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Authors

Buizert, Christo Fudge, TJ Roberts, William HG <u>et al.</u>

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PALEOCLIMATE

Antarctic surface temperature and elevation during the Last Glacial Maximum

Christo Buizert¹*, T. J. Fudge², William H. G. Roberts³, Eric J. Steig², Sam Sherriff-Tadano⁴, Catherine Ritz⁵, Eric Lefebvre⁵, Jon Edwards¹, Kenji Kawamura^{6,7,8}, Ikumi Oyabu⁶, Hideaki Motoyama⁶, Emma C. Kahle², Tyler R. Jones⁹, Ayako Abe-Ouchi⁴, Takashi Obase⁴, Carlos Martin¹⁰, Hugh Corr¹⁰, Jeffrey P. Severinghaus¹¹, Ross Beaudette¹¹, Jenna A. Epifanio¹, Edward J. Brook¹, Kaden Martin¹, Jérôme Chappellaz⁵, Shuji Aoki¹², Takakiyo Nakazawa¹², Todd A. Sowers¹³, Richard B. Alley¹³, Jinho Ahn¹⁴, Michael Sigl¹⁵, Mirko Severi^{16,17}, Nelia W. Dunbar¹⁸, Anders Svensson¹⁹, John M. Fegyveresi²⁰, Chengfei He²¹, Zhengyu Liu²¹, Jiang Zhu²², Bette L. Otto-Bliesner²², Vladimir Y. Lipenkov²³, Masa Kageyama²⁴, Jakob Schwander¹⁵

Water-stable isotopes in polar ice cores are a widely used temperature proxy in paleoclimate reconstruction, yet calibration remains challenging in East Antarctica. Here, we reconstruct the magnitude and spatial pattern of Last Glacial Maximum surface cooling in Antarctica using borehole thermometry and firn properties in seven ice cores. West Antarctic sites cooled ~10°C relative to the preindustrial period. East Antarctic sites show a range from ~4° to ~7°C cooling, which is consistent with the results of global climate models when the effects of topographic changes indicated with ice core air-content data are included, but less than those indicated with the use of water-stable isotopes calibrated against modern spatial gradients. An altered Antarctic temperature inversion during the glacial reconciles our estimates with water-isotope observations.

sing oxygen and hydrogen isotope ratios in ancient polar ice as records of past site temperature requires a calibration (1). Surface temperature and the isotopic composition of precipitation correlate spatially in Antarctica, with a regression coefficient α_{s} (spatial slope) of 0.80 per mil per kelvin (%K⁻¹) for δ^{18} O (the ratio of 18 O to 16 O) (2). Reconstructing past temperatures requires regression over time, and this temporal slope $\alpha_{\rm T}$ may differ from $\alpha_{\rm S}$. In East Antarctica, where the longest continuous ice core records, going back to 800 thousand years before present (ka BP), have been extracted (3), independent temperature estimates are not available, and the spatial slope is commonly used to convert isotopic ratios to temperature (1); this approach gives a surface temperature difference $\Delta T_{\rm S}$ of around -9° C between the Last Glacial Maximum (LGM) (26 to 18 ka BP) and the preindustrial period (1, 4, 5).

Antarctic LGM-preindustrial isotope changes depend on many factors, including hemispheric sea surface temperatures (β), sea ice extent (7), ice sheet elevation (β), vapor origin and transport, precipitation seasonality, and postdepositional isotopic exchange (9). Isotopeenabled general circulation models seek to capture these physical processes, making them an invaluable tool for studying isotopic variations. Such models simulate LGMpreindustrial $\alpha_{\rm T}$ ranging from 0.3 to 1.4 % K⁻¹ in central East Antarctica (implied $\Delta T_{\rm S}$ of -4° to -20°C), which implies that several aforementioned processes are poorly constrained (8, 10–12).

We distinguish three temperatures: (i) the climatic temperature T_{CLIM} at constant elevation (relative to the present-day geoid); (ii) the surface temperature T_{S} , which may differ from the climatic temperature because of changing ice sheet topography; and (iii) the vapor condensation temperature T_{C} , which is warmer than the surface because of the strong Antarctic inversion (2, 13).

In this study, we empirically reconstruct LGM surface temperature across Antarctica (Fig. 1) using two independent methods. We investigated five East Antarctic ice cores— EPICA (European Project for Ice Coring in Antarctica) Dome C (EDC), EPICA Dronning Maud Land (EDML), Dome Fuji (DF), Talos Dome (TAL), and South Pole (SP)—and two West Antarctic cores—West Antarctic Ice Sheet (WAIS) Divide (WD) and Siple Dome (SDM).

First, we estimated $\Delta T_{\rm S}$ at EDC and DF from the measured borehole temperature profiles (Fig. 2) using a method similar to that used recently at WD (14). Owing to the downward ice flow and low thermal diffusivity, the ice sheet maintains an imprint of its past surface temperature history. The large ice sheet thickness at EDC and DF is favorable for preserving past temperatures, yet the low accumulation rate is not. Consequently, the relative uncertainty in the EDC and DF borehole reconstructions is larger than that at WD. To constrain the problem better, we used downward ice velocities measured by means of phase-sensitive radio-echo sounding (EDC only) and accurate age constraints derived through volcanic synchronization to the layer-counted WD time scale.

We forced a one-dimensional heat transportice flow model at the surface boundary with a temperature history that is based on the δ^{18} O record scaled with a constant α_{T} value (10). Applying traditional isotope scaling ($\alpha_T \approx$ 0.7 ‰K⁻¹, yielding $\Delta T_{\rm S} = -9^{\circ}$ C at EDC and -7.5°C at DF) simulates temperature profiles that do not fit the borehole observations at either site (Fig. 2). At EDC, the model-data fit is optimized for $\alpha_{\rm T}$ = 1.14 ‰K⁻¹, which is consistent with $\Delta T_{\rm S}$ = -5.5°C (95% confidence range is −6.9° to −3.1°C). At DF, the optimal $\Delta T_{\rm S}$ is in the -2.0° to -5.4°C range; we provide a range without a best estimate because, at DF, there are no direct constraints on the downward ice velocity. In Fig. 1, the WD, EDC, and DF borehole estimates are marked "BH."

Second, we reconstructed past climate at all seven sites using the dependence of firm densification, the gradual transformation of polar snow to ice, on $T_{\rm S}$ and accumulation rate (*A*). Air bubbles are isolated from the atmosphere at the lock-in depth (50 to 120 m below the surface), an event preserved in two ice core signals (*15*): δ^{15} N of N₂ which records past firm column thickness by means of gravitational enrichment, and the gas age–ice age difference, or Δ age. The δ^{15} N and Δ age-isopleths are perpendicular in $T_{\rm S}$ -*A* space (Fig. 3A), meaning that if δ^{15} N and Δ age are independently known, a distinctive climatic ($T_{\rm S}$, *A*)

¹College of Earth Ocean and Atmospheric Sciences, Oregon State University, Corvallis, OR 97331, USA. ²Department of Earth and Space Science, University of Washington, Seattle, WA 98195, USA. ³Geographical and Environmental Sciences, Northumbria University, Newcastle, UK. ⁴Atmosphere and Ocean Research Institute, The University of Tokyo, Kashiwa 277-8568, Japan. ⁵Université Grenoble Alpes, CNRS, IRD, IGE, Grenoble, France. ⁶National Institute of Polar Research, Tachikawa, Tokyo, Japan. ⁷Department of Polar Science, The Graduate University of Colorado, Boulder, CO 80309, USA. ¹⁰British Antarctic Survey, Cambridge, UK. ¹¹Scripps Institution of Oceanography, University of California, San Diego, La Jolla, CA 92093, USA. ¹²Center for Atmospheric and Oceanic Studies, Graduate School of Science, Tohoku University, Sendai 980-8578, Japan. ¹³The Earth and Environmental Systems Institute, Pennsylvania State University, University Park, PA 16802, USA. ¹⁴School of Earth and Environmental Sciences, ISP. ^CNR, Venice-Mestre, Ital. ¹¹Scoupe Bern, 3012 Bern, Switzerland. ¹⁶Department of Chemate Science Department, New Mexico Tech, Socorro, NM 87801, USA. ¹⁹Niels Bohr Institute, University of Colorado, Colenabagen, Copenhagen, Denmark. ²⁰School of Earth and Sustainability, Northern Arizona University, Flagstaff, AZ 86011, USA. ²¹Department of Geography, Ohio State University, Columbus, OH 43210, USA. ²²National Center for Atmospheric Research, Boulder, CO 80307, USA. ²³Climate and Environmental Research Laboratory, Arctic and Attractic Research Institute, Science, ISP. ^CNR, Venice-Mestre, Ital. ²¹Laboratoire des Sciences du Climat et de l'Environment-IPSL, University Fars-Saclay, Gif-sur-Yvette, France. ***Corresponding author. Email: christo.buizert@oregonstate.edu**

solution exists (subject to the uncertainties of the firn model).

Synchronization using both volcanic deposits and globally synchronous abrupt atmospheric methane variations allowed us to estimate Δ age empirically for the Antarctic ice cores (10, 16). We used an inverse dynamical firm densification-heat transport model (17, 18) to reconstruct $T_{\rm S}$ and A histories that optimize the fit to Δ age and δ^{15} N data (Fig. 3, B and C). Reconstructed accumulation rates agree (within uncertainty) with independent estimates (fig. S8). Methodological biases and uncertainties are estimated by using a Monte Carlo approach (10). The histograms in Fig. 1 give the $\Delta T_{\rm S}$ distribution of the Δ age-based reconstruction.

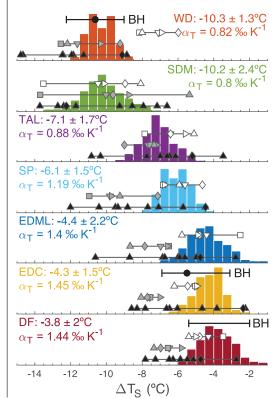
In East Antarctica, $\Delta T_{\rm S}$ ranges from $-3.8^{\circ} \pm$ $2.0^{\circ}C$ (DF) to $-7.1^{\circ} \pm 1.7^{\circ}C$ (TAL); at DF, EDC, and EDML, $\Delta T_{\rm S}$ is substantially lower than estimates from isotope scaling that use α_{s} . The two West Antarctic sites have similar $\Delta T_{\rm S}$ of $-10.2^{\circ} \pm 2.4^{\circ}$ C (SDM) to $-10.3^{\circ} \pm 1.3^{\circ}$ C (WD). The \triangle age- and borehole-based reconstruction methods agree within uncertainty at all sites (Fig. 1). Allowing for more flanklike ice flow at EDC during the glacial period (which would occur if the divide position were different from that at present) improves the agreement by changing the borehole estimate to around $-4.5^{\circ}C(10)$; we choose to report the -5.5°C value to keep both methods independent. PMIP4 (Paleoclimate Modeling Intercomparison Project phase 4) simulations (19) find a seven-site-mean $\Delta T_{\rm S}$ magnitude that is $1.2^{\circ} \pm 4.6^{\circ}$ C larger than our Δ age-based reconstructions (mean and spread of 10 climate models; Fig. 1).

We emphasize that the firn method is primarily constrained by the empirical ∆age estimates. Because $T_{\rm S}$ and A broadly covary via the saturation vapor pressure, the deglacial climatic changes run parallel to the δ^{15} Nisopleths (Fig. 3A). Therefore, $\delta^{15}N$ data alone do not constrain the magnitude of climate change meaningfully. The effects of $T_{\rm S}$ and A are additive in Δ age, however, making Δ age a sensitive proxy for climate change (Fig. 3D), as first noted by Schwander et al. (20). The empirical ∆age at 24 ka BP is larger than at 18 ka BP for all five cores where both are available, and coldest conditions in Antarctica occur ~27 to ~24 ka BP in our reconstructions (fig. S8h); this follows expectations from local insolation (21).

We propose that elevation changes explain the spatial differences in $\Delta T_{\rm S}$ (8). Let Δz be the LGM elevation anomaly relative to the present. We present WD and DF total aircontent data (fig. S12) and interpret them in terms of elevation change (22). These data suggest a 420-m (range, 280 to 590 m) contrast in Δz between WD and central East Antarctica (here, DF and EDC)—for example,

 $\Delta z = +300$ m at WAIS and $\Delta z = -120$ m in central East Antarctica (Fig. 4B). Our estimate is broadly in agreement with LGM ice sheet reconstructions that suggest a West-East Δz contrast between 160 and 560 m (10). Although the implied Δz at WAIS exceeds the observed highstand at ice margin nunataks (23), such data do not strongly constrain the elevation at WD more than 500 km away. The corresponding $\Delta T_{\rm S}$ contrast (WD $\Delta T_{\rm S}$ minus the average $\Delta T_{\rm S}$ at DF and EDC) is $-6.2^{\circ} \pm 2.3^{\circ}$ C in the Δ age-based reconstructions, $-6.0^{\circ} \pm 2.0^{\circ}$ C in the borehole reconstructions, and $-5.9^{\circ} \pm 2.7^{\circ}C$ in the PMIP4 model ensemble; the level of agreement suggests this is a robust feature of Antarctic LGM climate. This temperature contrast is thus plausibly linked to Δz through the (spatial) lapse rate in the interior of Antarctica of around -12° C km⁻¹ (2, 24).

To further assess the elevation impact on $\Delta T_{\rm S}$, we perform an atmosphere-ocean general circulation model (AOGCM) sensitivity study of Antarctic LGM climate using the MIROC (Model for Interdisciplinary Research on Climate) and HadCM3 (Hadley Centre Coupled Model version 3) models and a series of LGM topographic reconstructions (10). We first estimated climatic LGM cooling using full LGM boundary conditions (including LGM albedo) but preindustrial Antarctic topography; this yielded a seven-site average ΔT_{CLIM} of -4.7° and -7.0°C in the MIROC and HadCM3 models, respectively, but stronger albedo-driven cooling is found over the Ross and Weddell Seas due to ice growth onto the continental



shelf (Fig. 4A). Simulated climatic ΔT_{CLIM} is similar in interior West and East Antarctica in the absence of topographic change.

Next, we performed climate simulations with five Antarctic LGM topographic reconstructions. These reconstructions suggest Δz of +100 to +600 m in interior WAIS and down to -250 m in interior East Antarctica (Fig. 4B). These changes result in greater $\Delta T_{\rm S}$ in West than in central East Antarctica (Fig. 4C), in agreement with our reconstructions. By comparing the various topographic reconstructions, we find that $\Delta T_{\rm S}$ is closely linked to Δz in both models through the dry adiabatic lapse rate of -9.8°C km⁻¹ (Fig. 4D). Also, a fraction of the variance is due to the topography altering the atmospheric circulation around Antarctica, rather than the direct lapse-rate effect. We find a correlation r = 0.96 between the reconstructed and the simulated site $\Delta T_{\rm S}$ pattern (averaged across the five topographic reconstructions and both models); for the PMIP4 multimodel mean, this correlation is r = 0.95. We conclude that changes in LGM ice sheet topography plausibly explain the $\Delta T_{\rm S}$ spatial variability in our reconstruction (8).

Our findings have implications for the interpretation of water isotopes in Antarctic ice cores. We found α_T in the range of 0.9 to 1.4 ‰K⁻¹ in East Antarctica and, therefore, $\alpha_T > \alpha_S$, opposite to Greenland, where $\alpha_T < \alpha_S$ (*17*, 25). We compared our α_T with those from LGM and preindustrial simulations using the latest-generation isotope-enabled Community Earth System Model (iCESM) (Fig. 4E). The

Fig. 1. Summary of Antarctic LGM

cooling estimates. Black markers with horizontal error bars marked "BH" give borehole estimates; WD results are from (14). Histograms give distribution of Δ age-based temperature reconstructions from a Monte Carlo sampling (N = 1000) of model parameters: listed are mean and 2σ standard deviation of the distribution, as well as the implied temporal isotope slope α_{T} . ΔT_{S} is the LGM (18 to 21.4 ka BP) minus preindustrial (0.5 to 2.5 ka BP) condition. White (MIROC), grav (HadCM3), and black (PMIP4) show AOGCM-simulated $\Delta T_{\rm S}$, with symbols denoting different LGM topography reconstructions (10): Pollard and Deconto (downward triangle) (32); Whitehouse et al. (square) (33); Glac-1D (diamond) (29); Golledge et al. (rightward triangle) (34); and Ice-6G (upward triangle).

good agreement (r = 0.91; 0.06 ‰K⁻¹ mean offset) demonstrates that our reconstructed $\alpha_{\rm T}$ are consistent with isotope physics, yet the large intermodel spread in simulated $\alpha_{\rm T}$ [see section S3.5 in (10) for a review] prevents us from claiming that it validates our results. Although the $\alpha_{\rm T}$ agree well, iCESM simulated a $\Delta T_{\rm S}$ and LGM-preindustrial δ^{18} O change that are both too large (compared with our reconstructions and ice core data, respectively).

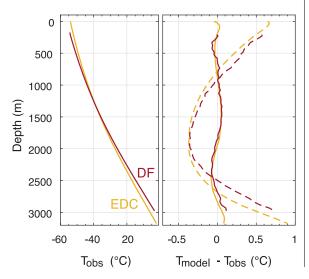
Fig. 2. Borehole temperature reconstruction for EDC and DF.

(Left) Site borehole temperature observations at EDC (yellow) and DF (red). At both sites, the ice-bedrock interface is at the pressure melting point (-2.2°C). (Right) Model-data mismatch at EDC (yellow) and DF (red) for an ice flow-heat transport model forced by the optimized temperature histories (solid lines, $\Delta T_{\rm S}$ of -5.5°C at EDC and -3.2°C at DF) and forced with water-isotope scaling of 0.7 ‰K⁻¹ (dashed lines, $\Delta T_{\rm S}$ of -9.0°C at EDC and -7.5°C at DF).

Fig. 3. ∆age-based temperature reconstructions.

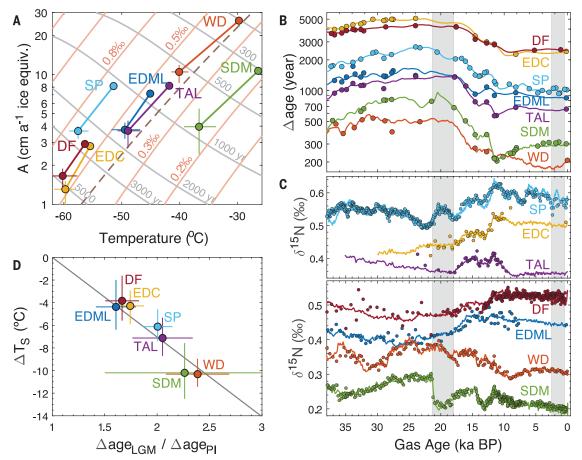
(**A**) Δ age and δ ¹⁵N-isopleths (gray and salmon, respectively) in the steady-state Herron-Langway firn densification model as a function of $T_{\rm S}$ and A. The dashed line shows accumulation scaling by means of the saturation vapor pressure at the site (ignoring the atmospheric inversion). Reconstructed preindustrial and LGM conditions at the seven sites are indicated. (B) Model fit to empirical Δ age constraints. Grav vertical bars denote the LGM (21.4 to 18 ka BP) and preindustrial (2.5 to 0.5 ka BP) periods. (C) Model fit to δ^{15} N data, divided over two panels to prevent overlapping curves. Data are shown on the WD2014 time scale (30, 31). (D) Reconstructed $\Delta T_{\rm S}$ versus ratio of LGM ∆age over preindustrial Δage (with linear fit) (Δage_{n}), showing the utility of Δ age as a climate proxy.

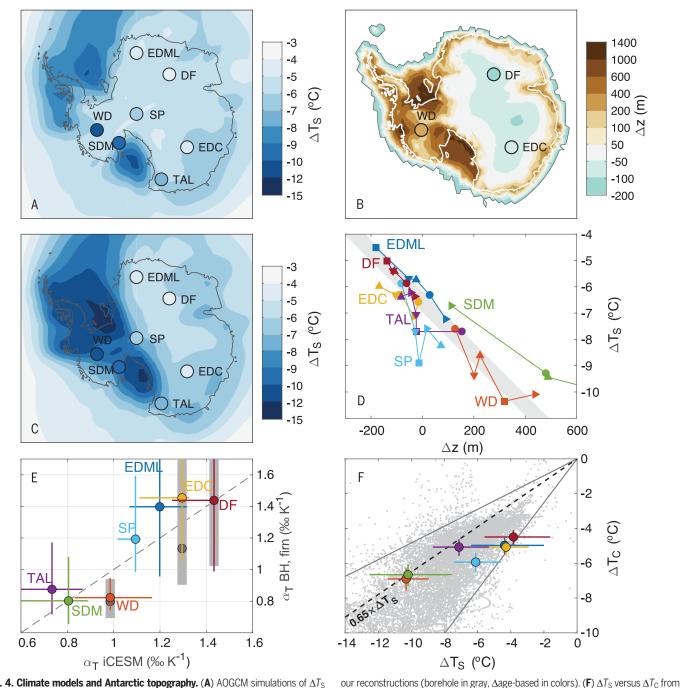
Last, we investigated changes to the strong surface-based inversion in the Antarctic boundary layer (Fig. 4F). The condensation temperature $T_{\rm C}$ is higher than $T_{\rm S}$, and they correlate spatially with a slope $dT_{\rm C}/dT_{\rm S}$ in the 0.63 to 0.67 range (2, 13, 26). $T_{\rm C}$ controls precipitation δ^{18} O, with a present-day spatial sensitivity of $d\delta^{18}$ O/ $dT_{\rm C} = d\delta^{18}$ O/ $dT_{\rm S} \times dT_{\rm S}/dT_{\rm C} \approx 0.80/0.65 = 1.23 \% {\rm K}^{-1}$. We now assume that, unlike $\Delta T_{\rm S}$, the LGM-preindustrial change $\Delta T_{\rm C}$



can be estimated by using this spatial slope via $\Delta T_{\rm C} = \Delta \delta^{18} {\rm O}/1.23$ (Fig. 4F). At WD and SDM, the $\alpha_T \approx \alpha_S$ assumption holds, suggesting that the ratio $\Delta T_{\rm C}/\Delta T_{\rm S}$ is close to the present-day ratio of 0.65; in central East Antarctica, the ratio $\Delta T_{\rm C}/\Delta T_{\rm S}$ exceeds 0.65, which is consistent with $\alpha_{\rm T} > \alpha_{\rm S}$. We plotted simulated $\Delta T_{\rm S}$ versus $\Delta T_{\rm C}$ across interior Antarctica from a wide range of AOGCMs and topographies; we found that the ratio $\Delta T_{\rm C}/\Delta T_{\rm S}$ ranges from 0.48 to 1.3 (95% interval, gray lines), with our empirical reconstructions falling within the model data cloud (Fig. 4F). In aggregate, these simulations find that $\Delta T_{\rm C}/\Delta T_{\rm S}$ tends to exceed the present-day ratio of 0.65 (~79% of model data points): such a change to the inversion structure would result in $\alpha_{\rm T} > \alpha_{\rm S}$ for $\Delta T_{\rm S}$. In the iCESM simulations, the $\Delta T_{\rm C}/\Delta T_{\rm S}$ and $\alpha_{\rm T}$ fields look similar, with the $\Delta T_{\rm C}/\Delta T_{\rm S} = 0.65$ contour line broadly aligning with the $\alpha_{\rm T}$ = $0.8~{\rm \% K^{-1}}$ contour line (fig. S11). We conclude that physically plausible changes to the inversion (27, 28) may reconcile our reconstructions with previous work on Antarctic LGM water isotopes.

Our reconstructions improve the LGM Antarctic temperature estimation and provide a benchmark for testing the ability of (isotope-enabled) climate models to simulate climate states radically different from the late





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Fig. 4. Climate models and Antarctic topography. (A) AOGCM simulations of $\Delta T_{\rm S}$ using preindustrial ice topography in Antarctica (average of MIROC and HadCM models), with Δ age-based $\Delta T_{\rm S}$ reconstructions for the seven sites. (B) Simulated LGM elevation anomaly (shaded, average of five topographies) with LGM elevation anomaly of +310, -80, and -140 m at WD, EDC, and DF (*10*). (C) As in (A), but using LGM ice topography in Antarctica (average of five LGM topographies and both MIROC and HadCM models). (D) Elevation change versus $\Delta T_{\rm S}$ in the AOGCM simulations (average of MIROC and HadCM models); symbols denote the different LGM topographic reconstructions (see Fig. 1 caption for legend). The gray bar shows the dry adiabatic lapse rate. (E) Temporal isotope slope $\alpha_{\rm T}$ from the iCESM model against

Holocene. For surface temperature, the spatial isotopic slope is not always a good approximation of the temporal slope, challenging the prevalent interpretation of ice core water isotopes in Antarctica.

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 Δ age-based $\Delta T_{\rm S}$ and isotope-based $\Delta T_{\rm C}$ (large dots with error bars) and from

LGM-preindustrial AOGCM simulations (small gray dots, gray lines enclose the central 95% of estimates); the black dashed line represents the modern spatial slope

(2). Models plotted are PMIP3 [except for one model that simulates $\Delta T_{\rm S} > 0^{\circ}$ C],

PMIP4 (all model output publicly available), and all iCESM, MIROC, and HadCM3

800 hPa); $T_{\rm C}$ is taken to be the annual mean troposphere temperature maximum

(typically ~500 hPa). The models have an average preindustrial spatial dT_{c}/dT_{s}

of 0.68 (range, 0.31 to 0.89) in interior Antarctica.

simulations used in this work; we show interior Antarctica (surface pressure <

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SUPPLEMENTARY MATERIALS

science.sciencemag.org/content/372/6546/1097/suppl/DC1 Materials and Methods Figs. S1 to S12 Tables S1 to S7 References (*36–160*) Data S1 18 June 2020; accepted 29 April 2021 10.1126/science.abd2897



Antarctic surface temperature and elevation during the Last Glacial Maximum

Christo Buizert, T. J. Fudge, William H. G. Roberts, Eric J. Steig, Sam Sherriff-Tadano, Catherine Ritz, Eric Lefebvre, Jon Edwards, Kenji Kawamura, Ikumi Oyabu, Hideaki Motoyama, Emma C. Kahle, Tyler R. Jones, Ayako Abe-Ouchi, Takashi Obase, Carlos Martin, Hugh Corr, Jeffrey P. Severinghaus, Ross Beaudette, Jenna A. Epifanio, Edward J. Brook, Kaden Martin, Jérôme Chappellaz, Shuji Aoki, Takakiyo Nakazawa, Todd A. Sowers, Richard B. Alley, Jinho Ahn, Michael Sigl, Mirko Severi, Nelia W. Dunbar, Anders Svensson, John M. Fegyveresi, Chengfei He, Zhengyu Liu, Jiang Zhu, Bette L. Otto-Bliesner, Vladimir Y. Lipenkov, Masa Kageyama and Jakob Schwander

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Antarctic paleotemperatures

It has been widely thought that East Antarctica was ~9°C cooler during the Last Glacial Maximum, close to the ~ 10°C difference between then and now determined independently for West Antarctica. Buizert *et al.* used borehole thermometry, firn density reconstructions, and climate modeling to show that the temperature in East Antarctica was actually only ~4° to 7°C cooler during the Last Glacial Maximum. This result has important consequences for our understanding of Antarctic climate, polar amplification, and global climate change. *Science*, abd2897, this issue p. 1097

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