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AN ACTIVE DISTURBANCE REJECTION BASED APPROACH TO VIBRATION SUPPRESSION IN TWO-INERTIA SYSTEMS

Shen Zhao and Zhiqiang Gao

I. INTRODUCTION

Vibration suppression is important in motion control applications because vibration causes dynamic stresses, energy wastes and performance degradations [1]. By law of physics, mechanical resonance is unavoidable in every system involving motion, but the natural frequencies of such systems are usually quite high and not excited during most common motion maneuvers, where a simple proportional-integral-derivative (PID) controller is often sufficient to

meet the design requirements. Control design becomes an issue, however, when the performance improvements push the loop bandwidth to its limit where the resonant modes come into play. The most common resonance seen in industry can be attributed to the compliant couplings, such as gear boxes, long shafts and belts, which can be treated as springs [2].

To deal with resonance, there are mechanical and electrical means. Since the resonance is caused by compliance, a stiffer transmission, i.e. a direct coupling in place of a belt, will be an obvious solution. Adding more mechanical damping will surely be helpful. In addition, increasing the motor inertia is found to be an effective way to alleviate the resonance [2]. These mechanical methods are costly, which leads us to electrical options, consist of low-pass filter, notch filter [3] and bi-quad filter [4], all for the objective of attenuating the loop gain amplitude at the resonant frequency so that the resonance is suppressed. Some of the electrical methods are equivalent mathematically

to the mechanical methods mentioned above. Active resonance damping control [4] actually increases the effective physical damping by adding a torque that is proportional to the speed difference between the motor and load. In [5], the active suspension indeed increases the effective damping or spring constant depends on the control design. Acceleration feedback control [4, 6], however, increases the motor inertia equivalently. There are still other control methods available, such as center of mass control [2, 4] and resonance ratio control [7, 8].

All of the above control methods predicate on the detailed mathematical model of the physical process that may or may not be readily available. Even if such a model is obtained at considerable cost, the parameters of the model often change during operation, which may lead to variations in the resonant frequency, leaving the notch filter approach, for example, vulnerable. The attempt to address this flaw leads to solutions such as the adaptive notch filter [9], which is designed to tune the filter parameters on the fly based on adaptive control theory, adding complexity and cost to the design, implementation, and tuning of the control system. It is in this background that an alternative solution is proposed in this paper.

To deal with the resonance problem in motion control, as described above, we resort to a rather novel control method that requires very little system model information and makes the control system tolerant of unknown changes in system dynamics. This method is known as active disturbance rejection control (ADRC) [10–15], based on the key concept of treating the unknown dynamics and disturbances in a physical process as the total disturbance, building a state observer, known as the extended state observer (ESO), to estimate it in real time, and then canceling its effect using a part of the control signal. In the context of the motion control, the resonant mode is not canceled out using a notch filter, but its effect to motion, the ripples in torque, is estimated and canceled in real time using the motor torque, after which the motion dynamic behaves largely like a rigid body. Note that a similar disturbance rejection method was shown in [16], where only the external disturbance is estimated using a state observer and transfer function combined design based on detailed model information.

This paper is organized as follows. The problem description based on two-inertia system model is given in Section II, followed by the main result in Section III, where the motion control problem is reformulated in the context of ADRC. Simulation results and comparison to existing methods are shown in Section IV. Hardware experiments are conducted to verify the

simulation results and are presented in Section V. Finally concluding remarks are included in Section VI.

II. PROBLEM DESCRIPTION AND EXISTING SOLUTION

The compliant resonance problem can be simplified and represented by the two-inertia system model [2, 8] as shown in Fig. 1.

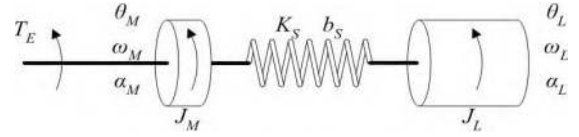


Fig. 1. Two-inertia system model.

Motor inertia (J_M) is connected to load inertia (J_L) by a spring (K_S) with damping (b_S). A torque (T_E) is applied on the motor side to drive the system. Angular acceleration, angular velocity and angular position of the motor and the load are denoted as α_M , ω_M , θ_M , and α_L , ω_L , θ_L , respectively. Through simple analysis, we can derive the transfer functions from input (T_E) to different outputs (ω_M , ω_L , θ_M , and θ_L). The transfer function from T_E to ω_M is

$$\frac{\omega_M}{T_E} = \frac{1}{(J_M + J_L)s} \frac{J_L s^2 + b_S s + K_S}{J_P s^2 + b_S s + K_S} \quad (1)$$

where $J_P = J_M J_L / (J_M + J_L)$. Similarly, we can get the other three transfer functions.

$$\frac{\theta_M}{T_E} = \frac{1}{(J_M + J_L)s^2} \frac{J_L s^2 + b_S s + K_S}{J_P s^2 + b_S s + K_S} \quad (2)$$

$$\frac{\omega_L}{T_E} = \frac{1}{(J_M + J_L)s} \frac{b_S s + K_S}{J_P s^2 + b_S s + K_S} \quad (3)$$

$$\frac{\theta_L}{T_E} = \frac{1}{(J_M + J_L)s^2} \frac{b_S s + K_S}{J_P s^2 + b_S s + K_S} \quad (4)$$

The first term of each transfer function is exactly the same as the transfer function for the rigid body model; the second term which contains resonance is introduced by the compliance. In both motor and load transfer functions, the denominators of the resonance term will produce a resonant frequency (ω_R), and the numerator of the resonance term in motor transfer functions will produce an anti-resonant frequency (ω_{AR}) [2]. They can be calculated by following equations.

$$\omega_R = \sqrt{K_S / J_P} \quad (5)$$

$$\omega_{AR} = \sqrt{K_S/J_L} \quad (6)$$

The Bode plots of velocity transfer functions of rigid body model and compliant model (two-inertia system model) are shown in Fig. 2 for comparison. At low frequency (below the anti-resonant frequency) the two models behave the same. The motor and load are connected as a whole just like the rigid body. As frequency goes higher, the motor and load become disconnected and behave differently. Around resonant frequency there is a 180 degree phase difference between the motor and load, which to some extent represents the resonance as well.

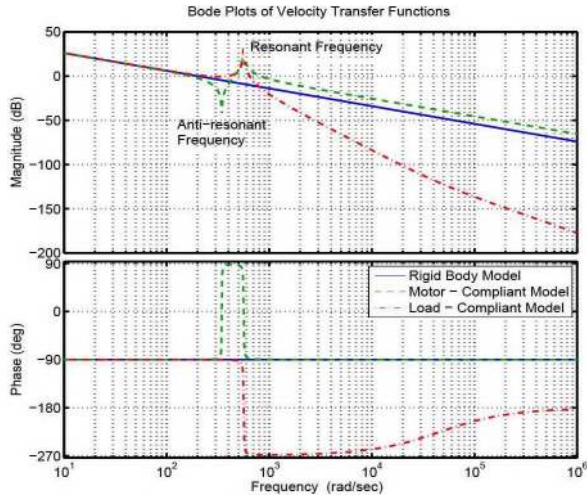


Fig. 2. Bode plots of velocity transfer functions - Rigid vs. Compliant.

Several existing methods are described in [4] that deal with the resonance. A notch filter in the form of

$$T_N(s) = \frac{s^2 + \omega_R^2}{s^2 + 2\zeta\omega_R s + \omega_R^2} \quad (7)$$

is often used to attenuate the open loop gain at the resonant frequency. The bi-quad filter

$$T_{BQ}(s) = \frac{s^2 + 2\zeta_R\omega_R s + \omega_R^2}{s^2 + 2\zeta_{AR}\omega_{AR} s + \omega_{AR}^2} \quad (8)$$

as another solution, not only attenuates the open loop gain at the resonant frequency but also increases the open loop gain at the anti-resonant frequency making it more like a rigid body system. The acceleration feedback method employs a rigid-body Luenberger observer to estimate the motor acceleration and uses it as a feedback for the purpose of increasing the motor inertia.

In a typical configuration of two-inertia system, the sensor is normally mounted at the motor end, where only the motion of the motor is measured and fed back. We denote this set up as motor feedback and this is the common practice in industry. In most cases seen in industry, however, the objective is to control the motion of the load. Consequently, we will also investigate the alternative where we mount the sensor at the load end and use the measurement of the load as feedback, which is denoted as the load feedback. Although the load feedback provides the direct information on how the load behaves, there is a considerable amount of phase lag, comparing to the motor feedback, which makes the control design more challenging. One may suspect that this might be a main reason why the motor feedback configuration is widely used in industry.

Different applications may have different design objectives. Some regulate velocity, others position. To show the generality of the proposed method, both velocity control and position control are addressed in this paper.

III. The Proposed Solution

As mentioned in Section I, active disturbance rejection control (ADRC) provides an alternative design paradigm for the resonance problem in motion control. The main idea of ADRC is to treat any unknown dynamics of the system together with external disturbance as a total disturbance, use an extended state observer (ESO) to estimate this total disturbance in real time, and then cancel it in the control law [10]. In this manner we do not have to know the exact system model in order to control it, and particularly in this application we can treat the resonance, no matter what the frequency is, as part of the total disturbance.

For completeness, we consider two types of motion control, velocity control and position control, and two feedback options, motor feedback and load feedback. Since the only difference between velocity control and position control is that the plant has one more integrator in position control, we will only present the problem reformulation for velocity control in the ADRC structure with both feedback options.

3.1. Velocity control with motor feedback

With $b_0 = 1/J_M$, $b_1 = b_S/(J_M J_L)$, $b_2 = K_S/(J_M J_L)$, $a_1 = b_S/J_P$, $a_2 = K_S/J_P$, and considering an external disturbance w , (1) can be rewritten as

$$\ddot{y}_m + a_1 \dot{y}_m + a_2 y_m = b_0 \ddot{u} + b_1 \dot{u} + b_2 u + w \quad (9)$$

where y_m is the motor velocity, and u is torque applied to the motor. Integrating (9) twice on both sides, the third-order system with a relative degree of one becomes a first-order system [12] as below

$$\begin{aligned}\dot{y}_m &= b_0 u + (-a_1 y_m - a_2 \int y_m \\ &\quad + b_1 \int u + b_2 \iint u + \iint w) \\ &= b_0 u + f(y_m, \int y_m, \int u, \iint u, \iint w) \quad (10)\end{aligned}$$

Here $f(\cdot)$, including both external disturbance and internal dynamics — the resonance, represents the “total disturbance” to be estimated and mitigated. For the first order system (10), the output y_m is defined as the first state x_1 as usual; additionally, the total disturbance f is defined as the extended state x_2 . Thus the state space representation of (10) is

$$\begin{cases} \dot{x} = A_2 x + b_0 B_2 u + E_2 f \\ y_m = C_2 x \end{cases} \quad (11)$$

where $A_2 = \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}$, $B_2 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, $E_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$, $C_2 = \begin{bmatrix} 1 & 0 \end{bmatrix}$.

A corresponding second order ESO is then designed as

$$\dot{z} = A_2 z + \hat{b}_0 B_2 u + L_2 e_o \quad (12)$$

where \hat{b}_0 is the estimated value of b_0 , $L_2 = [\beta_1 \ \beta_2]^T$ is the observer gain vector, and $e_o = x_1 - z_1$ is the observer error. The observer gains are selected based on the observer bandwidth defined and discussed in Section 4.1. Mathematical proof has been shown in [15, 17] that the observer error is bounded if the derivate of the total disturbance f is bounded, and the bound of the observer error is inversely proportional to the observer bandwidth. With appropriate selection of the observer gains, the observer states z_1 and z_2 will track y_m and f respectively.

With the total disturbance being estimated, the control law is then designed as

$$u = \frac{u_0 - z_2}{\hat{b}_0} \quad (13)$$

$$u_0 = k_p(r - y_m) \quad (14)$$

where k_p is the controller gain and r is the reference input. Substituting (13) into (10),

$$\dot{y}_m = b_0 \frac{u_0 - z_2}{\hat{b}_0} + f \approx u_0 \quad (15)$$

Here we can see clearly that the total disturbance is “cancelled” and the plant becomes a pure integrator which can be easily controlled using a proportional controller given in (14).

Compare to the method in [16], the ESO estimates the total disturbance directly, not just the external disturbance, and it uses only a simple, easy to implement and tune state observer, without the need for an additional filter. Furthermore, in the ADRC design less system information is required namely only the motor inertia J_M , whereas in [16] a full system model is needed.

3.2. Velocity control with load feedback

Considering an external disturbance w , (3) can be rewritten as

$$\ddot{y}_l + a_1 \dot{y}_l + a_2 y_l = b_1 \dot{u} + b_2 u + w \quad (16)$$

where y_l is the load velocity, and u is torque applied to the motor. Integrating (16) once on both sides, the third-order system with a relative degree of two becomes a second-order system

$$\begin{aligned}\ddot{y}_l &= b_1 u + (-a_1 \dot{y}_l - a_2 y_l + b_2 \int u + \int w) \\ &= b_1 u + f(\dot{y}_l, y_l, \int u, \int w) \quad (17)\end{aligned}$$

Similarly, for the second order system (17), define the states $x_1 = y_l$, $x_2 = \dot{y}_l$ and $x_3 = f$. The states representation of (17) is

$$\begin{cases} \dot{x} = A_3 x + b_1 B_3 u + E_3 f \\ y_l = C_3 x \end{cases} \quad (18)$$

where $A_3 = \begin{bmatrix} 0 & 1 & 0 \\ 0 & 0 & 1 \\ 0 & 0 & 0 \end{bmatrix}$, $B_3 = \begin{bmatrix} 0 & 1 & 0 \end{bmatrix}^T$, $E_3 = \begin{bmatrix} 0 & 0 & 1 \end{bmatrix}^T$, $C_3 = \begin{bmatrix} 1 & 0 & 0 \end{bmatrix}$.

A corresponding third order ESO is designed as

$$\dot{z} = A_3 z + \hat{b}_1 B_3 u + L_3 e_o \quad (19)$$

where \hat{b}_1 is the estimated value of b_1 , $L_3 = [\beta_1 \ \beta_2 \ \beta_3]^T$ is the observer gain vector, and $e_o = x_1 - z_1$ is the observer error. With appropriate selection of the observer gains, the observer states z_1 , z_2 and z_3 will track y_l , \dot{y}_l and f respectively.

The control law is similarly designed as

$$u = \frac{u_0 - z_3}{\hat{b}_1} \quad (20)$$

$$u_0 = k_p(r - y_l) + k_d(\dot{r} - \dot{y}_l) \quad (21)$$

where k_p and k_d are the controller gains. In this case z_3 is the extended state and a PD controller is designed for the double integrator plant.

For the more detailed derivation of the ADRC control law and recent mathematical analysis of this design approach, the readers are referred to [17–22]. The focus of this paper is on its possible application in motion control in the presence of resonant mode.

IV. SIMULATION RESULTS AND COMPARISON

In this section, the proposed method is tested in simulation and compared to the three existing methods described in [4], using the motor feedback configuration for velocity control as in [4].

4.1. Parameters and profile selection

The proposed method is tested in simulations using the same system parameters as those in [4], with $K_S = 372 \text{ N}\cdot\text{m}/\text{rad}$, $b_S = 0.008 \text{ N}\cdot\text{m}\cdot\text{s}/\text{rad}$, $J_M = 1.88 \times 10^{-3} \text{ kg}\cdot\text{m}^2$, $J_L = 3.13 \times 10^{-3} \text{ kg}\cdot\text{m}^2$, and $J_P = 1.17 \times 10^{-3} \text{ kg}\cdot\text{m}^2$. In this case, the anti-resonant frequency ω_{AR} is 345 rad/s (or 55 Hz), and resonant frequency ω_R is 563 rad/s (or 90 Hz). We also compare our method with those discussed in [4] applying their fine tuned parameters in velocity control with motor feedback. The comparison is not done for other cases because [4] only considers velocity control with motor feedback.

Using the parameterization technique proposed in [23], the observer gains and controller gains are selected such that all of the observer eigenvalues are placed at $-\omega_o$ and all of the controller eigenvalues are placed at $-\omega_c$. Specifically, in a second-order ADRC, $\beta_1 = 2\omega_o$, $\beta_2 = \omega_o^2$, $k_p = \omega_c$; in a third-order ADRC, $\beta_1 = 3\omega_o$, $\beta_2 = 3\omega_o^2$, $\beta_3 = \omega_o^3$, $k_p = \omega_c^2$, $k_d = 2\omega_c$; and in a fourth-order ADRC, $\beta_1 = 4\omega_o$, $\beta_2 = 6\omega_o^2$, $\beta_3 = 4\omega_o^3$, $\beta_4 = \omega_o^4$, $k_p = \omega_c^3$, $k_d = 3\omega_c^2$, $k_{dd} = 3\omega_c$. Above ω_o is the observer bandwidth and ω_c is the controller bandwidth. By fixing the ratio between the observer and controller bandwidth, ω_o becomes the only tuning parameter making the tuning process very easy and intuitive. In this paper we set $\omega_c = \omega_o/2$.

Observer and controller bandwidth are selected based on following considerations: 1) the controller bandwidth should be higher than the required bandwidth given in the specification; 2) the observer bandwidth should be two to five times higher than the controller bandwidth; 3) the observer bandwidth

should be five to ten times less than the sampling rate. Normally higher the bandwidth is, better the performance is; the cost is that the system is more susceptible to noise and has less robustness.

Step reference is a commonly used profile in simulations and real tests, but it is too aggressive and contains components with very broad bandwidth, which will excite the resonant mode of the system. So in industry the trapezoidal profile, which is less aggressive and also energy saving, is widely used instead of step reference.

Even if a trapezoidal profile is used, the rising time of the profile is still crucial to the system performance. The faster the rising time is, more possible the system is going to have resonance. In order to avoid the resonance, we choose our rising time between 0.05 s and 0.1 s in our simulations.

4.2. Observer performance

The proposed method is simulated with the rising time set to 100 ms (0.1 s), the profile starting time set to 0.5 s and a disturbance of 1 N·m applied to the motor at 1 s. Fig. 3 and Fig. 4 show the plots of the observer states versus the actual system states with the observer bandwidth set to 400 Hz. It is noticed that the error is bounded and converges to zero very quickly indicating very good observer performance.

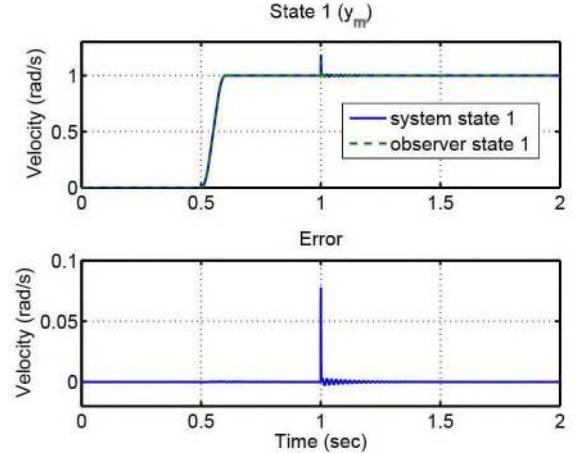


Fig. 3. Estimation of state 1.

4.3. Comparison

The proposed method is then compared to the notch filter, bi-quad filter, acceleration feedback methods with the same profile and disturbance as

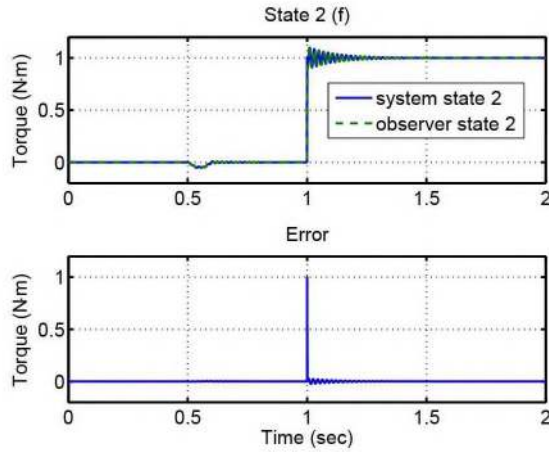


Fig. 4. Estimation of state 2.

Table 1. Motor responses : tracking performance

		Overshoot (%)	5% Settling Time (ms)
Notch Filter		4.2	133
Bi-quad Filter		1.3	115
Accel. Feedback		4.8	137
ADRC ω_o (Hz)	100	0.6	108
	200	0.2	97
	400	0.1	96

Table 2. Motor responses : disturbance rejection performance

		Max. Error (%)	5% Settling Time (ms)
Notch Filter		135	>1000
Bi-quad Filter		70	>1000
Accel. Feedback		72	72
ADRC ω_o (Hz)	100	58	66
	200	34	86
	400	18	94

described in the previous subsection. The results are shown in Table 1 and 2, as well as in Fig. 5.

It is observed that acceleration feedback has the biggest overshoot. Bi-quad filter has less overshoot because it cancels out both resonant and anti-resonant terms in the transfer function. ADRC has even less overshoot and the overshoot decreases as the bandwidth increases. The disturbance rejection ability of acceleration feedback is better than both notch filter and bi-quad filter, which have big errors and oscillate. But ADRC has the best disturbance rejection ability which increases as the bandwidth increases.

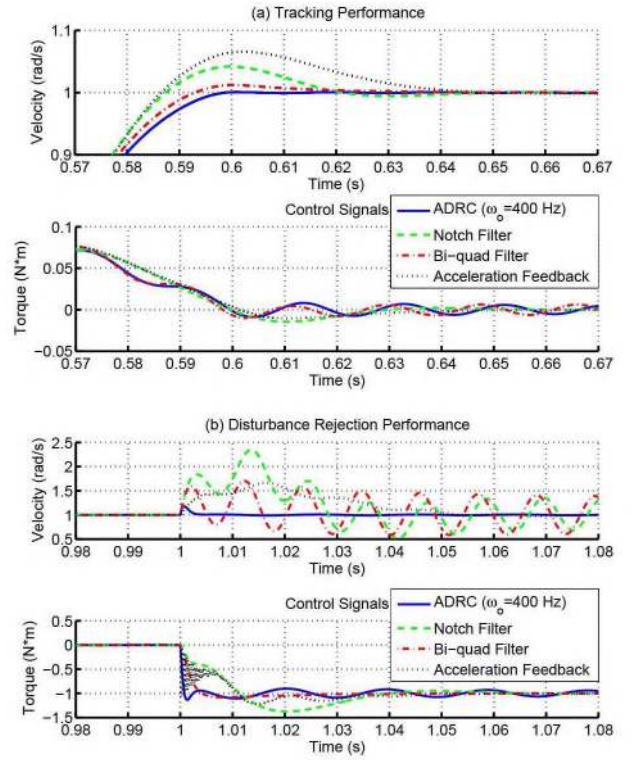


Fig. 5. Motor response comparison. (a) tracking response; (b) disturbance response.

Note that the bandwidth of ADRC can go well beyond the resonant frequency, which is quite difficult to achieve with other methods. As shown in [4] the closed-loop bandwidths associated with the notch filter, the bi-quad filter and acceleration feedback design are 32 Hz, 47 Hz and 37 Hz, respectively, well below the resonant frequency (90 Hz). Based on the frequency response analysis of ADRC [24], with ω_o set to 400 Hz, however, the closed-loop bandwidth of ADRC is found to be 192 Hz, which is well beyond the resonant frequency, unlike the existing methods.

The robustness of each controller is also tested by varying the load inertia without changing the controller parameters. The tests are performed with the load changing to 0.9, 1.1, 2 and 5 times of its original value. The bi-quad filter is found to be the most fragile, because the system becomes unstable for all four load changes. With the notch filter, the system is stable for the first two changes but becomes unstable for last two in the presence of external disturbances. Acceleration feedback and ADRC are stable for all four cases, but the former results in a bigger overshoot of 15%. The motor overshoot in ADRC remains mostly unchanged,

but the load oscillation becomes more pronounced with the increasing load.

4.4. Position control

In this subsection, some simulation results are provided to demonstrate the proposed method works for position control as well. The results are summarized in Table 3. The best performance is obtained at the medium bandwidth of 80 Hz; when the bandwidth goes beyond 150 Hz the system becomes unstable. Fig. 6 shows the response with $\omega_o = 80$ Hz.

Table 3. Tracking performance of position control

ω_o (Hz)	Overshoot (%)	5% Settling Time (ms)
40	0.4	157
60	0.3	133
80	0.1	116
100	0.2	105
120	0.3	102

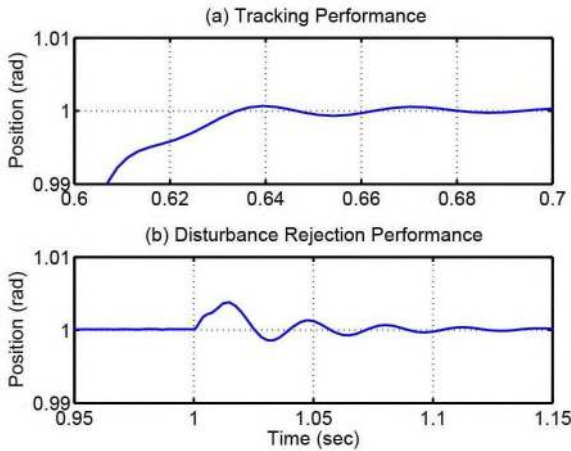


Fig. 6. Position control response. (a) tracking response; (b) disturbance response.

V. EXPERIMENT VERIFICATION

In addition to the simulation comparison with other methods, the proposed control solution to the vibration problem is also verified in hardware tests for the velocity control with motor feedback case. The experiments are conducted on the torsional apparatus Model 205 from Educational Control Products. For a fast validation, the control algorithm is implemented

using the MATLAB real-time workshop in this paper. For application purpose, the implementation of the proposed algorithm can be found in [25].

5.1. Test setup

The torsional apparatus Model 205 has a flexible vertical shaft connecting three disks (lower, middle and upper), with an encoder mounted on the lower disk for the purpose of position measurement. The lower disk is driven by a DC servo motor via belt and pulley system with 3 to 1 speed reduction ratio. In this experiment since we only consider the vibration in a two-inertia system, the upper disk is not used and the belt is tightened to provide a rigid connection that matches the simulation model. There are also brass weights that can be added to the middle disk to test the effect of changing the inertia of the load.

A personal computer (PC), with MATLAB real-time workshop installed, is used to implement the proposed control algorithm. A four-channel quadrature encoder input card (PCI-QUAD04) and a multi-function analog and digital I/O card (PCI-DAS1002), both from Measurement Computing, are install in the computer to interface with the torsional apparatus. A photo of the experimental system is shown in Fig. 7. A diagram is also given (see Fig. 8) to clearly show the mechanical and electrical connections of the system.



Fig. 7. Photo of the test setup.

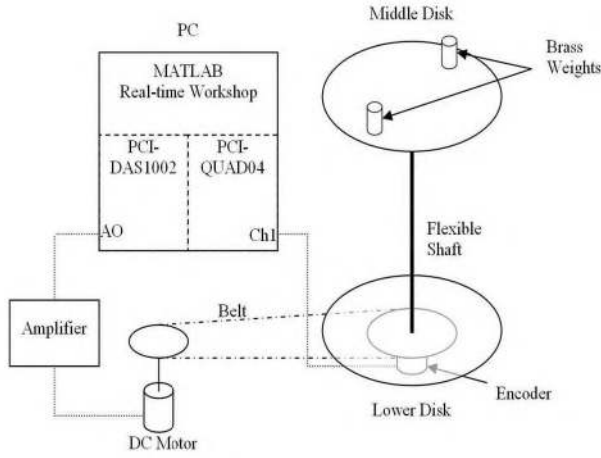


Fig. 8. Diagram of the test setup.

5.2. System parameters

The torque constant ($K_T = T_E/U$) of the motor is 0.058 N·m/V. The encoders generate 16000 pulses per round. Therefore the resolution for position measurement is 3.927×10^{-4} rad (6.25×10^{-5} round). The resolution for velocity measurement depends on the sampling rate, and is 0.196 rad/s (0.03125 round/s) at 500 Hz and 0.393 rad/s (0.0625 round/s) at 1 KHz, i.e. higher the sampling rate lower the resolution. To get a better resolution, a sampling rate of 500 Hz is adopted for velocity control.

To determine the parameters of the test equipment, a frequency sweep test is run by applying a chirp signal with an amplitude of 2 volts to the amplifier. The frequency changes from 0.1 Hz to 15 Hz in 30 seconds. Fig. 9 shows the motor velocity response. The anti-resonant frequency (ω_{AR}) and the resonant frequency (ω_R) are observed at 37.6 rad/s (or 5.99 Hz) and 48.1 rad/s (or 7.65 Hz) respectively from the test. The peak velocity at the resonant frequency is 3.08 round/s.

From Fig. 2 we can see that at low frequency the motor response and the load response are consistent and the whole system behaves like a rigid body. Thus another test is run with a 0.3 Hz sinusoid input to determine the total inertia ($J_T = J_M + J_L$) of the system. The gain at 0.3 Hz is found to be 107.76 round/s/N·m. From (1) J_T is calculated to be 4.92×10^{-3} kg·m². Together with the above frequency sweep test results, from (5) and (6), we get $J_M = 3.01 \times 10^{-3}$ kg·m², $J_L = 1.91 \times 10^{-3}$ kg·m², $J_P = 1.17 \times 10^{-3}$ kg·m², $K_S = 2.71$ N·m/rad, $b_S = 0.006$ N·m·s/rad.

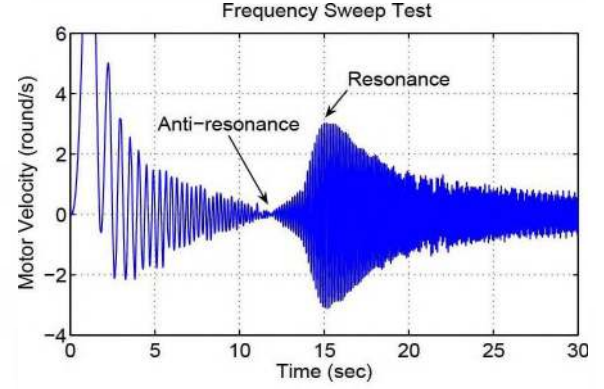


Fig. 9. Motor velocity response of the frequency sweep test.

According to the equipment manual the motor inertia, which includes the inertial of the DC motor, pulley and the lower disk, is around 2.65×10^{-3} kg·m² and the load inertia is around 2.00×10^{-3} kg·m², which matches the tests quite well.

5.3. Test results

A trapezoidal profile, as mentioned in subsection 4.1, with a magnitude of 8 round/s is used to run the tests. The rising time is chosen to be 0.5 seconds which is slower, due to a relative lower resonant frequency compared to the simulation case. The controller under test is described in subsection 3.1, with the observer bandwidth and the controller bandwidth set to 320 rad/s and 160 rad/s respectively. The results are shown in Fig. 10.

Both motor response and load response track the reference very well before the load change. A load with inertia of 3.29×10^{-3} kg·m² is added to the middle disk, which is equivalent to 2.7 times load change, to test the robustness of the control method. The motor velocity remains well controlled with the load change. But the load exhibits oscillations as expected, since resonant frequency is lowered with the load increase and the previous profile is a little fast compare to the new resonance. Test results show that decreasing the rising time to one second will greatly reduce the oscillations.

5.4. Frequency response analysis

Based on the system model, the open loop and closed-loop transfer functions are derived using the above system and controller parameters and the Bode plots are given in Fig. 11 and 12. From Fig. 11, the phase margin of the system is found to be 50 degrees. The closed-loop bandwidth is read from Fig. 12 to be

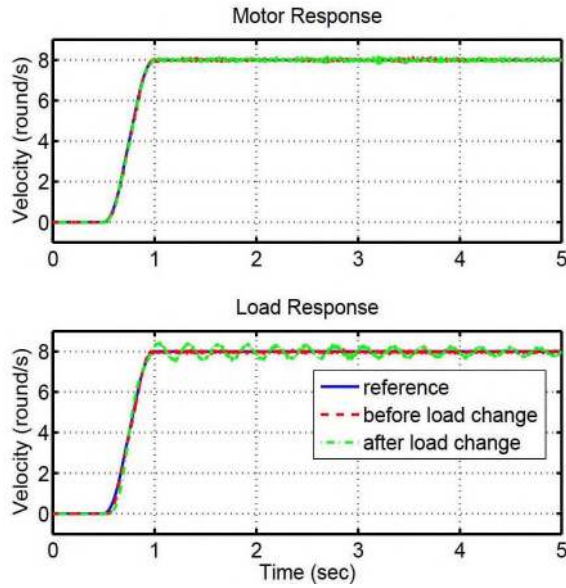


Fig. 10. Velocity control test results.

158 rad/s, which is well beyond the resonant frequency of the system (48.1 rad/s). The resonant mode of the system is attenuated by applying the proposed control method.

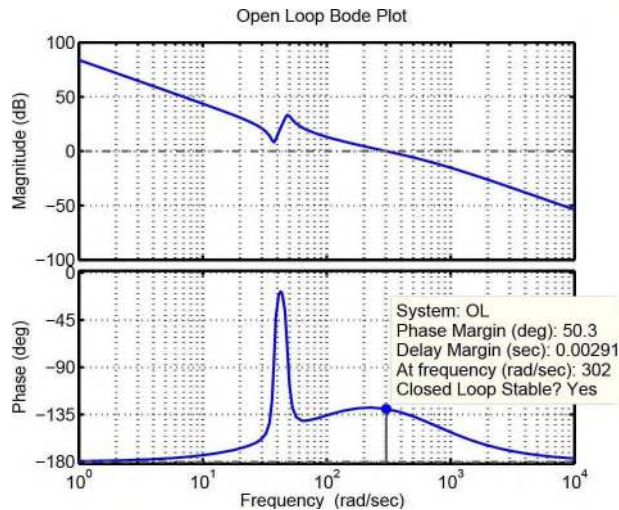


Fig. 11. Open loop Bode plot.

VI. CONCLUSIONS

A novel solution for resonance suppression in motion control is proposed. By reformulating the

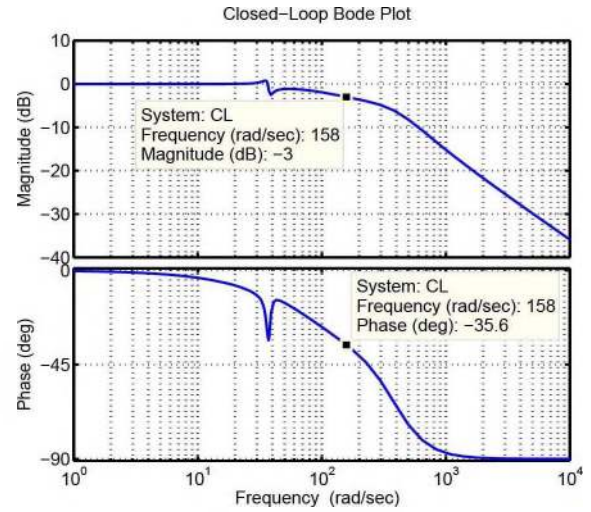


Fig. 12. Closed-loop Bode plot.

problem in the framework of active disturbance rejection control, solutions for both velocity control and position control are presented and compared with the existing methods favorably. It is shown that, with the proposed method, vibration can be eliminated even when the control bandwidth is pushed well beyond the resonant frequency, which is assumed unknown. Both simulation and hardware test results show that the proposed solution works quite well, making it a rather robust and practical solution for motion control.

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