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ANGULAR PROCESSES RELATED TO CAUCHY RANDOM WALKS

В статье изучается угловой процесс, связанный со случайным блужданиями в евклидовом и неевклидовом пространстве, шаги которых имеют распределение Коши.

Это приводит к различного рода нелинейным преобразованиям случайных величин с распределением Коши, сохраняющим плотность Коши. Мы приводим явный вид этих распределений для всевозможных комбинаций параметров масштаба и сдвига.

Анализируются непрерывные дроби, содержащие случайные величины с распределением Коши. Показано, что на *n*-м этапе случайные величины все еще имеют распределение Коши с параметрами, связанными с числами Фибоначчи. Это позволяет нам доказать, что последовательность сходится по распределению к золотому сечению.

Ключевые слова и фразы: гиперболическая тригонометрия, закон арксинуса, непрерывные дроби, числа Фибоначчи, нелинейные преобразования случайных величин.

1. Introduction. We consider a particle starting from the origin O of \mathbf{R}^2 which takes initially a horizontal step of length 1 and a vertical one, say \mathbf{C}_1 , with a standard Cauchy distribution. It reaches therefore the position $(1, \mathbf{C}_1)$. The line l_1 joining the origin with $(1, \mathbf{C}_1)$ forms a random angle Θ_1 with the horizontal axis (see Figure 1).

On l_1 the traveller repeats the same movement with a step of unit length (either forward or backward) along l_1 and a standard Cauchy distributed step, say \mathbf{C}_2 , on the line orthogonal to l_1 . The right triangle obtained with the last two displacements has a hypotenuse belonging to the line l_2 with random inclination Θ_2 on l_1 .

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Fig. 1. The angular process in the Euclidean plane. By $C_{(a_j,b_j)}$ we indicate the *j*-th random displacement with Cauchy distribution possessing scale parameter a_j and location parameter b_j

After *n* steps the sequence of random angles $\Theta_1, \ldots, \Theta_n$ describes the rotation of the moving particle around the starting point, their partial sums describe an angular random walk which can be written as

$$S_n = \Theta_1 + \dots + \Theta_n = \sum_{j=1}^n \operatorname{arctg} \mathbf{C}_j,$$
 (1.1)

where \mathbf{C}_j are independent standard Cauchy random variables. If the random steps of the planar random walk above were independent Cauchy random variables with scale parameter a_j and location parameter b_j , then the process (1.1) must be a little bit modified and rewritten as

$$S_n = \Theta_1 + \dots + \Theta_n = \sum_{j=1}^n \operatorname{arctg} C_j, \qquad (1.2)$$

where $C_j \sim C_{(a_j,b_j)}$. The model (1.2) can be extended also to the case where the first step has length d_j and the second one is Cauchy distributed with scale parameter a_j and location parameter b_j (see Figure 1), then

$$\operatorname{tg} \Theta_j = C_{(a_j/d_j, b_j/d_j)}.$$

The same random walk can be generated if the two orthogonal steps, at each displacement, are represented by two independent Gaussian random variables X_j and Y_j . In this case, for each right triangle, we can write

$$\operatorname{tg} \Theta_j = \frac{X_j}{Y_j}.$$

If X_j and Y_j are two standard independent Gaussian random variables then tg $\Theta_j = X_j/Y_j$ possesses standard Cauchy distribution and we get the model in (1.1). The model (1.2) can be obtained by considering orthogonal Gaussian steps with different variances and in this case the scale parameter of the random variables C_j is the ratio $a_j = \sigma_j^Y/\sigma_j^X$. The model (1.1) describing the angular random process has a hyperbolic counterpart. We consider a particle starting from the origin O of the Poincaré half-plane $\mathbf{H}_{+}^{2} = \{(x, y): y > 0\}$. At the *j*-th displacement, $j = 1, 2, \ldots$, the particle makes two steps of random hyperbolic length η_{j} and $\hat{\eta}_{j}$ on two orthogonal geodesic lines. The *j*-th displacement leads to a right triangle T_{j} with sides of length η_{j} and $\hat{\eta}_{j}$ and random acute angles Θ_{j} and $\hat{\Theta}_{j}$. In each triangle T_{j} the first step is taken on the extension of the hypotenuse of the triangle T_{j-1} (see Figure 2). From hyperbolic trigonometry (for basic results on hyperbolic geometry see, for example, [5]) we have that

$$\sin \Theta_j = \frac{\operatorname{sh} \widehat{\eta}_j}{\sqrt{\operatorname{ch}^2 \eta_j \operatorname{ch}^2 \widehat{\eta}_j - 1}}, \quad \cos \Theta_j = \frac{\operatorname{sh} \eta_j \operatorname{ch} \widehat{\eta}_j}{\sqrt{\operatorname{ch}^2 \eta_j \operatorname{ch}^2 \widehat{\eta}_j - 1}}.$$

From the above expressions we have that

$$\operatorname{tg}\Theta_j = \frac{\operatorname{th}\,\widehat{\eta}_j}{\operatorname{sh}\eta_j}$$

If we take independent random hyperbolic displacements η_j and $\hat{\eta}_j$ such that the random variables $E_j = \frac{\operatorname{th} \hat{\eta}_j}{\operatorname{sh} \eta_j}$ are standard Cauchy distributed, then $\Theta_j = \operatorname{arctg} \mathbf{C}_j$. If the triangles T_j were isosceles, then $\operatorname{tg} \Theta_j = 1/\operatorname{ch} \eta_j$ and $\Theta_j \in [-\pi/4, \pi/4]$ so that in this case the Cauchy distribution cannot be attributed to $\operatorname{tg} \Theta_j$.



Fig. 2. The angular random process in the Poincaré half-plane

In the model described here the random steps (and therefore the random angular windings Θ_j) are independent. If we consider the model of papers [2] and [3], where the displacements are taken orthogonally to the geodesic lines joining the origin O of \mathbf{H}^2_+ with the positions occupied at deviation instants, the angular displacements Θ_j must be such that

$$\sin \Theta_j = \frac{\operatorname{sh} \eta_j}{\sqrt{1 + \prod_{r=1}^j \operatorname{ch}^2 \eta_r}} = \operatorname{sh} \eta_j \cos \left(\operatorname{arctg} \prod_{r=1}^j \operatorname{ch}^2 \eta_r \right)$$

and therefore involve dependent random variables.

For the area of the random hyperbolic triangle T_j we note that

$$\operatorname{area}\left(T_{j}\right) = \frac{\pi}{2} - \Theta_{j} - \widehat{\Theta}_{j} = \frac{\pi}{2} - \left[\operatorname{arctg}\left(\frac{\operatorname{th}\,\widehat{\eta}_{j}}{\operatorname{sh}\,\eta_{j}}\right) + \operatorname{arctg}\left(\frac{\operatorname{th}\,\eta_{j}}{\operatorname{sh}\,\widehat{\eta}_{j}}\right)\right]$$
$$= \frac{\pi}{2} - \operatorname{arctg}\left(\frac{\operatorname{coth}\,\eta_{j}}{\operatorname{sh}\,\widehat{\eta}_{j}} + \frac{\operatorname{coth}\,\widehat{\eta}_{j}}{\operatorname{sh}\,\eta_{j}}\right) = \operatorname{arcctg}\left(\frac{\operatorname{coth}\,\eta_{j}}{\operatorname{sh}\,\widehat{\eta}_{j}} + \frac{\operatorname{coth}\,\widehat{\eta}_{j}}{\operatorname{sh}\,\eta_{j}}\right).$$

Since each acute angle inside T_j is linked to both sides of the triangle, the analysis of the random process $A_n = \sum_{j=1}^n \operatorname{area}(T_j)$ is much more complicated and we drop it.

Let $C_j \sim C_{(a_j,b_j)}$, j = 1, 2, ..., be independent Cauchy random variables, where a_j is the scale parameter and b_j is the location parameter. In the study of the angular random walk (1.1) and (1.2) we must analyze the distribution of the following nonlinear transformations of Cauchy random variables:

$$U = \frac{C_1 + C_2}{1 - C_1 C_2},\tag{1.3}$$

since

$$\operatorname{arctg} C_1 + \operatorname{arctg} C_2 = \operatorname{arctg} \frac{C_1 + C_2}{1 - C_1 C_2}.$$

Since the 1960 a wide class of nonlinear transformations of Cauchy random variables has been considered. Williams [10], Knight [6] and Letac [7] proved that transformations of the form

$$\varepsilon f(x) = kx + \alpha - \sum_{i=1}^{n} \frac{\mu_i}{x - \gamma_i},$$

where $\varepsilon = \pm 1$, $\alpha, \gamma_i \in \mathbf{R}$ and $k, \mu_i \ge 0$, preserve the Cauchy distribution.

In particular, in [10] the following characterization for Cauchy random variables is proved. The random variable X is a standard Cauchy if and only if (1 + bX)/(b - X) is a standard Cauchy for some constant b which is not the tangent of a rational multiple of π .

Knight [6] asserts that a random variable X is of Cauchy type if and only if the random variable (aX + b)/(cX + d) is still of Cauchy type, whenever $ad - bc \neq 0$.

Our problem is more strictly related to the results obtained by Pitman and Williams [8]. They proved that the standard Cauchy distribution is preserved under certain types of transformations represented by meromorphic functions whose poles are all real and simple. As a corollary they obtained that, if P and Q are two independent random variables uniformly distributed in $(0, \pi)$, then the random variables $X = \operatorname{tg} P$ and $Y = \operatorname{tg} Q$ are standard Cauchy and

$$\operatorname{tg}(P+Q) = \frac{X+Y}{1-XY} \stackrel{\operatorname{law}}{=} X.$$

We will show that the random variable (1.3) is endowed with Cauchy distribution in a much more general situation, namely when the random variables C_j , j = 1, 2, have nonzero location parameters b_j and scale parameters a_j . The scale and location parameters of U depend on both parameters a_j and b_j suitably combined.

In particular, if $b_1 = b_2 = 0$ and $a_1 = a_2 = 1$, then U is still distributed as a standard Cauchy distribution and therefore in (1.1) we have that

$$S_n \stackrel{\text{law}}{=} \operatorname{arctg} \mathbf{C}.$$

Since also $1/\mathbb{C}$ is a standard Cauchy (for basic properties of Cauchy random variables see, for example, [9, p. 105], from (1.3), a number of related random variables preserving the form of the Cauchy distribution can be considered. For example, the random variables

$$Z_1 = rac{\mathbf{C}_1 \mathbf{C}_2 + 1}{\mathbf{C}_1 - \mathbf{C}_2}, \quad Z_2 = rac{1 - \mathbf{C}_1 \mathbf{C}_2}{\mathbf{C}_1 + \mathbf{C}_2}, \quad Z_3 = rac{\mathbf{C}_1 + \mathbf{C}_2}{\mathbf{C}_1 \mathbf{C}_2 - 1},$$

also possess standard Cauchy distribution.

If $b_1 = b_2 = 0$ and the scale parameters a_1, a_2 are different, then (1.3) still preserves the Cauchy distribution but with scale parameter equal to $(a_1 + a_2)/(1 + a_1a_2)$ and location parameter equal to zero. This can be grasped by means of the following relationship:

$$\operatorname{arctg} C_1 + \operatorname{arctg} C_2 \stackrel{\text{law}}{=} \operatorname{arctg} \left\{ \frac{a_1 + a_2}{1 + a_1 a_2} \mathbf{C} \right\},$$
 (1.4)

where $C_j \sim C_{(a_j,0)}$. Result (1.4) is illustrated in Figure 3.



Fig. 3. The figure shows that shooting a ray with inclination Θ_1 , uniformly distributed, against the line at distance a_1 and then shooting a ray with a uniformly distributed angle Θ_2 on the line at distance a_2 is equivalent to shooting on the barrier at the distance $a = (a_1 + a_2)/(1 + a_1a_2)$ with a uniformly distributed angle Θ

By iterating the process (1.4) we arrive at the formula

$$\sum_{j=1}^{3} \operatorname{arctg} C_j \stackrel{law}{=} \operatorname{arctg} \left\{ \frac{\sum_{j=1}^{3} a_j + a_1 a_2 a_3}{1 + \sum_{i \neq j} a_i a_j} \mathbf{C} \right\}$$

which gives an insight into further extensions of the process outlined above.

Much more complicated are the cases, where the location parameters of the Cauchy distributions are different from zero. For the special case, where C_1 and C_2 are independent Cauchy such that $C_1 \sim C_{(1,b)}$ and $C_2 \sim C_{(1,b)}$, the random variable (1.3) still possesses Cauchy density with scale parameter $(2b^2 + 4)/(b^4 + 4)$ and location parameter $2b^3/(b^4 + 4)$.

We have obtained the general distribution of (1.3), where C_1 and C_2 are independent Cauchy such that $C_1 \sim C_{(a_1,b_1)}$ and $C_2 \sim C_{(a_2,b_2)}$, and also the distribution of

$$U = \frac{\gamma C_1 + \delta C_2}{\alpha - \beta C_1 C_2}$$

for arbitrary positive real numbers β , γ , δ and $\alpha + \beta \neq 0$. In particular, if C_1 and C_2 are independent standard Cauchy random variables, then U is Cauchy with scale parameter equal to $|(\gamma + \delta)/(\alpha + \beta)|$ and location parameter equal to zero.

In the last section we examine continued fractions involving Cauchy random variables. In particular, we have studied

$$V_n = \frac{1}{1 + \frac{1}{1 + \frac{1}{1 + \mathbf{C}}}} \tag{1.5}$$

and

$$U_n = \frac{1}{1 + \frac{1}{1 + \frac{1}{1 + \frac{1}{1 + C^2}}}} \tag{1.6}$$

which generalize the random variables $V_1 = 1/(1 + \mathbf{C})$ and $U_1 = 1/(1 + \mathbf{C}^2)$. Continued fractions involving random variables have been analyzed in [4] and more recently in [1]. The random variable U_1 has the arcsine distribution in [0, 1], while $U_t = tU_1$, with t > 0, has distribution

$$\mathbf{P}\{U_t \in \mathrm{d}s\} = \frac{\mathrm{d}s}{\pi\sqrt{s(t-s)}}, \qquad 0 < s < t.$$

For each $n \ge 1$, the random variables V_n are Cauchy distributed with scale parameter a_n and location parameter b_n that can be expressed in terms of Fibonacci numbers. This permits us to prove the monotonicity of a_n and b_n and that $\lim_{n\to\infty} a_n = 0$ and $\lim_{n\to\infty} b_n = \varphi - 1$, where $\varphi = (1 + \sqrt{5})/2$ is the golden ratio. Finally we obtain that the sequences of random variables $1 + V_n$ and $1 + U_n$, $n \ge 1$, converge in distribution to the number $\varphi = (1 + \sqrt{5})/2$. This should be expected, since it has the infinite fractional expansion

$$\frac{1+\sqrt{5}}{2} = 1 + \frac{1}{1+\frac{1}{1+\dots}} \tag{1.7}$$

which is related to (1.5) and (1.6).

2. Centered Cauchy random variables. In this section we study the distribution of the following random variable:

$$U = \frac{\gamma \mathbf{C}_1 + \delta \mathbf{C}_2}{\alpha - \beta \mathbf{C}_1 \mathbf{C}_2},\tag{2.1}$$

where \mathbf{C}_1 and \mathbf{C}_2 are independent, standard Cauchy. We assume, without restrictions, that β , γ , δ are nonnegative real numbers all different from zero (because of the symmetry of \mathbf{C}_j , j = 1, 2) and $\alpha + \beta \neq 0$. In the following theorem we prove that U is still Cauchy distributed.

Theorem 2.1. The random variable U in (2.1) possesses Cauchy distribution with scale parameter equal to $|(\gamma + \delta)/(\alpha + \beta)|$ and location parameter equal to zero. We can also restate the result in symbols as

$$U \stackrel{\text{law}}{=} \frac{\gamma + \delta}{\alpha + \beta} \mathbf{C}.$$

Proof. The density of the random variable U can be obtained by means of the transformation

$$u = rac{\gamma x + \delta y}{lpha - eta x y}, \quad v = y,$$

with $(x, y) \in \mathbf{R}^2$ and Jacobian equal to

$$|J| = \frac{\alpha \gamma + \delta \beta v^2}{(\gamma + \beta u v)^2}$$

The joint density g = g(u, v) reads

$$g(u,v) = \frac{1}{\pi^{2}(1+v^{2})} \frac{1}{1 + ((\alpha u - \delta v)/(\gamma + \beta uv))^{2}} \frac{\alpha \gamma + \delta \beta v^{2}}{(\gamma + \beta uv)^{2}} \\ = \frac{1}{\pi^{2}(1+v^{2})} \frac{\alpha \gamma + \delta \beta v^{2}}{v^{2}(\beta^{2}u^{2} + \delta^{2}) + 2uv(\beta \gamma - \delta \alpha) + \alpha^{2}u^{2} + \gamma^{2}}$$

where $(u, v) \in \mathbf{R}^2$, and must be integrated with respect to v in order to obtain the distribution of U. Therefore

$$\mathbf{P}\left\{\frac{\gamma \mathbf{C}_{1} + \delta \mathbf{C}_{2}}{\alpha - \beta \mathbf{C}_{1} \mathbf{C}_{2}} \in \mathrm{d}u\right\} \\
= \frac{\mathrm{d}u}{\pi^{2}} \int_{-\infty}^{\infty} \left[\frac{1}{1 + v^{2}} \frac{\alpha \gamma + \delta \beta v^{2}}{v^{2}(u^{2}\beta^{2} + \delta^{2}) + 2vu(\beta \gamma - \delta \alpha) + u^{2}\alpha^{2} + \gamma^{2}}\right] \mathrm{d}v \\
= \frac{\mathrm{d}u}{\pi^{2}} \int_{-\infty}^{\infty} \left[\frac{Av + B}{1 + v^{2}} + \frac{Cv + D}{v^{2}(u^{2}\beta^{2} + \delta^{2}) + 2vu(\beta \gamma - \delta \alpha) + u^{2}\alpha^{2} + \gamma^{2}}\right] \mathrm{d}v, \quad (2.2)$$

where

$$\begin{split} A &= \frac{2u(\beta\gamma - \alpha\delta)(\beta\delta - \alpha\gamma)}{[u^2(\alpha - \beta)^2 + (\gamma - \delta)^2][u^2(\alpha + \beta)^2 + (\gamma + \delta)^2]},\\ B &= \frac{(\alpha\gamma - \beta\delta)[(\gamma^2 - \delta^2) + u^2(\alpha^2 - \beta^2)]}{[u^2(\alpha - \beta)^2 + (\gamma - \delta)^2][u^2(\alpha + \beta)^2 + (\gamma + \delta)^2]},\\ C &= -A(u^2\beta^2 + \delta^2),\\ D &= \frac{(\beta\gamma - \alpha\delta)[u^4\alpha\beta(\beta^2 - \alpha^2) + u^2(\beta\gamma - \alpha\delta)(3\alpha\gamma - \beta\delta) + \gamma\delta(\gamma^2 - \delta^2)]}{[u^2(\alpha - \beta)^2 + (\gamma - \delta)^2][u^2(\alpha + \beta)^2 + (\gamma + \delta)^2]}.\end{split}$$

We start by evaluating the first part of the integral (2.2):

$$\begin{split} &\int_{-\infty}^{\infty} \left[\frac{Av}{1+v^2} + \frac{Cv}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \right] \mathrm{d}v \\ &= \frac{A}{2} \int_{-\infty}^{\infty} \left[\frac{2v}{1+v^2} - \frac{2v(u^2\beta^2+\delta^2)\pm 2u(\beta\gamma-\delta\alpha)}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \right] \mathrm{d}v \\ &= \frac{A}{2} \lim_{d\to\infty,c\to-\infty} \log\left(\frac{1+v^2}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \right) \Big|_{c}^{d} \\ &+ Au(\beta\gamma-\delta\alpha) \int_{-\infty}^{\infty} \frac{1}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \mathrm{d}v \\ &= Au(\beta\gamma-\delta\alpha) \int_{-\infty}^{\infty} \frac{1}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \mathrm{d}v \\ &= Au(\beta\gamma-\delta\alpha) \frac{\pi}{u^2\alpha\beta+\gamma\delta}, \end{split}$$
(2.3)

where the last integral is obtained by means of the change of variable

$$v\sqrt{u^2\beta^2+\delta^2}+\frac{u(\beta\gamma-\delta\alpha)}{\sqrt{u^2\beta^2+\delta^2}}=z\sqrt{u^2\alpha^2+\gamma^2-\frac{u^2(\beta\gamma-\delta\alpha)^2}{u^2\beta^2+\delta^2}}=z\frac{u^2\alpha\beta+\gamma\delta}{\sqrt{u^2\beta^2+\delta^2}}.$$

In view of result (2.3) and inserting the values of A, B and D we have that

$$\begin{split} \frac{\mathrm{d}u}{\pi^2} & \int_{-\infty}^{\infty} \left[\frac{Av+B}{1+v^2} + \frac{Cv+D}{v^2(u^2\beta^2+\delta^2)+2vu(\beta\gamma-\delta\alpha)+u^2\alpha^2+\gamma^2} \right] \mathrm{d}v \\ &= \frac{\mathrm{d}u}{\pi} \frac{1}{u^2\alpha\beta+\gamma\delta} \left[Au(\beta\gamma-\alpha\delta) + B(u^2\alpha\beta+\gamma\delta) + D \right] \\ &= \frac{\mathrm{d}u}{\pi} \left[\frac{(\gamma\delta+u^2\alpha\beta)(\beta\gamma-\alpha\delta)[u^2(\beta^2-\alpha^2)+(\gamma^2-\delta^2)]}{(u^2\alpha\beta+\gamma\delta)[u^2(\alpha-\beta)^2+(\gamma-\delta)^2][u^2(\alpha+\beta)^2+(\gamma+\delta)^2]} \right] \\ &+ \frac{(\gamma\delta+u^2\alpha\beta)(\alpha\gamma-\beta\delta)[u^2(\alpha^2-\beta^2)+(\gamma^2-\delta^2)]}{(u^2\alpha\beta+\gamma\delta)[u^2(\alpha-\beta)^2+(\gamma-\delta)^2][u^2(\alpha+\beta)^2+(\gamma+\delta)^2]} \\ &= \frac{\mathrm{d}u}{\pi} \frac{u^2(\beta^2-\alpha^2)(\beta-\alpha)(\gamma+\delta)+(\gamma^2-\delta^2)(\alpha+\beta)(\gamma-\delta)}{[u^2(\alpha-\beta)^2+(\gamma-\delta)^2][u^2(\alpha+\beta)^2+(\gamma+\delta)^2]} \\ &= \frac{\mathrm{d}u}{\pi} \frac{(\alpha+\beta)(\gamma+\delta)[u^2(\alpha-\beta)^2+(\gamma-\delta)^2]}{[u^2(\alpha-\beta)^2+(\gamma-\delta)^2][u^2(\alpha+\beta)^2+(\gamma+\delta)^2]} \end{split}$$

$$= \frac{\mathrm{d}u}{\pi} \frac{(\alpha + \beta)(\gamma + \delta)}{u^2(\alpha + \beta)^2 + (\gamma + \delta)^2}$$

Another approach is based on the conditional characteristic function,

$$\begin{split} \mathbf{E} \left\{ e^{i\beta U} \mid \mathbf{C}_2 = v \right\} &= \frac{1}{\pi} \int_{-\infty}^{\infty} e^{i\beta(\gamma u + \delta v)/(\alpha - \beta u v)} \frac{\mathrm{d}u}{1 + u^2} \\ &= \frac{1}{\pi} \int_{-\infty}^{\infty} e^{i\beta w} \frac{\alpha \gamma + \delta \beta v^2}{v^2(\beta^2 w^2 + \delta^2) + 2vw(\beta \gamma - \delta \alpha) + w^2 \alpha^2 + \gamma^2} \, \mathrm{d}w \end{split}$$

where $w = (\gamma u + \delta v)/(\alpha - \beta u v)$. The inverse Fourier transform gives the conditional density g(w|v) and thus we arrive again at the integral (2.2). Theorem 2.1 is proved.

R e m a r k 2.1. A special case implied by Theorem 2.1 concerns the random variable $G \to G$

$$V = \frac{C_1 + C_2}{1 - C_1 C_2},$$

where $C_j \sim C_{(a_j,0)}$, j = 1, 2, are independent Cauchy random variables with location parameter equal to zero and scale parameter $a_j > 0$. If we choose $\gamma = a_1, \, \delta = a_2, \, \alpha = 1$, and $\beta = a_1 a_2$, we conclude that

$$V \stackrel{\text{law}}{=} \frac{a_1 \mathbf{C}_1 + a_2 \mathbf{C}_1}{1 - a_1 a_2 \mathbf{C}_1 \mathbf{C}_2} \stackrel{\text{law}}{=} \frac{a_1 + a_2}{1 + a_1 a_2} \mathbf{C}.$$

In the same way we obtain that

$$\widehat{V} = \frac{\gamma C_1 + \delta C_2}{\alpha - \beta C_1 C_2} \stackrel{\text{law}}{=} \frac{\gamma a_1 \mathbf{C}_1 + \delta a_2 \mathbf{C}_2}{\alpha - \beta a_1 a_2 \mathbf{C}_1 \mathbf{C}_2} \stackrel{\text{law}}{=} \frac{\gamma a_1 + \delta a_2}{\alpha + \beta a_1 a_2} \mathbf{C},$$

where $\alpha + \beta a_1 a_2 \neq 0$.

R e m a r k 2.2. In view of Theorem 2.1 we can obtain by recurrence the distribution of the current angle after n steps for the angular random walk

$$S_n = \sum_{j=1}^n \operatorname{arctg} C_j$$

described in the introduction. We have

$$S_n = \sum_{j=1}^n \operatorname{arctg} C_{(a_j,0)} = \operatorname{arctg} C_{(\hat{a}_{n-1},0)} + \operatorname{arctg} C_{(a_n,0)},$$

where $\operatorname{arctg} C_{(\hat{a}_{n-1},0)}$ is the random variable S_{n-1} . In particular, if $a_j = 1$ for $j = 1, \ldots, n$, we have the following property of the standard Cauchy random variables \mathbf{C}_j :

$$\sum_{j=1}^{n} \operatorname{arctg} \mathbf{C}_{j} \stackrel{\text{law}}{=} \operatorname{arctg} \mathbf{C}.$$

R e m a r k 2.3. A simple byproduct of Theorem 2.1 is that

$$\mathbf{E} e^{i\beta U} = \frac{1}{\pi^2} \int_{\mathbf{R}^2} \exp\left\{i\beta \frac{x+y}{1-xy}\right\} \frac{a_1 a_2 \, \mathrm{d}x \, \mathrm{d}y}{(a_1^2+x^2)(a_2^2+y^2)} \\ = \frac{1}{\pi^2} \int_{-\pi/2}^{\pi/2} \int_{-\pi/2}^{\pi/2} \exp\left\{i\beta \frac{a_1 \, \mathrm{tg} \, \theta_1 + a_2 \, \mathrm{tg} \, \theta_2}{1-a_1 a_2 \, \mathrm{tg} \, \theta_1 \, \mathrm{tg} \, \theta_2}\right\} \mathrm{d}\theta_1 \, \mathrm{d}\theta_2 \\ = \exp\left\{-\frac{a_1 + a_2}{1+a_1 a_2} \, |\beta|\right\}.$$
(2.4)

In (2.4) we have used the transformations $x = a_1 \operatorname{tg} \theta_1$ and $y = a_2 \operatorname{tg} \theta_2$. In the special case $a_1 = a_2 = 1$ the relationship (2.4) yields

$$e^{-|\beta|} = \frac{1}{\pi^2} \int_{-\pi/2}^{\pi/2} \int_{-\pi/2}^{\pi/2} \exp\left\{i\beta \frac{\operatorname{tg}\theta_1 + \operatorname{tg}\theta_2}{1 - \operatorname{tg}\theta_1 \operatorname{tg}\theta_2}\right\} d\theta_1 d\theta_2$$

= $\frac{1}{\pi^2} \int_{-\pi/2}^{\pi/2} \int_{-\pi/2}^{\pi/2} e^{i\beta \operatorname{tg}(\theta_1 + \theta_2)} d\theta_1 d\theta_2 = \frac{2}{\pi^2} \int_0^{\pi} x \cos(\beta \operatorname{tg} x) dx.$ (2.5)

In the last step of (2.5) we have used the transformations $\theta_1 + \theta_2 = x$ and $\theta_2 = y$. The integral (2.5) shows that, if (Θ_1, Θ_2) is uniform in the square $S = \{(\theta_1, \theta_2): -\pi/2 < |\theta_i| < \pi/2, i = 1, 2\}$, then the random variable $W = \operatorname{tg}(\Theta_1 + \Theta_2)$ has characteristic function $e^{-|\beta|}$ because $\Theta_1 + \Theta_2$ is uniform and therefore W is Cauchy distributed.

R e m a r k 2.4. It is well known that for a planar Brownian motion $\{(B_1(t), B_2(t)), t > 0\}$ starting from (x, y) the random variable $B_1(T_y)$ possesses Cauchy distribution with parameters (x, y), where

$$T_y = \inf\{t > 0: B_2(t) = 0\}.$$

If the starting points of two planar Brownian motions $(B_1^i(t), B_2^i(t))$, for i = 1, 2, are located on the y axis as in Figure 4, then we have that

$$\begin{split} \Theta &= \Theta_1 + \Theta_2 \; = \; \operatorname{arctg} B_1^1(T_{a_1}) + \operatorname{arctg} B_1^2(T_{a_2}) \\ &\stackrel{\text{law}}{=} \; \operatorname{arctg} C_1 + \operatorname{arctg} C_2 \stackrel{\text{law}}{=} \; \operatorname{arctg} \frac{a_1 + a_2}{1 + a_1 a_2} \mathbf{C}, \end{split}$$

where C_1 and C_2 are two independent Cauchy random variables with scale parameters a_1 and a_2 , respectively, and location parameters equal to zero. Therefore if the starting point of a third Brownian motion has coordinates $(0, (a_1 + a_2)/(1 + a_1a_2))$ then $B(T_{(a_1+a_2)/(1+a_1a_2)})$ represents its hitting position on the x-axis. This point forms with (0, 1) and the origin a right triangle with an angle equal to $\Theta = \Theta_1 + \Theta_2$.



Fig. 4. The hitting position on the x-axis of a planar Brownian motion is Cauchy distributed. In the figure the random angles Θ_1 , Θ_2 , and $\Theta = \Theta_1 + \Theta_2$ are shown. For $a = (a_1 + a_2)/(1 + a_1a_2)$ the right-hand figure shows the hitting position $B_1(T_a)$

3. Noncentered Cauchy random variables. For independent Cauchy random variables C_1 and C_2 , with location parameters b_1 and b_2 and scale parameters a_1 and a_2 , the random variable U in (2.1) is still Cauchy distributed with both parameters affected by the values of the location parameters b_1 , b_2 and the scale parameters a_1 , a_2 .

Theorem 3.1. If C_i , i = 1, 2, are two independent Cauchy random variables with location parameters b_i and scale parameters a_i , then the random variable U is still Cauchy distributed with scale parameter

$$a_U = \frac{(a_1 + a_2)(1 + a_1a_2 - b_1b_2) + (b_1 + b_2)(a_1b_2 + a_2b_1)}{(1 + a_1a_2 - b_1b_2)^2 + (a_1b_2 + a_2b_1)^2}$$

and location parameter

$$b_U = \frac{(a_1 + a_2)(a_1b_2 + a_2b_1) - (b_1 + b_2)(1 + a_1a_2 - b_1b_2)}{(1 + a_1a_2 - b_1b_2)^2 + (a_1b_2 + a_2b_1)^2}.$$

P r o o **f**. We obtain the density function of the random variable U by observing that

$$\mathbf{P}\{U \in \mathrm{d}w\} = \mathbf{E}\left\{\mathbf{P}\left\{\frac{C_1 + C_2}{1 - C_1 C_2} \in \mathrm{d}w \,\middle|\, C_2\right\}\right\}$$

and remarking that

$$\mathbf{P}\left\{\frac{C_1+C_2}{1-C_1C_2} \in \mathrm{d}w \,\middle|\, C_2=y\right\} = \frac{a_1}{\pi[(w-y)/(1+wy)-b_1]^2+a_1^2} \frac{1+y^2}{(1+wy)^2} \,\mathrm{d}w.$$

Therefore

$$\mathbf{P}\{U \in \mathrm{d}w\} = \frac{a_1 a_2 \,\mathrm{d}w}{\pi^2} \int_{-\infty}^{\infty} \frac{1}{[(w-y)/(1+wy)-b_1]^2 + a_1^2} \frac{1+y^2}{(1+wy)^2} \frac{1}{(y-b_2)^2 + a_2^2} \,\mathrm{d}y \\
= \frac{a_1 a_2 \,\mathrm{d}w}{\pi^2} \int_{-\infty}^{\infty} \frac{1+y^2}{[w-y-b_1(1+wy)]^2 + a_1^2(1+wy)^2} \frac{1}{(y-b_2)^2 + a_2^2} \,\mathrm{d}y \\
= \frac{a_1 a_2 \,\mathrm{d}w}{\pi^2} \int_{-\infty}^{\infty} \frac{1+y^2}{[y^2 E - 2yF + G][y^2 - 2yH + K]} \,\mathrm{d}y,$$
(3.1)

where

$$E = 1 + 2b_1w + (a_1^2 + b_1^2)w^2, \quad F = -b_1 - (a_1^2 + b_1^2 - 1)w + b_1w^2,$$

$$G = a_1^2 + b_1^2 - 2b_1w + w^2, \quad H = b_2, \quad K = a_2^2 + b_2^2.$$

We rewrite the integral in (3.1) in the following form:

$$\mathbf{P}\{U \in dw\} = \frac{a_1 a_2 dw}{\pi^2} \int_{-\infty}^{\infty} \left[\frac{Ay + B}{y^2 E - 2yF + G} + \frac{Cy + D}{y^2 - 2yH + K} \right] dy \quad (3.2)$$

with

$$A + CE = 0, \quad -2AH + B - 2CF + DE = 1,$$

$$AK - 2HB + CG - 2DF = 0, \quad BK + DG = 1.$$
(3.3)

The integrals in (3.2) can be worked out by means of the changes of variables

$$y\sqrt{E} - \frac{F}{\sqrt{E}} = x\sqrt{G - \frac{F^2}{E}}, \quad y - H = x\sqrt{K - H^2}.$$

The first integral becomes

$$\int_{-\infty}^{\infty} \frac{Ay+B}{(y\sqrt{E}-F/\sqrt{E})^2+G-F^2/E} \,\mathrm{d}y$$
$$= \int_{-\infty}^{\infty} \frac{1}{E} \frac{BE+AF+xA\sqrt{GE-F^2}}{(x^2+1)\sqrt{GE-F^2}} \,\mathrm{d}x,$$

and the second one takes the form

$$\int_{-\infty}^{\infty} \frac{Cy+D}{(y-H)^2+K-H^2} \, \mathrm{d}y = \int_{-\infty}^{\infty} \frac{D+CH+xC\sqrt{K-H^2}}{(x^2+1)\sqrt{K-H^2}} \, \mathrm{d}x.$$

A substantial simplification can be obtained because

$$\sqrt{GE - F^2} = a_1(1 + w^2), \qquad \sqrt{H - K^2} = a_2.$$

If we turn back to the distribution (3.2), in view of the above calculations, we have that

$$\mathbf{P}\{U \in dw\} = \frac{a_2 dw}{\pi^2 (1+w^2)E} \int_{-\infty}^{\infty} \frac{BE + AF + xAa_1(1+w^2)}{x^2 + 1} dx + \frac{a_1 dw}{\pi^2} \int_{-\infty}^{\infty} \frac{D + CH + xCa_2}{1+x^2} dx.$$

We observe that, by the first equation of (3.3),

$$\frac{a_1 a_2 \operatorname{d} w}{\pi^2} \left(\frac{A}{E} + C\right) \int_{-\infty}^{\infty} \frac{x}{1 + x^2} \operatorname{d} x = 0$$

and

$$\mathbf{P}\{U \in \mathrm{d}w\} = \left\{a_2 \frac{BE + AF}{\pi E(1+w^2)} + a_1 \frac{D + CH}{\pi}\right\} \mathrm{d}w$$
$$= \left\{a_2 \frac{B - CF}{\pi(1+w^2)} + a_1 \frac{D + CH}{\pi}\right\} \mathrm{d}w.$$

The values of B, C, and D can be derived by solving the system (3.3), simplified as

$$B + 2C(HE - F) + DE = 1, -2BH + C(G - EK) - 2DF = 0,$$

 $BK + DG = 1.$

By means of cumbersome computations we arrive at the final result

$$\begin{split} \mathbf{P} \{ U \in \mathrm{d}w \} \\ &= \frac{1}{\pi} \Big\{ (1 + a_1 a_2 - b_1 b_2) [w(a_1 b_2 + a_2 b_1) + a_1 + a_2] \\ &- (a_1 b_2 + a_2 b_1) [w(1 + a_1 a_2 - b_1 b_2) - (b_1 + b_2)] \Big\} \\ &\times \Big\{ [w(1 + a_1 a_2 - b_1 b_2) - (b_1 + b_2)]^2 + [w(a_1 b_2 + a_2 b_1) + a_1 + a_2]^2 \Big\}^{-1} \\ &= \frac{(a_1 + a_2)(1 + a_1 a_2 - b_1 b_2) + (b_1 + b_2)(a_1 b_2 + a_2 b_1)}{(1 + a_1 a_2 - b_1 b_2)^2 + (a_1 b_2 + a_2 b_1)^2} \\ &\times \Big\{ \Big[w + \frac{(a_1 + a_2)(a_1 b_2 + a_2 b_1) - (b_1 + b_2)(1 + a_1 a_2 - b_1 b_2)}{(1 + a_1 a_2 - b_1 b_2)^2 + (a_1 b_2 + a_2 b_1)^2} \Big]^2 \\ &+ \Big[\frac{(a_1 + a_2)(1 + a_1 a_2 - b_1 b_2) + (b_1 + b_2)(a_1 b_2 + a_2 b_1)}{(1 + a_1 a_2 - b_1 b_2)^2 + (a_1 b_2 + a_2 b_1)^2} \Big]^2 \Big\}^{-1}. \end{split}$$

Theorem 3.1 is proved.

R e m a r k 3.1. In view of Theorem 3.1 it is possible to obtain the following particular cases.

(i) For $a_1 = a_2 = 1$ and $b_1 = b_2 = b$, we have that

$$a_U = \frac{2b^2 + 4}{b^4 + 4}, \quad b_U = \frac{2b^3}{b^4 + 4}.$$

This shows that U has center of symmetry on the positive half-line if b > 0and on the negative half-line if b < 0, therefore the nonlinear transformation U preserves the sign of the mode.

(ii) For
$$a_1 = a_2 = a$$
 and $b_1 = b_2 = b$ we have that

$$a_U = rac{2a(1+a^2+b^2)}{(1+a^2-b^2)^2+(2ab)^2}, \quad b_U = rac{2b(a^2+b^2-1)}{(1+a^2-b^2)^2+(2ab)^2}.$$

We note that a_U and b_U depend simultaneously from the scale and location parameters of the random variables involved in U.

4. Continued fractions. The property that the reciprocal of a Cauchy random variable has still a Cauchy distribution has a number of possible extensions which we deal with in this section.

We start by considering the sequence

$$V_1 = \frac{1}{1+\mathbf{C}}, \quad V_2 = \frac{1}{1+\frac{1}{1+\mathbf{C}}}, \quad \dots, \quad V_n = \frac{1}{1+\frac{1}{1+\dots+\frac{1}{1+\mathbf{C}}}}, \quad (4.1)$$

and prove the following theorem.

Theorem 4.1. The random variables defined in (4.1) have Cauchy distribution, $V_n \sim C_{(a_n,b_n)}$, where the scale parameters a_n and the location parameters b_n satisfy the recursive relationships

$$a_{n+1} = \frac{a_n}{(1+b_n)^2 + a_n^2}, \qquad n = 1, 2, \dots,$$
 (4.2)

$$b_{n+1} = \frac{b_n + 1}{(1+b_n)^2 + a_n^2}, \qquad n = 1, 2, \dots$$
 (4.3)

P r o o f. Let us assume that V_n possesses Cauchy density with parameters a_n and b_n . Then V_{n+1} writes

$$\mathbf{P}\{V_{n+1} < v\} = \mathbf{P}\left\{\frac{1}{1+V_n} < v\right\} = \mathbf{P}\left\{\frac{1}{1+a_n+b_n\mathbf{C}} < v\right\}.$$

After some computations the density of V_{n+1} can be written as

$$f_{V_{n+1}}(v) = \frac{a_n/((1+b_n)^2 + a_n^2)}{\pi[v - (b_n + 1)/((1+b_n)^2 + a_n^2)]^2 + a_n^2/[(1+b_n)^2 + a_n^2]^2}, \qquad v \in \mathbf{R}.$$

It can be ascertained directly that V_1 possesses Cauchy distribution with parameters $a_1 = 1/2$ and $b_1 = 1/2$.

Theorem 4.1 is proved.

R e m a r k 4.1. We have evaluated the following table of parameters a_n and b_n :

n	1	2	3	 10^{2}
a_n	1/2	1/5	1/13	 $5.77e^{-42}$
b_n	1/2	3/5	8/13	 0.618034

For n = 1, 2, 3 we can observe that the scale parameters a_n coincide with the inverse of the odd-indexed Fibonacci numbers while the sequence b_n has the numerators coinciding with the even-indexed Fibonacci numbers and the denominators correspond to the odd-indexed Fibonacci numbers.

In light of the previous considerations we can show that for $n \ge 1$

$$b_n = \frac{F_{2n}}{F_{2n+1}}, \quad a_n = \frac{1}{F_{2n+1}},$$
(4.4)

where F_n , $n \ge 0$, is the Fibonacci sequence. Recalling that the Fibonacci numbers admit the representation (it can easily be checked by induction)

$$F_n = \frac{\varphi^n - (1 - \varphi)^n}{\sqrt{5}},\tag{4.5}$$

where $\varphi = (1 + \sqrt{5})/2$ is the golden ratio, we now prove that if a_n and b_n have the representations (4.4), then also a_{n+1} and b_{n+1} can be expressed in the same form. From (4.2) and (4.3) we have

$$\begin{split} b_{n+1} &= \frac{F_{2n}/F_{2n+1}+1}{(F_{2n}/F_{2n+1}+1)^2+1/F_{2n+1}^2} = \frac{F_{2n+2}F_{2n+1}}{F_{2n+2}^2+1} \\ &= F_{2n+2} \bigg[\frac{\varphi^{2n+1}-(1-\varphi)^{2n+1}}{\varphi^{4n+4}+(1-\varphi)^{4n+4}-2\varphi^{2n+2}(1-\varphi)^{2n+2}+5} \bigg] \sqrt{5} \\ &= F_{2n+2} \bigg[\frac{\varphi^{2n+1}-(1-\varphi)^{2n+1}}{[\varphi^{2n+1}-(1-\varphi)^{2n+1}][\varphi^{2n+3}-(1-\varphi)^{2n+3}]} \bigg] \sqrt{5} = \frac{F_{2n+2}}{F_{2n+3}}. \end{split}$$

Similar calculations prove that $a_{n+1} = \frac{1}{F_{2n+3}}$. In view of representations (4.4) and (4.5), it is easy to show that

$$\lim_{n \to \infty} b_n = \lim_{n \to \infty} \frac{F_{2n}}{F_{2n+1}} = \lim_{n \to \infty} \frac{1 - ((1-\varphi)/\varphi)^{2n}}{\varphi - (1-\varphi)((1-\varphi)/\varphi)^{2n}} = \frac{1}{\varphi} = \varphi - 1,$$
$$\lim_{n \to \infty} a_n = \lim_{n \to \infty} \frac{1}{F_{2n+1}} = 0.$$

Otherwise, observing that the sequence b_n , $n \ge 1$, is increasing, because

$$\frac{b_{n+1}}{b_n} = \frac{F_{2n+2}F_{2n+1}}{F_{2n+3}F_{2n}} = \frac{\varphi^{4n+3} + (1-\varphi)^{4n+3} + 1}{\varphi^{4n+3} + (1-\varphi)^{4n+3} - 4} \ge 1,$$

and taking the limits in (4.2) and (4.3) we have that

$$L = \frac{L}{(1+H)^2 + L^2}, \quad H = \frac{H+1}{(1+H)^2 + L^2}, \tag{4.6}$$

where $H = \lim_{n \to \infty} a_n$ and $L = \lim_{n \to \infty} b_n$. From the relationships (4.6) we derive the equality

$$\frac{L}{H} = \frac{L}{H+1}$$

that implies L = 0. In fact, for $L \neq 0$, we arrive at the absurd that H = H + 1. Substituting L = 0 in the second formula of (4.6) we obtain

$$H = \frac{1}{H+1}.$$

Since H satisfies the algebraic equation $H^2 + H - 1 = 0$, it follows that $H = \varphi - 1$, where φ is the golden ratio (see Figure 5).



Fig. 5. In the left figure the densities of the Cauchy random variables V_1 , V_2 , V_3 , and V_4 are shown. In the right figure the densities of U_1 , U_2 , U_3 and U_4 are plotted

R e m a r k $\,$ 4.2. A slightly more general case concerns the sequence

$$W_{1} = \frac{1}{c_{1} + d_{1}C_{(a_{0},b_{0})}} = \frac{1}{c_{1} + a_{0}d_{1} + b_{0}d_{1}\mathbf{C}},$$
$$W_{2} = \frac{1}{c_{2} + d_{2}W_{1}}, \quad W_{3} = \frac{1}{c_{3} + d_{3}W_{2}}, \dots$$

By performing calculations similar to those of Theorem 4.1 we have that W_1 has Cauchy distribution with scale parameter a_1 and location parameter b_1 such that

$$a_1 = \frac{d_1 a_0}{(c_1 + d_1 b_0)^2 + d_1^2 a_0^2}, \quad b_1 = \frac{c_1 + d_1 b_0}{(c_1 + d_1 b_0)^2 + d_1^2 a_0^2}$$

Similarly, if $W_n \sim C_{(a_n,b_n)}$, then $W_{n+1} \sim C_{(a_{n+1},b_{n+1})}$, where

$$a_{n+1} = \frac{d_{n+1}a_n}{(c_{n+1} + d_{n+1}b_n)^2 + d_{n+1}^2 a_n^2},$$

$$b_{n+1} = \frac{c_{n+1} + d_{n+1}b_n}{(c_{n+1} + d_{n+1}b_n)^2 + d_{n+1}^2 a_n^2}$$
(4.7)

for every $n \ge 2$. The sequences in (4.7) for $c_n = d_n = 1$ coincide with (4.2) and (4.3).

Another sequence of continued fractions involving the Cauchy distribution is the following one:

$$U_1 = \frac{1}{1 + \mathbf{C}^2}, \quad U_2 = \frac{1}{1 + \frac{1}{1 + \mathbf{C}^2}}, \quad \dots, \quad U_n = \frac{1}{1 + \frac{1}{1 + \dots + \frac{1}{1 + \mathbf{C}^2}}}.$$
 (4.8)

It is well known that the random variable U_1 possesses the arcsine law. Unlike the sequence V_n studied above the sequence U_n , $n \ge 1$, has a density structure changing with n. Some calculations are sufficient to show that U_1, U_2, U_3, U_4 have density, respectively, equal to

$$f_{U_1}(u) = \frac{1}{\pi\sqrt{u(1-u)}}, \quad 0 < u < 1,$$

$$f_{U_2}(u) = \frac{1}{\pi u\sqrt{(1-u)(2u-1)}}, \quad \frac{1}{2} < u < 1,$$

$$f_{U_3}(u) = \frac{1}{\pi(1-u)\sqrt{(2u-1)(2-3u)}}, \quad \frac{1}{2} < u < \frac{2}{3},$$

$$f_{U_4}(u) = \frac{1}{\pi(2u-1)\sqrt{(2-3u)(5u-3)}}, \quad \frac{3}{5} < u < \frac{2}{3}$$

The general result concerning U_n is stated in the following theorem.

Theorem 4.2. For every $n \ge 1$ the distribution of the random variable U_n is given by

$$\mathbf{P}\{U_n \in \mathrm{d}u\} = \frac{1}{\pi[(-1)^{n+1}\alpha_n + (-1)^n\beta_n u]} \frac{1}{\sqrt{(-1)^n\beta_n + (-1)^{n+1}(\alpha_n + \beta_n)u}} \times \frac{1}{\sqrt{(-1)^{n+1}(\alpha_n + \beta_n) + (-1)^n(\alpha_n + 2\beta_n)u}} \,\mathrm{d}u, \qquad (4.9)$$

where $(-1)^n(\alpha_n + \beta_n)/(\alpha_n + 2\beta_n) < (-1)^n u < (-1)^n \beta_n/(\alpha_n + \beta_n)$, and $\alpha_n, \beta_n \in \mathbf{N}$ satisfy the recursive relationships $\alpha_n = \beta_{n-1}, \beta_n = \alpha_{n-1} + \beta_{n-1}$.

P r o o f. From (4.8) we have that $U_{n+1} = 1/(1 + U_n)$. Then proceeding by induction, i.e., assuming that U_n has distribution (4.9), we obtain that

$$\begin{split} \mathbf{P}\{U_{n+1} \in \mathrm{d}u\} &= \frac{\mathrm{d}}{\mathrm{d}u} \, \mathbf{P}\Big\{U_n > \frac{1-u}{u}\Big\} \mathrm{d}u = \frac{\mathrm{d}}{\mathrm{d}u} \int_{(1-u)/u}^{h(\alpha_n,\beta_n)} \mathbf{P}\{U_n \in \mathrm{d}u\} \\ &= \frac{1}{\pi} \frac{1}{u^2} \frac{1}{[(-1)^{n+1}\alpha_n + (-1)^n \beta_n (1-u)/u]} \\ &\times \frac{1}{\sqrt{(-1)^n \beta_n + (-1)^{n+1} (\alpha_n + \beta_n) (1-u)/u}} \\ &\times \frac{1}{\sqrt{(-1)^{n+1} (\alpha_n + \beta_n) + (-1)^n (\alpha_n + 2\beta_n) (1-u)/u}} \, \mathrm{d}u \end{split}$$

$$= \frac{1}{\pi[(-1)^{n}\beta_{n} + (-1)^{n+1}(\alpha_{n} + \beta_{n})u]} \times \frac{1}{\sqrt{(-1)^{n+1}(\alpha_{n} + \beta_{n}) + (-1)^{n}(\alpha_{n} + 2\beta_{n})u}} \times \frac{1}{\sqrt{(-1)^{n}(\alpha_{n} + 2\beta_{n}) + (-1)^{n+1}(2\alpha_{n} + 3\beta_{n})u}} \,\mathrm{d}u.$$

In the first integral the function $h(\alpha_n, \beta_n)$ represents the right boundary of the support of U_n . We conclude that U_{n+1} possesses distribution (4.9) by taking $\alpha_{n+1} = \beta_n$ and $\beta_{n+1} = \alpha_n + \beta_n$. Theorem 4.2 is proved.

R e m a r k 4.3. The sequence β_n is a Fibonacci sequence, since we have that $\beta_n = \beta_{n-1} + \alpha_{n-1} = \beta_{n-1} + \beta_{n-2}$. We note that the sequence of coefficients α_n and β_n are such that $\lim_{n\to\infty} \alpha_{n+1}/\beta_{n+1} = \lim_{n\to\infty} \beta_n/\beta_{n+1} = \varphi - 1$. On the base of arguments similar to those of Remark 4.1 it is possible to show that the sequence U_n , $n \ge 1$, converges in distribution to $\varphi - 1$. In this case the upper and lower bounds of the domain of definition of the densities $f_{U_n}(u)$, $n \ge 1$ are expressed as ratios of Fibonacci numbers (see Figure 5).

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