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*Technical Report 32-1350*

*A Program for Computing Shock-Tube  
Gasdynamic Properties*

*T. E. Horton*

*W. A. Menard*

**JET PROPULSION LABORATORY  
CALIFORNIA INSTITUTE OF TECHNOLOGY  
PASADENA, CALIFORNIA**

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**TECHNICAL REPORT 32-1350**

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## **Preface**

The work described in this report was performed by the Environmental Sciences Division of the Jet Propulsion Laboratory (JPL).

One of the authors, T. E. Horton, is currently on the staff of the University of Mississippi.



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## **Abstract**

A computer program for calculating the chemical equilibrium properties associated with moving, standing, and reflected normal shocks has been developed. The program calculates thermodynamic properties and chemical composition from basic spectroscopic data. Both dissociation and ionization over an unlimited temperature range can be considered. The chemical equilibrium and composition equations are developed from statistical mechanics, and the procedure for solving the coupled equilibrium and normal shock equations is presented. Flow diagrams, operating instructions, and listings of the program and input data are included. The program can treat initial mixtures consisting of up to ten gases. In this report, input data are listed for sixty-six species which are composed of carbon, nitrogen, oxygen, hydrogen, argon, neon, and helium.



# A Program for Computing Shock-Tube Gasdynamic Properties

## I. Introduction

Investigations of atomic and molecular phenomena utilizing shock tubes as well as studies of the feasibility of planetary entry at high velocities have created a need for rapid solution techniques for aerothermochemistry problems involving the high-temperature properties of gas mixtures. Although most problems in the field of aerothermodynamics are characterized by large force, chemical composition, and temperature gradients and thus by nonequilibrium phenomena, there are many processes of interest in this field for which either the gradients disappear or the reaction rates are sufficiently fast so that the assumption of chemical equilibrium is rational. This report, as in an earlier publication (Ref. 1), is confined to the problem of supplying thermodynamic properties and chemical composition of multi-component, high-temperature, real-gas mixtures which satisfy the conditions of chemical equilibrium.

The subject of equilibrium properties of dissociated and ionized gases is of interest in a variety of fields ranging from astrophysics to re-entry system design, and has been an active area of research as witnessed by the literature (Refs. 2-37). A review of this material is given in

Appendix A. These published works can be categorized as those dealing primarily with (1) the physics of the subject, (2) the computational or formulational aspects of the subject, and (3) the presentation of computed property data. Work in the third category can be further divided into those in which long range interactions between particles are either neglected or included in some manner. Some authors term the former ideal or perfect real-gas mixtures and the latter as non-ideal or imperfect real-gas mixtures.

A variety of gases have been treated. However, data are usually presented in the form of a Mollier chart or tabulation with composition data presented only for a limited number of pressure or density levels. When this information is used to reduce data taken in a shock-tube experiment, the resulting uncertainty and effort are undesirable. The JPL Thermochemistry and Normal-Shock Computer Program was originally designed to alleviate this difficulty by completely mechanizing the process of generating parameters for moving, standing, and reflected shock conditions as experienced in shock tubes. As a result of this ability to handle a variety of possible gas-mixture combinations, the program has been used extensively by the staff at JPL.

However, because of the method used in handling the species' thermochemistry data, the temperature range was limited. Also, because the computations could be made only at temperatures for which library data were available, the determination of standing bow shock and reflected shock gas compositions was not possible except by plotting Mollier data computed for neighboring temperatures.

To eliminate these difficulties, a revised program has been written, one in which our goal was to retain the versatility and reliability of the old program, but in which the species' thermochemistry data, i.e., specific enthalpy and free energy, are generated internally from spectroscopic data. The generation of the specific enthalpy and free energy was to be handled by a subroutine so that the model used in computing these could be changed. The subroutine and problems associated with the computation of the partition function and thermochemistry data are discussed in Ref. 38. Also, the output format was to be handled through a subroutine which could be modified with ease.

The result of this revision is the present program which is written in Fortran II for use on an IBM 7090 or 7094 computer. The program calculates chemical equilibrium composition and the parameters associated with traveling, standing, and reflected normal shocks. The initial gas mixture, temperature, and pressure can be varied with ease. The program can treat initial mixtures consisting of up to ten gases. Both dissociation and ionization over an unlimited temperature range can be considered. The computational capacity of the program is a mixture of 100 different species composed of 10 different elements. These numbers are limited only by the initial dimension statement and can be increased, providing computer storage is available. The output is a complete thermodynamic and chemical description of the gas which satisfies the Hugoniot equation for the important normal shock cases which are useful in both shock-tube experimentation and free-flight entry calculations.

The program can also be used to generate Mollier thermodynamic data at selected temperatures and densities if this option is desired. A mathematical statement of the criterion for chemical equilibrium is given prior to further discussion of the details of the computer program.

## II. Chemical Equilibrium

In an earlier report (Ref. 1), the mass action and mass balance equations, which yield the chemical composition

for chemical equilibrium, were developed within the context of classical thermodynamics as in (Refs. 39, 40, and 41). In this section these same equations are developed with statistical mechanical arguments. The equations are then compared to the classical expressions to show the well known relationship between equilibrium constant, free energy, and partition function.

The problem is one of determining the most probable stationary distribution of atoms and electrons among the energy levels of chemical components of the mixture, subject to the constraint that the total mass of each element specie remain constant. Thus, for a mixture composed of  $m$  number of elemental species which combine to form  $s$  number of compounds, the compounds will be identified by the numbers  $1, 2, \dots, s$  and the elements by the numbers  $s + 1, s + 2, s + 3, \dots, s + m$ . Further, let the symbols  $X^{(s+1)}, X^{(s+2)}, X^{(s+3)}, \dots, X^{(s+m)}$  be used to represent the usual symbols for the chemical elements so that the compounds and ions of the mixture can be written as

$$Y^{(p)} = X_{\alpha_{p, s+1}}^{(s+1)} X_{\alpha_{p, s+2}}^{(s+2)} X_{\alpha_{p, s+3}}^{(s+3)} \dots X_{\alpha_{p, s+m}}^{(s+m)}$$

where  $p = 1, 2, 3, \dots, s$ . The  $\alpha_{ij}$  indicates the number of atoms of the  $j$ th numbered element in the  $i$ th numbered specie. The  $\alpha_{ij}$ 's come from the set of positive and negative real integers which include zero. The array of  $(\alpha_{p, s+1}, \alpha_{p, s+2}, \alpha_{p, s+3}, \dots, \alpha_{p, s+m})$  which represents the essentials of the chemical formula for the  $p$ th compound is termed the formula vector for the  $p$ th specie. The introduction of this concept by Brinkley (Refs: 42, 43, and 44) allows one to use the theory of linear vector space to formulate problems in stoichiometry.

To determine the arrangement of elementary specie among the  $(s + m)$  different chemical combinations and the distribution of the chemical combinations among their  $q_k$  number of energy levels  $\epsilon_1^{(k)}, \epsilon_2^{(k)}, \epsilon_3^{(k)}, \dots, \epsilon_{q_k}^{(k)}$ , each energy level having degeneracy  $g_1^{(k)}, g_2^{(k)}, g_3^{(k)}, \dots, g_{q_k}^{(k)}$ , respectively, one must find an array of occupation numbers or concentrations  $N_1^{(k)}, N_2^{(k)}, N_3^{(k)}, \dots, N_{q_k}^{(k)}$  for  $k = 1, 2, 3, \dots, s + m$ , which represents their most probable arrangement. The most probable arrangement is the one that can be achieved in the greatest number of ways. If the assumptions are made that the long-range interactions between particles of the mixture are negligible so that the collision times are small compared to times between collisions and that the particles of a particular chemical specie are indistinguishable from each other, then from the theory of combinatory analysis we have the

number of ways  $W_i^{(k)}$  in which  $N_i^{(k)}$  objects can be distributed in  $g_i^{(k)}$  distinguishable categories when there is no restriction on the number in each category given by

$$W_i^{(k)} = \frac{(g_i^{(k)})(g_i^{(k)} + 1)(g_i^{(k)} + 2) \cdots (g_i^{(k)} + N_i^{(k)} - 1)}{N_i^{(k)}!} \quad (1)$$

For the usual chemical system,  $g_i^{(k)} \gg N_i^{(k)}$ ; so

$$W_i^{(k)} = \frac{(g_i^{(k)})^{N_i^{(k)}}}{N_i^{(k)}!} \quad (2)$$

Thus, for each specie the number of arrangements of its  $N^{(k)}$  particle with  $N_1^{(k)}, N_2^{(k)}, N_3^{(k)}, \dots, N_{q_k}^{(k)}$  in each energy level is

$$W^{(k)} = \prod_{i=1}^{q_k} \frac{(g_i^{(k)})^{N_i^{(k)}}}{(N_i^{(k)})!} \quad (3)$$

And the number of arrangements for the mixture is

$$W = \prod_{k=1}^{s+m} W^{(k)} \quad (4)$$

The natural logarithm of the above expression can be written as

$$\ln W = \sum_{k=1}^{s+m} \sum_{i=1}^{q_k} N_i^{(k)} \left( \ln \frac{g_i^{(k)}}{N_i^{(k)}} + 1 \right) \quad (5)$$

with the aid of Stirling's approximation, i.e.,

$$\ln N! = N(\ln N - 1)$$

for  $N \gg 1$ . For the most probable arrangement, the above expression must be an extremum for variations in the  $N_i^{(k)}$ 's subject to the constraints ① that the total number or mass of elemental species is a constant, i.e.,

$$\sum_{j=1}^{q_e} dN_j^{(e)} + \sum_{k=1}^s \sum_{i=1}^{q_k} \alpha_{ke} dN_i^{(k)} = 0 \quad (6)$$

for  $e = s+1, s+2, s+3, \dots, s+m$  and ② that the total energy of the system remain constant, i.e.,

$$\sum_{k=1}^{s+m} \sum_{i=1}^{q_k} \epsilon_i^{(k)} dN_i^{(k)} = 0 \quad (7)$$

(Note that the above expression assumes that the energies of the individual states are independent of variations in the concentrations. This is valid only when the long-range interactions are negligible.) For  $W$  to have an extreme value, the differential of Eq. (5) must be zero; so

$$\sum_{k=1}^{s+m} \sum_{i=1}^{q_k} \ln \frac{g_i^{(k)}}{N_i^{(k)}} dN_i^{(k)} = 0 \quad (8)$$

Now using the technique of Lagrangian multipliers, we find from the combination of Eqs. (7) and (8) with the  $m$  equations of the type of Eq. (6) that

$$N_i^{(k)} = g_i^{(k)} \prod_{e=s+1}^{s+m} (\exp \gamma_e)^{\alpha_{ke}} (\exp \beta \epsilon_i^{(k)}) \quad (9)$$

for  $i = 1, 2, 3, \dots, q_k$  and  $k = 1, 2, 3, \dots, s$  and that

$$N_j^{(e)} = g_j^{(e)} \exp(\gamma_e) \exp(\beta \epsilon_j^{(e)}) \quad (10)$$

for  $j = 1, 2, 3, \dots, q_e$  and  $e = s+1, s+2, \dots, s+m$ . Summing the above expressions over all energy states gives

$$N^{(k)} = \sum_{i=1}^{q_k} N_i^{(k)} = \prod_{e=s+1}^{s+m} (\exp \gamma_e)^{\alpha_{ke}} \sum_{i=1}^{q_k} g_i^{(k)} \exp(\beta \epsilon_i^{(k)}) \quad (11)$$

for  $k = 1, 2, 3, \dots, s$  and

$$N^{(e)} = \sum_{j=1}^{q_e} N_j^{(e)} = \exp \gamma_e \sum_{j=1}^{q_e} g_j^{(e)} \exp(\beta \epsilon_j^{(e)}) \quad (12)$$

for  $e = s+1, s+2, \dots, s+m$ . The partition function  $Q_l$  for the  $l$ th specie is defined as

$$Q_l = \sum_{j=1}^{q_l} g_j^{(l)} \exp(\beta \epsilon_j^{(l)}) \quad (13)$$

where the undetermined Lagrangian multiplier  $\beta$  is  $(-1/kT)$ . Expressions of the type of Eqs. (11) and (12) can be combined to eliminate the remaining undetermined Lagrangian multipliers  $\gamma_{s+1}, \gamma_{s+2}, \gamma_{s+3}, \dots, \gamma_{s+m}$  and yield

$$\frac{N^{(i)}}{\prod_{j=s+1}^{s+m} (N^{(j)})^{\alpha_{ij}}} = \frac{Q_i}{\prod_{j=s+1}^{s+m} (Q_j)^{\alpha_{ij}}} \quad (14)$$

or for

$$n^j = \frac{N^j}{AV}$$

where  $V$  is the volume of the mixture and  $A$  is Avagadro's number.

$$\frac{n^{(i)}}{\prod_{j=s+1}^{s+m} (n^{(j)})^{\alpha_{ij}}} = \frac{\left(\frac{Q_i}{VA}\right)}{\prod_{j=s+1}^{s+m} \left(\frac{Q_j}{VA}\right)^{\alpha_{ij}}} \quad (15)$$

for  $i = 1, 2, 3, \dots, s$ . Equations of the type of Eq. (15) are identical to the "equations of mass action" given in Ref. 1 providing

$$-\frac{f_i^0}{RT} = \ln(RT) + \ln\left(\frac{Q_i}{VA}\right) \quad (16)$$

Furthermore, from the integral of Eq. (7), the specific enthalpy can be shown to be

$$\frac{h_i^0}{RT} = 1 + \frac{\sum_{j=1}^{g_i} \frac{\epsilon_i}{KT} \exp\left(-\frac{\epsilon_i}{KT}\right)}{Q_i} = 1 + KT \left( \frac{\partial \ln(Q_i)}{\partial KT} \right)_v \quad (17)$$

The other constraints on the concentrations supplied by the conservation of mass of chemical elements require that

$$n_j = \sum_{i=1}^s \alpha_{ij} n_i = B_j \rho / \rho_o \quad (18)$$

for  $j = s+1, s+2, s+3, \dots, s+m$  where

$B_j$  = concentration of the  $j$ th element at standard temperature and pressure; and

$\rho_o$  = density of the mixture at standard temperature and pressure.

Equation (18) is frequently referred to as the "equation of mass balance." The solution of the  $s$  mass action equations and the  $m$  mass balance equations for a given temperature, density, and concentration of elements at standard conditions is a set of  $s+m$  concentrations which satisfy the criterion of chemical equilibrium.

After the equilibrium concentrations have been determined, the thermodynamic properties of the gaseous mixture are calculated by the following relations:

$$\begin{aligned} n &= \sum_{i=1}^{s+m} n_i \\ P &= nRT \\ H &= \frac{\sum_{i=1}^{s+m} n_i h_i^0}{\rho} \\ Z &= \frac{\frac{P}{\rho RT}}{\frac{P_o}{\rho_o RT_o}} = \frac{M_o}{M} \end{aligned} \quad (19)$$

$$\frac{SM_o}{R} = \frac{M_o}{\rho} \sum_{i=1}^{s+m} n_i \left( \frac{h_i^0 - f_i^0}{RT} - \ln P - \ln \frac{n_i}{n} \right)$$

The quantities which appear above are defined as:

$h_i^0$  specific enthalpy of the  $i$ th constituent at one atmosphere

$M_i$  molecular weight; whether in lb/lb-mole or g/g-mole, the values are identical

$M_o$  molecular weight of mixture at STP

$n$  total concentration of mixture, moles/liter

$n_i$  concentration of the  $i$ th constituent

$R$  universal gas constant  $\approx 1.9860$  Btu/ $^{\circ}$ R-lb-mole  
 $\approx 1.98726$  cal/ $^{\circ}$ K-g-mole

$S$  entropy

$\frac{SM_o}{R}$  dimensionless entropy

$T$  temperature

$T_o$  standard temperature  $\approx 273.15^{\circ}$ K

$\rho$  mass density

### III. Normal Shock Problem

From the aerothermodynamic point of view, one of the most important yet simplest applications of the chemical equilibrium data described above is the normal shock problem. The equations governing the variation of properties through a normal shock, when the coordinate system is chosen so the flow is steady and one dimensional,

and when the effects of body forces, diffusion, and radiative transfer are negligible, are:

*continuity:*

$$\frac{d}{dx}(\rho U) = 0 \quad (20)$$

*linear momentum:*

$$\rho U \frac{dU}{dx} + \frac{dP}{dx} = \frac{1}{2} \frac{d}{dx} \left( \eta \frac{dU}{dx} \right) \quad (21)$$

*and energy:*

$$\rho U \frac{dE}{dx} + P \frac{dU}{dx} = \frac{d}{dx} \left( K \frac{dT}{dx} \right) + \frac{1}{2} \eta \left( \frac{dU}{dx} \right)^2 \quad (22)$$

Integration of these equations with respect to  $x$  yields

$$\rho U = C_1 \quad (23)$$

$$\rho U^2 + P = \frac{1}{2} \eta \frac{dU}{dx} + C_2 \quad (24)$$

$$\rho U \left( H + \frac{U^2}{2} \right) = \left( K \frac{dT}{dx} + \frac{1}{2} \eta U \frac{dU}{dx} \right) + C_2 \quad (25)$$

where  $C_1$ ,  $C_2$ , and  $C_3$  are constants. If these equations are applied across a normal shock, and if either: (1) the fluid is assumed to be inviscid ( $\eta = 0$ ) and nonheat-conducting ( $K = 0$ ), or (2) the fluid is assumed to be in thermodynamic equilibrium (i.e., the gradients are zero), then the equations reduce to the form

$$\rho_A U_A = \rho_B U_B \quad (26)$$

$$\rho_A U_A^2 + P_A = \rho_B U_B^2 + P_B \quad (27)$$

$$H_A + \frac{U_A^2}{2} = H_B + \frac{U_B^2}{2} \quad (28)$$

Therefore, Eqs. (26), (27), and (28) relate the thermodynamic properties and velocities in the up-stream and down-stream equilibrium regions of the flow field associated with a standing normal shock as shown in Fig. 1.

Three special applications of Eqs. (26), (27), and (28) to the equilibrium regions of shock flows are (1) the moving shock problem, (2) the standing shock problem, and (3) the reflected shock problem.

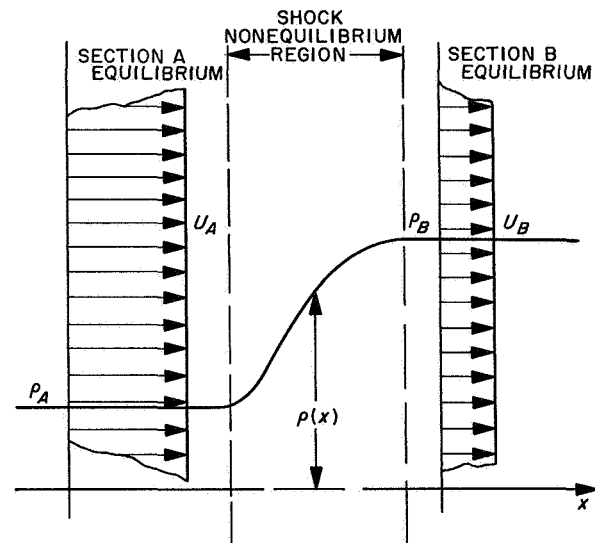


Fig. 1. Flow field associated with a normal shock

#### A. Free-Flight or Moving Shock Case

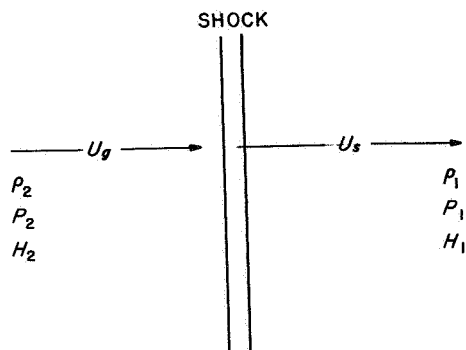
The free-flight or moving shock problem is of interest for determining the velocities and thermodynamic properties behind the normal portion of the bow shock associated with a high velocity blunt body. The free-flight problem is equivalent to the moving shock problem in a shock tube. Physically, the problem is the determination of the properties behind a shock which is propagating in quiescent gas at a specified velocity. Mathematically, the problem is to satisfy the Hugoniot equation,

$$H_B - H_A = \frac{1}{2} (P_B - P_A) \left( \frac{1}{\rho_A} + \frac{1}{\rho_B} \right) \quad (29)$$

which is obtained from a suitable manipulation of Eqs. (26), (27), and (28). The initial properties  $H_A$ ,  $P_A$ , and  $\rho_A$  are known, and the chemical equilibrium computation yields relations for  $H_B(\rho_B, T_B)$  and  $P_B(\rho_B, T_B)$ . Thus, by selecting one property in region (B), the other properties are determined by Eq. (29). The Hugoniot equation and the auxiliary equations are shown in Figs. 2 and 3 with subscripts appropriate to the free-flight and moving shock problems.

#### B. Standing Shock Case

The standing shock situation considered here results from the flow of the high-velocity, high-enthalpy gas associated with the above mentioned moving normal shock case over a stationary blunt body. The flow around the body is such that a stationary bow shock develops.



EQUATIONS TO BE SATISFIED:

$$H_2 - H_1 = \frac{1}{2} (P_2 - P_1) \left( \frac{1}{\rho_1} + \frac{1}{\rho_2} \right) \quad (30)$$

$$U_s = \sqrt{\frac{P_2}{\rho_1} \left( \frac{P_2 - P_1}{P_2 - P_1} \right)} = \text{SHOCK VELOCITY, (31)}$$

$$U_g = U_s \left( 1 - \frac{\rho_1}{\rho_2} \right) = \text{GAS VELOCITY (32)}$$

Fig. 2. Moving normal shock

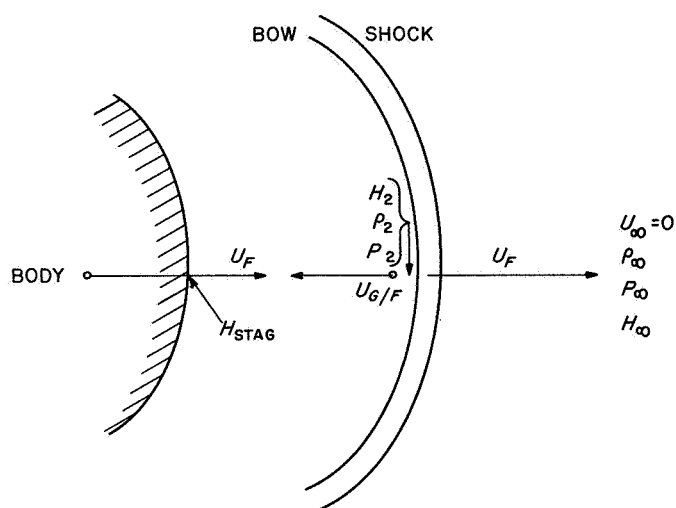
Physically, the problem is one of determining the properties and velocity in the equilibrium region immediately behind the bow shock. The important thermodynamic properties and velocities are indicated in Fig. 4.

Also shown in this figure are the appropriate forms of Eqs. (26), (27), and (28) which for known  $U_g$ ,  $\rho_2$ ,  $H_2$ ,  $P_2$  and the implicit chemical equilibrium property relations  $H_3(T_3, \rho_3)$  and  $P_3(T_3, \rho_3)$  allow the velocities and state of the gas behind the standing shock to be determined. The results of this type of calculation are useful in shock tube simulation of entry conditions.

### C. Reflected Shock Case

The third problem of interest, the reflected shock case, results from the stagnation of the high-velocity, high-enthalpy gas of a moving shock by an immovable plane surface. The gas comes to rest by means of a reflected shock which propagates into the gas flowing toward the surface.

The thermodynamic properties and velocities are defined in Fig. 5 along with the equations to be satisfied. Other than a slightly different form of the equations, the mathematical problem here is equivalent to the standing shock case described earlier. The results of this type of calculation are useful in interpreting experiments performed at the end of a shock tube.



EQUATIONS TO BE SATISFIED:

$$H_2 - H_{\infty} = \frac{1}{2} (P_2 - P_{\infty}) \left( \frac{1}{\rho_{\infty}} + \frac{1}{\rho_2} \right) \quad (33)$$

$$U_F = \sqrt{\frac{P_2}{\rho_{\infty}} \left( \frac{P_2 - P_{\infty}}{P_2 - P_{\infty}} \right)} = \text{FLIGHT VELOCITY (34)}$$

$$U_{G/F} = \sqrt{\frac{P_{\infty}}{\rho_2} \left( \frac{P_2 - P_{\infty}}{P_2 - P_{\infty}} \right)} = \text{VELOCITY OF GAS RELATIVE TO THE BODY (35)}$$

$$H_{\text{STAG}} = H_2 + \frac{1}{2} U_{G/F}^2 = \text{STAGNATION ENTHALPY (36)}$$

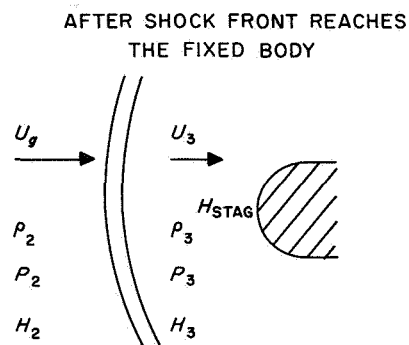
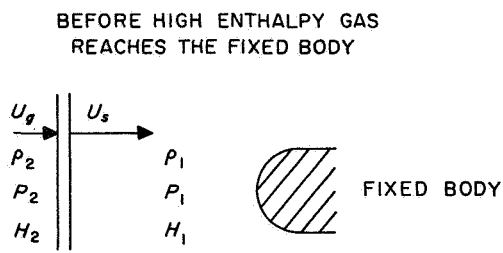
Fig. 3. Free-flight bow shock

Before considering the numerical techniques used to satisfy the chemical equilibrium and shock equations, it is important to emphasize that the above equations closely approximate shock-tube and free-flight conditions only where the shock heated gas is in chemical equilibrium and the effect of wall shear and heat transfer is small.

## IV. Numerical Procedure

### A. Chemical Equilibrium

To obtain the  $s + m$  constituents of a gaseous mixture in chemical equilibrium requires the simultaneous solution of the  $m$  mass balance equations and the  $s$  nonlinear mass action equations. The Newton-Raphson method is well suited for solving such a system of simultaneous algebraic equations with a high-speed digital computer. The Newton-Raphson method applied to the present problem allows a solution to be obtained by an iterative



EQUATIONS TO BE SATISFIED

$$H_3 = H_2 + \frac{U_g^2}{2} \left[ 1 - \left( \rho_2 / \rho_3 \right)^2 \right] \quad (37)$$

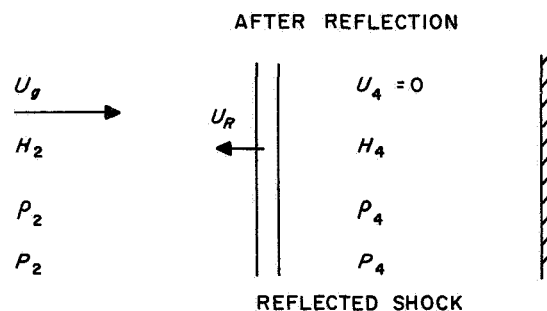
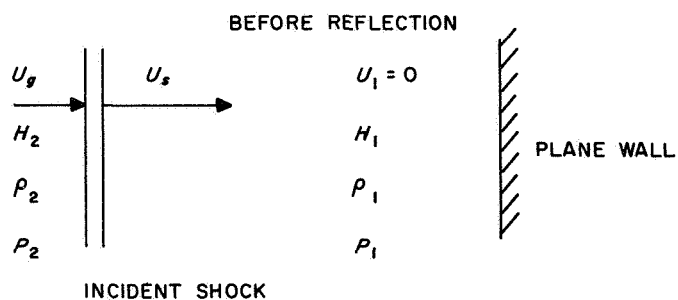
$$p_3 = p_2 + \rho_2 \frac{U_g^2}{2} \left[ 1 - \rho_2 / \rho_3 \right] \quad (38)$$

$$U_3 = \left( \rho_2 / \rho_3 \right) U_g \quad (39)$$

$$H_{STAG} = H_3 + \frac{U_3^2}{2} \quad (40)$$

Fig. 4. Standing shock case

procedure which involves solving a linearized form of Eqs. (15) and (18). The method assumes that a point defined by a set of concentrations which nearly satisfy the mass action and mass balance conditions is known. By expanding the mass action and mass balance equations in a Taylor series about the near solution point and neglecting terms of the series which contain derivatives higher than the first order, a system of linear equations is obtained which approximate the mass action and mass balance equations. The solution of the linear system of equations results in a new set of concentrations which are nearer the actual solution to the mass action and the mass balance equations. Clearly, by repeating this procedure, a better set of concentrations can be found so that by iterating enough times a solution with any desired



EQUATIONS TO BE SATISFIED:

$$H_4 = H_2 + \frac{U_g^2}{2} \left( \frac{\rho_4 / \rho_2 + 1}{\rho_4 / \rho_2 - 1} \right) \quad (41)$$

$$p_4 = p_2 + \rho_2 \frac{U_g^2}{2} \left( \frac{\rho_4 / \rho_2}{\rho_4 / \rho_2 - 1} \right) \quad (42)$$

$$U_R = \frac{U_g^2}{\rho_4 / \rho_2 - 1} = \text{REFLECTED SHOCK VELOCITY} \quad (43)$$

Fig. 5. Reflected shock case

accuracy can be found. With a high-speed digital computer this procedure is quite simple, fast, and accurate.

The linearized form of the mass action (Eq. 15) and mass balance (Eq. 18), obtained from the series expansion with the temperature, density, and initial molar concentration constant, is of the following form:

(1) Mass action.

$$(Z_k)_{p+1} - \sum_{j=s+1}^{s+m} \alpha_{kj} (Z_j)_{p+1} + (D_k)_p = 0 \quad (44)$$

where

$$k = 1, 2, 3, \dots, s$$

$(\ )_p$  = indicates value from the  $p$ th iteration

$$(D_k)_p = \ln(n_k)_p + \ln RT + \frac{f_k^o}{RT}$$

$$- \sum_{j=s+1}^{s+m} \alpha_{kj} \left[ \ln(n_j)_p + \ln RT + \frac{f_j^o}{RT} \right]$$

(2) Mass balance.

$$(n_j)_p (Z_j)_{p+1} + \sum_{k=1}^s \alpha_{kj} (n_k)_p (Z_k)_{p+1} + (D_j)_p = 0 \quad (45)$$

where

$$j = s+1, s+2, s+3, \dots, s+m$$

and

$$(D_j)_p = (n_j)_p + \sum_{k=1}^s \alpha_{kj} (n_k)_p - \frac{B_j \rho}{\rho_o}$$

In the above expressions,

$$(Z_i)_{p+1} = \ln(n_i)_{p+1} - \ln(n_i)_p$$

The use of the natural logarithm of the concentration as the variable in the series expansion was suggested in Ref. 40 and was found to give good results.

This linear system of  $s+m$  equations can be simplified to  $m$  equations by using Eq. (44) to eliminate  $Z_1, Z_2, Z_3, \dots, Z_s$  from Eq. (45). The result of this substitution is a system of  $m$  linear equations of the form:

$$\sum_{j=1}^m (Y_{ij})_p (Z_{s+j})_{p+1} = (C_i)_p \quad (46)$$

where  $i = 1, 2, 3, \dots, m$ . The quantities  $Y_{ij}$  and  $C_i$  are given by

$$(Y_{ij})_p = \delta_{ij} (n_{s+i})_p + \sum_{k=1}^s (n_k)_p (\alpha_{k, s+1}) (\alpha_{k, s+1})$$

$$\delta_{ij} = 0 \quad \text{for } i \neq j$$

$$\delta_{ij} = 1 \quad \text{for } i = j$$

and

$$(C_i)_p = - \left[ (D_{s+i})_p - \sum_{k=1}^s (n_k)_p (D_k)_p (\alpha_{k, s+1}) \right]$$

To use Eq. (46) in obtaining a solution to Eqs. (15) and (18), the following procedure is used:

- (1) Using initial guesses,  $(n_1)_o, (n_2)_o, (n_3)_o, \dots, (n_{s+m})_o$ , for the concentrations at a selected temperature, density, and concentration of chemical elements at standard conditions, the coefficients  $(D_i)_o$  and  $(Y_{ij})_o$  are evaluated.
- (2) A solution to the system of Eq. (46) is found using the Crout method. The solution

$$(Z_{s+1})_1, (Z_{s+2})_1, (Z_{s+3})_1, \dots, (Z_{s+m})_1$$

yields values for

$$(n_{s+1})_1, (n_{s+2})_1, (n_{s+3})_1, \dots, (n_{s+m})_1$$

- (3) The  $(n_1)_1, (n_2)_1, (n_3)_1, \dots, (n_s)_1$  are evaluated using Eq. (44).
- (4) The new values of the concentration  $(n_i)_1$  with  $i = 1, 2, 3, \dots, s+m$  replace the initial guesses, and the above procedure is repeated  $k$  times until

$$1 - \left| \frac{(n)_k}{(n)_{k+1}} \right| < 10^{-6}$$

and

$$1 - \left| \frac{(\rho)_k}{(\rho)_{k+1}} \right| < 10^{-6}$$

and

$$1 - \left| \frac{(\bar{H})_k}{(\bar{H})_{k+1}} \right| < 10^{-6}$$

where

$$(\ )_{k+1} = \text{value from the } (k+1) \text{ iteration}$$

$$n = \text{total concentration}$$

$$\rho = \text{density of mixture}$$

$$\bar{H} = \text{enthalpy per unit volume}$$

This convergence criterion has been found to result in a set of concentrations  $(n_1)_{k+1}, (n_2)_{k+1}, (n_3)_{k+1}, \dots, (n_{s+m})_{k+1}$  which agree with  $(n_1)_k, (n_2)_k, (n_3)_k, \dots, (n_{s+m})_k$  to at least three significant figures.

## B. Moving and Free-Flight Normal Shock Case

The moving shock problem, as discussed earlier, consists of finding a thermodynamic state defined by  $T_2$  and



$\rho_2$  which satisfies the Hugoniot equation for specified initial conditions. The initial conditions are specified by the pressure, temperature, and concentrations. To find a final state which satisfies the Hugoniot equation, a final temperature  $T_2$  is specified and the value of  $\rho_2$  is varied until values of  $p_2$  and  $H_2$  are found which satisfy the Hugoniot equations. After the properties which satisfy the Hugoniot equations are found, the Eqs. (31), (32), (34), (35), and (36) are evaluated.

The systematic procedure for finding the correct value of  $\rho_2$  is based on a Newton-Raphson iteration scheme in which the Hugoniot equation

$$F(v_2, T_2) = [H_2(v_2, T_2) - H_1] - \frac{1}{2} [P_2(v_2, T_2) - P_1] \left[ \frac{1}{\rho_1} + v_2 \right] \quad (47)$$

is expanded in  $v_2$  at some near solution point with  $T_2$  held constant. The resulting iteration relation for the  $n$ th iteration is

$$(v_2)_{n+1} = \left( \frac{1}{\rho_2} \right)_{n+1} = \left( \frac{1}{\rho_2} \right)_n - \frac{F((v_2)_n, T_2)}{\left( \frac{\partial F}{\partial v_2} \right)_{T_2, n}} \quad (48)$$

where the partial derivative is evaluated by

$$\left( \frac{\partial F}{\partial v_2} \right)_{T_2, n} = \frac{F((v_2)_n + (\Delta v_2)_n, T_2) - F((v_2)_n, T_2)}{\Delta v_2}$$

with

$$(\Delta v_2)_n = 10^{-6} (v_2)_n$$

The Hugoniot equation is considered satisfied after  $n$  iterations if:

$$\left| \frac{(1/\rho_2)_{n+1} - (1/\rho_2)_n}{(1/\rho_2)_n} \right| \leq 10^{-3}$$

### C. Reflected and Standing Normal Shock Cases

The procedures followed for the reflected shock case and the standing shock case are similar. Therefore, only the reflected shock case will be discussed below.

The problem in the reflected shock case is to find a  $T_4$  and  $\rho_4$  associated with values of  $H_4$  and  $P_4$  which satisfy Eqs. (41) and (42). The values of  $H_2$ ,  $\rho_2$ ,  $P_2$ , and  $U_g$  are known from the moving normal shock solution.

The solution technique is based upon a series expansion of the thermodynamic properties  $H$  and  $P$  about some point  $(T_a, \rho_a)$  which is near  $(T_4, \rho_4)$ . The truncated series are of the form

$$H = H_a + \left. \frac{\partial H}{\partial \rho} \right|_a \delta \rho + \left. \frac{\partial H}{\partial T} \right|_a \delta T \quad (49)$$

$$P = P_a + \left. \frac{\partial P}{\partial \rho} \right|_a \delta \rho + \left. \frac{\partial P}{\partial T} \right|_a \delta T \quad (50)$$

with

$$\delta \rho = \rho - \rho_a$$

$$\delta T = T - T_a$$

Equations (41) and (42) are functions of  $\rho_4$ . By evaluating these equations at  $\rho_4 = \rho_a$ , approximate values of  $H_4$  and  $P_4$  result. Substituting these approximate values ( $H'_4$  and  $P'_4$ ) into the left-hand side of Eqs. (49) and (50), and combining so that  $\delta T$  is eliminated, yields

$$\delta \rho = \frac{(H'_4 - H_a) \left. \frac{\partial P}{\partial T} \right|_a - (P'_4 - P_a) \left. \frac{\partial H}{\partial T} \right|_a}{\left. \frac{\partial P}{\partial T} \right|_a \left. \frac{\partial H}{\partial \rho} \right|_a - \left. \frac{\partial P}{\partial \rho} \right|_a \left. \frac{\partial H}{\partial T} \right|_a} \quad (51)$$

The partial derivatives are evaluated by the finite difference relations

$$\begin{aligned} \left. \frac{\partial P}{\partial T} \right|_a &= \frac{P(T_a + \Delta T, \rho_a) - P(T_a, \rho_a)}{\Delta T} \\ \left. \frac{\partial H}{\partial T} \right|_a &= \frac{H(T_a + \Delta T, \rho_a) - H(T_a, \rho_a)}{\Delta T} \\ \left. \frac{\partial H}{\partial \rho} \right|_a &= \frac{H(T_a, \rho_a + \Delta \rho) - H(T_a, \rho_a)}{\Delta \rho} \\ \left. \frac{\partial P}{\partial \rho} \right|_a &= \frac{P(T_a, \rho_a + \Delta \rho) - P(T_a, \rho_a)}{\Delta \rho} \end{aligned}$$

where  $\Delta \rho / \rho = 10^{-4}$  and  $\Delta T = 500^\circ \text{K}$ . The  $500^\circ \text{K}$  increment was found to give the best results.

Now, the value of  $\delta \rho$  obtained from the evaluation of Eq. (51) yields a density  $(\rho)_{n+1} = (\rho)_n + \delta \rho$  which is closer to the desired value  $\rho_4$ . By reevaluating the quantities in the right-hand side of Eq. (51) at the new value of density, further improvements result. When

$$\frac{(\delta \rho)_n}{(\rho_4)_n} \leq 10^{-3}$$

the value of  $(\rho_4)_{n+1}$  is considered to be a solution. Using this value for  $\rho_4$ , the values of  $H_4$  and  $P_4$  are evaluated using Eqs. (41) and (42), and the value of  $T_4$  is obtained from Eq. (49). Once convergence is achieved, equilibrium computations are performed at the values of  $T_4$  and  $\rho_4$  to obtain the actual chemical concentrations and properties at the solution point.

## V. Program Description

The numerical solution procedures discussed earlier have been programmed so that the following combinations of problem cases can be considered:

- (1) *The Mollier data case* in which the program calculates the chemical equilibrium properties and composition of a gaseous mixture of specified temperature, density, and initial composition. The computation takes about 2 s\* per temperature-density point.
- (2) *The moving normal shock case* in which the program calculates the properties of a moving shock when the temperature ( $T_2$ ), initial temperature ( $T_1$ ), pressure ( $P_1$ ), and chemical composition are specified. The computer requires about 5 s for each solution.
- (3) *A combination case* in which the normal moving shock case is first solved and then a solution for the resulting *standing shock* is obtained. The shock parameters and chemical compositions for both situations make up the output for this case. The time required for this case is 20 s per moving shock.
- (4) *A combination case* in which the normal moving shock case is first solved and then a solution for the resulting *reflected shock* is obtained. The shock parameters and chemical compositions for both situations makes up the output data. The time requirement for this case is identical to that of (3) above.
- (5) *A combination case* in which the normal moving shock solution is first obtained, then *both standing and reflected solutions* are generated. The shock parameters and chemical compositions for all three conditions of thermodynamic equilibrium are printed out. This case takes 45 s per moving normal shock.

\*The computation time is dependent upon the number of species. The times quoted in the five cases are for problems having 20 species in the final composition. As an example of this time dependence in the moving-standing-reflected shock combination case, the computation takes 15 s when 10 species are used, 45 s for 20 species, and 2 min for 40 species.

The program has sufficient capacity to store in its library data up to 130 chemical species. Of these, 32 can be atomic or atomic ions, 20 can be diatomic species, and 75 can be polyatomic species. The other species are the double-ionized helium, the proton for ionized hydrogen, and the electron. These above limitations are somewhat arbitrary and may be modified by simply changing dimension statements in the program providing storage capacity is available in the computer system. The program library consists primarily of spectroscopic data from which the energy levels of the various species can be deduced.

These library data are used by a thermochemical property subroutine to compute the partition functions and subsequently the specific free energies and enthalpies. In this way the program can be updated as more precise energy-level or heat-of-formation data become available. Also, the use of a subroutine allows the procedure used in computing the partition functions to be changed with ease. The subroutine and the problems in computing partition functions are discussed in an accompanying report (Ref. 38).

Also, to allow easy modification of the output format, the output data for all five cases are handled by a common subroutine.

The gross features and logic of the program are shown in Fig. 6. For a detailed description, the flow diagram and program listing in Appendices B and D, respectively, should be helpful.

## VI. Program Input

The input data for the program is divided into two decks. The first deck, referred to as "library data," contains the data from which thermochemical properties (specific enthalpy and free energy) for each chemical species can be computed using the techniques of statistical mechanics. The entire library data deck is read into the machine following the program deck.

The second deck of input data, referred to as "case data," is used to identify the particular case (i.e., Mollier data, moving normal shock, or combined normal shock), to specify which species of the library data constitute the mixture, and to specify certain initial and final state conditions. This division of the input data into two parts, allows the thermochemistry data in the library to be used without rearrangement or modification, while making it accessible for modification if better values become available. The order of the program deck, library data deck, and case data decks is shown in Fig. 7.

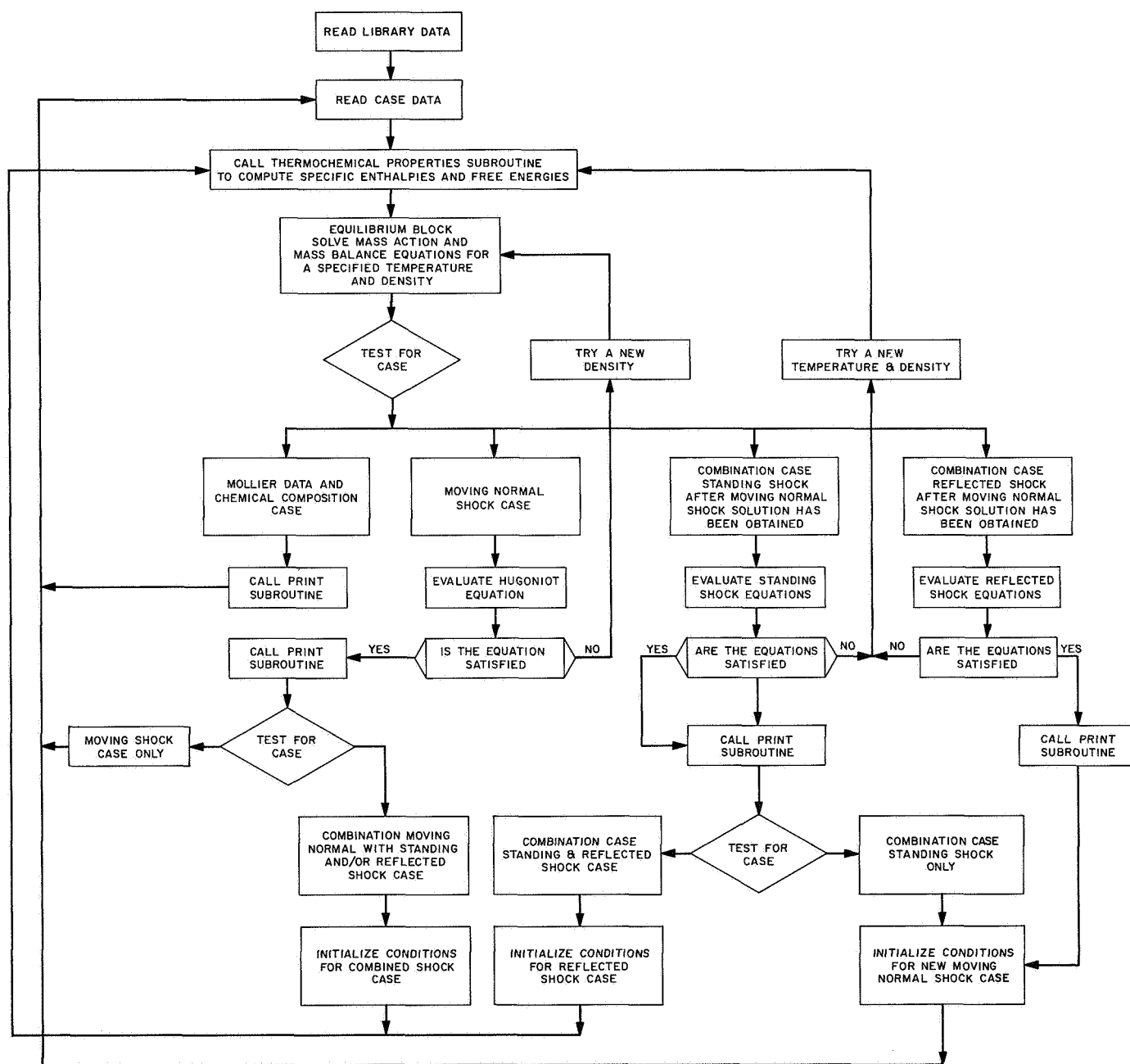


Fig. 6. Flow diagram for the program

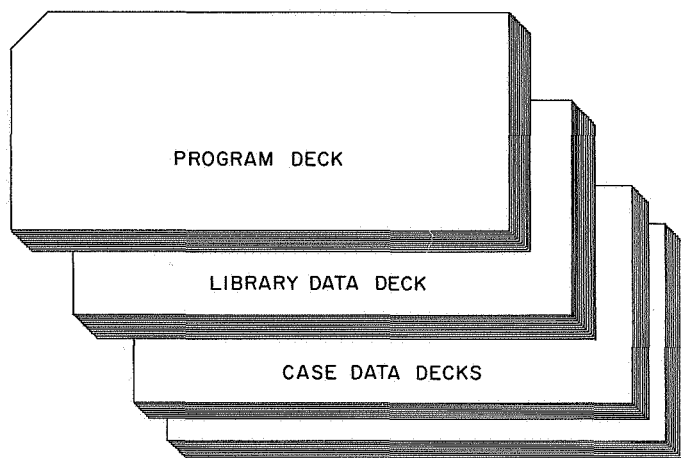


Fig. 7. Order of the program and data decks

#### A. Library Data

The library data deck consists of packets of cards for each species which contain the formula vector, molecular weight, heat of formation, and sufficient information about the structure and energy states to allow one to compute the partition function and the specific enthalpy and

free energy. Because the information used to represent the energy states and structure as well as the details of the above computation for atomic species is quite different from that for diatomic species, each library data packet takes one of five different forms. The form or format for four of these—(1) atomic or atomic ion, (2) diatomic molecule or ion, (3) linear polyatomic molecule, and (4) non-linear polyatomic molecule—is given in Table 1. The fifth type of species contains those particles which lack internal structure. For the energy range considered here, these are the electron, the proton or hydrogen positive ion, and the helium double-positive ion. Each of the packets is assigned a library identification number which corresponds to the position of the packet in the deck. The identification number is for the user's convenience as the computer does not read it but assigns identification numbers based upon the order of the packet in the deck. Thus, if the order of deck is changed by adding or deleting or rearranging packets, the packets should be renumbered so that they can be properly identified. The only constraint on the order of the packets is that those species used as "components" of the formula vectors be listed last and in an order such that the matrix of their formula vectors is a unitary matrix. For illustrative purposes, we

Table 1. Example data packets

#### A. Atomic and atomic ions

11	C++	1	0	0	0	0-2	12.01005	01	20	0.0	Species identification card
C++	HEAT FORMATION	0.991689E	03								Heat of formation card
CIII		1.00				0.0					Energy level block
CIII		5.00				52394.797					
CIII		4.00				52338.000					
CIII		3.00				102351.375					
CIII		9.00				137450.500					
CIII		5.00				145875.062					
CIII		1.00				182520.187					
CIII		4.00				238160.687					
CIII		12.00				259672.062					
CIII		15.00				269962.875					
CIII		5.00				276482.687					
CIII		12.00				308264.750					
CIII		16.00				317748.000					
CIII		78.00				321964.687					
CIII		49.00				333395.000					
CIII		116.00				346526.000					
CIII		90.00				358639.000					2
CIII		29.00				365585.000					
CIII		39.00				373748.000					
CIII		37.00				381958.000					

Table 1 (contd)

## B. Diatomic

21	N2+	0	2	0	0	0-1	28.016	02	4	2.0	Specie identification card
N2+	HEAT FORMATION	0.359306E	03								Heat of formation card
N2+	X DOUBLET SIGMA	0.0	2207.23	16.22	1.9322	0.0202	2.0				Energy level block
N2+	A DOUBLET PI	9020.0	1902.84	14.91	1.722	0.018	4.0				
N2+	B DOUBLET SIGMA	25570.0	2419.84	23.19	2.083	0.0195	2.0				
N2+	C DOUBLET SIGMA	64550.0	2050.0	14.92	1.65	0.05	2.0				

## C. Linear polyatomic

07	C02	1	0	2	0	0	0	0	0	0	0	44.011	03	4	0.0	Specie identification card
C02	HEAT FORMATION	-	939643E	+02												Heat of formation card
C02	SINGLET SIGMA	2.01.0	0.3906	0.0022	4											Energy level block
	667.33															
	667.33															
	1342.86															
	2349.3															

## D. Non-linear polyatomic

05	N02	0	1	2	0	0	0	0	0	0	0	46.008	04	3	0.0	Specie identification card
N02	HEAT FORMATION	0.876600E	+01													Heat of formation card
N02		2.02.0	3.47500E	-4063.54500E	-4067.02000E	-4003										Energy level block
N02		.756.8														
N02		1357.8														
N02		1665.5														

will consider an arrangement of the library data packets as given in Table 2. A more extensive list of library data is given in Appendix E.

The library data are arranged with Card No. 1 first, followed by the species data packets. The species data packets consist of a first card which is called the species identification card, a heat of formation card, and the energy level block.

## 1. Card No. 1

The information on Card No. 1 is in format 2I2.\*\* The first word on this card gives the number of chemical components or elements, plus the electron, which are listed in the library. The second word is the total number of species in the library. For the sample listing in Table 2, the first card would contain 05 in the first two columns,

since C, N, O, e, A are present, and 36 in the second two columns.

## 2. Data Packets

*Specie Identification Card:* The first card of each packet contains in the following order:

- (1) The specie identification number.
- (2) The chemical formula.
- (3) The  $\alpha_{ij}$ 's of the formula vector.
- (4) The molecular weight.
- (5) The specie type number (01 for atomic or atomic ions, 02 for diatomic species, 03 for linear polyatomic species, 04 for non-linear polyatomic species, and 05 for electrons, H<sup>+</sup>, and He<sup>++</sup>).

\*\*Pertinent format notation is given in Appendix C.

Table 2. Sample library data list

ID No. <i>i</i>	Formula	Formula vector				
		$\alpha_{i,32}$	$\alpha_{i,33}$	$\alpha_{i,34}$	$\alpha_{i,35}$	$\alpha_{i,36}$
1	CN	1	1	0	0	0
2	NO <sup>+</sup>	0	1	1	-1	0
3	CO	1	0	1	0	0
4	NO	0	1	1	0	0
5	NO <sub>2</sub>	0	1	2	0	0
6	N <sub>2</sub> O	0	2	1	0	0
7	CO <sub>2</sub>	1	0	2	0	0
8	N <sub>2</sub>	0	2	0	0	0
9	O <sub>2</sub>	0	0	2	0	0
10	C <sup>+++</sup>	1	0	0	-3	0
11	C <sup>++</sup>	1	0	0	-2	0
12	C <sup>+</sup>	1	0	0	-1	0
13	C <sup>-</sup>	1	0	0	1	0
14	CO <sub>+</sub>	1	0	1	-1	0
15	C <sub>2</sub> N <sub>2</sub>	2	2	0	0	0
16	C <sub>2</sub>	2	0	0	0	0
17	C <sub>3</sub>	3	0	0	0	0
18	N <sup>+++</sup>	0	1	0	-3	0
19	N <sup>++</sup>	0	1	0	-2	0
20	N <sup>+</sup>	0	1	0	-1	0
21	N <sub>2</sub> <sup>+</sup>	0	2	0	-1	0
22	O <sup>+++</sup>	0	0	1	-3	0
23	O <sup>++</sup>	0	0	1	-2	0
24	O <sup>+</sup>	0	0	1	-1	0
25	O <sup>-</sup>	0	0	1	1	0
26	O <sub>2</sub> <sup>-</sup>	0	0	2	1	0
27	O <sub>2</sub> <sup>+</sup>	0	0	2	-1	0
28	O <sub>3</sub>	0	0	3	0	0
29	A <sup>+++</sup>	0	0	0	-3	1
30	A <sup>++</sup>	0	0	0	-2	1
31	A <sup>+</sup>	0	0	0	-1	1
32	C	1	0	0	0	0
33	N	0	1	0	0	0
34	O	0	0	1	0	0
35	e	0	0	0	1	0
36	A	0	0	0	0	1

- (6) The number of energy levels (for atomic and ionic species, this will be the number of energy level cards which make up the packet and cannot exceed 70; for diatomic species, this will be the number of electronic state cards which make up the packet and can not exceed 10; for linear polyatomic or nonlinear polyatomic species, this will be the number of vibrational frequency cards which make up the packet and cannot exceed 12; for the electron H<sup>+</sup> and He<sup>++</sup>, this number is zero).
- (7) The last number on the card is zero for all species except the diatomic forms. For homo-nuclear diatomic molecules, the value is 2.0; for unlike nuclei, the value is 1.0. This is a symmetry number for diatomic species.

The format of this card is 12X, A6, 10I2, 1F10.6, I2, I4, 1F4.1. The column positions of the items listed above are shown in Table 1.

*Heat of Formation Card:* The second card of each packet contains the value of the heat of formation,  $h_o^\circ$ , in kcal/mole for the species associated with the packet. The format of this card is 20X, 1E12.6.

*Energy Level Block Cards of Atomic Species:* The energy level block for atomic or ionic species consists of a tabulation of the electronic energy levels of the ground and excited states together with the degeneracy of these levels. One degeneracy,  $g_i$ , and one energy level,  $E_i$ , in units of cm<sup>-1</sup> are listed on each card. The cards are arranged with the ground state first and with successive excited states following. The format for each card is 19X, 1F6.2, 3X, 1F10.3. The number of these cards in each packet should correspond to the number listed in the next to last position on the species identification card for the species. The total number of these cards in each atomic or ionic species data packet should not exceed 70. Provision has been made for 32 atomic and atomic ion data packets.

*Energy Level Block Cards of Diatomic Species:* The energy level block for diatomic species consists of a tabulation of the electronic energy levels of the ground and excited states together with the degeneracy and the vibrational and rotational constants for these levels. The first 25 columns of each card are not read by the machine so they can be used for identification. Each card represents an electronic level and contains the electronic energy term,  $T_e$ ; the vibrational frequency term,  $\omega_e$ ; the anharmonic term,  $\omega_e x_e$ ; the rotational term,  $B_e$ ; the vibrational-rotational coupling term,  $\alpha_e$ ; and last, the degeneracy of the electronic energy level,  $p$ . The term values are in units of cm<sup>-1</sup>. The cards are arranged with the ground level first followed by the successive excited states. The format of each card is 25X, 1F11.4, 1F9.3, 1F10.5, 1F10.6, 1F9.6, 1F6.2. The number of cards in each packet should correspond to the number listed in the next to last position on the species identification card for the packet. The total number of these cards in each diatomic species data packet should not exceed 10. Provision has been made for 20 diatomic data packets.

*Energy Level Block Cards of Linear Polyatomic Species:* The energy level block for linear polyatomic species contains the information necessary to describe the rotational-vibrational energy states of the ground electronic configuration of the molecule. The first 25 columns

of the first card of this block are not read by the computer and may be used for identification. The format of this card is 25X, 2F3.1, 2F8.5. The first word on the card is the symmetry number for the molecule,  $\sigma$ , followed by the degeneracy of the electronic ground state,  $p$ ; the rotational terms  $B_e$  and  $\alpha_e$  are in units of  $\text{cm}^{-1}$ . The last item on this card which is not read is the number of vibrational degrees of freedom for the molecule for which frequencies,  $\omega$ , ( $\text{cm}^{-1}$ ) are tabulated on successive cards in 1F8.2 format. The above number should also be listed in the next to last position on the species identification card for the packet. The total number of vibrational degrees of freedom should not exceed 12. The total combined number of linear and nonlinear polyatomic atomic data packets can not exceed 75.

**Energy Level Block Cards of Non-Linear Polyatomic Species:** The energy level block for non-linear polyatomic molecules contains the information necessary to describe the rotational-vibrational energy states of the ground electronic configuration. The first 25 columns of the first card of this block are not read and can be used for identification. The format is 25X, 2F3.1, 1F8.5, 4X, 1F8.5, 4X, 1F8.5. The first two words are the symmetry number,  $\sigma$ , and degeneracy of the electronic ground state,  $p$ . The other items on the card are the three moments of inertia ( $I_x, I_y, I_z$ ) for the molecule. These should be in units of  $\text{gm-cm}^2 \times 10^{40}$  as the "E - 40" part of the word is not processed by the computer. The above card is followed by cards which contain the frequencies ( $\omega$ ) of the vibrational degrees of freedom of the molecule in 25X, 1F8.2 format. The number of frequency cards should be indicated in the next to last position on the species identification card. The total number of vibrational degrees of freedom should not exceed 12. The total combined number of linear and nonlinear polyatomic data packets cannot exceed 75.

For the data packets for the electrons,  $\text{H}^+$  ion, and  $\text{He}^{++}$  ion, no energy level block is required.

## B. Case Data

The case-data deck communicates to the computer both control data and thermodynamic condition data. The control data are used to inform the computer as to which "problem case" is to be considered and when "new data" are to be read into the machine. The "new data" may be either a change of initial pressure, initial temperature, initial mole fraction composition, or a complete change of the problem case and initial mixture constituents. Proper use of the control data will allow the user to run any combination of problem cases while varying the

initial thermodynamic state and chemical composition at will.

The arrangement of the case data cards is summarized in Fig. 8. The format for the case data deck will be illustrated by the three following examples.

**Example 1.** Determine the thermodynamic properties for densities ( $\rho/\rho_A$ ) of 0.1 and 0.01 and temperatures between 3000 and 7000°K. The symbol  $\rho_A$  is the density of air at standard conditions.

**Example 2.** Determine the thermodynamic properties and velocities associated with a normal shock in an initial mixture composed of 30%  $\text{CO}_2$  and 70%  $\text{N}_2$ . The initial temperature is 300°K, and the initial pressure is 0.004 atm. The temperatures in the shock-heated gas are 4500°K, 5000°K, 5500°K, and 6000°K. Also, at the same final temperatures as listed above, consider the properties of shocks in a mixture of 50%  $\text{CO}_2$  and 50%  $\text{N}_2$ , with an initial temperature of 300°K and an initial pressure of 0.004 atm.

**Example 3.** Determine the properties of both reflected and standing shocks resulting from traveling incident shocks in an air-like mixture composed of 78.08%  $\text{N}_2$ , 20.95%  $\text{O}_2$ , and 0.97% A. The initial temperature of the gaseous mixture is 300°K. The initial pressures considered are 0.001 and 0.01 atm. The temperatures behind the incident shock range from 4000 to 8000°K.

The form of case data for these examples is given in Tables 3, 4, and 5. The explanation of the format of the cards follows.

### 1. First Card

The first card of the case data deck supplies the total number of chemical species (compounds, ions, atoms, and the electron) and the total number of chemical elements (including the electron as a chemical element) which comprise the high temperature chemical equilibrium mixture. The information on this card is in format 2I2. The first word is the total number of elements in the mixture. The second word is the total number of chemical species.

In Example 1, the possible constituents of the high temperature mixture are  $\text{N}_2$ ,  $\text{N}^{+++}$ ,  $\text{N}^{++}$ ,  $\text{N}^+$ ,  $\text{N}_2^+$ ,  $\text{N}$ , and  $e^-$ . The seven chemical species are composed of chemical elements N and e-. Thus, the first word is 02, and the second is 07.

In Example 2, the chemical species are all those combinations of C, N, O, and  $e^-$  which are listed in Table 2. There are 32 of these. The number of elements plus electrons is 04.

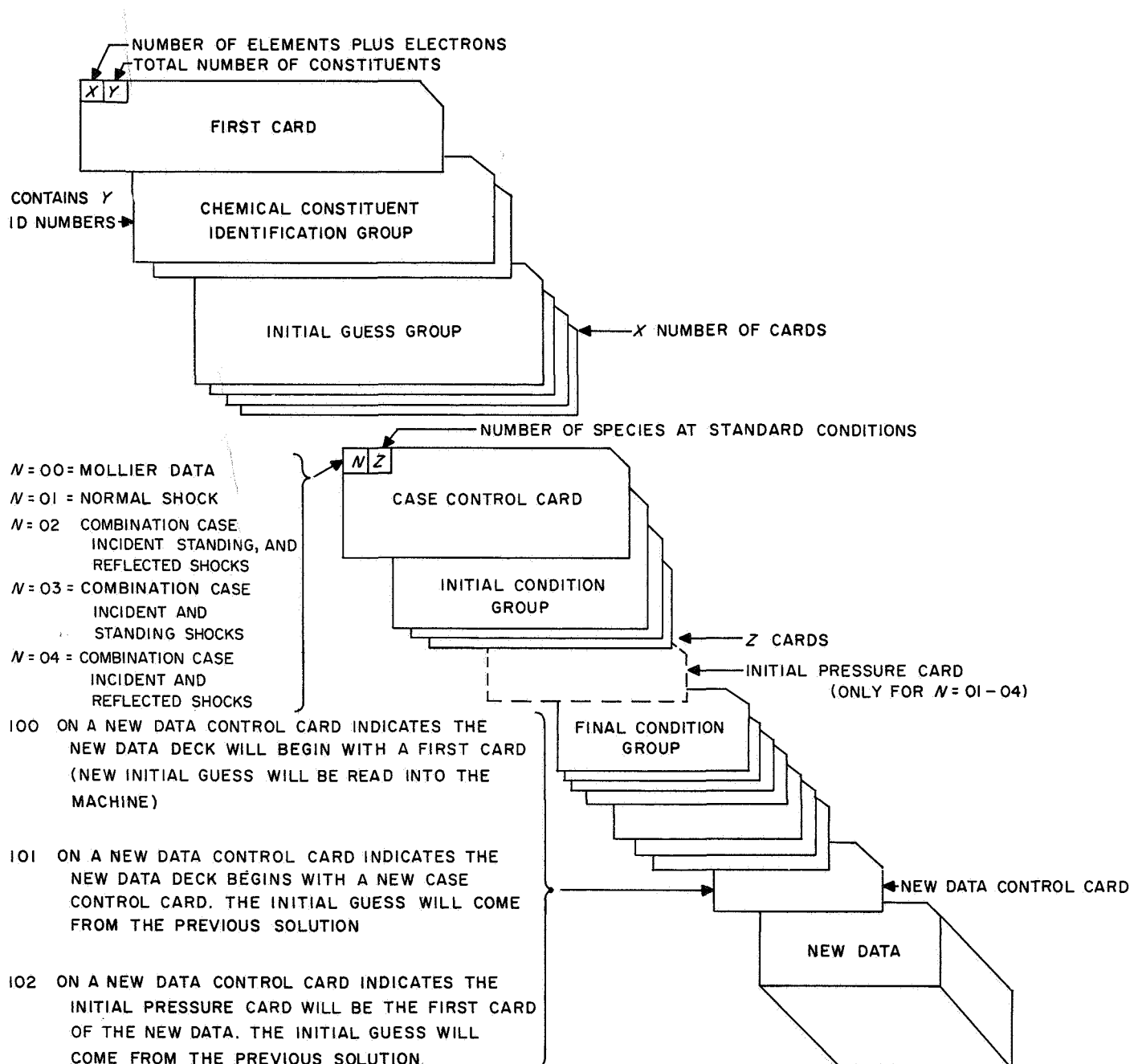


Fig. 8. The case data deck



**Table 3. Example 1: Mollier data case**[illegible]

**Table 4. Example 2: Moving normal shock case (30 % CO<sub>2</sub> – 70 % N<sub>2</sub> and 50 % CO<sub>2</sub> – 50 % N<sub>2</sub>)**

[illegible]

In Example 3, the constituent species are the combinations of N, O, e<sup>-</sup>, and A which are listed in Table 2. There are 24 species which contain these four chemical elements.

## 2. Chemical Constituent Identification Group

This group consists of up to two cards which list the identification numbers of the constituent chemical species of the high temperature mixture. The identification numbers are found in Table 2. Each number is listed with an I4 format. With this format, there are a maximum of 20 identification numbers on each card. The identification numbers should be in order so that the chemical elements occur in the last positions. The number of identification numbers must equal the number in the second word position of the "first card."

For Example 1, the 7 constituents  $N_2$ ,  $N^{+++}$ ,  $N^{++}$ ,  $N^+$ ,  $N_2^+$ ,  $N$ , and  $e^-$  have the identification numbers 8, 18, 19, 20, 21, 33, and 35, respectively. The seven values can be

listed on one card. Two cards are required to list the 32 constituent identification numbers for Example 2 and two are required for identification numbers for Example 3.

### 3. Initial Guess Cards

There is an initial concentration guess card for each chemical element in the mixture. The number of initial guess cards is equal to the number appearing in the first word position of the "first card." The format for an initial guess card is 4X, 1E12.6. The first word position contains the chemical symbol for the element. The second word position contains the approximate concentration (mole/l) of the element at the first equilibrium state for which a solution is sought. The order of the initial guess cards should be the same as the order of the elements in Table 2.

As the convergence of numerical procedure for the chemical equilibrium computation is dependent upon the knowledge of a set of concentrations which nearly satisfy

**Table 5. Example 3: Combination normal shock case (air mixture)**[illegible]

the mass action and mass balance equations, the initial guess values are very important. If the initial guess values are in error by more than an order of magnitude, the convergence of the chemical equilibrium calculation is doubtful. The procedure for selecting the initial guesses is somewhat of an art which develops with the use of the program.

The program automatically uses the concentrations from the previous equilibrium calculation as the initial guess values for the next chemical equilibrium computation. Thus, the set of final and initial condition states should be arranged in an orderly progression to ensure convergence of the numerical procedure.

In the case of Example 1, the concentrations of  $6 \times 10^{-8}$  mole/l for N and  $1 \times 10^{-14}$  mole/l for the electrons are the approximate concentrations at a temperature of 3000°K and a density of 0.129 g/l. For Example 2, the concentrations (in mole/l) of

C	$1.76 \times 10^{-5}$
N	$3.85 \times 10^{-4}$
O	$6.44 \times 10^{-4}$
e <sup>-</sup>	$4.74 \times 10^{-7}$

are associated with the temperature of 6000°K and the density of about 0.05 g/l. For Example 3, the concentrations of

N	$1.8 \times 10^{-6}$
O	$1.6 \times 10^{-4}$
e <sup>-</sup>	$3.0 \times 10^{-9}$
A	$4.2 \times 10^{-6}$

are for a temperature of 4000°K and a density of about 0.01 g/l.

#### 4. Case Control Card

The case control card contains two numbers in I2 format. The first number identifies the case of the problem to be considered. If this number is 00, the indication is that Mollier data is to be generated. If this number is 01, moving normal shock problems are to be solved. The number 03 indicates a combination case is desired consisting of an incident shock solution and a standing shock solution. For the number 04, one obtains a combination case in which incident shock and reflected shock properties are the result. With the number 02, a combination case consisting of incident, standing, and reflected shock conditions results. No other values are valid for the first word position on the case control card.

The second number on the card identifies the number of compounds present in the mixture at standard conditions. This number must equal the number of cards in the initial condition group.

#### 5. Initial Condition Group

The format for the initial condition control card is I2, A6, 1F7.5, 1E12.6. There is one card for each of the initial constituents of the mixture. The number of cards in this group must be equal to the number in the second word position of the control card.

The first word is the species identification number taken from Table 2. The second word is the chemical symbol for the species. The third word is the mole-fraction of the species in the mixture at standard conditions. For the Mollier data case, this is all the information required on an initial condition group card, and any information in the fourth position is read into the computer but is not utilized in the computations.

For the moving and combination normal shock cases, the value in the fourth word position is the initial temperature,  $T_1$ .

The order of the cards within the initial condition group is not important; however, the best procedure would be to arrange the identification numbers which appear on them in order. The data on these cards are not only used to calculate the thermodynamic properties of the gas at standard and initial conditions, but also are used with the library data to calculate the values of  $B_i$  which appear in the mass action equations.

For Example 1, the 100%  $N_2$ -Mollier data case, there is only one card in the initial condition group. The first

word is 08 (which agrees with the specie identification number for  $N_2$  in Table 2). The position and format of the chemical formula and mole-fraction are shown in Table 3. Normally, the fourth position would be vacant for a Mollier data case. The value listed for the temperature will not effect the operation of the program as it is not used in Mollier data computations.

For Example 2, the first card of the initial condition group corresponds to  $CO_2$  and contains the specie ID number of 07, the mole-fractional concentration at standard conditions of 0.3, and the initial temperature of 300°K. The second card corresponds to the compound  $N_2$  and contains the ID number of 08, the mole-fraction of 0.7, and the initial temperature of 300°K.

For Example 3, the cards list:

ID No.	Chemical symbol	Mole-fraction @ STP	Initial temperature, °K
08	$N_2$	0.7808	300
09	$O_2$	0.2095	300
36	A	0.0097	300

#### 6. Initial Pressure Card

The initial pressure card is not included when the Mollier data case is to be run. The card contains the pressure of the gas in atmospheres before the passage of the shock. The pressure value is in a format E12.6.

In Example 1, the initial pressure card is not necessary.

In Example 2, the initial pressure is 0.04 atm. Farther down in the new case data, the initial pressure of 0.04 atm is repeated for the 50%- $CO_2$  and 50%- $N_2$  mixture.

In Example 3, the first initial pressure is 0.001 atm. In the new data, the initial pressure is changed to 0.01 atm.

#### 7. Final Condition Group

The final condition group is composed of an unlimited number of cards. There is one card for each thermodynamic or shock condition. The thermodynamic condition associated with the first card of this group should be compatible with the concentrations listed on the initial guess cards.

For the Mollier data cases, each card specifies the temperature and density of the thermodynamic state for which the chemical equilibrium composition and properties are to be calculated. The format for a Mollier data case, final condition group card is 3X, 2E12.6. The first three spaces are blank, followed by the first word which is the temperature in °K. The second word is the ratio of density at the solution point to density of air at standard conditions. Experience has shown that the most efficient way of calculating data over a range of temperatures and densities is to calculate a series of temperatures while holding the density constant, then change the density by not more than an order of magnitude and calculate through the series of temperatures in reverse order, and so on.

For the shock cases, the format of the final condition cards is 3X, 1E12.6. The first three spaces are left blank, followed by the temperature in °K. The first card of a final condition group should correspond to the concentrations listed on the initial guess cards. The temperatures listed on the final condition group cards are not limited in any way.

### 8. New Data Control Card

A new data control card is used for the following options:

- (1) To change initial pressure in normal shock and combined shock cases.
- (2) To change initial mixture mole-fractions and/or to change initial temperature and/or change cases.
- (3) To completely change the problem by reading in a completely new deck of case data beginning with the "first card."

If the computations are complete when the last card of the final condition group has been processed, then the new data card is not used. However, if as shown in Table 5, an option such as (1) is considered, with the initial pressure changed from  $10^{-3}$  to  $10^{-2}$  atm, then a new data card with 102 must be used to preface the new data. In this option, the new data consists of a new pressure card followed by a new final condition group. Note that the last temperature and first temperature of the new data are the same to ensure convergence at the new pressure.

If an option such as (2) is desired, the new data control card should contain a 101. This will cause the machine to read in a new set of data which should consist of new case control card, initial condition group, initial pressure card, and final condition group. An example illustrating the use of the 101 new data control card is given in Table 4, where the card is used to change the mole-fraction composition of the  $\text{CO}_2\text{-N}_2$  mixture for a moving normal shock case. Because the program uses the concentrations from the previous calculation as the initial guesses for the new calculations, the change in initial state and concentration should not be large.

The option (3) is achieved by using the 100 new data control card. After this card, a completely new set of case data, beginning with "first card," should follow. A 100 new data control card used between each of the example data of Tables 3, 4, and 5 would allow the three problems to be solved in succession.

## VII. Program Output

The output for the first solution of each of the three example cases is shown in Tables 6, 7, and 8.

Table 6. Chemical equilibrium composition and properties of a complex mixture

INITIAL CONDITIONS				MOLE FRACTION OF N2				AT S.T.P.---1.00000			
FINAL CONDITIONS											
	CONCENTRATION				MOLE FRACTION						
	PARTICLES/CC		MOLE/LITER								
N2	0.2779E 19	0.4613E-02	1.0000E 00								
N+++	0.6023E-14	1.0000E-35	0.2168E-32								
N++	0.6023E-14	1.0000E-35	0.2168E-32								
N+	0.4229E 04	0.7020E-17	0.1522E-14								
N2+	0.5718E 07	0.9492E-14	0.2058E-11								
N	0.3611E 14	0.5996E-07	0.1300E-04								
E-	0.5722E 07	0.9500E-14	0.2059E-11								
TEMPERATURE (DEG K)-----				0.3000E 04	P/PO-----		0.1136E 01	ENTROPY (SMO/R)-----		0.3196E 02	
ENTHALPY (HMO/RTO)-----				0.4463E 02	RHO/RH00----		0.1000E-00	RH00-----		0.1292E 01	
MIXTURE COLD MOL WT-----				0.2801E 02	Z-----		0.1000E 01				

Table 7. Real gas normal shock-wave properties

```

INITIAL CONDITIONS
TEMPERATURE (DEG K)----- 0.3000E 03  RHO1/RHO0---0.4126E-02  RHO0 (GM/LITER)-----0.1292E 01
ENTHALPY (HMO/RT0)-----0.4799E 02  P1/P0-----0.4000E-02  MIXTURE COLD MOL WT-----0.3281E 02
                                MOLE FRACTION OF CO2  AT S.T.P.---0.30000
                                MOLE FRACTION OF N2   AT S.T.P.---0.70000

FINAL CONDITIONS
                                CONCENTRATION      MOLE FRACTION
                                PARTICLES/CC  MOLE/LITER
CN      0.1363E 16  0.2264E-05  0.7475E-03
NO+     0.2582E 15  0.4286E-06  0.1416E-03
CO      0.3807E 18  0.6321E-03  0.2087E-00
NO      0.1149E 17  0.1908E-04  0.6300E-02
NO2     0.6960E 11  0.1155E-09  0.3816E-07
N2O     0.6417E 12  0.1065E-08  0.3518E-06
CO2     0.7981E 14  0.1325E-06  0.4376E-04
N2      0.7937E 18  0.1318E-02  0.4351E-00
O2      0.3218E 15  0.5343E-06  0.1764E-03
C+++    0.6023E-14  1.0000E-35  0.3302E-32
C++     0.7776E-01  0.1291E-21  0.4263E-19
C+      0.1807E 14  0.3000E-07  0.9908E-05
C-      0.5384E 10  0.8938E-11  0.2952E-08
CO+     0.9704E 13  0.1611E-07  0.5320E-05
C2N2    0.4430E 10  0.7355E-11  0.2429E-08
C2      0.4501E 13  0.7473E-08  0.2468E-05
C3      0.1058E 10  0.1756E-11  0.5800E-09
N+++    0.6023E-14  1.0000E-35  0.3302E-32
N++     0.1648E-05  0.2736E-26  0.9035E-24
N+      0.2420E 13  0.4017E-08  0.1326E-05
N2+     0.1344E 13  0.2231E-08  0.7366E-06
O+++    0.6023E-14  1.0000E-35  0.3302E-32
O++     0.2608E-09  0.4330E-30  0.1430E-27
O+      0.4931E 13  0.8187E-08  0.2704E-05
O-      0.5673E 12  0.9419E-09  0.3110E-06
O2-     0.2388E 09  0.3964E-12  0.1309E-09
O2+     0.1584E 12  0.2630E-09  0.8685E-07
O3      0.8013E 08  0.1330E-12  0.4393E-10
C       0.1062E 17  0.1763E-04  0.5821E-02
N       0.2327E 18  0.3864E-03  0.1276E-00
O       0.3924E 18  0.6514E-03  0.2151E-00
E-      0.2942E 15  0.4885E-06  0.1613E-03
SHOCK VELOCITY (FPS)----- 0.1813E 05  RHO2/RHO0-- 0.5520E-01  VEL BEHIND SHOCK (FPS)--0.1678E 05
TEMPERATURE (DEG K)----- 0.6000E 04  RHO2/RHO1-- 0.1338E 02  ENTROPY (SMO/R)-----0.4649E 02
ENTHALPY (HMO/RT0)----- 0.1714E 03  P2/P0----- 0.1491E 01  Z-----0.1393E 01
STAG ENTHALPY (HMO/RT0)- 0.1726E 03  P2/P1----- 0.3727E 03

```

For a Mollier data case, the format is shown in Table 6. Listed are the initial conditions, where the composition of the mixture at standard temperature and pressure is given. The chemical concentration of each specie in the heated mixture is listed under "Final Conditions." The first column is the concentration in particles per cubic centimeter, and the second column is in mole per liter. The third column lists the mole fraction concentration.

The tabulation of specie concentration is followed by a display of the thermodynamic properties in dimensionless form. These values are; the temperature in °K; the dimensionless enthalpy per unit mass ( $HMO/RT_0$ ); the cold mixture molecular weight; the pressure in atmospheres ( $P_0 = 1$  atm); the dimensionless density;  $Z$  as

defined in Eq. (19); the entropy in dimensionless form ( $SMO/R$ ); and  $RHO0 = \rho_0$ , which is the density of air in g/l at standard conditions. The  $R$  is the universal gas constant 1.98726 cal/mole-°K;  $M_0$  is the cold molecular weight, and  $T_0$  is 273.15°K.

The forms of the output data for the moving normal shock case and the combination normal shock case are shown in Tables 7 and 8. The two are similar except for the additional listings in Table 8 of the properties of the standing and reflected normal shocks. For the combination cases 03 and 04, the output consists of listings of the moving and standing shocks, and moving and reflected shocks, respectively. The description of the concentration and thermodynamic property display given

Table 8. Combination normal shock solutions

```

(6) REAL GAS NORMAL SHOCK WAVE PROPERTIES
INITIAL CONDITIONS
TEMPERATURE (DEG K)----- 0.3000E 03  RHO1/RHO0---0.9104E-03  RHO0 (GM/LITER)-----0.1292E 01
ENTHALPY (HMO/RT0)----- 0.3835E 01  P1/P0-----1.0000E-03  MIXTURE COLD MOL WT-----0.2896E 02
                                MOLE FRACTION OF  N2  AT S.T.P.---0.78080
                                MOLE FRACTION OF  O2  AT S.T.P.---0.20950
                                MOLE FRACTION OF   A  AT S.T.P.---0.00970

FINAL CONDITIONS
CONCENTRATION      MOLE FRACTION
PARTICLES/CC  MOLE/LITER
NO+  0.1784E 13  0.2962E-08  0.5675E-05
NO   0.6641E 16  0.1102E-04  0.2112E-01
NO2  0.1607E 12  0.2669E-09  0.5113E-06
N2O  0.1345E 12  0.2233E-09  0.4277E-06
N2   0.2025E 18  0.3361E-03  0.6440E 00
O2   0.2431E 16  0.4037E-05  0.7734E-02
V+++ 0.6023E-14  1.0000E-35  0.1916E-31
V++  0.6023E-14  1.0000E-35  0.1916E-31
N+   0.7968E 06  0.1323E-14  0.2534E-11
N2+  0.7639E 07  0.1268E-13  0.2430E-10
O+++ 0.6023E-14  1.0000E-35  0.1916E-31
O++  0.6023E-14  1.0000E-35  0.1916E-31
O+   0.2184E 09  0.3626E-12  0.6947E-09
O-   0.6597E 10  0.1095E-10  0.2099E-07
O2-  0.3635E 08  0.6035E-13  0.1156E-09
O2+  0.9947E 09  0.1651E-11  0.3164E-08
J3   0.3434E 09  0.5702E-12  0.1092E-08
AR+++0.6023E-14  1.0000E-35  0.1916E-31
AR++ 0.6023E-14  1.0000E-35  0.1916E-31
AR+  0.1289E 06  0.2141E-15  0.4101E-12
N    0.1075E 16  0.1785E-05  0.3421E-02
O    0.9921E 17  0.1647E-03  0.3156E-00
E-   0.1779E 13  0.2953E-08  0.5657E-05
AR   0.2563E 16  0.4255E-05  0.8153E-02
SHOCK VELOCITY (FPS)----- 0.1319E 05  RHO2/RHO0-- 0.9832E-02  VEL BEHIND SHOCK (FPS)---0.1197E 05
TEMPERATURE (DEG K)----- 0.4000E 04  RHO2/RHO1-- 0.1080E 02  ENTROPY (SMO/R)-----0.3997E 02
ENTHALPY (HMO/RT0)----- 0.1060E 03  P2/P0----- 0.1713E-00  Z-----0.1190E 01
STAG ENTHALPY (HMO/RT0)- 0.1069E 03  P2/P1----- 0.1713E 03

```

above for the Mollier data case is applicable to the shock solution cases. Under the heading "Initial Conditions" are the temperature, dimensionless enthalpy, density and pressure, mixture molecular weight, and composition of the initial unshocked gas. The heading "Final Conditions" is equivalent to subscripting the properties with 2 for the moving shock, 3 for the standing shock, and 4 for the reflected shock.

The quantities listed below the concentrations of the moving normal shock are: velocity,  $U_s$ , or free-flight velocity,  $U_f$ , in ft/sec; temperature,  $T_2$ ; dimensionless enthalpy,  $H_2 M_o / RT_o$ ; dimensionless stagnation enthalpy on a free-flight body; the density ratios  $\rho_2 / \rho_o$  and  $\rho_2 / \rho_1$ ; pressure ratios  $P_2 / P_o$  and  $P_2 / P_1$ ; the velocity  $U_g$  associated with the moving normal shock; the entropy,  $S_2 M_o / R$ ; and  $Z_2 = M_o / M_2$ .

The final conditions listed under the heading "Real Gas Standing Normal Shock Wave Properties" are: the con-

centrations behind the standing wave; the moving shock velocity; temperature,  $T_3$ ; dimensionless enthalpy,  $H_3 M_o / RT_o$ ; "cold" molecular weight,  $M_o$ ; pressure,  $P_3$ , in atmospheres; pressure ratio,  $P_3 / P_2$ ; density ratios  $\rho_3 / \rho_o$  and  $\rho_3 / \rho_2$ ; standard density,  $\rho_o$ ; velocity behind the standing shock,  $U_3$ , in ft/sec; dimensionless stagnation enthalpy on the model; dimensionless entropy,  $S_3 M_o / R$ ; and  $Z_3$ .

The final conditions listed under the heading "Real Gas Reflected Normal Shock Wave Properties" are: the concentrations behind the reflected wave; the shock velocity,  $U_s$ ; temperature,  $T_4$ ; dimensionless enthalpy,  $H_4 M_o / RT_o$ ; "cold" molecular weight,  $M_o$ ; the pressure,  $P_4$ , in atmospheres; the pressure ratio,  $P_4 / P_2$ ; the density ratios,  $\rho_4 / \rho_o$  and  $\rho_4 / \rho_2$ ; the reflected shock velocity in ft/sec; dimensionless entropy,  $S_4 M_o / R$ ;  $Z_4$  and  $\rho_o$ .

The program prints one shock solution per page. For the combination case shown in Table 8, the solutions for

Table 8 (contd)

{b} REAL GAS STANDING NORMAL SHOCK WAVE PROPERTIES									
INITIAL CONDITIONS		TEMP DEG K--0.3000E 03				P/P0--1.0000E-03			
		MOLE FRACTION OF N2				AT S.T.P.---0.78080			
		MOLE FRACTION OF O2				AT S.T.P.---0.20950			
		MOLE FRACTION OF A				AT S.T.P.---0.00970			
FINAL CONDITIONS									
		CONCENTRATION		MOLE FRACTION					
		PARTICLES/CC	MOLE/LITER						
N0+	0.3763E 15	0.6248E-06	0.2038E-03						
N0	0.1787E 17	0.2967E-04	0.9678E-02						
N02	0.1494E 12	0.2480E-09	0.8090E-07						
N20	0.1106E 13	0.1836E-08	0.5988E-06						
N2	0.9560E 18	0.1587E-02	0.5177E 00						
O2	0.6224E 15	0.1033E-05	0.3370E-03						
N+++	0.6023E-14	1.0000E-35	0.3262E-32						
N++	0.6023E-14	1.0000E-35	0.3262E-32						
N+	0.3084E 13	0.5120E-08	0.1670E-05						
N2+	0.1744E 13	0.2895E-08	0.9442E-06						
O+++	0.6023E-14	1.0000E-35	0.3262E-32						
O++	0.6023E-14	1.0000E-35	0.3262E-32						
O+	0.7421E 13	0.1232E-07	0.4019E-05						
O-	0.1044E 13	0.1733E-08	0.5652E-06						
O2-	0.5906E 09	0.9806E-12	0.3199E-09						
O2+	0.3041E 12	0.5049E-09	0.1647E-06						
O3	0.2230E 09	0.3702E-12	0.1208E-09						
AR+++	0.6023E-14	1.0000E-35	0.3262E-32						
AR++	0.6023E-14	1.0000E-35	0.3262E-32						
AR+	0.3573E 11	0.5932E-10	0.1935E-07						
N	0.2831E 18	0.4700E-03	0.1533E-00						
O	0.5744E 18	0.9536E-03	0.3111E-00						
E-	0.3879E 15	0.6440E-06	0.2101E-03						
AR	0.1375E 17	0.2282E-04	0.7446E-02						
SHOCK VELOCITY (FPS)----		0.1319E 05	P3/P0-----0.1526E 01	VEL BEHIND ST SHOCK-----0.2231E 04					
TEMPERATURE (DEG K)-----		0.6066E 04	P3/P2-----0.8907E 01	STAG ENTHALPY (HMO/RT0)-0.1900E 03					
ENTHALPY (HMO/RT0)-----		0.1870E 03	RHO3/RHO0---0.5275E-01	ENTROPY (SMO/R)-----0.4157E 02					
MO-MIXTURE COLD MOL WT--		0.2896E 02	RHO3/RHO2---0.5364E 01	Z-----0.1303E 01					
			RHO0 (GM/L)-0.1292E 01						

the moving, standing and reflected shocks were printed on three pages. To prevent the shock solution of a complicated mixture, involving many species, from listing the output on more than one page, the program will list the specie concentrations in two columns when the number of species is larger than 35. An example of this is shown in the moving shock solution of Table 9. This solution is for a mixture of 94% air and 6% water vapor at an initial pressure of 1 mm Hg (0.00131579 atm) and a final temperature of 10,000°K. Since the number of species is larger than 35, the concentrations are listed in two sets.

In addition to the above "scheduled" output, the program has been designed to halt all computations when the number of iteration loops through the equilibrium block exceeds 50, or when the number of iteration loops through the moving normal shock equations exceeds 20. After 50 equilibrium block iterations, the program prints "equilibrium state not reached in 50 iterations" followed by a tabulation of the concentrations for the last two

iterations. After 20 normal shock iterations, the machine prints "no Hugoniot solution in 20 iterations" and halts.

### VIII. Discussion

The accuracy of the program has been checked in a number of ways. First, the accuracy of the matrix inversion portion of the program was checked and found to give correct solutions to arbitrary systems of 10 simultaneous linear equations accurate to seven figures. Second, through comparisons with hand computations the accuracy of the equilibrium solutions have been proven. The basic convergence criteria has been shown to be reliable by the comparison of successive interactions. Perhaps one should realize that the concentrations of electrons and ions can be in error at low temperature, below 2000°K or 3000°K, due to the fact that when such a small concentration is present, a change in their concentration results in a negligible change in gross properties. In every case encountered in four years of



Table 8 (contd)

(c) REAL GAS REFLECTED NORMAL SHOCK WAVE PROPERTIES										
INITIAL CONDITIONS		TEMP DEG K--0.3000E 03				P/P0--1.0000E-03				
		MOLE FRACTION OF N2				AT S.T.P.---0.78080				
		MOLE FRACTION OF O2				AT S.T.P.---0.20950				
		MOLE FRACTION OF A				AT S.T.P.---0.00970				
FINAL CONDITIONS										
	CONCENTRATION			MOLE FRACTION						
	PARTICLES/CC		MOLE/LITER							
NO+	0.7230E 15	0.1200E-05	0.3002E-03							
NO	0.1982E 17	0.3290E-04	0.8227E-02							
NO2	0.1473E 12	0.2445E-09	0.6115E-07							
N2O	0.1391E 13	0.2309E-08	0.5774E-06							
N2	0.1095E 19	0.1818E-02	0.4545E-00							
O2	0.5366E 15	0.8908E-06	0.2228E-03							
N+++	0.6023E-14	1.0000E-35	0.2501E-32							
N++	0.6023E-14	1.0000E-35	0.2501E-32							
N+	0.2012E 14	0.3340E-07	0.8354E-05							
N2+	0.7846E 13	0.1303E-07	0.3258E-05							
O+++	0.6023E-14	1.0000E-35	0.2501E-32							
O++	0.6023E-14	1.0000E-35	0.2501E-32							
O+	0.2772E 14	0.4601E-07	0.1151E-04							
O-	0.1972E 13	0.3274E-08	0.8189E-06							
O2-	0.8511E 09	0.1413E-11	0.3534E-09							
O2+	0.6363E 12	0.1056E-08	0.2642E-06							
O3	0.2182E 09	0.3623E-12	0.9061E-10							
AR+++	0.6023E-14	1.0000E-35	0.2501E-32							
AR++	0.6023E-14	1.0000E-35	0.2501E-32							
AR+	0.1748E 12	0.2903E-09	0.7259E-07							
N	0.5545E 18	0.9206E-03	0.2302E-00							
O	0.7202E 18	0.1196E-02	0.2990E-00							
E-	0.7775E 15	0.1291E-05	0.3228E-03							
AR	0.1717E 17	0.2851E-04	0.7130E-02							
SHOCK VELOCITY (FPS)----				0.1319E 05	P4/P0-----	0.2128E 01	REF SHOCK VEL (FPS)-----			0.2100E 04
TEMPERATURE (DEG K)-----				0.6487E 04	P4/P2-----	0.1242E 02	ENTROPY (SMO/R)-----			0.4255E 02
ENTHALPY (HMO/RT0)-----				0.2196E 03	RHO4/RHO0---	0.6589E-01	Z-----			0.1360E 01
MO-MIXTURE COLD MOL WT--				0.2896E 02	RHO4/RHO2---	0.6700E 01	RHO0-----			0.1292E 01

experience with the JPL Thermochemistry and Normal Shock Program where errors were encountered, they are in species of negligible concentration. Mostly, this is because, in any machine computation, we reach a point where the numbers computed are either so small or large that they exceed the capacity of the computer. In some cases as in this program, it is not advisable to call these small quantities zero. To alleviate this difficulty in our equilibrium computations, we have considered all concentrations less than  $10^{-35}$  moles/liter, which corresponds to approximately  $6 \times 10^{-15}$  particles/cm<sup>3</sup> to be equal to  $10^{-35}$  moles/liter. Concentrations of this magnitude are obviously negligible, but are included in the output.

Finally the accuracy of the shock solutions have been proven by spot comparison with a variety of shock parameter tabulations including Refs. 12, 30, and 45. Also, the results of the present program have been compared with those of its predecessor. In all cases where comparable thermochemistry was used, the shock parameter and concentrations were in agreement to 3 or 4 figures. As some

of these computations were performed using only 4 significant figures, this agreement would seem remarkable. However, comparison of the result of the program computed, using 4 significant figures with those using 8 figures, reveals agreement for concentrations to 3 to 4 figures; thus, the round off errors are not severe when using only 4 significant figures in the computations.

In computing properties of a variety of mixtures over a wide range of conditions, not a single case has been encountered in which the program has failed to achieve a solution. All failures of the program to converge to a solution or to yield obviously erroneous results have been associated with the failure of the user to exercise proper care in treating the input data.

A word of caution about using the program in extreme ranges of temperature and pressure is called for at this point. As has been pointed out intentionally in the comments on the literature given in the Appendix A, caution must be used at a low temperature when the pressure

Table 9. Moving normal shock solution for an air-water vapor mixture

REAL GAS NORMAL SHOCK WAVE PROPERTIES									
INITIAL CONDITIONS									
TEMPERATURE (DEG K)-----	0.3000E 03	RHO1/RHO0----	0.1171E-02	RHO0 (GM/LITER)-----	0.1292E 01				
ENTHALPY (HMO/RT0)-----	0.2443E 01	P1/P0-----	0.1316E-02	MIXTURE COLD MOL WT-----	0.2831E 02				
	MOLE FRACTION OF	N2	AT S.T.P.----	0.73400					
	MOLE FRACTION OF	O2	AT S.T.P.----	0.19700					
	MOLE FRACTION OF	H2O	AT S.T.P.----	0.06000					
	MOLE FRACTION OF	AR	AT S.T.P.----	0.00900					
FINAL CONDITIONS									
	CONCENTRATION		MOLE FRACTION		CONCENTRATION		MOLE FRACTION		
	PARTICLES/CC	MOLE/LITER			PARTICLES/CC	MOLE/LITER			
N2	0.3273E 16	0.5434E-05	0.3267E-02	O2	0.2417E 13	0.4013E-08	0.2413E-05		
H2	0.3756E 12	0.6235E-09	0.3749E-06	NO	0.1253E 15	0.2079E-06	0.1250E-03		
NH	0.8646E 13	0.1435E-07	0.8629E-05	OH	0.2753E 13	0.4570E-08	0.2748E-05		
N2+	0.3881E 14	0.6443E-07	0.3873E-04	NO-	0.8265E 09	0.1372E-11	0.8249E-09		
N0+	0.1113E 15	0.1848E-06	0.1111E-03	O2+	0.3786E 12	0.6285E-09	0.3778E-06		
O2-	0.3358E 08	0.5575E-13	0.3352E-10	OH+	0.1493E 12	0.2479E-09	0.1490E-06		
OH-	0.2338E 08	0.3881E-13	0.2333E-10	H2+	0.1394E 11	0.2315E-10	0.1391E-07		
N2O	0.7034E 09	0.1168E-11	0.7021E-09	NO2	0.5101E 08	0.8469E-13	0.5092E-10		
NH3	0.3053E 02	0.5068E-19	0.3047E-16	O3	0.1793E 06	0.2976E-15	0.1789E-12		
H2O	0.8986E 07	0.1492E-13	0.8968E-11	N-	0.4738E 13	0.7867E-08	0.4729E-05		
N+	0.1568E 17	0.2603E-04	0.1565E-01	N++	0.2847E 07	0.4727E-14	0.2842E-11		
N+++	0.1435E-12	0.2382E-33	0.1432E-30	O-	0.3045E 13	0.5056E-08	0.3039E-05		
O+	0.3197E 16	0.5308E-05	0.3191E-02	O++	0.3320E 04	0.5513E-17	0.3314E-14		
O+++	0.6023E-14	1.0000E-35	0.6012E-32	H-	0.2798E 12	0.4645E-09	0.2793E-06		
H+	0.9660E 15	0.1604E-05	0.9641E-03	AR+	0.6625E 14	0.1100E-06	0.6612E-04		
AR++	0.3084E 06	0.5120E-15	0.3078E-12	AR+++	0.9068E-10	0.1505E-30	0.9051E-28		
N	0.6828E 18	0.1134E-02	0.6815E 00	O	0.2147E 18	0.3564E-03	0.2143E-00		
H	0.5668E 17	0.9410E-04	0.5657E-01	AR	0.4258E 16	0.7069E-05	0.4250E-02		
E-	0.2005E 17	0.3329E-04	0.2001E-01						
SHOCK VELOCITY (FPS)----	0.3246E 05	RHO2/RHO0--	0.1747E-01	VEL BEHIND SHOCK (FPS)---	0.3028E 05				
TEMPERATURE (DEG K)-----	0.1000E 05	RHO2/RHO1--	0.1492E 02	ENTROPY (SMO/R)-----	0.5841E 02				
ENTHALPY (HMO/RT0)-----	0.6048E 03	P2/P0-----	0.1365E 01	Z-----	0.2085E 01				
STAG ENTHALPY (HMO/RT0)-	0.6076E 03	P2/P1-----	0.1037E 04						

is high, i.e., over 100 atm and at high temperature. There is a limited temperature-pressure range in which the ideal or perfect real-gas assumptions used in this program are valid. At low temperatures, the boundary is determined by the "van der Waals" forces which become important at high densities. At high pressure and tem-

perature, the problem of computing gas properties depends upon the shift in ionization potential and energy levels due to the coulombic or Debye type of micro-field associated with a plasma. This problem of defining the range of application of the results of this program are discussed further in Ref. 38.

## Appendix A

### Review of Some of the Literature on Computation of Equilibrium Properties of Mixtures

(1966) Koopman (Ref. 2, 1966): An iterative procedure has been developed and programmed in MAD computer language for the calculation of the states attained behind incident and reflected shocks in fully dissociated and partially ionized gases. The calculations include the effects of truncation of the partition function sum and lowering of the ionization potential, effects which the author states "have only a modest effect on the determination of the final electron density under conditions encountered in shock tubes." The final equilibrium mixtures can include only atomic and ionic forms of neon, argon, hydrogen, oxygen, sulfur, carbon, and nitrogen.

(1966) Zelenik and Gordon (Ref. 3, 1966): An iterative method for calculating the equilibrium composition and thermodynamic properties of a mixture of atomic and ionic species is presented in which both the Debye-Hückel and the ionization potential lowering effects are considered.

(1965) Drellishak, Aeschliman, and Cambel (Ref. 4, 1965) (Ref. 14, 1963); Drellishak (Ref. 13, 1964): The physical problems in computing partition function and chemical composition of a gas mixture consisting of atoms, molecules, and ions, including the effect of the Debye cutoff criterion and lowering of the ionization potential are discussed. For the molecular species, different degrees of sophistication have been examined in computing ground-state rotational-vibrational contribution to the partition functions. Results in the form of chemical composition and internal partition functions of argon, nitrogen, and oxygen plasma are presented in graphical form up to 35,000°K for pressures between 0.01 and 5 atm.

(1965) Lewis and Burges (Ref. 5, 1965); Lewis and Neel (Ref. 10, 1964) (Ref. 12, 1964): A compilation of normal shock-wave parameters is presented, which is based upon the early imperfect real-gas thermodynamic properties of air computed by Hilsenrath and Klein. The altitude-velocity charts extend up to velocities of 50,000 fps and up to altitudes of 400,000 ft. The shock tables present incident, standing, and reflected shock param-

eters in air for incident shock Mach numbers up to 30. In a comparison of imperfect and perfect specific heats and speeds of sound, the differences were found to be "modest" for temperatures below 15,000°K and pressures below 100 atm.

(1965) Hilsenrath and Klein (Ref. 6, 1965): An extensive tabulation is presented of thermodynamic properties and chemical composition (28 species) of air from 1500°K to 15,000°K in 100°K increments and at close density intervals. The effects of second virial corrections at high pressures and Debye-Hückel corrections at high electron concentrations were included. The appendix of this report contains an excellent review by A. R. Hochstim of early calculations of thermodynamic properties of air.

(1964) Hester and Sewell (Ref. 11, 1964); Sewell (Ref. 19, 1962); (Ref. 31, 1961); Fenter and Gibbons (Ref. 33, 1961): These studies are tabulations of chemical compositions and equilibrium properties of nitrogen, argon, air, and helium for a wide range of densities and temperatures up to 100,000°K. These data are based upon ideal or perfect real-gas analysis with the exception of helium, in which the influence of truncating the partition function sum at different principal quantum numbers was examined. In this case allowing  $n_{\min}$  to vary and holding  $n_{\max} = 4$  fixed, this results in only 2% deviation in the thermodynamic properties.

(1964) McChesney and Jones (Ref. 7, 1964): The equilibrium shock parameters for xenon are calculated for initial pressures between 0.5 and 750 mm Hg and shocked gas temperatures up to 20,000°K. Also, the problem of determining the electronic partition function of the atomic and ionic species is discussed. By means of computed examples, the role of both the missing energy levels and the Debye-Hückel effects is examined. The authors' conclusion is that for the normal range of shock tube experiments, the differences arising from the Debye-Hückel theory and the missing energy levels do not seriously modify the properties of a shock ionized plasma. Because of its low ionization potentials, one would expect the effects of partition function truncation procedures, ionization potential lowering, and missing energy levels to be more pronounced in xenon than in the other noble gases.

(1964) Pinkus (Ref. 9, 1964): This investigation presented in tabular form for the properties of air at relatively low temperatures (subcritical to 1500°K) but at high pressures (up to 1000 atm) where real-gas phenomena are dominated by "van der Waals" type forces. The computations are based upon an equation of state proposed by Hirschfelder. Also, the plenum conditions for both isothermal and isentropic expansion conditions are given.

(1964) Kubin and Presley (Ref. 8, 1964): These authors have made computations on hydrogen based on the assumption of ideal mixture behavior with only 6 electronic states used in the electronic partition function evaluation, no lowering of the ionization potential, and treatment of molecular hydrogen as a rigid rotation-harmonic oscillation in its electronic ground state. The results are within a few percent of the "more rigorous" computations of Rosenbaum and Levitt.

(1963) Vasiliu (Ref. 15, 1963): This report presents tables of air properties at high pressures obtained from a computer program developed for this purpose. The approximations to the thermodynamic properties have been calculated using "air" as a single substance and adding the second and third virial corrections to existing or extrapolated values of these properties. Tables are given for the temperature range of 200 to 6000°K and for values of the common logarithm of the relative density up to 2.80, with (1) temperature and logarithm density and (2) temperature and entropy, as the independent variables.

(1962) Rosenbaum and Levitt (Ref. 27, 1962): This report presents hydrogen properties based upon a modified Altman truncation scheme which is an extension of Bethe's original proposal of truncating the partition function at a point where the classical electron orbit diameter approaches the distance between atoms. The above scheme takes into consideration the influence of the azimuthal quantum number upon the mean volume of the excited atom.

(1962) Hochstim and Terhune (Ref. 23, 1962): Properties of argon plasma to 100,000°K are presented. The computations were performed by an iterative technique so that the effect of Debye-Hückel corrections could be included to account for lowering of the ionization potential and termination of the partition function series.

(1962) Lick and Emmons (Ref. 17, 1962): Computations are presented of helium properties using the Unsöld theory to account for the partition function series truncation and ionization potential lowering. The results ran in pressures from  $10^{-4}$  to  $10^3$  atm for temperatures up to 50,000°K and are presented both in tabular and chart form. The authors compare the Unsöld, Bethe, and Debye truncation criterion. They conclude that the effect of ionization potential lowering and series truncation is less than a few percent, and that enthalpy, entropy, and composition would be accurate to within 10% if only ground states were considered. However, computations in which ionization potential lowering was examined indicated an error of about 30% in  $\text{He}^+$  concentrations at 30,000°K and a pressure of 100 atm. The error is reduced at higher and lower temperatures and is much smaller at lower pressures.

(1962) Raymond (Ref. 20, 1962): Calculation of equilibrium composition and Mollier data for  $\text{CO}_2$  based upon ideal real-gas analysis for a pressure range of  $10^{-4}$  to  $10^2$  atm and temperature up to 24,000°K is presented.

(1962) Vale (Ref. 25, 1962) and Barnard and Hawkins (Ref. 24, 1962): These papers deal with computations of the composition of mixtures which involve multiphase products.

(1962) Michels and Schneidern (Ref. 22, 1962): A method is presented for computing the high-temperature properties of mixtures in which the virial equation of state rather than the ideal equation of state is used to represent the components of the mixture. Comparisons of real and ideal behavior are presented for air, water gas and the  $\text{H}_2\text{-O}_2$  system. The imperfect behavior is most pronounced at high pressure, i.e., about 100 to 1000 atm. Below this, the ideal or perfect real-gas assumption is valid.

(1962) Boynton (Ref. 21, 1962): An analysis is presented of the problem of calculating the equilibrium composition of an imperfect single-phase mixture. Computational techniques derived by the method of White, Johnson, and Dantzig and by the Newton-Raphson method are compared for gases obeying the virial equation of state. For the non-ideal or imperfect mixtures, the conclusion is that the Newton-Raphson method is simpler.

(1962) Rudin and Ragent (Ref. 26, 1962): A review of the state of the art, sources of discrepancy, and accuracy of thermodynamic computations for air is presented.

(1962) Marrone (Ref. 18, 1962): The (Ref. 36, 1959) work at NBS forms the basis for this tabulation of shock parameters in air.

(1961) Strahl (Ref. 32, 1961): A computation of thermodynamic properties and composition is presented for a model Venus atmosphere 85% CO<sub>2</sub>-15% N<sub>2</sub> which is based upon perfect real-gas computations similar to the early air work of Gilmore at the Rand Corporation. The maximum temperature was 24,000°K.

(1961) Hochstim (Ref. 30, 1961): Ideal real-gas tabulations are presented of thermodynamic properties, chemical composition, and normal shock parameters for air with additional carbon, O<sub>2</sub>, N<sub>2</sub>, and CO<sub>2</sub>. The thermochemistry is based upon early NBS tabulations. The maximum temperature was 10,000°K.

(1961) Stupochenko, Stakhanov, Samuilov, Pleshanov, and Rozhdestvenskii (Ref. 28, 1961): The authors present properties of air below 12,000°K over a sizable range of pressure. Particularly noteworthy in this work is the detailed analysis of the vibrational-rotational partition function and the estimates of those effects which were neglected. The authors used an ideal mixture assumption.

(1961) Sherman and Martinek (Ref. 29, 1961): This paper which is primarily concerned with the electrical transport properties of mercury presents some data on the thermodynamic property of this ionized gas. The method of Ecker and Weizel was used to account for the ionization potential lowering and electronic partition function truncation due to the micro field associated with the plasma. The authors present a comparison of the degree of ionization with a fixed arbitrary cut off and with the systematic treatment of the partition function,

truncation, and ionization potential lowering. At 15,000°K the maximum difference of 3% occurs at the highest pressure considered, i.e., one atmosphere.

(1960) Wachman, Linevsky, and Lyon (Ref. 34, 1960): Computations were made using a perfect real-gas analysis to determine the properties of air-carbon mixtures. The equilibrium mixture contained 1 to 30% carbon by weight. Tabulations of equilibrium composition for a temperature range 4000° to 10,000°K and for pressure ranging from 0.1 to 100 atm are presented.

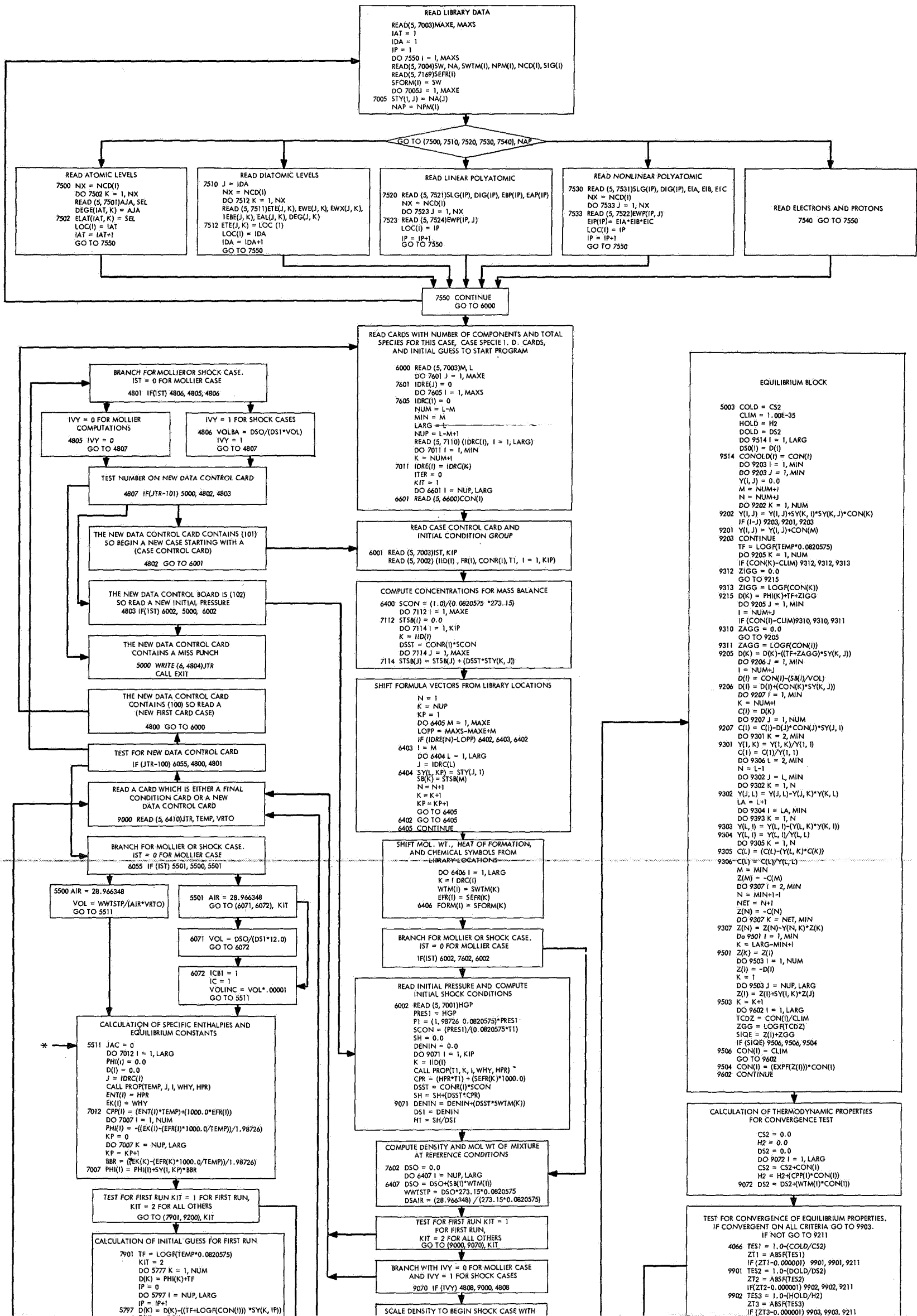
(1960) McGee and Heller (Ref. 16, 1960): Enthalpy and entropy tabulations as function of temperature and pressure are presented for hydrogen, helium, and lithium for temperatures up to 50,000°K. Their analysis takes into consideration imperfect gas behavior through the inclusion of the configuration integral term in the partition function. The potential used in its evaluation was based upon Debye-Hückel theory. Ionization potential lowering was based upon the work of Ecker and Weizel which is similar to the Unsöld theory. Truncation of the partition function sum is based upon the Inglis and Teller theory.

(1960) Brandmaier and Harnett (Ref. 35, 1960): A good survey of the techniques used for computing equilibrium composition of mixtures is presented.

(1959) Hilsenrath, Klein, and Woolley (Ref. 36, 1959): This earlier work at NBS treated air as an ideal or perfect real gas.

(1958) Rosher (Ref. 37, 1958): The author discussed the implicit and explicit assumptions made in computing the equilibrium composition of a gaseous mixture using the methods of statistical mechanics from available spectroscopic data.

**Appendix B**  
**Program Flow Diagram**







## Appendix C

### Key to Format Statements

I-format is used to enter fixed point numbers. The prefix numeral indicates the number of words. The suffix numeral indicates the number of spaces per word. For example, 2I3 is used to enter the numbers 5 and 416 in the following key punch form:

Word No. 1	Word No. 2
0 0 5	4 1 6

F-format is used to enter floating point numbers for which the exponent notation is not needed to locate the decimal position. The prefix indicates the number of words. The suffix, which is a decimal, indicates the number of spaces per word and the number of digits from the right to the decimal position. For example, 3F5.3 is used to enter the numbers 0.72, 4.95 and 0.012 in the following key punch form:

Word No. 1	Word No. 2	Word No. 3
0 . 7 2 0	4 . 9 5 0	0 . 0 1 2

E-format is used to enter floating point numbers for which the exponent notation is used to locate the decimal position. The prefix indicates the number of words. The suffix, which is a decimal number, indicates the number of spaces per word and the number of digits following the decimal. For example, 2E12.6 is used to enter the numbers  $6.32 \times 10^5$  and .072 in the following key punch form:

Word No. 1
0 . 6 3 2 0 0 0 E + 0 6 0 . 7 2 0 0 0 0 E - 0 1

A-format is used to enter alpha-numeric data. The suffix is used to indicate the number of spaces per word. For example, A4 is used to enter CO<sub>2</sub> in the following key punch form:

C 0 2
-------

X-format is used for spacing. No values are read or printed.

Blanks in the key punch fields are read as zeros.

## Appendix D

### Program Listing and Description of Some of the Symbols and Fields Used in the Program

```

ODIMENSION NA(10),SWTM(100),NPM(100),NCD(100),SEFR(100),SFORM(100),
1STY(100,10),IDRC(100),IDRE(10),CON(100),IID(9),FR(9),CONR(9),
2STSB(10),SY(100,10),Y(10,10),SB(100),WTM(100),EFR(100),FORM(100),
3ENT(100),EK(100),CPP(100),PHI(100),D(100),C(100),Z(100),DSO(100),
4CONOLD(100),DELDH(2),SFCON(100),SIG(100),ELAT(32,70),DEGE(32,70),
5LOC(100),SLG(75),DIG(75),EBP(75),
6EAP(75),EIP(75),EWP(75,12),ETE(20,10),EWE(20,10),
7EWX(20,10),EBE(20,10),EAL(20,10),DEG(20,10)
COMMON NPM, NCD,SIG,ELAT,DEGE,EWE,EWX,EBE,EAL,DEG,LOC,SLG,DIG,EBP
1,EAP,EIP,EWP,ETE,SWTM,FORM,FR,CONR,CON,BRPT2,BRPT5,BRPT6,BRPT7,
2BRPT8,CS2,DIN,DSAIR,DS2,ENTHA,ENTRO,HOT,PRESR,PRES1,PRES2,PRPT1,
3RHOR,RRTT2,RRTT5,RRTT6,RRTT7,STAGE,TEMP,T1,VGAS,VRTO,WWTSTP,ZSA,
4KIP,LARG
4804 FORMAT(15H THE VALUE JTR=,I4,2X,12HIS NOT VALID)
6410 FORMAT(1I3,2E12.6)
6600 FORMAT(4X,1E12.6)
7001 FORMAT(1E12.6)
7002 FORMAT(I2,A6,1F7.5,1F12.6)
7003 FORMAT(2I2)
7004 FORMAT(12X,A6,10I2,1F10.6,I2,I4,1F4.1)
7110 FORMAT(20I4)
7169 FORMAT(20X,1E12.6)
9601 FORMAT(3X,A4,2X,2E12.6)
7501 FORMAT(17X,1F8.2,3X,1F10.3)
7511 FORMAT(25X,1F11.4,1F9.3,1F10.5,1F10.6,1F9.6,1F6.2)
7521 FORMAT(25X,2F3.1,2F8.5)
7522 FORMAT(25X,1F8.2)
7524 FORMAT(1F8.2)
7531 FORMAT(25X,2F3.1,1F8.5,4X,1F8.5,4X,1F8.5,4X)
C
C      READ LIBRARY DATA
C
      READ INPUT TAPE 5,7003,MAXE,MAXS
      IAT=1
      IDA=1
      IP=1
      DO 7550 I=1,MAXS
      READ INPUT TAPE 5,7004,SW,NA,SWTM(I),NPM(I),NCD(I),SIG(I)
      READ INPUT TAPE 5,7169,SEFR(I)
      SFORM(I)=SW
      DO7005J=1,MAXE
7005  STY(I,J)=NA(J)
      NAP=NPM(I)
      GO TO(7500,7510,7520,7530,7540),NAP
C
C      READ ATOMIC LEVELS
C
7500  NX=NCD(I)
      DO 7502 K=1,NX
      READ INPUT TAPE 5,7501,AJA,SEL
      DEGE(IAT,K)=AJA
7502  ELAT(IAT,K)=SEL
      LOC(I)=IAT
      IAT=IAT+1
      GO TO 7550
C
C      READ DIATOMIC LEVELS
C

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7510 J=IDA
      NX=NCD(I)
      DO 7512 K=1,NX
        READ INPUT TAPE 5,7511,ETE(J,K),EWE(J,K),EWX(J,K),
        1EBE(J,K),EAL(J,K),DEG(J,K)
7512 ETE(J,K)=ETE(J,K)+(EWE(J,K)-EWE(J,1))*0.5-(EWX(J,K)-EWX(J,1))
        1*0.25
        LOC(I)=IDA
        IDA=IDA+1
        GO TO 7550
C
C      READ LINEAR POLYATOMIC
C
7520 READ INPUT TAPE 5,7521,SLG(IP),DIG(IP),EBP(IP),EAP(IP)
      NX=NCD(I)
      DO 7523 J=1,NX
7523 READ INPUT TAPE 5,7524,EWP(IP,J)
      LOC(I)=IP
      IP=IP+1
      GO TO 7550
C
C      READ NONLINEAR POLYATOMIC
C
7530 READ INPUT TAPE 5,7531,SLG(IP),DIG(IP),EIA,EIB,EIC
      NX=NCD(I)
      DO 7533 J=1,NX
7533 READ INPUT TAPE 5,7522,EWP(IP,J)
      EIP(IP)=EIA*EIB*EIC
      LOC(I)=IP
      IP=IP+1
      GO TO 7550
C
C      READ ELECTRONS AND PROTONS
C
7540 GO TO 7550
7550 CONTINUE
      GO TO 6000
C
C      READ CARDS WITH NUMBER OF COMPONENTS AND TOTAL SPECIES FOR THIS
C      CASE ,CASE SPECIE I. D. CARDS, AND INITIAL GUESS TO START PROGRAM
C
6000 READ INPUT TAPE 5,7003,M,L
      DO 7601 J=1,MAXE
7601 IDRE(J)=0
      DO 7605 I=1,MAXS
7605 IDRC(I)=0
      NUM=L-M
      MIN=M
      LARG=L
      NUP=L-M+1
      READ INPUT TAPE 5,7110,(IDRC(I),I=1,LARG)
      DO 7011 I=1,MIN
      K=NUM+I
7011 IDRE(I)=IDRC(K)
      ITER=0
      KIT=1
      DO 6601 I=NUP,LARG
6601 READ INPUT TAPE 5,6600,CON(I)
C

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```

C      READ CASE CONTROL CARD AND INITIAL CONDITION GROUP
C
6001 READ INPUT TAPE 5,7003,IST,KIP
    READ INPUT TAPE 5,7002,(IID(I),FR(I),CONR(I),T1,I=1,KIP)
C
C      COMPUTE CONCENTRATIONS FOR MASS BALANCE
C
6400 SCON=(1.0)/(0.0820575*273.15)
    DO 7112 I=1,MAXE
7112 STSB(I)=0.0
    DO 7114 I=1,KIP
        K=IID(I)
        DSST=CONR(I)*SCON
        DO 7114 J=1,MAXE
7114 STSB(J)=STSB(J)+(DSST*STY(K,J))
C
C      SHIFT FORMULA VECTORS FROM LIBRARY LOCATIONS
C
        N=1
        K=NUP
        KP=1
        DO 6405 M=1,MAXE
            LOPP=MAXS-MAXE+M
            IF (IDRE(N)-LOPP) 6402,6403,6402
6403 I=M
            DO 6404 L=1,LARG
                J=IDRC(L)
6404 SY(L,KP)=STY(J,I)
                SB(K)=STSB(M)
                N=N+1
                K=K+1
                KP=KP+1
            GO TO 6405
6402 GO TO 6405
6405 CONTINUE
C
C      SHIFT MOL. WT., HEAT OF FORMATION , AND CHEMICAL SYMBOLS FROM
C      LIBRARY LOCATIONS
C
        DO 6406 I=1,LARG
            K=IDRC(I)
            WTM(I)=SWTM(K)
            EFR(I)=SEFR(K)
6406 FORM(I)=SFORM(K)
C
C      BRANCH FOR MOLLIER OR SHOCK CASE. IST=0 FOR MOLLIER CASE
C
        IF(IST) 6002,7602,6002
C
C      READ INITIAL PRESSURE AND COMPUTE INITIAL SHOCK CONDITIONS
C
6002 READ INPUT TAPE 5,7001,HGP
    PRES1=HGP
    P1=(1.98726/0.0820575)*PRES1
    SCON=(PRES1)/(0.0820575*T1)
    SH=0.0
    DENIN=0.0
    DO 9071 I=1,KIP
        K=IID(I)

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      CALL PROP(T1,K,I,WHY,HPR)
      CPR=(HPR*T1)+(SEFR(K)*1000.0)
      DSST=CONR(I)*SCON
      SH=SH+(DSST*CPR)
9071 DENIN=DENIN+(DSST*SWTM(K))
      DS1=DENIN
      H1=SH/DS1
C
C      COMPUTE DENSITY AND MOL WT OF MIXTURE AT REFERENCE CONDITIONS
C
7602 DS0=0.0
      DO6407 I=NUP,LARG
6407 DS0=DS0+(SB(I)*WTM(I))
      WWTSTP=DS0*273.15*0.0820575
      DSAIR=(28.966348)/(273.15*0.0820575)
C
C      TEST FOR FIRST RUN KIT=1 FOR FIRST RUN,KIT=2 FOR ALL OTHERS
C
      GO TO(9000,9070),KIT
C
C      BRANCH WITH IVY=0 FOR MOLLIER CASE AND IVY=1 FOR SHOCK CASES
C
9070 IF(IVY) 4808,9000,4808
C
C      SCALE DENSITY TO BEGIN SHOCK CASE WITH DENSITY RATIO FROM PREVIOUS
C      SOLUTION
C
4808 VOL=DS0/(DS1*VOLBA)
      GO TO 9000
C
C      READ A CARD WHICH IS EITHER A FINAL CONDITION CARD OR A NEW DATA
C      CONTROL CARD
C
9000 READ INPUT TAPE 5, 6410,JTR,TEMP,VRTD
C
C      TEST FOR NEW DATA CONTROL CARD
C
      IF(JTR-100)6055,4800,4801
C
C      THE NEW DATA CONTROL CARD CONTAINS (100)SO READ A(NEW FIRST CARD CASE)
C
4800 GO TO 6000
C
C      BRANCH FOR MOLLIER OR SHOCK CASE. IST=0 FOR MOLLIER CASE
C
4801 IF(IST) 4806,4805,4806
C
C      IVY=0 FOR MOLLIER COMPUTATIONS
C
4805 IVY=0
      GO TO 4807
C
C      IVY=1 FOR SHOCK CASES
C
4806 VOLBA=DS0/(DS1*VOL)
      IVY=1
      GO TO 4807
C
C      TEST NUMBER ON NEW DATA CONTROL CARD

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C
4807 IF(JTR-101) 5000,4802,4803
C
C      THE NEW DATA CONTROL CARD CONTAINS(101) SO BEGIN A NEW CASE STARTING
C      WITH A (CASE CONTROL CARD)
C
4802 GO TO 6001
C
C      THE NEW DATA CONTROL CARD IS (102) SO READ A NEW INITIAL PRESSURE
C
4803 IF (IST) 6002,5000,6002
C
C      THE NEW DATA CONTROL CARD CONTAINS A MISS PUNCH
C
5000 WRITE OUTPUT TAPE 6,4804,JTR
      CALL EXIT
C
C      BRANCH FOR MOLLIER OR SHOCK CASE. IST=0 FOR MOLLIER CASE
C
6055 IF(IST)5501,5500,5501
5500 AIR=28.966348
      VOL=WWTSTP/(AIR*VRT0)
      GO TO 5511
5501 AIR=28.966348
      GO TO(6071,6072),KIT
6071 VOL=DS0/(DS1*12.0)
      GO TO 6072
6072 ICB1=1
      IC=1
      VOLINC=VOL*.00001
      GO TO 5511
C
C      CALCULATION OF SPECIFIC ENTHALPIES AND EQUILIBRIUM CONSTANTS
C
5511 JAC=0
      DO 7012 I=1,LARG
      PHI(I)=0.0
      D(I)=0.0
      J=IDRC(I)
      CALL PROP(TEMP,J,I,WHY,HPR)
      ENT(I)=HPR
      EK(I)=WHY
7012 CPP(I)=(ENT(I)*TEMP)+(1000.0*EFR(I))
      DO 7007 I=1,NUM
      PHI(I)= -((EK(I)-(EFR(I)*1000.0/TEMP))/1.98726)
      KP=0
      DO 7007 K=NUP,LARG
      KP=KP+1
      BBR=((EK(K)-(EFR(K)*1000.0/TEMP))/1.98726)
7007 PHI(I)=PHI(I) +SY(I,KP)*BBR
C
C      TEST FOR FIRST RUN KIT=1 FOR FIRST RUN,KIT=2 FOR ALL OTHERS
C
      GO TO(7901,9200),KIT
C
C      CALCULATION OF INITIAL GUESS FOR FIRST RUN
C
7901 TF=LOGF(TEMP*0.0820575)
      KIT=2

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      DO 5777 K=1,NUM
      D(K)=PHI(K)+TF
      IP=0
      DO 5797 I=NUP,LARG
      IP=IP+1
5797 D(K)=D(K)-((TF+LOGF(CON(I)))*SY(K,IP))
      D(K)=-D(K)
5777 CON(K)=EXP(-D(K))
C
C      BEGINNING OF EQUILIBRIUM ITERATION LOOP
C
9200 JAC=JAC+1
C
C      TEST FOR NUMBER OF ITERATIONS
C
      IF(JAC-50) 5003, 5003, 5991
C
C      EQUILIBRIUM BLOCK NOT CONVERGENT
C
5002 FORMAT(49H EQUILIBRIUM STATE NOT REACHED IN 50 ITERATIONS)
5991 WRITE OUTPUT TAPE 6,5002
      DO 5004 I=1,LARG
      WRITE OUTPUT TAPE 6,9601,FORM(I),CON(I),D(I)
5004 WRITE OUTPUT TAPE 6,9601,FORM(I),CONOLD(I),DSO(I)
      CALL EXIT
C
C      EQUILIBRIUM BLOCK
C
5003 COLD=CS2
      CLIM=1.00E-35
      HOLD=H2
      DOLD=DS2
      DO 9514 I=1,LARG
      DSO(I)=D(I)
9514 CONOLD(I)=CON(I)
      DO 9203 I=1,MIN
      DO 9203 J=1,MIN
      Y(I,J)=0.0
      M=NUM+I
      N=NUM+J
      DO 9202 K=1,NUM
9202 Y(I,J)=Y(I,J)+SY(K,I)*SY(K,J)*CON(K)
      IF(I-J) 9203,9201,9203
9201 Y(I,J)=Y(I,J)+CON(M)
9203 CONTINUE
      TF=LOGF(TEMP*0.0820575)
      DO 9205 K=1,NUM
      IF(CON(K)-CLIM)9312,9312,9313
9312 ZIGG=LOGF(CLIM)
      GO TO 9215
9313 ZIGG=LOGF(CON(K))
9215 D(K)=PHI(K)+TF+ZIGG
      DO 9205 J=1,MIN
      I=NUM+J
      IF(CON(I)-CLIM)9310,9310,9311
9310 ZAGG=LOGF(CLIM)
      GO TO 9205
9311 ZAGG=LOGF(CON(I))
9205 D(K)=D(K)-((TF+ZAGG)*SY(K,J))

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DO 9206 J=1,MIN
I=NUM+J
D(I)=CON(I)-(SB(I)/VOL)
DO 9206 K=1,NUM
9206 D(I)=D(I)+(CON(K)*SY(K,J))
DO 9207 I=1,MIN
K=NUM+I
C(I)=D(K)
DO 9207 J=1,NUM
9207 C(I)=C(I)-D(J)*CON(J)*SY(J,I)
DO 9301 K=2,MIN
9301 Y(1,K)=Y(1,K)/Y(1,1)
C(1)=C(1)/Y(1,1)
DO 9306 L=2,MIN
N=L-1
DO 9302 J=L,MIN
DO 9302 K=1,N
9302 Y(J,L)=Y(J,L)-Y(J,K)*Y(K,L)
LA=L+1
DO 9304 I=LA,MIN
DO 9303 K=1,N
9303 Y(L,I)=Y(L,I)-(Y(L,K)*Y(K,I))
9304 Y(L,I)=Y(L,I)/Y(L,L)
DO 9305 K=1,N
9305 C(L)=C(L)-(Y(L,K)*C(K))
9306 C(L)=C(L)/Y(L,L)
M=MIN
Z(M)=-C(M)
DO 9307 I=2,MIN
N=MIN+1-I
NET=N+1
Z(N)=-C(N)
DO 9307 K=NET,MIN
9307 Z(N)=Z(N)-Y(N,K)*Z(K)
DO 9501 I=1,MIN
K=LARG-MIN+I
9501 Z(K)=Z(I)
DO 9503 I=1,NUM
Z(I)=-D(I)
K=1
DO 9503 J=NUP,LARG
Z(I)=Z(I)+SY(I,K)*Z(J)
9503 K=K+1
DO 9602 I=1,LARG
TCDZ=CON(I)/CLIM
ZGG=LOGF(TCDZ)
SIQE=Z(I)+ZGG
IF(SIQE) 9506,9506,9504
9506 CON(I)=CLIM
GO TO 9602
9504 CON(I)=(EXP(Z(I)))*CON(I)
9602 CONTINUE
C
C      CALCULATION OF THERMODYNAMIC PROPERTIES FOR CONVERGENCE TEST
C
CS2=0.0
H2=0.0
DS2=0.0
DO 9072 I=1,LARG

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      CS2=CS2+CON(I)
      H2=H2+(CPP(I)*CON(I))
9072 DS2=DS2+(WTM(I)*CON(I))
C
C      TEST FOR CONVERGENCE OF EQUILIBRIUM PROPERTIES.  IF CONVERGENT ON ALL
C      CRITERIA GO TO 9903.  IF NOT GO TO 9211
C
4066 TES1=1.0-(COLD/CS2)
      ZT1=ABSF(TES1)
      IF(ZT1-0.000001) 9901, 9901, 9211
9901 TES2=1.0-(DOLD/DS2)
      ZT2=ABSF(TES2)
      IF(ZT2-0.000001) 9902, 9902, 9211
9902 TES3=1.0-(HOLD/H2)
      ZT3=ABSF(TES3)
      IF(ZT3-0.000001) 9903, 9903, 9211
C
C      AVERAGE CONCENTRATIONS FROM CONSECUTIVE ITERATIONS FOR FIRST 4
C      ITERATIONS TO IMPROVE RATE OF CONVERGENCE
C
9211 IF(JAC-4) 9221,9221,9200
9221 DO 9220 I=1,LARG
9220 CON(I)=(CONOLD(I)+CON(I))/2.0
      GO TO 9200
C
C      COMPUTE FINAL THERMODYNAMIC PROPERTIES
C
9903 H2=H2/DS2
      T2=TEMP
      P2=(1.98726*T2*CS2)
      PRES2=(T2*CS2*0.0820575)
      ZSA=(PRES2*273.15*VOL)/T2
      ENTHA=H2/(1.98726*273.15)*WWTSTP
      ENTRO=-LOGF(PRES2)
      DO 9695 I=1,LARG
      AFR=CON(I)/CS2
      IF(CON(I))9694,9696,9694
9696 TOM=0.0
      GO TO 9695
9694 TOM=LOGF(AFR)
9695 ENTRO=ENTRO+(AFR*((ENT(I)+EK(I))/1.98726)-TOM)
      ENTRO=ENTRO*(CS2/DS2)*WWTSTP
C
C      BRANCH FOR MOLLIER OR SHOCK CASE.  IST=0 FOR MOLLIER CASE
C
      IF(IST) 9800,9090,9800
C
C      BRANCH POINT FOR MOVING, STANDING, AND REFLECTED SHOCK CASES
C
9800 GO TO (9093,9820,9830,9840,9844,9850,9860,9870,9880),ICB1
Q
C      BEGIN MOVING SHOCK CALCULATIONS
C
9093 DHP=H2-H1
      DHH=((P2-P1)*(DS2+DS1))/(2.0*DS1*DS2)
      DELDH(IC)=DHP-DHH
      VPP=DS2/DSAIR
      GO TO ( 9414, 9404),IC
C

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C          INCREMENT INVERSE DENSITY AND COMPUTE EQUILIBRIUM PROPERTIES
C
9414 IC=2
    VOL=VOL+VOLINC
    GO TO 9400
9401 IC=1
9400 JAC=0
    VOLINC=VOL*.00001
    GO TO 9200
C
C          TEST FOR CONVERGENCE OF MOVING SHOCK SOLUTION
C
9404 DVOL=((VOLINC)/(1.0-(DELDH(1)/DELDH(2))))
    HTES=ABSF(DVOL/VOL)
    VOL=VOL-DVOL
    IF(HTES-0.001 )9405,9405,9415
C
C          TEST FOR THE NUMBER OF MOVING SHOCK ITERATIONS
C
9415 CONTINUE
    ITER=ITER+1
    IF(ITER-20)9401,9401,5551
C
C          MOVING SHOCK SOLUTION NOT CONVERGENT IN 20 ITERATIONS
C
5550 FORMAT (41H NO HUGONIOT SOLUTION IN 20 ITERATIONS )
5551 WRITE OUTPUT TAPE 6, 5550
    CALL EXIT
C
C          COMPUTE THE MOVING SHOCK PARAMETERS
C
9405 VELO=((1000.027*1013250.0)*(DS2/DS1)*(PRES1-PRES2))/(DS1-DS2)
    VEL=SQRTF(VELO)
    VEL=VEL/(12.0*2.54001)
    HDT=H1/(1.98726*273.15)*WWTSTP
    DIN=DS1/DS0
    PRESR=PRES2/PRES1
    RHOR=DS2/DS1
    VGAS=VEL*(1.0-(1.0/RHOR))
    BEXAR=(VEL*12.0*2.54001)/RHOR
    STAGE=ENTHA+((BEXAR*BEXAR)/(2.0*273.15*0.831433E+08))*WWTSTP
C
C          MOVING SHOCK SOLUTION PRINT
C
9420 KDT=2
    PRPT1=VEL
    CALL PRIN(KDT)
    ITER=0
C
C          BRANCH FOR DIFFERENT SHOCK CASES (IST=1 MOVING SHOCK CASE)
C          (IST=2 COMBINATION OF MOVING ,STANDING ,AND REFLECTED SHOCK CASES)
C          (IST=3 COMBINATION OF MOVING AND STANDING SHOCK CASES)
C          (IST=4 COMBINATION OF MOVING AND REFLECTED SHOCK CASES)
C
    GO TO(9000,9810,9810,9810),IST
C
C          MOLLIER CASE SOLUTION PRINT
C
9090 KDT=1

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      CALL PRIN(KDT)
      GO TO 9000
C
C      STORE QUANTITIES FROM MOVING SHOCK SOLUTION FOR NEXT CASE
C
9810 TEMPSF=TEMP
      DO 9811 IKE=1,LARG
9811 SFCON(IKE)=CON(IKE)
      VOLSF=VOL
      FK1=(12.0*2.54001)
      FK2=(FK1*FK1)/(1.01325E+09)
      FK3=(0.4184)*(2.7315)*(1.98726E+10)
      FK4=(FK1*FK1)/FK3
      CAPS=(VGAS/VOLSF)**2
      UHB1=(ENTHA)+(FK4/2.0)*(VGAS*VGAS)*(WWTSTP)
      UPB1= PRES2 +(FK2*DS2*VGAS*VGAS)
      DENS=6.0/VOL
      VOL=(1.0/DENS)
C
C      BRANCH FOR COMBINATION SHOCK, STANDING SHOCK, OR REFLECTED SHOCK
C
      ITBT=IST-1
      GO TO(9812,9812,9813),ITBT
C
C      COMPUTE QUANTITIES FOR STANDING SHOCK
C
9812 CAPP=(FK2)*(CAPS)*(DS0)*(VOL)
      CAPH=(FK4)*(WWTSTP)*(CAPS)*(VOL)*(VOL)
      ICB1=2
      GO TO 9400
C
C      COMPUTE QUANTITIES FOR REFLECTED SHOCK
C
9813 ICB1=6
      DDTT=VOLSF*DENS-1.0
      CAPH=(WWTSTP*VGAS*VGAS*FK4)/DDTT
      CAPP=CAPH/VOLSF
      GO TO 9400
C
C      START OF STANDING SHOCK ITERATION BLOCK. THE TEMPERATURE IS INCREASED
C      BY 500 DEG INCREMENTS UNTIL AN ENTHALPY GREATER THAN THAT GIVEN BY THE
C      STANDING SHOCK EQUATION IS FOUND. THIS IS POINT(A)
C
9820 HEAT3= UHB1-(CAPH/2.0)
      THH=HEAT3-ENTHA
      IF(THH) 9822,9821,9821
9821 TEMP=TEMP+500.0
      GO TO 5511
C
C      COMPUTE PROPERTIES AT A 500 DEG INCREMENT FROM POINT (A) HOLDING
C      THE DENSITY CONSTANT THIS IS POINT (B)
C
9822 HA=ENTHA
      PA=PRES2
      TA=TEMP
      DRA=(1.0/VOL)
      TEMP=TEMP-500.0
      ICB1=3
      GO TO 5511

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C
C      COMPUTE PROPERTIES AT AN INCREMENT IN DENSITY FROM POINT (B) BUT
C      AT THE SAME TEMPERATURE. THIS IS POINT (C)
C
9830 HB=ENTHA
    PB=PRES2
    TB=TEMP
    DRB=(1.0/VOL)
    DRC=DRB*(1.0001)
    VOL=1.0/DRC
    ICB1=4
    GO TO 9400
C
C      BY LINEAR EXTRAPOLATION COMPUTE A NEW DENSITY AND TEMPERATURE
C      SOLUTION POINT
C
9840 HC=ENTHA
    PC=PRES2
    TC=TEMP
    DRC=(1.0/VOL)
    PTEX=UPB1-CAPP
    HTEX=UHB1-(CAPH/2.0)
    PDPT=(PA-PB)/500.0
    PDHT=(HA-HB)/500.0
    PDHR=(HC-HB)/(DRC-DRB)
    PDPR=(PC-PB)/(DRC-DRB)
    DENOM=(PDHT*PDPR)-(PDHR*PDPT)
    DUME1=(PTEX-PB)*PDHT
    DUME2=(HTEX-HB)*PDPT
    DRNEW=(DUME1-DUME2)/DENOM
    DNEW=DRB+DRNEW
    VOL=(1.0/DNEW)
    T3=TB+(HTEX-HB-(PDHR*DRNEW))/PDHT
    ADT=(DNEW/DRB)-1.0
    AVD=ABSF(ADT)
    CAPP=(FK2)*(CAPS)*(DSO)*(VOL)
    CAPH=(FK4)*(WWTSTP)*(CAPS)*(VOL)*(VOL)
C
C      TEST FOR CONVERGENCE OF STANDING SHOCK SOLUTION
C
    IF(AVD-0.001) 9843,9843,9842
C
C      SET NEW INVERSE DENSITY VALUE AND BEGIN STANDING SHOCK
C      ITERATION PROCESS AGAIN
C
9842 VOL=(1.0/DNEW)
    TEMP=TB
    ICB1=2
    GO TO 9400
C
C      COMPUTE CONCENTRATIONS FOR STANDING SHOCK CONDITION
C
9843 TEMP=T3
    VOL=1.0/DNEW
    ICB1=5
    GO TO 5511
C
C      COMPUTE STANDING SHOCK PARAMETERS AND PRINT
C

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9844 XOST=(18.81*WWTSTP)/(2.0*847500.0*273.15*1.9873)
      BRPT2=WWTSTP/(28.96*VOL)
      BRPT8=VOLSF/VOL
      BRPT7=PRES2/(PRES1*PRESR)
      BRPT6=VGAS/BRPT8
      BRPT5=ENTHA+(BRPT6*BRPT6*XOST)
      KDT=3
      CALL PRIN(KDT)
C
C      BRANCH POINT AFTER STANDING SHOCK SOLUTION. STANDING SHOCK CASE
C      ONLY 9847 OR TO BEGIN REFLECTED SHOCK CASE 9846
C
      GO TO (9846,9847),ITBT
C
C      STARTING POINT FOR REFLECTED SHOCK CASE WHEN STANDING SHOCK CASE HAS
C      BEEN COMPLETED
C
9846 ICB1=6
      DDTT=VOLSF/VOL-1.0
      CAPH=(WWTSTP*VGAS*VGAS*FK4)/DDTT
      CAPP=CAPH/VOLSF
      GO TO 5511
C
C      START OF REFLECTED SHOCK ITERATION BLOCK. THE TEMPERATURE IS INCREASED
C      BY 500 DEG INCREMENTS UNTIL AN ENTHALPY GREATER THAN THAT GIVEN BY THE
C      REFLECTED SHOCK EQUATION IS FOUND. THIS IS POINT(A)
C
9850 HEAT4=UHB1+CAPH
      TSHH=HEAT4-ENTHA
      IF(TSHH) 9852,9851,9851
9851 TEMP=TEMP+500.0
      GO TO 5511
C
C      COMPUTE PROPERTIES AT A 500 DEG INCREMENT FROM POINT (A) HOLDING
C      THE DENSITY CONSTANT THIS IS POINT (B)
C
9852 HA=ENTHA
      PA=PRES2
      TA=TEMP
      DRA=(1.0/VOL)
      TEMP=TEMP-500.0
      ICB1=7
      GO TO 5511
C
C      COMPUTE PROPERTIES AT AN INCREMENT IN DENSITY FROM POINT (B) BUT
C      AT THE SAME TEMPERATURE. THIS IS POINT (C)
C
9860 HB=ENTHA
      PB=PRES2
      TB=TEMP
      DRB=(1.0/VOL)
      DRC=DRB*(1.0001)
      VOL=1.0/DRC
      ICB1=8
      GO TO 9400
C
C      BY LINEAR EXTRAPOLATION COMPUTE A NEW DENSITY AND TEMPERATURE
C      SOLUTION POINT
C

```

```

9870 HC=ENTHA
      PC=PRES2
      TC=TEMP
      DRC=(1.0/VOL)
      PTEX=UPB1+CAPP
      HTEX=UHB1+CAPH
      PDPT=(PA-PB)/500.0
      PDHT=(HA-HB)/500.0
      PDHR=(HC-HB)/(DRC-DRB)
      PDPR=(PC-PB)/(DRC-DRB)
      DENOM=(PDHT*PDPR)-(PDHR*PDPT)
      DUME1=(PTEX-PB)*PDHT
      DUME2=(HTEX-HB)*PDPT
      DRNEW=(DUME1-DUME2)/DENOM
      DNEW=DRB+DRNEW
      VOL=(1.0/DNEW)
      T3=TB+(HTEX-HB-(PDHR*DRNEW))/PDHT
      ADT=(DNEW/DRB)-1.0
      AVD=ABSF(ADT)
      DDTT=(VOLSF/VOL)-1.0
      CAPH=(WWTSTP*VGAS*VGAS*FK4)/DDTT
      CAPP=(WWTSTP*VGAS*VGAS*FK4)/(DDTT*VOLSF)
C
C      TEST FOR CONVERGENCE OF REFLECTED SHOCK SOLUTION
C
      IF(AVD-0.001) 9873,9873,9872
C
C      SET NEW INVERSE DENSITY RATIO AND BEGIN REFLECTED SHOCK
C      ITERATION PROCESS AGAIN
C
9872 VOL=(1.0/DNEW)
      TEMP=TB
      ICB1=6
      GO TO 9400
C
C      COMPUTE CONCENTRATIONS FOR REFLECTED SHOCK CONDITION
C
9873 TEMP=T3
      VOL=1.0/DNEW
      ICB1=9
      GO TO 5511
C
C      COMPUTE REFLECTED SHOCK PARAMETERS AND PRINT
C
9880 RRTT2=WWTSTP/(28.96*VOL)
      RRTT6=VOLSF/VOL
      RRTT5=PRES2/(PRES1*PRESR)
      RRTT7=VGAS/(RRTT6-1.0)
      KDT=4
      CALL PRIN(KDT)
C
C      END POINT FOR REFLECTED AND STANDING CASES. BEGIN A NEW MOVING
C      SHOCK COMPUTATION
C
9847 VOL=VOLSF
      TEMP=TEMSF
      DO 9879 I=1,LARG
9879 CON(I)=SFCON(I)
      ITER=0

```

```

GO TO 9000
END
C      SUBROUTINE FOR THERMOCHEMICAL PROPERTIES
      SUBROUTINE PROP(TDC,I,K,PEK,PENT)
      DIMENSION NPM(100),NCD(100),EAP(75),EIP(75),EWP(75,12),ETE(20,10),
      IEWE(20,10),EWX(20,10),EBE(20,10),EAL(20,10),DEG(20,10),SIG(100),
      2DEGE(32,70),LOC(100),SLG(75),DIG(75),EBP(75),ELAT(32,70),SWTM(100)
      COMMON NPM, NCD,SIG,ELAT,DEGE,EWE,EWX,EBE,EAL,DEG,LOC,SLG,DIG,EBP
      1,EAP,EIP,EWP,ETE,SWTM
C      TDC=TEMPERATURE
C      I=I.D.NUMBER FOR SPECIE IN INPUT LIST
C      K=I.D.NUMBER FOR THE ORDER OF THE SPECIE IN EQUILIBRIUM COMPUTATION
C      PENT=(H-H0)/T
C      PEK=-(F-H0)/T
C      DEG(I,J)=DEGENERACY OF THE J,TH ENERGY LEVEL FOR I,TH SPECIE
C      EL(I,J)=J,TH ENERGY LEVEL FOR I,TH SPECIE
      THA=(1.3804*TDC)/(2.99793*0.66252)
      PGR=1.98726
      IBB=NPM(I)
      GO TO(5711,5721,5731,5741,5751),IBB
C      COMPUTATIONFOR ATOMIC THERMOCHEMICAL PROPERTIES
5711 BLB=20.0
      ELL=2.3026*BLB*THA
      J=1
      L=LOC(I)
      PAR=0.0
      HNT=0.0
5700 IF(ELAT(L,J)-ELL) 5710,5710,5720
5710 VTV=DEGE(L,J)*EXP(-ELAT(L,J)/THA)
      PAR=PAR+VTV
      HNT=HNT+(ELAT(L,J)*VTV)
      J=J+1
      IF(J-NCD(I))5700,5700,5720
5720 PENT =PGR*(2.5+((HNT)/(PAR*THA)))
      AB=LOGF(PAR)
      BB=LOGF(SWTM(I))
      CB=LOGF(TDC)
      PEK =PGR*((2.5*CB)+(1.5*BB)+AB)-7.28353
      GO TO 5790
C      COMPUTATION FOR DIATOMIC PROPERTIES
5721 A=SWTM(I)
      L=LOC(I)
      W=EWE(L,1)-(2.0*EWX(L,1))
      BBB=EBE(L,1)-(0.5*EAL(L,1))
      B=(2.99793*0.66252*BBB)/(1.3804)
      G= DEG(L,1)/SIG(I)
      AA=LOGF(A)
      CC=LOGF(G)
      BB=LOGF(B)
      E=1.5*AA-BB+CC
      ENF=(PGR*E)-7.28353
      T=LOGF(TDC)
      TR=3.5*PGR*T
      ENF=ENF+TR
      U=W/THA
      EU=EXP(U)
      EUD=EU-1.0
      TED=EU/EUD
      TL=LOGF(TED)

```

```

PEK=ENF+(PGR*TL)
PENT=PGR*(3.5+(U/EUD))
CEK=1.0
CEN=0.0
NX=NCD(I)
DO 5722 IH=2,NX
G=EXP(-ETE(L,IH)/THA)
TL=(DEG(L,IH)*G)/DEG(L,1)
CEK=CEK+TL
5722 CEN=CEN+(TL*ETE(L,IH))
AA=LOGF(CEK)
BB=CEN/CEK
PEK=PEK+(AA*PGR)
PENT=PENT+((BB*PGR)/THA)
A=EBE(L,1)/EWE(L,1)
B=EAL(L,1)/EBE(L,1)
AA=EWX(L,1)/EWE(L,1)
BB=(8.0*A)/U
BF2=B/EUD
BF3=(2.0*AA*U)/(EUD*EUD)
BE2=((U*B*EU)-(2.0*U*AA))/(EUD*EUD)
BE3=(4.0*AA*U*U*EU)/(EUD*EUD*EUD)
CEK=PGR*(BB+BF2+BF3)
CEN=PGR*(BB+BE2+BE3)
PEK=PEK+CEK
PENT=PENT+CEN
GO TO 5790
C      COMPUTATION FOR LINEAR POLYATOMIC PROPERTIES
5731 EN=3.5*PGR
ENF=-7.28353
A=SWTM(I)
L=LOC(I)
G=DIG(L)/SLG(L)
BB=LOGF(G)
AB=LOGF(A)
AA=PGR*((1.5*AB)+BB)
ENF=ENF+AA
BBB=EBP(L)-(0.5*EAP(L))
B=(2.99793*0.66252*BBB)/(1.3804)
BB=LOGF(B)
AA=PGR*BB
ENF=ENF-AA
ET=LOGF(TDC)
AB=ET*PGR*3.5
ENF=ENF+AB
N=NCD(I)
DO 7532 J=1,N
U=ERP(L,J)/THA
EU=EXP(U)
EUD=EU-1.0
EP=EU/EUD
EH=U/EUD
EL=LOGF(EP)
EHR=PGR*EH
ELR=PGR*EL
EN=EN+EHR
7532 ENF=ENF+ELR
PEK=ENF
PENT=EN

```



```

      GO TO 5790
C      COMPUTATION FOR NONLINEAR POLYATOMIC PROPERTIES
5741 EN=4.0
      ENF=0.0
      A=SWTM(I)
      L=LOC(I)
      B=EIP(L)
      G=DIG(L)/SLG(L)
      CB=LOGF(G)
      BB=LOGF(B)
      AB=LOGF(A)
      AA=(1.5*AB)+(0.5*BB)+CB
      CAN=(PGR*AA)-17.1614
      N=NCD(I)
      DO 7542 J=1,N
      U=EWL(L,J)/THA
      EU=EXPF(U)
      EUD=EU-1.0
      EP=EU/EUD
      EH=U/EUD
      EL=LOGF(EP)
      EN=EN+EH
7542 ENF=ENF+EL
      BNT=LOGF(TDC)
      PEK=CAN+PGR*((4.0*BNT)+ENF)
      PENT=PGR*(EN)
      GO TO 5790
C      COMPUTATION FOR ELECTRONS AND PROTONS
5751 PENT=(2.50*PGR)
      AB=LOGF(TDC)
      BB=LOGF(SWTM(I))
      TEE=1.0-SWTM(I)
      IF(TEE)5753,5753,5752
5752 CB=LOGF(2.0)
      GO TO 5754
5753 CB=0.0
5754 PEK=PGR*((2.5*AB)+(1.5*BB)+CB)-7.28353
      GO TO 5790
5790 RETURN
      END
      SUBROUTINE PRIN(KDT)
      ODIMENSION NA(10),SWTM(100),NPM(100),NCD(100),SEFR(100),SFORM(100),
      1STY(100,10),IDRC(100),IDRE(10),CON(100),IID(9),FR(9),CONR(9),
      2STSB(10),SY(100,10),Y(10,10),SB(100),WTM(100),EFR(100),FORM(100),
      3ENT(100),EK(100),CPP(100),PHI(100),D(100),C(100),Z(100),DSO(100),
      4CONOLD(100),DELDH(2),SFCON(100),SIG(100),ELAT(32,70),DEGE(32,70),
      5LOC(100),SLG(75),DIG(75),EBP(75),
      6EAP(75),EIP(75),EWP(75,12),ETE(20,10),EWE(20,10),
      7EWX(20,10),EBE(20,10),EAL(20,10),DEG(20,10)
      COMMON NPM, NCD,SIG,ELAT,DEGE,EWE,EWX,EBE,EAL,DEG,LOC,SLG,DIG,EBP
      1,EAP,EIP,EWP,ETE,SWTM,FORM,FR,CONR,CON,BRPT2,BRPT5,BRPT6,BRPT7,
      2BRPT8,CS2,DIN,DSAIR,DS2,ENTHA,ENTRO,HOT,PRESR,PRES1,PRES2,PRPT1,
      3RHOR,RRTT2,RRTT5,RRTT6,RRTT7,STAGE,TEMP,T1,VGAS,VRTO,WWTSTP,ZSA,
      4KIP,LARG
      300 FORMAT(1H1 9X,70H CHEMICAL EQUILIBRIUM COMPOSITION AND PROPERTIES
      1OF A COMPLEX MIXTURE /)
      301 FORMAT(1H1 9X,38H REAL GAS NORMAL SHOCK WAVE PROPERTIES /)
      3020FORMAT(1H1 9X,49H REAL GAS STANDING NORMAL SHOCK WAVE PROPERTIES
      1/)

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3030FORMAT(1H1 9X,49H REAL GAS REFLECTED NORMAL SHOCK WAVE PROPERTIES
1/)
304 FORMAT(34H INITIAL CONDITIONS TEMP DEG K--1E10.4,4X,8H P/P0--
1 1E10.4)
3050FORMAT(40H MOLE FRACTION OF A6,13H AT S.T.P
1.---1F7.5)
306 FORMAT(49H CONCENTRATION MOLE FRACTION )
307 FORMAT(36H PARTICLES/CC MOLE/LITER )
308 FORMAT(3X,A6,3E12.4)
309 FORMAT(100H CONCENTRATION MOLE FRACTION
1 CONCENTRATION MOLE FRACTION )
310 FORMAT(82H PARTICLES/CC MOLE/LITER
1 PARTICLES/CC MOLE/LITER )
311 FORMAT(3X,A6,3E12.4,3X,A6,3E12.4)
3120FORMAT(26H TEMPERATURE (DEG K)-----1E11.4,14H RHO1/RHO0---1E10.4
1,26H RHO0 (GM/LITER)-----1E10.4)
3130FORMAT(26H ENTHALPY (HMO/RT0)-----1E11.4,14H P1/P0-----1E10.4
1,26H MIXTURE COLD MOL WT-----1E10.4)
3140FORMAT(26H SHOCK VELOCITY (FPS)-----1E11.4,14H RHO2/RHO0-- 1E10.4
1,26H VEL BEHIND SHOCK (FPS)--1E10.4)
3150FORMAT(26H TEMPERATURE (DEG K)-----1E11.4,14H RHO2/RHO1-- 1E10.4
1,26H ENTROPY (SMO/R)-----1E10.4)
3160FORMAT(26H ENTHALPY (HMO/RT0)-----1E11.4,14H P2/P0----- 1E10.4
1,26H Z-----1E10.4)
3170FORMAT(26H STAG ENTHALPY (HMO/RT0)-1E11.4,14H P2/P1----- 1E10.4
1)
318 FORMAT(22H INITIAL CONDITIONS )
319 FORMAT(20H FINAL CONDITIONS )
3200FORMAT(26H SHOCK VELOCITY (FPS)-----1E11.4,14H P3/P0-----1E10.4
1,26H VEL BEHIND ST SHOCK-----1E10.4)
3210FORMAT(26H TEMPERATURE (DEG K)-----1E11.4,14H P3/P2-----1E10.4
1,26H STAG ENTHALPY (HMO/RT0)-1E10.4)
3220FORMAT(26H ENTHALPY (HMO/RT0)-----1E11.4,14H RHO3/RHO0---1E10.4
1,26H ENTROPY (SMO/R)-----1E10.4)
3230FORMAT(26H MO-MIXTURE COLD MOL WT--1E11.4,14H RHO3/RHO2---1E10.4
1,26H Z-----1E10.4)
3240FORMAT(50H RHO0 (GM/L)-1E10.4
1)
3250FORMAT(26H SHOCK VELOCITY (FPS)-----1E11.4,14H P4/P0-----1E10.4
1,26H REF SHOCK VEL (FPS)-----1E10.4)
3260FORMAT(26H TEMPERATURE (DEG K)-----1E11.4,14H P4/P2-----1E10.4
1,26H ENTROPY (SMO/R)-----1E10.4)
3270FORMAT(26H ENTHALPY (HMO/RT0)-----1E11.4,14H RHO4/RHO0---1E10.4
1,26H Z-----1E10.4)
3280FORMAT(26H MO-MIXTURE COLD MOL WT--1E11.4,14H RHO4/RHO2---1E10.4
1,26H RHO0-----1E10.4)
3290FORMAT(26H TEMPERATURE (DEG K)-----1E11.4,14H P/P0-----1E10.4
1,26H ENTROPY (SMO/R)-----1E10.4)
3300FORMAT(26H ENTHALPY (HMO/RT0)-----1E11.4,14H RHO/RHO0---1E10.4
1,26H RHO0-----1E10.4)
3310FORMAT(26H MIXTURE COLD MOL WT--1E11.4,14H Z-----1E10.4
1/)
GO TO(1300,1310,1320,1330),KDT
1300 WRITE OUTPUT TAPE 6,300
WRITE OUTPUT TAPE 6,318
GO TO 1340
1310 WRITE OUTPUT TAPE 6,301
WRITE OUTPUT TAPE 6,318
VPP=(DIN*WWTSTP)/28.9662

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        WRITE OUTPUT TAPE 6,312,T1,VPP,DSAIR
        WRITE OUTPUT TAPE 6,313,HOT,PRES1,WWTSTP
        GO TO 1340
1320 WRITE OUTPUT TAPE 6,302
        GO TO 1331
1330 WRITE OUTPUT TAPE 6,303
1331 WRITE OUTPUT TAPE 6,304,T1,PRES1
1340 DO 1341 I=1,KIP
1341 WRITE OUTPUT TAPE 6,305,FR(I),CONR(I)
        WRITE OUTPUT TAPE 6,319
        IF(LARG-35) 1350,1370,1370
1350 WRITE OUTPUT TAPE 6,306
        WRITE OUTPUT TAPE 6,307
        DO 1351 I=1,LARG
        CPMOL=CON(I)*0.602338E+21
        CONFR=CON(I)/CS2
1351 WRITE OUTPUT TAPE 6,308,FORM(I),CPMOL,CON(I),CONFR
        GO TO 1376
1370 WRITE OUTPUT TAPE 6,309
        WRITE OUTPUT TAPE 6,310
        I=1
1371 J=I+1
        IF(J-LARG) 1373,1374,1375
1373 CONF1=CON(I)/CS2
        CONF2=CON(J)/CS2
        CPML1=CON(I)*0.602338E+21
        CPML2=CON(J)*0.602338E+21
        WRITE OUTPUT TAPE 6,311,FORM(I),CPML1,CON(I),CONF1,FORM(J),CPML2,
1CON(J),CONF2
        I=I+2
        GO TO 1371
1374 CONF1=CON(I)/CS2
        CONF2=CON(J)/CS2
        CPML1=CON(I)*0.602338E+21
        CPML2=CON(J)*0.602338E+21
        WRITE OUTPUT TAPE 6,311,FORM(I),CPML1,CON(I),CONF1,FORM(J),CPML2,
1CON(J),CONF2
        GO TO 1376
1375 CONF1=CON(I)/CS2
        CPML1=CON(I)*0.602338E+21
        WRITE OUTPUT TAPE 6,308,FORM(I),CPML1,CON(I),CONF1
1376 CONTINUE
        GO TO (1380,1385,1390,1395),KDT
1380 WRITE OUTPUT TAPE 6,329,TEMP,PRES2,ENTRO
        WRITE OUTPUT TAPE 6,330,ENTHA,VRTO,DSAIR
        WRITE OUTPUT TAPE 6,331,WWTSTP,ZSA
        GO TO 1399
1385 VPP=DS2/DSAIR
        WRITE OUTPUT TAPE 6,314,PRPT1,VPP,VGAS
        WRITE OUTPUT TAPE 6,315,TEMP,RHOR,ENTRO
        WRITE OUTPUT TAPE 6,316,ENTHA,PRES2,ZSA
        WRITE OUTPUT TAPE 6,317,STAGE,PRESR
        GO TO 1399
1390 WRITE OUTPUT TAPE 6,320,PRPT1,PRES2,BRPT6
        WRITE OUTPUT TAPE 6,321,TEMP,BRPT7,BRPT5
        WRITE OUTPUT TAPE 6,322,ENTHA,BRPT2,ENTRO
        WRITE OUTPUT TAPE 6,323,WWTSTP,BRPT8,ZSA
        WRITE OUTPUT TAPE 6,324,DSAIR
        GO TO 1399
1395 WRITE OUTPUT TAPE 6,325,PRPT1,PRES2,RRTT7
        WRITE OUTPUT TAPE 6,326,TEMP,RRTT5,ENTRO
        WRITE OUTPUT TAPE 6,327,ENTHA,RRTT2,ZSA
        WRITE OUTPUT TAPE 6,328,WWTSTP,RRTT6,DSAIR
1399 RETURN
        END

```

## Nomenclature for Library Data

MAXE = Number of components required to represent the formula vectors of all chemical species considered in the library. This is the number of elements plus the electron.

MAXS = Total number of chemical species considered in the library.

NA = Temporary storage for formula vector.

SWTM = Library storage location for molecular weights.

NPM = Identification number for the different types of chemical species :  
(atomic and atomic ion species NPM = 01), (diatomic species NPM = 02), (linear polyatomic NPM = 03), (nonlinear polyatomic NPM = 04), (Protons, HE III, Electrons NPM = 05).

NCD = Number of data cards which make up a library packet for each specie:

- (1) For atomic and atomic ion species. This is the number of energy level cards in the library and  $NCD \leq 70$ .
- (2) For diatomic species. This is the number of electronic energy level cards in the library and  $NCD \leq 10$ .
- (3) For polyatomic species. This is the number of vibrational frequencies in the library and  $NCD \leq 12$ . Only ground electronic levels are considered.

SEFR = Library storage location of heats of formation (kcal/mole).

SFORM = Library storage location of chemical symbols for species.

STY = Library storage location for the components of the formula vector for each specie.

DEGE = Atomic specie degeneracy  $\Sigma (2j + 1)$  for an energy level, g.  $\left. \begin{array}{l} \text{ } \end{array} \right\} \begin{array}{l} 32 \text{ species;} \\ 70 \text{ levels.} \end{array}$

ELAT = Atomic specie energy levels,  $\epsilon$ .

LOC = Number locating energy level library data within the respective blocks for atomic, diatomic and polyatomic species.

<p>ETE = Diatomic <math>T_e</math>, electronic term values</p> <p>EWE = Diatomic <math>\omega_e</math></p> <p>EWX = Diatomic <math>\omega_e x_e</math></p> <p>EBE = Diatomic <math>B_e</math></p> <p>EAL = Diatomic <math>\alpha_e</math></p> <p>DEG = Diatomic <math>p</math>, degeneracy</p> <p>SIG = Diatomic <math>\sigma</math> symmetry parameter</p>	$\left. \begin{array}{l} \text{ } \end{array} \right\}$	<p>vibrational parameters</p> <p>rotational parameters</p>	$\left. \begin{array}{l} \text{ } \end{array} \right\}$	<p>Dimensioned for a maximum of 20 species with 10 electronic levels.</p>
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## Nomenclature for Library Data (contd)

SLG = Linear and nonlinear polyatomic $\sigma$ symmetry parameter	$\left. \begin{array}{l} \text{EBP} = \text{Linear polyatomic } Be \\ \text{EAP} = \text{Linear polyatomic } \alpha_e \\ \text{EIP} = \text{Nonlinear polyatomic } I_A \times I_B \times I_C \\ \text{EWP} = \text{Nonlinear and linear polyatomic } \omega \end{array} \right\} \text{rotational parameters}$	Dimensions for 75 polyatomic species with a max of 12 vibration frequencies each.
DIG = Linear and nonlinear polyatomic $p$ degeneracy		
EBP = Linear polyatomic $Be$		
EAP = Linear polyatomic $\alpha_e$		
EIP = Nonlinear polyatomic $I_A \times I_B \times I_C$ product of moments of inertia		

## Nomenclature for Case Input

- LARG = Total number of species in the high temperature case under consideration.
- MIN = Total number of formula vector components required to represent the species in the high temperature mixture; the number of elements plus the electron.
- IDRC = Identification number used to select the high temperature product species from the library list.
- IDRE = Identification number used to select the high temperature components (electrons plus elements) from the library list.
- CON = Storage locations of the high temperature species concentrations.
- ITER = Number of iterations through the equilibrium block (i.e., mass action and mass balance equations).
- KIT = A number which is "1" when externally supplied initial guesses of component concentrations are used to start the computation. For all successive computations in which the concentrations from the previous solution are used to start the concentration computation, the value is "2."

## Nomenclature for Case Input (contd)

IST = Case control number:

- (1) IST = 00; Equilibrium solution or Mollier data case
- (2) IST = 01; Moving shock case
- (3) IST = 02; Combination moving, standing, and reflected cases
- (4) IST = 03; Combination of moving and standing cases
- (5) IST = 04; Combination of moving and reflected cases

KIP = Number of species in the initial or STP mixture.

IID = Identification number used to select the reactants data from the library list.

FR = Chemical symbol for reactant.

CONR = Fraction of reactant in mixture at STP.

STSB =  $B_i$ —concentration of the  $i$ th element at STP, moles/liter.

SY = Storage location of modified formula vector used in the case computation.

SB = Concentration "STSB" or  $B_i$  used in the case computation.

WTM = Storage location of molecular weights used in the case computation.

EFR = Storage location of heat of formation used in the case computation.

FORM = Storage location of chemical symbol used in the case computation.

ENT = Computed  $(H - H_o)/T$ .

EK = Computed  $-(F - H_o)/T$ .

CPP =  $H$ , cal/mole.

PHI = Computed equilibrium constant.

TEMP = Equilibrium temperature, °K.

HGP = Initial pressure  $P_1$  in atmospheres (atm).

T1 = Initial temperature,  $T_1$  in °K.

VRTO = Equilibrium density ratio,  $\rho/\rho_o$ .

## Appendix E

### Library Data

0866

```

1      C2      2 0 0 0 0 0 0 0 0 0 24.07230 02   9 2.0
C2 HEAT FORMATION 0.196000E 03
C2 LC A SINGLET SIGMA      0.0      1854.7      13.34      1.820      0.01760      1.0
C2 X TRIPLET SIGMA      716.2      1641.3      11.67      1.633      0.01640      6.0
C2 TRIPLET PI      6435.1      1470.4      11.19      1.499      0.01630      3.0
C2 LC B SINGLET PI      8391.1      1608.3      12.08      1.616      0.01690      2.0
C2 SINGLET DELTA      10186.3      1481.0      11.14      1.486      0.01500      2.0
C2 SINGLET SIGMA      14191.5      1470.0      10.06      1.475      0.013      1.0
C2 TRIPLET SIGMA      14282.1      1891.0      14.23      1.897      0.01800      3.0
C2 A TRIPLET PI      19950.3      1788.2      16.44      1.753      0.0161      6.0
C2 C SINGLET PI      34277.4      1809.1      15.81      1.733      0.018      2.0
2      N2      0 2 0 0 0 0 0 0 0 0 28.01340 02   2 2.0
N2 HEAT FORMATION 0.      E 00
N2 X SINGLET SIGMA      0.0      2358.1      14.19      1.990      0.0170      1.0
N2 A TRIPLET SIGMA      50206.0      1460.4      13.89      1.430      0.0130      3.0
3      O2      0 0 2 0 0 0 0 0 0 0 31.99880 02   4 2.0
O2 HEAT FORMATION 0.      E 00
O2 X TRIPLET SIGMA      0.0      1580.4      12.07      1.446      0.0158      3.0
O2 LC A SINGLET DELTA      7918.1      1509.0      12.00      1.426      0.0171      2.0
O2 LC B SIGMA      13195.2      1432.7      13.95      1.400      0.0182      1.0
O2 A TRIPLET SIGMA      36096.0      819.0      22.50      1.050      3.0
4      H2      0 0 0 2 0 0 0 0 0 0 2.01594 02   1 2.0
H2 HEAT FORMATION .000000E 00
H2 X SINGLET SIGMA      0.0      4395.2      117.9      60.81      2.99000      1.0
5      CN      1 1 0 0 0 0 0 0 0 0 26.01785 02   3 1.0
CN HEAT FORMATION 0.109000E 03
CN X DOUBLET SIGMA      0.0      2068.8      13.19      1.900      0.0180      2.0
CN A DOUBLET PI      9246.0      1812.6      12.70      1.720      0.0170      4.0
CN B DOUBLET SIGMA      25752.1      2163.7      20.10      1.970      0.0220      2.0
6      CO      1 0 1 0 0 0 0 0 0 0 28.01055 02   2 1.0
CO HEAT FORMATION -.271990E 02
CO X SINGLET SIGMA      0.0      2169.8      13.29      1.931      0.0175      1.0
CO LC A TRIPLET PI      48687.4      1743.6      14.47      1.691      0.0195      6.0
7      CH      1 0 0 1 0 0 0 0 0 0 13.01912 02   5 1.0
CH HEAT FORMATION .141600E 03
CH X DOUBLET PI      0.0      2861.6      64.30      14.460      0.53400      2.0
CH X DOUBLET PI      17.9      2861.6      64.30      14.460      0.53400      2.0
CH A DOUBLET DELTA      23150.0      2921.0      90.40      14.910      0.67000      4.0
CH B DOUBLET SIGMA      25993.0      1808.0      201.50      7.170      0.53000      2.0
CH C DOUBLET SIGMA      31828.0      2073.4      57.00      7.880      0.28000      2.0
8      NO      0 1 1 0 0 0 0 0 0 0 30.00610 02   4 1.0
NO HEAT FORMATION 0.214500E 02
NO X DOUBLET PI      0.0      1903.7      13.97      1.705      0.0178      2.0
NO X DOUBLET PI      121.1      1903.4      13.97      1.705      0.0178      2.0
NO LC A QUAD PI      38406.0      1019.0      12.80      1.220      0.0198      8.0
NO A DOUBLET SIGMA      43966.0      2371.0      14.50      1.995      0.0164      2.0
9      NH      0 1 0 1 0 0 0 0 0 0 15.01467 02   4 1.0
NH HEAT FORMATION +.790000E+02
NH X TRIPLET SIGMA      0.0      3315.0      94.7      16.67      0.65      3.0
NH LC A SINGLET DELTA      14970.0      3186.0      29.0      16.54      0.17      2.0
NH LC B SINGLET SIGMA      23415.0      3480.0      16.40      1.0
NH A TRIPLET PI      29819.0      3232.5      99.2      16.69      0.74      6.0
10     OH      0 0 1 1 0 0 0 0 0 0 17.00737 02   2 1.0
OH HEAT FORMATION .925000E 01
OH X DOUBLET PI      0.0      3738.4      81.80      18.850      0.70000      4.0
OH A DOUBLET SIGMA      32687.0      3179.4      103.30      17.230      0.78000      2.0
11     CO+     1 0 1 0 0 0 0 -1 0 0 28.01000 02   3 1.0
CO+ HEAT FORMATION 0.295900E 03

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CO+	X DOUBLET SIGMA	0.0	2214.3	15.16	1.977	0.0190	2.0
CO+	A DOUBLET PI	20733.0	1562.1	13.53	1.589	0.01940	4.0
CO+	B DOUBLET SIGMA	45877.0	1734.2	27.93	1.800	0.03020	2.0
12	CH+	1 0 0 1 0 0 0-1	0 0 13.01857	02	2	0.0	
CH+	HEAT FORMATION	+.398100E+03					
CH+	X SINGLET SIGMA	0.0	2739.7		14.18	0.49	1.0
CH+	A SINGLET PI	24171.3	1840.1	98.5	11.89	0.92	2.0
13	N2+	0 2 0 0 0 0 0-1	0 0 28.01285	02	3	2.0	
N2+	HEAT FORMATION	0.359297E 03					
N2+	X DOUBLET SIGMA	0.0	2207.6	16.90	1.9300	0.02000	2.0
N2+	A DOUBLET PI	9168.3	1902.8	14.60	1.7400	0.01800	4.0
N2+	B DOUBLET SIGMA	25452.9	2441.1	31.20	2.0850	0.02500	2.0
14	NO-	0 1 1 0 0 0 0 1	0 0 30.00665	02	1	1.0	
NO-	HEAT FORMATION	+.140000E+02					
NO-	TRIPLET SIGMA	0.0	1600.0	12.0	1.4	0.016	3.0
15	NO+	0 1 1 0 0 0 0-1	0 0 30.008	02	2	1.0	
NO+	HEAT FORMATION	.235170E 03					
NO+	X SINGLET SIGMA	0.0	2377.1	16.35	2.0020	0.02020	1.0
NO+	LC A TRIPLET SIGMA	40332.0	1572.0				3.0
16	O2+	0 0 2 0 0 0 0-1	0 0 31.99825	02	4	2.0	
O2+	HEAT FORMATION	0.278178E+03					
O2+	X DOUBLET PI	0.0	1876.4	16.53	1.6722	0.01984	2.0
O2+	X DOUBLET PI	200.0	1876.4	16.53	1.6722	0.01984	2.0
O2+	LC,A QUAD PI	31900.0	1035.7	10.39	1.1047	0.01575	8.0
O2+	A DOUBLET PI	38795.0	900.0	13.40	1.0617	0.01906	4.0
17	O2-	0 0 2 0 0 0 0 1	0 0 31.99935	02	2	2.0	
O2-	HEAT FORMATION	-.133700E 02					
O2-	DOUBLET PI LC G	0.0	1170.0	8.50	1.247	0.0130	4.0
O2-	DOUBLET PI LC U	29400.0					4.0
18	OH+	0 0 1 1 0 0 0-1	0 0 17.00682	02	2	1.0	
OH+	HEAT FORMATION	+.316000E+03					
OH+	X TRIPLET SIGMA	0.0	3229.0	131.0	16.74	0.73	3.0
OH+	TRIPLET PI	28459.0	2202.0	108.0	13.75	0.86	6.0
19	OH-	0 0 1 1 0 0 0 1	0 0 17.00792	02	1	1.0	
OH-	HEAT FORMATION	-.329000E+02					
OH-	X SINGLET SIGMA	0.0	3735.0	74.7	18.87	0.65	1.0
20	H2+	0 0 0 2 0 0 0-1	0 0 2.01539	02	1	2.0	
H2+	HEAT FORMATION	+.355740E+03					
H2+	X DOUBLET SIGMA	0.0	2297.0	62.0	29.80	1.4	2.0
21	C3	3 0 0 0 0 0 0 0 0	0 0 36.03345	03	4	0.0	
C3	HEAT FORMATION	0.188000E+03					
C3	SINGLET SIGMA	2.01.0 0.4280	0.0	4			
	500.0						
	500.0						
	1300.0						
	2200.0						
22	C2N2	2 2 0 0 0 0 0 0 0	0 0 52.03570	03	7	0.0	
C2N2	HEAT FORMATION	0.734000E+02					
C2N2	SINGLET SIGMA	2.01.0 0.1605	0.0	7			
	226.0						
	226.0						
	506.0						
	506.0						
	848.0						
	2149.0						
	2322.0						
23	CO2	1 0 2 0 0 0 0 0 0	0 0 44.00995	03	4	0.0	
CO2	HEAT FORMATION	-.939630E+02					
CO2	SINGLET SIGMA	2.01.0 0.39021	0.0022	4			



1388.17  
 667.40  
 667.40  
 2349.16  
 24 C2H 2 0 0 1 0 0 0 0 0 25.03027 03 4 0.0  
 C2H HEAT FORMATION .116700E+03  
 C2H DOUBLET PI 1.04.0 1.4482 0.0 4  
 1923.4  
 3216.7  
 712.0  
 712.0  
 25 C2H2 2 0 0 2 0 0 0 0 0 26.03824 03 7 0.0  
 C2H2 HEAT FORMATION .543300E+02  
 C2H2 SINGLET SIGMA 2.01.0 1.17649 0.0 7  
 3372.0  
 1973.0  
 3288.0  
 729.0  
 729.0  
 612.0  
 612.0  
 26 N2O 0 2 1 0 0 0 0 0 0 44.01280 03 4 0.0  
 N2O HEAT FORMATION +.204350E+02  
 N2O SINGLET SIGMA 1.01.0 0.41901 0.0040 4  
 2223.7  
 588.8  
 588.8  
 1284.9  
 27 CH2 1 0 0 2 0 0 0 0 0 14.02709 04 3 0.0  
 CH2 HEAT FORMATION .950000E+02  
 CH2 SINGLET A 2.01.0 1.39300E-40 2.57100E-40 3.96400E-4003  
 CH2 SINGLET A 3105.0  
 CH2 SINGLET A 3019.0  
 CH2 SINGLET A 1570.0  
 28 CH3 1 0 0 3 0 0 0 0 0 15.03506 04 6 0.0  
 CH3 HEAT FORMATION .334000E+02  
 CH3 DOUBLET A 6.02.0 5.85500E-40 2.92800E-40 2.92800E-4006  
 CH3 DOUBLET A 2920.0  
 CH3 DOUBLET A 447.0  
 CH3 DOUBLET A 1228.0  
 CH3 DOUBLET A 1228.0  
 CH3 DOUBLET A 3000.0  
 CH3 DOUBLET A 3000.0  
 29 CH4 1 0 0 4 0 0 0 0 0 16.04303 04 9 0.0  
 CH4 HEAT FORMATION -.159900E+02  
 CH4 SINGLET A 12.01.0 5.33000E-40 5.33000E-40 5.33000E-4009  
 CH4 SINGLET A 2914.2  
 CH4 SINGLET A 1526.0  
 CH4 SINGLET A 1526.0  
 CH4 SINGLET A 1300.2  
 CH4 SINGLET A 1300.2  
 CH4 SINGLET A 1300.2  
 CH4 SINGLET A 3020.3  
 CH4 SINGLET A 3020.3  
 CH4 SINGLET A 3020.3  
 30 CH0+ 1 0 1 1 0 0 0 -1 0 0 29.01797 04 3 0.0  
 CH0+ HEAT FORMATION .205000E+03  
 CH0+ SIGMA 1.02.018.72920E-4012.51470E-4019.92120E-4003  
 CH0+ SIGMA 2700.0

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CH0+  SIGMA          1820.2
CH0+  SIGMA          1083.0
31      CH20    1 0 1 2 0 0 0 0 0 0 30.02649 04    6 0.0
      CH2OHEAT FORMATION -.269000E+02
CH20  SIGMA          2.01.0 2.97500E-4021.60800E-4024.67000E-4006
CH20  SIGMA          2780.0
CH20  SIGMA          2843.0
CH20  SIGMA          1746.0
CH20  SIGMA          1247.0
CH20  SIGMA          1500.0
CH20  SIGMA          1167.0
32      NO2      0 1 2 0 0 0 0 0 0 0 46.00550 04    3 0.0
      NO2 HEAT FORMATION +.860000E+01
      NO2          2.02.068.20000E-4064.55000E-40 3.48600E-40
      NO2          1319.7
      NO2          749.8
      NO2          1617.75
33      NH3      0 1 0 3 0 0 0 0 0 0 17.03061 04    6 0.0
      NH3 HEAT FORMATION -.934000E+01
      NH3          3.01.0 2.81400E-40 2.81400E-40 4.51730E-40
      NH3          3336.2
      NH3          949.87
      NH3          3443.9
      NH3          3443.6
      NH3          1627.4
      NH3          1626.1
34      O3       0 0 3 0 0 0 0 0 0 0 47.99820 04    3 0.0
      O3 HEAT FORMATION 0.346390E+02
      O3          2.01.0 7.87490E-4062.84400E-4070.88800E-4003
      O3          705.0
      O3          1043.0
      O3          1110.0
35      H2O      0 0 1 2 0 0 0 0 0 0 18.01534 04    3 0.0
      H2O HEAT FORMATION -.571020E+02
      H2O SINGLET SIGMA 2.01.0 1.00400E-40 1.92900E-40 3.01500E-4003
      H2O          3657.0
      H2O          1594.7
      H2O          3755.7
36      HCO      1 0 1 1 0 0 0 0 0 0 29.01852 04    3 0.0
      HCO HEAT FORMATION .000000E+00
      HCO DOUBLET A 1.02.018.72920E-40 1.25147E-4019.92120E-4003
      HCO DOUBLET A 2700.0
      HCO DOUBLET A 1820.2
      HCO DOUBLET A 1083.0
37      C-       1 0 0 0 0 0 0 1 0 0 12.01170 01    51 0.0
      C- HEAT FORMATION .144000E 03
      C-          4.00          0.0
      C-          10.00         19223.000
      C-          6.00         28840.000
      C-          12.00         83366.000
      C-          6.00         86223.187
      C-          12.00         88109.500
      C-          2.00         93582.250
      C-          20.00         94883.062
      C-          12.00         95533.187
      C-          14.00         96864.187
      C-          6.00         97805.750
      C-          10.00         99665.000
      C-          12.00        103736.750

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C-	6.00	104227.375
C-	34.00	104767.000
C-	32.00	104882.687
C-	24.00	105020.000
C-	2.00	106478.562
C-	6.00	106780.062
C-	20.00	106870.687
C-	10.00	107039.000
C-	12.00	109927.875
C-	2.00	110029.187
C-	8.00	110196.000
C-	32.00	110248.000
C-	38.00	110304.000
C-	16.00	110470.500
C-	10.00	110545.750
C-	6.00	112320.750
C-	2.00	112565.875
C-	10.00	112682.562
C-	20.00	112799.000
C-	60.00	112862.000
C-	16.00	112947.500
C-	6.00	114072.000
C-	62.00	114160.000
C-	40.00	114275.000
C-	6.00	114890.000
C-	32.00	114988.000
C-	24.00	115004.000
C-	18.00	115103.000
C-	18.00	115483.000
C-	44.00	115535.000
C-	18.00	115618.000
C-	44.00	115887.000
C-	36.00	115902.000
C-	58.00	116164.000
C-	22.00	116259.000
C-	58.00	116367.000
C-	28.00	116441.000
C-	28.00	116625.000
38	C+	1 0 0 0 0 0 0 -1 0 0 12.0106 01 36 0.0
C+	HEAT FORMATION	0.429537E 03
CII	2.00	0.0
CII	4.00	64.000
CII	2.00	43000.199
CII	4.00	43021.797
CII	6.00	43050.699
CII	6.00	74930.875
CII	4.00	74933.187
CII	2.00	96494.062
CII	6.00	110666.250
CII	2.00	116537.875
CII	6.00	131735.750
CII	4.00	142024.375
CII	10.00	145551.437
CII	10.00	150462.750
CII	2.00	157234.375
CII	6.00	162524.562
CII	6.00	166988.437
CII	6.00	167033.375
CII	16.00	168124.312

CII	14.00	168979.000
CII	2.00	173348.125
CII	6.00	175295.187
CII	6.00	178220.750
CII	10.00	178494.750
CII	14.00	178956.437
CII	2.00	181258.000
CII	20.00	181770.437
CII	6.00	182044.500
CII	10.00	184064.875
CII	14.00	184376.187
CII	4.00	184688.687
CII	12.00	186463.750
CII	10.00	188612.687
CII	2.00	194571.875
CII	28.00	195812.250
CII	20.00	196580.750
39	C++	1 0 0 0 0 0 0 -2 0 0 12.01005 01 20 0.0
C++ HEAT FORMATION 0.991689E 03		
CIII	1.00	0.0
CIII	5.00	52394.797
CIII	4.00	52338.000
CIII	3.00	102351.375
CIII	9.00	137450.500
CIII	5.00	145875.062
CIII	1.00	182520.187
CIII	4.00	238160.687
CIII	12.00	259672.062
CIII	15.00	269962.875
CIII	5.00	276482.687
CIII	12.00	308264.750
CIII	16.00	317748.000
CIII	78.00	321964.687
CIII	49.00	333395.000
CIII	116.00	346526.000
CIII	90.00	358639.000
CIII	29.00	365585.000
CIII	39.00	373748.000
CIII	37.00	381958.000
40	C+++	1 0 0 0 0 0 0 -3 0 0 12.0095 01 22 0.0
C+++HEAT FORMATION 0.209569E 04		
CIV	2.00	0.0
CIV	2.00	64484.199
CIV	4.00	64591.297
CIV	2.00	302847.875
CIV	6.00	320080.000
CIV	10.00	324890.875
CIV	2.00	401346.687
CIV	6.00	408322.187
CIV	10.00	410338.187
CIV	14.00	410434.062
CIV	2.00	445366.062
CIV	6.00	448861.000
CIV	42.00	449948.375
CIV	2.00	468765.000
CIV	6.00	470763.000
CIV	64.00	471407.875
CIV	2.00	482659.000
CIV	6.00	483931.000

CIV	10.00	484309.000	
CIV	54.00	484346.875	
CIV	20.00	492743.000	
CIV	60.00	492745.000	
41	N-	0 1 0 0 0 0 0 1 0 0	14.00725 01 32 0.0
N-	HEAT FORMATION	.115990E 03	
N-	5.00	0.0	
N-	3.00	158.500	
N-	1.00	226.500	
N-	5.00	15867.699	
N-	1.00	33792.398	
N-	5.00	73767.750	
N-	3.00	76794.687	
N-	15.00	86631.000	
N-	9.00	88630.812	
N-	5.00	95476.375	
N-	3.00	96225.500	
N-	45.00	97420.187	
N-	15.00	99094.500	
N-	9.00	99680.375	
N-	15.00	101135.000	
N-	5.00	102116.187	
N-	3.00	102411.625	
N-	5.00	102661.625	
N-	25.00	102865.062	
N-	15.00	102908.125	
N-	9.00	103869.375	
N-	8.00	105019.000	
N-	40.00	105385.250	
N-	9.00	105911.250	
N-	8.00	106545.062	
N-	40.00	106751.187	
N-	8.00	107445.375	
N-	40.00	107573.062	
N-	8.00	108021.375	
N-	40.00	108105.687	
N-	48.00	108470.187	
N-	48.00	108731.500	
42	N+	0 1 0 0 0 0 0 -1 0 0	14.00615 01 54 0.0
N+	HEAT FORMATION	0.447564E 03	
NII	1.00	0.0	
NII	3.00	49.100	
NII	5.00	131.300	
NII	5.00	15315.699	
NII	1.00	32687.098	
NII	5.00	47167.699	
NII	15.00	92237.875	
NII	9.00	109218.187	
NII	5.00	144189.062	
NII	12.00	149077.312	
NII	3.00	155129.875	
NII	3.00	164611.562	
NII	15.00	166679.437	
NII	3.00	166765.687	
NII	3.00	168893.000	
NII	9.00	170667.000	
NII	6.00	174212.875	
NII	21.00	186653.312	
NII	5.00	187092.187	

NII	15.00	187492.687
NII	9.00	188858.062
NII	7.00	189336.000
NII	3.00	190121.125
NII	21.00	209825.250
NII	5.00	209926.875
NII	15.00	210301.875
NII	9.00	210705.375
NII	7.00	211104.750
NII	3.00	211335.500
NII	15.00	220717.000
NII	7.00	221137.562
NII	15.00	206108.687
NII	9.00	211828.750
NII	9.00	196712.125
NII	3.00	197859.250
NII	9.00	214385.250
NII	3.00	214828.000
NII	25.00	224169.250
NII	20.00	226055.187
NII	3.00	202169.875
NII	15.00	202862.000
NII	9.00	203259.687
NII	3.00	203532.750
NII	5.00	205350.687
NII	1.00	206327.500
NII	7.00	211030.875
NII	21.00	211061.000
NII	27.00	211390.750
NII	9.00	211402.875
NII	15.00	211411.250
NII	5.00	211491.125
NII	21.00	221074.250
NII	27.00	221302.187
NII	9.00	221312.062
43	N++	0 1 0 0 0 0 0 -2 0 0 14.00560 01 43 0.0
N++ HEAT FORMATION 0.113022E 04		
NIII	2.00	0.0
NIII	4.00	174.500
NIII	2.00	57192.098
NIII	4.00	57252.000
NIII	6.00	57333.199
NIII	10.00	101023.750
NIII	2.00	131003.500
NIII	2.00	145876.062
NIII	4.00	145986.500
NIII	4.00	186802.250
NIII	10.00	203072.187
NIII	2.00	221302.375
NIII	6.00	230408.562
NIII	6.00	245701.687
NIII	10.00	267244.375
NIII	12.00	287713.875
NIII	6.00	297263.062
NIII	2.00	301088.187
NIII	26.00	309856.687
NIII	6.00	311716.062
NIII	4.00	314224.000
NIII	12.00	317402.250

NIII	10.00	317781.750
NIII	14.00	320287.500
NIII	10.00	321065.750
NIII	2.00	327056.750
NIII	28.00	330396.687
NIII	20.00	332860.000
NIII	12.00	334568.875
NIII	12.00	336213.375
NIII	14.00	339855.687
NIII	10.00	341947.875
NIII	20.00	342752.000
NIII	18.00	343116.000
NIII	10.00	354517.000
NIII	14.00	354955.687
NIII	18.00	355214.000
NIII	12.00	368704.750
NIII	10.00	373376.000
NIII	6.00	374805.250
NIII	20.00	376953.250
NIII	16.00	377970.750
NIII	16.00	379405.000
44	N+++	0 1 0 0 0 0 0 -3 0 0 14.00505 01 23 0.0
N+++ HEAT FORMATION 0.222425E 04		
NIV	1.00	0.0
NIV	4.00	67199.562
NIV	5.00	67343.750
NIV	3.00	130695.000
NIV	4.00	175536.687
NIV	5.00	175661.500
NIV	5.00	188885.000
NIV	1.00	235370.000
NIV	3.00	377206.000
NIV	1.00	388858.000
NIV	12.00	405944.375
NIV	15.00	419979.375
NIV	5.00	429158.000
NIV	9.00	465463.375
NIV	3.00	473032.000
NIV	21.00	484525.000
NIV	40.00	499851.000
NIV	34.00	506292.000
NIV	53.00	516650.000
NIV	7.00	521868.000
NIV	45.00	554419.000
NIV	15.00	574940.000
NIV	20.00	593704.000
45	0-	0 0 1 0 0 0 0 1 0 0 15.99995 01 24 0.0
0- HEAT FORMATION .252000E 02		
0-	4.00	0.0
0-	2.00	404.000
0-	6.00	102406.500
0-	4.00	102681.240
0-	2.00	102841.200
0-	4.00	104731.860
0-	2.00	105057.100
0-	6.00	115918.700
0-	4.00	116041.690
0-	2.00	116144.390
0-	8.00	116988.210

0-	6.00	117164.830
0-	6.00	117310.000
0-	10.00	128219.920
0-	22.00	128516.000
0-	6.00	132596.260
0-	4.00	132745.770
0-	6.00	133000.000
0-	2.00	133224.100
0-	18.00	133546.000
0-	10.00	133606.390
0-	18.00	133924.000
0-	14.00	137603.000
0-	10.00	138708.000
46	0+	0 0 1 0 0 0 0 -1 0 0 15.99885 01 41 0.0
0+ HEAT FORMATION 0.372942E 03		
OII	4.00	0.0
OII	10.00	26808.398
OII	6.00	40466.898
OII	6.00	119837.687
OII	6.00	120001.062
OII	10.00	165987.687
OII	6.00	185340.625
OII	6.00	185499.187
OII	2.00	188888.375
OII	4.00	189068.312
OII	2.00	195710.375
OII	2.00	203942.187
OII	30.00	207002.500
OII	12.00	208484.187
OII	10.00	211712.625
OII	10.00	212593.187
OII	6.00	214229.437
OII	26.00	228746.875
OII	28.00	231530.250
OII	52.00	232959.250
OII	6.00	233430.062
OII	10.00	234454.437
OII	12.00	238892.937
OII	6.00	240516.250
OII	26.00	246028.937
OII	16.00	248185.250
OII	6.00	250251.000
OII	14.00	251220.875
OII	18.00	252607.687
OII	16.00	253048.312
OII	52.00	254590.687
OII	128.00	255977.500
OII	42.00	256136.187
OII	28.00	259286.187
OII	42.00	261180.000
OII	172.00	265925.000
OII	18.00	267783.375
OII	24.00	274782.375
OII	50.00	275841.250
OII	38.00	276109.062
OII	10.00	278140.000
47	0++	0 0 1 0 0 0 0 -2 0 0 15.99830 01 25 0.0
0++ HEAT FORMATION 0.118260E 04		
OIII	1.00	0.0



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49      H-      0 0 0 1 0 0 0 1 0 0 1.00852 01 16 0.0
H-      HEAT FORMATION .344000E 02
      H-      1.00      0.0

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H-	3.00	159850.312					
H-	1.00	166271.687					
H-	12.00	169081.062					
H-	3.00	183231.062					
H-	1.00	184859.000					
H-	9.00	185558.875					
H-	23.00	186095.875					
H-	4.00	190292.437					
H-	60.00	191446.562					
H-	13.00	193795.062					
H-	51.00	193915.750					
H-	68.00	195256.812					
H-	188.00	196950.312					
H-	340.00	197398.562					
H-	78.00	198056.500					
50	H+	0 0 0 1 0 0 0 -1 0 0	1.00742	05	0 0.0		
H+	HEAT FORMATION	.365030E 03					
51	AR+	0 0 0 0 1 0 0 -1 0 0	39.94745	01	28 0.0		
A+	HEAT FORMATION	0.363345E 03					
AII		4.00	0.0				
AII		2.00	1433.090				
AII		2.00	108809.062				
AII		26.00	132432.125				
AII		6.00	135198.250				
AII		6.00	138353.500				
AII		18.00	142301.125				
AII		10.00	143225.437				
AII		6.00	145783.062				
AII		12.00	147993.187				
AII		18.00	149299.750				
AII		16.00	151211.000				
AII		12.00	155165.562				
AII		14.00	157361.750				
AII		12.00	158856.000				
AII		6.00	159523.187				
AII		4.00	160371.125				
AII		6.00	161177.312				
AII		2.00	167439.250				
AII		14.00	170664.062				
AII		16.00	172471.062				
AII		10.00	173534.375				
AII		6.00	174549.062				
AII		12.00	181742.312				
AII		24.00	183820.375				
AII		10.00	184133.125				
AII		18.00	185238.187				
AII		32.00	187038.250				
52	AR++	0 0 0 0 1 0 0 -2 0 0	39.94690	01	41 0.0		
A++	HEAT FORMATION	0.100028E 04					
AIII		5.00	0.0				
AIII		3.00	1112.400				
AIII		1.00	1570.200				
AIII		5.00	14010.000				
AIII		1.00	33267.000				
AIII		5.00	113800.687				
AIII		3.00	114797.562				
AIII		1.00	115328.375				
AIII		3.00	144023.000				
AIII		25.00	144907.000				

AIII	15.00	156917.562
AIII	5.00	174375.000
AIII	3.00	180679.000
AIII	9.00	186402.125
AIII	12.00	186657.187
AIII	8.00	187823.000
AIII	16.00	188714.000
AIII	15.00	196679.750
AIII	18.00	204797.312
AIII	9.00	207233.062
AIII	9.00	209151.812
AIII	7.00	210212.250
AIII	8.00	211004.812
AIII	9.00	213950.812
AIII	15.00	225402.562
AIII	21.00	226646.000
AIII	9.00	231341.750
AIII	18.00	240291.625
AIII	9.00	243424.937
AIII	25.00	246046.562
AIII	5.00	250712.250
AIII	18.00	252289.000
AIII	12.00	266877.500
AIII	36.00	267895.812
AIII	15.00	269000.750
AIII	9.00	271507.875
AIII	18.00	272250.875
AIII	22.00	281473.812
AIII	8.00	282099.125
AIII	15.00	283919.750
AIII	9.00	286009.187
53	AR+++	0 0 0 0 1 0 0-3 0 0 39.94635 01 27 0.0
A+++HEAT FORMATION 0.194349E 04		
ARIV	4.00	0.0
ARIV	4.00	21090.000
ARIV	6.00	21219.000
ARIV	2.00	34854.000
ARIV	4.00	35035.000
ARIV	6.00	117564.000
ARIV	4.00	118515.000
ARIV	2.00	119044.000
ARIV	10.00	146000.000
ARIV	4.00	166356.000
ARIV	2.00	167444.000
ARIV	2.00	177833.000
ARIV	6.00	250906.562
ARIV	6.00	251972.000
ARIV	2.00	256093.250
ARIV	4.00	257348.875
ARIV	10.00	268151.375
ARIV	2.00	285960.125
ARIV	10.00	286751.625
ARIV	8.00	287555.812
ARIV	12.00	289834.625
ARIV	4.00	290256.437
ARIV	10.00	291667.687
ARIV	6.00	295806.750
ARIV	2.00	299563.187
ARIV	14.00	304399.875

ARIV		10.00	306236.250						
54	NE+	0 0 0 0 0 1 0-1 0 0	20.18245 01	26	0.0				
NE+ HEAT FORMATION		.497186E 03							
NEII		4.00	0.0						
NEII		2.00	782.000						
NEII		2.00	217050.000						
NEII		6.00	219133.000						
NEII		4.00	219650.750						
NEII		2.00	219949.875						
NEII		6.00	224089.250						
NEII		22.00	246396.500						
NEII		20.00	249110.750						
NEII		10.00	251013.250						
NEII		6.00	252956.000						
NEII		6.00	254167.000						
NEII		14.00	274411.250						
NEII		8.00	276278.562						
NEII		10.00	277346.062						
NEII		20.00	279139.062						
NEII		52.00	280174.375						
NEII		18.00	281173.500						
NEII		12.00	282000.000						
NEII		6.00	283323.687						
NEII		46.00	302905.750						
NEII		64.00	303475.687						
NEII		30.00	305366.187						
NEII		20.00	306689.750						
NEII		16.00	307992.187						
NEII		10.00	327954.687						
55	NE++	0 0 0 0 0 1 0-2 0 0	20.18190 01	25	0.0				
NE++ HEAT FORM		.144435E 04							
NEIII		5.00	0.0						
NEIII		3.00	650.000						
NEIII		1.00	927.000						
NEIII		5.00	25841.000						
NEIII		1.00	55747.000						
NEIII		5.00	204292.000						
NEIII		3.00	204879.000						
NEIII		1.00	205204.000						
NEIII		3.00	289479.000						
NEIII		5.00	314148.000						
NEIII		3.00	319444.875						
NEIII		15.00	352745.875						
NEIII		15.00	353148.000						
NEIII		9.00	356776.500						
NEIII		5.00	357930.000						
NEIII		9.00	374434.000						
NEIII		3.00	379834.000						
NEIII		8.00	389069.312						
NEIII		7.00	389139.000						
NEIII		21.00	391450.250						
NEIII		45.00	398946.937						
NEIII		4.00	399082.562						
NEIII		15.00	409847.500						
NEIII		12.00	412320.187						
NEIII		75.00	436561.312						
56	NE+++	0 0 0 0 0 1 0-3 0 0	20.18135 01	29	0.0				
NE+++ HEAT FORM		.291405E 04							
NEIV		4.00	0.0						

NEIV	10.00	40950.000							
NEIV	6.00	62167.000							
NEIV	6.00	183860.000							
NEIV	4.00	184477.000							
NEIV	2.00	184799.000							
NEIV	10.00	253807.000							
NEIV	2.00	299351.000							
NEIV	6.00	319751.000							
NEIV	12.00	479651.000							
NEIV	6.00	484623.000							
NEIV	6.00	488917.000							
NEIV	10.00	511411.000							
NEIV	16.00	525017.000							
NEIV	2.00	551712.000							
NEIV	34.00	579375.000							
NEIV	22.00	580095.000							
NEIV	24.00	605417.000							
NEIV	8.00	612781.000							
NEIV	18.00	634413.000							
NEIV	36.00	642934.000							
NEIV	10.00	664124.000							
NEIV	56.00	672799.000							
NEIV	36.00	697855.000							
NEIV	0.0	697855.000							
NEIV	16.00	709460.000							
NEIV	10.00	724690.000							
NEIV	14.00	740607.000							
NEIV	10.00	754597.000							
57	HE+	0 0 0 0 0 0 1-1 0 0	4.00205	01	5	0.0			
	HE+ HEAT FORMATION	.566930E+03							
	HEII	2.00 0.0							
	HEII	8.00 329184.937							
	HEII	18.00 390142.875							
	HEII	32.00 411478.000							
	HEII	50.00 421353.062							
58	HE++	0 0 0 0 0 0 1-2 0 0	4.00150	05	0	0.0			
	HE++HEAT FORMATION	.182102E+04							
59	C	1 0 0 0 0 0 0 0 0	12.01115	01	41	0.0			
	C HEAT FORMATION	0.169990E 03							
	CI	1.00 0.0							
	CI	3.00 16.400							
	CI	5.00 43.500							
	CI	5.00 10193.699							
	CI	1.00 21648.398							
	CI	5.00 33735.199							
	CI	4.00 60353.000							
	CI	8.00 60393.520							
	CI	15.00 64088.559							
	CI	3.00 68858.000							
	CI	15.00 69744.375							
	CI	3.00 70744.250							
	CI	9.00 71385.687							
	CI	5.00 72611.000							
	CI	1.00 73976.187							
	CI	9.00 75256.250							
	CI	5.00 77680.500							
	CI	14.00 78199.312							
	CI	31.00 78250.187							
	CI	3.00 78338.000							

CI	7.00	78531.000
CI	3.00	78727.875
CI	9.00	79311.062
CI	18.00	80222.687
CI	12.00	81344.437
CI	6.00	81770.312
CI	51.00	83761.000
CI	30.00	84986.187
CI	11.00	86187.000
CI	46.00	86326.875
CI	12.00	86504.000
CI	26.00	87713.000
CI	37.00	87807.000
CI	21.00	88547.000
CI	34.00	88624.000
CI	21.00	89081.000
CI	31.00	89155.000
CI	21.00	89450.000
CI	22.00	89517.000
CI	15.00	89779.000
CI	15.00	89968.375
60	N	0 1 0 0 0 0 0 0 0 0 14.0067 01 51 0.0
N	HEAT FORMATION	0.112507E 03
NI	4.00	0.0
NI	10.00	19223.000
NI	6.00	28840.000
NI	12.00	83366.000
NI	6.00	86223.187
NI	12.00	88109.500
NI	2.00	93582.250
NI	20.00	94883.062
NI	12.00	95533.187
NI	14.00	96864.187
NI	6.00	97805.750
NI	10.00	99665.000
NI	12.00	103736.750
NI	6.00	104227.375
NI	34.00	104767.000
NI	32.00	104882.687
NI	24.00	105020.000
NI	2.00	106478.562
NI	6.00	106780.062
NI	20.00	106870.687
NI	10.00	107039.000
NI	12.00	109927.875
NI	2.00	110029.187
NI	8.00	110196.000
NI	32.00	110248.000
NI	38.00	110304.000
NI	16.00	110470.500
NI	10.00	110545.750
NI	6.00	112320.750
NI	2.00	112565.875
NI	10.00	112682.562
NI	20.00	112799.000
NI	60.00	112862.000
NI	16.00	112947.500
NI	6.00	114072.000
NI	62.00	114160.000

[illegible]

AI	1.00	0.0
AI	5.00	93143.750
AI	3.00	93750.625
AI	1.00	94553.687
AI	3.00	95399.812
AI	3.00	104102.125
AI	12.00	105462.750
AI	8.00	106237.562
AI	13.00	107289.687
AI	4.00	111818.062
AI	21.00	112750.187
AI	20.00	113716.562
AI	8.00	114641.000
AI	19.00	114821.937
AI	15.00	116942.812
AI	8.00	117183.625
AI	1.00	117563.000
AI	15.00	118469.062
AI	22.00	119023.687
AI	5.00	119444.875
AI	12.00	119566.062
AI	6.00	119847.750
AI	8.00	120188.625
AI	48.00	120207.312
AI	10.00	120619.062
AI	7.00	120753.500
AI	7.00	121068.750
AI	23.00	121165.375
AI	30.00	121653.375
AI	17.00	122036.125
AI	12.00	122160.187
AI	12.00	122329.687
AI	6.00	122514.250
AI	68.00	122695.687
AI	23.00	123205.812
AI	9.00	123372.937
AI	13.00	123557.437
AI	16.00	123653.187
AI	35.00	123832.500
AI	84.00	124046.625
AI	23.00	124349.000
AI	5.00	124554.937
AI	37.00	124609.875
AI	75.00	124860.625
AI	28.00	125054.062
AI	17.00	125150.000
AI	39.00	125219.875
AI	68.00	125388.625
AI	53.00	125482.687
AI	36.00	125631.687
AI	78.00	125750.375
AI	61.00	125922.500
AI	76.00	126135.375
AI	80.00	126295.750
AI	84.00	126517.375
AI	12.00	127130.000
AI	12.00	127410.000
AI	24.00	127760.000
AI	16.00	127970.000



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64      NE      0 0 0 0 0 1 0 0 0 0 20.18300 01 25 0.0
NE HEAT FORMATION .000000E+00
NE      1.00      0.0
NE      5.00     134043.750
NE      3.00     134461.187
NE      1.00     134820.562
NE      3.00     135890.625
NE      3.00     148259.687
NE     12.00     149659.000
NE      8.00     150317.812
NE      9.00     150860.437
NE      3.00     151040.375
NE      1.00     152972.687
NE      8.00     158603.062
NE      4.00     159536.562
NE     40.00     161592.250
NE     35.00     162832.625
NE     20.00     163710.562
NE      9.00     165830.125
NE      8.00     166977.312
NE     148.00     167062.500
NE     20.00     168926.625
NE     124.00     169536.250
NE     148.00     170879.937
NE     216.00     171689.937
NE     428.00     172575.812
NE     188.00     173280.000
65      HE      0 0 0 0 0 0 1 0 0 0 4.00260001 16 0.0
HE HEAT FORMATION .000000E 00
HE      1.00      0.0
HE      3.00     159850.312
HE      1.00     166271.687
HE     12.00     169081.062
HE      3.00     183231.062
HE      1.00     184859.000
HE      9.00     185558.875
HE     23.00     186095.875
HE      4.00     190292.437
HE     60.00     191446.562
HE     13.00     193795.062
HE     51.00     193915.750
HE     68.00     195256.812
HE     188.00     196950.312
HE     340.00     197398.562
HE      78.00     198056.500
66      E-      0 0 0 0 0 0 0 1 0 0 0.00054905 0 0.0
E- HEAT FORMATION 0.      E 00

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## Nomenclature

$B_i$	concentration of the $i$ th element at STP, moles/l.	$P_i$	partial pressure of the $i$ th specie.
$C_i$	quantity defined in Eq. (46).	$R$	universal gas constant, 1.98726 cal/mole — °K.
$D_i$	constant defined with Eqs. (44) and (45).	$S$	entropy per unit mass.
$f_i^o$	partial molal free energy of the $i$ th specie evaluated at a pressure of one atmosphere, cal/mole.	$T$	temperature, °K.
$g_i^{(k)}$	degeneracy of the $i$ th energy level of the $k$ th specie.	$U$	velocity.
$H$	enthalpy per unit mass.	$v$	$= 1/\rho$ specific volume of gas, l/g
$h_i^o$	partial molal enthalpy of the $i$ th specie evaluated at a pressure of one atmosphere, cal/mole.	$X^{(j)}$	the chemical symbol for the $j$ th numbered chemical element.
$h_{oi}^o$	energy of formation for one mole of the $i$ th specie at a pressure of one atmosphere and a temperature of 0°K, cal/mole.	$Y^{(i)}$	the chemical symbol for the $i$ th numbered compound.
$K$	thermal conductivity.	$Y_{ij}$	matrix coefficient defined in Eq. (46).
$M_o$	molecular weight at standard conditions (cold molecular weight).	$Z$	$= \frac{M_o}{M}$
$N_i^{(k)}$	number of particles of the $k$ th constituent.	$\alpha_{ij}$	the number of atoms of the $j$ th numbered element in the molecule of the $i$ th numbered specie.
$n_i$	concentration of the $i$ th specie, mole/l.	$\epsilon_i^{(k)}$	energy of the $i$ th level for the $k$ th specie.
$N_i^{(k)}$	number of particles of the $k$ th constituent in the $i$ th energy level.	$\eta$	viscosity.
$P$	total pressure of mixture.	$\rho$	mass density.
		$\rho_o$	density of mixture at STP.
		$( )_o$	property of mixture at STP.

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