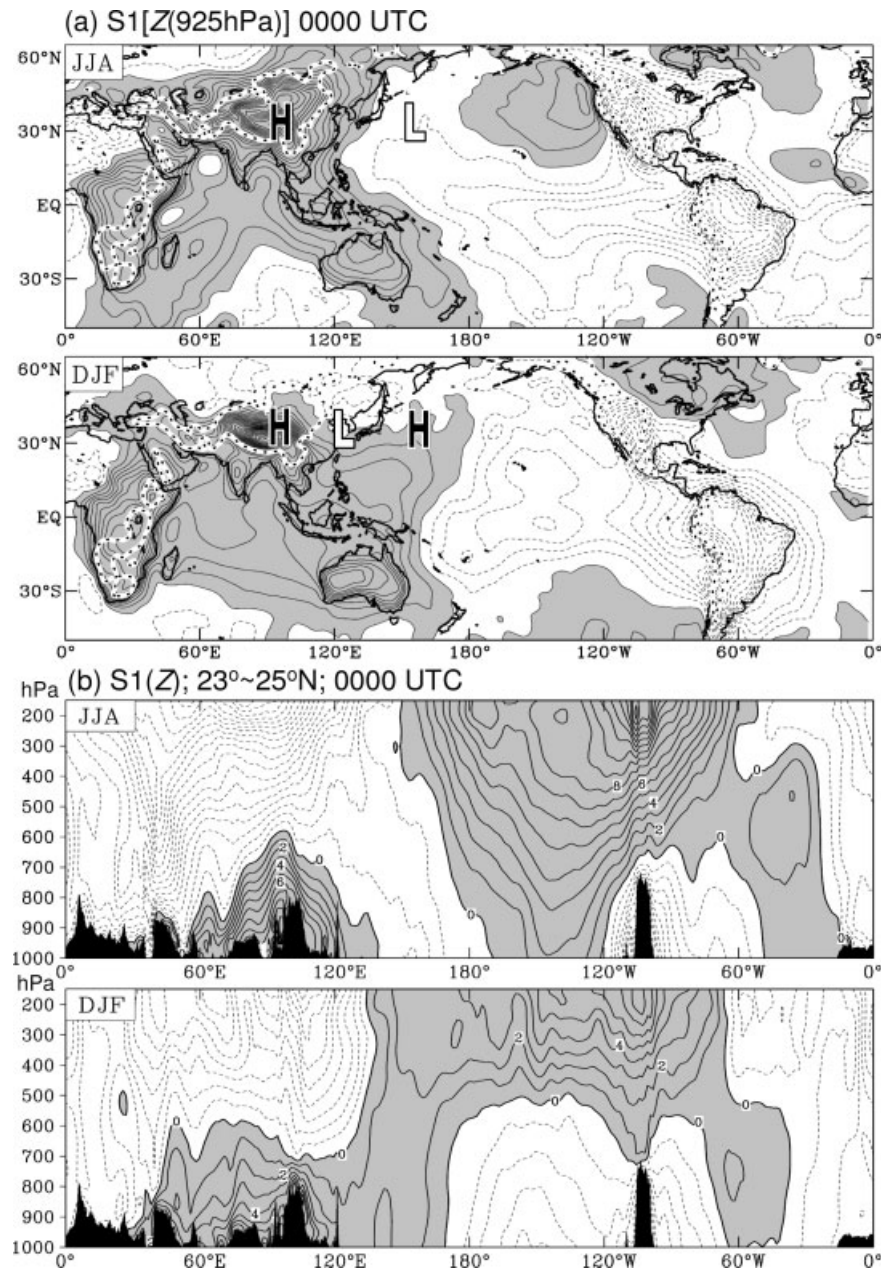
#### 4. Seasonal Variation of the LSB-like Circulation and Its Influence on Diurnal Rainfall

Because the thermally-driven global migrating tide follows the movement of the sun, seasonal variations in this migrating tidal wave should occur. Such a change implies a seasonal variation in  $S1(\text{SLP})$  and subsequently a seasonal variation in the planetary-scale LSB-like circulation. It is noted that the basic structure of  $S1(\text{SLP})$  (Figure 2(b)) is close to that of the diurnal harmonic of geopotential height (denoted as  $S1(Z)$ ) at 925 hPa (Figure 4(a); JJA). To explore the seasonal variation of large-scale diurnal circulation, we examine both the differences of  $S1[Z(925 \text{ hPa})]$  (Figure 4(a))

and its related vertical cross-section of  $S1(Z)$  at 23°–25°N (Figure 4(b)) between summer and winter. A large dimension of high–low land–sea dipole structure appears in  $S1[Z(925 \text{ hPa})]$  over the EAWNP region during northern summer but this dipole disappears in winter (Figure 4(a)). The PGF related to the giant high–low land–sea dipole of  $S1[Z(925 \text{ hPa})]$  in summer can induce a planetary-scale land-breeze-like circulation over the EAWNP region, as suggested by Figure 3. However, due to the small dimension of high–low land–sea dipole of  $S1[Z(925 \text{ hPa})]$  over the EAWNP region in winter, only a very weak land-breeze-like circulation is observed (not shown).

Vertical cross-sections of  $S1(Z)$  averaged between 23° and 25°N during summer and winter are compared (Figure 4(b)) to explain further the seasonal difference of  $S1[Z(925 \text{ hPa})]$ . The basic structure of  $S1(Z)$  is a predominantly wave-number-1 pattern at the middle and upper troposphere, consistent with the major characteristics of global migrating tide as described in Chapman and Lindzen (1970). From summer to winter, this global migrating  $S1(Z)$ , which can propagate downward to the surface over the WNP ocean, is shifted longitudinally  $\sim 60^\circ$  westward (e.g. Chen *et al.*, 2001). Such a seasonal shift results in a sign change (from negative to positive) of  $S1(Z)$  over the WNP ocean from summer to winter. In the winter (summer) morning over the WNP ocean, the positive (negative) value of  $S1(Z)$  indicating a high (low) pressure system appears to limit (promote) the eastward extension of the land-breeze-like circulation from the coastline of EA to the WNP ocean. In contrast to the WNP ocean which is dominated by the



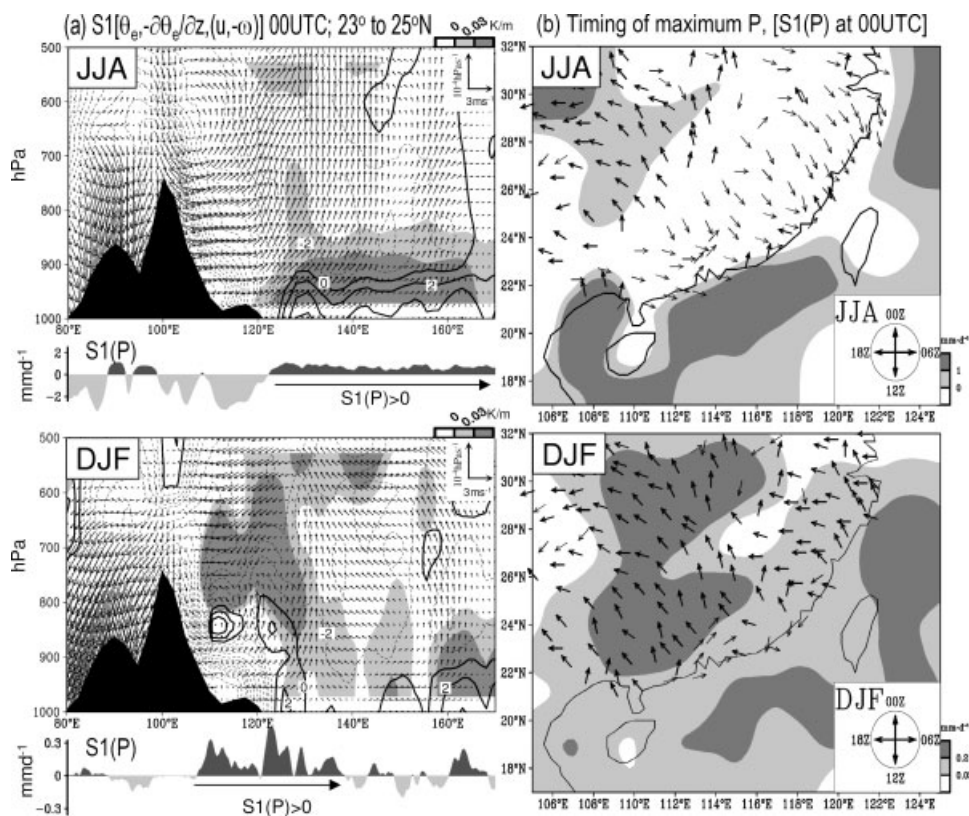


**Figure 4.** (a) The diurnal harmonic component of 925 hPa geopotential height  $S1[Z(925\text{ hPa})]$  at 0000 UTC for the 1989–2009 summer (JJA) and winter (DJF) periods. (b) The vertical cross-section of  $S1(Z)$  averaged between 23°N and 25°N. The contour intervals for  $S1(Z)$  in (a) and (b) are 1 metre. Positive areas of  $S1(Z)$  are stippled.

migrating tidal mode, the lower-tropospheric  $S1(Z)$  over the EA continent is controlled by the non-migrating tidal mode induced by the regional thermal forcings (Hsu and Hoskins, 1989; Chen *et al.*, 2001). The seasonally varying upper-troposphere downward propagating migrating tidal component combined with the lower-troposphere non-migrating tidal component forms a seasonally varying high–low land–sea dipole structure of  $S1[Z(925\text{ hPa})]$  over the EAWNP region (Figure 4(a)) and, in turn, induces the change of the horizontal extent of the giant LSB-like circulation.

Next, we investigate the impact of the planetary-scale LSB-like circulation on the diurnal rainfall variation in EA by examining the variation of equivalent potential temperature ( $\theta_e$ ) because it consists of temperature and moisture changes that are important for exciting the global pressure tide (e.g. Forbes and Garrett, 1979). As seen from  $S1(\theta_e)$  in Figure 5(a)

(contoured), the lower tropospheric air over the ocean in the morning (0000 UTC) is noticeably warmer than that over the land. This land–sea difference of  $S1(\theta_e)$  in the lower atmosphere agrees with the land–sea difference of  $S1[Z(925\text{ hPa})]$  shown in Figure 4 because the warmer air lowers the atmospheric pressure. The warmer air depicted by positive values of  $S1(\theta_e)$  covers a larger area of the WNP ocean in summer than in winter. Such a difference leads to the seasonal variation of  $S1(\text{SLP})$  which causes a smaller LSB-like circulation in winter than in summer. Despite the seasonal difference, the unstable environmental situation indicated by positive values of  $S1(-\partial\theta_e/\partial z)$  (shading in Figure 5(a)) over the ocean agrees with the fact that the precipitation over the sea tends to peak during the hours from late night to early morning (Kraus, 1963; Yang and Smith, 2006; Dai *et al.*, 2007). This feature of early morning maximum precipitation over the WNP ocean can be seen



**Figure 5.** (a) The vertical cross-section of diurnal harmonic of equivalent potential temperature ( $\theta_e$ ; contoured),  $(-\partial\theta_e/\partial z)$ ; shaded), and  $((u, -\omega)$ ; vector) at  $23^\circ$ – $25^\circ\text{N}$ , 0000 UTC for the 1989–2009 summer and winter periods. The contour interval of  $\theta_e$  is  $5 \times 10^{-2}$  K. The diurnal harmonic of precipitation (i.e.  $S1(P)$ ) at  $23^\circ$ – $25^\circ\text{N}$ , 0000 UTC from TRMM 34B2 data production is shown in the bottom panel of (a). (b) The phase clock for the timing of maximum  $P$  within a day observed from surface stations is denoted by vectors and the positive areas of  $S1(P)$  at 0000 UTC are stippled with a scale given at right.

from the  $S1$  component of precipitation,  $S1(P)$ , at 0000 UTC as shown in Figure 5(a) (bottom panel) and Figure 5(b) (shading).

A seasonal difference of the diurnal cycle of precipitation in southern China, which was first noted in Li *et al.* (2008), can be revealed from both TRMM precipitation and surface observations (Figure 5(b)). This seasonal difference of  $S1(P)$  appears to be coherent with  $S1(-\partial\theta_e/\partial z)$ , the diurnal mode of moist static instability. As seen in Figure 5(a), the unstable areas with positive  $S1(-\partial\theta_e/\partial z)$ , which is distributed over the WNP ocean in summer, is shifted westward to the EA continent in winter and, in turn, causes the seasonal change of the LSB-like circulation. This westward shifting of  $S1(-\partial\theta_e/\partial z)$  in the lower atmosphere causes the morning air over the EA continent to become more unstable and moist in the winter than in the summer. Consistent with the distribution of positive  $S1(-\partial\theta_e/\partial z)$ ,  $S1(P)$  over the EA continent reaches its maximum (minimum) in the winter (summer) morning. Such a seasonal timing difference of diurnal precipitation is not observed over inner China between  $100^\circ$  and  $110^\circ\text{E}$  where the effect of the shifting of the LSB-like anomalies circulation is small. These features of the diurnal precipitation timing as revealed in Figure 5(b) agree well with those observed previously (Zhao *et al.*, 2005; Yu *et al.*, 2007a, 2007b; Li *et al.*, 2008; Zhou *et al.*, 2008). Our finding on the difference of the planetary-scale LSB-like circulation between summer and winter therefore provides a dynamical explanation of the timing difference of diurnal precipitation between summer and winter in this region.

A more direct evidence for clarifying how an LSB-like circulation modulates the diurnal rainfall variation can be

revealed by diagnosing the following water vapour budget equation:

$$P = E + (-\nabla \cdot Q) + \left(-\frac{\partial W}{\partial t}\right), \quad (4)$$

where  $W$ ,  $\nabla \cdot Q$ ,  $P$  and  $E$  are the total precipitable water, the convergence or divergence of vertically integrated water vapour flux, the precipitation and the evaporation. Chen (2005) suggested that rainfall over EA is maintained by the water vapour supply through the convergence of water vapour flux (i.e.  $P \sim (-\nabla \cdot Q)$ ). An examination of Eq. (4) at 1200 UTC over southeast Asia (SEA; boxed areas in Figure 6(a)), where there is larger diurnal rainfall variability, indicates that  $S1(P)_{\text{SEA}}$  is mainly contributed by  $S1(-\nabla \cdot Q)_{\text{SEA}}$  (Figure 6(b)). On the other hand, it is found that the correlation coefficient between the temporal variations of  $S1(P)_{\text{SEA}}$  and  $S1(-\nabla \cdot Q)_{\text{SEA}}$  is very high ( $\sim 0.94$  in summer and  $\sim 0.92$  in winter), suggesting that the evolution of  $S1(P)_{\text{SEA}}$  follows that of  $S1(-\nabla \cdot Q)_{\text{SEA}}$ . Because the variation of  $S1(-\nabla \cdot Q)$  is led by that of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  (e.g. Sui *et al.*, 1998), a sea-breeze-like circulation in Figure 1(b) likely results in the convergence of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$ , as well as the convergence of  $S1(-\nabla \cdot Q)$ , over the EA continent for supporting the positive values of  $S1(P)$  at 1200 UTC (Figure 6(c)).

Based on such a close relationship between  $S1(P)$ ,  $S1(-\nabla \cdot Q)$  and  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$ , further estimation for the contribution of seasonally varying  $S1(\text{SLP})$  on affecting  $S1(P)$  over south China is made through estimating the contribution of seasonally varying  $S1(\text{PGF})$  to the



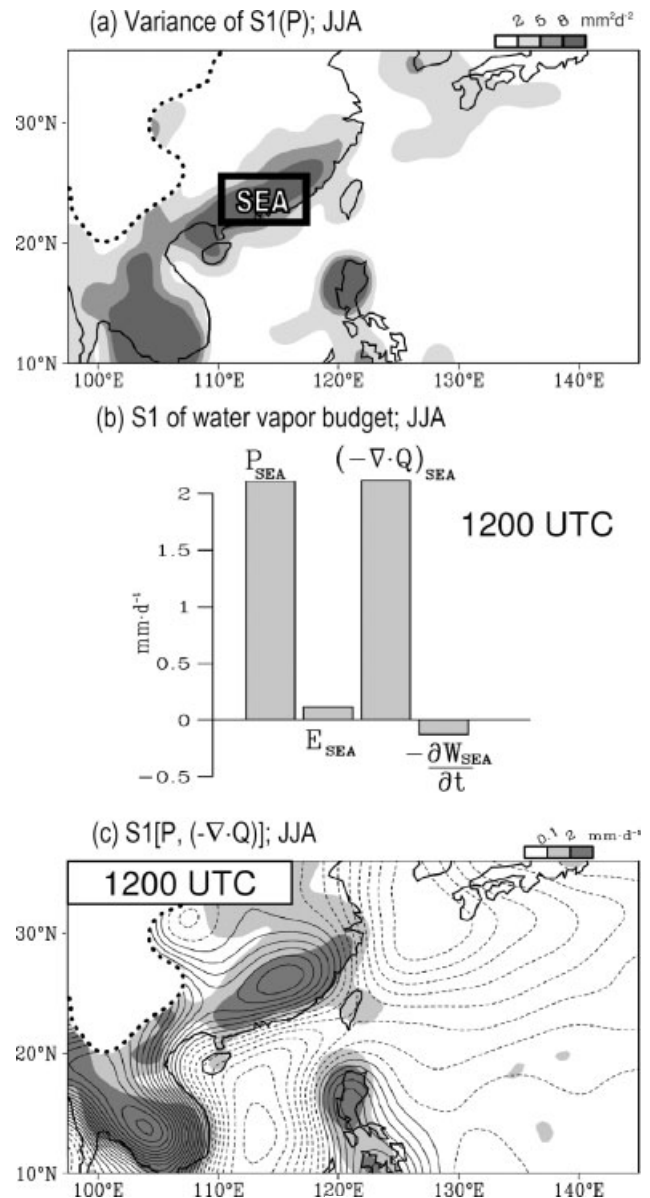
change of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$ . Using Eqs (2)–(3), wind induced by different dynamic forcing terms are reconstructed for the calculation of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$ . As seen from Figure 7,  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  at 0000 UTC is primarily contributed by  $S1(\text{PGF})$  and secondarily contributed by  $S1(\text{CF})$  and  $S1(\text{VDIF}+\text{RES})$  during both summer and winter. Corresponding to the seasonal change of the LSB-like circulation,  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  over south China changes from divergence in summer (Figure 7(b)) to convergence in winter (Figure 7(g)). This feature explains why in Figure 5 the  $S1(P)$  has an elevated structure over south China in winter but not in summer. Because  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  induced from  $S1(\text{PGF})$  only affects  $S1(P)$ , the highest percentage for the seasonal variation of  $S1(\text{SLP})$  impacts on  $\Delta P$  over south China should not exceed the percentage of  $S1(P)$  on explaining the variability of  $\Delta P$  ( $\sim 70\%$ ; Yin *et al.*, 2009).

Dai and Deser (1999) noted that the large-scale surface divergence/convergence is a major factor for modulating the diurnal rainfall over the oceans. Because a local LSB only appears over coastal regions or over big islands, there must be another reason different to the local LSB for causing the divergence/convergence over most of the WNP ocean where the local LSB does not exist. As seen from Figure 7, it is possible that the wind divergence/convergence induced by a large-scale LSB-like circulation leads to the diurnal modulation of rainfall over the WNP ocean.

In this study, a large-scale LSB-like circulation is suggested to be mainly induced by the PGF of  $S1(\text{SLP})$  which involves a non-migrating component caused by local thermal forcing (Lieberman and Leovy, 1995). Because  $S1(P)$  is led by  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$ , the contribution of  $S1(P)$  from a local LSB that is induced by local thermal forcing might be estimated as:

$$\left[ \frac{\text{Variance of } S1(-\nabla \cdot V) \text{ induced by the non-migrating component of } S1(\text{SLP})}{\text{Variance of total } S1(-\nabla \cdot V)} \right]. \quad (5)$$

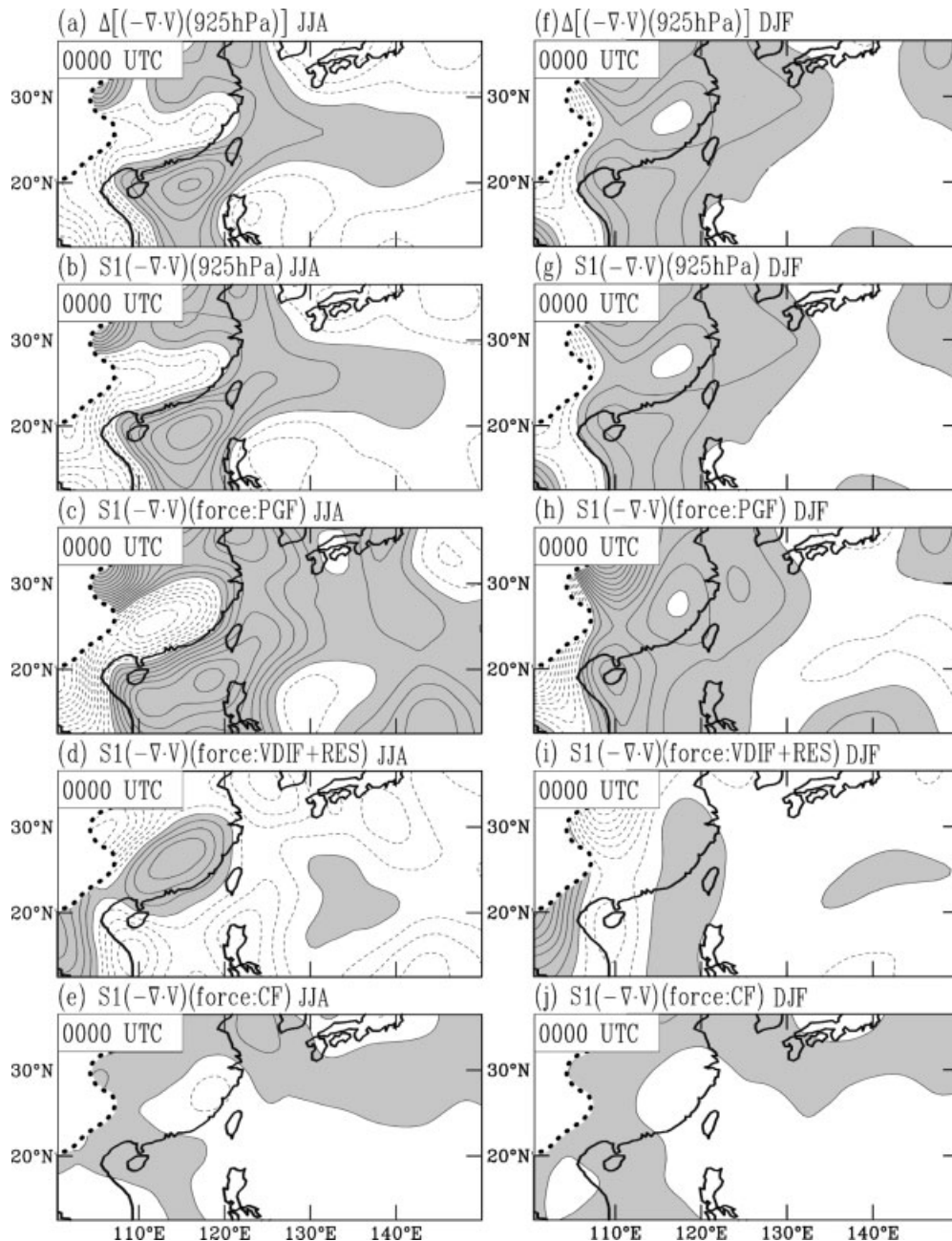
Haurwitz and Cowley (1973) indicated that the wave-number-1 component of  $S1(\text{SLP})$  can be seen as migrating tidal wave, whereas the other wave numbers are called non-migrating tides. Based on Haurwitz and Cowley (1973), we extract the non-migrating component of  $S1(\text{SLP})$  by removing its wave-number-1 component through a wave separation analysis. Then, similar to Figure 7, we calculate  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  induced by the non-migrating component of  $S1(\text{SLP})$  for Eq. (5). An estimation of Eq. (5) area-averaged over SEA, an area marked in Figure 6(a), indicates that about 70% of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  in summer is coming from its non-migrating component, consistent with Chen *et al.* (2001) that the non-migrating tidal wave is dominant over a major land mass. Such a result suggests that the LSB induced by local thermal forcing is important for modulating the diurnal rainfall over the EA continent. This suggestion does not contradict our earlier discussion regarding the role of LSB-like circulation in affecting  $S1(P)$  because the local LSB along the coast of EA is also part of the large-scale LSB-like circulation (see Figure 1). As for another selected area ( $125^\circ\text{--}140^\circ\text{E}$ ,  $21^\circ\text{--}25^\circ\text{N}$ ) over the WNP ocean, about 40% of  $S1[(-\nabla \cdot V)(925 \text{ hPa})]$  is contributed by its non-migrating component, consistent with Chen *et al.* (2001) that the migrating tidal wave is dominant over oceans.



**Figure 6.** (a) Variance of 3-hourly  $S1(P)$  during the summer months. (b) Diagnosis of the water vapour budget extracted from GEOS5 for the maintenance of  $S1(P)$  at 1200 UTC area-averaged over southeast Asia (SEA,  $110^\circ\text{--}118^\circ\text{E}$ ,  $21^\circ\text{--}25^\circ\text{N}$ ), boxed areas in (a) in summer. The evaporation, precipitable water, and convergence of water vapour flux are denoted by  $E$ ,  $W$ , and  $(-\nabla \cdot Q)$  in (b). (c) The JJA mean of  $S1(P)$  superimposed with  $S1(-\nabla \cdot Q)$  at 1200 UTC. The mountain areas in (b) are blocked. The contour interval of  $(-\nabla \cdot Q)$  in (c) is  $0.75 \text{ mm} \cdot \text{d}^{-1}$ .

## 5. Conclusion

The diurnal wind variations over the EAWNP region consist of a mountain–valley breeze over the slopes of the Tibetan Plateau and an LSB near the coast, both of which are local features. A typical LSB is generally to be considered with a horizontal scale of less than 100 km (Neumann and Mahrer, 1971). Using the global reanalysis data, we find that these local diurnal circulations are apparently coupled with the global-scale diurnal tide and together produce a planetary-scale LSB-like circulation variation covering  $\sim 1000 \text{ km}$  over the EAWNP region. The analysis of momentum budget indicates that the PGF produced from the diurnal harmonic of pressure tidal wave, i.e.  $S1(\text{SLP})$ , is the major dynamic forcing inducing this planetary-scale diurnal circulation



**Figure 7.** The JJA mean of (a) anomaly and (b) diurnal harmonic of 925 hPa wind convergence, i.e.  $(-\nabla \bullet V)$ , at 0000 UTC. The reconstructed  $S1[(-\nabla \bullet V)(925 \text{ hPa})]$  induced by (c) PGF, (d) VDIF+RES, and (e) CF. (f)–(j) are similar to (a)–(e), but for DJF. The contour interval in all panels is  $5 \times 10^{-7} \text{ s}^{-1}$ .

feature. Moreover, we suggest that the thermal forcing exciting the migrating component of  $S1(\text{SLP})$  (e.g. Lindzen, 1967; Groves and Wilson, 1982), in addition to the land–sea differential heating, helps generate this large-scale LSB-like circulation.

Pronounced seasonal variation of this planetary-scale LSB-like circulation between summer and winter is also found, with summer featuring broader coverage of the circulation. Such a seasonal difference is mainly caused by seasonal variations of the diurnal tidal wave, as inferred from the difference of  $S1(Z)$  between summer and winter. Examinations of equivalent potential temperature and water vapour budget further indicate that the seasonal varying  $\theta_e$  causes the seasonal difference of LSB-like circulation which generates the surface convergence and leads to the seasonal change of water vapour convergence. This feature assists

in the phase change of diurnal precipitation over the EA continent from the late afternoon peak in summer to the morning peak in winter, as noted in Li *et al.* (2008). Because the low-level circulation is important in representing the physical processes acting above the surface and controls the timing of the precipitation, this planetary-scale LSB-like circulation and the impact of such a circulation on the seasonal difference of diurnal precipitation timing deserve attention for future modelling research.

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