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Institutions: Purdue University
Published on: 01 Apr 1998 - SIAM Journal on Numerical Analysis (Society for Industrial and Applied Mathematics)
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Peters, Jörg and Reif, Ulrich, "Analysis of Algorithms Generalizing B- Spline Subdivision" (1996).
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# ANALYSIS OF ALGORITHMS GENERALIZING 

 B-SPLINE SUBDIVISION
## Jorg Peters

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## CSD-TR 96-066

October 1996

# Analysis of Algorithms Generalizing B-Spline Subdivision 

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May 1, 1996


#### Abstract

A new set of tools for verifying smoothness of surfaces generated by stationary subdivision algorithms is presented. The main challenge here is the verification of injectivity of the characteristic map. The tools are sufficiently versatile and easy to wield to allow, as an application, a full analysis of algorithms generalizing biquadratic and bicubic B-spline subdivision. In the case of generalized biquadratic subdivision the analysis yields a hitherto unknown sharp bound strictly less than one on the second largest eigenvalue of any smoothly converging subdivision.


## 1 Introduction

The idea of generating smooth free-form surfaces of arbitrary topology by iterated mesh refinement dates back to 1978, when two papers [CC78], [DS78] appeared back to back in the same issue of Computer Aided Design. Named after their inventors, the Doo-Sabin and the Catmull-Clark algorithm represent generalizations of the subdivision schemes for biquadratic and bicubic B-splines, respectively. By combining a construction principle of striking simplicity with high fairness of the generated surfaces, both algorithms have since become standard tools in Computer Aided Geometric Design. However, despite a number of attempts [DS78], [BS86], [BS88], so far the convergence to smooth limit surfaces could not be proven rigorously so far.

The proof techniques and actual proofs to be presented here are based on the concept of the characteristic map as introduced in [Rei95a]. The characteristic map is a smooth map from some compact domain $U$ to $\mathrm{R}^{2}$ which can be assigned to stationary linear subdivision schemes. It depends only on the structure of the algorithm and not on the

[^0]data. If this map is both regular and injective, then the corresponding algorithm generates $C^{1}$-limit surfaces. It is shown in this paper that on the other hand non-injectivity at an interior point of the map implies non-smoothness of the limit surfaces. Further, we establish two sufficient criteria for regularity and injectivity of the characteristic map which allow a straightforward verification. The weaker one, however still applicable in many cases, only requires the sign of one partial derivative of one segment of the characteristic map to be positive.

A careful analysis of the Doo-Sabin and the Catmull Clark algorithm yields the following results:

- The Doo-Sabin algorithm in its general form uses weights $\alpha=\left[\alpha^{0}, \ldots, \alpha^{n-1}\right]$ for computing a new $n$-gon from an old one can. Affine invariance and symmetry, i.e. $\sum_{j} \alpha^{j}=1, \alpha^{j}=\alpha^{n-j}, j \in Z_{n}$ imply that the discrete Fourier transform of $\alpha$ is real and of type $\hat{\alpha}=\left[1, \hat{\alpha}^{1}, \hat{\alpha}^{2}, \ldots, \hat{\alpha}^{2}, \hat{\alpha}^{1}\right]$. If $\lambda:=\hat{\alpha}^{1}$ is greater in modulus than the other entries except for 1 and if

$$
\begin{equation*}
1 / 4<\lambda<\lambda_{\max }(n) \tag{1.1}
\end{equation*}
$$

for certain values $\lambda_{\max }(n)<1$ then the limit surface is smooth. The bound $\lambda_{\max }(n)$ can be computed explicitly, see Table 1. If $1>\lambda>\lambda_{\max }(n)$ then the limit is a continuous, yet non-smooth surface.

- In particual, the Doo-Sabin algorithm in its original form (5.1) complies with the conditions, hence generates smooth limit surfaces.
- The Catmull-Clark algorithm in its general form uses three weights $\alpha, \beta, \gamma$ summing up to one for computing the new location of an extraordinary vertex from its predecessor and the centers of its neighbors. If

$$
\begin{equation*}
2\left|4 \alpha-1 \pm \sqrt{(4 \alpha-1)^{2}+8 \beta-4}\right|<c_{n}+5+\sqrt{\left(c_{n}+9\right)\left(c_{n}+1\right)} \tag{1.2}
\end{equation*}
$$

with $c_{n}:=\cos (2 \pi / n)$, then the limit surface is smooth. If one of the both values on the left hand side exceeds the right hand side, then the limit surface is not smooth.

- In particular, the Catmull-Clark algorithm in its original form (6.2) complies with the conditions and generates smooth limit surfaces.


## 2 Generalized subdivision and the characteristic map

In this section we briefly outline the results of subdivision analysis as developed in [Rei95a], and establish a new necessary condition for $C^{1}$-subdivision schemes.

Generalized $B$-spline subdivision generates a sequence $\mathrm{C}_{m}$ of finer and finer control polyhedra converging to some limit surface $y$. On the regular part of the mesh, standard B-spline subdivision is used for refinement, whereas special rules apply near extraordinary mesh points. Since all subdivision masks considered here are of fixed finite size, we can restrict ourselves to analyzing meshes with a single extraordinary mesh point of valence $n \neq 4$. The regular parts of the control polygons $\mathrm{C}_{m}$ correspond to B -spline surfaces $y_{m}$ that form an ascending sequence

$$
\begin{equation*}
\mathrm{y}_{0} \subset \mathrm{y}_{1} \subset \mathrm{y}_{2} \subset \cdots \tag{2.1}
\end{equation*}
$$

converging to the limit surface,

$$
\begin{equation*}
y=\bigcup_{m \in \mathrm{~N}} \mathrm{y}_{m} \tag{2.2}
\end{equation*}
$$

With the prolongation of $\mathbf{y}_{m}$ defined by

$$
\begin{equation*}
\mathbf{x}_{m}:=\operatorname{closure}\left(\mathbf{y}_{m+1} \backslash \mathbf{y}_{m}\right) \tag{2.3}
\end{equation*}
$$

the limit surface is the essentially disjoint ${ }^{1}$ union

$$
\begin{equation*}
\mathbf{y}=y_{0} \cup \bigcup_{m \in \mathrm{~N}} \mathbf{x}_{m} \tag{2.4}
\end{equation*}
$$

The $\mathbf{x}_{m}$ are ring-shaped surface layers which can be parametrized conveniently over a common domain $U \times \mathrm{Z}_{n}, \mathrm{Z}_{\mathrm{n}}:=\mathrm{Z} \backslash n$, consisting of $n$ copies of the compact set

$$
\begin{equation*}
U:=[0,2]^{2} \backslash[0,1)^{2}, \tag{2.5}
\end{equation*}
$$

see Figure 1. Each surface layer $\mathbf{x}_{m}$ can be parametrized in terms of control points $\mathbf{B}_{m}^{k} \in \mathrm{R}^{3}$ and piecewise polynomial functions $N^{\ell}$ according to

$$
\begin{equation*}
\mathbf{x}_{m}: U \times Z_{n} \ni(u, v, j) \mapsto \mathbf{x}_{m}^{j}(u, v)=\sum_{\ell=0}^{L} N^{\ell}(u, v, j) \mathbf{B}_{m}^{\ell} \tag{2.6}
\end{equation*}
$$

Without loss of generality, we may assume that the functions $N^{\ell}$ are linearly independent. Otherwise, the setup can be reduced without altering the properties of the scheme. The $n$ parts $\mathbf{x}_{m}^{0}, \ldots, \mathbf{x}_{m}^{n-1}$ forming $\mathbf{x}_{m}$ are referred to as segments. Collecting the functions $N^{\ell}$ in a row vector $N$ and the control points in a column vector $\mathbf{B}_{m}$ yields the vector notation

$$
\begin{equation*}
\mathbf{x}_{m}(u, v, j)=\mathbf{x}_{m}^{j}(u, v)=N(u, v, j) \mathbf{B}_{m} \tag{2.7}
\end{equation*}
$$

The schemes to be considered here are linear and stationary, i.e. there exists a square subdivision matrix $A$ with

$$
\begin{equation*}
\mathbf{B}_{m+1}=A \mathbf{B}_{m}, \quad \mathbf{B}_{m}=A^{m} \mathbf{B}_{0} \tag{2.8}
\end{equation*}
$$

[^1]


Figure 1: Domain $U$ (left) and structure of surface layers $\mathrm{x}_{\mathrm{m}}$ (right).
Definition 2.1 Let the eigenvalues $\lambda_{0}, \ldots, \lambda_{L}$ of $A$ be ordered by modulus,

$$
\begin{equation*}
\left|\lambda_{0}\right| \geq\left|\lambda_{1}\right| \geq \cdots \geq\left|\lambda_{L}\right|, \tag{2.9}
\end{equation*}
$$

and denote by $\psi_{0}, \ldots, \psi_{L}$ the corresponding generalized real eigenvectors. If $\left|\lambda_{0}\right|>\left|\lambda_{1}\right|=$ $\left|\lambda_{2}\right|>\left|\lambda_{3}\right|$ then the characteristic map of the subdivision algorithm is defined by

$$
\begin{equation*}
\Psi: U \times \mathrm{Z}_{n} \mapsto N\left[\psi_{1}, \psi_{2}\right] \in \mathrm{R}^{2}, \tag{2.10}
\end{equation*}
$$

or in complex form by

$$
\begin{equation*}
\Psi_{*}: U \times \mathrm{Z}_{n} \mapsto N \psi_{*} \in \mathrm{C}, \quad \psi_{*}:=\psi_{1}+i \psi_{2} . \tag{2.11}
\end{equation*}
$$

Remark i) $\left[\psi_{1}, \psi_{2}\right]$ is a $(L+1) \times 2$-matrix. Its rows play the role of 2 D control points. ii) Throughout, the subscript $*$ will indicale that the we refer to the complexification of a two-dimensional real variable or function. We will switch complex and real representation without further notice.

On the left hand side, Figure 2 shows a typical example of a characteristic map for $n=5$ as obtained for example by the Doo-Sabin algorithm. In order to guarantee affine invariance of the algorithm, the rows of $A$ must sum up to 1 . Thus, $[1, \ldots, 1]$ is always an eigenvector of $A$ to the eigenvalue 1. The following Theorem establishes a sufficient criterion for subdivision algorithms generating smooth limit surfaces.
Theorem 2.1 If $\lambda:=\lambda_{1}=\lambda_{2}, 1>\lambda>\left|\lambda_{3}\right|$, is a real eigenvalue with geometric multiplicity 2 , and if the characteristic map is regular and injective, then the limit surface y is a regular $C^{1}$-manifold for almost every choice of initial data $\mathbf{B}_{0}$.


Figure 2: Injective (left) and non-injective (right) characteristic map
A proof of this Theorem can be found in [Rei95a]. Generalizations, though not required here, are provided in [Rei95b] and [PR96]. Subsequently, it will be assumed that the eigenvalues of $A$ satisfy the assumptions of Theorem 2.1.

The following Theorem states a necessary criterion for the convergence of a subdivision scheme to smooth limit surfaces.

Theorem 2.2 If the characteristic map of a subdivision scheme is non-injective, i.e. there exist $(u, v, j) \neq\left(u^{\prime}, v^{\prime}, j^{\prime}\right)$ such that

$$
\begin{equation*}
\Psi(u, v, j)=\Psi\left(u^{\prime}, v^{\prime}, j^{\prime}\right) \tag{2.12}
\end{equation*}
$$

and if $\Psi(u, v, j)$ is an interior point of $\Psi\left(U, \mathrm{Z}_{\mathrm{n}}\right)$, then the limit surface y is not a regular $C^{1}$-manifold for almost every choice of initial data $\mathbf{B}_{0}$.

Proof Choose an $\varepsilon$-neighborhood $V_{\varepsilon}(\Psi(u, v, j))$ such that $V_{\varepsilon}(\Psi(u, v, j)) \subset \Psi\left(U, \mathrm{Z}_{n}\right)$. Then there exist neighborhoods $V$ and $V^{\prime}$ of $(u, v, j)$ and $\left(u^{\prime}, v^{\prime}, j^{\prime}\right)$, respectively, with $\Psi(V)=\Psi\left(V^{\prime}\right)=V_{\varepsilon}(\Psi(u, v, j))$. If $\tilde{\Psi}$ is a continuous map sufficiently close to $\Psi$, i.e. $\|\Psi-\tilde{\Psi}\|_{\infty}<\varepsilon / 2$, then $\Psi(V) \cap \tilde{\Psi}\left(V^{\prime}\right) \neq \emptyset$ and $\tilde{\Psi}$ is also not injective. Now, express $\mathbf{B}_{m}$ in terms of the generalized eigenvectors $\psi_{\ell}$,

$$
\begin{equation*}
\mathbf{B}_{0}=\sum_{\ell=0}^{L} \psi_{\ell} \mathbf{b}_{\ell}, \quad \mathbf{B}_{m}=\mathbf{b}_{0}+\lambda^{m}\left(\psi_{1} \mathbf{b}_{1}+\psi_{2} \mathbf{b}_{2}\right)+o\left(|\lambda|^{m}\right) . \tag{2.13}
\end{equation*}
$$

Then for almost every choice of initial data $\mathbf{B}_{0}$, the coefficients $\mathbf{b}_{1}$ and $\mathbf{b}_{2}$ are linearly independent, and we can choose coordinates such that $b_{0}=0$ is the origin and $b_{1}=$
$e_{1}, b_{2}=e_{2}$ are the first two unit vectors. A rescaling of the surface layers yields

$$
\begin{equation*}
\tilde{\mathbf{x}}_{m}:=\lambda^{-m} \mathbf{x}_{m}=\mathbf{e}_{1} \psi_{1}+\mathbf{e}_{2} \psi_{2}+o(1)=[\Psi, 0]+o(1) \quad \text { as } \quad m \rightarrow \infty \tag{2.14}
\end{equation*}
$$

Now, assume that $y$ is a regular $C^{1}$-manifold. Then the latter equation implies that the langent plane at the origin is the $x y$-plane. The projection $\tilde{\psi}_{m}$ of $\tilde{\mathbf{x}}_{m}$ on the $x y$-plane is converging to $\Psi$, so $\bar{\psi}_{m}$ is non-injective for $m$ sufficiently large. Consequently, the projection of the layers $\mathbf{x}_{m}$ to the $x y$-plane are non-injective near the origin for almost all $m$. This contradicts the assumption, since the projection of a regular $C^{1}$-manifold on its tangent plane is locally injective.

Finally, let us state two basic properties of characteristic maps. The first one is derived from the fact that $\Psi$ and $A \Psi=\lambda \Psi$ join smoothly,

$$
\begin{equation*}
\Psi^{j}(1, t)=\lambda \Psi^{j}(2,2 t), \quad \Psi^{j}(t, 1)=\lambda \Psi^{j}(2 t, 2), \quad t \in[0,1] . \tag{2.15}
\end{equation*}
$$

The second one expresses continuity between segments,

$$
\begin{equation*}
\Psi^{j}(0, t)=\Psi^{j+1}(t, 0), \quad t \in[1,2] \tag{2.16}
\end{equation*}
$$

## 3 Symmetry and Fourier Analysis

This section examines the special structure of the characteristic map for subdivision schemes obeying generic symmetry assumptions, namely that subdivision is independent of the particular labeling of control points used for refining the control mesh. According to the split of $\mathrm{x}_{m}$ into $n$ segments, the vectors $\mathrm{B}_{m}$ of control points can be divided into $n$ equally structured blocks,

$$
\begin{equation*}
\mathbf{B}_{m}:=\left[\mathbf{B}_{m}^{0} ; \ldots ; \mathbf{B}_{m}^{n-\mathbf{1}}\right] \tag{3.1}
\end{equation*}
$$

and $A$ is partitioned into $n \times n$ square blocks $A^{j, j^{\prime}}, j, j^{\prime} \in Z_{n}$.
Definition 3.1 A subdivision algorithm is called symmetric if it is invariant both under a shift $S$ and a reflection $R$ of the labeling of the vector $\mathbf{B}_{m}$ of control points. $S$ and $R$ are permutation matrices characterized by

$$
\begin{align*}
N(u, v, j+1) \mathbf{B}_{m} & \equiv N(u, v, j) S \mathbf{B}_{m}  \tag{3.2}\\
N(v, u,-j) \mathbf{B}_{m} & \equiv N(u, v, j) R \mathbf{B}_{m} \tag{3.3}
\end{align*}
$$

Invariance of the subdivision algorithm means that the subdivision matrix $A$ commutes both with $R$ and $S$, i.e.

$$
\begin{equation*}
S A=A S, \quad R A=A R \tag{3.4}
\end{equation*}
$$

With $S^{j, j^{\prime}}$ the segments of $S$ and $E$ the identity matrix of the same size as $S^{j, j^{\prime}}$, the shift matrix $S$ is given by

$$
\begin{equation*}
S^{j, j^{t}}=\delta_{j, j^{\prime}+1} E \tag{3.5}
\end{equation*}
$$

Comparison of $S A$ and $A S$ shows that the subdivision matrix of a symmetric scheme is block-cyclic, i.e.

$$
\begin{equation*}
A^{j}:=A^{j, 0}=A^{j+j^{\prime}, j^{\prime}}, \quad j, j^{\prime} \in \mathrm{Z}_{n} \tag{3.6}
\end{equation*}
$$

Thus, (2.8) becomes

$$
\begin{equation*}
\mathbf{B}_{m+1}^{j}=\sum_{j^{\prime}=0}^{n-1} A^{j^{\prime}} \mathbf{B}_{m}^{j^{\prime}+j}, \quad j \in \mathrm{Z}_{n} \tag{3.7}
\end{equation*}
$$

With

$$
\begin{equation*}
\omega_{n}:=c_{n}+i s_{n}:=\exp (2 \pi i / n) \tag{3.8}
\end{equation*}
$$

define the discrete Fourier transform by

$$
\begin{equation*}
\hat{p}^{k}:=\sum_{j=0}^{n-1} \omega_{n}^{-j k} p^{j}, \quad k \in Z_{n} \tag{3.9}
\end{equation*}
$$

Here, $p=\left[p^{0}, \ldots, p^{n-1}\right]$ is an $n$-vector in the generalized sense that its entries $p^{j}$ can be either scalars, or vectors, or matrices. It will always become clear from the context, what is meant. Applying the discrete Fourier transform to (3.7) yields

$$
\begin{equation*}
\hat{\mathbf{B}}_{m+1}^{k}=\hat{A}^{k} \hat{\mathbf{B}}_{m}^{k} \tag{3.10}
\end{equation*}
$$

see [Lip81] for a comprehensive introduction to the Fourier analysis of cyclic systems.
Theorem 3.1 The characteristic map is non-injective unless the dominant eigenvalue $\lambda$ is an eigenvalue of $\hat{A}^{1}$ and $\hat{A}^{n-1}$.

Proof $\lambda$ is an eigenvalue of $A$ if and only if it is an eigenvalue of $\hat{A}^{k}$ for some $k \in$ $\{0, \ldots, n-1\}$. If $\lambda$ is an eigenvalue of $\hat{A}^{k}$ then it is also an eigenvalue of $\hat{A}^{n-k}$, since $A$ is real and $\hat{A}^{n-k}=\overline{\hat{A}}^{k}$. Let $\hat{A}^{k} \hat{\psi}=\lambda \hat{\psi}$, then

$$
\begin{equation*}
\psi_{*}:=\left[\omega_{n}^{0} \hat{\psi}, \omega_{n}^{k} \hat{\psi}, \ldots, \omega_{n}^{k(n-1)} \hat{\psi}\right] \tag{3.11}
\end{equation*}
$$

is a complex eigenvector of $A$. Consequently, the segments $\Psi^{j}$ of the complex characteristic map satisfy

$$
\begin{equation*}
\Psi_{*}^{j}=\omega_{n}^{j k} \Psi_{*}^{0} \tag{3.12}
\end{equation*}
$$

Now, the winding number of the closed curve

$$
\begin{equation*}
\alpha_{*}(t):[0, \pi / 2) \times \mathrm{Z}_{n} \ni(t, j) \mapsto \Psi_{\star}^{j}(2 \cos (t), 2 \sin (t)) \tag{3.13}
\end{equation*}
$$

is either $k$ or $n-k$. So, if $k \notin\{1, n-1\}$ the curve $\alpha_{\text {* }}$ has self-intersections implying that $\Psi_{\star}$ is not injective.

The effect of a dominant eigenvalue stemming from the wrong Fourier component is depicted in Figure 2 on the right hand side. It shows the non-injective characteristic map for the Doo-Sabin algorithm for $n=5$ with weights chosen such that $\lambda$ is an eigenvalue of $\hat{A}^{2}$ and $\hat{A}^{4}$. As a consequence of Theorem 3.1 , it will be assumed that $\lambda$ is an eigenvalue of $\hat{A}^{1}$ and $\hat{A}^{n-1}$, subsequently. So, (3.12) becomes

$$
\begin{equation*}
\Psi_{*}^{j}=\omega_{n}^{j} \Psi_{*}^{0} \tag{3.14}
\end{equation*}
$$

The following Lemma is the key to reducing the analysis of the characteristic map $\Psi$ to the examination of a single segment, say $\Psi^{0} . \Psi$ is called normalized if $\hat{\psi}$ is scaled such that $\Psi^{0}(2,2)=(d, 0)$ with $d>0$. Note that normalization is always possible if $\Psi$ is injective, since then $\Psi^{0}(2,2) \neq \Psi^{0}(1,1)=\lambda \Psi^{0}(2,2)$.

Lemma 3.1 If $\Psi_{.}$is a normalized characteristic map of a symmetric scheme, then

$$
\begin{equation*}
\Psi_{\star}^{j}(u, v)=\bar{\Psi}^{-j}(v, u), \tag{3.15}
\end{equation*}
$$

and in particular

$$
\begin{equation*}
\Psi^{0}(u, v)=\bar{\Psi}_{*}^{0}(v, u) \tag{3.16}
\end{equation*}
$$

Proof (3.2) and (3.3) yield $N(u, v, j) S^{-1} R=N(u, v, j) R S$ implying that

$$
\begin{equation*}
S^{-1} R=R S \tag{3.17}
\end{equation*}
$$

since the functions forming $N$ are assumed to be linearly independent. From (3.4) one concludes that $R \psi_{*}$ is an eigenvector of $A$ to the eigenvalue $\lambda$, i.e. it can be written as

$$
\begin{equation*}
R \psi_{*}=a \psi_{*}+b \bar{\psi}_{*}, \quad a, b \in \mathrm{C} \tag{3.18}
\end{equation*}
$$

On using $S \psi_{=}=\omega_{n} \psi_{m}$, one obtains from the latter two equations

$$
\begin{equation*}
S^{-1} R \psi_{m}=a \bar{\omega}_{n} \psi_{m}+b \omega_{n} \bar{\psi}_{m}=a \omega_{n} \psi_{m}+b \omega_{n} \bar{\psi}_{m}=R S \psi_{m} \tag{3.19}
\end{equation*}
$$

Since $\psi_{*}$ and $\bar{\psi}_{*}$ are linearly independent, this implies $a=0$, hence $R \psi_{*}=b \overline{\psi_{*}}$. In order to determine $b$, consider $\Psi^{0}(2,2)$. By (3.3),

$$
\begin{equation*}
d=N(2,2,0) \psi_{*}=N(2,2,0) R \psi_{*}=b N(2,2,0) \bar{\psi}_{*}=b d \tag{3.20}
\end{equation*}
$$

thus $b=1$ and $R \psi_{*}=\bar{\psi}_{*}$. Finally, we obtain

$$
\begin{equation*}
\Psi_{\mathbf{z}}^{j}(u, v)=N(u, v, j) \psi_{-}=N(v, u,-j) \bar{\psi}_{m}=\bar{\Psi}_{\pi}^{-j}(v, u) \tag{3.21}
\end{equation*}
$$

## 4 Criteria for regularity and injectivity

In this section we derive a sequence of Lemmas resulting in two sufficient criteria for the regularity and injectivity of the characteristic map that can be verified efficiently. Throughout, it will be assumed that $\Psi$ is a normalized characteristic map of a symmetric subdivision scheme.

The first Lemma states that for regular functions injectivity is equivalent to injectivity at the boundary.

Lemma 4.1 Denote by $\partial U$ the boundary of $U$ and by $\Psi_{a}^{0}$ the restriction of $\Psi^{0}$ to $\partial U$. If $\Psi^{0}=\left[\Psi_{1}^{0}, \Psi_{2}^{0}\right]$ is regular, i.e.

$$
D \Psi^{0}:=\left(\begin{array}{cc}
\Psi_{1, u}^{0} & \Psi_{1, v}^{0}  \tag{4.1}\\
\Psi_{2, u}^{0} & \Psi_{2, v}^{0}
\end{array}\right), \quad J^{0}:=\operatorname{det} D \Psi^{0} \neq 0
$$

then $\Psi^{0}$ is injective if and only if $\Psi_{\partial}^{0}$ is injective.
Proof Assume that $\Psi^{0}$ is regular and $\Psi_{a}^{0}$ is injective. By the Inverse Function Theorem (IFT), points in the interior $\stackrel{\circ}{U}$ of $U$ are mapped to points in the interior of $\Psi^{0}(U)$, i.e.

$$
\begin{equation*}
\partial \Psi^{0}(U) \cap \Psi^{0}(\stackrel{\circ}{U})=\emptyset \tag{4.2}
\end{equation*}
$$

Define the function $\mu$ assigning the number of pre-images to the points in $\Psi^{0}(U)$,

$$
\begin{equation*}
\mu: \Psi^{0}(U) \ni(x, y) \mapsto \#\left\{(u, v) \in U: \Psi^{0}(u, v)=(x, y)\right\} \in \mathrm{N} . \tag{4.3}
\end{equation*}
$$

Injectivity of $\Psi_{a}^{0}$ and (4.2) imply $\mu\left(\partial \Psi^{0}(U)\right)=1 . \mu$ is upper semi-continuous by the IFT, hence $\mu\left(\Psi^{0}(U)\right)=1$, i.e. $\Psi$ is injective.

The second Lemma gives a sufficient condition for $\Psi_{\infty}^{0}$ being located in a sector of angle $2 \pi / n$ in the complex plane.

Lemma 4.2 If $\Psi^{0}$ is regular and

$$
\begin{equation*}
\Psi_{1, v}^{0}(1, t)>0, \quad \Psi_{2, v}^{0}(1, t)>0 \tag{4.4}
\end{equation*}
$$

for all $t \in[0,1]$, then

$$
\begin{equation*}
-\pi / n \leq \arg \Psi^{0}(u, v) \leq \pi / n \tag{4.5}
\end{equation*}
$$

for all $(u, v) \in U$.

Proof By Lemma 3.1, we have

$$
\begin{equation*}
\Psi_{1}^{0}(u, v)=\Psi_{1}^{0}(v, u), \quad \Psi_{2}^{0}(u, v)=-\Psi_{2}^{0}(v, u) \tag{4.6}
\end{equation*}
$$

and in particular $\Psi_{2}^{0}(t, t)=0, t \in[1,2]$. Let $p_{\mathbf{a}}(t):=p_{1}(t)+i p_{2}(t):=\Psi_{\mathbf{a}}^{0}(1, t), t \in[0,1]$, then $p_{2}$ is monotonically increasing because $\Psi_{2, v}^{0}>0$ and

$$
\begin{equation*}
\Psi_{2}^{0}(1,0)=p_{2}(0)<p_{2}(t)<p_{2}(1)=\Psi_{2}^{0}(1,1)=0, \quad t \in(0,1) \tag{4.7}
\end{equation*}
$$

Now, consider $\Psi^{0}(t, 0), t \in[1,2]$. By (3.12), (2.16) and (3.16) we obtain for $t \in[1,2]$

$$
\begin{equation*}
\omega_{n} \Psi_{\star}^{0}(t, 0)=\Psi_{\star}^{1}(t, 0)=\Psi_{*}^{0}(0, t)=\bar{\Psi}_{*}^{0}(t, 0) . \tag{4.8}
\end{equation*}
$$

This implies either $\arg \Psi_{*}^{0}(t, 0)=-\pi / n$ or $\arg \Psi_{*}^{0}(0, t)=\pi-\pi / n$. The second case contradicts (4.7), thus

$$
\begin{equation*}
\arg \Psi_{( }^{0}(t, 0)=-\pi / n, \quad \arg \Psi_{*}^{0}(0, t)=\pi / n \tag{4.9}
\end{equation*}
$$

This means that $\Psi^{0}(t, 0)$ is a part of the straight half line $h(r)=r \exp (-i \pi / n), r>$ 0 . Since $p_{\infty}$ is monotonically increasing in both real and imaginary part, it has no intersections with $h$ except for $p_{*}(0)$, hence

$$
\begin{equation*}
-\pi / n=\arg p_{*}(0)<\arg p_{*}(t)<\arg p_{*}(1)=0, \quad t \in(0,1) . \tag{4.10}
\end{equation*}
$$

Using the scaling property (2.15) and symmetry with respect to the real axis, the latter two equations imply that (4.5) holds for all $(u, v) \in \partial U$. By the IFT, we have $\partial \Psi_{u}^{0}(U) \subset$ $\Psi^{0}(\partial U)$, i.e. $-\pi / n \leq \arg \partial \Psi^{0}(U) \leq \pi / n$. Since $\Psi^{0}(U)$ is compact, this implies $-\pi / n \leq$ $\arg \Psi^{\mathrm{o}}(U) \leq \pi / n$ as asserted.

The third Lemma provides a sufficient condition on the partial derivatives of $\Psi^{0}$ that imply injectivity.

Lemma 4.3 If $\Psi^{0}$ is regular and $\Psi_{1, v}^{0}(1, t), \Psi_{2, v}^{0}(1, t)>0, t \in[0,1]$, then $\Psi^{0}$ is injective.
Proof By Lemma 4.1 it suffices to show that the restriction $\Psi_{\partial}^{0}$ of $\Psi^{0}$ to the boundary of $U$ is injective. Let

$$
\begin{array}{lll}
p_{\star}^{1}(t):=\Psi_{m}^{0}(1, t), & p_{\approx}^{2}(t):=\Psi_{-}^{0}(1+t, 0), & p_{m}^{3}(t):=\Psi_{*}^{0}(2,2 t) \\
p_{m}^{4}(t):=\Psi_{m}^{0}(t, 1), & p_{*}^{5}(t):=\Psi_{\approx}^{0}(0,1+t), & p_{m}^{6}(t):=\Psi_{*}^{0}(2 t, 2) \tag{4.12}
\end{array}
$$

for $t \in[0,1]$, see Figure 3, then

$$
\begin{gather*}
p_{*}=p_{*}^{1}=\bar{p}_{\pi}^{4}=\lambda p_{\pi}^{3}=\lambda \bar{p}_{*}^{6}  \tag{4.13}\\
p_{*}^{2}=\bar{p}_{\pi}^{5} \tag{4.14}
\end{gather*}
$$



Figure 3: Curves $p_{*}^{1}, \ldots, p_{*}^{\mathbf{6}}$.
with $p_{*}=p_{1}+i p_{2}$ defined as in the Proo[ of Lemma 4.2. $p_{*}^{2}$ and $p_{*}^{5}$ do not intersect, since

$$
\begin{equation*}
\arg p_{n}^{2}=-\pi / n, \quad \arg p_{n}^{5}=\pi / n \tag{4.15}
\end{equation*}
$$

by (4.9). Both curves also do not have self-intersections, since they are regularly parametrized and parts of straight lines. Next, we show that $\arg p_{-}(t)$ is monotonically increasing in $t$. By (4.7) and (4.9),

$$
\begin{equation*}
p_{2}(0)<0, \quad p_{2}(1)=0, \quad p_{1}(0)=p_{2}(0) \sin (-\pi / n)>0 . \tag{4.16}
\end{equation*}
$$

By assumption, $p_{1}$ and $p_{2}$ are monotone increasing, thus $p_{1}>0$ and $p_{2} \leq 0$. This implies

$$
\begin{equation*}
\frac{d}{d t}\left(\arg p_{z}\right)=\frac{p_{1} p_{2}^{\prime}-p_{1}^{\prime} p_{2}}{p_{1}^{2}+p_{2}^{2}}>0 \tag{4.17}
\end{equation*}
$$

as announced. Monotonicity of $\arg p_{*}=\arg p^{1}=\arg p^{3}=-\arg p^{4}=-\arg p^{6}$ has the following consequences: First, it guarantees that $p_{*}^{1}, p_{*}^{3}, p_{*}^{4}, p_{\pi}^{6}$ do not have selfintersections. Second, it excludes intersections of $p_{*}^{1}$ and $p_{*}^{3}$, since $p_{-}^{1}=\lambda p_{m}^{3}, \lambda \neq 1$. Analogously, $p_{*}^{4}$ and $p_{*}^{6}$ are disjoint. Third, the only intersections of $p_{*}^{1}$ and $p_{*}^{3}$ with $p_{m}^{2}$ are

$$
\begin{align*}
& \Psi_{\star}^{0}(1,0)=p_{\pi}^{1}(0)=p_{*}^{2}(0)  \tag{4.18}\\
& \Psi_{\approx}^{0}(2,0)=p_{\pi}^{3}(0)=p_{\approx}^{2}(1) \tag{4.19}
\end{align*}
$$

and analogously for $p_{m}^{4}, p_{m}^{5}, p_{m}^{6}$. Fourth, the only intersections of $p_{\star}^{1} \cup p_{*}^{2} \cup p_{*}^{3}$ and $p_{*}^{4} \cup p_{*}^{5} \cup p_{*}^{6}$ are

$$
\begin{align*}
& \Psi_{\approx}^{0}(1,1)=p_{\star}^{1}(1)=p_{\approx}^{4}(1)  \tag{4.20}\\
& \Psi_{\approx}^{0}(2,2)=p_{\pi}^{3}(1)=p_{\pi}^{6}(1) \tag{4.21}
\end{align*}
$$

and the proof is complete.

The following Theorem establishes a sufficient criterion for the partial derivatives of $\Psi^{0}$ which guarantees regularity and injectivity of the characteristic map. Its usefulness is due to the fact that it requires only estimates for the partial derivatives of a single segment $\Psi^{0}$. Since for generalized B-spline subdivision schemes the functions in questions are piecewise polynomial, the conditions can be verified numerically or even analytically using B-spline representations and the convex hull property.

Theorem 4.1 If $\Psi^{0}$ is regular and $\Psi_{1, v}^{0}(1, t), \Psi_{2, v}^{0}(1, t)>0, t \in[0,1]$, then the characteristic map $\Psi$ is regular and injective.

Proof By Lemma 4.3, $\Psi^{0}$ is regular and injective. (3.14) says, that $\Psi^{j}$ is obtained from $\Psi^{0}$ by a $2 \pi j / n$-rotation about the origin. So, each $\Psi^{j}, j \in Z_{n}$ is regular and injective. Further, the segments $\Psi^{j}$ do not overlap, since Lemma 4.2 yields

$$
\begin{equation*}
(2 j-1) \pi / n \leq \arg \Psi_{=}^{j} \leq(2 j+1) \pi / n, \quad j \in Z_{n} . \tag{4.22}
\end{equation*}
$$

The assumptions of the following Corollary are stronger than those of Theorem 4.1, but can be verified with less effort since no products of partial derivatives are involved.

Corollary 4.1 If $\Psi_{1, v}^{0}(u, v), \Psi_{2, v}^{0}(u, v)>0$ for all $(u, v) \in U$, then the characteristic map $\Psi$ is regular and injective.

Proof The symmetry relation (4.6) yields

$$
D \Psi^{0}(u, v)=\left(\begin{array}{rr}
\Psi_{1, v}^{0}(v, u) & \Psi_{1, v}^{0}(u, v)  \tag{4.23}\\
-\Psi_{2, v}^{0}(v, u) & \Psi_{2, v}^{0}(u, v)
\end{array}\right),
$$

so $J^{0}>0$ if $\Psi_{1, v}^{0}, \Psi_{2, v}^{0}>0$.


Figure 4: Mesh refinement by the Doo-Sabin algorithm.


Figure 5: Masks for the Doo-Sabin algorithm.

## 5 The Doo-Sabin algorithm

### 5.1 Algorithm

The Doo-Sabin algorithm is a generalization of the subdivision scheme for biquadratic tensor product B -splines. For each $n$-gon of the original mesh, a new, smaller $n$-gon is created and connected suitably with its neighbors, see Figure 4. Figure 5 shows the mask for generating a new $n$-gon from an old one for the regular case $n=4$ (left), and the general case (right). The weights suggested by Doo and Sabin in [DS78] are

$$
\begin{equation*}
\alpha^{j}=\frac{\delta_{j, 0}}{4}+\frac{3+2 \cos (2 \pi j / n)}{4 n} \tag{5.1}
\end{equation*}
$$

Below we analyze more general schemes assuming beforehand nothing but affine invariance and symmetry,

$$
\begin{equation*}
\sum_{j=0}^{n-1} \alpha_{n}^{j}=1, \quad \alpha^{j}=\alpha^{n-j} \tag{5.2}
\end{equation*}
$$



Figure 6: Labeling of control points for the Doo-Sabin algorithm

### 5.2 Characteristic map

Each of the $n$ segments $\mathbf{x}_{m}^{j}, j \in Z_{n}$, of the surface layers generated by the Doo-Sabin algorithm consists of 3 biquadratic B -spline patches. Accordingly, the $n$ blocks $\mathbf{B}_{m}^{j}$ forming the vector of control points $\mathbf{B}_{m}$ consist of 9 elements, each. The labeling is shown in Figure 6. The $9 \times 9$-matrices $\hat{A}^{k}, k \in Z_{n}$, as introduced in (3.10) have the following structure,

$$
\hat{A}^{k}=\left(\begin{array}{ccc}
\hat{\alpha}^{k} & 0 & 0  \tag{5.3}\\
\hat{A}_{1,0}^{k} & \hat{A}_{1,1}^{k} & 0 \\
\hat{A}_{2,0}^{k} & \hat{A}_{2,1}^{k} & 0
\end{array}\right) .
$$

With $p:=9 / 16, q:=3 / 16, r:=1 / 16$, the sub-matrices are given by

$$
\begin{array}{ll}
\hat{A}_{10}^{k}=\left(\begin{array}{c}
p+\bar{\omega}_{n} q \\
p \\
p+\omega_{n} q
\end{array}\right), & \hat{A}_{11}^{k}=\left(\begin{array}{ccc}
q & 0 & \bar{\omega}_{n} r \\
q & r & q \\
\omega_{n} r & 0 & q
\end{array}\right) \\
\hat{A}_{20}^{k}=\left(\begin{array}{c}
q+\bar{\omega}_{n} r \\
q \\
r \\
q \\
q+\omega_{n} r
\end{array}\right), & \hat{A}_{21}^{k}=\left(\begin{array}{ccc}
p & 0 & \bar{\omega}_{n} q \\
p & q & r \\
q & p & q \\
r & q & p \\
\omega_{n} q & 0 & p
\end{array}\right) . \tag{5.4}
\end{array}
$$

The matrix $A_{1, \mathrm{I}}^{k}$ has eigenvalues $1 / 4,1 / 8,1 / 16$, hence each of them is an $n$-fold eigenvalue of the subdivision matrix $A$. Further, $A$ has a $5 n$-fold eigenvalue 0 stemming from the $5 \times 5$-zero submatrix of $\hat{A}^{k}$. Due to their high multiplicity, these eigenvalues do not come into account for playing the role of the dominant eigenvalue $\lambda$. The only eigenvalues left are the upper left entries $\hat{\alpha}^{k}, k \in Z_{n}$, of $\hat{A}^{k}$ obtained by applying the discrete Fourier transform to the vector $\left[\alpha^{0}, \ldots, \alpha^{n-1}\right]$ of weights for the $n$-gon. Since the $\alpha^{j}$ sum up to 1 , we have $\lambda_{0}:=\hat{\alpha}_{n}^{0}=1$. Due to symmetry, the remaining eigenvalues are real and occur in pairs according to $\hat{\alpha}_{n}^{k}=\hat{\alpha}_{n}^{n-k}$. From the theory developed in the preceding sections we know that

$$
\begin{equation*}
\lambda:=\hat{\alpha}_{n}^{1}=\hat{\alpha}_{n}^{n-1} \tag{5.5}
\end{equation*}
$$

must satisfy

$$
\begin{equation*}
1>\lambda>\max \left\{1 / 4,\left|\hat{\alpha}_{n}^{2}\right|, \ldots,\left|\hat{\alpha}_{n}^{n-2}\right|\right\} \tag{5.6}
\end{equation*}
$$

The eigenvector of the Matrix $\hat{A}^{1}$ corresponding to $\lambda$ is

$$
\psi=\left(\begin{array}{c}
2 \lambda(16 \lambda-1)(8 \lambda-1)(4 \lambda-1)  \tag{5.7}\\
6 \lambda(16 \lambda-1)\left(6 \lambda-1+2 \bar{\omega}_{n} \lambda\right) \\
18 \lambda\left(32 \lambda^{2}-1+4 c_{n} \lambda\right) \\
6 \lambda(16 \lambda-1)\left(6 \lambda-1+2 \omega_{n} \lambda\right) \\
(16 \lambda-1)\left(12 \lambda^{2}+18 \lambda-3+\bar{\omega}_{n}\left(4 \lambda^{2}+12 \lambda-1\right)\right) \\
6 \lambda\left(32 \lambda^{2}+64 \lambda-12+c_{n}(20 \lambda+1)-i s_{n}(16 \lambda-1)\right) \\
64 \lambda^{3}+512 \lambda^{2}-46 \lambda-8+36 c_{n} \lambda(2 \lambda+1) \\
6 \lambda\left(32 \lambda^{2}+64 \lambda-12+c_{n}(20 \lambda+1)+i s_{n}(16 \lambda-1)\right) \\
(16 \lambda-1)\left(12 \lambda^{2}+18 \lambda-3+\omega_{n}\left(4 \lambda^{2}+12 \lambda-1\right)\right)
\end{array}\right) .
$$

Note that the characteristic map depends only on the values of $\lambda$ and $n$. That is, all masks $\alpha$ with identical first Fourier component yield the same characteristic map.

### 5.3 Verification

Let us start with briefly discussing the particular case $\lambda=1 / 2$ as obtained in particular for the weights in (5.1). Here we obtain the eigenvector

$$
\psi_{*}=3\left(\begin{array}{ccc}
7 & 14+7 \omega_{n} & 21+14 \omega_{n}  \tag{5.8}\\
14+7 \bar{\omega}_{n} & 21+6 c_{n} & 28+2 \bar{\omega}_{n}+9 \omega_{n} \\
21+14 \bar{\omega}_{n} & 28+2 \omega_{n}+9 \bar{\omega}_{n} & 35+12 c_{n}
\end{array}\right)
$$

The segment $\Psi_{*}^{0}$ of the characteristic map consists of three bi-quadratic patches, which can be expressed in Bernstein-Bézier form with the following coefficients,

Computing the partial derivative $\Psi_{u, v}^{0}=\Psi_{1, v}^{0}+i \Psi_{2, v}^{0}$ with respect to $v$ yields three quadratic-linear patches with coefficients


Both the real and the imaginary part of the coefficients are positive. So, by the convex hull property and Corollary 4.1 the algorithm is verified to generate smooth limit surfaces.

The situation for general $\lambda$ is more subtle, in particular as $\lambda \rightarrow 1$. First, Corollary 4.1 turns out to be insufficient. Second, there exists a limit value $\lambda_{\max }(n)<1$ depending on $n$ such that even the assumptions of Theorem 4.1 are not fulfilled for $1>\lambda>$ $\lambda_{\max }$. It will be shown in the next subsection that this is due to an actual loss of smoothness as $\lambda$ passes the bound. All formulas required here were derived using a computer algebra system. They are partially rather lengthy and will not be stated explicitly unless necessary. Rather, we depict the crucial results graphically.

In order to apply Theorem 4.1, we have to compute $J^{0}$, i.e. the determinant of the Jacobian of $\Psi^{0} . J^{0}$ is a continuous, piecewise bi-cubic function over $U$ thath can be expressed in Bernstein-Bézier form with $3 \times 16$ coefficients $J_{\mu}^{0}, \mu=1, \ldots, 48$ depending on $n$ and $\lambda$. Explicit calculation shows that all coefficients $J_{\mu}^{0}$ are of type

$$
\begin{equation*}
J_{\mu}^{0}\left(c_{n}, \lambda\right)=s_{n}(16 \lambda-1)\left(P_{\mu}(\lambda)+c_{n} Q_{\mu}(\lambda)\right) \tag{5.11}
\end{equation*}
$$

with $P_{\mu}, Q_{\mu}$ polynomials of degree $\leq 6$ in $\lambda$. We give the coefficient corresponding to $J^{0}(1,1)$ the index $\mu=1$, i.e.

$$
\begin{equation*}
J^{0}(1,1)=J_{1}^{0}=s_{n}(16 \lambda-1)\left(P_{1}+c_{n} Q_{1}\right) \tag{5.12}
\end{equation*}
$$

The polynomials $P_{1}$ and $Q_{1}$ are

$$
\begin{align*}
P_{1}(\lambda) & :=96 \lambda^{3}\left(-128 \lambda^{3}+128 \lambda^{2}-7 \lambda-2\right)  \tag{5.13}\\
Q_{1}(\lambda) & :=864 \lambda^{4} . \tag{5.14}
\end{align*}
$$

In order to apply analytic tools, it is convenient to consider $c_{n}$ as a free variable varying in the interval $c_{n} \in[-1 / 2,1]$ corresponding to $n \geq 3$. For fixed $\lambda \in(1 / 4,1)$ there is at most one value $\tilde{c}_{n}$ where $J_{\mu}^{0}$ changes sign,

$$
\begin{equation*}
\tilde{c}_{n}=R_{\mu}(\lambda):=-P_{\mu}(\lambda) / Q_{\mu}(\lambda) \tag{5.15}
\end{equation*}
$$

Figure 7 shows a plot of all these functions as well as a magnification of the significant region. From the analysis of the case $\lambda=1 / 2$ we know that $J_{\mu}^{0}\left(c_{n}, 1 / 2\right)>0$ for all $\mu$. Thus, $J^{0}$ is positive as long as ( $\lambda, c_{n}$ ) lies in the shaded region, which is bounded by $R_{1}(\lambda)$. More precisely, the feasible set for ( $\lambda, c_{n}$ ) providing positivity of $J^{0}$ is

$$
\begin{equation*}
\Lambda:=\left\{\left(\lambda, c_{n}\right) \in(1 / 4,1) \times[-1 / 2,1]: J_{1}^{0}\left(c_{n}, \lambda\right)>0\right\} \tag{5.16}
\end{equation*}
$$

For verifying the assumptions of Theorem 4.1 it remains to show that $\Psi_{1, v}^{0}(1, t), \Psi_{2, v}^{0}(1, t)>$ 0 for $t \in[0, I]$ and $\left(c_{n}, \lambda\right) \in \Lambda$. Note that both functions are linear in $t$. So, it suffices to check positivity for $t \in\{0,1\}$, which follows immediately from

$$
\begin{align*}
& \Psi_{1, v}^{0}(1,0)=2 \lambda(16 \lambda-1)(4 \lambda-1)(4 \lambda+1)\left(1-c_{n}\right)  \tag{5.17}\\
& \Psi_{2, v}^{0}(1,0)=2 \lambda(16 \lambda-1)(4 \lambda-1)(4 \lambda+1) s_{n}  \tag{5.18}\\
& \Psi_{2, v}^{0}(1,1)=12 s_{n} \lambda^{2}(16 \lambda-1)  \tag{5.19}\\
& \Psi_{1, v}^{0}(1,1)=J_{1}^{0} / 2 \Psi_{2, v}^{0}(1,1) . \tag{5.20}
\end{align*}
$$

Finally, let us summarize the results derived in this section.
Theorem 5.1 Let $\alpha_{n}^{0}, \ldots, \alpha_{n}^{n-1}$ be symmetric weights for the Doo-Sabin algorithm. If $\lambda:=\hat{\alpha}_{n}^{1}=\hat{\alpha}_{n}^{n-1}$ satisfies

$$
\begin{gather*}
1>\lambda>\max \left\{1 / 4,\left|\hat{\alpha}_{n}^{2}\right|, \ldots,\left|\hat{\alpha}_{n}^{n-2}\right|\right\}  \tag{5.21}\\
128 \lambda^{2}(1-\lambda)-7 \lambda-2+9 \lambda \cos (2 \pi / n)>0 \tag{5.22}
\end{gather*}
$$

then the limit surface y is a regular $\mathrm{C}^{1}$-manifold for almost every choice of initial data $B_{0}$.


Figure 7: Feasible set $\Lambda$ and functions $R_{\mu}(\lambda)$.

### 5.4 Failure beyond the bound

In contrast to the lower bound $\lambda>1 / 4$, which appears naturally, the existence of an upper bound for $\lambda$ may surprise. It is not an artifact of the particular type of sufficient conditions in Theorem 4.1, but a sharp bound beyond which the Doo-Sabin algorithm provably fails. If $J_{1}^{0}<0$ then

$$
\begin{align*}
& \Psi_{2, v}^{0}(1,1)=12 s_{n} \lambda^{2}(16 \lambda-1)>0  \tag{5.23}\\
& \Psi_{1, v}^{0}(1,1)=J_{1}^{0} / 2 \Psi_{2, v}^{0}(1,1)<0 \tag{5.24}
\end{align*}
$$

Consider the curve $\left[g_{1}(t), g_{2}(t)\right]:=\Psi^{0}(t, t)$. Symmetry with respect to the $x$-axis implies $g_{2}(t) \equiv 0$. For the first component we obtain

$$
\begin{equation*}
g_{1}^{\prime}(1)=2 \Psi_{1, v}(1,1)<0, \quad g_{1}(1)=\Psi_{1}^{0}(1,1)<\lambda \Psi_{1}^{0}(2,2)=g_{1}(2) \tag{5.25}
\end{equation*}
$$

hence for each sufficiently small $\varepsilon>0$ there exists an $\varepsilon^{\prime}>\varepsilon$ such that $g_{1}(1+\varepsilon)=$ $g_{1}\left(1+\varepsilon^{\prime}\right)$. This implies the non-injectivity of the characteristic map $\Psi$,

$$
\begin{equation*}
\Psi^{0}(1+\varepsilon, 1+\varepsilon)=\Psi^{0}\left(1+\varepsilon^{\prime}, 1+\varepsilon^{\prime}\right), \quad \varepsilon \neq \varepsilon^{\prime} \tag{5.26}
\end{equation*}
$$

Moreover, for $\varepsilon$ sufficiently small, $J^{0}(1+\varepsilon, 1+\varepsilon)<0$ by continuity. So, $\Psi^{0}(1+\varepsilon, 1+\varepsilon)$ is an interior point of $\Psi\left(U, \mathrm{Z}_{n}\right)$ by the IFT, and the assumptions of Theorem 2.2 are fulfilled proving sharpness of the bound. Figure 8 shows a magnification of the characteristic map for $n=3$ in a vicinity of $\Psi^{0}(1,1)$ for $\lambda=1 / 2$ (left) and $\lambda=0.95$ (right). The latter


Figure 8: Characteristic map for $n=3$ and $\lambda=0.5$ (left) and $\lambda=0.95$ (right).
case corresponds to weights $\alpha^{0}=0.9667, \alpha^{1}=\alpha^{2}=0.01667$. Six layers of a subdivision surface generated by these weights are shown in Figure 9. The magnification on the right hand side is non-proportional, i.e. the 'height' of the surface has been expanded in order to depict its wavy shape. Let us conclude the discussion of the Doo-Sabin algorithm with a brief description of the qualitative and quantitative behavior of $\lambda_{\max }(n)$. As $n \rightarrow \infty, \lambda_{\max }(n)$ is increasing monotonically towards 1 . The asymptotic behavior for large $n$ is

$$
\begin{equation*}
\lambda_{\max }(n)=1-\frac{\pi^{2}}{7 n^{2}}+O\left(n^{-4}\right) \tag{5.27}
\end{equation*}
$$

The lowest bound occurs for $n=3$, namely

$$
\begin{equation*}
\lambda_{\max }(3)=\frac{\sqrt{187}}{24} \cos \left(\frac{1}{3} \arctan \left(\frac{27 \sqrt{5563}}{1576}\right)\right)+\frac{1}{3} \approx 0.8773 \tag{5.28}
\end{equation*}
$$

Table 1 lists the values of $\lambda_{\text {max }}$ for $n=3, \ldots, 12$.

## 6 The Catmull-Clark algorithm

### 6.1 Algorithm

The Catmull-Clark algorithm is a generalization of the subdivision scheme for bicubic tensor product B-splines. Each $n$-gon of the original mesh is subdivided into $n$ quadrilaterals thus generating a purely quadrilateral mesh after the first step. There are three masks for subdividing such a mesh, namely one for computing a new centroid, one for


Figure 9: Non-smooth surface generated by the Doo-Sabin algorithm with $\lambda=0.95$.

| $n$ | $\lambda_{\max }(n)$ |
| :---: | :---: |
| 3 | 0.8772872432 |
| 4 | 0.9223409885 |
| 5 | 0.9478340134 |
| 6 | 0.9628618416 |
| 7 | 0.9723152935 |


| $n$ | $\lambda_{\max }(n)$ |
| ---: | :---: |
| 8 | 0.9786066941 |
| 9 | 0.9829902941 |
| 10 | 0.9861607079 |
| 11 | 0.9885250887 |
| 12 | 0.9903339733 |

Table 1: Values of the bound $\lambda_{\max }(n)$ for $n=3, \ldots, 12$.


Figure 10: Mesh refinement by the Catmull-Clark algorithm.


Figure 11: Masks for the Catmull-Clark algorithm.
a new edge point, and one for the new location of a former vertex, see Figure 11. So, the variables at disposal are the weights

$$
\begin{equation*}
\alpha, \beta, \gamma, \quad \alpha+\beta+\gamma=1 \tag{6.1}
\end{equation*}
$$

In [CC78], Catmull and Clark suggest

$$
\begin{equation*}
\alpha=1-\frac{7}{4 n}, \beta=\frac{3}{2 n}, \gamma=\frac{1}{4 n} . \tag{6.2}
\end{equation*}
$$

### 6.2 Characteristic map

Each of the $n$ segments $x_{m}^{j}, j \in J_{n}$, of the surface layers generated by the Catmull-Clark algorithm consists of 3 bicubic $B$-spline patches. Accordingly, the $n$ blocks $B_{m}^{j}$ forming the vector of control points $\mathbf{B}_{m}$ consist of 13 elements, each. The labeling used here is shown in Figure 12. Note that the centroid

$$
\begin{equation*}
\mathbf{M}_{m}:=\mathbf{B}_{m}^{0}=\cdots=\mathbf{B}_{m}^{n-1} \tag{6.3}
\end{equation*}
$$

is replaced by $n$ identical copies in order to achieve the desired periodic structure. For all masks involving $\mathbf{M}_{m}$ we substitute

$$
\begin{equation*}
\mathbf{M}_{m}=\frac{1}{n} \sum_{j=0}^{n-1} \mathbf{B}_{m}^{j} \tag{6.4}
\end{equation*}
$$

The $13 \times 13$-matrices $\hat{A}^{k}, k \in J_{n}$, turn out to have the following structure,

$$
\hat{A}^{k}=\left(\begin{array}{ccc}
\hat{A}_{0,0}^{k} & 0 & 0  \tag{6.5}\\
\hat{A}_{1,0}^{k} & \hat{A}_{1,1}^{k} & 0 \\
\hat{A}_{2,0}^{k} & \hat{A}_{2,1}^{k} & 0
\end{array}\right)
$$



Figure 12: Labeling of control points for the Catmull-Clark algorithm.
With

$$
\begin{equation*}
p_{1}:=1 / 64, p_{2}:=3 / 32, p_{3}:=9 / 16, q_{1}:=1 / 16, q_{2}:=3 / 8, r:=1 / 4 \tag{6.6}
\end{equation*}
$$

the sub-matrices are given by

$$
\begin{gather*}
\hat{A}_{0,0}^{k}=\left(\begin{array}{cccc}
\alpha \delta_{k, 0} & \beta \delta_{k, 0} & \gamma \delta_{k, 0} \\
q_{2} \delta_{k, 0} & 2 q_{1} c_{n}+q_{2} & q_{1}\left(1+\bar{\omega}_{n}\right) \\
r \delta_{k, 0} & r\left(1+\omega_{n}\right) & r
\end{array}\right)  \tag{6.7}\\
\left(\begin{array}{cc}
\hat{A}_{\hat{1}, 0}^{k} & \hat{A}_{1,1}^{k} \\
\hat{A}_{2,0}^{k} & \hat{A}_{2,1}^{k}
\end{array}\right)=\left(\begin{array}{ccc|cccc}
q_{2} \delta_{k, 0} & q_{1}+q_{2} \omega_{n} & q_{1} & q_{1} & q_{1} \omega_{n} & 0 & 0 \\
p_{2} \delta_{k, 0} & 2 p_{1} c_{n}+p_{3} & p_{2}\left(1+\bar{\omega}_{n}\right) & p_{1} \bar{\omega}_{n} & p_{2} & p_{1} & 0 \\
q_{1} \delta_{k, 0} & q_{1} \omega_{n}+q_{2} & q_{2} & 0 & q_{1} & q_{1} & 0 \\
p_{1} \delta_{k, 0} & p_{2}\left(1+\omega_{n}\right) & p_{3} & p_{2} & p_{1}\left(1+\omega_{n}\right) & p_{2} & p_{1} \\
\hline 0 & q_{2} & q_{1}\left(1+\bar{\omega}_{n}\right) & q_{1} \bar{\omega}_{n} & q_{2} & q_{1} & 0 \\
0 & r & r & 0 & r & r & 0 \\
0 & q_{1} & q_{2} & q_{1} & q_{1} & q_{2} & q_{1} \\
0 & 0 & r & r & 0 & r & r \\
0 & q_{1} \omega_{n} & q_{2} & q_{2} & q_{1} \omega_{n} & q_{1} & q_{1} \\
0 & r \omega_{n} & r & r & T \omega_{n} & 0 & 0
\end{array}\right)
\end{gather*}
$$

The eigenvalues $1 / 8,1 / 16,1 / 32,1 / 64$ of the sub-matrix $\hat{A}_{1,1}^{k}$ are $n$-fold eigenvalues of $A$. Other non-zero eigenvalues come only from $\hat{A}_{0,0}^{k}$. For $k=0$ we obtain the obligatory
eigenvalue $\lambda_{0}=1$ and, letting $\gamma:=1-\alpha-\beta$,

$$
\begin{equation*}
\lambda_{1,2}^{0}=\left(4 \alpha-1 \pm \sqrt{(4 \alpha-1)^{2}+8 \beta-4}\right) / 8 \tag{6.8}
\end{equation*}
$$

which might be either both real or complex conjugate. For $k \neq 0$, the non-zeros eigenvalues of $\hat{A}_{0,0}^{k}$ are

$$
\begin{align*}
& \lambda_{1}^{k}=\left(c_{n, k}+5+\sqrt{\left(c_{n, k}+9\right)\left(c_{n, k}+1\right)}\right) / 16  \tag{6.9}\\
& \lambda_{2}^{k}=\left(c_{n, k}+5-\sqrt{\left(c_{n, k}+9\right)\left(c_{n, k}+1\right)}\right) / 16 \tag{6.10}
\end{align*}
$$

where $c_{n, k}:=\cos (2 \pi k / n)$. Let

$$
\begin{equation*}
\lambda:=\lambda_{1}^{1}=\lambda_{1}^{n-1}=\left(c_{n}+5+\sqrt{\left(c_{n}+9\right)\left(c_{n}+1\right)}\right) / 16 \tag{6.11}
\end{equation*}
$$

then straightforward calculus shows that for all $n$

$$
\begin{array}{rlrl}
1>\lambda>1 / 4>\lambda_{2}^{k}>1 / 8, & & k=1, \ldots, n-1 \\
& \lambda>\lambda_{1}^{k}>1 / 4, & & k=2, \ldots, n-2 \tag{6.13}
\end{array}
$$

Consequently, $\lambda$ is dominant if $\alpha, \beta, \gamma$ are chosen such that

$$
\begin{equation*}
\lambda>\max \left\{\left|\lambda_{1}^{0}\right|,\left|\lambda_{2}^{0}\right|\right\} \tag{6.14}
\end{equation*}
$$

In particular, this inequality holds for the original weights of Catmull-Clark (6.2), as can be verified by inspection. A characterization of feasible positive weights can be found in [BS88] ${ }^{2}$. For computing the characteristic map, the eigenvector $\hat{\psi}$ of $\hat{A}^{1}$ is partitioned into three blocks, $\hat{\psi}=\left[\hat{\psi}_{0}, \hat{\psi}_{1}, \hat{\psi}_{2}\right]$ according to the special structure of $\hat{A}^{1}$. Then $\hat{A}^{1} \hat{\psi}=\lambda \hat{\psi}$ is equivalent to

$$
\begin{align*}
\left(\hat{A}_{0,0}^{1}-\lambda\right) \hat{\psi}_{0} & =0 \\
\left(\hat{A}_{1,1}^{1}-\lambda\right) \hat{\psi}_{1} & =-\hat{A}_{1,0}^{1} \hat{\psi}_{0}  \tag{6.15}\\
\hat{\psi}_{2} & =\left(\hat{A}_{2,0}^{1} \hat{\psi}_{0}+\hat{A}_{2,1}^{1} \hat{\psi}_{1}\right) / \lambda
\end{align*}
$$

Now, $\hat{\psi}$ can be computed conveniently slarting from

$$
\begin{equation*}
\hat{\psi}_{0}:=\left[1+\bar{\omega}_{n}, 16 \lambda-2 c_{n}-6\right]^{T} \tag{6.16}
\end{equation*}
$$

which solves the first eigenvector equation. Note that the characteristic map depends only on $n$, and not on the particular choice of weights $\alpha, \beta, \gamma$ provided that ( 6.11 ) holds.

[^2]
### 6.3 Verification

Corollary 4.1 is sufficient for verifying the algorithm. One proceeds as follows:

1. For given $n \geq 3$, compute the dominant eigenvalue $\lambda$ according to (6.14) and the corresponding eigenvector $\hat{\psi}$ according to (6.15).
2. Express the three patches of the segment $\Psi^{0}$ of the characteristic map in BernsteinBézier form.
3. Compute the forward differences $\Delta_{\mu}, \mu=1, \ldots, 36$, of Bézier coefficients corresponding to the partial derivative with respect to $v$.
4. If all $\Delta_{\mu}, \mu=1, \ldots, 36$ are positive in both components, then by the convex hull property of the Bernstein-Bézier form the assumptions of Corollary 4.1 are fulfilled and the characteristic map is regular and injective.

This procedure can be run on a computer algebra system, but the resulting expressions are rather length, and discussing them is not very instructive. A numerical treatment is more convenient and yields equally reliable results, since only a finite number of quantities has to be checked for sign. The findings are summarized on Figure 13. The left and right hand side correspond to the two components of $\Delta_{\mu}$. The top row shows the values of all $\Delta_{\mu}$ for $n=3, \ldots, 20$. The bottom row shows the minimum of the $\Delta_{\mu}$ on a double-logarithmic scale for $n=10, \ldots, 10,000$, which should cover most cases of practical relevance. The positivity of all differences is evident. By Corollary 4.1, this proves smooth convergence of the Catmull-Clark algorithm provided that the inequality (6.11) holds.

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Figure 13: Bézier coefficients of the partial derivatives $\Delta_{\mu}$.
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[^0]:    -NSF National Young Investigator Award $9457806-\mathrm{CCR}$
    ${ }^{\dagger}$ BMBF Projekt 03-HO7STU-2

[^1]:    ${ }^{1}$ The intersection consists exclusively of points on the common boundary curve, which are identified.

[^2]:    ${ }^{2}$ The result in the reference is incorrect for $n=3$ and certain $\alpha, \beta$ yielding complex eigenvalues $\lambda_{1,2}^{0}$.

