

ANALYSIS OF THE HETEROGENEOUS MULTISCALE METHOD FOR ELLIPTIC HOMOGENIZATION PROBLEMS

WEINAN E, PINGBING MING, AND PINGWEN ZHANG

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1. INTRODUCTION AND MAIN RESULTS

1.1. **General methodology.** Consider the classical elliptic problem

$$(1.1) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x)\nabla u^\varepsilon(x)) = f(x) & x \in D \subset \mathbb{R}^d, \\ u^\varepsilon(x) = 0 & x \in \partial D. \end{cases}$$

Here ε is a small parameter that signifies explicitly the multiscale nature of the coefficient $a^\varepsilon(x)$. Several classical multiscale methodologies have been developed for the numerical solution of this elliptic problem, the most well known among which is the multigrid technique [8]. These classical multiscale methods are designed to resolve the details of the fine scale problem (1.1) and are applicable for general problems, i.e., no special assumptions are required for the coefficient $a^\varepsilon(x)$. In contrast modern multiscale methods are designed specifically for recovering partial information about u^ε at a sublinear cost, i.e., the total cost grows sublinearly with the cost of solving the fine scale problem [18]. This is only possible by exploring the special features that $a^\varepsilon(x)$ might have, such as scale separation. The simplest example is when

$$(1.2) \quad a^\varepsilon(x) = a\left(x, \frac{x}{\varepsilon}\right),$$

where $a(x, y)$ can either be periodic in y , in which case we assume the period to be $I = [-1/2, 1/2]^d$, or random but stationary under shifts in y , for each fixed $x \in D$. In both cases, it has been shown that [5, 36]

$$(1.3) \quad \|u^\varepsilon(x) - U(x)\|_{L^2(D)} \rightarrow 0,$$

where $U(x)$ is the solution of a homogenized equation:

$$(1.4) \quad \begin{cases} -\operatorname{div}(\mathcal{A}(x)\nabla U(x)) = f(x) & x \in D, \\ U(x) = 0 & x \in \partial D. \end{cases}$$

The homogenized coefficient $\mathcal{A}(x)$ can be obtained from the solutions of the so-called cell problem. In general, there are no explicit formulas for $\mathcal{A}(x)$, except in one dimension.

Several numerical methods have been developed to deal specifically with the case when $a(x, y)$ is periodic in y . References [3, 4, 7] propose to solve the homogenized equations as well as the equations for the correctors. Schwab et al. [29, 38] use multiscale test functions of the form $\varphi(x, x/\varepsilon)$ where $\varphi(x, y)$ is periodic in y to extract the leading order behavior of $u^\varepsilon(x)$, extending an idea that was used analytically in the work of [2, 15, 34, 44] for the homogenization problems. These methods have the feature that their cost is independent of ε , hence sublinear as $\varepsilon \rightarrow 0$, but so far they are restricted to the periodic homogenization problem. An alternative proposal for more general problems but with much higher cost is found in [20, 25].

1.2. **Heterogeneous multiscale method.** HMM [16, 17, 18] is a general methodology for designing sublinear algorithms by exploiting scale separation and other special features of the problem. It consists of two components: selection of a macroscopic solver and estimating the missing macroscale data by solving locally the fine scale problem.

For (1.1) the macroscopic solver can be chosen as a conventional \mathcal{P}_k finite element method on a triangulation \mathcal{T}_H of element size H which should resolve the macroscale features of $a^\varepsilon(x)$. The missing data is the effective stiffness matrix at

this scale. This stiffness matrix can be estimated as follows. Assuming that the effective coefficient at this scale is $\mathcal{A}_H(x)$, if we knew $\mathcal{A}_H(x)$ explicitly, we could have evaluated the quadratic form

$$\int_D \nabla V(x) \cdot \mathcal{A}_H(x) \nabla V(x) dx$$

by numerical quadrature: For any $V \in X_H$, the finite element space,

$$(1.5) \quad A_H(V, V) \simeq \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell (\nabla V \cdot \mathcal{A}_H \nabla V)(x_\ell),$$

where $\{x_\ell\}$ and $\{\omega_\ell\}$ are the quadrature points and weights in K , $|K|$ is the volume of K . In the absence of explicit knowledge of $\mathcal{A}_H(x)$, we approximate $(\nabla V \cdot \mathcal{A}_H \nabla V)(x_\ell)$ by solving the problem:

$$(1.6) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x) \nabla v_\ell^\varepsilon(x)) = 0 & x \in I_\delta(x_\ell), \\ v_\ell^\varepsilon(x) = V_\ell(x) & x \in \partial I_\delta(x_\ell), \end{cases}$$

where $I_\delta(x_\ell)$ is a cube of size δ centered at x_ℓ , and V_ℓ is the linear approximation of V at x_ℓ . We then let

$$(1.7) \quad (\nabla V \cdot \mathcal{A}_H \nabla V)(x_\ell) \simeq \frac{1}{\delta^d} \int_{I_\delta(x_\ell)} \nabla v_\ell^\varepsilon(x) \cdot a^\varepsilon(x) \nabla v_\ell^\varepsilon(x) dx.$$

(1.5) and (1.7) together give the needed approximate stiffness matrix at the scale H . For convenience, we will define the corresponding bilinear form: For any $V, W \in X_H$

$$A_H(V, W) := \sum_{K \in \mathcal{T}_H} \frac{|K|}{\delta^d} \sum_{x_\ell \in K} \omega_\ell \int_{I_\delta(x_\ell)} \nabla v_\ell^\varepsilon(x) \cdot a^\varepsilon(x) \nabla w_\ell^\varepsilon(x) dx,$$

where w_ℓ^ε is defined for $W \in X_H$ in the same way that v_ℓ^ε in (1.6) was defined for V .

In order to reduce the effect of the imposed boundary condition on $\partial I_\delta(x_\ell)$, we may replace (1.7) by

$$(1.8a) \quad (\nabla V \cdot \mathcal{A}_H \nabla V)(x_\ell) \simeq \frac{1}{(\delta')^d} \int_{I_{\delta'}(x_\ell)} \nabla v_\ell^\varepsilon(x) \cdot a^\varepsilon(x) \nabla v_\ell^\varepsilon(x) dx,$$

where $\delta' < \delta$. For example, we may choose $\delta' = \delta/2$. In (1.6), we used the Dirichlet boundary condition. Other boundary conditions are possible, such as Neumann and periodic boundary conditions. In the case when $a^\varepsilon(x) = a(x, x/\varepsilon)$ and $a(x, y)$ is periodic in y , one can take $I_\delta(x_\ell)$ to be $x_\ell + \varepsilon I$, i.e., $\delta = \varepsilon$ and use the boundary condition that $v_\ell^\varepsilon(x) - V_\ell(x)$ is periodic on I_δ .

So far the algorithm is completely general. The savings compared with solving the full fine scale problem comes from the fact that we can choose $I_\delta(x_\ell)$ to be smaller than K . The size of $I_\delta(x_\ell)$ is determined by many factors, including the accuracy and cost requirement, the degree of scale separation, and the microstructure in $a^\varepsilon(x)$. One purpose for the error estimates that we present below is to give guidelines on how to select $I_\delta(x_\ell)$. As mentioned already, if $a^\varepsilon(x) = a(x, x/\varepsilon)$ and $a(x, y)$ is periodic in y , we can simply choose $I_\delta(x_\ell)$ to be $x_\ell + \varepsilon I$, i.e., $\delta = \varepsilon$. If $a(x, y)$ is random, then δ should be a few times larger than the local correlation

length of a^ε . In the former case, the total cost is independent of ε . In the latter case, the total cost depends only weakly on ε (see [31]).

The final problem is to solve

$$(1.8) \quad \min_{V \in X_H} \frac{1}{2} A_H(V, V) - (f, V).$$

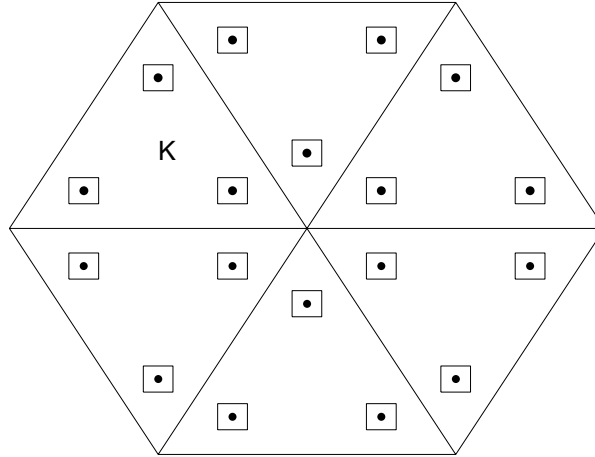


FIGURE 1. Illustration of HMM for solving (1.1). The dots are the quadrature points. The little squares are the microcell $I_\delta(x_\ell)$.

To summarize, HMM has the following features:

- (1) It gives a framework that allows us to maximally take advantage of the special features of the problem such as scale separation. For periodic homogenization problems, the cost of HMM is comparable to the special techniques discussed in [3, 7, 29, 35]. However HMM is also applicable for random problems and for problems whose coefficient $a^\varepsilon(x)$ does not have the structure of $a(x, x/\varepsilon)$. For problems without scale separation, we may consider other possible special features of the problem such as local self-similarity, which is considered in [19].
- (2) For problems without any special features, HMM becomes a fine scale solver by choosing an H that resolves the fine scales and letting $\mathcal{A}_H(x) = a^\varepsilon(x)$.

Some related ideas exist in the literature. Durlafsky [14] proposed an up-scaling method, which directly solves some local problems for obtaining the effective coefficients [33, 40, 41]. Oden and Vemaganti [35] proposed a method that aims at recovering the oscillations in ∇u^ε locally by solving a local problem with some given approximation to the macroscopic state U as the boundary condition. This idea is sometimes used in HMM to recover the microstructural information. Other numerical methods that use local microscale solvers to help extract macroscale behavior are found in [26, 27].

The numerical performance of HMM including comparison with other methods is discussed in [31].

This paper will focus on the analysis of HMM. We will estimate the error between the numerical solutions of HMM and the solutions of (1.4). We will also discuss

how to construct better approximations of u^ε from the HMM solutions. Our basic strategy is as follows. First we will prove a general statement that the error between the HMM solution and the solution of (1.4) is controlled by the standard error in the macroscale solver plus a new term, called $e(\text{HMM})$, due to the error in estimating the stiffness matrix. We then estimate $e(\text{HMM})$. This second part is only done for either periodic or random homogenization problems, since concrete results are only possible if the behavior of u^ε is well understood. We believe that this overall strategy will be useful for analyzing other multiscale methods.

We will always assume that $a^\varepsilon(x)$ is smooth, symmetric and uniformly elliptic:

$$(1.9) \quad \lambda I \leq a^\varepsilon \leq \Lambda I$$

for some $\lambda, \Lambda > 0$. We will use the summation convention and standard notation for Sobolev spaces (see [1]). We will use $|\cdot|$ to denote the absolute value of a scalar quantity, the Euclidean norm of a vector and the volume of a set K . For the quadrature formula (1.5), we will assume the following accuracy conditions for k th-order numerical quadrature scheme [11]:

$$(1.10) \quad \int_K p(x) dx = \frac{1}{|K|} \int_K p(x) dx = \sum_{\ell=1}^L \omega_\ell p(x_\ell) \quad \text{for all } p(x) \in \mathcal{P}_{2k-2}.$$

Here $\omega_\ell > 0$, for $\ell = 1, \dots, L$. For $k = 1$, we assume the above formula to be exact for $p \in \mathcal{P}_1$.

1.3. Main results. Our main results for the linear problem are as follows.

Theorem 1.1. *Denote by U_0 and U_{HMM} the solution of (1.4) and the HMM solution, respectively. Let*

$$e(\text{HMM}) = \max_{\substack{x_\ell \in K \\ K \in \mathcal{T}_H}} \|\mathcal{A}(x_\ell) - \mathcal{A}_H(x_\ell)\|,$$

where $\|\cdot\|$ is the Euclidean norm. If U_0 is sufficiently smooth and (1.10) holds, then there exists a constant C independent of ε, δ and H , such that

$$(1.11) \quad \|U_0 - U_{\text{HMM}}\|_1 \leq C(H^k + e(\text{HMM})),$$

$$(1.12) \quad \|U_0 - U_{\text{HMM}}\|_0 \leq C(H^{k+1} + e(\text{HMM})).$$

If there exists a constant C_0 such that $e(\text{HMM})|\ln H| < C_0$, then there exists a constant H_0 such that for all $H \leq H_0$,

$$(1.13) \quad \|U_0 - U_{\text{HMM}}\|_{1,\infty} \leq C(H^k + e(\text{HMM}))|\ln H|.$$

At this stage, no assumption on the form of $a^\varepsilon(x)$ is necessary. U_0 can be the solution of an arbitrary macroscopic equation with the same right-hand side as in (1.1). Of course for U_{HMM} to converge to U_0 , i.e., $e(\text{HMM}) \rightarrow 0$, U_0 must be chosen as the solution of the homogenized equation, which we now assume exists. To obtain quantitative estimates on $e(\text{HMM})$, we must restrict ourselves to more specific cases.

Theorem 1.2. *For the periodic homogenization problem, we have*

$$e(\text{HMM}) \leq \begin{cases} C\varepsilon & \text{if } I_\delta(x_\ell) = x_\ell + \varepsilon I, \\ C\left(\frac{\varepsilon}{\delta} + \delta\right) & \text{otherwise.} \end{cases}$$

In the first case, we replace the boundary condition in (1.6) by a periodic boundary condition: $v_\ell^\varepsilon - V_\ell$ is periodic with period εI . For the second result we do not need to assume that the period of $a(x, \cdot)$ is a cube: In fact it can be of arbitrary shape as long as its translation tiles up the whole space.

Another important case for which a specific estimate on $e(\text{HMM})$ can be obtained is the random homogenization. In this case, using results in [43], we have

Theorem 1.3. *For the random homogenization problem, assuming that (A) in the Appendix holds (see [43]), we have*

$$\mathbb{E} e(\text{HMM}) \leq \begin{cases} C(\kappa) \left(\frac{\varepsilon}{\delta}\right)^\kappa & d = 3, \\ \text{remains open} & d = 2, \\ C(\kappa) \left(\frac{\varepsilon}{\delta}\right)^{1/2} & d = 1, \end{cases}$$

where

$$\kappa = \frac{6 - 12\gamma}{25 - 8\gamma}$$

for any $0 < \gamma < 1/2$. By choosing γ small, κ can be arbitrarily close to $6/25$.

The probabilistic set-up will be given in the Appendix. To prove this result, we assume that (1.8a) is used with $\delta' = \delta/2$.

1.4. Recovering the microstructural information. In many applications, the microstructure information in $u^\varepsilon(x)$ is very important. U_{HMM} by itself does not give this information. However, this information can be recovered using a simple post-processing technique. For the general case, having U_{HMM} , one can obtain locally the microstructural information using an idea in [35]. Assume that we are interested in recovering u^ε and ∇u^ε only in the subdomain $\Omega \subset D$. Consider the following auxiliary problem:

$$(1.14) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x)\nabla \tilde{u}^\varepsilon(x)) = f(x) & x \in \Omega_\eta, \\ \tilde{u}^\varepsilon(x) = U_{\text{HMM}}(x) & x \in \partial\Omega_\eta, \end{cases}$$

where Ω_η satisfies $\Omega \subset \Omega_\eta \subset D$ and $\operatorname{dist}(\partial\Omega, \partial\Omega_\eta) = \eta$. We then have

Theorem 1.4. *There exists a constant C such that*

$$(1.15) \quad \left(\int_\Omega |\nabla(u^\varepsilon - \tilde{u}^\varepsilon)|^2 dx \right)^{1/2} \leq \frac{C}{\eta} \left(\|U_0 - U_{\text{HMM}}\|_{L^\infty(\Omega_\eta)} + \|u^\varepsilon - U_0\|_{L^\infty(\Omega_\eta)} \right).$$

For the random problem, the last term was estimated in [43].

A much simpler procedure exists for the periodic homogenization problem. Consider the case when $k = 1$ and choose $I_\delta = x_K + \varepsilon I$, where x_K is the barycenter of K . Here we have assumed that the quadrature point is at x_K .

Let \tilde{u}^ε be defined piecewise as follows:

- (1) $\tilde{u}^\varepsilon|_{I_\delta} = v_K^\varepsilon$, where v_K^ε is the solution of (1.6) with the boundary condition that $v_K^\varepsilon - U_{\text{HMM}}$ is periodic with period εI and $\int_{I_\delta} (\tilde{u}^\varepsilon - U_{\text{HMM}})(x) dx = 0$.
- (2) $(\tilde{u}^\varepsilon - U_{\text{HMM}})|_K$ is periodic with period εI .

For this case, we can prove

Theorem 1.5. *Let \tilde{u}^ε be defined as above. Then*

$$(1.16) \quad \left(\sum_{K \in \mathcal{T}_H} \|\nabla(u^\varepsilon - \tilde{u}^\varepsilon)\|_{0,K}^2 \right)^{1/2} \leq C(\sqrt{\varepsilon} + H).$$

Similar results with some modification hold for nonlinear problems. The details are given in §5.

1.5. Some technical background. In this subsection, we will list some general results that will be frequently referred to later on.

Given a triangulation \mathcal{T}_H , it is called *regular* if there is a constant σ such that

$$\frac{H_K}{\rho_K} \leq \sigma \quad \text{for all } K \in \mathcal{T}_H$$

and if the quantity

$$H = \max_{K \in \mathcal{T}_H} H_K$$

approaches zero, where H_K is the diameter of K and ρ_K is the diameter of the largest ball inscribed in K . \mathcal{T}_H satisfies an *inverse assumption* if there exists a constant ν such that

$$\frac{H}{H_K} \leq \nu \quad \text{for all } K \in \mathcal{T}_H.$$

A regular family of triangulation of \mathcal{T}_H satisfying the inverse assumption is called *quasi-uniform*.

The following interpolation result for the Lagrange finite element is adapted from [10]. Here and in what follows, for any $k \geq 2$, $\nabla^k v$ is understood in a piecewise manner.

Theorem 1.6 ([10]). *Let Π be k th-order Lagrange interpolate operator, and assume that the following inclusions hold:*

$$(1.17) \quad W^{k+1,p}(\hat{K}) \hookrightarrow \mathcal{C}^0(\hat{K}) \quad \text{and} \quad W^{k+1,p}(\hat{K}) \hookrightarrow W^{m,q}(\hat{K}).$$

Then

$$(1.18) \quad |v - \Pi v|_{m,q,K} \leq C|K|^{1/q-1/p} \frac{H_K^{k+1}}{\rho_K^m} |v|_{k+1,p,K}.$$

If \mathcal{T}_H is regular, we have the global estimate

$$(1.19) \quad |v - \Pi v|_{m,q,D} \leq CH^{k+1-m+\min\{0,d(1/q-1/p)\}} |v|_{k+1,p,D}.$$

Inequality (1.18) is proven in [10, Theorem 3.1.6], and (1.19) is a direct consequence of (1.18) and the *inverse inequality* below.

Using (1.19) with $p = q = 2$ and $m = 2, k = 1$, we have $\|v - \Pi v\|_{2,D} \leq C\|v\|_{2,D}$. Hence

$$(1.20) \quad \|\Pi v\|_{2,D} \leq \|v - \Pi v\|_{2,D} + \|v\|_{2,D} \leq C\|v\|_{2,D}.$$

We will also need the following form of the inverse inequality.

Theorem 1.7 ([10, Theorem 3.2.6]). *Assume that \mathcal{T}_H is regular, and assume also that the two pairs (l, r) and (m, q) with $l, m \geq 0$ and $r, q \in [0, \infty]$ satisfy*

$$l \leq m \quad \text{and} \quad \mathcal{P}_k(\hat{K}) \subset W^{l,r}(\hat{K}) \cap W^{m,q}(\hat{K}).$$

Then there exists a constant $C = C(\sigma, \nu, l, r, m, q)$ such that

$$(1.21) \quad |v|_{m,q,K} \leq CH_K^{l-m+d(1/q-1/r)} |v|_{l,r,K}$$

for any $v \in \mathcal{P}_k(\hat{K})$.

If in addition \mathcal{T}_H satisfies the inverse assumption, then there exists a constant $C = C(\sigma, \nu, l, r, m, q)$ such that

$$(1.22) \quad \left(\sum_{K \in \mathcal{T}_H} |v|_{m,q,K}^q \right)^{1/q} \leq CH^{l-m+\min\{0, d(1/q-1/r)\}} \left(\sum_{K \in \mathcal{T}_H} |v|_{l,r,K}^r \right)^{1/r}$$

for any $v \in X_H$ and $r, q < \infty$, with

$$\begin{aligned} \max_{K \in \mathcal{T}_H} |v|_{m,\infty,K} \quad \text{replacing} \quad & \left(\sum_{K \in \mathcal{T}_H} |v|_{m,q,K}^q \right)^{1/q}, \quad \text{if } q = \infty, \\ \max_{K \in \mathcal{T}_H} |v|_{l,\infty,K} \quad \text{replacing} \quad & \left(\sum_{K \in \mathcal{T}_H} |v|_{l,r,K}^r \right)^{1/r}, \quad \text{if } r = \infty. \end{aligned}$$

The following simple result will be used repeatedly.

Lemma 1.8. Let $A_1(x)$ and $A_2(x)$ be symmetric matrices satisfying (1.9). Let φ_1 be the solution of

$$(1.23) \quad -\operatorname{div}(A_1(x)\nabla\varphi_1(x)) = \operatorname{div}(\tilde{A}_1(x)\nabla F_1(x)) \quad x \in \Omega,$$

with either the Dirichlet or periodic boundary condition on $\partial\Omega$. Let φ_2 be a solution of (1.23) with A_1, \tilde{A}_1 and F_1 replaced by A_2, \tilde{A}_2 and F_2 , respectively, and let φ_2 satisfy the same boundary condition as φ_1 . Then

$$(1.24) \quad \begin{aligned} \lambda \|\nabla(\varphi_1 - \varphi_2)\|_{0,\Omega} &\leq \max_{x \in \Omega} |(\tilde{A}_1 - \tilde{A}_2)(x)| \|\nabla F_1\|_{0,\Omega} + \max_{x \in \Omega} |(A_1 - A_2)(x)| \|\nabla \varphi_2\|_{0,\Omega} \\ &+ \max_{x \in \Omega} |\tilde{A}_2(x)| \|\nabla(F_1 - F_2)\|_{0,\Omega}. \end{aligned}$$

Proof. Inequality (1.24) is a direct consequence of

$$\lambda \|\nabla(\varphi_1 - \varphi_2)\|_{0,\Omega}^2 \leq \int_{\Omega} \nabla(\varphi_1 - \varphi_2) \cdot ((\tilde{A}_2 - \tilde{A}_1)\nabla F_1 + (A_2 - A_1)\nabla \varphi_2 + \tilde{A}_2\nabla(F_2 - F_1)).$$

□

The following simple result underlies the stability of HMM for problem (1.1).

Lemma 1.9. Let φ be the solution of

$$(1.25) \quad \begin{cases} -\operatorname{div}(a\nabla\varphi) = 0 & \text{in } \Omega \subset \mathbb{R}^d, \\ \varphi = V_\ell & \text{on } \partial\Omega, \end{cases}$$

where V_ℓ is a linear function and $a = (a_{ij})$ satisfies

$$\lambda I \leq a \leq \Lambda I.$$

Then we have

$$(1.26) \quad \|\nabla V_\ell\|_{0,\Omega} \leq \|\nabla\varphi\|_{0,\Omega} \quad \text{and} \quad \left(\int_{\Omega} \nabla\varphi \cdot a\nabla\varphi \right)^{1/2} \leq \left(\int_{\Omega} \nabla V_\ell \cdot a\nabla V_\ell \right)^{1/2}.$$

Proof. Notice that $\varphi = V_\ell$ on the edges of Ω , using the fact that ∇V_ℓ is a constant in Ω , and integration by parts leads to

$$\int_{\Omega} \nabla(\varphi - V_\ell)(x) \nabla V_\ell(x) \, dx = 0,$$

which implies

$$\int_{\Omega} |\nabla \varphi(x)|^2 \, dx = \int_{\Omega} |\nabla V_\ell(x)|^2 \, dx + \int_{\Omega} |\nabla(\varphi - V_\ell)(x)|^2 \, dx.$$

This gives the first result in (1.26). Multiplying (1.25) by $\varphi(x) - V_\ell(x)$ and integrating by parts, we obtain

$$\begin{aligned} \int_{\Omega} \nabla \varphi(x) \cdot a \nabla \varphi(x) \, dx + \int_{\Omega} \nabla(\varphi - V_\ell)(x) \cdot a \nabla(\varphi - V_\ell)(x) \, dx \\ = \int_{\Omega} \nabla V_\ell(x) \cdot a \nabla V_\ell(x) \, dx. \end{aligned}$$

This gives the second part of (1.26). □

Remark 1.10. For this result, the coefficient $a = (a_{ij})$ may depend on the solution, i.e., (1.25) may be nonlinear.

Remark 1.11. The same result holds if we use instead a periodic boundary condition: $\varphi - V_\ell$ is periodic with period Ω .

2. GENERALITIES

Here we prove Theorem 1.1. We will let $U_H = U_{\text{HMM}}$ for convenience.

Since U_H is the numerical solution associated with the quadratic form A_H , U_0 is the exact solution associated with the quadratic form A , defined for any $V \in H_0^1(D)$ as

$$A(V, V) = \int_D \nabla V(x) \cdot \mathcal{A}(x) \nabla V(x) \, dx.$$

To estimate $U_0 - U_H$, we view A_H as an approximation to A , and we use *Strang's first lemma* [10].

Using (1.26) with $\Omega = I_\delta(x_\ell)$ and (1.9), for any $V \in X_H$, we have

$$\begin{aligned} A_H(V, V) &\geq \lambda \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell \int_{I_\delta(x_\ell)} |\nabla V_\ell(x)|^2 \, dx \\ &= \lambda \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell |\nabla V(x_\ell)|^2 \\ (2.1) \qquad &= \lambda \|\nabla V\|_0^2. \end{aligned}$$

Similarly, for any $V, W \in X_H$, we obtain

$$\begin{aligned}
|A_H(V, W)| &\leq \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell \left(\int_{I_\delta(x_\ell)} \nabla V_\ell \cdot a^\varepsilon \nabla V_\ell \right)^{\frac{1}{2}} \left(\int_{I_\delta(x_\ell)} \nabla W_\ell \cdot a^\varepsilon \nabla W_\ell \right)^{\frac{1}{2}} \\
&\leq \Lambda \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell |\nabla V(x_\ell)| |\nabla W(x_\ell)| \\
&= \Lambda \sum_{K \in \mathcal{T}_H} \int_K |\nabla V(x)| |\nabla W(x)| dx \\
(2.2) \quad &\leq \Lambda \|\nabla V\|_0 \|\nabla W\|_0.
\end{aligned}$$

The existence and the uniqueness of the solutions to (1.8) follow from (2.1) and (2.2) via the Lax-Milgram lemma and the *Poincaré inequality*.

To streamline the proof of Theorem 1.1, we introduce the following auxiliary bilinear form \hat{A}_H .

$$\hat{A}_H(V, W) = \sum_{K \in \mathcal{T}_H} \hat{A}_K(V, W) \quad \text{with} \quad \hat{A}_K(V, W) = |K| \sum_{x_\ell \in K} \omega_\ell (\nabla W \cdot \mathcal{A} \nabla V)(x_\ell).$$

Classical results on numerical integration [11, Theorem 6] give for any $V, W \in X_H$,

$$(2.3) \quad \left| \hat{A}_K(V, W) - \int_K \nabla W \cdot \mathcal{A} \nabla V dx \right| \leq CH^m \|V\|_{m,K} \|\nabla W\|_{0,K} \quad 1 \leq m \leq k.$$

Moreover, for any $V, W \in X_H$, if $\|V\|_{k+1}$ and $\|W\|_2$ are bounded, we have [11, Theorem 8],

$$(2.4) \quad |\hat{A}_H(V, W) - A(V, W)| \leq CH^{k+1} \|V\|_{k+1} \|W\|_2.$$

Proof of Theorem 1.1. Using the *first Strang lemma* [10, Theorem 4.1.1], we have

$$\|U_0 - U_H\|_1 \leq C \inf_{V \in X_H} \left(\|U_0 - V\|_1 + \sup_{W \in X_H} \frac{|A_H(V, W) - A(V, W)|}{\|W\|_1} \right).$$

Let $V = \Pi U_0$ and using (1.19) with $m = 1, p = q = 2$, we have

$$(2.5) \quad \inf_{V \in X_H} \|U_0 - V\|_1 \leq \|U_0 - \Pi U_0\|_1 \leq CH^k.$$

It remains to estimate $|A_H(V, W) - A(V, W)|$ for $V = \Pi U_0$ and $W \in X_H$. Using (2.3), we get

$$\begin{aligned}
|A_H(V, W) - A(V, W)| &\leq |A_H(V, W) - \hat{A}_H(V, W)| + |\hat{A}_H(V, W) - A(V, W)| \\
(2.6) \quad &\leq (e(\text{HMM})) \|\nabla V\|_0 + CH^k \|V\|_k \|\nabla W\|_0.
\end{aligned}$$

This gives (1.11)

To get the L^2 estimate, we use the Aubin-Nitsche dual argument [10]. To this end, consider the following auxiliary problem: Find $w \in H_0^1(D)$ such that

$$(2.7) \quad A(v, w) = (U_0 - U_H, v) \quad \text{for all } v \in H_0^1(D).$$

The standard regularity result reads [24]

$$(2.8) \quad \|w\|_2 \leq C \|U_0 - U_H\|_0.$$

Putting $v = U_0 - U_H$ into the right-hand side of (2.7), we obtain

$$\begin{aligned}
 \|U_0 - U_H\|_0^2 &= A(U_0 - U_H, w - \Pi w) + (A_H(U_H, \Pi w) - A(U_H, \Pi w)) \\
 &= A(U_0 - U_H, w - \Pi w) \\
 &\quad + (A_H(U_H - \Pi U_0, \Pi w) - A(U_H - \Pi U_0, \Pi w)) \\
 (2.9) \quad &\quad + (A_H(\Pi U_0, \Pi w) - A(\Pi U_0, \Pi w)).
 \end{aligned}$$

Using (2.6) with $k = 1$, we bound the first two terms in the right-hand side of the above identity as

$$|A(U_0 - U_H, w - \Pi w)| \leq C \|U_0 - U_H\|_1 \|w - \Pi w\|_1 \leq CH \|U_0 - U_H\|_1 \|w\|_2$$

and

$$|A_H(U_H - \Pi U_0, \Pi w) - A(U_H - \Pi U_0, \Pi w)| \leq (e(\text{HMM}) + CH) \|U_0 - U_H\|_1 \|\Pi w\|_1.$$

The last term in the right-hand side of (2.9) may be decomposed into

$$\begin{aligned}
 A_H(\Pi U_0, \Pi w) - A(\Pi U_0, \Pi w) &= (A_H(\Pi U_0, \Pi w) - \hat{A}_H(\Pi U_0, \Pi w)) \\
 &\quad + (\hat{A}_H(\Pi U_0, \Pi w) - A(\Pi U_0, \Pi w)).
 \end{aligned}$$

It follows from (2.4) that

$$|\hat{A}_H(\Pi U_0, \Pi w) - A(\Pi U_0, \Pi w)| \leq CH^{k+1} \|U_0\|_{k+1} \|w\|_2.$$

By definition of $e(\text{HMM})$ and using (1.20), we get

$$|A_H(\Pi U_0, \Pi w) - \hat{A}_H(\Pi U_0, \Pi w)| \leq Ce(\text{HMM}) \|\nabla \Pi U_0\|_0 \|w\|_2.$$

Combining the above estimates and using (2.8) lead to (1.12).

It remains to prove (1.13). As in [37], for any point $z \in D$, we define the regularized Green's function $G^z \in H_0^1(D)$ and the discrete Green's function $G_H^z \in X_H$ as

$$\begin{aligned}
 A(G^z, V) &= (\delta_z, \partial V) \quad \text{for all } V \in H_0^1(D), \\
 (2.10) \quad A(G_H^z, V) &= (\delta_z, \partial V) \quad \text{for all } V \in X_H,
 \end{aligned}$$

where δ_z is the regularized Dirac- δ function defined in [37]. It is well known that

$$(2.11) \quad \|G^z - G_H^z\|_{1,1} \leq C \quad \text{and} \quad \|G_H^z\|_{1,1} \leq C |\ln H|.$$

A proof for (2.11) can be obtained by using the weighted-norm technique [37]. We refer to [9, Chapter 7] for details. Using the definition of G^z and G_H^z , a simple manipulation gives

$$\begin{aligned}
 \partial(U_0 - U_H)(z) &= A(G^z, U_0 - \Pi U_0) + A(G^z, \Pi U_0 - U_H) \\
 &= A(G^z - G_H^z, U_0 - \Pi U_0) + A(G_H^z, U_0 - U_H) \\
 &= A(G^z - G_H^z, U_0 - \Pi U_0) + A_H(U_H, G_H^z) - A(U_H, G_H^z) \\
 &= A(G^z - G_H^z, U_0 - \Pi U_0) + (A_H(\Pi U_0, G_H^z) - A(\Pi U_0, G_H^z)) \\
 &\quad + (A_H(U_H - \Pi U_0, G_H^z) - A(U_H - \Pi U_0, G_H^z)).
 \end{aligned}$$

Using (2.11), we obtain

$$\begin{aligned}
 \|U_0 - U_H\|_{1,\infty} &\leq C \|U_0 - \Pi U_0\|_{1,\infty} + |A(\Pi U_0, G_H^z) - A_H(\Pi U_0, G_H^z)| \\
 &\quad + |A(U_H - \Pi U_0, G_H^z) - A_H(U_H - \Pi U_0, G_H^z)|.
 \end{aligned}$$

Using (2.6), we get

$$\begin{aligned} |A(\Pi U_0, G_H^z) - A_H(\Pi U_0, G_H^z)| &\leq (e(\text{HMM}) + CH^k) \sum_{K \in \mathcal{T}_H} \|\Pi U_0\|_{k,K} \|\nabla G_H^z\|_{0,K} \\ &\leq C(e(\text{HMM}) + H^k) \sum_{K \in \mathcal{T}_H} \|\Pi U_0\|_{k,\infty,K} \|\nabla G_H^z\|_{L^1(K)} \\ &\leq C(e(\text{HMM}) + H^k) |\ln H| \|U_0\|_{k+1,\infty}, \end{aligned}$$

where we have used the inverse inequality (1.21).

Similarly, we have

$$\begin{aligned} &|A(U_H - \Pi U_0, G_H^z) - A_H(U_H - \Pi U_0, G_H^z)| \\ &\leq (e(\text{HMM}) + CH) \sum_{K \in \mathcal{T}_H} \|U_H - \Pi U_0\|_{1,K} \|\nabla G_H^z\|_{0,K} \\ &\leq C(e(\text{HMM}) + H) \sum_{K \in \mathcal{T}_H} \|U_H - \Pi U_0\|_{1,\infty,K} \|\nabla G_H^z\|_{0,1,K} \\ &\leq C(e(\text{HMM}) + H) |\ln H| \|U_0 - U_H\|_{1,\infty} \\ &\quad + C(e(\text{HMM}) + H) |\ln H| H^k \|U_0\|_{k+1,\infty}. \end{aligned}$$

A combination of the above three estimates yields

$$\begin{aligned} \|U_0 - U_H\|_{1,\infty} &\leq CH^k + C(e(\text{HMM}) + H) |\ln H| \|U_0 - U_H\|_{1,\infty} \\ &\quad + C(e(\text{HMM}) + H^k) |\ln H| \|U_0\|_{k+1,\infty}. \end{aligned}$$

If $e(\text{HMM})|\ln H| < C_0 := 1/(2C)$, then there exists a constant H_0 such that for all $H \leq H_0$,

$$C(e(\text{HMM}) + H) |\ln H| \leq 1/2 + CH |\ln H| < 1.$$

We thus obtain (1.13) and this completes the proof. \square

Combining the foregoing proof for the L^2 and $W^{1,\infty}$ estimates, using the Green's function defined in [39], we obtain

Remark 2.1. Under the same condition for the $W^{1,\infty}$ estimate in Theorem 1.1, we have

$$\|U_0 - U_H\|_{L^\infty} \leq C(e(\text{HMM}) + H^{k+1}) |\ln H|^2.$$

3. ESTIMATING $e(\text{HMM})$

In this section, we estimate $e(\text{HMM})$ for problems with locally periodic coefficients. The estimate of $e(\text{HMM})$ for problems with random coefficients can be found in the Appendix.

We assume that $a^\varepsilon(x) = a(x, x/\varepsilon)$, where a^ε is smooth in x and periodic in y with period I . Define $\kappa = \lfloor \delta/\varepsilon \rfloor$, and we introduce \hat{V}_ℓ as

$$(3.1) \quad \hat{V}_\ell(x) = V_\ell(x) + \varepsilon \chi^k \left(x_\ell, \frac{x}{\varepsilon} \right) \frac{\partial V_\ell}{\partial x_k}(x),$$

where $\{\chi^j\}_{j=1}^d$ is defined as: For $j = 1, \dots, d$, $\chi^j(x, y)$ is periodic in y with period I and satisfies

$$(3.2) \quad -\frac{\partial}{\partial y_i} \left(a_{ik} \frac{\partial \chi^j}{\partial y_k} \right) (x, y) = \frac{\partial}{\partial y_i} a_{ij}(x, y) \quad \text{in } I, \quad \int_I \chi^j(x, y) dy = 0.$$

Given $\{\chi^j\}_{j=1}^d$, the homogenized coefficient $\mathcal{A} = (\mathcal{A}_{ij}(x))$ is given by

$$\mathcal{A}_{ij}(x) = \int_I \left(a_{ij} + a_{ik} \frac{\partial \chi^j}{\partial y_k} \right) (x, y) dy.$$

Note that $\{\chi^j\}_{j=1}^d$ is smooth and bounded in all norms.

First let us consider the case when $I_\delta(x_\ell) = x_\ell + \varepsilon I$, and (1.6) is solved with the periodic boundary condition. Denote by \hat{v}_ℓ^ε the solution of (1.6) with the coefficients $a^\varepsilon(x)$ replaced by $a(x_\ell, x/\varepsilon)$. \hat{v}_ℓ^ε may be viewed as a perturbation of v_ℓ^ε . Using Lemma 1.8, we get

$$(3.3) \quad \|\nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon)\|_{0, I_\varepsilon} \leq C\varepsilon \|\nabla V_\ell\|_{0, I_\varepsilon}.$$

Observe that $\hat{v}_\ell^\varepsilon = \hat{V}_\ell$. A direct calculation yields

$$\begin{aligned} (\nabla W \cdot (\mathcal{A}_H - \mathcal{A}) \nabla V)(x_\ell) &= \int_{I_\varepsilon} \nabla w_\ell^\varepsilon \cdot \left[a\left(x, \frac{x}{\varepsilon}\right) - a\left(x_\ell, \frac{x}{\varepsilon}\right) \right] \nabla v_\ell^\varepsilon dx \\ &\quad + \int_{I_\varepsilon} \nabla w_\ell^\varepsilon \cdot a\left(x_\ell, \frac{x}{\varepsilon}\right) \nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) dx. \end{aligned}$$

Using (3.3), we get

$$(3.4) \quad e(\text{HMM}) \leq C\varepsilon.$$

Next we consider the more general case when I_δ is a cube of size δ not necessarily equal to ε . The following analysis applies equally well to the case when the period of $a(x, \cdot)$ is of general and even nonpolygonal shape. This situation arises in some examples of composite materials [30]. We will show that if δ is much larger than ε , then the averaged energy density for the solution of (1.6) closely approximates the energy density of the homogenized problem. We begin with the following observation:

$$\begin{aligned} (\nabla W \cdot \mathcal{A} \nabla V)(x_\ell) &= \nabla W_\ell(x) \cdot \mathcal{A}(x_\ell) \nabla V_\ell(x) \\ (3.5) \quad &= \int_{I_{\kappa\varepsilon}(x_\ell)} \nabla W_\ell \cdot a\left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \hat{V}_\ell dx. \end{aligned}$$

We first establish some estimates on the solution of the cell problem (1.6). We will write I_δ instead of $I_\delta(x_\ell)$ if there is no risk of confusion.

Lemma 3.1. *There exists a constant C independent of ε and δ such that for each ℓ ,*

$$(3.6) \quad \|\nabla v_\ell^\varepsilon\|_{0, I_\delta \setminus I_{\kappa\varepsilon}} \leq C \left(\left(\frac{\varepsilon}{\delta}\right)^{1/2} + \delta \right) \|\nabla V_\ell\|_{0, I_\delta}.$$

Proof. We still denote by \hat{v}_ℓ^ε the solution of (1.6) with the coefficient $a^\varepsilon(x)$ replaced by $a(x_\ell, x/\varepsilon)$. Using Lemma 1.8, we get

$$(3.7) \quad \|\nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon)\|_{0, I_\delta} \leq C\delta \|\nabla V_\ell\|_{0, I_\delta}.$$

Define $\theta_\ell^\varepsilon = \hat{v}_\ell^\varepsilon - \hat{V}_\ell$, which obviously satisfies

$$(3.8) \quad \begin{cases} -\operatorname{div}\left(a\left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \theta_\ell^\varepsilon(x)\right) = 0 & x \in I_\delta(x_\ell), \\ \theta_\ell^\varepsilon(x) = -\varepsilon \chi^k\left(x_\ell, \frac{x}{\varepsilon}\right) \frac{\partial V_\ell}{\partial x_k}(x) & x \in \partial I_\delta(x_\ell). \end{cases}$$

Note that θ_ℓ^ε is simply the *boundary layer correction* for the cell problem (1.6) [5]. It is proved in [45, (1.51) in §1.4], using the rescaling $x' = x/\delta$ over I_δ and $\varepsilon' = \varepsilon/\delta$.

$$(3.9) \quad \|\nabla\theta_\ell^\varepsilon\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2} \|\nabla V_\ell\|_{0,I_\delta}.$$

This together with (3.7) gives

$$(3.10) \quad \|\nabla(v_\ell^\varepsilon - \hat{V}_\ell)\|_{0,I_\delta} \leq C\left(\left(\frac{\varepsilon}{\delta}\right)^{1/2} + \delta\right) \|\nabla V_\ell\|_{0,I_\delta}.$$

A straightforward calculation gives

$$(3.11) \quad \|\nabla\hat{V}_\ell\|_{0,I_\delta \setminus I_{\kappa\varepsilon}} \leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2} \|\nabla V_\ell\|_{0,I_\delta},$$

which together with (3.10) leads to

$$\begin{aligned} \|\nabla v_\ell^\varepsilon\|_{0,I_\delta \setminus I_{\kappa\varepsilon}} &\leq \|\nabla\hat{V}_\ell\|_{0,I_\delta \setminus I_{\kappa\varepsilon}} + \|\nabla(v_\ell^\varepsilon - \hat{V}_\ell)\|_{0,I_\delta \setminus I_{\kappa\varepsilon}} \\ &\leq \|\nabla\hat{V}_\ell\|_{0,I_\delta \setminus I_{\kappa\varepsilon}} + \|\nabla(v_\ell^\varepsilon - \hat{V}_\ell)\|_{0,I_\delta} \\ &\leq C\left(\left(\frac{\varepsilon}{\delta}\right)^{1/2} + \delta\right) \|\nabla V_\ell\|_{0,I_\delta}. \end{aligned}$$

This gives (3.6). □

As in (3.6), we also have

$$(3.12) \quad \|\nabla v_\ell^\varepsilon\|_{0,I_\delta \setminus I_{(\kappa-2)\varepsilon}} \leq C\left(\left(\frac{\varepsilon}{\delta}\right)^{1/2} + \delta\right) \|\nabla V_\ell\|_{0,I_\delta}.$$

Theorem 3.2.

$$(3.13) \quad e(\text{HMM}) \leq C\left(\frac{\varepsilon}{\delta} + \delta\right).$$

Proof. Note that $v_\ell^\varepsilon = (v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) + \theta_\ell^\varepsilon + \hat{V}_\ell$. We have

$$(\nabla W \cdot (\mathcal{A}_H - \mathcal{A})\nabla V)(x_\ell) = :I_1 + I_2 + I_3,$$

where

$$\begin{aligned} I_1 &= \int_{I_\delta} \nabla w_\ell^\varepsilon \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) dx, \quad I_2 = \int_{I_\delta} \nabla w_\ell^\varepsilon \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla\theta_\ell^\varepsilon dx, \\ I_3 &= \int_{I_\delta} \nabla w_\ell^\varepsilon \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla\hat{V}_\ell dx - \nabla W_\ell \cdot \mathcal{A}(x_\ell) \nabla V_\ell. \end{aligned}$$

Using (3.7) and (2.2), we bound I_1 as

$$\begin{aligned} |I_1| &\leq \Lambda\delta^{-d} \|\nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon)\|_{0,I_\delta} \|\nabla w_\ell^\varepsilon\|_{0,I_\delta} \\ &\leq C\delta^{1-d} \|\nabla V_\ell\|_{0,I_\delta} \|\nabla W_\ell\|_{0,I_\delta} = C\delta |\nabla V_\ell| |\nabla W_\ell|. \end{aligned}$$

Using the symmetry of a^ε , $I_2 = \int_{I_\delta} \nabla\theta_\ell^\varepsilon \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla w_\ell^\varepsilon dx$ and

$$\begin{aligned} I_2 &= \int_{I_\delta} \nabla\left(\theta_\ell^\varepsilon + \varepsilon\chi^k\left(x_\ell, \frac{x}{\varepsilon}\right) \frac{\partial V_\ell}{\partial x_k}(1 - \rho^\varepsilon)\right) \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla w_\ell^\varepsilon dx \\ &\quad - \int_{I_\delta} \nabla\left(\varepsilon\chi^k\left(x_\ell, \frac{x}{\varepsilon}\right) \frac{\partial V_\ell}{\partial x_k}(1 - \rho^\varepsilon)\right) \cdot a\left(x, \frac{x}{\varepsilon}\right) \nabla w_\ell^\varepsilon dx, \end{aligned}$$

where $\rho^\varepsilon(x) \in C_0^\infty(I_\delta)$, $|\nabla \rho^\varepsilon| \leq C/\varepsilon$, and

$$(3.14) \quad \rho^\varepsilon(x) = \begin{cases} 1 & \text{if } \text{dist}(x, \partial I_\delta) \geq 2\varepsilon, \\ 0 & \text{if } \text{dist}(x, \partial I_\delta) \leq \varepsilon. \end{cases}$$

Using (1.6) and $\theta_\ell^\varepsilon + \varepsilon \chi^k(x_\ell, \frac{x}{\varepsilon}) \frac{\partial V_\ell}{\partial x_k} (1 - \rho^\varepsilon) \in H_0^1(I_\delta)$, integrating by parts makes the first term in the right-hand side of I_2 vanish; therefore we write I_2 as

$$I_2 = - \int_{I_\delta} a_{ij} \left(x, \frac{x}{\varepsilon}\right) \frac{\partial w_\ell^\varepsilon}{\partial x_i} \frac{\partial \chi^k}{\partial y_j} \frac{\partial V_\ell}{\partial x_k} (1 - \rho^\varepsilon) dx + \varepsilon \int_{I_\delta} a_{ij} \left(x, \frac{x}{\varepsilon}\right) \frac{\partial w_\ell^\varepsilon}{\partial x_i} \chi^k \frac{\partial V_\ell}{\partial x_k} \frac{\partial \rho^\varepsilon}{\partial x_j} dx.$$

Using (3.12), we bound I_2 as

$$|I_2| \leq C \delta^{-d} \|\nabla w_\ell^\varepsilon\|_{0, I_\delta \setminus I_{(\kappa-2)\varepsilon}} \|\nabla V_\ell\|_{0, I_\delta \setminus I_{(\kappa-2)\varepsilon}} \leq C \left(\frac{\varepsilon}{\delta} + \delta^2\right) |\nabla W_\ell| |\nabla V_\ell|.$$

Using (3.2) and integrating by parts, we obtain

$$\int_{I_\delta} \nabla w_\ell^\varepsilon \cdot a \left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \hat{V}_\ell dx = \int_{I_\delta} \nabla W_\ell \cdot a \left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \hat{V}_\ell dx,$$

which together with (3.5) gives

$$I_3 = \int_{I_\delta} \nabla w_\ell^\varepsilon \left[a \left(x, \frac{x}{\varepsilon}\right) - a \left(x_\ell, \frac{x}{\varepsilon}\right) \right] \nabla \hat{V}_\ell dx + \frac{1}{\delta^d} \int_{I_\delta \setminus I_{\kappa\varepsilon}} \nabla W_\ell \cdot a \left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \hat{V}_\ell dx + (|\kappa\varepsilon/\delta|^d - 1) \nabla W_\ell \cdot \mathcal{A}(x_\ell) \nabla V_\ell.$$

The last term of I_3 is bounded by

$$| |\kappa\varepsilon/\delta|^d - 1 | |\nabla W_\ell \cdot \mathcal{A}(x_\ell) \nabla V_\ell| \leq C \frac{\varepsilon}{\delta} |\nabla V_\ell| |\nabla W_\ell|,$$

where we have used $| |\kappa\varepsilon/\delta|^d - 1 | \leq C\varepsilon/\delta$. Using (3.11), we get

$$\begin{aligned} \delta^{-d} \left| \int_{I_\delta \setminus I_{\kappa\varepsilon}} \nabla W_\ell \cdot a \left(x_\ell, \frac{x}{\varepsilon}\right) \nabla \hat{V}_\ell dx \right| &\leq C \frac{1}{\delta^d} \|\nabla \hat{V}_\ell\|_{0, I_\delta \setminus I_{\kappa\varepsilon}} \|\nabla W_\ell\|_{0, I_\delta \setminus I_{\kappa\varepsilon}} \\ &\leq C \left(\frac{\varepsilon}{\delta} + \delta^2\right) |\nabla V_\ell| |\nabla W_\ell|. \end{aligned}$$

Consequently, we obtain

$$\begin{aligned} |I_3| &\leq C \delta^{1-d} \|\nabla w_\ell^\varepsilon\|_{0, I_\delta} \|\nabla \hat{V}_\ell\|_{0, I_\delta} + C \left(\frac{\varepsilon}{\delta} + \delta^2\right) |\nabla V_\ell| |\nabla W_\ell| \\ &\leq C \left(\frac{\varepsilon}{\delta} + \delta\right) |\nabla V_\ell| |\nabla W_\ell|. \end{aligned}$$

Combining the estimates for I_1, I_2 and I_3 gives the desired result (3.13). □

Remark 3.3. An explicit expression for v_ℓ^ε is available in one dimension, from which we may show that the bound for $e(\text{HMM})$ is sharp.

4. RECONSTRUCTION AND COMPRESSION

4.1. **Reconstruction procedure.** Next we consider how to construct better approximations to u^ε from U_H . We will restrict ourselves to the case when $k = 1$.

Proof of Theorem 1.4. Subtracting (1.1) from (1.14), we obtain

$$\begin{cases} -\operatorname{div}(a^\varepsilon(x)\nabla(\tilde{u}^\varepsilon - u^\varepsilon)(x)) = 0 & x \in \Omega_\eta, \\ (\tilde{u}^\varepsilon - u^\varepsilon)(x) = U_H(x) - u^\varepsilon(x) & x \in \partial\Omega_\eta. \end{cases}$$

Using classical interior estimates for elliptic equation [24], we have

$$\|\nabla(\tilde{u}^\varepsilon - u^\varepsilon)\|_{0,\Omega} \leq \frac{C}{\eta} \|\tilde{u}^\varepsilon - u^\varepsilon\|_{0,\Omega_\eta}.$$

Using the Hopf maximum principle, we get

$$\begin{aligned} \frac{1}{\eta^2} \int_{\Omega_\eta} |(\tilde{u}^\varepsilon - u^\varepsilon)(x)|^2 dx &\leq \frac{C}{\eta^2} \|\tilde{u}^\varepsilon - u^\varepsilon\|_{L^\infty(\Omega_\eta)}^2 \leq \frac{C}{\eta^2} \|u^\varepsilon - U_H\|_{L^\infty(\partial\Omega_\eta)}^2 \\ &\leq \frac{C}{\eta^2} \left(\|U_0 - U_H\|_{L^\infty(\Omega_\eta)}^2 + \|u^\varepsilon - U_0\|_{L^\infty(\Omega_\eta)}^2 \right). \end{aligned}$$

A combination of the above two results implies Theorem 1.4. □

Proof of Theorem 1.5. Denote $I_\varepsilon(x_K) = x_K + \varepsilon I$ and define \hat{u}^ε as the solution of

$$(4.1) \quad -\operatorname{div}\left(a\left(x_K, \frac{x}{\varepsilon}\right)\nabla\hat{u}^\varepsilon(x)\right) = 0 \quad \text{in } I_\varepsilon(x_K),$$

with the boundary condition that $\hat{u}^\varepsilon - U_H$ is periodic on $\partial I_\varepsilon(x_K)$ and

$$\int_{I_\varepsilon(x_K)} (\hat{u}^\varepsilon - U_H) dx = 0,$$

where x_K is the barycenter of K .

It is easy to verify that \hat{u}^ε takes the explicit form

$$(4.2) \quad \hat{u}^\varepsilon(x) = U_H(x) + \varepsilon\chi^k\left(x_K, \frac{x}{\varepsilon}\right) \frac{\partial U_H}{\partial x_k}(x).$$

Note that the periodic extension of $\hat{u}^\varepsilon - U_H$ is still $\varepsilon\chi^k\left(x_K, \frac{x}{\varepsilon}\right) \frac{\partial U_H}{\partial x_k}(x)$. This means that \hat{u}^ε is also well defined for the whole of K and takes the same explicit form as (4.2).

Using $\int_I \chi^k(x_K, y) dy = 0$ for $k = 1, \dots, d$ and that ∇U_H is a piecewise constant on K , we obtain

$$(4.3) \quad \int_{I_\varepsilon(x_K)} (\hat{u}^\varepsilon - U_H)(x) dx = \int_{I_\varepsilon(x_K)} \varepsilon\chi^k\left(x_K, \frac{x}{\varepsilon}\right) \frac{\partial U_H}{\partial x_k}(x) dx = 0.$$

As in (3.7), we have

$$\|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0,I_\varepsilon(x_K)} \leq C\varepsilon\|\nabla U_H\|_{0,I_\varepsilon(x_K)}.$$

From the construction of \tilde{u}^ε , we have for any $x_1 \in K$,

$$\|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0,I_\varepsilon(x_1)} = \|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0,I_\varepsilon(x_K)}.$$

Since ∇U_H is constant over K , we get

$$(4.4) \quad \|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0,K} \leq C\varepsilon\|\nabla U_H\|_{0,K}.$$

Adding up for all $K \in \mathcal{T}_H$ and using the a priori estimate $\|\nabla U_H\|_0 \leq C\|f\|_0$, we obtain

$$(4.5) \quad \left(\sum_{K \in \mathcal{T}_H} \|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0,K}^2 \right)^{1/2} \leq C\varepsilon\|\nabla U_H\|_0 \leq C\varepsilon.$$

Using (4.2), a straightforward calculation gives

$$\frac{\partial \hat{u}^\varepsilon}{\partial x_i} = \frac{\partial U_H}{\partial x_i} + \frac{\partial \chi^k}{\partial y_i} \left(x_K, \frac{x}{\varepsilon} \right) \frac{\partial U_H}{\partial x_k}.$$

Define the first order approximation of u^ε as

$$u_1^\varepsilon(x) = U_0 + \varepsilon \chi^k \left(x, \frac{x}{\varepsilon} \right) \frac{\partial U_0}{\partial x_k},$$

where $\{\chi^k\}_{k=1}^d$ is the solutions of (3.2). Obviously,

$$\frac{\partial u_1^\varepsilon}{\partial x_i} = \frac{\partial U_0}{\partial x_i} + \left(\varepsilon \frac{\partial \chi^k}{\partial x_i} + \frac{\partial \chi^k}{\partial y_i} \right) \left(x, \frac{x}{\varepsilon} \right) \frac{\partial U_0}{\partial x_i} + \varepsilon \chi^k \left(x, \frac{x}{\varepsilon} \right) \frac{\partial^2 U_0}{\partial x_i \partial x_k}.$$

A combination of the above estimates leads to

$$\begin{aligned} \|\nabla(\hat{u}^\varepsilon - u_1^\varepsilon)\|_{0,K} &\leq C\|\nabla(U_H - U_0)\|_{0,K} + C\varepsilon|U_0|_{1,K} \\ &\quad + \left\| \left(\frac{\partial \chi^k}{\partial y_i} \left(x, \frac{x}{\varepsilon} \right) - \frac{\partial \chi^k}{\partial y_i} \left(x_K, \frac{x}{\varepsilon} \right) \right) \frac{\partial U_0}{\partial x_k} \right\|_{0,K} + C\varepsilon|U_0|_{2,K} \\ &\leq C|U_0 - U_H|_{1,K} + C(\varepsilon + H)\|U_0\|_{2,K}. \end{aligned}$$

Summing up for all $K \in \mathcal{T}_H$ and using Theorem 1.1 for $k = 1$ and Theorem 1.2 for the case $I_\delta = x_K + \varepsilon I$, we get

$$\left(\sum_{K \in \mathcal{T}_H} \|\nabla(\hat{u}^\varepsilon - u_1^\varepsilon)\|_{0,K}^2 \right)^{1/2} \leq C(\varepsilon + H),$$

which together with (4.5) and the classical estimate for $u^\varepsilon - u_1^\varepsilon$ [5, 32, 45], i.e.,

$$\|u^\varepsilon - u_1^\varepsilon\|_1 \leq C\sqrt{\varepsilon},$$

gives

$$\begin{aligned} \left(\sum_{K \in \mathcal{T}_H} \|\nabla(u^\varepsilon - \tilde{u}^\varepsilon)\|_{0,K}^2 \right)^{1/2} &\leq \|u^\varepsilon - u_1^\varepsilon\|_1 + \left(\sum_{K \in \mathcal{T}_H} \|\nabla(u_1^\varepsilon - \hat{u}^\varepsilon)\|_{0,K}^2 \right)^{1/2} \\ &\quad + \left(\sum_{K \in \mathcal{T}_H} \|\nabla(\hat{u}^\varepsilon - \tilde{u}^\varepsilon)\|_{0,K}^2 \right)^{1/2} \\ &\leq C(\sqrt{\varepsilon} + H). \end{aligned}$$

□

Corollary 4.1.

$$(4.6) \quad \|\tilde{u}^\varepsilon - u^\varepsilon\|_0 \leq C(\varepsilon + H^2).$$

Proof. Using the definition of \tilde{u}^ε , we have $\int_{I_\varepsilon(x_K)} (\tilde{u}^\varepsilon - U_H)(x) dx = 0$. Together with (4.3), we have

$$\int_{I_\varepsilon(x_K)} (\tilde{u}^\varepsilon - \hat{u}^\varepsilon)(x) dx = 0.$$

An application of the *Poincaré inequality* gives

$$\|\tilde{u}^\varepsilon - \hat{u}^\varepsilon\|_{0, I_\varepsilon(x_K)} \leq C\varepsilon \|\nabla(\tilde{u}^\varepsilon - \hat{u}^\varepsilon)\|_{0, I_\varepsilon(x_K)} \leq C\varepsilon^2 \|\nabla U_H\|_{0, I_\varepsilon(x_K)}.$$

As before for any $I_\varepsilon(x_1)$, $\|\tilde{u}^\varepsilon - \hat{u}^\varepsilon\|_{0, I_\varepsilon(x_1)} = \|\tilde{u}^\varepsilon - \hat{u}^\varepsilon\|_{0, I_\varepsilon(x_K)}$, note that ∇U_H is a constant on K . We obtain

$$(4.7) \quad \|\tilde{u}^\varepsilon - \hat{u}^\varepsilon\|_{0, K} \leq C\varepsilon^2 \|\nabla U_H\|_{0, K}.$$

On each element K , we have

$$\|\hat{u}^\varepsilon - U_H\|_{0, K} \leq C\varepsilon \|\nabla U_H\|_{0, K}.$$

Combining the above and summing up for all $K \in \mathcal{T}_H$, we get

$$\|\tilde{u}^\varepsilon - U_H\|_0 \leq C\varepsilon \|\nabla U_H\|_0 \leq C\varepsilon,$$

which together with

$$\|u^\varepsilon - U_H\|_0 \leq \|u^\varepsilon - U_0\|_0 + \|U_0 - U_H\|_0 \leq C(\varepsilon + H^2)$$

leads to (4.6), where we have used the estimate for U_0 [5, 32, 45], i.e.,

$$\|u^\varepsilon - U_0\|_0 \leq C\varepsilon. \quad \square$$

4.2. Compression operator. The compression operator (denoted by Q) maps the microvariables to the macrovariables [16]. It plays an important role in the general framework of HMM, even though for the present problem HMM can be formulated without explicitly specifying the compression operator beforehand. Typically the compression operator is some spatial/temporal averaging, or projection to some slow manifolds. It is of interest to consider the error bound for $Qu^\varepsilon - U_H$. We first list some natural properties of the *compression operator*.

- For any $\phi \in X$, $Q\phi \in X_H$.
- There exists a constant C such that

$$\|Q\phi\|_0 \leq C\|\phi\|_0.$$

- For any $k \geq 1$, if $\phi \in H^{k+1}(\Omega) \cap H_0^1(\Omega)$, then

$$\|\phi - Q\phi\|_0 \leq CH^{k+1}|\phi|_{k+1}.$$

Theorem 4.2. *Assume that Q satisfies all three requirements and $U_0 \in H^{k+1}(D)$ for any $k \geq 1$. Then*

$$(4.8) \quad \|Qu^\varepsilon - U_H\|_0 \leq C(\varepsilon + H^{k+1}).$$

Moreover, if \mathcal{T}_H is quasi-uniform, then

$$(4.9) \quad \|Qu^\varepsilon - U_H\|_1 \leq C\left(\frac{\varepsilon}{H} + H^k\right).$$

Proof. We decompose $Qu^\varepsilon - U_H$ into

$$(4.10) \quad Qu^\varepsilon - U_H = Q(u^\varepsilon - U_0) + (QU_0 - U_0) + (U_0 - U_H).$$

Using the fact that Q is bounded in L_{norm}^2 , we obtain

$$\|Q(u^\varepsilon - U_0)\|_0 \leq C\|u^\varepsilon - U_0\|_0 \leq C\varepsilon.$$

Using the third property of Q , we have

$$\|QU_0 - U_0\|_0 \leq CH^{k+1}.$$

Using Theorem 1.1 and the first estimate in Theorem 1.2, we have

$$\|U_0 - U_H\|_0 \leq C(\varepsilon + H^{k+1}).$$

A combination of these three estimates implies (4.8), which together with the inverse inequality (cf. Theorem 1.7) leads to (4.9). \square

It remains to give some examples of the *compression operator*. The following two types of operators meet all three requirements:

- the L^2 -projection operator onto X_H ,
- the Clément-type interpolation operator [12].

Remark 4.3. Notice that in one dimension, the standard Lagrange interpolant does not meet the second requirement. However, it is still possible to derive (4.9) via another approach. Moreover, a careful study of one dimensional examples shows that the term ε/H in (4.9) is sharp.

5. NONLINEAR HOMOGENIZATION PROBLEMS

5.1. **Algorithms and main results.** We consider the following nonlinear problem which has been discussed in [6, 23]:

$$(5.1) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x, u^\varepsilon(x)) \nabla u^\varepsilon(x)) = f(x) & x \in D, \\ u^\varepsilon(x) = 0 & x \in \partial D. \end{cases}$$

In this section, we define $X := W_0^{1,p}(D)$ with $p > 1$ and X_H is defined as the \mathcal{P}_k finite element subspace of X .

We assume that $a^\varepsilon(x, u^\varepsilon)$ satisfies

$$\lambda |\xi|^2 \leq a_{ij}^\varepsilon \xi_i \xi_j \leq \Lambda |\xi|^2 \quad \text{for all } \xi \in \mathbb{R}^d,$$

with $0 < \lambda \leq \Lambda$. Moreover, we assume that $a^\varepsilon(x, z)$ is equi-continuous in z uniformly with respect to x and ε .

The homogenized problem, if it exists, is of the following form:

$$\begin{cases} \mathcal{L}U_0 := -\operatorname{div}(\mathcal{A}(x, U_0(x)) \nabla U_0(x)) = f(x) & x \in D, \\ U_0(x) = 0 & x \in \partial D. \end{cases}$$

If we let

$$A(v, w) = (\mathcal{A}(x, v) \nabla v, \nabla w) \quad \text{for all } v, w \in X,$$

then

$$(5.2) \quad A(U_0, v) = (f, v) \quad \text{for all } v \in X',$$

where X' is the dual space of X .

The linearized operator of \mathcal{L} at U_0 is defined for any $v \in H_0^1(D)$ by

$$\mathcal{L}_{\text{lin}}(U_0)v = -\operatorname{div}(\mathcal{A}(x, U_0) \nabla v + \mathcal{A}_p(x, U_0) \nabla U_0 v),$$

where $\mathcal{A}_p(x, z) := \nabla_z \mathcal{A}(x, z)$. \mathcal{L}_{lin} induces a bilinear form through

$$\hat{A}(u; v, w) := (\mathcal{A}(x, u) \nabla v, \nabla w) + (\mathcal{A}_p(x, u) \nabla u v, \nabla w) \quad \text{for all } v, w \in H_0^1(D).$$

Our basic assumption is that the linearized operator \mathcal{L}_{lin} is an isomorphism from $H_0^1(D)$ to $H^{-1}(D)$, so U_0 must be an isolated solution of (5.2).

To formulate HMM, for each quadrature point x_ℓ , define v_ℓ^ε to be the solutions of

$$(5.3) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x, v_\ell^\varepsilon) \nabla v_\ell^\varepsilon(x)) = 0 & x \in I_\delta(x_\ell), \\ v_\ell^\varepsilon(x) = V_\ell(x) & x \in \partial I_\delta(x_\ell). \end{cases}$$

We can define w_ℓ^ε similarly.

For any $V, W \in X_H$, define

$$\nabla W(x_\ell) \cdot \mathcal{A}_H(x_\ell, V(x_\ell)) \nabla V(x_\ell) = \int_{I_\delta(x_\ell)} \nabla w_\ell^\varepsilon(x) \cdot a^\varepsilon(x, v_\ell^\varepsilon(x)) \nabla v_\ell^\varepsilon(x) dx$$

and

$$A_H(V, W) := \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell \nabla W(x_\ell) \cdot \mathcal{A}_H(x_\ell, V(x_\ell)) \nabla V(x_\ell).$$

The HMM solution is given by the problem:

Problem 5.1. Find $U_H \in X_H$ such that

$$(5.4) \quad A_H(U_H, V) = (f, V) \quad \text{for all } V \in X_H.$$

For any $v, v_H, w \in X$, define

$$(5.5) \quad R(v, v_H, w) := A(v_H, w) - A(v, w) - \hat{A}(v; v_H - v, w).$$

It is easy to see that for any v and v_H satisfying $\|v\|_{1, \infty} + \|v_H\|_{1, \infty} \leq M$,

$$(5.6) \quad |R(v, v_H, w)| \leq C(M)(\|e_H\|_{0, 2p}^2 + \|e_H \nabla e_H\|_{0, p}) \|\nabla w\|_{0, q}$$

for $e_H := v - v_H$ and $\frac{1}{p} + \frac{1}{q} = 1, p, q \geq 1$ (see [42, Lemma 3.1] for a similar result). Therefore we have

Lemma 5.2. $U_H \in X_H$ is the solution of Problem 5.1 if and only if

$$(5.7) \quad \begin{aligned} \hat{A}(U_0; U_0 - U_H, V) &= R(U_0, U_H, V) \\ &+ A_H(U_H, V) - A(U_H, V) \quad \text{for all } V \in X_H. \end{aligned}$$

For any $V, W \in X_H$, define

$$(5.8) \quad E(V, W) := \nabla W(x_\ell) \cdot (\mathcal{A}_H - \mathcal{A})(x_\ell, V(x_\ell)) \nabla V(x_\ell).$$

Define $e(\text{HMM})$ as

$$(5.9) \quad e(\text{HMM}) = \max_{\substack{x_\ell \in K, K \in \mathcal{T}_H \\ V \in X_H \cap W^{1, \infty}(D), W \in X_H}} \frac{|E(V, W)|}{|\nabla V_\ell| |\nabla W_\ell|}.$$

The existence and uniqueness of the solution of Problem 5.1 are proved in the following lemma.

Lemma 5.3. Assume that $U_0 \in W^{2, p}(D)$ with $p > d$ and \mathcal{L}_{lin} is an isomorphism from $H_0^1(D)$ to $H^{-1}(D)$. If $e(\text{HMM})$ is uniformly bounded and there exist constants $H_0 > 0$ and $\mathcal{M}_1 > 0$ such that for $0 < H \leq H_0$ and

$$(5.10) \quad e(\text{HMM})^{1/2} |\ln H| \leq \mathcal{M}_1,$$

then Problem 5.1 has a solution U_H satisfying

$$(5.11) \quad \begin{aligned} \|U_H - P_H U_0\|_{1, \infty} &\leq e(\text{HMM})^{1/2} + H^{1-d/p}, \\ \|U_0 - U_H\|_{1, \infty} &\leq C(e(\text{HMM})^{1/2} + H^{1-d/p}), \end{aligned}$$

where $P_H U_0 \in X_H$ is defined as

$$(5.12) \quad \hat{A}(U_0; P_H U_0, V) = \hat{A}(U_0; U_0, V) \quad \text{for all } V \in X_H.$$

Moreover, if there exists a constant $\eta(M)$ with $0 < \eta(M) < 1$ such that

$$(5.13) \quad \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell |E(V, Z) - E(W, Z)| \leq \eta(M) \|V - W\|_1 \|Z\|_1$$

for all $V, W \in X_H \cap W^{1,\infty}(D)$ and $Z \in X_H$, satisfying $\|V\|_{1,\infty} + \|W\|_{1,\infty} \leq M$, then there exists a constant $H_1 > 0$ such that for $0 < H \leq H_1$, the HMM solution U_H satisfying (5.11)₁ is locally unique.

Proof. Since \mathcal{L}_{in} is an isomorphism from $H_0^1(D)$ to $H^{-1}(D)$, there exists a constant C such that

$$\sup_{W \in H_0^1(D)} \frac{\hat{A}(U_0; V, W)}{\|W\|_1} \geq C \|V\|_1 \quad \text{for all } V \in H_0^1(D).$$

Using [42, Lemma 2.2], we conclude that there exists a constant $H_2 > 0$ such that for $0 < H \leq H_2$,

$$(5.14) \quad \sup_{W \in X_H} \frac{\hat{A}(U_0; V, W)}{\|W\|_1} \geq C \|V\|_1 \quad \text{for all } V \in X_H.$$

Therefore there is a unique solution $P_H U_0 \in X_H$ satisfying (5.12) and

$$(5.15) \quad \|U_0 - P_H U_0\|_{1,\infty} \leq C H^{1-d/p}.$$

Moreover, let \hat{G}_H^z be the finite element approximation of the regularized Green's function associated with $\hat{A}(U_0; \cdot, \cdot)$. Using [42, equation 2.11], or using (5.14), similarly to (2.11), we have

$$(5.16) \quad \|\hat{G}_H^z\|_{1,1} \leq C |\ln H|.$$

Define a nonlinear mapping $T: X_H \rightarrow X_H$ by

$$\hat{A}(U_0; T(V), W) = \hat{A}(U_0; U_0, W) - R(U_0, V, W) + A(V, W) - A_H(V, W),$$

for any $W \in X_H$. Obviously T is continuous due to (5.14) and (5.6).

Let

$$B := \{ V \in X_H \mid \|V - P_H U_0\|_{1,\infty} \leq e(\text{HMM})^{1/2} + H^{1-d/p} \}.$$

We next prove that there exists a constant $H_0 > 0$ such that for all $0 < H \leq H_0$, $T(B) \subset B$.

Notice that

$$\hat{A}(U_0; T(V) - P_H U_0, W) = -R(U_0, V, W) + A(V, W) - A_H(V, W).$$

Taking $W = \hat{G}_H^z$ in the above equation, using (5.16) and (5.6), we obtain, for $\|V\|_{1,\infty} \leq M$,

$$\begin{aligned} \|T(V) - P_H U_0\|_{1,\infty} &\leq C(M) \|U_0 - V\|_{1,\infty}^2 |\ln H| + C(e(\text{HMM}) + H) |\ln H| \|V\|_{1,\infty} \\ &\leq C(M) (\|U_0 - P_H U_0\|_{1,\infty}^2 + \|P_H U_0 - V\|_{1,\infty}^2) |\ln H| + CM(e(\text{HMM}) + H) |\ln H| \\ &\leq C(M) (e(\text{HMM}) + H^{2-2d/p}) |\ln H| + CM(e(\text{HMM}) + H) |\ln H|. \end{aligned}$$

Since $V \in B$ and $e(\text{HMM})$ is uniformly bounded, e.g., $e(\text{HMM}) \leq M_1$, we have

$$(5.17) \quad \begin{aligned} \|V\|_{1,\infty} &\leq \|V - P_H U_0\|_{1,\infty} + \|P_H U_0\|_{1,\infty} \leq C(U_0) + e(\text{HMM})^{1/2} \\ &\leq C(U_0) + M_1^{1/2} =: M_0. \end{aligned}$$

Combining the above two estimates, we obtain

$$\|T(V) - P_H U_0\|_{1,\infty} \leq C(M_0)(e(\text{HMM}) + H^{2-2d/p} + H)|\ln H|.$$

Define $\mathcal{M}_1 := 1/C(M_0)$. Using (5.10), we obtain

$$\|T(V) - P_H U_0\|_{1,\infty} \leq e(\text{HMM})^{1/2} + C(M_0)(H^{2-2d/p} + H)|\ln H|.$$

Therefore there exists a constant H_3 such that for $0 < H \leq H_3$, we have

$$\|T(V) - P_H U_0\|_{1,\infty} \leq e(\text{HMM})^{1/2} + H^{1-d/p}.$$

Let $H_0 := \min(H_2, H_3)$. Then for $0 < H \leq H_0$, we have $T(B) \subset B$. An application of Brouwer's fixed point theorem gives the existence of a $U_H \in B$ such that $T(U_H) = U_H$. By definition, U_H satisfies (5.11)₁. Together with (5.15) it yields (5.11)₂.

To prove uniqueness, assume that both U_H and \hat{U}_H are solutions of (5.4) satisfying (5.11)₁. Using (5.14), we obtain

$$\begin{aligned} C\|U_H - \hat{U}_H\|_1 &\leq \sup_{W \in X_H} \frac{\int_0^1 \hat{A}(U_H^t; U_H - \hat{U}_H, W) dt}{\|W\|_1} \\ &\leq \sup_{W \in X_H} \frac{|A(U_H, W) - A(\hat{U}_H, W)|}{\|W\|_1}, \end{aligned}$$

where $U_H^t = (1-t)\hat{U}_H + tU_H$. Note that

$$\begin{aligned} A(U_H, W) - A(\hat{U}_H, W) &= (A(U_H, W) - A_H(U_H, W)) \\ &\quad - (A(\hat{U}_H, W) - A_H(\hat{U}_H, W)). \end{aligned}$$

Since both U_H and \hat{U}_H sit in the set B , we can use (5.17) to get $\|U_H\|_{1,\infty} + \|\hat{U}_H\|_{1,\infty} \leq 2M_0$. Using (5.13), we obtain

$$\|U_H - \hat{U}_H\|_1 \leq (\eta(2M_0) + C_1 H)\|U_H - \hat{U}_H\|_1.$$

If we choose H_1 such that

$$\eta(2M_0) + C_1 H_1 < 1,$$

then if $H < H_1$, we have $U_H = \hat{U}_H$. Therefore the HMM solution is locally unique. \square

From here on, when we talk about the HMM solution, we are referring to this particular solution that satisfies the condition in Lemma 5.3.

Based on the above lemma, we prove a nonlinear analog of Theorem 1.1.

Theorem 5.4. *Under the assumptions in Lemma 5.3, let U_0 and U_H be solutions of (5.2) and (5.4), respectively. Assume in addition that $U_0 \in W^{k+1,\infty}(D)$. Then there exist constants H_0 and \mathcal{M}_1^* such that if $0 < H < H_0$ and $\mathcal{M}_1 < \mathcal{M}_1^*$, then*

$$(5.18) \quad \|U_0 - U_H\|_1 \leq C(H^k + e(\text{HMM})),$$

$$(5.19) \quad \|U_0 - U_H\|_{1,\infty} \leq C(H^k + e(\text{HMM}))|\ln H|.$$

Proof. Note that $U_0 \in W^{k+1,\infty}(D)$ and from (5.14), we have

$$(5.20) \quad \|U_0 - P_H U_0\|_1 \leq C H^k, \quad \|U_0 - P_H U_0\|_{1,\infty} \leq C H^k.$$

Using (5.7) with $V = P_H U_0 - U_H$ and (5.14) and (2.3), we obtain

$$(5.21) \quad \|P_H U_0 - U_H\|_1 \leq C \|U_0 - U_H\|_{1,4}^2 + C(H^k + e(\text{HMM})).$$

Using the interpolation inequality

$$\|U_0 - U_H\|_{1,4}^2 \leq \|U_0 - U_H\|_1 \|U_0 - U_H\|_{1,\infty}$$

together with (5.20) and (5.17) gives

$$\begin{aligned} \|P_H U_0 - U_H\|_1 &\leq C_1(e(\text{HMM})^{1/2} + H^{1-d/p}) \|P_H U_0 - U_H\|_1 \\ &\quad + C(H^k + e(\text{HMM})). \end{aligned}$$

Let $V = \hat{G}_H^z$ in (5.7). Using (5.16), we obtain

$$\begin{aligned} \|P_H U_0 - U_H\|_{1,\infty} &\leq C_2(e(\text{HMM}) + H) |\ln H| \|P_H U_0 - U_H\|_{1,\infty} \\ &\quad + C(\|U_0 - U_H\|_{1,\infty}^2 + e(\text{HMM}) + H^k) |\ln H|. \end{aligned}$$

Since (5.10) holds, using (5.11)₂ and (5.20)₂, we obtain

$$\begin{aligned} \|P_H U_0 - U_H\|_{1,\infty} &\leq C_2(e(\text{HMM})^{1/2} + H^{1-d/p}) |\ln H| \|P_H U_0 - U_H\|_{1,\infty} \\ &\quad + C(e(\text{HMM}) + H^k) |\ln H|. \end{aligned}$$

Now we choose

$$\mathcal{M}_1^* := \min\left(\frac{|\ln H|}{2C_1}, \frac{1}{2C_2}\right)$$

and H_0 such that $e(\text{HMM})^{1/2} |\ln H| \leq M_1^*$, and therefore

$$\begin{aligned} C_1(e(\text{HMM})^{1/2} + H^{1-d/p}) &\leq \frac{1}{2} + C_1 H_0^{1-d/p} < 1, \\ C_2(e(\text{HMM})^{1/2} + H^{1-d/p}) |\ln H| &\leq \frac{1}{2} + C_2 H^{1-d/p} |\ln H| < 1. \end{aligned}$$

Thus we obtain

$$\begin{aligned} \|P_H U_0 - U_H\|_1 &\leq C(H^k + e(\text{HMM})), \\ \|P_H U_0 - U_H\|_{1,\infty} &\leq C(H^k + e(\text{HMM})) |\ln H|. \end{aligned}$$

Using (5.20) once again gives (5.18) and (5.19). □

5.2. Estimating $e(\text{HMM})$. It remains to estimate $e(\text{HMM})$ and verify assumptions (5.10) and (5.13). We assume that $a^\varepsilon(x, u^\varepsilon) = (a_{ij}(x, x/\varepsilon, u^\varepsilon))$, and for $1 \leq i, j \leq d$, the coefficients $a_{ij}^\varepsilon(x, y, z)$ are smooth in x, z and periodic in y with period I . These types of problems, among others, have been considered in [5, 6, 23]. The homogenized coefficient $\mathcal{A} = (\mathcal{A}_{ij}(x, p))$ is given for any $p \in \mathbb{R}$ by

$$\mathcal{A}_{ij}(x, p) = \int_I \left(a_{ij} + a_{ik} \frac{\partial \chi^j}{\partial y_k} \right) (x, y, p) dy,$$

where $\{\chi^k\}_{k=1}^d$ is defined for any $p \in \mathbb{R}$ by

$$(5.22) \quad -\frac{\partial}{\partial y_i} \left(a_{ik} \frac{\partial \chi^j}{\partial y_k} \right) (x, y, p) = \frac{\partial}{\partial y_i} a_{ij}(x, y, p),$$

with the periodic boundary condition in y and $\int_I \chi^k(x, y, p) dy = 0$. It is clear that $\mathcal{A}(x, p)$ is also smooth in x and p and satisfies [6, Proposition 3.5]

$$(5.23) \quad \lambda I \leq \mathcal{A} \leq \frac{\Lambda^2}{\lambda} I.$$

Using Lemma 1.9 and Remark 1.10, for the solution of (5.3), we have

$$(5.24) \quad \|\nabla V_\ell\|_{0, I_\delta} \leq \|\nabla v_\ell^\varepsilon\|_{0, I_\delta} \leq \sqrt{\Lambda/\lambda} \|\nabla V_\ell\|_{0, I_\delta}.$$

This gives a bound for \mathcal{A}_H :

$$\lambda I \leq \mathcal{A}_H \leq \frac{\Lambda^2}{\lambda} I,$$

which together with (5.23) implies

$$(5.25) \quad e(\text{HMM}) \leq 2\Lambda^2/\lambda.$$

This shows that $e(\text{HMM})$ is uniformly bounded.

To simplify the presentation, we will show how to estimate $e(\text{HMM})$ when (5.3) is changed slightly to

$$(5.26) \quad \begin{cases} -\operatorname{div}(a^\varepsilon(x, V(x_\ell)) \nabla v_\ell^\varepsilon(x)) = 0 & x \in I_\delta(x_\ell), \\ v_\ell^\varepsilon(x) = V_\ell(x) & x \in \partial I_\delta(x_\ell) \end{cases}$$

and $A_H(V, W)$ is changed to

$$A_H(V, W) = \sum_{K \in \mathcal{T}_H} |K| \sum_{x_\ell \in K} \omega_\ell \int_{I_\delta(x_\ell)} \nabla w_\ell^\varepsilon(x) \cdot a^\varepsilon(x, V(x_\ell)) \nabla v_\ell^\varepsilon(x) dx.$$

If $\delta = \varepsilon$, we replace the Dirichlet boundary condition in (5.26) by the periodic boundary condition, i.e., $v_\ell^\varepsilon(x) - V_\ell(x)$ is periodic on $\partial I_\varepsilon(x_\ell)$.

Theorem 5.5. *If $(\sqrt{\varepsilon/\delta} + \delta)|\ln H|$ is sufficiently small, then (5.10) and (5.13) hold and*

$$(5.27) \quad e(\text{HMM}) \leq C \left(\left(\frac{\varepsilon}{\delta} \right)^{1/2} + \delta \right).$$

In the case of $\delta = \varepsilon$, if $\varepsilon|\ln H|$ is sufficiently small, then (5.10) and (5.13) hold and

$$(5.28) \quad e(\text{HMM}) \leq C\varepsilon.$$

In what follows, we concentrate on the first case. The second case when $\delta = \varepsilon$ will be commented on.

Let us first fix more notation. Denote by \hat{v}_ℓ^ε the solutions of (5.26) with the coefficient $a(x, x/\varepsilon, V(x_\ell))$ replaced by $a(x_\ell, x/\varepsilon, V(x_\ell))$. Similarly we define w_ℓ^ε to be the solution of (5.26) with V replaced by $W \in X_H$. Also, \hat{w}_ℓ^ε can be defined in the same way and w_ℓ^ε and \hat{w}_ℓ^ε can be viewed as the perturbations of v_ℓ^ε and \hat{v}_ℓ^ε , respectively. Moreover, we define

$$\begin{aligned} a_V(x_\ell) &= a(x_\ell, \frac{x}{\varepsilon}, V(x_\ell)), & a_W(x_\ell) &= a(x_\ell, \frac{x}{\varepsilon}, W(x_\ell)), \\ a_V(x) &= a(x, \frac{x}{\varepsilon}, V(x_\ell)), & a_W(x) &= a(x, \frac{x}{\varepsilon}, W(x_\ell)). \end{aligned}$$

Observe that \hat{v}_ℓ^ε and \hat{w}_ℓ^ε also satisfy (5.24), and using Lemma 1.8, we have

$$(5.29) \quad \|\nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon)\|_{0, I_\delta} \leq C\delta \|\nabla V_\ell\|_{0, I_\delta}, \quad \|\nabla(w_\ell^\varepsilon - \hat{w}_\ell^\varepsilon)\|_{0, I_\delta} \leq C\delta \|\nabla W_\ell\|_{0, I_\delta}.$$

Lemma 5.6. *We have*

$$(5.30) \quad \begin{aligned} \|\nabla(\hat{v}_\ell^\varepsilon - \hat{w}_\ell^\varepsilon)\|_{0,I_\delta} &\leq C(|(V - W)(x_\ell)|(\|\nabla V_\ell\|_{0,I_\delta} + \|\nabla W_\ell\|_{0,I_\delta}) \\ &\quad + \|\nabla(V_\ell - W_\ell)\|_{0,I_\delta}). \end{aligned}$$

Proof. Observe that

$$\begin{aligned} -\operatorname{div}(a_V(x_\ell)\nabla(\hat{v}_\ell^\varepsilon - V_\ell)) &= \operatorname{div}(a_V(x_\ell)\nabla V_\ell), \\ -\operatorname{div}(a_W(x_\ell)\nabla(\hat{w}_\ell^\varepsilon - W_\ell)) &= \operatorname{div}(a_W(x_\ell)\nabla W_\ell). \end{aligned}$$

Both $\hat{v}_\ell^\varepsilon - V_\ell$ and $\hat{w}_\ell^\varepsilon - W_\ell$ vanish on $\partial I_\delta(x_\ell)$. Using Lemma 1.8, we obtain

$$\begin{aligned} \lambda\|\nabla(\hat{v}_\ell^\varepsilon - V_\ell - \hat{w}_\ell^\varepsilon + W_\ell)\|_{0,I_\delta} &\leq \Lambda\|\nabla(V_\ell - W_\ell)\|_{0,I_\delta} \\ &\quad + \max_{x \in I_\delta} |(a_V - a_W)(x_\ell)|(\|\nabla V_\ell\|_{0,I_\delta} + \|\nabla(\hat{w}_\ell^\varepsilon - W_\ell)\|_{0,I_\delta}). \end{aligned}$$

Using (5.24) for \hat{w}_ℓ^ε , we obtain

$$\|\nabla(\hat{w}_\ell^\varepsilon - W_\ell)\|_{0,I_\delta} \leq \|\nabla\hat{w}_\ell^\varepsilon\|_{0,I_\delta} + \|\nabla W_\ell\|_{0,I_\delta} \leq C(1 + \sqrt{\Lambda/\lambda})\|\nabla W_\ell\|_{0,I_\delta},$$

which together with $\max_{x \in I_\delta} |(a_V - a_W)(x_\ell)| \leq C|(V - W)(x_\ell)|$ gives (5.30). \square

Next we establish the estimate for $(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) - (w_\ell^\varepsilon - \hat{w}_\ell^\varepsilon)$. Let $\psi_\ell^\varepsilon := v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon$ and $\tilde{\psi}_\ell^\varepsilon := w_\ell^\varepsilon - \hat{w}_\ell^\varepsilon$. Clearly, $\psi_\ell^\varepsilon, \tilde{\psi}_\ell^\varepsilon$ vanish on $\partial I_\delta(x_\ell)$ and satisfy

$$\begin{aligned} -\operatorname{div}(a_V(x)\nabla\psi_\ell^\varepsilon) &= \operatorname{div}((a_V(x) - a_V(x_\ell))\nabla\hat{v}_\ell^\varepsilon) \quad x \in I_\delta(x_\ell), \\ -\operatorname{div}(a_W(x)\nabla\tilde{\psi}_\ell^\varepsilon) &= \operatorname{div}((a_W(x) - a_W(x_\ell))\nabla\hat{w}_\ell^\varepsilon) \quad x \in I_\delta(x_\ell). \end{aligned}$$

Lemma 5.7. *We have*

$$(5.31) \quad \begin{aligned} \|\nabla(\psi_\ell^\varepsilon - \tilde{\psi}_\ell^\varepsilon)\|_{0,I_\delta} &\leq C\delta(|(V - W)(x_\ell)|(\|\nabla V_\ell\|_{0,I_\delta} + \|\nabla W_\ell\|_{0,I_\delta}) \\ &\quad + \|\nabla(V_\ell - W_\ell)\|_{0,I_\delta}). \end{aligned}$$

Proof. Using Lemma 1.8, we have

$$\begin{aligned} \lambda\|\nabla(\psi_\ell^\varepsilon - \tilde{\psi}_\ell^\varepsilon)\|_{0,I_\delta} &\leq \max_{x \in I_\delta} |a_V(x) - a_V(x_\ell) - a_W(x) + a_W(x_\ell)| \|\nabla\hat{v}_\ell^\varepsilon\|_{0,I_\delta} \\ &\quad + \max_{x \in I_\delta} |a_V(x) - a_W(x)| \|\nabla\tilde{\psi}_\ell^\varepsilon\|_{0,I_\delta} \\ &\quad + \max_{x \in I_\delta} |a_W(x) - a_W(x_\ell)| \|\nabla(\hat{v}_\ell^\varepsilon - \hat{w}_\ell^\varepsilon)\|_{0,I_\delta}. \end{aligned}$$

Using

$$(5.32) \quad \max_{x \in I_\delta} |a_V(x) - a_V(x_\ell) - a_W(x) + a_W(x_\ell)| \leq C\delta|(V - W)(x_\ell)|,$$

from (5.29)₂, we have $\|\nabla\tilde{\psi}_\ell^\varepsilon\|_{0,I_\delta} \leq C\delta\|\nabla W_\ell\|_{0,I_\delta}$. Collecting the above estimates and using (5.30), we obtain (5.31). \square

Define

$$\begin{aligned} \hat{V}_\ell(x) &= V_\ell(x) + \varepsilon\chi_\ell^k\left(x_\ell, \frac{x}{\varepsilon}, V(x_\ell)\right)\frac{\partial V_\ell}{\partial x_k}(x), \\ \hat{W}_\ell(x) &= W_\ell(x) + \varepsilon\chi_\ell^k\left(x_\ell, \frac{x}{\varepsilon}, W(x_\ell)\right)\frac{\partial W_\ell}{\partial x_k}(x), \end{aligned}$$

where $\{\chi_\ell^k\}_{k=1}^d$ are the solutions of (5.22) with coefficient replaced by $a_{ij}(x_\ell, y, p)$.

Denote $\theta_\ell^\varepsilon = \hat{v}_\ell^\varepsilon - \hat{V}_\ell$ and $\tilde{\theta}_\ell^\varepsilon = \hat{w}_\ell^\varepsilon - \hat{W}_\ell$. Observe that

$$\begin{cases} -\operatorname{div}(a_V(x_\ell)\nabla\theta_\ell^\varepsilon(x)) = 0 & x \in I_\delta(x_\ell), \\ \theta_\ell^\varepsilon(x) = -\varepsilon\chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, V(x_\ell))\frac{\partial V_\ell}{\partial x_k} & x \in \partial I_\delta(x_\ell), \end{cases}$$

and

$$\begin{cases} -\operatorname{div}(a_W(x_\ell)\nabla\tilde{\theta}_\ell^\varepsilon(x)) = 0 & x \in I_\delta(x_\ell), \\ \tilde{\theta}_\ell^\varepsilon(x) = -\varepsilon\chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, W(x_\ell))\frac{\partial W_\ell}{\partial x_k} & x \in \partial I_\delta(x_\ell). \end{cases}$$

Similarly to (3.9), we have

$$(5.33) \quad \|\nabla\theta_\ell^\varepsilon\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2}\|\nabla V_\ell\|_{0,I_\delta}, \quad \|\nabla\tilde{\theta}_\ell^\varepsilon\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2}\|\nabla W_\ell\|_{0,I_\delta}.$$

Let ρ^ε be defined as in (3.14) and define

$$\begin{aligned} \varphi_\ell^\varepsilon &= \theta_\ell^\varepsilon + \varepsilon\chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, V(x_\ell))\frac{\partial V_\ell}{\partial x_k}(1 - \rho^\varepsilon), \\ \tilde{\varphi}_\ell^\varepsilon &= \tilde{\theta}_\ell^\varepsilon + \varepsilon\chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, W(x_\ell))\frac{\partial W_\ell}{\partial x_k}(1 - \rho^\varepsilon). \end{aligned}$$

Observe that φ_ℓ^ε and $\tilde{\varphi}_\ell^\varepsilon$ vanish on $\partial I_\delta(x_\ell)$ and satisfy

$$\begin{aligned} -\operatorname{div}(a_V(x_\ell)\nabla\varphi_\ell^\varepsilon(x)) &= \operatorname{div}(a_V(x_\ell)\nabla(\varphi_\ell^\varepsilon - \theta_\ell^\varepsilon)(x)) & x \in I_\delta(x_\ell), \\ -\operatorname{div}(a_W(x_\ell)\nabla\tilde{\varphi}_\ell^\varepsilon(x)) &= \operatorname{div}(a_W(x_\ell)\nabla(\tilde{\varphi}_\ell^\varepsilon - \tilde{\theta}_\ell^\varepsilon)(x)) & x \in I_\delta(x_\ell). \end{aligned}$$

Lemma 5.8. *We have*

$$(5.34) \quad \begin{aligned} \|\nabla(\theta_\ell^\varepsilon - \tilde{\theta}_\ell^\varepsilon)\|_{0,I_\delta} &\leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2}\left(\|\nabla(V_\ell - W_\ell)\|_{0,I_\delta} \right. \\ &\quad \left. + |(V - W)(x_\ell)|(\|\nabla V_\ell\|_{0,I_\delta} + \|\nabla W_\ell\|_{0,I_\delta})\right). \end{aligned}$$

Proof. Using Lemma 1.8, we obtain

$$\begin{aligned} \lambda\|\nabla(\varphi_\ell^\varepsilon - \tilde{\varphi}_\ell^\varepsilon)\|_{0,I_\delta} &\leq \max_{x \in I_\delta} |a_V(x_\ell) - a_W(x_\ell)|(\|\nabla(\varphi_\ell^\varepsilon - \theta_\ell^\varepsilon)\|_{0,I_\delta} + \|\nabla\tilde{\varphi}_\ell^\varepsilon\|_{0,I_\delta}) \\ &\quad + \max_{x \in I_\delta} |a_W(x_\ell)|\|\nabla(\varphi_\ell^\varepsilon - \theta_\ell^\varepsilon - \tilde{\varphi}_\ell^\varepsilon + \tilde{\theta}_\ell^\varepsilon)\|_{0,I_\delta}. \end{aligned}$$

A direct computation gives that

$$\|\nabla(\varphi_\ell^\varepsilon - \theta_\ell^\varepsilon)\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{\frac{1}{2}}\|\nabla V_\ell\|_{0,I_\delta} \quad \text{and} \quad \|\nabla(\tilde{\varphi}_\ell^\varepsilon - \tilde{\theta}_\ell^\varepsilon)\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{\frac{1}{2}}\|\nabla W_\ell\|_{0,I_\delta},$$

which together with (5.33) gives

$$\|\nabla\tilde{\varphi}_\ell^\varepsilon\|_{0,I_\delta} \leq \|\nabla\tilde{\theta}_\ell^\varepsilon\|_{0,I_\delta} + \|\nabla(\tilde{\varphi}_\ell^\varepsilon - \tilde{\theta}_\ell^\varepsilon)\|_{0,I_\delta} \leq C\left(\frac{\varepsilon}{\delta}\right)^{1/2}\|\nabla W_\ell\|_{0,I_\delta}.$$

Note that

$$\begin{aligned} \varphi_\ell^\varepsilon - \theta_\ell^\varepsilon - \tilde{\varphi}_\ell^\varepsilon + \tilde{\theta}_\ell^\varepsilon &= \varepsilon(1 - \rho^\varepsilon) \left(\chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, V(x_\ell)) - \chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, W(x_\ell)) \right) \frac{\partial V_\ell}{\partial x_k} \\ &\quad + \varepsilon(1 - \rho^\varepsilon) \chi_\ell^k(x_\ell, \frac{x}{\varepsilon}, W(x_\ell)) \frac{\partial}{\partial x_k} (V_\ell - W_\ell). \end{aligned}$$

Using the continuity of $\{\chi_\ell^k\}_{k=1}^d$, a direct computation gives

$$\|\nabla(\varphi_\ell^\varepsilon - \theta_\ell^\varepsilon - \tilde{\varphi}_\ell^\varepsilon + \tilde{\theta}_\ell^\varepsilon)\|_{0, I_\delta} \leq C \left(\frac{\varepsilon}{\delta} \right)^{1/2} \left(|(V - W)(x_\ell)| \|\nabla V_\ell\|_{0, I_\delta} + \|\nabla(V_\ell - W_\ell)\|_{0, I_\delta} \right).$$

Adding these up, we obtain (5.34). \square

In the next lemma, we shall prove that $E(V, W)$ has certain continuity with respect to V .

Lemma 5.9. *For any $V, W, Z \in X_H$ satisfying $\|V\|_{1, \infty} + \|W\|_{1, \infty} \leq M$, there exists a constant $C(M)$ such that (5.13) holds with $\eta(M) = C(M)(\sqrt{\varepsilon/\delta} + \delta)$.*

Proof. Using the definition of v_ℓ^ε and noticing that $v_\ell^\varepsilon = (v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) + \theta_\ell^\varepsilon + \hat{V}_\ell$, we obtain

$$\begin{aligned} \int_{I_\delta} \nabla z_\ell^\varepsilon \cdot a_V(x) \nabla v_\ell^\varepsilon \, dx &= \int_{I_\delta} \nabla Z_\ell \cdot a_V(x) \nabla v_\ell^\varepsilon \, dx \\ &= \nabla Z_\ell \cdot \left(\int_{I_\delta} a_V(x) (\nabla(v_\ell - \hat{v}_\ell^\varepsilon) + \nabla\theta_\ell^\varepsilon + \nabla\hat{V}_\ell) \, dx \right), \end{aligned}$$

which together with $\int_{I_{\kappa\varepsilon}} \nabla Z_\ell \cdot a_V(x_\ell) \nabla \hat{V}_\ell \, dx = \nabla Z(x_\ell) \cdot \mathcal{A}(x_\ell, V(x_\ell)) \nabla V(x_\ell)$ gives

$$E(V, Z) - E(W, Z) = :I_1 + I_2 + I_3,$$

with

$$\begin{aligned} I_1 &= \nabla Z_\ell \cdot \left[\int_{I_\delta} (a_V(x) \nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) - a_W(x) \nabla(w_\ell^\varepsilon - \hat{w}_\ell^\varepsilon)) \, dx \right], \\ I_2 &= \nabla Z_\ell \cdot \left[\int_{I_\delta} a_V(x) \nabla\theta_\ell^\varepsilon - a_W(x) \nabla\tilde{\theta}_\ell^\varepsilon \, dx \right], \end{aligned}$$

and

$$\begin{aligned} I_3 &= \nabla Z_\ell \cdot \left[\int_{I_\delta} a_V(x) \nabla\hat{V}_\ell \, dx - \int_{I_{\kappa\varepsilon}} a_V(x_\ell) \nabla\hat{V}_\ell \, dx \right. \\ &\quad \left. - \int_{I_\delta} a_W(x) \nabla\hat{W}_\ell \, dx + \int_{I_{\kappa\varepsilon}} a_W(x_\ell) \nabla\hat{W}_\ell \, dx \right]. \end{aligned}$$

I_1 can be decomposed into

$$I_1 = \nabla Z_\ell \cdot \left[\int_{I_\delta} (a_V(x) - a_W(x)) \nabla(v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) \, dx + \int_{I_\delta} a_W(x) \nabla(\psi_\ell^\varepsilon - \tilde{\psi}_\ell^\varepsilon) \, dx \right].$$

Using (5.29) and (5.31), we bound I_1 as

$$|I_1| \leq C\delta \left(|\nabla(V_\ell - W_\ell)| + |(V - W)(x_\ell)| (|\nabla V_\ell| + |\nabla W_\ell|) \right) |\nabla Z_\ell|.$$

Similarly, using Lemma 5.8, we bound I_2 as

$$|I_2| \leq C \left(\frac{\varepsilon}{\delta} \right)^{1/2} \left(|\nabla(V_\ell - W_\ell)| + |(V - W)(x_\ell)| (|\nabla V_\ell| + |\nabla W_\ell|) \right) |\nabla Z_\ell|.$$

I_3 can be rewritten as

$$\begin{aligned} I_3 &= \nabla Z_\ell \cdot \int_{I_\delta} \left((a_V(x) - a_V(x_\ell)) \nabla \hat{V}_\ell - (a_W(x) - a_W(x_\ell)) \nabla \hat{W}_\ell \right) dx \\ &\quad - \nabla Z_\ell \cdot \left(\frac{1}{|\kappa \varepsilon|^d} - \frac{1}{\delta^d} \right) \int_{I_\delta} (a_V(x_\ell) \nabla \hat{V}_\ell - a_W(x_\ell) \nabla \hat{W}_\ell) dx \\ &\quad + \frac{1}{|\kappa \varepsilon|^d} \nabla Z_\ell \cdot \int_{I_\delta \setminus I_{\kappa \varepsilon}} (a_V(x_\ell) \nabla \hat{V}_\ell - a_W(x_\ell) \nabla \hat{W}_\ell) dx. \end{aligned}$$

As in I_1 and using (5.32), we bound I_3 as

$$|I_3| \leq C \left(\frac{\varepsilon}{\delta} + \delta \right) \left(|\nabla(V_\ell - W_\ell)| + |(V - W)(x_\ell)| (|\nabla V_\ell| + |\nabla W_\ell|) \right) |\nabla Z_\ell|.$$

Adding these up, we get

$$\begin{aligned} |E(V, Z) - E(W, Z)| &\leq C \left(\frac{\varepsilon}{\delta} + \delta + \left(\frac{\varepsilon}{\delta} \right)^{1/2} \right) \left(|\nabla(V_\ell - W_\ell)| \right. \\ (5.35) \quad &\quad \left. + |(V - W)(x_\ell)| (|\nabla V_\ell| + |\nabla W_\ell|) \right) |\nabla Z_\ell|. \end{aligned}$$

Consequently we obtain

$$\begin{aligned} |K| \sum_{x_\ell \in K} \omega_\ell |E(V, Z) - E(W, Z)| &\leq C \left(\delta + \left(\frac{\varepsilon}{\delta} \right)^{1/2} \right) \|\nabla(V - W)\|_{0,K} \|\nabla Z\|_{0,K} \\ &\quad + C \left(\delta + \left(\frac{\varepsilon}{\delta} \right)^{1/2} \right) \|V - W\|_{L^\infty(K)} (\|\nabla V\|_{0,K} + \|\nabla W\|_{0,K}) \|\nabla Z\|_{0,K}. \end{aligned}$$

Using the inverse inequality (1.21) on each element, we obtain

$$\begin{aligned} \|V - W\|_{L^\infty(K)} \|\nabla V\|_{0,K} &\leq C H_K^{-d/2} \|V - W\|_{0,K} H_K^{d/2} \|\nabla V\|_{L^\infty(K)} \\ &= C \|V - W\|_{0,K} \|\nabla V\|_{L^\infty(K)}. \end{aligned}$$

Similarly $\|V - W\|_{L^\infty(K)} \|\nabla W\|_{0,K} \leq C \|V - W\|_{0,K} \|\nabla W\|_{L^\infty(K)}$. A combination of the above estimates gives (5.13) with $\eta(M) = C(M)(\sqrt{\varepsilon/\delta} + \delta)$. \square

Proof of the first case in Theorem 5.5. Let $W = 0$ in (5.35) and note that $V \in W^{1,\infty}(D)$. We obtain (5.27). Therefore if $(\sqrt{\varepsilon/\delta} + \delta)|\ln H|$ is sufficiently small, then $e(\text{HMM})^{1/2}|\ln H|$ can be smaller than any given threshold; this verifies (5.10).

Next let $\eta(M) = C(M)(\sqrt{\varepsilon/\delta} + \delta)$. Then if $(\sqrt{\varepsilon/\delta} + \delta)|\ln H|$ is sufficiently small, we have $\eta(M) < 1$; this verifies (5.13). \square

Remark 5.10. Compared to the linear case, the upper bound for $e(\text{HMM})$ for the case when $\delta/\varepsilon \notin \mathbb{Z}$ degrades to $\sqrt{\varepsilon/\delta}$. This is due to the fact that A_H is nonsymmetric.

In the case of $\delta = \varepsilon$, note that $\hat{v}_\ell^\varepsilon = \hat{V}_\ell$ and $\hat{w}_\ell^\varepsilon = \hat{W}_\ell$. So a direct calculation gives Lemma 5.6 for this case. Lemma 5.7 is also valid with δ replaced by ε . We also have $\theta_\ell^\varepsilon = 0$ and $\tilde{\theta}_\ell^\varepsilon = 0$. Observing that for any $V, Z \in X_H$,

$$\int_{I_\varepsilon(x_\ell)} \nabla z_\ell^\varepsilon \cdot a_V(x_\ell) \nabla \hat{v}_\ell^\varepsilon dx = \int_K \nabla Z_\ell \cdot \mathcal{A}(x_\ell, V(x_\ell)) \nabla V_\ell dx,$$

we may rewrite $E(V, Z)$ as

$$E(V, Z) = \int_{I_\varepsilon} \nabla z_\ell^\varepsilon \cdot a_V(x) \nabla (v_\ell^\varepsilon - \hat{v}_\ell^\varepsilon) dx + \int_{I_\varepsilon} \nabla z_\ell^\varepsilon \cdot (a_V(x) - a_V(x_\ell)) \nabla \hat{v}_\ell^\varepsilon dx.$$

Consequently, (5.13) holds with $\eta(M) = C(M)\varepsilon$.

Proof of the second case in Theorem 5.5. Let $W = 0$ in (5.35) and note that $V \in W^{1,\infty}(D)$. We obtain (5.28). Therefore if $\varepsilon|\ln H|$ is sufficiently small, then $e(\text{HMM})^{1/2}|\ln H|$ can be smaller than any given threshold; this verifies (5.10).

Next let $\eta(M) = C(M)\varepsilon$. Then if $\varepsilon|\ln H|$ is sufficiently small, we have $\eta(M) < 1$. This proves (5.13). \square

APPENDIX A. ESTIMATING $e(\text{HMM})$ FOR PROBLEMS WITH RANDOM COEFFICIENTS

Here we estimate $e(\text{HMM})$ for the random case. Our basic strategy follows that of [43].

Denote a probability space by (Ω, \mathcal{F}, P) and let $a(y, \omega) = (a_{ij}(y, \omega))$ be a random field whose statistics are invariant under integer shifts and which satisfies the uniform ellipticity condition that there exist constants λ and Λ such that

$$\lambda|\xi|^2 \leq a_{ij}(y, \omega)\xi_i\xi_j \leq \Lambda|\xi|^2,$$

for almost all $y \in \mathbb{R}^d$ and $\omega \in \Omega$. For $j = 1, \dots, d$, denote by $\varphi_j(y, \omega)$ the solutions of the cell problem:

$$(A.1) \quad \mathcal{L}_y \varphi_j := -\operatorname{div}_y (a_{ij}(y, \omega) \nabla_y \varphi_j) = \operatorname{div}_y (a_{ij} \cdot e_j),$$

where $\{e_j\}_{j=1}^d$ are the standard basis in \mathbb{R}^d . $\nabla \varphi_j$ is required to be stationary under integer shift. The existence of φ_j is proved in [28, 36]. In general φ_j is not stationary. The homogenized coefficient \mathcal{A} [28, 36] is given by

$$\mathcal{A} = \langle a(I + \nabla \varphi) \rangle.$$

Here and in the following, we use the notation

$$\langle f \rangle = \mathbb{E} \int_{[0,1]^d} f(y) dy$$

and

$$[f; m] = \frac{1}{m^d} \int_{[0,m]^d} f(y) dy,$$

where \mathbb{E} denotes the expectation in the probability space (Ω, \mathcal{F}, P) .

As in [43], we will consider the following auxiliary problem:

$$(A.2) \quad \mathcal{L}u + \rho u = f,$$

for any $\rho > 0$, where f is of the form

$$f = \sum_{j=1}^d D_j g_j + h,$$

with $g_j, h \in \rho_G$ which is defined as

$$\rho_G := \{\psi \mid \langle |\psi|^2 \rangle \leq G^2\},$$

and ψ is a random field whose statistics are stationary with respect to integer shifts.

The solution of (A.2) can be expressed with the help of a diffusion process η generated by the operator \mathcal{L} .

For each fixed realization of $\{a(y, \cdot)\}$, denote by η_x the diffusion process generated by \mathcal{L}_1 and starting from x at $t = 0$, and denote by M_x the expectation with respect to η_x . Let

$$\Gamma(\tau) = \int_0^\tau f(\eta(t))e^{-\rho t} dt.$$

Then it is well known [21] that the solution of (A.2) is given by

$$u(x) = M_x \Gamma(\infty).$$

The following results are either standard or proved in [28, 43].

Lemma A.1. *If u is the solution of (A.2), then there exists a constant C independent of ρ such that*

$$(A.3) \quad \langle |\nabla u|^2 \rangle + \rho \langle u^2 \rangle \leq C \left(\langle g^2 \rangle + \frac{1}{\rho} \langle h^2 \rangle \right),$$

$$(A.4) \quad \langle (M_x \Gamma(\infty))^2 \rangle^{1/2} \leq C \frac{G^2}{\rho},$$

$$(A.5) \quad \langle M_x (\Gamma(\infty) - \Gamma(s))^2 \rangle \leq C \frac{G^2}{\rho} e^{-2s\rho}.$$

Because of the lowest order term ρu , the Green's function associated with the operator $\mathcal{L}_1 + \rho I$ decays exponentially with rate $\mathcal{O}(\sqrt{\rho})$. To make this statement precise, we define the norm $\|x\| := \max_i |x_i|$, and

$$Q_\rho := \left\{ x \in \mathbb{R}^d \mid \|x\| \leq \left(\frac{1}{\rho} \right)^{\frac{1}{2} \lceil \ln(1/\rho) \rceil^{1/2}} \right\}.$$

Let τ be the first exit time of Q_ρ starting at $x \in Q_\rho$. Let $\hat{\varphi}_\rho(x) = M_x \Gamma(\tau)$.

Lemma A.2 ([43]). *If ρ is sufficiently small, then*

$$(A.6) \quad \mathbb{E} \int_{\|x\| \leq 10} |\varphi_\rho(x) - \hat{\varphi}_\rho(x)|^2 dx \leq C G^2 e^{-C |\ln(1/\rho)|^2},$$

$$(A.7) \quad \mathbb{E} \int_{\|x\| \leq 1} |\nabla \varphi_\rho(x) - \nabla \hat{\varphi}_\rho(x)|^2 dx \leq C G^2 e^{-C |\ln(1/\rho)|^2}.$$

To prepare for the discussion on the consequence of the mixing condition, we mention

Lemma A.3 ([43]). *Let $\{a_{ij}, g_j\}$ and $\{\tilde{a}_{ij}, \tilde{g}_j\}$ be two sets of data such that*

$$\{a_{ij}(y), g_j(y)\} = \{\tilde{a}_{ij}(y), \tilde{g}_j(y)\}$$

for $y \notin B$, where B is a domain in \mathbb{R}^d , and let φ_ρ and $\hat{\varphi}_\rho$ be the solutions of (A.2) associated with $\{a_{ij}, g_j\}$ and $\{\tilde{a}_{ij}, \tilde{g}_j\}$, respectively (with $h = 0$). Then

$$(A.8) \quad \int_{\mathbb{R}^d} |\varphi_\rho(x) - \tilde{\varphi}_\rho(x)|^2 dx \leq \frac{C}{\rho} \int_{\mathbb{R}^d} (G^2 + |\nabla \varphi_\rho(x)|^2) I_B(x) dx,$$

where I_B is the indicator function of the domain B .

Now we introduce the crucial mixing condition [22]. Let B be a domain in \mathbb{R}^d . Denote by $\mathcal{F}(B)$ the σ -algebra generated by $\{a_{ij}(y, \omega), y \in B\}$. Let ξ, η be two random variables that are measurable with respect to $\mathcal{F}(B_1)$ and $\mathcal{F}(B_2)$, respectively. Then

$$(A) \quad \frac{|\mathbb{E}\xi\eta - \mathbb{E}\xi\mathbb{E}\eta|}{(\mathbb{E}\xi^2)^{1/2}(\mathbb{E}\eta^2)^{1/2}} \leq e^{-\lambda q},$$

where $q = \text{dist}(B_1, B_2)$, $\lambda > 0$ is a fixed constant.

Lemma A.4. *Under condition (A), we have*

$$(A.9) \quad \mathbb{E}[\varphi_\rho; m]^2 \leq C \left(\frac{G^2}{\rho} \left(\frac{|\ln(1/\rho)|^2}{\rho^{1/2}m} \right)^d + e^{-c(\ln(1/\rho))^2} \right).$$

Proof. For $\ell = (\ell_1, \dots, \ell_d) \in \mathbb{Z}^d \cap [0, m]^d := \mathbb{Z}_m^d$, denote by I^ℓ the cube of size 1 centered at $\ell + \frac{1}{2} = (\ell_1 + \frac{1}{2}, \dots, \ell_d + \frac{1}{2})$, and let $\varphi^\ell = \int_{I^\ell} \varphi_\rho(x) dx$. Then

$$[\varphi_\rho; m] = \frac{1}{m^d} \sum_{\ell \in \mathbb{Z}_m^d} \varphi^\ell.$$

We first estimate $\mathbb{E}(\varphi^\ell \varphi^k)$. If $|\ell - k| \leq C\rho^{-1/2}|\ln(1/\rho)|^2$, then

$$\mathbb{E}(\varphi^\ell \varphi^k) \leq \langle \varphi_\rho^2(x) \rangle^{1/2} \langle \varphi_\rho^2(x) \rangle^{1/2} \leq C \frac{G^2}{\rho}.$$

If $|\ell - k| \geq C\rho^{-1/2}|\ln(1/\rho)|^2$, then let $B_1 := \ell + B(\rho^{-1/2}|\ln(1/\rho)|^2)$ and $B_2 := k + B(\rho^{-1/2}|\ln(1/\rho)|^2)$, where $B(s)$ is a ball of size s in the norm $\|\cdot\|$. Denote by $\tilde{\varphi}_1(x)$ the solution of (A.2) in which the coefficient $(a_{ij}(y, \omega))$ is modified in $\mathbb{R}^d \setminus B_1$ such that it is measurable with respect to $\mathcal{F}(B_1)$, and similarly denote by $\tilde{\varphi}_2(x)$ the solution of (A.2) in which the coefficient $(a_{ij}(y, \omega))$ is modified in $\mathbb{R}^d \setminus B_2$ such that it is measurable with respect to $\mathcal{F}(B_2)$. The modified coefficients $(\tilde{a}_{ij}(y, \omega))$ should still satisfy the condition on a_{ij} listed in the beginning of this subsection. From $\tilde{\varphi}_1$ and $\tilde{\varphi}_2$, we can similarly define $\tilde{\varphi}_1^\ell$ and $\tilde{\varphi}_2^k$. Using (A.6), we have

$$\begin{aligned} \mathbb{E}(\tilde{\varphi}_1^\ell - \tilde{\varphi}^\ell)^2 &\leq CG^2 e^{-C|\ln(1/\rho)|^2}, \\ \mathbb{E}(\tilde{\varphi}_1^k - \tilde{\varphi}^k)^2 &\leq CG^2 e^{-C|\ln(1/\rho)|^2}. \end{aligned}$$

Since

$$\begin{aligned} \mathbb{E}(\varphi^\ell \varphi^k) &= \mathbb{E}(\tilde{\varphi}^\ell \tilde{\varphi}^k) + \mathbb{E}(\varphi^\ell - \tilde{\varphi}_1^\ell) \tilde{\varphi}_2^k + \mathbb{E} \tilde{\varphi}_1^\ell (\varphi^k - \tilde{\varphi}_2^k) \\ &\quad + \mathbb{E}(\varphi^\ell - \tilde{\varphi}_1^\ell) (\varphi^k - \tilde{\varphi}_2^k) \end{aligned}$$

and

$$\begin{aligned} |\mathbb{E} \tilde{\varphi}_1^\ell|, |\mathbb{E} \tilde{\varphi}_2^k| &\leq CGe^{-C|\ln(1/\rho)|^2}, \\ |\mathbb{E} \tilde{\varphi}_1^\ell \tilde{\varphi}_2^k| &\leq C \left(\frac{G^2}{\rho} e^{-C|\ell-k|} + G^2 e^{-C|\ln(1/\rho)|^2} \right), \end{aligned}$$

we thus have

$$|\mathbb{E} \varphi^\ell \varphi^k| \leq C \left(\frac{G^2}{\rho} e^{-C|\ell-k|} + G^2 e^{-C|\ln(1/\rho)|^2} \right).$$

Hence

$$\begin{aligned}
\mathbb{E}[\varphi_\rho; m]^2 &= \frac{1}{m^{2d}} \sum_{\ell, k \in \mathbb{Z}_m^d} \mathbb{E} \varphi^\ell \varphi^k \\
&= \frac{1}{m^{2d}} \left(\sum_{|\ell-k| \leq \rho^{-1/2} |\ln(1/\rho)|^2} \mathbb{E} \varphi^\ell \varphi^k + \sum_{|\ell-k| > \rho^{-1/2} |\ln(1/\rho)|^2} \mathbb{E} \varphi^\ell \varphi^k \right) \\
&\leq \frac{1}{m^{2d}} \left(\frac{G^2}{\rho} m^d \left(\rho^{-1/2} |\ln(1/\rho)|^2 \right)^d + m^{2d} G^2 e^{-C |\ln(1/\rho)|^2} \right) \\
&\quad \times \frac{G^2}{\rho} \sum_{|\ell-k| \geq \rho^{-1/2} |\ln(1/\rho)|^2} e^{-C |\ell-k|} \\
&\leq C \frac{G^2}{\rho} \left(\frac{\ln(1/\rho)}{\rho^{1/2} m} \right)^d + e^{-C |\ln(1/\rho)|^2}.
\end{aligned}$$

□

Proceeding along the same line as in [43, Theorem 2.1], using condition **(A)**, we have

Lemma A.5. *For any $0 < \gamma < 1/2$, under condition **(A)**, there exists a constant C such that*

$$(A.10) \quad |\mathcal{A} - \langle a(I + \nabla \chi_\rho) \rangle| \leq C \rho^{\frac{d-2-2\gamma}{4+d}},$$

where $\chi_\rho = \{\chi_{k,\rho}\}_{k=1}^d$, and $\chi_{k,\rho}$ is the solution of (A.1) with $g = (a_{k1}, \dots, a_{kd})$.

Now we are ready to estimate $e(\text{HMM})$. Define $m = \frac{\delta}{2\varepsilon}$ and denote by φ_j^m the solution of (A.1) on $I_m = [0, m]^d$ with the boundary condition that $\varphi_j^m(y) = 0$ on ∂I_m , and let $\varphi^m = (\varphi_1^m, \dots, \varphi_d^m)$. Then

$$e(\text{HMM}) = |\mathcal{A} - [a(I + \nabla \varphi^m)]_m|.$$

Define $\varphi_\rho, \varphi_\rho^m$ similarly. We have

$$e(\text{HMM}) \leq E_1 + E_2 + E_3,$$

where

$$\begin{aligned}
E_1 &= |\mathcal{A} - [a(I + \nabla \varphi_\rho)]_m|, \\
E_2 &= |[a(I + \nabla \varphi_\rho) - a(I + \nabla \varphi_\rho^m)]_m|, \\
E_3 &= |[a \nabla (\varphi_\rho^m - \varphi^m)]_m|.
\end{aligned}$$

Obviously,

$$E_1 = |\mathcal{A} - \langle a(I + \nabla \varphi_\rho) \rangle + [\tilde{\psi}]_m|,$$

with $\tilde{\psi} = \langle a(I + \nabla \varphi_\rho) \rangle - a(I + \nabla \varphi_\rho)$. It was proved in [43, Lemma 2.5] that

$$\mathbb{E} |[\tilde{\psi}]_m| \leq (\mathbb{E} [\tilde{\psi}]_m^2)^{1/2} \leq C \left(\frac{|\ln(1/\rho)|^2}{\rho^{1/2} m} \right)^{d/2}.$$

The above inequality together with Lemma A.5 gives

$$\mathbb{E} E_1 \leq C \left(G \rho^{\frac{d-2-2\gamma}{4+d}} + \left(\frac{|\ln(1/\rho)|^2}{\rho^{1/2} m} \right)^{d/2} \right).$$

To estimate E_2 , denote by τ_m the first exit time for the domain I_{2m} . Then $\varphi_\rho^{2m} = M_x \Gamma(\tau_m)$. For any $s > 0$,

$$\begin{aligned} |\varphi_\rho - \varphi_\rho^{2m}| &= |M_x(\Gamma(\infty) - \Gamma(\tau_m))| \\ &\leq M_x\{|\Gamma(\infty)| + |\Gamma(\tau_m)|; \tau_m \leq s\} \\ &\quad + M_x\{e^{-s\rho} M_{\rho(s)}|\Gamma(\infty) - \Gamma(\tau_m)|; \tau_m > s\} \\ &\leq C\left(M_x((\Gamma(\infty))^2 + \Gamma(\tau_m)^2)\right)^{1/2} \{P_x(\tau_m \leq s)\}^{1/2} + e^{-s\rho}. \end{aligned}$$

Since

$$P_x\{\tau_m \leq s\} \leq e^{-Cm^2/s},$$

we get

$$\mathbb{E}[|\varphi_\rho - \varphi_\rho^{2m}|^2]_{2m} \leq C \frac{G^2}{\rho} \left(e^{-Cm^2/s} + e^{-s\rho}\right)^2.$$

Optimizing in s , we get

$$\mathbb{E}[|\varphi_\rho - \varphi_\rho^{2m}|^2]_{2m} \leq C \frac{G^4}{\rho^2} e^{-Cm\rho^{1/2}}.$$

Using standard interior estimates, we have

$$\mathbb{E} E_2 \leq \left(\mathbb{E}[|\nabla(\varphi_\rho - \varphi_\rho^m)|^2]_m\right)^{1/2} \leq \frac{C}{m} \left(\mathbb{E}[|\varphi_\rho - \varphi_\rho^{2m}|^2]_{2m}\right)^{1/2} \leq \frac{CG^2}{m\rho} e^{-Cm\rho^{1/2}}.$$

As for E_3 , proceeding along the same line as in the estimate of E_1 , we get

$$\mathbb{E} E_3 \leq C \left(G^2 \rho^{\frac{d-2-2\gamma}{4+d}} + \left(\frac{|\ln(1/\rho)|^2}{\rho^{1/2}m}\right)^{d/2}\right).$$

To sum up, we have

$$\mathbb{E} e(\text{HMM}) \leq C \left(G^2 \rho^{\frac{d-2-2\gamma}{4+d}} + \left(\frac{|\ln(1/\rho)|^2}{\rho^{1/2}m}\right)^{d/2} + \frac{CG^2}{m\rho} e^{-Cm\rho^{1/2}}\right).$$

Optimizing in ρ with respect to the first two terms, we get $\rho_0 = m^{-\frac{2d}{d+4\alpha}}$ with $\alpha = (d - 2 - 2\gamma)/(d + 4)$. Hence

$$(A.11) \quad \mathbb{E} e(\text{HMM}) \leq C \left(\frac{|\ln m|^d}{m^\kappa} + \frac{CG^2}{m\rho_0} e^{-Cm\rho_0^{1/2}}\right) \leq C \frac{|\ln m|^d}{m^\kappa},$$

with

$$\kappa = \frac{d/2}{1 + \frac{d(d+4)/4}{d-2-2\gamma}}, \quad 0 < \gamma < \frac{1}{2}.$$

Obviously, the $|\ln m|^d$ factor in (A.11) can be absorbed into the factor $m^{-\kappa}$. This proves Theorem 1.3 for $d = 3$.

Remark A.6. The estimate (A.11) is unlikely to be optimal. If $d = 1$, a direct calculation gives

$$\mathbb{E} e(\text{HMM}) \leq C \left(\frac{\varepsilon}{\delta}\right)^{1/2},$$

whereas the estimate (A.11) does not apply for $d = 1$ and $d = 2$. We may use the techniques in [13] to derive improved bounds for $e(\text{HMM})$ if the magnitude of the oscillation in the coefficients (a_{ij}^ε) is sufficiently small.

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DEPARTMENT OF MATHEMATICS AND PACM, PRINCETON UNIVERSITY, PRINCETON, NEW JERSEY 08544 AND SCHOOL OF MATHEMATICAL SCIENCES, PEKING UNIVERSITY, BEIJING 100871, PEOPLE'S REPUBLIC OF CHINA

E-mail address: `weinan@Princeton.EDU`

NO. 55, ZHONG-GUAN-CUN EAST ROAD, INSTITUTE OF COMPUTATIONAL MATHEMATICS, ACADEMY OF MATHEMATICS AND SYSTEMS SCIENCE, CHINESE ACADEMY OF SCIENCES, BEIJING 100080, PEOPLE'S REPUBLIC OF CHINA

E-mail address: `mpb@lsec.cc.ac.cn`

SCHOOL OF MATHEMATICAL SCIENCES, PEKING UNIVERSITY, BEIJING 100871, PEOPLE'S REPUBLIC OF CHINA

E-mail address: `pzhang@pku.edu.cn`