

# Anomaly polynomial of general 6d SCFTs

Kantaro Ohmori

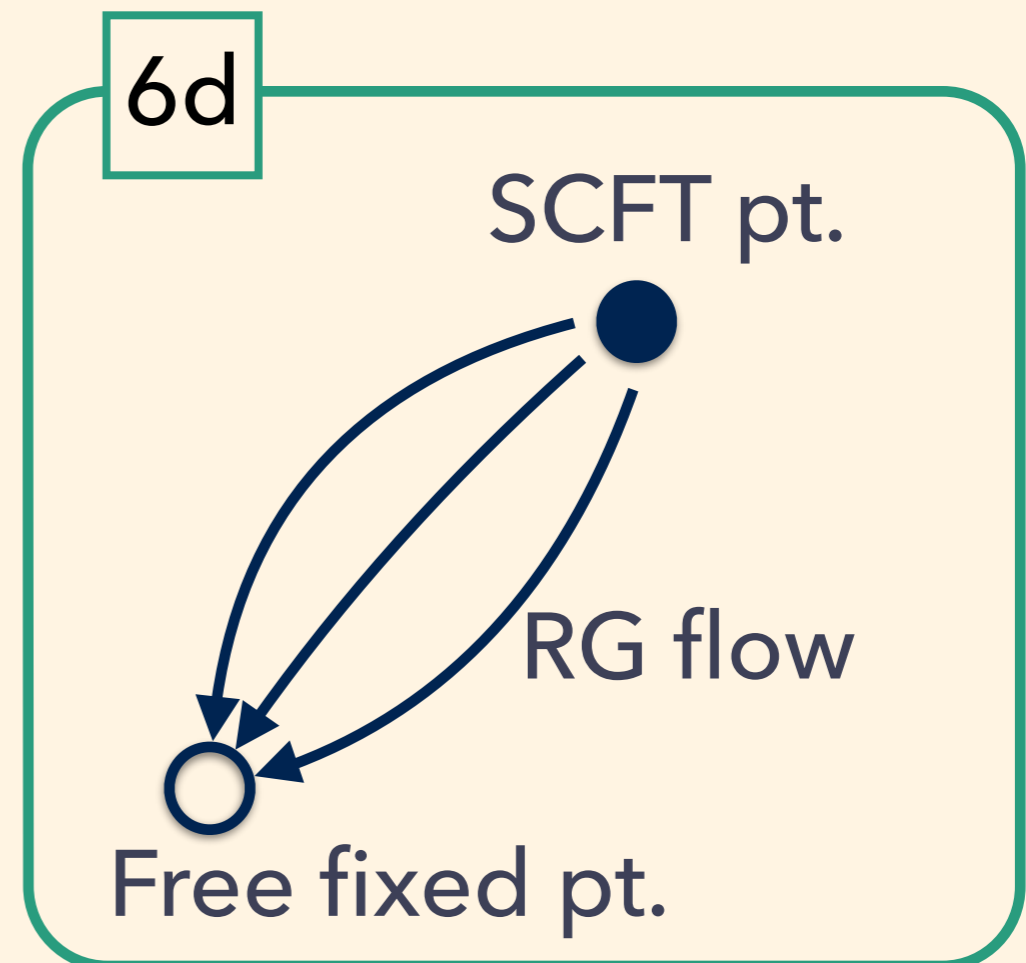
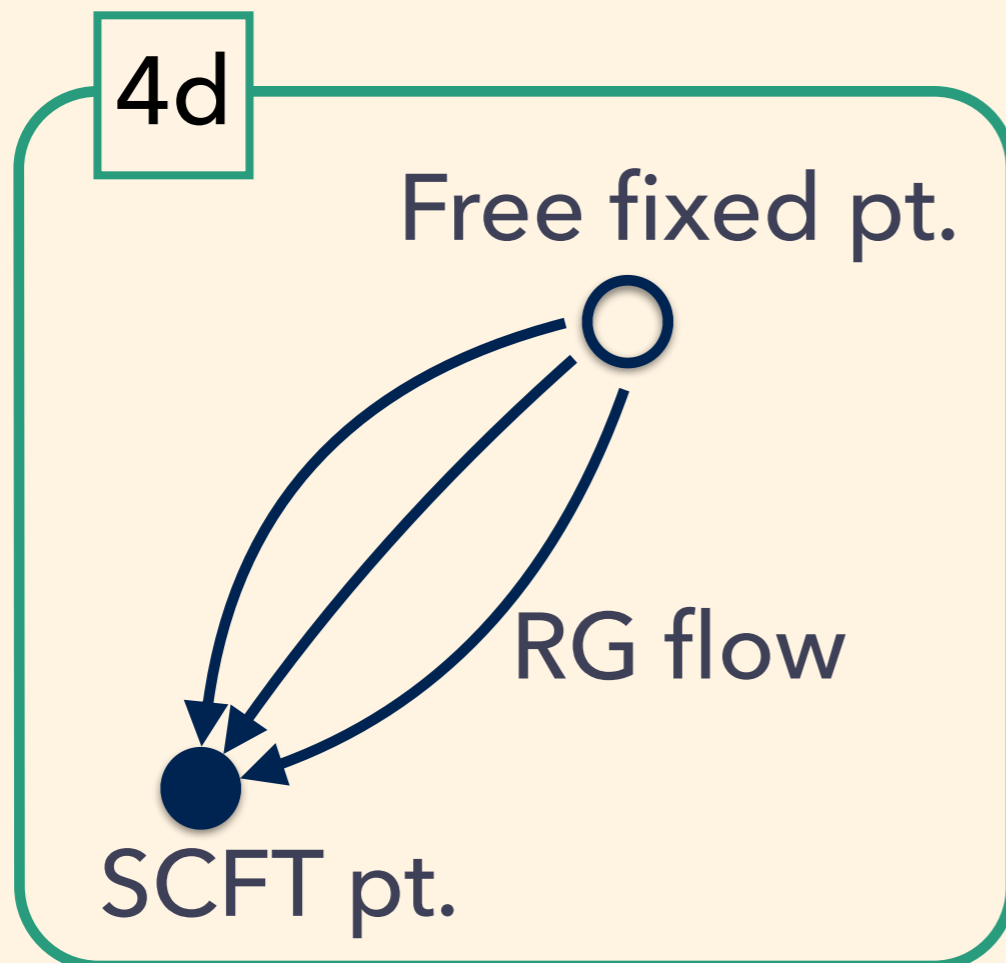
w/ Hiroyuki Shimizu, Yuji Tachikawa, Kazuya Yonekura

arXiv:1408.5572, 1????????

# Introduction

# Interacting 6d SCFTs

- $\mathcal{N}=(2,0)$  or  $\mathcal{N}=(1,0)$  (chiral theories)
- UV fixed point



# Why 6d SCFTs?

- Strongly coupled system  
: no known Lagrangian
- Compactification  $\Rightarrow$  low dimensional systems
- Might control low dimensional dualities  
e.g.  $\mathcal{N}=(2,0)$  theories of ADE type

# What's done?

- String (M,F-) theoretical constructions

[Witten '96],[Ganor,Hanany '96],[Seiberg '97],[Intrilligator Blum '97]

[Brunner,Karch '97],[Hanany,Zaffaroni '97],[Aspinwall,Morrison '97] etc.

[Heckman,Morrison,Vafa '13][Gaiotto Tomasiello '14],

[Del-Zotto,Heckman,Tomasiello,Vafa '14]

- Enormous works for  $\mathcal{N}=(2,0)$  theories

[Gaiotto '09]...

- Anomaly polynomials

for above  $\mathcal{N}=(1,0)$  theories

- Inflow:[Freed, Harvey, Minasian, Moore, '98],[KO,Shimizu,Tachikawa '14]

- Tensor branch:[Intrilligator '14],[KO,Shimizu,Tachikawa,Yonekura '14]

# Plan

- Brane constructions of 6d SCFTs
- Anomaly polynomial
- Tensor branch anomaly matching
- Torus compactifications (if time permits)

# Brane constructions of 6d SCFTs

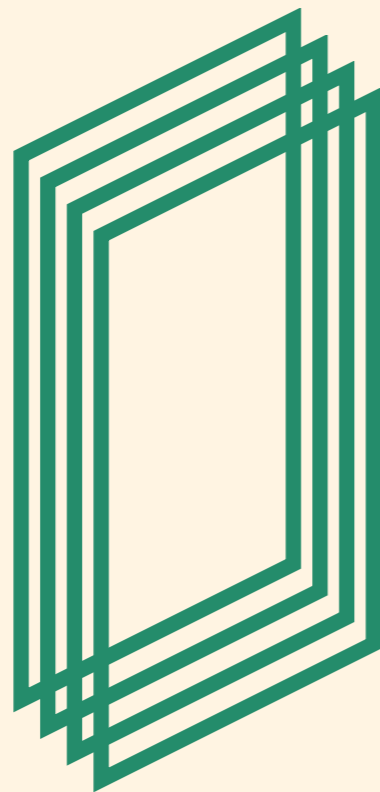
# 6d $\mathcal{N}=(1,0)$ Supermultiplets

- 8 supercharges,  $SU(2)$  R-symmetry
- Tensor multiplet:  $(B_{\mu\nu}^+, \psi^+, a)$   
 $a \in \mathbb{R}$  : "tensor branch" vev (preserves R-sym)
- Vector multiplet:  $(A_\mu, \lambda^-)$   
No scalar
- Hyper multiplet:  $(\phi_i, \psi^+)$   $i = 1, 2, 3, 4$   
 $\phi \in \mathbb{R}^4$  : "Higgs branch" vev (breaks R-sym)



# $\mathcal{N}=(2,0)$ theory of $A$ -type

$Q$  coincident M5-branes

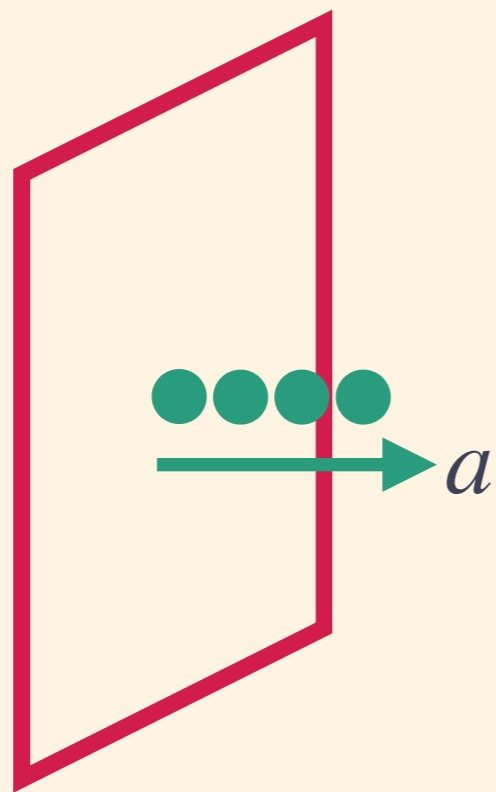


$\Rightarrow A_{Q-1}$ -type  $(2,0)$  theory  
+ center of mass mode

# E-string theory

$Q$  M5-branes on

“End-of-the-world” brane (10d)



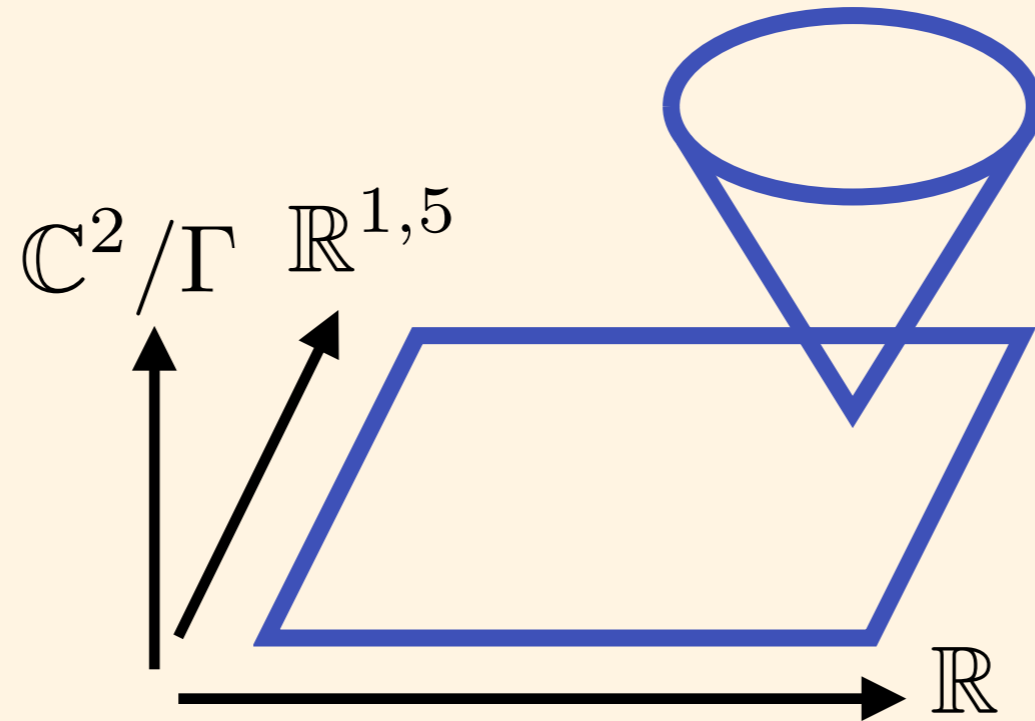
$\Rightarrow$  E-string theory of rank  $Q$   
with  $E_8$  flavor symmetry  
+ free hyper  
(center of mass mode)

# M5-branes on $\mathbb{C}^2/\Gamma_{A,D,E}$

$\Gamma_G \subset SU(2)$  : Finite Subgroup

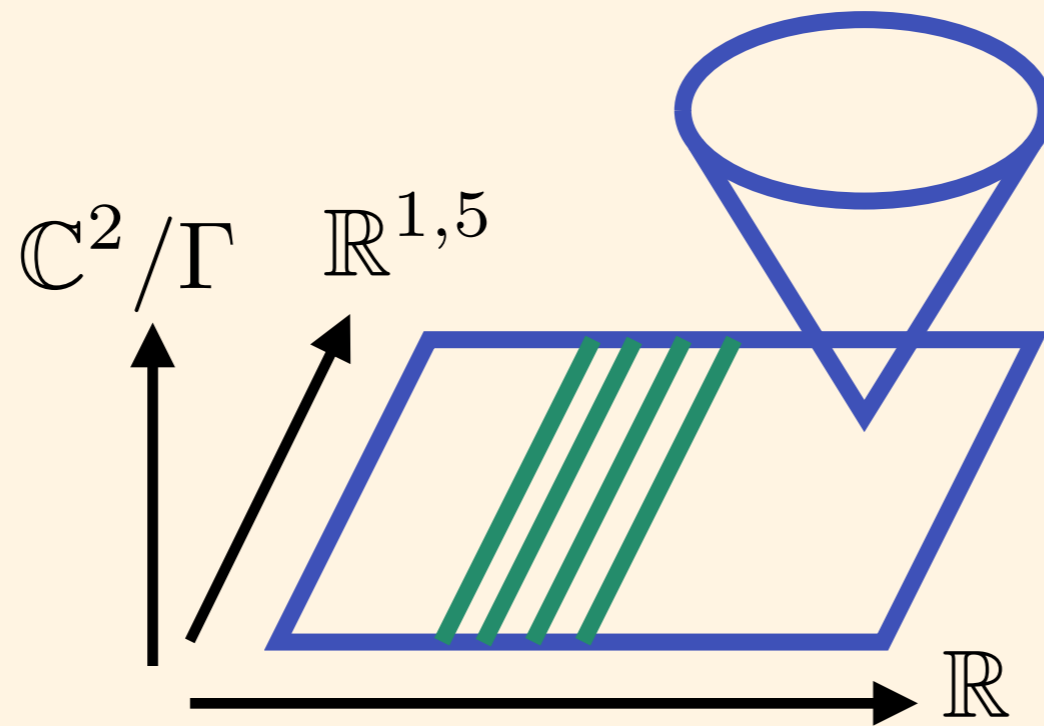
$\Rightarrow$  7d Vector mult. with gauge group  $G$

On singular locus of  $\mathbb{C}^2/\Gamma \times \mathbb{R}^{1,6}$



# M5-branes on $\mathbb{C}^2/\Gamma_{A,D,E}$

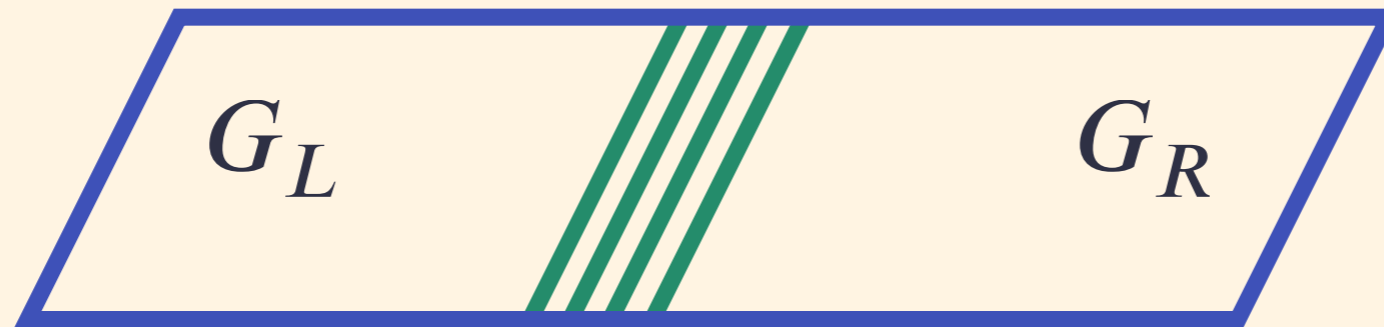
$Q$  M5-branes on  
Singular locus of  $\mathbb{C}^2/\Gamma \times \mathbb{R}^{1,6}$



$G \times G$  flavor symmetry

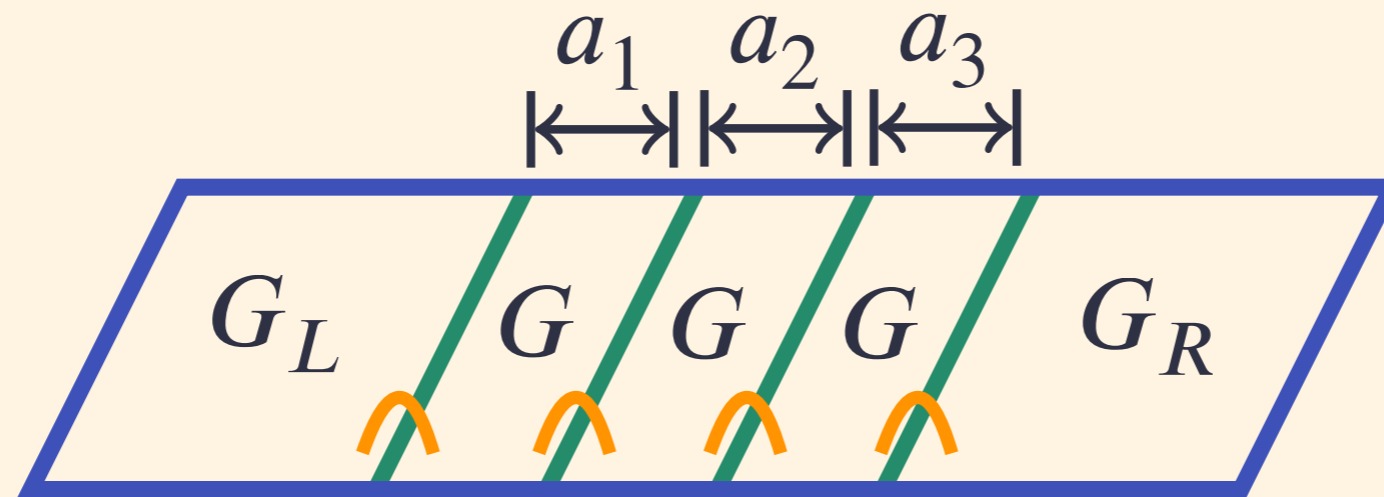
# Tensor branch theory

## $A_k$ -type singularity case



# Tensor branch theory

## $A_k$ -type singularity case



$$G = SU(k + 1) \text{ + bifundamentals}$$

$Q-1$  dynamical vector mult.s of  $SU(k + 1)$

+ bifundamental hypers of neighboring  $SU(k+1)$ 's

+  $Q-1$  Tensor multiplets with tensor vev.s  $a_i$  (dynamical)

+ massive string (M2 branes)

# Tensor branch theory

## $D_k$ -type singularity case



$$G = SO(2k)$$

# Tensor branch theory

## $D_k$ -type singularity case

“Fractional M5”



$$G = SO(2k)$$



# Tensor branch theory

## $D_k$ -type singularity case

“Fractional M5”



$$G = SO(2k) \quad +\text{half bifundamentals}$$
$$G' = USp(2k - 8)$$

$SO(2k)$ - $USp(2k-8)$  alternating quiver theory

Gauge couplings are governed by tensor vev.s (dynamical)

This system is also realized by D6-O6-NS5 system in IIA

# Tensor branch theory

## $D_k$ -type singularity case

$$k = 4$$



$$G = SO(8)$$
$$G' = USp(0)$$

# Tensor branch theory

## $D_k$ -type singularity case

$$k = 4$$



$G = SO(8)$  +rank 1 E-string theories

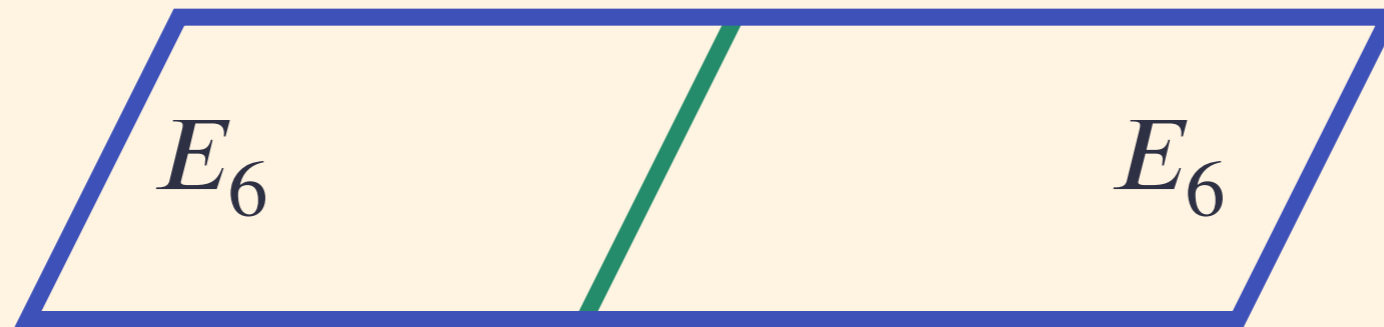
“( $SO(8), SO(8)$ ) conformal matter” = E-string

$$SO(8) \times SO(8) \subset E_8$$

[Del-Zotto, Heckman, Tomasiello, Vafa '14]

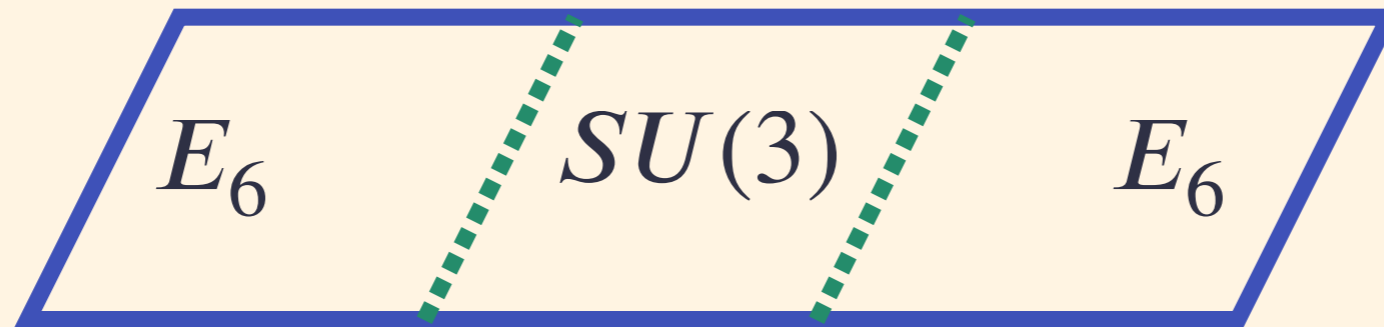
# Tensor branch theory

## $E_6$ -type singularity case



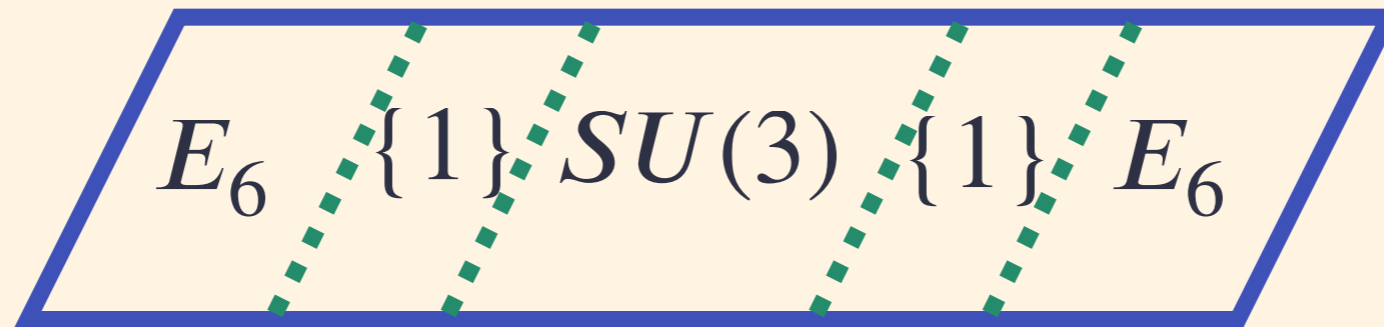
# Tensor branch theory

## $E_6$ -type singularity case



# Tensor branch theory

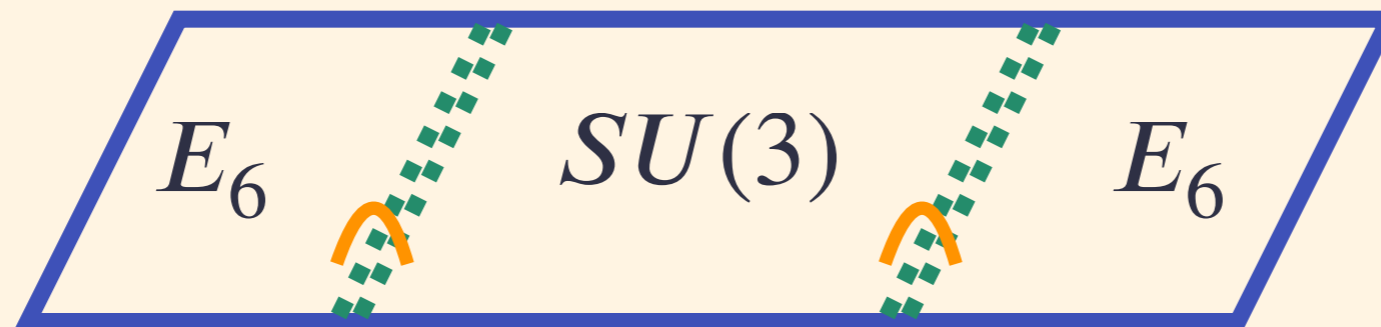
## $E_6$ -type singularity case



# Tensor branch theory

## $E_6$ -type singularity case

[Del-Zotto, Heckman, Tomasiello, Vafa '14]



rank 1 E-string theories

“( $SU(3), E_6$ ) conformal matter” = E-string

$$SU(3) \times E_6 \subset E_8$$

# Other 6d SCFTs

- End-of-the-world brane + ALE singularity  
+ (multiple) M5 branes  
[Del-Zotto, Heckman, Tomasiello, Vafa '14]
- F-theory construction  
[Heckman, Morrison, Vafa '13]
- Lagrangian constructions  
[Smilga '07], [Samtleben, E. Sezgin, and R. Wimmer '11, '12]  
[Ho, Matsuo '14] for (2,0) theories



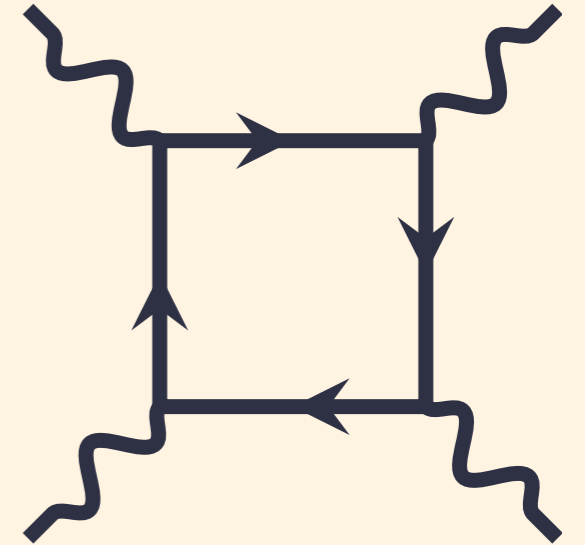
# Anomaly polynomials

# Anomaly in 6d SCFTs

- Anomaly polynomial 8-form  $I_8$

$$\begin{aligned} I_8 \supset & (\text{Tr } F_R^2)^2, & SU(2)_R^4 \\ & (\text{Tr } R^2)^2, (\text{Tr } R^4), & \text{grav.}^4 \\ & (\text{Tr } F_G^2)^2, & G^4 \\ & \text{Tr } F_R^2 \text{Tr } R^2, \text{Tr } F_R^2 \text{Tr } F_G^2, \dots & \text{mixed} \end{aligned}$$

- Anomaly polynomial should exist even for non-perturbative SCFTs.



# 6d Green-Schwarz term

- $\int B \wedge X_4, X_4 \supset \text{Tr} F_R^2, \text{Tr} F_G^2, \text{Tr} R^2$

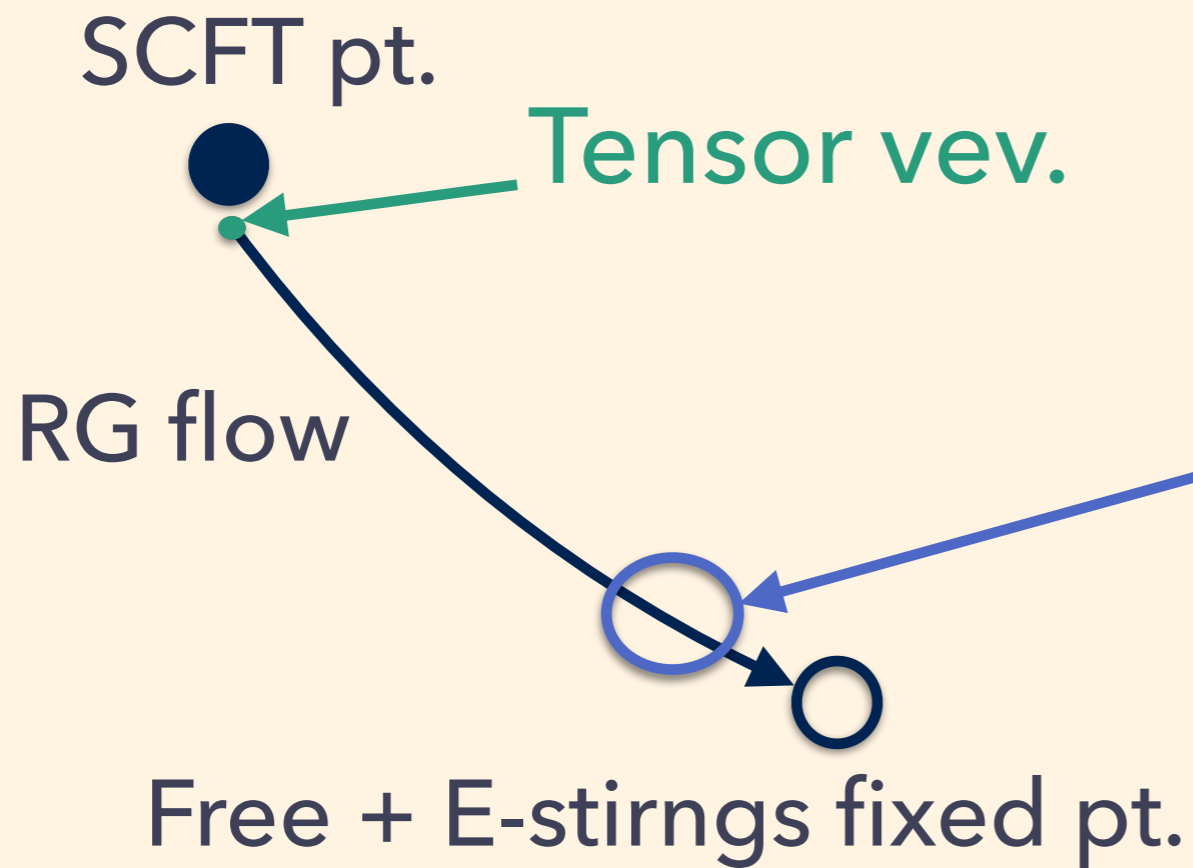
induces modified Bianchi identity  $dH \propto X_4$   
and additional anomaly  $\delta I_8 = \frac{1}{2} X_4^2$  [Monnier '13]

- This contribution is important for calculating anomaly polynomial of 6d SCFTs.

# **Tensor branch anomaly matching**

# Basic idea

[Intrilligator '14],[KO,Shimizu,Tachikawa,Yonekura '14]



vector multiplets,  
"bifundamentals",  
( = hyper or E-string)  
tensor multiplets  
+ corrections  
from massive string

GS terms  $\int B \wedge X$

$$I^{UV} = I^{\text{Naive}} + I^{\text{GS}} \quad \leftarrow \quad I^{\text{GS}} = \frac{1}{2} X^2$$

# Tensor branch anomaly matching

- Vector, tensor, hyper mult. and E-strings generate  $I^{\text{Naive}} \supset (\text{Tr } F_{\text{gauge}}^2)^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } R^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } F_R^2, \text{Tr } F_{\text{gauge}}^2 \text{Tr } F_{\text{flavor}}^2$
- Anomaly for E-strings are calculated from anomaly inflow: [KO, Shimizu, Tachikawa '14]
- Because the UV theory are superconformal,  $I^{UV} = I^{\text{Naive}} + I^{\text{GS}}$  should not contain  $F_{\text{gauge}}$

# Tensor branch anomaly matching

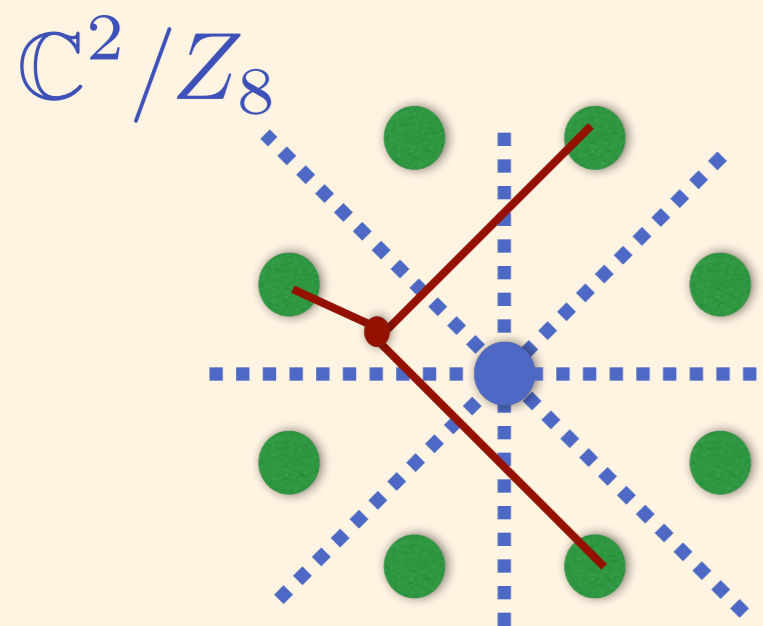
- Example:  $I^{\text{Naive}} \supset -\alpha(\text{Tr} F_{\text{gauge}}^2)^2 - \beta \text{Tr} F_{\text{gauge}}^2 \text{Tr} R^2$
- To cancel this by  $I^{\text{GS}} = \frac{1}{2} X^2$   
 $X$  should contain  
$$X \supset \sqrt{2\alpha} \text{Tr} F_{\text{gauge}} + \frac{\beta}{\sqrt{2\alpha}} \text{Tr} R^2$$
- We can fix all of the GS terms  $\int B_i \wedge X^i$   
in this manner. (If # of  $X^i = \#$  of gauge fields)
- $I^{UV} = I^{\text{Naive}} + I^{\text{GS}}$  : Do square completion!

# Tensor branch anomaly matching

- Result for  $Q$  M5's on Singularities:

$$I^{\text{UV}} \supset \frac{Q^3 |\Gamma|^2}{96} (\text{Tr} F_R^2)^2$$

- The leading behavior can be understood from Higgs branch and **M2 junction**:



$Q^3$  behavior  $\leftarrow$  **M2 junction**.  
 $Q^3 |\Gamma|^3$  ways of suspending junction  
 divided by  $\Gamma$

M5 branes and mirrors



# Torus compactifications

(on going with the same collaborators )

# Torus compactifications

- What is the 4d theory of M5 branes on torus-compactified singularities?

	$\mathbb{R}^{1,3}$	$T^2$	$\mathbb{R}$	$\mathbb{C}^2/\Gamma$
<b>M5</b>	•	•		

- Is that SCFT?
- Can we find candidates in Class S theories?
- How to check?

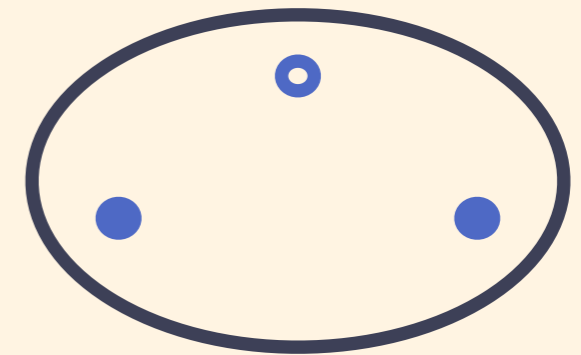
# Known examples and Class S

- Single M5 on  $\Gamma = A_k$  singularity  
⇒ bifundamental both in 6d and 4d.
- Single M5 on  $\Gamma = D_4$  singularity  
⇒ Rank 1 E-string theory in 6d  
⇒ Minahan-Nemeshansky  $E_8$  theory in 4d

# Known examples and Class S

- Single M5 on  $\Gamma = A_k$  singularity  
 $\Rightarrow$  bifundamental both in 4d.

= Class S of type  $A_k$  def'ed by

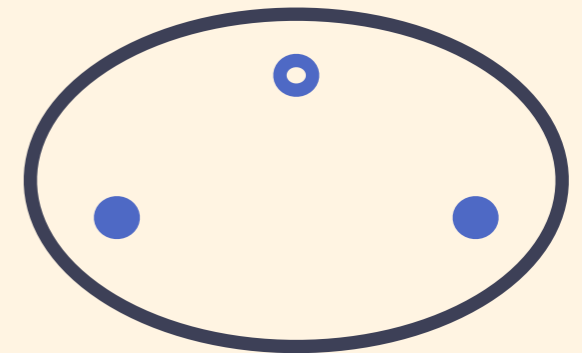


- : “full” puncture (full symmetry)
- : “simple” puncture ( $U(1)$  or no sym.)

# Known examples and Class S

- Single M5 on  $\Gamma = D_4$  singularity  
⇒ Rank 1 E-string theory in 6d  
⇒ Minahan-Nemeshansky  $E_8$  theory in 4d

= Class S of type  $D_4$  def'ed by

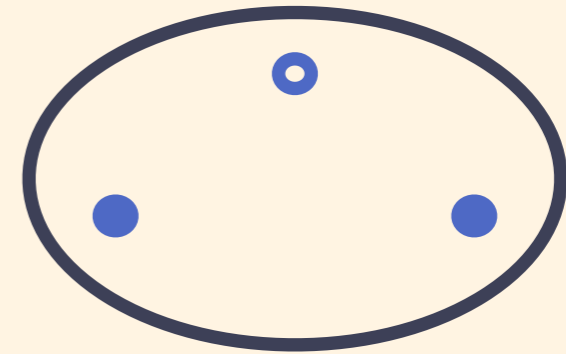


- : "full" puncture (full symmetry)
- : "simple" puncture ( $U(1)$  or no sym.)

# Guess for $Q=1$

Single M5 on  $\mathfrak{g}$ -sing.  $\times T^2$  without center of mass  
(`(G,G) conformal matter`)

?  
= Class S of type  $\mathfrak{g}$  def'ed by



- : “full” puncture (full symmetry)
- : “simple” puncture ( $U(1)$  or no sym.)

“4d conformal matter”

# Checks

- Dim.s of Higgs/Coulomb branches matches
- The geometry of Higgs branch =  $\mathbb{C}^2/\Gamma_{\mathfrak{g}}$
- With plausible assumption, we can calculate the conformal anomalies and flavor levels for one M5 on  $\mathfrak{g}$ -sing.  $\times T^2$   
 $\Rightarrow$  matches with the class S conjecture.
- M/IIB duality chain? Not so clear.

# Conclusion

- The anomaly polynomials of 6d SCFTs can be calculated by adding tensor branch contributions and GS contributions.
- For “6d (G,G) conformal matter”, the candidate for the torus compactified theories can be found in the Class S theories.



# Further directions

- Comactifications of 6d (1,0) SCFTs
  - 4d anomaly polynomials for  $T^2$  compactified theories?
  - What and how the 4d theories are?
- 6d a-theorems?

# Appendix

# Tensor branch

## anomaly matching wo/ vector

- $N=(2,0)$  and E-string theories :  
there is no vector in tensor branch theory
- We know 5d theories obtained by circle compactifications for those theories.
- In 5d coulomb branch corresponds to 6d tensor branch, we have  $U(1)$  vector  $A_i$  come from 6d self-dual tensor  $B_i$  and massive mode because of coulomb vev.

# Tensor branch

## anomaly matching wo/ vector

- In 5d coulomb branch corresponds to 6d tensor branch, we have U(1) vector  $A_i$  come from 6d self-dual tensor  $B_i$  and massive mode because of coulomb vev.
- Those massive modes generates 5d CS terms  $\int A_i \wedge X^i$ , which is calculable.
- This should come from 6d GS term  $\int B_i \wedge X^i$