

A Significant Role of MoO3 on the Optical, Thermal, and Radiation Shielding Characteristics of B2O3-P2O5-Li2O Glasses

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A Significant Role of MoO₃ on the Optical, Thermal, and Radiation Shielding

Characteristics of B2O3-P2O5 -Li2O Glasses

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ABSTRACT

Glasses based on borophosphate with the formula $42.5P_2O_5 - 42.5B_2O_3 - (15-x) Li_2O - xMoO_3$

mol% where $x = (0 \le x \ge 15)$ were manufactured using the melt-quenching methodology.

The status of prepared samples was identified by (XRD). The temperature of the glass transition

T_g, the temperature of onset glass crystallisation T_c and the temperature of the crystallisation

 T_p were evaluated using a differential thermal analyser (DTA). The energy gap (E_{opt}) , Urbach

 (E_u) , and parameters of dispersion were calculated through the data of optical spectra. Physical

properties were determined and calculated, such as molar refractivity, metallization, electron

polarizability, electronegativity, loss of reflection and dispersion parameters. Raising MoO₃ at

the expense of Li₂O was used to assess the level of protection. For radiation protection

applications, the glasses under investigation had superior characteristics.

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1

1. Introduction

B₂O₃-P₂O₅ glasses with superior efficiency can be used in a variety of different settings. Solid-state batteries, and nonlinear optics borophosphate glasses are appropriate. Due to its obvious advantages, lithium borophosphate is classical glass that has become recognized in storage batteries. These glasses are used as storage batteries in optical and electronic instruments. The addition of modifiers such as Li₂O influences of these characteristics. Li₂O will be combined instead of B-change BO₄ to BO₃ [1-8]. The characteristics of the combined glass (B₂O₃+P₂O₅) networks vary from those of the single glass B₂O₃ and P₂O₅ networks.

Transition metal oxides (TMOs) are a fascinating group of semiconductor materials because of their technological advantages for use in microelectronics and display systems. MoO₃ among (TMOs), due to its excellent use in optical materials and electrochemical devices has received increasing attention in recent years. There is a goal that support in manufacturing these glasses regarding the MoO₃ used in these implementations. Different MoO₃ preparation glasses were developed and investigated in response to this broad range of applications [9-16].

Due to the existence of MoO₄ and MoO₆ in the glass network, MoO₃ appears as the former non-conventional network. The existence of MoO₃ in glass systems does have a modifier impact on UV spectra [9-16]. There are increasingly diverse technologies for molybdenum borophosphate-based glasses, such as laser host fibers and superconducting switches.

Both scientifically and technologically, the significant advances of alkaline borophosphate glasses are considerable. The existence of PO₄ & BO₄ structural units, this consequence approaches from the structural issues connected with covalent links. In several fields, B₂O₃-P₂O₅ - Li₂O - MoO₃ glasses possess high applications because of their radiation shielding and good FT-IR transmission [17-33]. The object of this study is to assist you in the

preparation of B₂O₃-P₂O₅ - Li₂O - MoO₃ glasses and investigating their optical and neutron shielding using Phy-X/PSD [**34**] properties.

2. Methodology

Table 1 shows how we formed the glasses in our published articles using the melt-quenching methodology. By melting together of B₂O₃ in its H₃BO₄ (Merck), Li₂O in it Li₂CO₃ (Aldrich), P₂O₅ in it (NH₄)₂HPO₄ (Merck) and MoO₃ (Merck) in an open ceramic crucible. With the evaporation of CO₂, NH₃ and H₂O, Li₂CO₃, (NH₄)₂HPO₄ and H₃BO₄ and are converted into Li₂O, P₂O₅ and B₂O₃. The furnace temperature was changed at a melting temperature of 1050 °C. At 350 °C the prepared samples are annealed.

The Philips X-ray diffractometer (model PW / 1710) checked the condition of these glasses and glass-ceramics. The spectrophotometer was used to measure the optical spectra of the investigated glass system (type JASCO V- 670). The thermal investigation was carried out with a DTA-50 (Shimadzu-Japan). Phy-X / PSD can calculate a variety of shielding considerations [34]. Electron density (N_{eff}) was predictable as: $N_{eff} = N \frac{Z_{eff}}{\sum_i F_i A_i}$,. Effective cross-section of removal (Σ_R) projected as: $\left(\frac{\Sigma_R}{\rho}\right) = \sum_i w_i \left(\frac{\Sigma_R}{\rho}\right)_i$ and $R = \sum_i \rho_i \left(\frac{R}{\rho}\right)_i$, G-P fitting parameters has been predictable as $P = \frac{P1(\log Z2 - \log Zeq) + Z2(\log Zeq - \log Z1)}{\log Z2 - \log Z1}$. EABF and EBF were predictable using G-P fitting $B(E,X) = 1 + \frac{b-1}{K-1}(K^X - 1)$ for $K \neq 1$, B(E,X) = 1 + (b-1)x K= 1 where $K(E,X) = cx^a + d\frac{tanh(\frac{x}{Xk} - 2) - tanh(-2)}{1 - tanh(-2)}$

3. Results and Discussion

3.1. XRD

Figure 1 depicts the glass system's X-ray features. In the glass samples, XRD revealed no discrete lines or sharp peaks, indicating a high degree of glass status.

3.2 Optical spectra

The absorption (A), transmittance (T), and reflectance (R) of glass samples are shown in Figures 2 and 3. Spectral UV according to reports, increasing. As a result, MoO₃ is to blame for the slight increase in BO [35-44]. Figure 4 shows the absorption coefficient of the glasses.

3.2.1 band gap E_{opt}

Glass spectrum in the UV and VIS areas were used for the estimated of the band gap energy E_{opt} is estimated by $(\alpha.hv)^{1/2} = B(hv - E_{opt})$ where B is an energy independent constant and hv is photon energy. By intrigue the $(\alpha.hv)^{1/2}$ versus hv as Fig.5. Plot of $(\alpha.hv)^{1/2}$ against photon energy (hv) to evaluate the indirect E_{opt} from the intercept. E_{opt} increases with increasing MoO₃, as revealed in Table 2, due to oxygen bridges (BO) form and bind excited electrons more tightly than non-bridging oxygen electrons (NBO). Urbach energy has been calculated $\propto_0 exp\left(\frac{hv}{E_u}\right)$, Fig. 6 and table 2 show that their E_{opt} values have an inverse relationship. The values of E_{opt} and E_u was shown in Fig.7.

The refractive index was calculated using: $n = \frac{(1-R)^2 + k^2}{(1+R)^2 + k^2}$ where $k = \alpha \lambda/4\pi$. The refractive index presented in Fig.8 for fabricated glasses. It has already been stated that density increase, the refractive index of these samples increased. As a result, it can be directly compared to reflectance and density, and opposite to molar volume.

3.2.2 Dispersion parameters

As approximated, molar polarization and polarizability of glasses were computed:

$$R_m = \langle n^2 - 1 | n^2 + 2 \rangle V m$$
, $\alpha_m = (3|4\pi N) R_m$, and $\alpha_0^{2-} = \frac{\left[\frac{Vm}{2.52} \left(\frac{n^2 - 1}{n^2 + 2}\right) - \sum \alpha_{cat}\right]}{N_o^{2-}}$ [45-49].

Polarization was linked with the optical basicity; $\Lambda = 1.67 \left(1 - \frac{1}{\alpha_0^2}\right)$. Figs. 9,10& 11 exemplifies the polarizabilities, Molar Polarizability α_m and optical basicity separately of the

samples. The refractive index is trending in the same direction with MoO₃ content has been reported to enhance.

The molar refractivity as $E_{opt.}$ $R_m = Vm \left(1 - \sqrt{Eg/20}\right)$ and molar polarizability $(\alpha_m) \alpha_m = \left(\frac{3}{4\pi N}\right) R_m$. Reflection loss $R_L = \left(\frac{R_m}{Vm}\right)$. Because the molar volume decreases with Mo⁺, these values of (R_m) (α_m) and (R_L) decline. The criterion for metallization is predicted as $M = 1 - \frac{R_m}{Vm}$, the metallization value increase with Mo⁺. The electronegativity (χ) is predicted as $\chi = 0.2688 E_{opt.}$ where $E_{opt.}$ bandgap. Thus, with Mo⁺ increasing, the electronegativity (χ) values increase. The electron polarizability is predicted as, $\alpha = -0.9 \chi + 3.5$ and optical basicity $\alpha = -0.5 \chi + 1.7$. $\alpha = 0.5 \chi + 1.7$. $\alpha = 0.5$

The dispersion was calculated by Wemple and Didomenico E_0 and E_d [50-53]. The hypnosis designates explained by $n^2-1=\frac{E_0}{E_0^2-E^2}$ [64-70]. The plotting of, $(n^2-1)^{-1}$ with $(h\upsilon)^2\,E_0$ and E_d predictable from the slope and intercept as shown in Figs. 12 & 13. It mentioned that with increasing MoO₃, E_0 and E_d were slightly enhanced. The optical energy E_{opt} that represent $Eopt=\frac{E_d}{2}$. Refractive Static index at infinite wavelength (n_0) was estimated by $n_0=\sqrt{1+\frac{E_d}{E_0}}$ and the static dielectric $\varepsilon_\infty=n_0^2$. The oscillator's wavelength (λo) and strength (S_0) were calculated using the following formula $n^2-1=\frac{S_0\,\lambda_0^2}{1-(\frac{\lambda_0}{\lambda})^2}$. These items are obtainable in Table 3.

3.3 DTA

The thermal analysis (DTA) of glass samples demonstrated in Fig. 14. The temperature of the glass transition, Tg, is 493-532 ± 3 °C. The temperature of the glass crystallisation T_c starts at 537-580 ± 3 °C. The temperature of the glass crystallisation T_c ends at 606-645 ± 3 °C. According to DTA observations, Tg increases from 493 into 532 °C, Tc increases from 537

into 580 °C and Tp increases from 606 into 645 °C with the increase of MoO₃ content. The transformation of Li-O to Mo-O linkages is significantly associated with this behaviour. Hence, Li-Li (137.3±6.3 KJ/mol) dissociation energy is weaker than Mo-Mo (449.4±1 KJ/mol) dissociation energy and adding MoO₃ variations the basic B units from BO₃ to BO₄. Thermal stability estimated by $\Delta T = (Tc - Tg)$, weighted thermal stability $Hg = \Delta T/Tg$ and $S = (Tp - Tc) \Delta T/Tg$. It observed that all thermal stability of samples improved as MoO₃. The T_g , T_c , T_p and thermal stability values are obtainable in Table 4.

3.4 Photon Shielding Features

The level of protection was assessed in this article by increasing MoO₃ at the expense of Li₂O with the composition $42.5B_2O_3 - 42.5P_2O_5 - (15 - x)Li_2O - xMoO_3$, $(0 \le x \ge 15)$. The mean free path (MFP) is depicted in Fig.15. It was stated that as photon energy increases, the values of (MFP) increase. This insight revealed that as the photon's energy increases, it becomes capable of transmitting samples on purpose. Because the lower value of the (MFP) sample has a higher MoO3 content, good radiation attenuation glasses are available. [54-71].

Figure 16 demonstrations the (N_{eff}) of fabricated glasses. It is demonstrated that (N_{eff}) decreases and then rises as energy increases. This significant decrease can be attributed to the interaction of Compton scattering. The effect of forming pairs at higher energy levels as MoO₃ is linked to the increasing in (N_{eff}) .

The ASC of fabricated glasses are presented in Fig. 17. The ASC and ESC values are expected to decrease as energy rates increase. This decrease occurs due to the Compton scattering interaction. The C_{eff} of fabricated glasses depicted in Figure 18. With the increase in photon energy, it is predicted that C_{eff} will decrease. The impact of pair-creation was reflected in the increase in C_{eff} .

The EBF & EABF of fabricated glasses were characterized by Figs. 19&20. EBF and EABF values are determined by the lower energy and concentration of the glass samples. At lower energy levels, EBF and EABF values are low, but they rise as energy levels rise. After that, gradually decrease as energy level rises. So, we can divide the energy scale into three parts low, medium, and high. The first part (low energy): the typical phase is the photoelectric effect, and the relationship is reversed with light, and the glasses will absorb the energy photons. The photons are therefore not allowed to build-up. In the second part (medium energy): the common process is the Compton scattering, the values of EBF and, EABF is increased in all samples independent from the MFP. Through this part, the photons stay in the samples for a long time, as these processes lead to high accumulation value due to multiple scattering processes. Third parts (high energy): the common process is the pair production. In this process, EBF and, EABF is decreased with energy. Therefore, these data helped in the determination of maximum radiation intensity occur. In this research, highest radiation occurs on the surface of the sample. In Fig. 21, fast cross section neutron removal (FNRCS) is shown. It was stated that MoO3 improved FNRCS.

4. Conclusions

In the existing research, molybdenum lithium borophosphate glasses $42.5P_2O_5 - 42.5B_2O_3 - (15-x)$ Li₂O - xMoO₃ where $x = (0 \le x \ge 15)$ were fabricated with conventional melt-quenching procedures. The optical, thermal, and shielding factors were observed. The findings showed the following objects:

- 1- Because of the increase in MoO₃, the metallization of these glasses was improved.
- 2- The 2.23 for G 1, 2.32 for G 2, 2.38 for G 3, 2.41 for G 4, and 2.49 for G 5 were identified as the indirect optical bands that were collected.
- 3- Urbach energies of these samples were reduces as the content of MoO₃ increase.
- 4- As the density of the investigated glasses rises, the refractive index rises as well.

- 5- These glasses were investigated for molar polarization, polarizability, and optical basicity.
- 6- T_g , T_c , T_p and thermal stability values are enhanced with MoO₃.
- 7- The fabricated samples' gamma shielding features were predictable. The impact of adding MoO₃ to the glasses on their shielding ability was mentioned.
- 8- Lower value of the (MFP) sample has more MoO₃ are good radiation attenuation glasses are available.
- 9- As the concentration of MoO₃ increased, these glasses have a high cross section neutron removal rate.

The findings discovered that as MoO₃ increase the glass system can result in significant improvements in attenuation and optical characteristics. Furthermore, it is possible to use this glass in optoelectronic, optical devices, and a radiation shield for use in x-ray centers.

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Figures

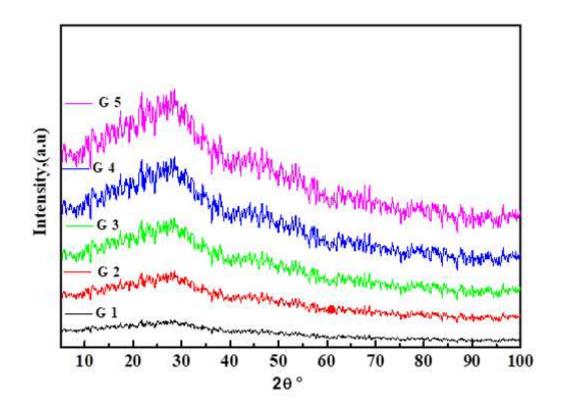


Figure 1

XRD of the studied glasses.

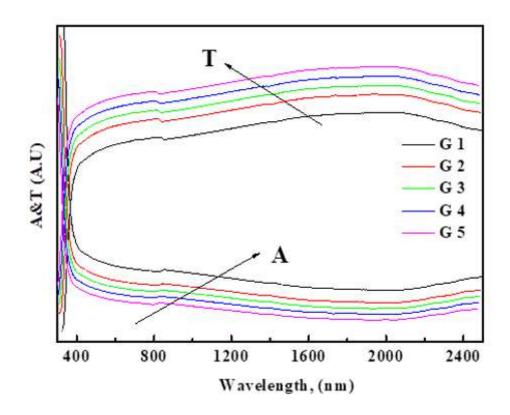


Figure 2

The absorbance (A) and Transmittance (T) of the prepared glasses.

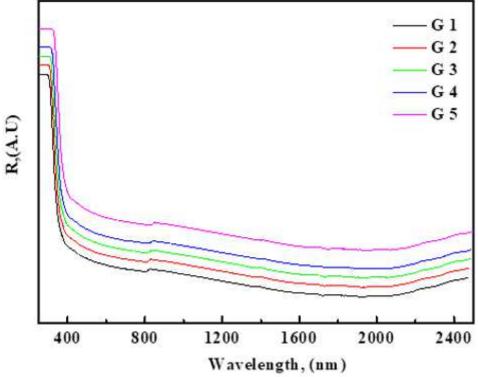


Figure 3

The reflectance (R) of the prepared glasses.

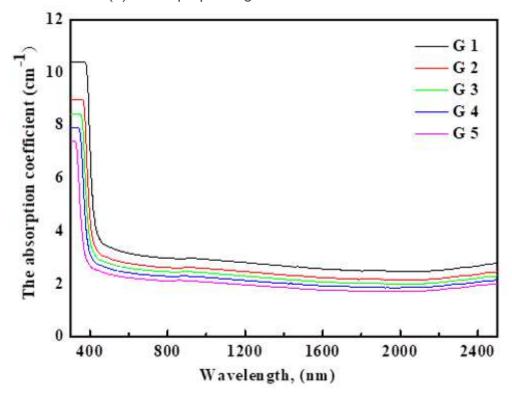


Figure 4

The absorption coefficient of the prepared glasses.

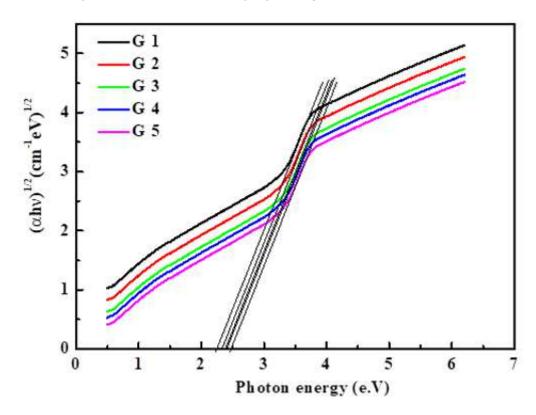


Figure 5

Plot of $(\alpha h \upsilon)1/2$ against photon energy $(h \upsilon)$ to calculate the indirect optical band gap from the intercept of the curves.

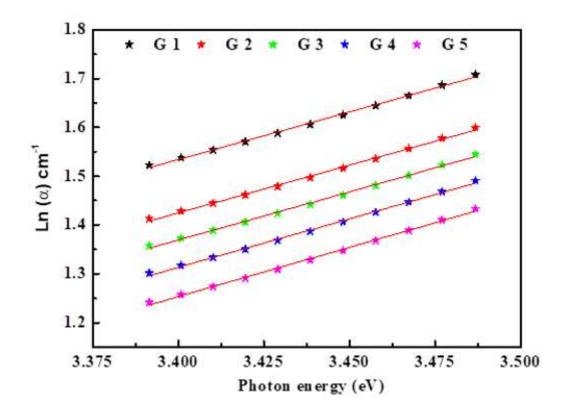


Figure 6 $\label{eq:continuous}$ Dependence of $ln(\alpha)$ upon the photon energy (hu) for the prepared glasses.

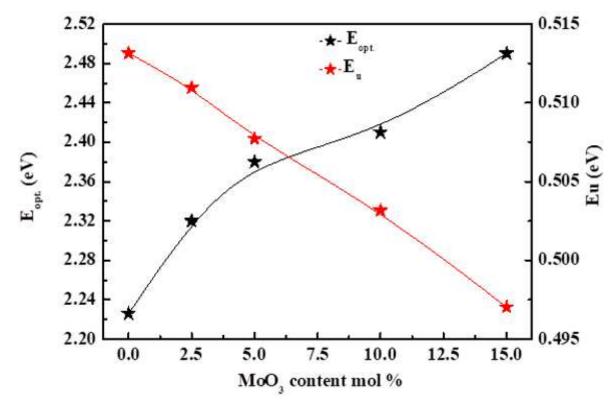


Figure 7

Optical band gap and Urbach energy versus concentration of MoO3.

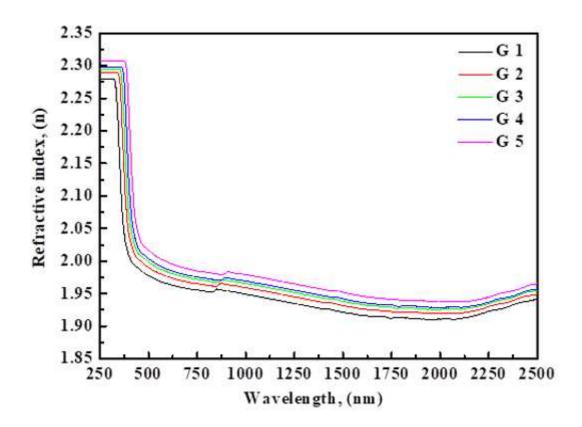


Figure 8

Refractive index of the prepared glasses.

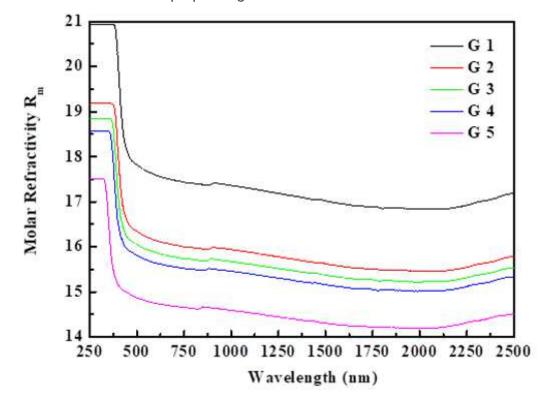


Figure 9

Molar refractivity of the prepared glasses.

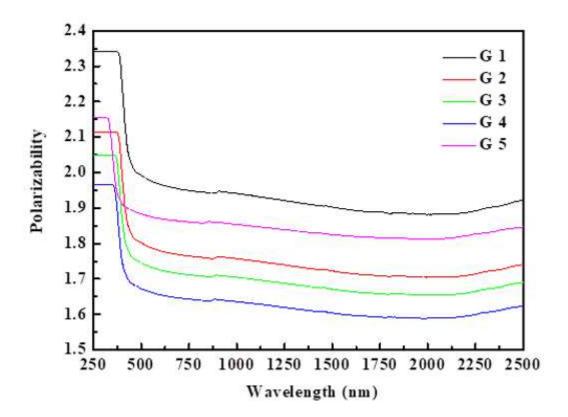


Figure 10

Electronic polarizability of the prepared glasses.

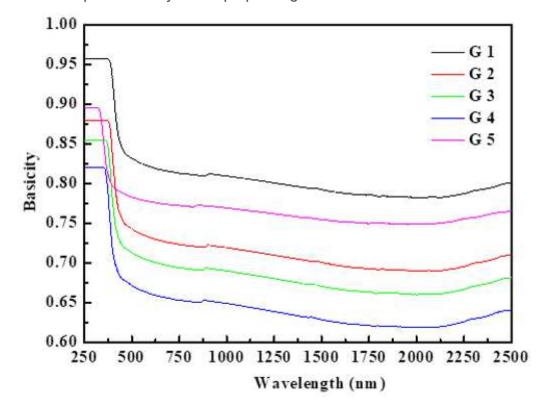


Figure 11

Optical basicity of the prepared glasses.

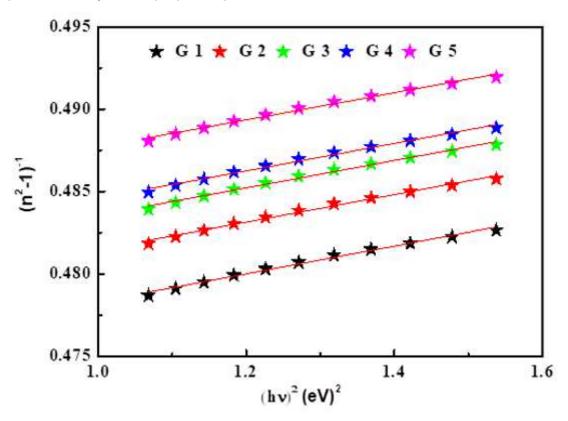


Figure 12 $\label{eq:Variation} \mbox{Variation of (n2-1)-1 with (hu)2 for the prepared glasses.}$

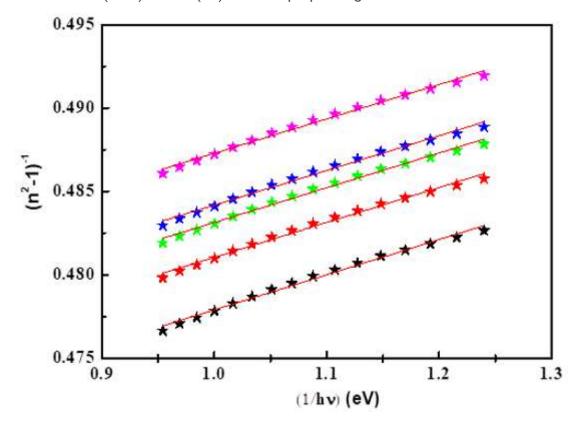


Figure 13 $\label{eq:Variation} \mbox{Variation of (n2-1)-1 with 1/(hu) for the prepared glasses.}$

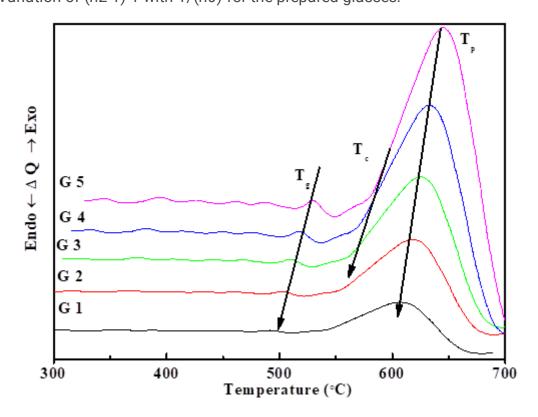


Figure 14

DTA of the prepared glasses.

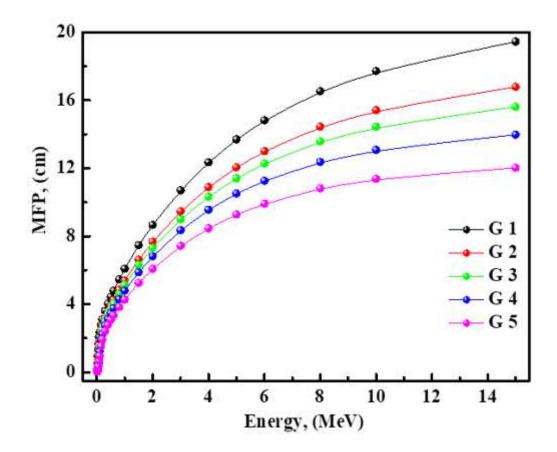


Figure 15

The MFP for the prepared glasses as a function of photon energy.

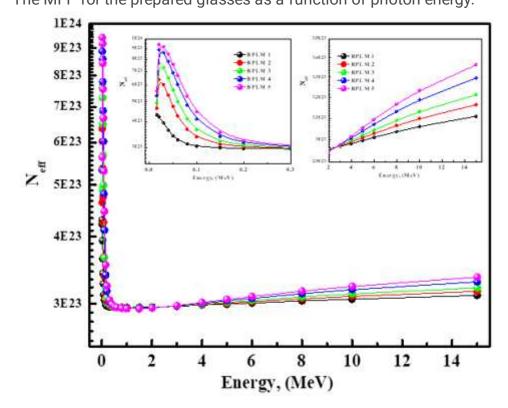


Figure 16

The (Neff) for the prepared glasses as a function of photon energy.

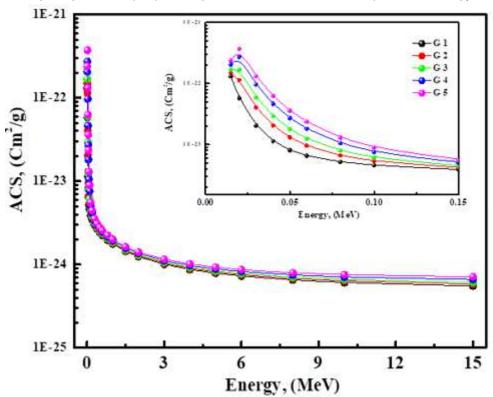


Figure 17

The ASC for the prepared glasses as a function of photon energy

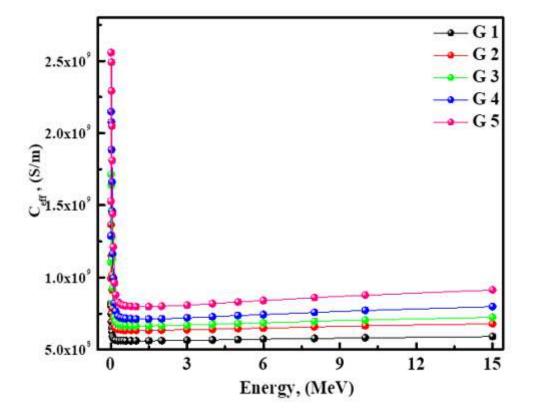


Figure 18

The Ceff for the prepared glasses as a function of photon energy.

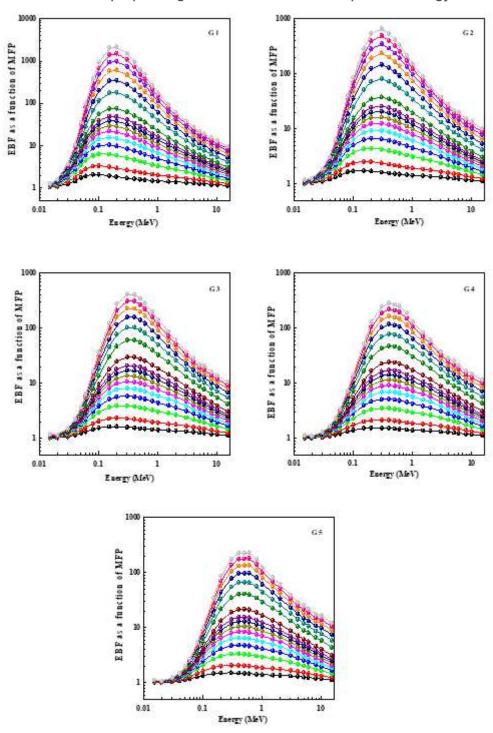


Figure 19

Variation of EBF versus the gamma ray energy for the prepared glasses as a function of photon energy.

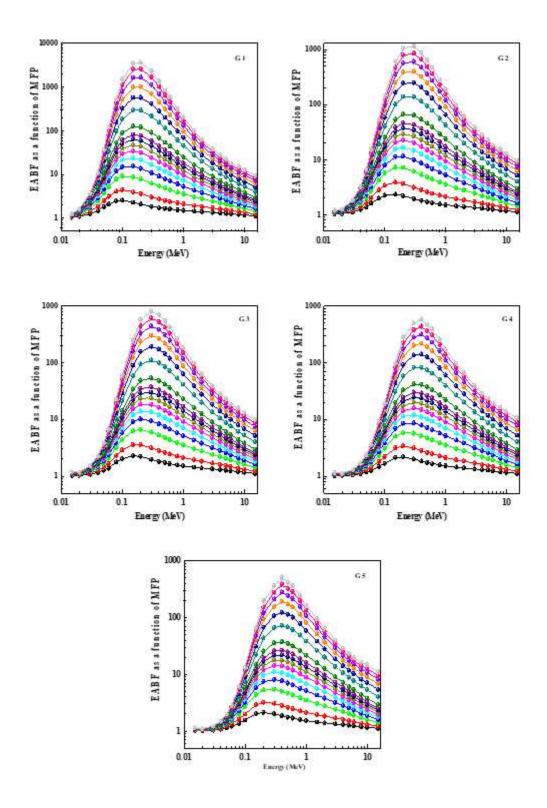


Figure 20

Variation of EABF versus the gamma ray energy for the prepared glasses as a function of photon energy.

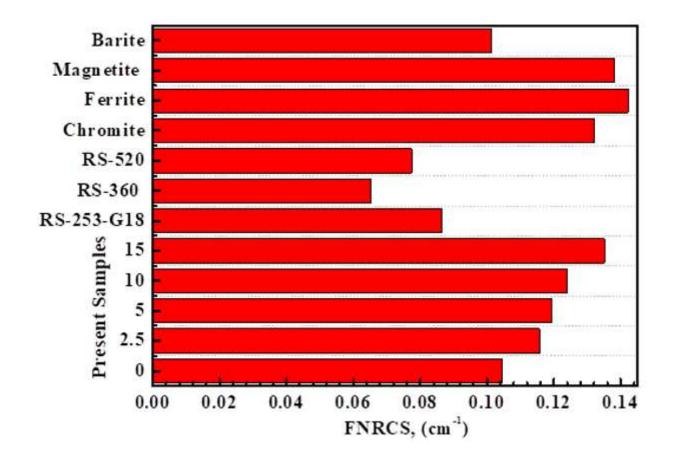


Figure 21

FNRCS for the prepared glasses comparison with standard materials.