

## **A THREE DIMENSIONAL MOVIE OF STRUCTURAL CHANGES IN BACTERIORHODOPSIN**

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**One sentence summary:** Time-resolved serial crystallography at an X-ray free electron laser reveals structural changes in bacteriorhodopsin.

## **ABSTRACT**

Bacteriorhodopsin (bR) is a light-driven proton pump and a model membrane transport protein. Here we use time-resolved serial femtosecond crystallography at an X-ray free electron laser to visualize conformational changes in bR from nanoseconds to milliseconds following photoactivation. An initially twisted retinal chromophore displaces a conserved tryptophan residue of transmembrane helix F on the cytoplasmic side of the protein while dislodging a key water molecule on the extracellular side. The resulting cascade of structural changes throughout the protein show how motions are choreographed as bR transports protons uphill against a transmembrane concentration gradient.

Energy coupled membrane transport proteins are ubiquitous in biology. The basic framework underpinning unidirectional membrane transport is called the *alternating access model* and was proposed by Jardetzky half a century ago (1). This theory postulates that a high-affinity binding site which is initially accessible to one side of the membrane is converted through structural changes into a low-affinity binding site that is accessible to the other side of the membrane (fig. S1). The recent advent of time-resolved serial femtosecond crystallography (TR-SFX) at an X-ray free electron laser (XFEL) (2-4) provides an opportunity to examine this framework.

Bacteriorhodopsin (bR) harvests the energy content of light to drive conformational changes leading to unidirectional proton transport. Energy stored within a transmembrane proton concentration gradient is converted by ATP-synthase into ATP or is coupled to other transport processes. Considerable effort has been made to understand how structural changes in bR transport a proton uphill against a transmembrane potential (5-7). Crystallographic structures derived from data recorded at cryogenic temperature reveal both light and mutation-induced structural changes in bR (6, 7) but this work is controversial because the reported structures show considerable variation (7), structures trapped at cryogenic temperature or by mutation do not correlate directly with time-dependent structural changes in wild-type bR, and it is argued that X-ray induced radiation damage may mislead the interpretation of trapped intermediate structures (8-10). Here we circumvent these concerns by recording a 3D movie of structural

changes in bR at room-temperature at 2.1 Å resolution using TR-SFX (2-4). The principle of “diffraction before destruction” (11) ensures that crystallographic structures are effectively free from the influence of X-ray induced radiation damage (12) using 10 fs long XFEL pulses at the SPring-8 Angstrom Compact Free Electron Laser (SACLA).

### **Structure of the bR resting conformation**

bR is a seven transmembrane  $\alpha$ -helix protein containing a buried *all-trans* retinal chromophore that is covalently attached through a protonated Schiff base (SB) to Lys216 of helix G. Light photo-excites the chromophore which isomerizes with high quantum yield to a 13-*cis* configuration (Fig. 1A) and thereby initiates a sequence of spectral (fig. S2) and structural changes that facilitate spontaneous proton exchange between acidic and basic amino acid residues (Fig. 1B). Electron density for the resting state bR structure (Fig. 1C, table S1) shows how the protonated SB forms a H-bond interaction to Wat402 (13) which is critical for attaining the remarkably high SB proton affinity with a pKa of 13.3 (14). Asp85, the primary proton acceptor, also forms a H-bond interaction with Wat402 and additional H-bond interactions with Thr89 and Wat401 ensure its low resting state pKa of 2.2 (15). A difference of eleven orders of magnitude between the proton affinities of the primary donor and acceptor prevents the leakage of protons from the extracellular (EC) medium to the cytoplasm (CP) but raises the question of what brings these proton affinities close together in order to facilitate spontaneous proton

exchange? It is also puzzling why it takes microseconds for the primary proton transfer to occur when the SB and Asp85 are initially separated by only 4 Å and a water mediated proton exchange pathway between these two groups is seen in the resting state (13). Moreover, protons are pumped from the CP to the EC yet retinal isomerization redirects the SB proton away from the EC and towards the CP (Fig. 1A), which appears to contradict Jardetzky's framework.

### **Time-resolved spectroscopy studies of bR microcrystals**

Photoactivated bR progresses through a sequence of spectral intermediates labeled K, L, M<sub>1</sub> (early M), M<sub>2</sub> (late M), N and O prior to returning to the bR resting state (fig. S2). A proton is transferred from the SB to Asp85 in the L to M transition and from Asp96 to the SB from M to N. Time-resolved difference absorption spectra from bR microcrystals suspended in a lipidic cubic phase (LCP) matrix (fig. S3A) show a photo-cycle turnover that is similar to that of wild-type bR in the purple membrane. Spectral decomposition reveals that for  $\Delta t \leq 10 \mu\text{s}$  the photoactivated population in microcrystals is dominated by the L intermediate but with traces of K, whereas the M-intermediate population increases to 50 % at  $\Delta t = 19 \mu\text{s}$  and falls to 50 % at  $\Delta t = 9 \text{ ms}$  (fig. S3B,C). Measurements from four independent crystal batches give the half-rise for the M-state as  $15 \pm 7 \mu\text{s}$  and the half-decay as  $8 \pm 3 \text{ ms}$ .

### **Overview of structural rearrangements within the bR photocycle**

In TR-SFX a continuous stream of microcrystals are injected across a focused XFEL beam and the delay between sample photoactivation and the arrival of an XFEL pulse is controlled electronically (2-4) (fig. S4). TR-SFX data were recorded to 2.1 Å resolution from light-adapted bR microcrystals following photoactivation by a ns laser pulse for the time-delays  $\Delta t = 16$  ns, 40 ns, 110 ns, 290 ns, 760 ns, 2  $\mu$ s, 5.25  $\mu$ s, 13.8  $\mu$ s, 36.2  $\mu$ s, 95.2  $\mu$ s, 250  $\mu$ s, 657  $\mu$ s and 1.725 ms (table S1), which are evenly spaced on a logarithmic scale. A global overview of the evolution of difference electron density reveals how structural changes first emerge near the active site of the protein and become stronger around Lys216 of helix G and Trp182 of helix F prior to cascading towards the EC side of the protein along helix C (fig. S5; Movie S1).

### **Early structural changes in the bR photocycle**

For  $\Delta t = 16$  ns, which corresponds to essentially pure K-intermediate, paired negative and positive difference electron density features adjacent to the C20 methyl group reveal that the retinal is initially tilted towards helix G in response to photo-isomerization (Fig. 2A). Combined quantum mechanics and molecular mechanics computations also favor a twisted retinal geometry (fig. S6) and indicate that approximately 17 kcal/mol of energy is initially stored within this distorted active site configuration, which equates to 32 % of the absorbed photon's energy and is compatible with earlier estimates (16, 17). By  $\Delta t = 290$  ns the positive difference density feature associated with the C20 methyl is much weaker and moves into the plane of the



retinal whereas paired negative and positive difference electron density features associated with Trp182 grow stronger (Fig. 2B), showing how Trp182 is displaced towards the CP as the retinal straightens (Fig. 2C). Negative and positive difference electron density features also indicate a movement of the side-chain and backbone of Lys216 has begun at  $\Delta t = 16$  ns (Movies S2-3).

A strong negative difference electron density peak visible at  $\Delta t = 16$  ns on Wat402 (Fig. 2A) shows that this water molecule is rapidly disordered by retinal isomerization. This feature doubles in strength as the photocycle evolves (Fig. 3A, Movies S2-3, table S2) from which we infer that the freedom of Wat402 to move is initially constrained by the limited size of the cavity in which it is buried, yet this water becomes increasingly mobile as the retinal is displaced towards the CP. Low-temperature trapping studies have suggested that Wat402 disorders upon retinal photo-isomerization (18) but this conclusion was challenged in light of similar observations due to X-ray induced radiation damage (8-10). Because no effects of radiation damage are visible when using XFEL pulses  $\leq 10$  fs in duration at an X-ray dose of 12 MGy (12, 19) these TR-SFX data bring closure to this debate. It is noteworthy that one trapped K-intermediate showed Wat402 disordering (18) but did not visualize a twist of the retinal C20 methyl towards helix G (Fig. 2A,C) whereas another captured this twist but kept Wat402 in the K-state model (8) and a third K-state structure showed neither change (20).

## Conformational changes associated with proton transfer

The key step in achieving unidirectional proton transport by bR is the primary transfer event from the SB to Asp85 (the L-to-M transition) because only the mutation of Asp85 (21) or the removal of retinal completely stops proton pumping. Our data reveal a smooth evolution of electron density changes on the CP side of helix F and G prior to proton transfer (Fig. 3B, Movies S1, S3). Paired positive and negative density features also show that the side-chain of Leu93 is displaced towards the CP and a weaker positive difference electron density peak arises between the retinal, Leu93 and Thr89 from  $\Delta t = 40$  ns until  $13.8 \mu\text{s}$  (Fig. 4A) and is strongest at  $\Delta t = 760$  ns (Fig. 3C, table S2). We modeled this feature as a transient water molecule (Wat452, Fig. 4B) which shows weak H-bond interactions with the SB ( $3.35 \text{ \AA}$ ) and the backbone carbonyl oxygen of Thr89 ( $3.35 \text{ \AA}$ ). Its crystallographic distance to O $\gamma$  of Thr89 of  $3.95 \text{ \AA}$  may also be classified as a weak H-bond interaction (22). Structural refinement shows that for  $\Delta t = 16$  ns, O $\gamma$  of Thr89 is separated from the SB by  $3.97 \text{ \AA}$  and the SB nitrogen is pointing towards helix G, whereas for  $\Delta t = 760$  ns the SB nitrogen rotates almost  $180^\circ$  towards helix C and reduces the SB-O $\gamma$  separation to  $3.30 \text{ \AA}$ . Fluctuations about these crystallographic position could therefore create a transient proton-transport pathway linking the SB to Asp85 through Wat452 and Thr89 (Fig. 4B).

A water molecule was built at the Wat452 location in one cryo-trapped L-intermediate structure

(23) but was not observed in any other L-state study (24-26). FTIR spectroscopy observations have inferred that a transient H-bond interaction of the SB to a water molecule arises in the L-state (27, 28) and Wat452 is the only plausible candidate for this interaction. Computer simulations also establish that a water molecule at this position can channel a proton to Asp85 through Thr89 (29). Another transient water molecule (Wat453) which forms H-bond interactions with the carbonyl oxygen of Lys216 and the backbone nitrogen of Gly220 was also predicted from FTIR spectroscopy studies (30) and is observed in the difference electron density.

On the EC side of the retinal, the disordering of Wat402 triggers a structural cascade visible from  $\Delta t \geq 13.8 \mu\text{s}$  as significant ( $\geq 4\sigma$ ) difference electron density peaks that show the disordering of Wat400 and Wat401 and the ordering of a new water molecule, Wat451, between Asp85 and Asp212 (Fig. 5, table S2). Concomitant with these active-site water rearrangements is an increase in the strength of paired negative and positive difference density features along the EC portions of helix C (residues 82 to 89, Fig. 3A) which are modeled as a concerted inwards movement of helix C towards helix G (Fig. 5C). These observations are similar to electron density changes recorded in low-temperature trapping studies of the bR L-state (7, 24, 25) but TR-SFX data show that this inwards flex of helix C approaches its maximum amplitude as the SB is deprotonated (*ie.* as the M-state population grows, Fig. 3A). Protons exchange between consecutive water molecules in the model channel gramicidin A on a sub-picosecond

time-scale (31) and the primary proton transfer should be very fast once a structure facilitating this exchange is attained. Consequently the time required for bR to evolve to a conformation with helix C bent towards helix G is the rate limiting step that controls the primary proton transfer and explains why it takes microseconds for the SB to be deprotonated.

### **Breaking the EC connectivity after proton transfer**

A key conceptual issue underpinning membrane transport is the nature of the “switch”, or more specifically how the SB breaks access to the EC side of bR after the primary proton transfer to Asp85. This is central to the mechanism of proton pumping because structural changes must prevent Asp85 from reprotonating the SB or otherwise the photocycle would be futile. A switch may be achieved by breaking the pathway for the reverse proton transfer, or by perturbing the chemical environment of key groups to shift the equilibrium in favor of the forward proton transfer reaction.

Difference electron density features associated with movements of the retinal, helix F (Trp182, Thr178 and Arg175) and helix G (Lys216, Ala215 and Ser214) all show a stepwise increase in strength from  $\Delta t = 5.25 \mu\text{s}$  to  $36.2 \mu\text{s}$ , which coincides with the primary proton transfer event (Fig. 3B). We therefore conclude that proton transfer from the SB to Asp85 affects a structural change within the retinal which we model as a displacement of the SB towards the CP by 0.15

Å, and of the C20 methyl by 0.21 Å, over the interval  $760 \text{ ns} \leq \Delta t \leq 36.2 \text{ } \mu\text{s}$ . These changes may be due to the retinal being subtly straighter when deprotonated (32) or due to proton transfer neutralizing the mutual electrostatic attraction between the SB and Asp85 (6). Both geometric and electrostatic effects increase the separation between the SB and Asp85 and would therefore increase the barrier for the reverse proton transfer.

Protein structural changes also hinder the reverse proton transfer from Asp85 to the SB. Wat452, which may help mediate the primary proton transfer through Thr89 (Fig. 4B), is observed for  $40 \text{ ns} \leq \Delta t \leq 13.8 \text{ } \mu\text{s}$  but is not visible after a proton is transferred to Asp85 (Fig. 3C, table S2). An H-bond interaction connecting Thr89 to Asp85 is also lost after the primary proton transfer event, with significant ( $\geq 4 \text{ } \sigma$ ) negative difference electron density becoming visible between O $\gamma$  of Thr89 and O $\delta$  of Asp85 for  $\Delta t \geq 36.2 \text{ } \mu\text{s}$  (Fig. 5D, Movie S2, table S2). These structural changes break the reverse proton transfer pathway from Asp85 to the SB through Thr89 and Wat452 while simultaneously increasing the pKa of Asp85. Thr89 moves closer to the retinal in the M-state with the H-bond from O $\gamma$  of Thr89 to the SB decreasing from 3.30 Å at 760 ns to 3.04 Å at 36.2  $\mu\text{s}$  and to 2.86 Å at 1.725 ms. This tighter H-bond donor interaction to the deprotonated SB locks Thr89 in a conformation that allows the Asp85 side chain to rotate and break its connectivity to Thr89 while forming a new H-bond interaction with Wat451. In contrast with TR-SFX, FTIR spectroscopy studies at low temperature have predicted a tight

Thr89-Asp85 H-bond interaction in the M-state (33), but the complex dynamics of these residue's H-bond interactions may preclude a unique interpretation of spectroscopic observations.

Synchronous with the M-state spectral transition is a displacement of the positively charged head group of Arg82 towards the EC (Figs. 3A, 5C), which implies that this movement is triggered by the protonation of Asp85. This rearrangement displaces a positive charge away from Asp85 and consequently increases the pKa of the primary proton acceptor, hindering the reverse proton transfer reaction (34). These structural perturbations also trigger a rearrangement of Glu194 and Glu204 which were modeled from  $\Delta t \geq 250 \mu\text{s}$ . Similar structural changes have been observed for trapped M-intermediates and were argued to influence the release of a proton to the EC medium (34, 35). Overall, this sequence highlights structural changes that favor Asp85 being protonated and reveal how the SB accessibility to the EC side of the protein is broken after the SB is deprotonated, thereby functionally coupling the primary proton transfer event to the switch in the bR photocycle.

### **Structural changes on the CP side of bR**

Correlated movements are observed on the CP side of helix F as paired negative and positive difference electron density stacking through Trp182, Thr178 and Arg175 (Fig. 6A,B; Movie

S3). As with electron density changes associated with the retinal and near the center of helix G, difference density features on CP portions of helix F arise rapidly but plateau once the SB is deprotonated (Fig. 3B). Structural refinement indicates that the retinal's C20 methyl group is displaced by 1.12 Å towards the CP over the time-sequence sampled here, with a corresponding displacement of Trp182 of 0.96 Å. The largest C $\alpha$  displacement is 1.13 Å and occurs for Lys216 at  $\Delta t = 1.725$  ms. In helix F the C $\alpha$  atoms of residues 174 to 178 are displaced by only 0.4 Å and therefore, as with cryo-trapped structures of M-intermediates (34-37), TR-SFX does not reveal the large movements of helices E and F that were predicted from bR triple-mutant structures (32, 38). This discrepancy may be reconciled by noting that motions of these helices are severely restricted in the P6<sub>3</sub> crystal form because residues 165 and 166 of the E-F loop participate in crystal contacts with residues 232 and 234 of the C-terminus (fig. S7). Although larger motions of the CP portions of helices E and F arise later in the bR photocycle (39, 40) these structural changes cannot be critical for the mechanism of proton pumping because they are suppressed in 3D crystals yet the photocycle is similar to that of bR in the purple membrane (fig. S3) and the D96G/F171C/F219L bR triple mutant is constitutively open to the CP yet is able to pump protons (41).

### **Mechanistic overview**

Retinal isomerization reorients the SB proton into a hydrophobic cavity while breaking its H-

bond to Wat402, both of which lower the proton affinity of the SB (14). An initially twisted retinal becomes planar within 290 ns, causing Trp182 and Leu93 to be displaced towards the cytoplasm and allowing a water molecule to order between Leu93, Thr89 and the SB in the L-state. H-bond interactions from the protonated SB to Wat452 or Thr89 create a pathway for proton transfer to Asp85 (Fig. 4B) and explain how the SB makes contact with Asp85 despite been turned towards the CP by photo-isomerization. A steric clash between C $\epsilon$  of Lys216 and Wat402 dislodges this water molecule, triggering the collapse of the water mediated H-bond network on the EC side of bR. This allows helix C to bend towards helix G approximately 10  $\mu$ s after photoactivation and raises the pKa of Asp85 to the point where it may spontaneously accept a proton from the SB. Once a proton is transferred, the Asp85-Thr89 H-bond is lost (Fig. 5D) which breaks the SB connectivity to the EC side of the protein. Consequently diffraction data spanning five orders of magnitude in time reveal how structural changes in bR achieve unidirectional membrane transport a half century after Jardetzky first proposed the alternating access framework obeyed by all membrane transporters (1).



## CAPTIONS TO FIGURES:

**Figure 1:** Structure and function of bR. (A) Schematic illustrating retinal covalently bound to Lys216 through a protonated Schiff base in an all-*trans* and 13-*cis* configuration. (B) Proton exchange steps (arrows) achieving proton pumping by bR. The primary proton transfer step is from the SB to Asp85 and corresponds to the spectroscopic L to M transition. (C)  $2mF_{\text{obs}} - DF_{\text{calc}}$  electron density for the bR active site in its resting conformation. Electron density (grey) is contoured at  $\pm 1.5 \sigma$  ( $\sigma$  is the root mean square electron density of the map).

**Figure 2:** Early structural changes in the bR photocycle. (A) View of the  $|F_{\text{obs}}|^{\text{light}} - |F_{\text{obs}}|^{\text{dark}}$  difference Fourier electron density map near the retinal for  $\Delta t = 16$  ns. Blue is positive difference electron density and yellow is negative difference electron density (contoured at  $\pm 3.5 \sigma$ ;  $\sigma$  is the root mean square electron density of the map). The resting state bR model (purple) was used for phases when calculating this map. Paired negative and positive difference electron densities indicate a sideways movement of the retinal's C20 methyl. (B) Identical representation but for the time-point  $\Delta t = 290$  ns. (C) Crystallographic models deriving from partial occupancy refinement for  $\Delta t = 16$  ns (blue) and  $\Delta t = 290$  ns (red) superimposed upon the resting bR structure (purple, partially transparent).

**Figure 3:** Relative changes of electron density amplitudes in bR with time. **(A)** Time-dependent root mean square difference electron density amplitudes associated with Wat402 (blue), helix C (green, backbone of Arg82, Ala84 and Asp85), Waters 400, 401 and 451 (red, EC waters) and the side chain of Arg82 (purple). **(B)** Time-dependent amplitudes associated with the retinal (blue), the side chain of Lys216 (green), helix G (red, backbone of Ser214, Ala215 and Lys216) and helix F (purple, side-chains of Trp178 and Trp182 and backbone of Arg175). **(C)** Time-dependent amplitudes associated with Leu93 (blue), the side chain of Thr89 (green), Wat452 (red) and the side chain of Asp85 (purple). Amplitudes were averaged over the last four time-points and scaled to 1 (except Leu93, Wat452 (42)). Amplitudes for  $\Delta t = 95.2 \mu\text{s}$ ,  $657 \mu\text{s}$  and  $1.725 \text{ ms}$  are scaled by 2.0, 1.5 and 0.8 respectively (42). Linear decomposition of electron density amplitudes show how these motions are coordinated (fig. S8). The M-state population measured from microcrystalline slurries is indicated for comparison (dashed black line).

**Figure 4:** Pathway for proton transfer from the SB to Asp85. **(A)** View of the  $|F_{\text{obs}}|^{\text{light}} - |F_{\text{obs}}|^{\text{dark}}$  difference Fourier electron density map (contoured at  $\pm 3.0 \sigma$ ) near the retinal for  $\Delta t = 760 \text{ ns}$ . A positive difference feature (arrow) suggests the ordering of a water molecule (Wat452). **(B)** Crystallographic model for the time-point  $\Delta t = 760 \text{ ns}$  (red) superimposed upon the resting state model (purple, partially transparent).

**Figure 5:** bR conformation controlling the primary proton transfer event. **(A,B)** View of the Difference Fourier electron density map immediately to the EC side of the retinal for **(A)**  $\Delta t = 2 \mu\text{s}$  and **(B)**  $\Delta t = 36.2 \mu\text{s}$ . All maps are contoured at  $3.5 \sigma$ . Difference Fourier electron density maps from this viewpoint are shown for all thirteen time-points in Movie S2. **(C)** Crystallographic structural models deriving from partial occupancy refinement are superimposed upon the resting bR structure (purple, partially transparent) for  $\Delta t = 16 \text{ ns}$  (blue),  $760 \text{ ns}$  (red),  $36.2 \mu\text{s}$  (orange) and  $1.725 \text{ ms}$  (yellow). **(D)** Close-up view of difference electron density across the Asp85-Thr89 H-bond for the time-points  $\Delta t = 2 \mu\text{s}$ ,  $5.25 \mu\text{s}$ ,  $13.8 \mu\text{s}$ ,  $36.2 \mu\text{s}$ ,  $95.2 \mu\text{s}$  and  $250 \mu\text{s}$ .

**Figure 6:** Conformational changes on the CP side of bR. **(A-B)** Close up view of the difference Fourier electron density map immediately to the CP side of the retinal for **(A)**  $\Delta t = 760 \text{ ns}$  and **(B)**  $\Delta t = 1.725 \text{ ms}$ . All maps are contoured at  $\pm 3.5 \sigma$ . Difference Fourier electron density maps from this viewpoint are shown for all thirteen time-points in Movie S3. **(C)** Crystallographic structural models deriving from partial occupancy refinement are superimposed upon the resting bR structure (purple, partially transparent) for  $\Delta t = 16 \text{ ns}$  (blue),  $760 \text{ ns}$  (red),  $36.2 \mu\text{s}$  (orange) and  $1.725 \text{ ms}$  (yellow).

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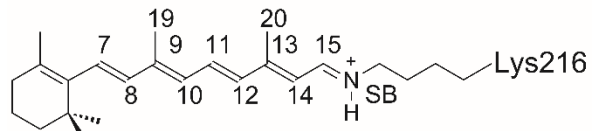
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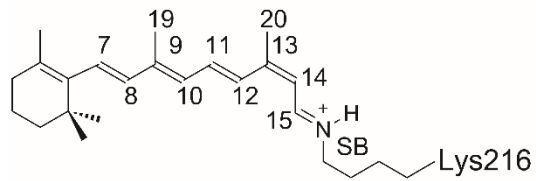
paper with input from all authors. Coordinates and structure factors have been deposited in the Protein Data Bank with accession numbers: 5B6V (bR resting state); 5B6W ( $\Delta t = 16$  ns); 5H2H ( $\Delta t = 40$  ns); 5H2I ( $\Delta t = 110$  ns); 5H2J ( $\Delta t = 290$  ns); 5B6X ( $\Delta t = 760$  ns); 5H2K ( $\Delta t = 2$   $\mu$ s); 5H2L ( $\Delta t = 5.25$   $\mu$ s); 5H2M ( $\Delta t = 13.8$   $\mu$ s); 5B6Y ( $\Delta t = 36.2$   $\mu$ s); 5H2N ( $\Delta t = 95.2$   $\mu$ s); 5H2O ( $\Delta t = 250$   $\mu$ s); 5H2P ( $\Delta t = 657$   $\mu$ s); and 5B6Z ( $\Delta t = 1.725$  ms). Raw diffraction images have been deposited in the CXIDB (accession ID 53).



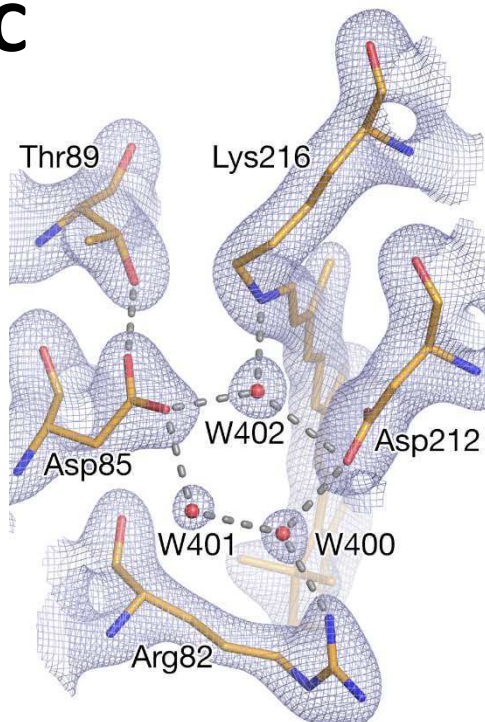
**A** all-*trans* retinal



13-*cis* retinal

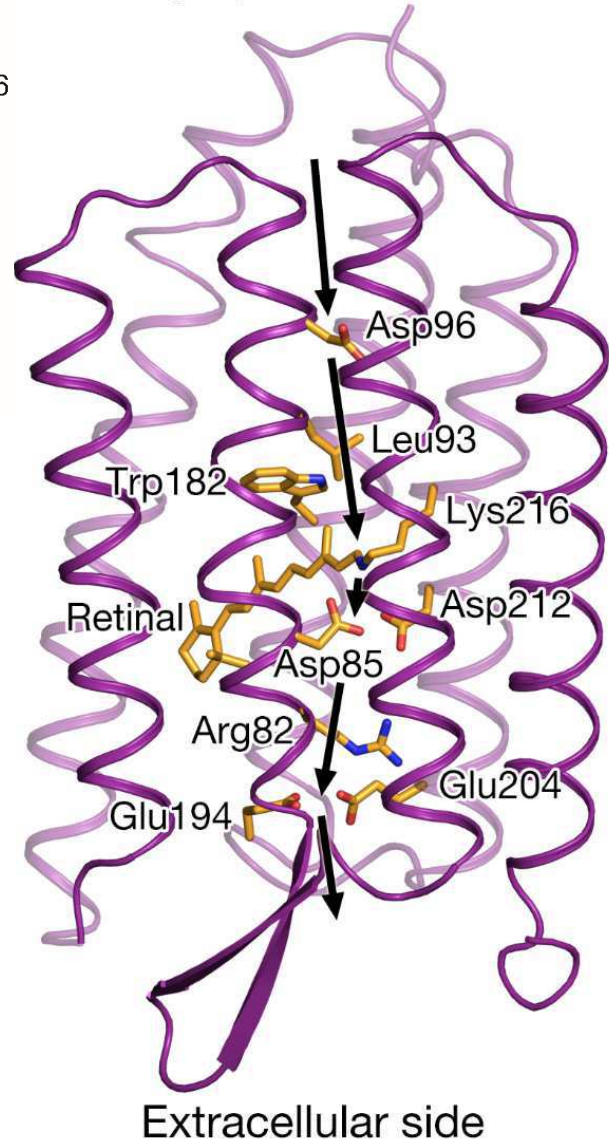


**C**

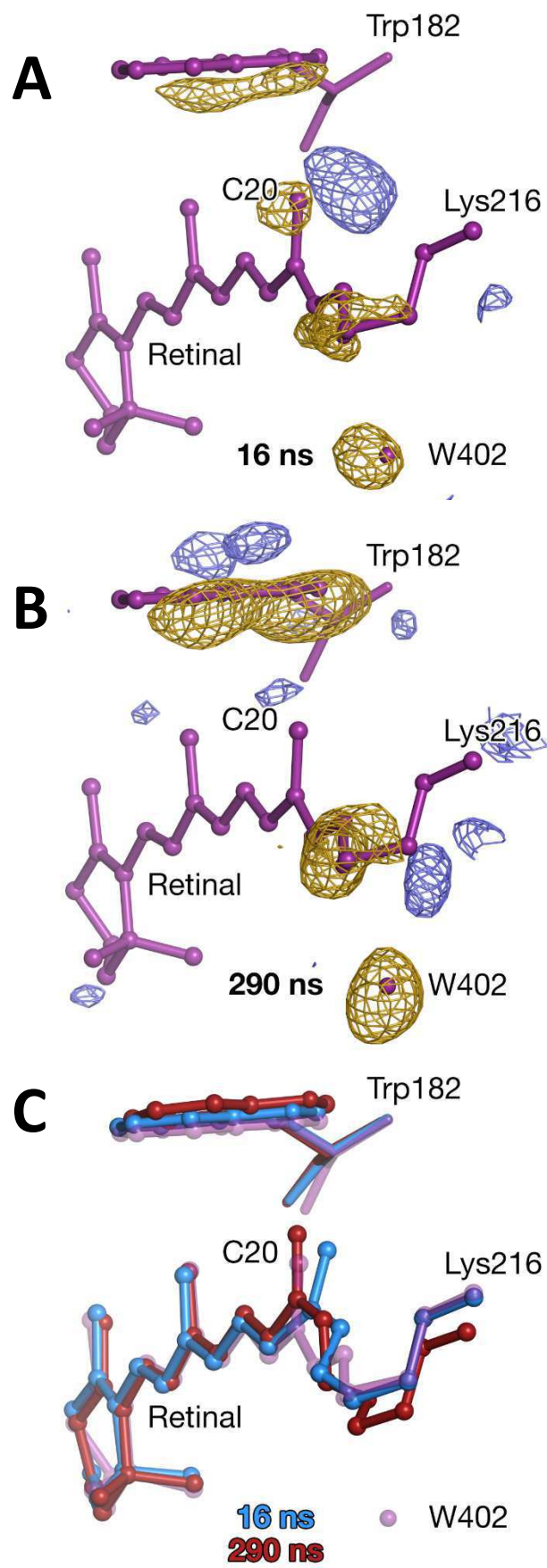


**B**

Cytoplasmic side



**Figure 1: Nango *et al.***



**Figure 2: Nango *et al.***

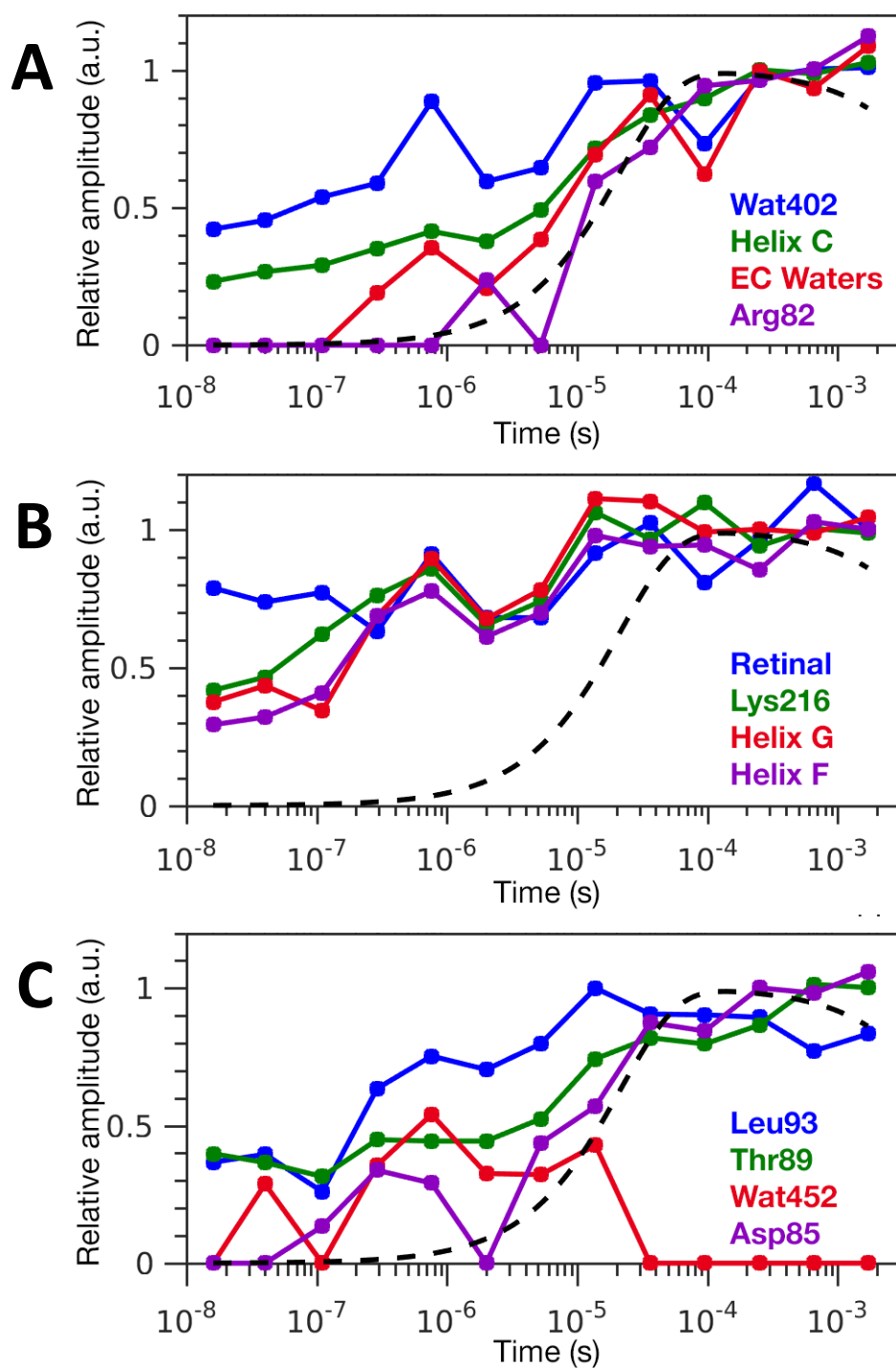


Figure 3: Nango *et al.*

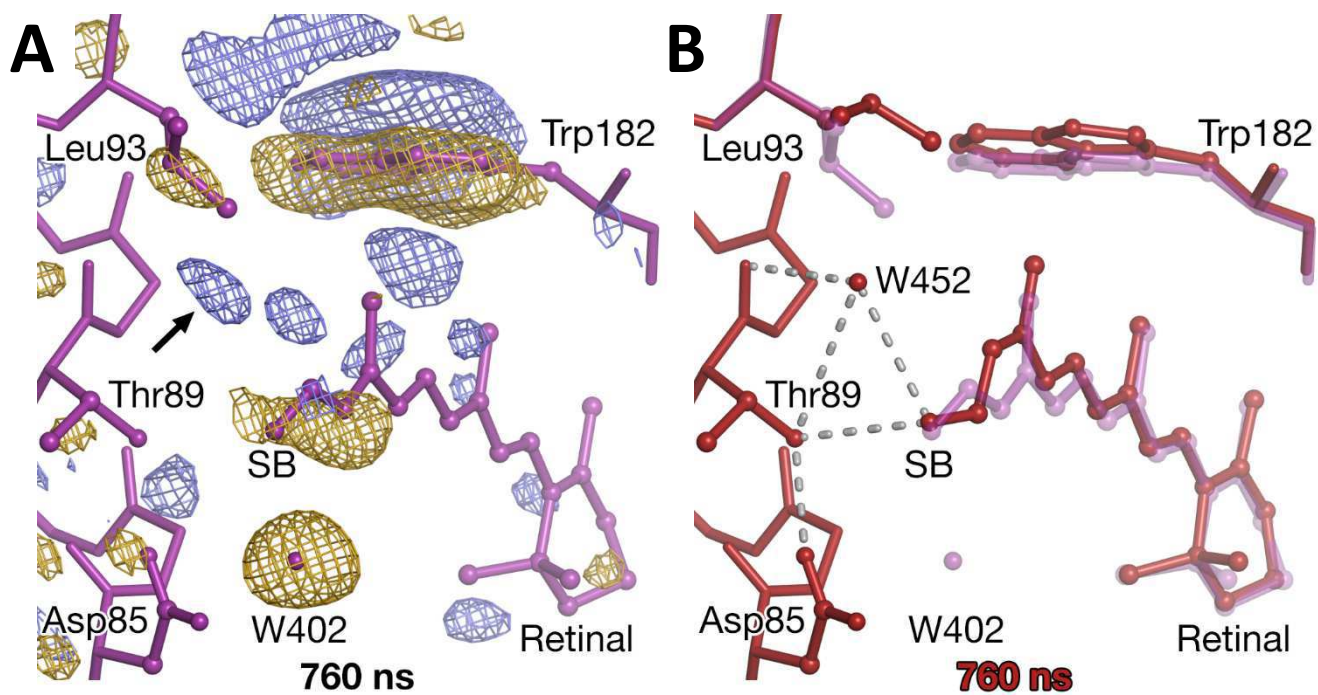
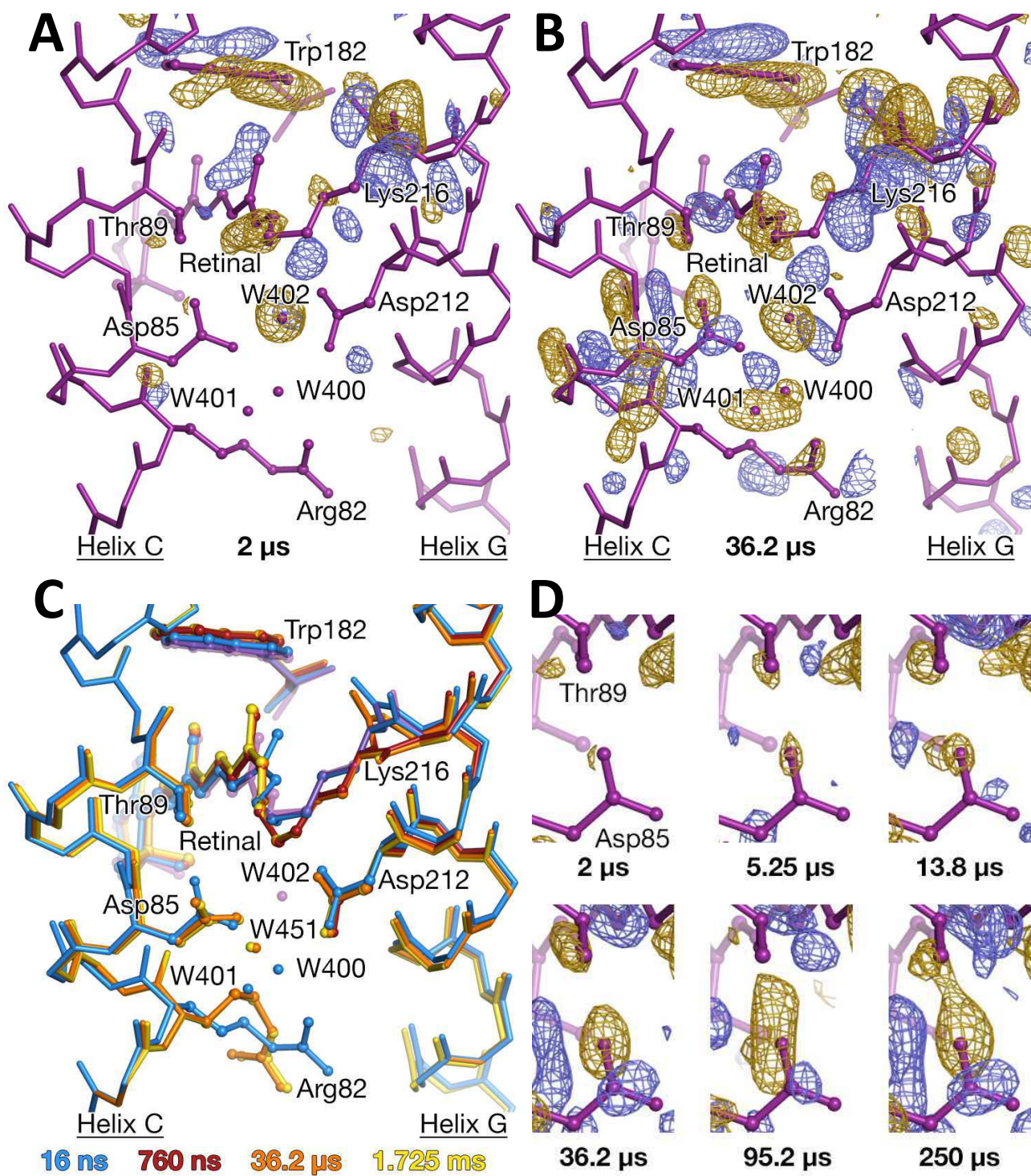
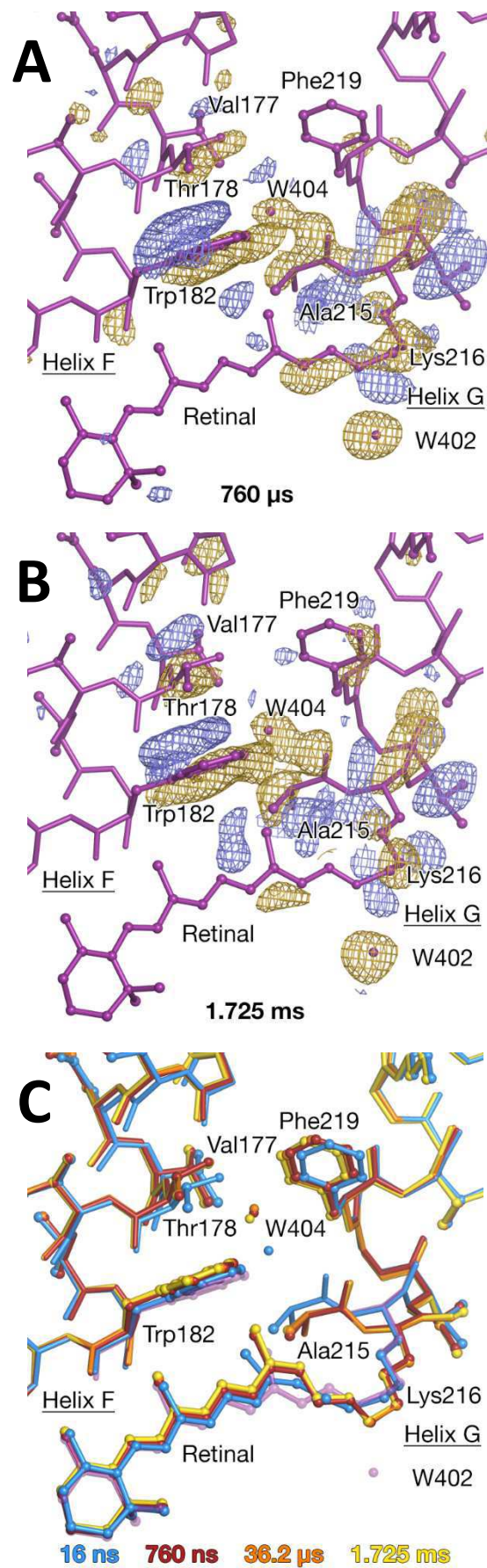


Figure 4: Nango *et al.*





**Figure 5: Nango *et al.***



**Figure 6: Nango *et al.***