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Acoustic higher-order topological insulator on a Kagome lattice

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Higher-order topological insulators (TIs)[1-5] are a family of recently-predicted topological phases of matter obeying an extended topological bulk-boundary correspondence principle. For example, a two-dimensional (2D) second-order TI does not exhibit gapless one-dimensional (1D) topological edge states, like a standard 2D TI, but instead has topologically-protected zero-dimensional (0D) corner states. The first prediction of a second-order TI[1], based on quantized quadrupole polarization, has been demonstrated in classical mechanical[6] and electromagnetic[7,8] metamaterials. Here, we experimentally realize a second-order TI in an acoustic metamaterial, based on a "breathing" Kagome lattice[9], that has zero quadrupole polarization but nontrivial bulk topology characterized by quantized Wannier centers (WCs)[2,9,10]. Unlike previous higher-order TI realizations, the corner states depend not only on the bulk topology but also on the corner shape; we show experimentally that they exist at acute-angled corners of the Kagome lattice, but not at obtuse-angled corners. This shape dependence allows corner states to act as topologically-protected but reconfigurable local resonances.

In a *d*-dimensional TI, the bulk-boundary correspondence principle[11] states that a topologically nontrivial bulk bandstructure implies the existence of (*d*-1)-dimensional boundary states. In the quantum Hall effect, for example, the nontrivial 2D bulk is characterized by nonzero Chern numbers, implying the existence of topologically-protected states on each one-dimensional (1D) edge. Recent theoretical work has led to the prediction of a new class of "higher-order TIs" obeying a generalization of the standard bulk-boundary correspondence[1-5,9,10,12-16]. A second-order TI in *d* dimensions lacks topologically-protected gapless (*d*-1)-dimensional boundary states, but instead exhibits (*d*-2)-dimensional topological states on the "boundaries of boundaries". Each (*d*-1)-dimensional boundary can itself be treated as a first-order TI. Likewise, a third-order TI in *d* dimensions supports (d-3)-dimensional topological states, and their (d-2)-dimensional boundaries are second-order TIs. So far, a few second-order 2D TIs have been realized, using classical mechanical[6] and electromagnetic[7,8] metamaterials. These realizations utilized square lattices with topological properties based on the quantization of quadrupole moments[1,2].

Here, we report on the experimental realization of a 2D higher-order TI on an acoustic Kagome lattice. This lattice has several distinctive features compared to previously-studied square lattice higher-order TIs. First, whereas the topological phases of previous higher-order TIs were characterized by quantized lattice quadrupole moments[1,6-8], the present lattice exhibits quantized *dipole* moments. The well-known 1D Su-Schrieffer-Heeger (SSH) model[17] (which has long been studied in the framework of conventional bulk-boundary correspondence[11]) exhibits similar quantized dipole polarizations, and our lattice can be used for the realization of a higher-order TI generalizing these features to 2D. Second, the quantized dipole moments manifest as acoustic corner states that depend not only on the bulk topology, but also on the corner shape; certain corners never support corner states, even when the bulk is topologically nontrivial. This behaviour can be explained using a topological invariant based on quantized WCs[9,15]. Third, although acoustics has been gaining increasing attention as a flexible platform for studying topological phases[18-26], all studies of acoustic TIs[19,20,23,24] in the emerging field of topological acoustics have, until now, been limited to first-order TIs; our work extends the use of acoustic platforms towards higher-order TIs.

The Kagome lattice is shown in Fig. 1a. Each unit cell consists of three atoms, and the nearest-neighbour couplings on the upward- and downward-pointing triangles are t_1 and t_2 respectively. This tight-binding model is an extension of the 1D SSH model[17], and its bulk topology can be characterized by the polarization[27,28], expressed as

$$p_i = -\frac{1}{S}\iint_{BZ} A_i d^2k \tag{1}$$

where $A_i = -i\langle u|\partial k_i|u\rangle$ with i=x, y is the Berry connection of the lowest band, and S is the area of first Brillouin zone. The polarization (p_x, p_y) is identical to the WC. Mirror symmetries restrict the WC to two positions within each unit cell, corresponding to the two topologically distinct phases of the bulk. We refer to these as topologically 'trivial' and 'nontrivial' phases. Previous theoretical studies[9] have shown that (p_x, p_y) is entirely determined by the ratio t_1/t_2. In the present experimental scenario, we only consider positive values of t_1/t_2. For $t_1/t_2>2$, the system is topologically trivial and the WC lies at (0,0), defined as the center of the upward-pointing triangle (indicated in blue in Fig. 1a). For $0< t_1/t_2 <1/2$, the system is topologically nontrivial, and the WC lies at (-1/2, -1/2√3), the center of the downward-pointing triangle (indicated in yellow in Fig. 1a). Note that even though the values of (p_x, p_y) depend on the choice of unit cell, the WC positions within the lattice are unambiguous.

We implement this Kagome lattice model using acoustic resonators, shown schematically in Fig.1b. Similar coupled-resonator structures have previously been used to study Weyl points and Landau levels in acoustics[21,22,25,26]. Each resonator is an air-filled cylindrical cavity with metal walls, of height H = 41 mm and radius r = 20 mm. The surfaces of the cavity are acoustic hard boundaries. For an isolated resonator, the resonant acoustic mode of interest is shown in Fig. 1c; the acoustic pressure varies sinusoidally in the axial (z) direction and is homogenous in the xy plane. The coupling between each pair of nearest-neighbour resonators is provided by two identical thin cylindrical connecting waveguides, placed at heights H/4 and 3H/4. The coupling strength is tunable by varying the radius of the connecting waveguides, with radius r_{c1} (r_{c2}) corresponding to the coupling strength t_1 (t_2) in Fig. 1a. For r_{c1}= r_{c2} = 5.2 mm and lattice constant a = 108 mm, numerical simulations produce the bulk bandstructure shown in Fig. 1d, which has two dispersive bands that meet at linear band-crossing points, with an additional flat band above.

To open a gap, we vary the coupling strengths t_1 and t_2. Upon decreasing r_{c1} to 2.08 mm and increasing r_{c2} to 8.32 mm, we achieve $t_1 /t_2 = 0.1$ (estimated by fitting simulation results to the tight-binding model). In this phase, the WCs are located at the centers of the downward-pointing triangles, marked by red stars in Fig. 2a. When a large triangle-shaped section is cut from the lattice, along the three red dashed lines depicted in Fig. 2a, the boundary runs through the downward-pointing triangles, and hence induces a separation of the charge associate with the WC. We therefore expect the corners of the large triangular section to host corner states. By the same token, the charges associated with the WCs along the edges also experience separation, giving rise to edge states.

The numerically calculated eigenfrequencies and eigenmodes are shown in Figs. 2b and c-f. As expected, three degenerate corner states are found at 4197.3 Hz, within the bulk bandgap. Fig. 2c shows the eigenmode of one of the corner states, showing that the acoustic pressure is highly localized at a corner resonator; there are two other degenerate corner states, localized at the other two corners. The intensity distribution of the corner states is distinct from the edge states (Fig. 2d) and bulk states (Figs. 2e-f). When we switch the values of r_{c1} and r_{c2}, so that $t_1/t_2> 2$, the system becomes trivial and there are no corner states (see Supplementary Information for detailed analysis and an experimental demonstration).

Our experimental sample, shown in Fig. 3a, was fabricated by drilling holes in three pieces of aluminum, and stacking them together between two organic glass sheets (see Supplementary Information for details). The acoustic measurement is conducted by exciting a resonator through a small hole in the bottom organic glass sheet, and then measuring the acoustic pressure of another resonator through a small hole in the upper organic glass sheet. We measure the bulk transmission by exciting and measuring the two resonators marked '1' and '2' in Fig. 3a; the results are shown as the black curve in Fig. 3b. Two peaks are observed at 3950 Hz and 4400 Hz, corresponding to the two bulk bands observed in the simulations of Fig. 2b, separated by a bandgap. The 4400 Hz peak is higher because of a higher density of states.

We then measure the edge transmission spectrum by exciting and measuring the resonators marked '3' and '4' in Fig. 3a. The measured transmission, indicated by the blue curve in Fig. 3b, shows a peak at around 4080Hz, corresponding to the edge states (see simulation in Fig. 2b). There is another peak at around 4400 Hz, coincident with the higher bulk band, but no peak corresponding to the lower bulk band was observed. This seems to be because the bulk states in the upper band have significant spatial overlap with the lattice edges, whereas those in the lower band have negligible spatial overlap (see Figs. 2e and f). Next, we measure the response of the lower left corner resonator by exciting and measuring from the same resonator. As shown by the red curve in Fig. 3b, the resulting spectrum shows a strong peak at around 4200 Hz, consistent with the frequency of the corner states predicted in the simulation of Fig. 2b. To demonstrate the robustness of corner states, we introduce disorder by placing small metal cylinders with random heights into all resonators, except the three corners (see Supplementary Information). The measured spectrum for the same lower left corner resonator (green curve in Fig. 3b) exhibits a similar resonance peak at around 4200 Hz, verifying the robustness of corner states against bulk disorder.

To further characterize the corner states, we map the distribution of acoustic pressure by exciting each resonator and measuring the acoustic pressure of the same resonator. As shown in Fig. 3c, at around 4200 Hz, the measured acoustic pressures at the three corners are much higher than at other points of the lattice. Figs. 3d-f also show the spatial distributions of the edge, lower bulk and upper bulk states. All these results match the simulation results in Fig. 2.

The aforementioned triangular sample had only one type of corner. We constructed an additional parallelogram-shaped sample that has three different types of corners, denoted by A, $B_{1,2}$, and C in Fig. 4a. This structure can be considered as being cut from the infinite lattice through the green dashed lines in Fig. 2a. Based on the preceding theoretical analysis, we expect a corner state at A (similar to the corners of the previous triangular sample), and edge states on the edges adjacent to that corner. However, the left and top edges are different: they do not pass through WCs, so we expect no corner states at C, B_1 and B_2, and no edge states on the left and top edges. The numerically-calculated eigenstates, shown in Fig. 4b, confirms this reasoning. Our experimental results, based on the same protocol described in the previous paragraph, are shown in Fig. 4c, and agrees well with the theoretical and numerical predictions. In Fig. 4d, we plot the acoustic spectra measured at the four corners. Corners B_1 and B_2 exhibit a peak around 4080 Hz because of the edge states, and another peak around 4450 Hz resulted from the higher-frequency bulk states. The corner C only has two peaks around 4000 Hz and 4450 Hz, for the two bulk bands. Over the whole frequency range of interest, only corner A possesses a peak around 4200Hz, which corresponds to the corner state predicted in Fig. 4b and observed in Fig. 4c.

By exchanging the values of t_1 and t_2, we can switch between the two topologically distinct phases, which transfers the corner state originally at corner A to corner C. (This is equivalent to simply rotating the structure by 180°.) However, corners B_1 and B_2 remain isolated from any corner state, and never exhibit corner states.

The above results demonstrate the acoustic analogue of a second-order TI on a Kagome lattice. Our structure is simple to realize and can serve as a basis for further studies. For example, the acoustic structure can be extended to three dimensions to build higher-order TIs with corner or hinge states. The establishment of quantized WCs as a topological invariant may stimulate more studies in predicting and characterizing higher-order TIs. The shape-dependence of corner states provides an extra degree of freedom, apart from the bulk topology, to switch on and off the topologically protected local resonances. These acoustic topological corner states may have useful applications in, for example, biomedical microfluidic devices, enabling the robust acoustic trapping and manipulation of cells or drug particles, and in high-precision acoustic sensors to selectively measure vibrational signals in a small region.

After the submission of our manuscript, several recent experimental papers on higher-order TIs were brought to our attention[29-32]. We note that the use of WCs as a topological invariant for predicting and charactering higher-order TIs is a distinctive feature of the present work.

References

1 Benalcazar, W. A., Bernevig, B. A. & Hughes, T. L. Quantized electric multipole insulators. *Science* **357**, 61-66 (2017).

2 Benalcazar, W. A., Bernevig, B. A. & Hughes, T. L. Electric multipole moments, topological multipole moment pumping, and chiral hinge states in crystalline insulators. *Phys. Rev. B* **96**, 245115 (2017).

3 Schindler, F. *et al.* Higher-order topological insulators. *Sci. Adv.* **4**, eaat0346 (2018).

4 Langbehn, J., Peng, Y., Trifunovic, L., von Oppen, F. & Brouwer, P. W. Reflection-Symmetric Second-Order Topological Insulators and Superconductors. *Phys. Rev. Lett.* **119**, 246401 (2017).

5 Song, Z., Fang, Z. & Fang, C. (d-2)-Dimensional Edge States of Rotation Symmetry Protected Topological States. *Phys. Rev. Lett.* **119**, 246402 (2017).

6 Serra-Garcia, M. *et al.* Observation of a phononic quadrupole topological insulator. *Nature* **555**, 342-345 (2018).

7 Peterson, C. W., Benalcazar, W. A., Hughes, T. L. & Bahl, G. A quantized microwave quadrupole insulator with topologically protected corner states. *Nature* **555**, 346-350 (2018).

8 Imhof, S. *et al*. Topolectrical-circuit realization of topological corner modes. *Nat. Phys.* **14**, 925-929 (2018).

9 Ezawa, M. Higher-Order Topological Insulators and Semimetals on the Breathing Kagome and Pyrochlore Lattices. *Phys. Rev. Lett.* **120**, 026801 (2018).

10 Ezawa, M. Minimal models for Wannier-type higher-order topological insulators and phosphorene. *Phys. Rev. B* **98**, 045125 (2018).

11 Hasan, M. Z. & Kane, C. L. Topological insulators. *Rev. Mod. Phys.* **82**, 3045–3067 (2010).

12 Lin, M. & Hughes, T. L. Topological Quadrupolar Semimetals. Preprint at https://arxiv.org/abs/1708.08457. (2017)

13 Ezawa, M. Magnetic second-order topological insulators and semimetals. *Phys. Rev. B* **97**, 155305 (2018).

14 Geier, M. *et* al. Second-order topological insulators and superconductors with an order-two crystalline symmetry. *Phys. Rev. B* **97**, 205135 (2018).

15 Khalaf, E. Higher-order topological insulators and superconductors protected by inversion symmetry. *Phys. Rev. B* **97**, 205136 (2018).

16 Ezawa, M. Strong and weak second-order topological insulators with hexagonal symmetry and Z_3 index. *Phys. Rev. B* **97**, 241402(R) (2018).

17 Su, W. P., Schrieffer, J. R. & Heeger, A. J. Solitons in Polyacetylene. *Phys. Rev. Lett.* **42**, 1698-1701 (1979).

18 Xiao, M. *et al*. Geometric phase and band inversion in periodic acoustic systems. *Nat. Phys.* **11**, 240-244 (2015).

19 Yang, Z. *et al.* Topological acoustics. *Phys. Rev. Lett.* **114**, 114301 (2015).

20 Khanikaev, A. B., Fleury, R., Mousavi, S. H. & Alu, A. Topologically robust sound propagation in an angular-momentum-biased graphene-like resonator lattice. *Nat. Commun.* **6**, 8260 (2015).

21 Xiao, M., Chen, W.-J., He, W.-Y. & Chan, C. T. Synthetic gauge flux and Weyl points in acoustic systems. *Nat. Phys.* **11**, 920-924 (2015).

22 Yang, Z. & Zhang, B. Acoustic Type-II Weyl Nodes from Stacking Dimerized Chains. *Phys. Rev. Lett.* **117**, 224301 (2016).

23 He, C. *et al.* Acoustic topological insulator and robust one-way sound transport. *Nat. Phys.* **12**, 1124-1129 (2016).

24 Lu, J. *et al.* Observation of topological valley transport of sound in sonic crystals. *Nat. Phys.* **13**, 369-374 (2016).

25 Yang, Z., Gao, F., Yang, Y. & Zhang, B. Strain-Induced Gauge Field and Landau Levels in Acoustic Structures. *Phys. Rev. Lett.* **118**, 194301 (2017).

26 Li, F. *et al.* Wely points and Fermi arcs in a chiral phononic crystal. *Nat. Phys.* **14,** 30-34 (2017).

27 Vanderbilt, D. & King-Smith, R. D. Electric polarization as a bulk quantity and its relation to surface charge. *Phys. Rev. B* **48**, 4442-4455 (1993).

28 King-Smith, R. D. & Vanderbilt, D. Theory of polarization of crystalline solids. *Phys. Rev. B* **47**, 1651(R)–1654(R) (1993).

29 Noh, J. *et al*. Topological protection of photonic mid-gap defect modes. *Nat. Photon.* **12**, 408-415 (2018).

30 Schindler, F. *et al.* Higher-order topology in bismuth. *Nat. Phys.* **14**, 918–924 (2018).

31 Zhang, X. *et al*. Observation of second-order topological insulators in sonic crystals. Preprint at https://arxiv.org/abs/1806.10028 (2018).

32 Ni, X. *et al*. Observation of Bulk Polarization Transitions and Higher-Order Embedded Topological Eigenstates for Sound. Preprint at https://arxiv.org/abs/1807.00896 (2018).

Methods

Fabrication and simulation. The aluminum plates are fabricated using mechanical machining. All the simulations are performed using finite element solver COMSOL Multiphysics (Pressure Acoustics module), with the walls modelled as acoustic hard-wall boundaries.

Data Availability.

The data that support the plots within this paper and other findings of this study are available from the corresponding author upon reasonable request.

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Author Contributions

All authors contributed extensively to this work. H.X. and Y.Y. fabricated structures and performed measurements. H. X., Y.Y. and F. G. performed simulation. Y.C. and B.Z. supervised the project.

Competing financial interests

The authors declare no competing financial interests.

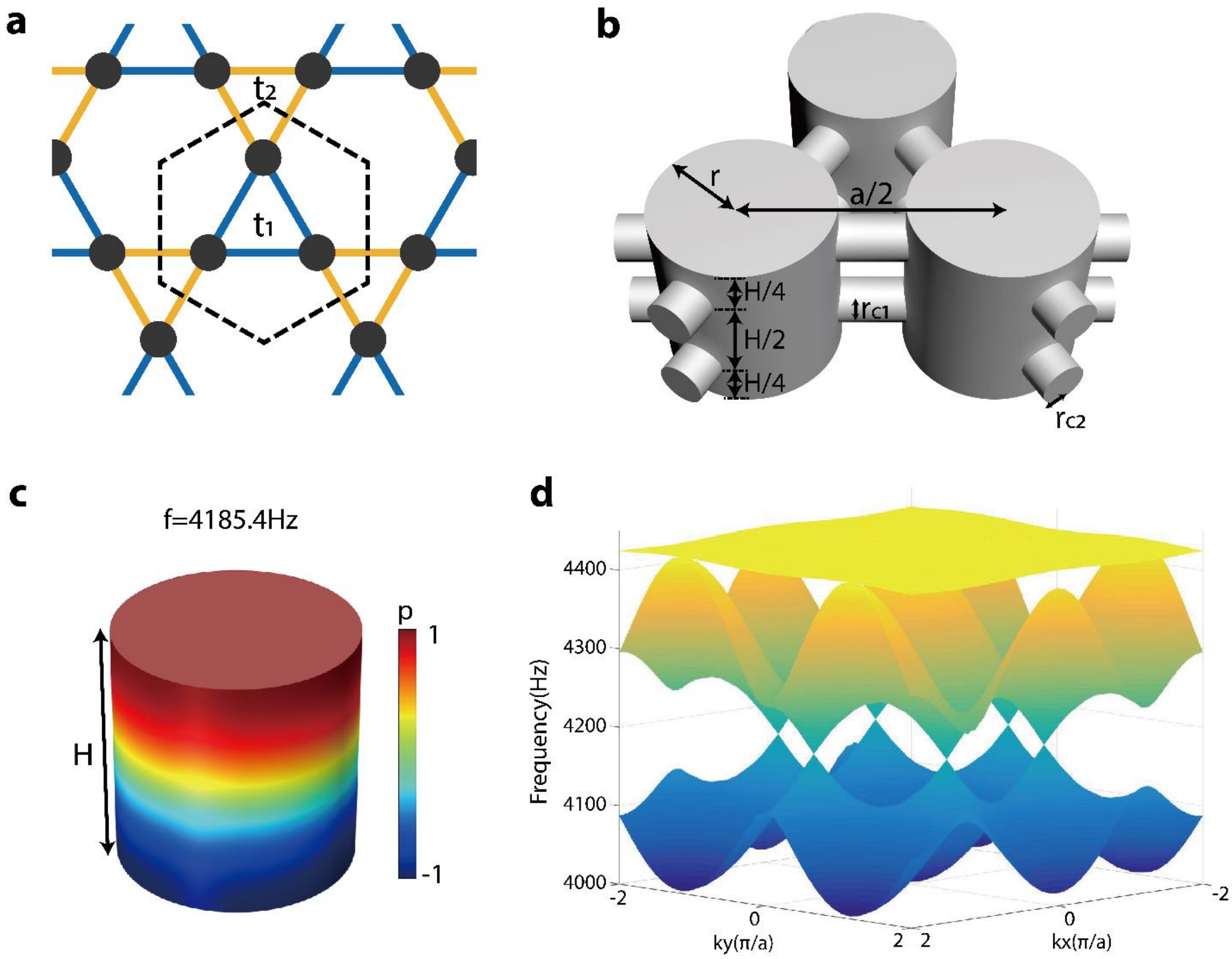

Figure 1 | Kagome lattice and its acoustic implementation. a, Tight-binding model for the Kagome lattice. The dashed hexagon denotes the unit cell. The blue (yellow) lines denote nearest-neighbour couplings of strength t_1 (t_2), which form the sides of upward- (downward-) pointing triangles. **b**, Unit cell of the acoustic Kagome lattice, with a cylindrical resonator at each site joined by thin waveguides at heights $H/4$ and $3H/4$. The connecting waveguides have radii r_{c1} or r_{c2}, corresponding to the t_1 and t_2 coupling strengths. **c**, Real part of the acoustic eigenpressure field for a single acoustic resonator at 4185.4Hz. **d**, Numerically-computed bulk bands for the acoustic Kagome lattice shown in **b**, at the critical point $r_{c1} = r_{c2}$.

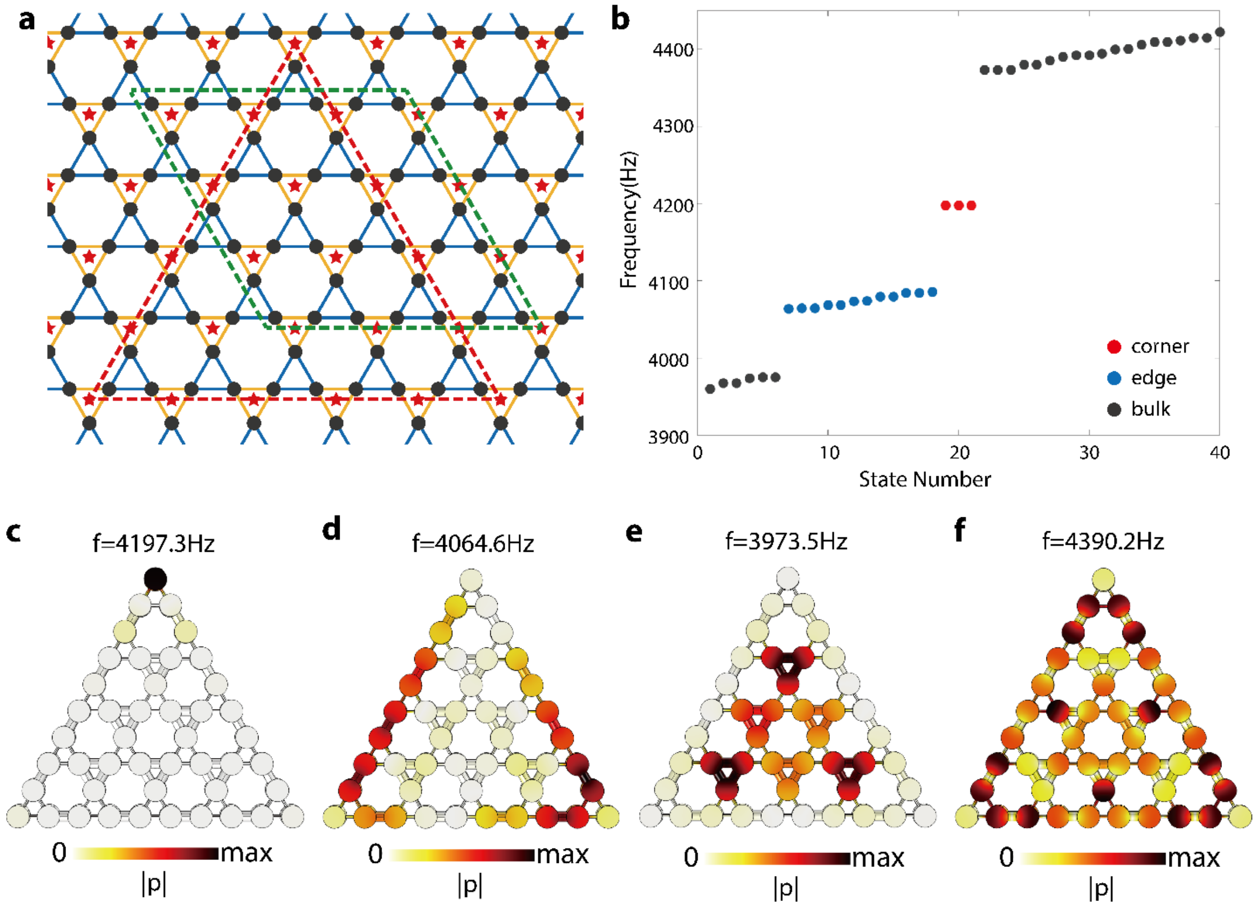

Figure 2 | Eigenmode simulations of a triangular acoustic structure. **a**, Lattice schematic, with red stars indicating the Wannier center positions in the topological nontrivial phase $t_1/t_2<1/2$. Red and green dashes indicate the edges for finite triangular and parallelogram-shaped samples. **b**, Numerically-computed eigenfrequencies for a triangular sample cut along the red dashed lines in **a**. Gray, blue and red dots denote bulk, edge and corner states, respectively. Three degenerate corner states occur at 4197.3 Hz. **c**-**f**, Typical acoustic eigen pressure fields of corner (**c**), edge (**d**) and bulk (**e** and **f**) states.

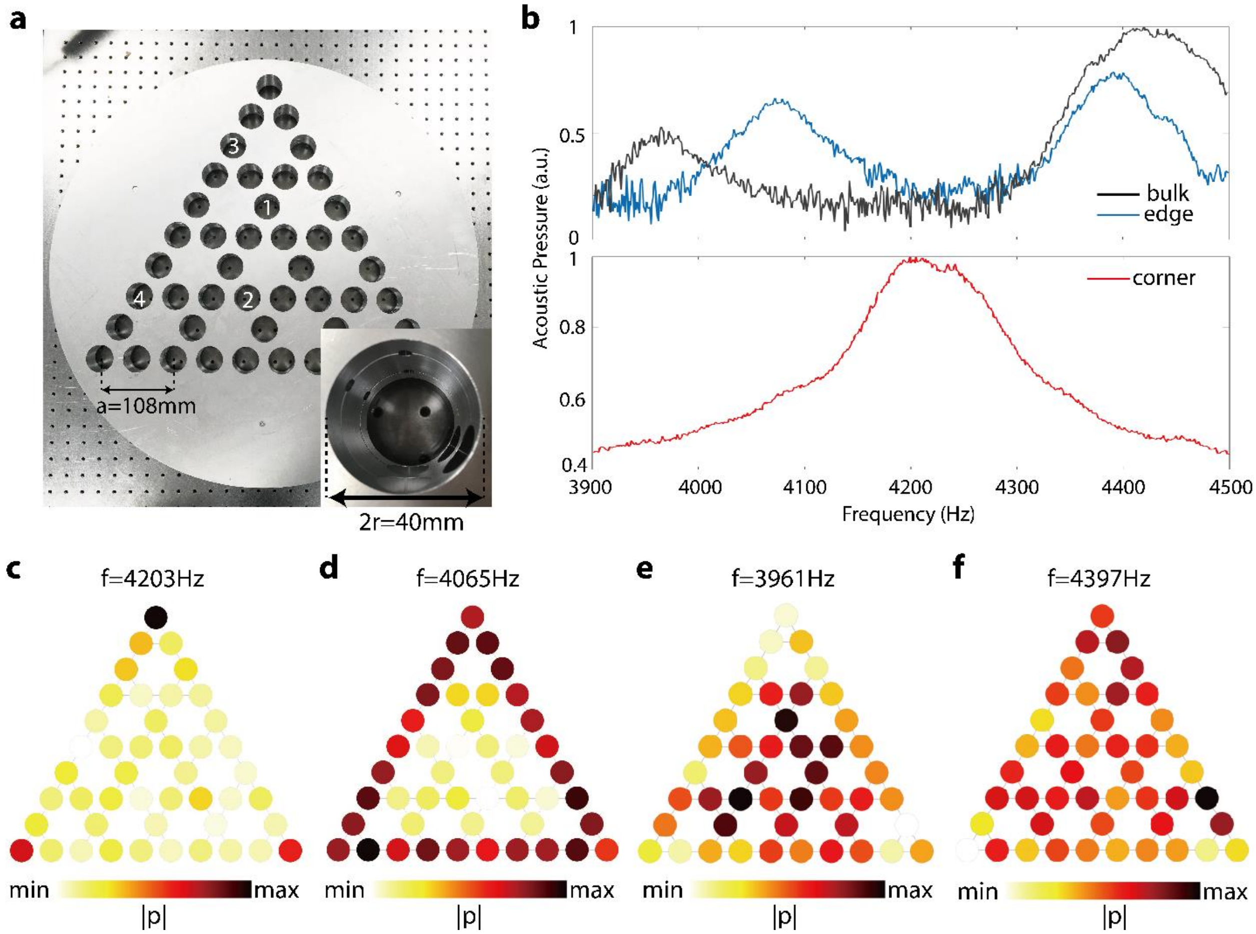

Figure 3 | Observation of topological corner states in a triangle-shaped finite acoustic structure. **a**, Photograph of the fabricated triangle-shaped structure. **b**, Upper: Measured bulk (black) and edge (blue) transmission spectra. Lower: Measured corner spectra for lattice with (green) and without (red) disorder. **c-f**, Measured acoustic pressure distributions in the nontrivial phase at typical frequencies for corner (**c**), edge (**d**) and bulk (**e** and **f**) states.

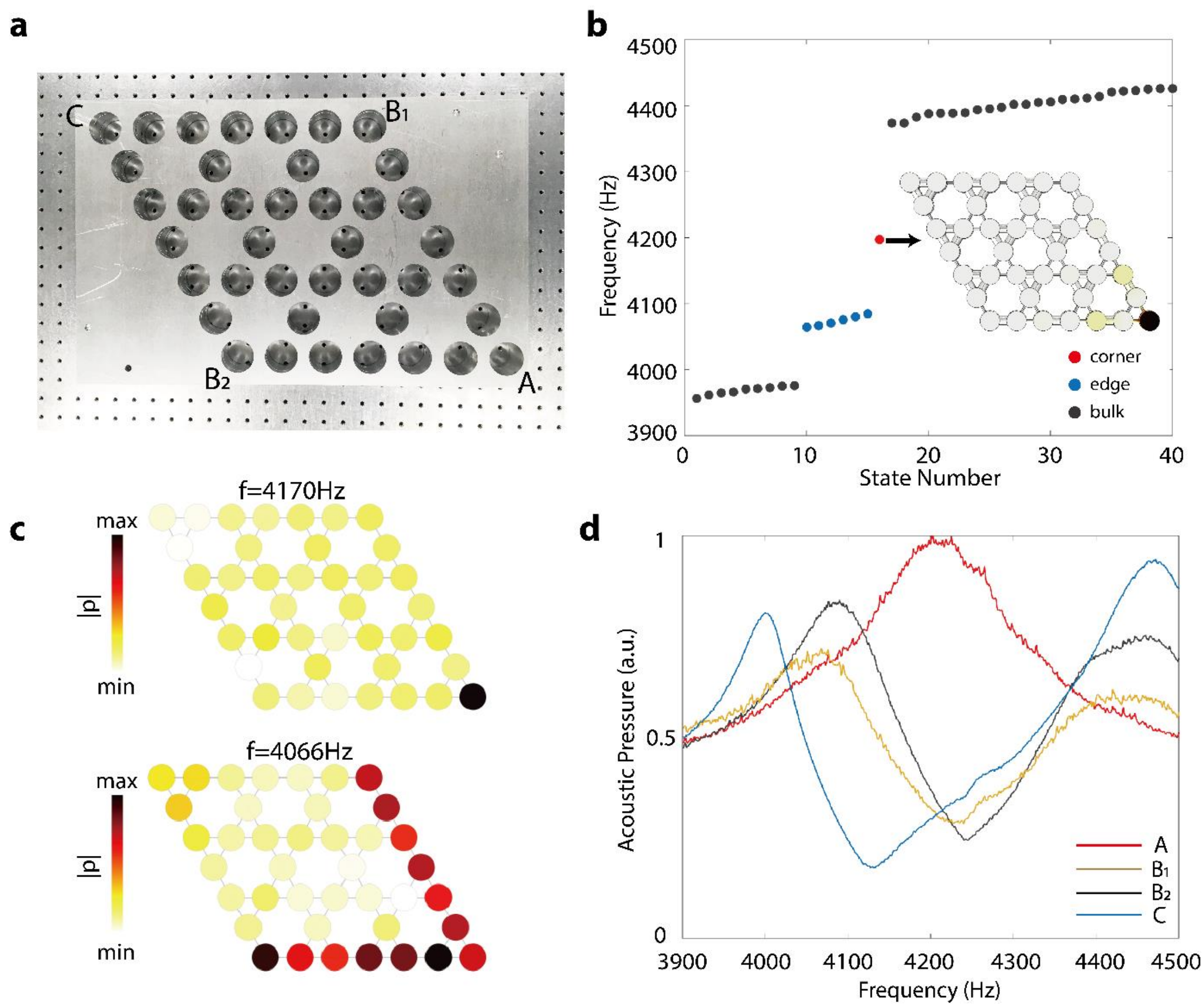

Figure 4 | Observation of topological corner states in a parallelogram-shaped finite acoustic structure. **a**, Photograph of the fabricated parallelogram-shaped structure. **b**, Numerically-computed eigenfrequencies of the structure. Gray, blue and red dots denote bulk, edge and corner states, respectively. There is a single non-degenerate corner state, localized at corner A. **c**, Measured acoustic pressure distributions at 4170 Hz and 4066 Hz. **d**, Measured spectra for the four corners A, B_1, B_2 and C.