1 2 3	DOI: 10.1002/ ((please add manuscript number)) Article type: Full Paper
4 5 6 7	Associative Enhancement of Time Correlated Response to Heterogeneous Stimuli in a Neuromorphic Nanowire Device
8 9 10	Curtis J. O'Kelly1,3, Jessamyn A. Fairfield1,3, David McCloskey2,3, John F. Donegan2,3 and John J. Boland1,3*
10 11 12 13	Dr. Curtis J. O'Kelly 1,3, Dr. Jessamyn A. Fairfield 1,3, Dr. David McCloskey 2,3, Prof. John F. Donegan 2,3 and Prof. John J. Boland 1,3
14 15 16 17 18	1 School of Chemistry, Trinity College Dublin, Dublin 2, Ireland 2 School of Physics, Trinity College Dublin, Dublin 2, Ireland 3 CRANN Institute, Trinity College Dublin, Dublin 2, Ireland *E-mail: jboland@tcd.ie. Phone: +353 1-896-3140.
19 20 21 22 22	Keywords: Memory, Learning, Associative Memory, Neuromorphic hardware, Pulse Stimuli, Synapse
23 24	In spite of the strong interest in brain-like or neuromorphic computation, relatively few
25	devices have emerged whose neuromorphic behavior is embedded in the hardware itself and
26	not reliant on external programming of synaptic weights. We describe here a neuromorphic
27	device based on a TiO ₂ nanowire that exhibits an associative memory response to the time
28	correlation between voltage and optical stimuli. Memristive characteristics are also observed
29	with current-voltage sweeps showing hysteresis loops and continuum resistance levels. The
30	nanowire device responds to heterogeneous voltage and optical pulse stimuli with spike-like
31	neuromorphic outputs. Moreover, uncorrelated pulses produce a weak response, consistent
32	with the interaction of coincident pulses with adsorbed and bulk oxygen in the surface
33	depletion region, leading to a nonlinear enhancement in conductance. The strength of this
34	learned enhancement depends on the both the time correlation and number of pulse stimuli,
35	consistent with spike timing dependent plasticity. The nanowire devices presented have neural
36	synapse-like properties that could serve as a building block for neuromorphic computation.
37	

1. Introduction

39

The human brain's ability to temporally associate multiple complex stimuli is a key 40 requirement for neuromorphic learning.^[1,2] Although many biologically-inspired silicon-based 41 technologies have reproduced basic sensory functions from sight to hearing,^[3, 4] mimicking 42 43 neural cognition and functionality in silicon has thus far relied on externally-set weighted 44 outputs to demonstrate learning, an approach which suffers from large device footprints, memory-computation lag and high power consumption.^[5-7] In an ideal neuromorphic device, 45 the synaptic weights between elements are modulated by electrical pulses from neighboring 46 47 devices and sensitive to a range of heterogeneous stimuli. These synaptic weights are 48 strengthened when pulses arrive at nearly the same time which produces a phenomenon known as spike timing dependent plasticity (STDP) and which leads to learning.^[2] Recent studies 49 seeking to emulate this behavior have focused on an emerging class of resistance switching 50 51 devices called memristors, whose non-volatile continuum conductance states closely imitate the analog response of individual biological synapses.^[8-11] Neuromorphic hardware based on 52 53 ionic channels and resistive switching materials has also been shown to emulate synaptic weight augmentation in response to external stimuli.^[10-14] However, the devices reported to date 54 respond to just a single type of stimulus, namely voltage pulses.^[15] 55

56 Here we show associative, time-dependent correlation between voltage and optical pulses in 57 a nanoscale neuromorphic device comprised of an Au-contacted TiO2 nanowire. The observed 58 behavior is similar to synaptic weight augmentation, but for physically heterogeneous stimuli, 59 demonstrating novel associative memory not previously shown for memristive devices. 60 Uncorrelated stimuli produce a weak electrical response, while coincident stimuli create a 61 temporary modification of the device properties that leads to an enhanced conductance. The 62 enhancement level depends on the time correlation and number of pulse stimuli, emulating 63 STDP and demonstrating neuromorphic learning in a solid state device. This device has the potential to broaden not only the memristor concept but also to form a novel neuromorphic 64

- hardware platform sensitive to multiple physically dissimilar stimuli, and which may also have
 applications in optoelectronic computation.^[16, 17]
- 67

68 2. Results and Discussion

69

70 The response of single TiO₂ nanowire devices to voltage and optical pulse stimuli was studied using the experimental setup in Figure 1a. As previously shown, the device is initially activated 71 by applying a large voltage (typically +10 V) until the current attains a steady-state value to 72 73 create a steady state population of charged dopants at the Au-TiO₂ interface that facilitates charge injection into the wire.^[18] The magnitude of charge injection is modulated by the 74 75 population of oxygen vacancies at the metal interface lowering the tunneling barrier. The effect 76 produces a memristive diode-like response linked directly to the history of voltage pulses 77 applied to the device. The current response to a half-wave sinusoidal voltage input is shown in Figure 1b and characterized by memristive hysteresis loops that evolve to higher conductance 78 79 levels with each voltage sweep. The corresponding *I-V* curve in Figure 1c is consistent with a neuromorphic resistive switch, displaying a continuum of conductance states. Similar devices 80 81 have also displayed evolving conductance levels in response to a repeated voltage pulse stimulus, tunable memory levels, and a unique reset state.^[18] An important aspect of the device 82 83 operation is the presence of a threshold for the onset of current rectification in either direction. This occurs below ± 3 V and is highlighted in the inset of Figure 1c. Now, building on the 84 85 previously reported memory properties, optical pulse stimuli are added in addition toto the 86 voltage stimuli previously reported.

The photocurrent ($I_{pc} = I_{Measured} - I_{Dark}$) recorded during one 20 s laser pulse is shown in **Figure 2**a. The photocurrent response is initially linear but then saturates above 5 V. The response of the device to uncorrelated 20 s voltage and optical pulses is shown in Figure 2b, in addition to a combined pulse where the voltage and optical stimuli are coincident. The hold

voltage is +2 V and was maintained for 300 s prior to the voltage pulse ($\Delta V = 0.75$ V) to 91 92 establish the current baseline. Critically, both the hold voltage and peak pulse voltage are below 93 the threshold for modifying device conductance. Whereas the responses to the uncorrelated 94 pulses is modest, the coincident pulse elicits an exceptional response that is clearly larger than 95 the sum of the responses to the individual pulses. The current enhancement can be computed from the additional charge dQ = I dt generated during each voltage, optical and coincident 96 pulse, which are 14 pC, 6 pC and 87 pC, respectively. Similar enhancement levels were 97 98 observed in over 10 individual devices produced in separate fabrication runs, with 99 representative data reported in the Supporting Information.

100 To quantify the enhanced conduction, we initially assume the properties of the material 101 remain unchanged during the coincident pulse. The current density generated by a voltage pulse 102 is determined by the density of electron (holes) n (p), their respective mobilities μ_n and μ_p , 103 and the local electric field E, which scales as V/L, where V is the applied voltage and L the inter-104 electrode separation. During the coincident pulse the increase in current may be expressed as a 105 sum of the current generated during a voltage pulse ΔI_{el} and an optical pulse ΔI_{pc} as:

106

107
$$\Delta I_{co}(V \to V + \Delta V, optical) = \Delta I_{el}(V \to V + \Delta V) + \Delta I_{pc}(V + \Delta V)^{\alpha}$$
 (1)

108

109 where alpha α is a scalar describing the degree of current nonlinearity during the coincident 110 pulse. (See Supporting Information for a detailed derivation.) The non-linearity in the I-V 111 response is expressed as $E \propto V^{\alpha}$. Assuming fixed geometry and no change in material 112 properties the coincident pulse response should be the sum of the voltage and the optical pulse 113 responses in Figure 2b, with the latter rescaled by the factor $I_{el}(V + \Delta V)/I_{el}(V)$, which is 114 estimated to be a factor of ~2 on average for the devices measured based on the data in Figure 115 1b (inset and caption). Applying Equation (1) to the data in Figure 2b, *i.e.*, 16 pA \neq (2.5 + 1 x

116 2) pA, demonstrates that there is a nonlinear increase in current and total charge (87 pC \neq (14 117 + 6 x 2) pC) during coincident pulses. This enhancement is observed over many devices and is 118 likely due to a modification of the material properties, as discussed below.

119 We now show that the degree of modification depends on the number and time correlation of 120 the pulse stimuli. Figure 3a shows an alternating train of voltage and optical pulses with a time 121 delay $\Delta t = 20$ s. Although each pulse is identical to that employed in Figure 2c, there is a weak 122 interaction between the pulses that produces an enhancement in the current during a steady train 123 of voltage-only pulses with the same time spacing. The enhancement factor is 1.5-2 when $\Delta t =$ 124 20 s. When $\Delta t = 5$ s, this factor increases to ~5 and the enhancement is observed to gradually 125 increase as the current response reaches a steady state value of 5.5 pA (Figure 3b). A train of 126 coincident pulses in Figure 3c produce the largest enhancements (> 20), which remains for a 127 certain time beyond the pulse duration directly analogous to synaptic weight change in 128 neuromorphic systems. Figure 3d is a plot of the current enhancements as a function of pulse 129 time separation, in which the curve has been mirrored through $\Delta t = 0$ to reflect the fact that the 130 enhancement phenomenon is not dependent on the order in which the stimuli are delivered. 131 This sensitivity to time correlation, even for heterogeneous stimuli, is strikingly similar to the STDP phenomenon found in biological systems,^[19] and that reported earlier for homogeneous 132 voltage impulses in memristors.^[8, 20] In contrast to biological systems, however, we do not 133 134 observe inhibition when the order of the stimuli is reversed

We now explore the response of the device to probe pulses (optical and voltage) after the device has experienced different numbers of coincident pulses. **Figure 4** shows a pulse train comprised of interleaved voltage (red) and optical (blue) pulses that allows for the direct overlap of voltage and optical stimulus during every third pulse (orange). The initial individual voltage and optical pulses produce a very modest response. The response to the first coincident pulse is strong, but significantly the response to subsequent voltage and optical probe pulses (i.e. after the coincident pulse) are also strengthened, similar to neural short term plasticity.^[21] Crucially,

each additional coincident pulse produces a growing response and is followed by an enhanced response to the subsequent voltage and optical probe pulses. These data are consistent with a learning behavior in which the device output grows following each coincident pulse. The enhanced response to the subsequent voltage and optical probe pulses, the levels of which are shown at the bottom of Figure 4a, demonstrates short-term associative memory – an effect observed in all fabricated devices.

148 **3. Mechanism**

The time correlated enhancement in the device performance can be understood from the known properties of TiO₂ and its interaction with light.^[22, 23] We begin by considering the band alignment across the single nanowire device structure as shown in Figure 1a. The Au contact leads have a work function of 5.1 eV and form Schottky barrier contacts with the wide bandgap TiO₂ nanowire, where the conduction band bends upwards due to the formation of a depletion region as shown in Figure 4b. This Schottky barrier height can be further modified by oxygen vacancy generation and adsorption of O₂ onto the wire.

During device activation, oxygen vacancies readily form at the TiO₂ interface with the Au metal contact under positive bias.^[24-28] The anodic conditions cause lattice oxygen O_0^{\times} to react, forming positively charged oxygen vacancies $V_0^{\bullet\bullet}$ and releasing electrons into the conduction band:^[29]

160

161
$$O_0^{\times} \rightleftharpoons V_0^{\bullet \bullet} + 2e' + \frac{1}{2}O_2(g)$$
 Reaction (1)

162

163 The additional screening produced by these electrons reduces the width of the depletion 164 region and increases the transmission through the contact, resulting in activation of the 165 nanowire. However, the presence of O_2 in the ambient and generated at the Au contact creates 166 an additional Schottky barrier due to reaction with electrons in the conduction band that 167 results in the adsorption of O_2^- species.^[30-33]

168

169 $e' + O_2(g) \rightarrow O_2^-(ads)$ Reaction (2) 170 This reaction occurs in wire regions with large populations of free electrons and hence 171 predominantly at the Au-TiO₂ interface. Charge depletion due to these adsorbed O₂⁻ species 172 increases the Schottky barrier at this interface. In the presence of light, however, 173 photogenerated holes are driven by the upward band bending towards the interface where they 174 are captured by the adsorbed O_2^{-} species, resulting in the liberation of neutral O_2 . Based on our device structure (see Figure 1) it is likely that most of the liberated O₂ via reaction (1) 175 176 remains trapped at the Au-TiO₂ interface where it can in turn recombine with photogenerated 177 electrons via reaction (2). Thus the electrode-nanowire interface is an effective region for 178 recombination of photogenerated electron-hole pairs. This can be iunderstood by considering 179 the Schottky barrier hieight and width under diufferent bias conditions. For the forward-180 biased Schottky barrier configuration considered here, the effective barrier height and the depletion width decrease as the voltage is increased.^[34] We therefore suggest that the large 181 182 band bending and broad depletion width present under low bias favours the recombination of 183 photogenerated carriers by sweeping holes towards the surface to react with O2- produced by 184 the adsorption of O2 with photogenerated electrons. This recombination process is suppressed 185 at higher applied biases and hence is responsible for the short term enhancement in 186 conductivity to coincident voltage and optical stimuli, which was observed in all devices. 187 Repeated coincident stimuli result in a steady build-up in the population of photogenerated 188 carriers resulting in the learning response described in Figure 4a. 189 The role of adsorbed oxygen at the metal/nanowire interface for the nonlinear conduction 190 enhancement mechanism and persistence with time can also be explored by the introduction 191 of a device capping layer. Repeating the same experiment as described in Figure 2b on a 192 device capped with spincast SU-8 photoresist results in the data shown in Figure 5. Current levels are increased in response to all pulse stimuli (voltage: $\Delta I_{el} = 2.75$ pA, $\Delta Q_{el} = 16$ pC; 193

optical: $\Delta I_{pc} = 5.32$ pA, $\Delta Q_{pc} = 25$ pC; coincident: $\Delta I_{Co} = 65$ pA, $\Delta Q_{Co} = 330$ pC) but notably 194 195 the current response to the optical stimulus is considerably larger and longer lived when 196 compared to the actual pulse width. The pulse width is designated by the pale purple strip that 197 defines the actual time duration of the optical pulse. Isolating the device from atmospheric 198 oxygen improves the lifetime of the associative memory effect, enabling longer term 199 potentiation. The lifetime of the effect also depends on current magnitude and recombination 200 dynamics, as seen in Figures S2 and S3, and can be further extended via operation in vacuum 201 (Figure S4).

202 Mechanistically, the capping layer limits the supply of atmospheric oxygen from the ambient 203 and hence limits the recombination of the carriers generated during a coincident pulse. 204 Oxygen liberated from the material during the pulse remains trapped however and free to re-205 adsorb on the surface. In contrast, devices without capping layers but measured under vacuum 206 conditions have little adsorbed oxygen and no capacity to replenish adsorbed oxygen from the 207 ambient, and in each case show a dramatically increased decay times for the photogenerated 208 current. The capping and vacuum experiments clearly show that the presence and availability 209 of oxygen dramatically affects the response of these devices. These oxygen mediated surface 210 reactions change the material functionality for a limit time period, enhancing the associative 211 memory response. These data strongly suggest that by controlling the surface environment 212 and/or functionalization it may be possible to improve neuromorphic performance.

213

214

215 **4.** Conclusion

216

We have introduced for the first time a single TiO_2 nanowire device that is capable of processing heterogeneous physical stimuli into an associative memory response. The device also demonstrates time correlation that mimics the STDP learning processes found in biological systems. Although the memory is volatile in the present configuration, modifying surface

environment is shown to extend decay lifetimes. The enhancement effect due to the time 221 222 correlation of the voltage and optical stimuli is robust and reproduced over several devices with 223 the single-nanowire structure and is amenable to extending the number of inputs or feed 224 forwarding the output of one device into another. The non-linear enhancement is postulated to 225 arise from the increased concentration of photogenerated carriers during the voltage pulse and 226 the mediating role adsorbed oxygen species have on the effective Schottky barrier height and 227 carrier lifetime, These observations are supported by data taken on devices with a capping layer 228 and those measured in vacuum. Given the simplicity of the device structure and the underlying 229 nature of the Schottky barrier properties that control its operation, it will likely be possible to 230 extend the range of physical stimuli (such as pressure, temperature, and magnetic field) that can 231 be detected with temporal correlation. This is the first demonstration of a material which can 232 process time-correlated heterogeneous inputs for neuromorphic applications, and the 233 functionality developed here will likely prove useful in the development of neuromorphic 234 sensing platforms.

- 235
- 236

237 **5. Experimental Section**

238 TiO₂ nanowire devices were fabricated as reported previously using UV lithography, spray deposition and e-beam lithography.^[18] UV lithography is used to define large area contact 239 240 pads before spray deposition of dilute solutions of TiO₂ in deionized water. TiO₂ wires 241 (EMFUTUR) 50-100 nm in diameter and 5-20 µm in length were used to prepare devices. 80 242 nm of Au metal was deposited via electron beam evaporation to form contacts to the 243 individual wires, with no adhesion layer used to ensure direct Au-wire contact. A Thor Labs 244 SC10 shutter controller and SH1 beam shutter was used to modulate a 405 nm wavelength, 245 4.5 mW collimated laser diode that provides 0.5 mW of power at the wire. Voltage pulse and optical shutter sequences were controlled using a Keithley 4200-SCS parameter analyser, 246 which also performs the current measurement. Vacuum measurements (10^{-6} torr) were 247

248	performed in an ARS DE-204NI cryostat at room temperature. Wire devices are coated in Su-		
249	8 photoresist (Microchem) to form a UV transparent capping layer to isolate the device from		
250	atmosphere. The SU-8 is dropped on the device, spun at low rpm (2000) and baked at 150°C		
251	for 5 minutes to cure the resist.		
252 253 254 255	Supporting Information Supporting Information is available from the Wiley Online Library or from the author.		
256 257	Acknowledgements C.J.O'K. designed and executed the experiments with the help of J.F. and J.B. C.J.O'K. and		
258	J.F. analyzed the data and co-wrote the manuscript. D.McC. contributed to the development		
259	and optimization of the experimental setup as well as providing insight into data modeling		
260	and analysis. J.D. provided fruitful discussion. J.B. led the project, oversaw its development		
261	and contributed to the manuscript. The authors wish to acknowledge funding from the		
262	European Research Council under Advanced Grant 321160. This publication has emanated		
263	from research supported in part by a research grant from Science Foundation Ireland (SFI)		
264	AMBER Centre under Grant Number SFI/12/RC/2278.		
265 266 267 268 269	Received: ((will be filled in by the editorial staff)) Revised: ((will be filled in by the editorial staff)) Published online: ((will be filled in by the editorial staff))		
270	References		
271 272 273 274 275 276 277 278	 W. Gerstner, R. Kempter, J. L. van Hemmen, H. Wagner, Nature 1996, 383, 76. Gq. Bi, Mm. Poo, Annual Review of Neuroscience 2001, 24, 139. J. Costas-Santos, T. Serrano-Gotarredona, R. Serrano-Gotarredona, B. Linares- Barranco, Circuits and Systems I: Regular Papers, IEEE Transactions on 2007, 54, 1444. W. Liu, A. G. Andreou, M. H. Goldstein, Jr., Neural Networks, IEEE Transactions on 1992, 3, 477. G. Indiveri, B. Linares-Barranco, T. J. Hamilton, A. van Schaik, R. Etienne-Cummings, T. Delbruck, SC. Liu, P. Dudek, P. Häfliger, S. Renaud, J. Schemmel, G. Cauwenberghs, J. 		

Arthur, K. Hynna, F. Folowosele, S. SAÏGHI, T. Serrano-Gotarredona, J. Wijekoon, Y. Wang,
K. Boahen, Frontiers in Neuroscience 2011, 5.

- [6] B. V. Benjamin, G. Peiran, E. McQuinn, S. Choudhary, A. R. Chandrasekaran, J. M.
 Bussat, R. Alvarez-Icaza, J. V. Arthur, P. A. Merolla, K. Boahen, Proceedings of the IEEE 2014,
 102, 699.
- 284 [7] P. A. Merolla, J. V. Arthur, R. Alvarez-Icaza, A. S. Cassidy, J. Sawada, F. Akopyan, B.
- L. Jackson, N. Imam, C. Guo, Y. Nakamura, B. Brezzo, I. Vo, S. K. Esser, R. Appuswamy, B.
- Taba, A. Amir, M. D. Flickner, W. P. Risk, R. Manohar, D. S. Modha, Science 2014, 345, 668.
- 287 [8] S. H. Jo, T. Chang, I. Ebong, B. B. Bhadviya, P. Mazumder, W. Lu, Nano Letters 2010, 288 10, 1297.
- 289 [9] B. Linares-Barranco, T. Serrano-Gotarredona, L. A. Camuñas-Mesa, J. A. Perez-290 Carrasco, C. Zamarreño-Ramos, T. Masquelier, Frontiers in Neuroscience 2011, 5.
- [10] T. Hasegawa, T. Ohno, K. Terabe, T. Tsuruoka, T. Nakayama, J. K. Gimzewski, M.
 Aono, Advanced Materials 2010, 22, 1831.
- [11] S. Ambrogio, S. Balatti, F. Nardi, S. Facchinetti, D. Ielmini, Nanotechnology 2013, 24,
 384012.
- 295 [12] S. La Barbera, D. Vuillaume, F. Alibart, ACS Nano 2015, 9, 941.
- 296 [13] S. Kyungah, K. Insung, J. Seungjae, J. Minseok, P. Sangsu, P. Jubong, S. Jungho, P. B.
- Kuyyadi, K. Jaemin, L. Kwanghee, L. Byounghun, H. Hyunsang, Nanotechnology 2011, 22,
 254023.
- 299 [14] M. D. Pickett, G. Medeiros-Ribeiro, R. S. Williams, Nat Mater 2013, 12, 114.
- 300 [15] Y. V. Pershin, M. D. Ventra, Neural Netw. 2010, 23, 881.
- 301 [16] C. T. Phare, Y.-H. Daniel Lee, J. Cardenas, M. Lipson, Nat Photon 2015, 9, 511.
- 302 [17] G. T. Reed, G. Mashanovich, F. Y. Gardes, D. J. Thomson, Nat Photon 2010, 4, 518.
- 303 [18] C. O'Kelly, J. A. Fairfield, J. J. Boland, ACS Nano 2014, 8, 11724.
- 304 [19] D. O. Hebb, *The Organization of Behavior: A Neuropsychological Theory*, Taylor &
 305 Francis Group, 2012.
- 306 [20] Y. V. Pershin, M. Di Ventra, Neural Networks 2010, 23, 881.
- 307 [21] T. Ohno, T. Hasegawa, T. Tsuruoka, K. Terabe, J. K. Gimzewski, M. Aono, Nat Mater308 2011, 10, 591.
- 309 [22] M. A. Henderson, Surface Science Reports 2011, 66, 185.
- 310 [23] G. Liu, N. Hoivik, X. Wang, S. Lu, K. Wang, H. Jakobsen, Electrochimica Acta 2013,311 93, 80.
- 312 [24] D. S. Jeong, H. Schroeder, U. Breuer, R. Waser, Journal of Applied Physics 2008, 104,
 313 123716.
- J. Joshua Yang, F. Miao, M. D. Pickett, D. A. Ohlberg, D. R. Stewart, C. N. Lau, R. S.
 Williams, Nanotechnology 2009, 20, 215201.
- 316 [26] D. H. Kwon, K. M. Kim, J. H. Jang, J. M. Jeon, M. H. Lee, G. H. Kim, X. S. Li, G. S.
- 317 Park, B. Lee, S. Han, M. Kim, C. S. Hwang, Nature nanotechnology 2010, 5, 148.
- 318 [27] U. Diebold, Surf. Sci. Rep. 2003, 48, 53.
- 319 [28] N. Kazuki, Y. Takeshi, K. Masaki, O. Keisuke, K. Annop, R. Sakon, M. Gang, H. Mati,
- 320 X. Bo, Z. Fuwei, H. Yong, K. Tomoji, Jpn. J. Appl. Phys. 2012, 51, 11PE09.
- 321 [29] D. S. Jeong, H. Schroeder, R. Waser, Phys. Rev. B 2009, 79, 195317.
- [30] M. S. T. Berger, O. Diwald, E. Knozinger, D. Panayotov, T. L. Thompson, J. T. Yates,
 Journal of Physical Chemistry B 2005, 6061.
- J. Zou, Q. Zhang, K. Huang, N. Marzari, The Journal of Physical Chemistry C 2010,
 114, 10725.
- 326 [32] M. D. Rasmussen, L. M. Molina, B. Hammer, The Journal of Chemical Physics 2004,327 120, 988.
- M. A. Henderson, W. S. Epling, C. L. Perkins, C. H. F. Peden, U. Diebold, The Journal
 of Physical Chemistry B 1999, 103, 5328.
- 330 [34] S. M. Sze, K. K. Ng, *Physics of Semiconductor Devices*, Wiley, 2006.
- 331



333





Figure 1. (a) An SEM image of a TiO₂ nanowire (NW) with Au contacts (scale bar 1 µm), 336 and a schematic of the experimental setup depicting a laser with a shutter linked to a controller and source measure unit. (b) Half wave sinusoidal voltage input produces an 337 338 evolving current as a function of the number of input half waves. (c) The increase in conductance at one polarity may be selectively decreased or reset using opposite polarity 339 340 inputs. The onset for current rectification occurs above ± 3 V, as shown in inset (c), the dotted 341 red lines are used to obtain the current at 2 and 2.75 V during sweep 1.

- 342
- 343
- 344 345





347 Figure 2. (a) Photocurrent produced in the nanowire device in response to a 20 s laser input at 348 various holding voltage levels, inset shows the linear increase in photocurrent ($\pm 20\%$ absolute value) up to 5 V with the photocurrent saturating at higher voltages. (b) Current 349 response from separate heterogeneous voltage and optical pulses followed by a pulse 350 351 consisting of both the optical and voltage pulse coincident on the nanowire simultaneously (2 352 V hold, pulse $\Delta 0.75$). Each pulse is separated by a long dwell time at a hold voltage of 2 V to eliminate temporal coupling. The magnitude of current ($\Delta I_{co} = 16 \text{ pA}$) and charge ($\Delta Q_{co} = 87$ 353 pC) produced during the coincident pulse is not simply the linear addition of the current or 354 charge produced by the individual voltage ($\Delta I_{el} = 2.5 \text{ pA}, \Delta Q_{el} = 14 \text{ pC}$) and optical pulses 355 356 $(\Delta I_{pc} = 1 \text{ pA}, \Delta Q_{pc} = 6 \text{ pC}).$ 357

- 358 359



360 361

Figure 3. Pulse timing is critical to the interaction of separate pulse stimuli. For voltage and optical pulses separated by (a) 20 seconds, (b) 5 seconds, and (c) 0 seconds, the device conductance is augmented with a maximum augmentation for coincident pulses (2 V holding bias, pulse $\Delta 0.75$). (d) A pseudo-plasticity strengthening graph analogous to STDP displaying the enhancement of device conductance as the time separation between pulses decreases.





Figure 4. (a) A plot showing the device response to a series of coincident pulses that results in a growing and heightened current output, consistent with a learning phenomenon (2 V holding bias, pulse $\Delta 0.75$). The lower panel shows the enhanced response to voltage (red) and optical (blue) probe pulses following each coincident pulse. The enhancement level

- 375 increases with the number of coincident pulses and demonstrates an associative memory
- 376 response, which is not observed following uncorrelated voltage and optical pulses. (b)
- 377 Schematic energy band diagram of the device interface at (i) 2 V hold condition and (ii)
- 378 during excitation with coincident stimuli comprised of a voltage (ΔV) and optical ($\hbar\omega$) pulse.
- 379 Oxygen adsorbs at the TiO₂/metal interface, scavenging conduction band electrons and
- 380 augmenting the barrier height. Under the coincident pulse stimulus photogenerated holes
- 381 combine with and release the adsorbed oxygen (see text) and when coupled with the increased
- 382 voltage enhances charge injection into the electrode.
- 383





386 Figure 5. Isolating the device from atmospheric oxygen via a polymer capping layer has a profound impact on photogenerated current lifetime and the overall current levels produced 387 388 by the device relative to the uncapped device, increasing the long term memory effect in 389 response to pulse stimuli (2 V holding bias, pulse $\Delta 0.75$).



390 Table of Contents Entry:

391

392 A single TiO2 nanowire device demonstrates associative memory between two different

393 pulse stimuli. Voltage pulses and optical UV laser light pulses both stimulate higher device 394 current, but when the two pulses temporally coincide a nonlinear current enhancement occurs, 395 similar to spike-timing dependent plasticity and caused by oxygen reactions within the 396 dominant of the device

- 396 depletion region of the device.
- 397

Keywords: Memory, Learning, Associative Memory, Neuromorphic hardware, Pulse
 Stimuli, Nanowire

- 400
- 401 Curtis J. O'Kelly1,3, Jessamyn A. Fairfield1,3, David McCloskey2,3, John F. Donegan2,3
- 402 and John J. Boland1,3*
- 403
- 404 Associative Enhancement of Time Correlated Response to Heterogeneous Stimuli in a
- 405 Neuromorphic Nanowire Device



- 408 Copyright WILEY-VCH Verlag GmbH & Co. KGaA, 69469 Weinheim, Germany, 2013. 409 410 **Supporting Information** 411 412 413 Associative Enhancement of Time Correlated Response to Heterogeneous Stimuli in a 414 **Neuromorphic Nanowire Device** 415 Curtis J. O'Kelly^{1,3}, Jessamyn A. Fairfield^{1,3}, David McCloskey^{2,3}, John F. Donegan^{2,3} and 416 John J. Boland^{1,3}* 417 418
- 419 S1. Additional Voltage and Optical Enhancement Data



Figure S1. Coincident pulses increase current in all similar devices in a non linear addition of the separate voltage and optical pulses. Three separate devices are shown to display the same reproducible behavior in response to coincident pulses composed of voltage and optical pulses applied at the same time.

425

426

427 **Table S1.** The relative increase in current produced during separate voltage pulses and optical428 pulses are shown below.

	Voltage Scaling	Optical Scaling
Device 1	13.96	12.47
Device 2	9.21	6.74
Device 3	7.84	4.93

429 430

430

431 S2. Non-Linear Enhancement Derivation

432 We initially assume the properties of the system remain unchanged during the coincident pulse, that current generated under the voltage pulse is due to the change in local electrical field 433 434 and the photocurrent generated is from the separation of electrons and holes in the depletion 435 region. The current density generated by a voltage pulse is determined by the density of electron (holes) n(p), their respective mobilities μ_n and μ_p , and the local electric field E, which scales 436 as V/L, where V is the applied voltage and L the inter-electrode separation. For our fixed device 437 438 geometry we can account for non-linearity in the current-voltage response by expressing $E \propto$ 439 V^{α} , where α is a scalar describing the degree of nonlinearity. The geometry-normalized current 440 increase (ΔI_{el}) (units: ampere per length) generated by a voltage pulse ΔV can be written as:

441
$$\Delta I_{el}(V \to V + \Delta V) = e \left[n \,\mu_n + p \,\mu_p \right] \cdot \left[(V + \Delta V)^{\alpha} - V^{\alpha} \right]$$
(S1)

442 The current response to an optical pulse (ΔI_{pc}) at a fixed voltage *V* results in the generation of 443 additional electrons (Δn) and holes (Δp) that increase the photocurrent by:

444
$$\Delta I_{pc}(V) = e \left[\Delta n \,\mu_n + \Delta p \,\mu_p \right] V^{\alpha} \tag{S2}$$

445 During the coincident pulse, these photogenerated carriers experience an increased electric field 446 due to the simultaneous voltage pulse. The increase in current can then be written as the sum of 447 the equations S1 and S2 at the voltage pulse ΔV :

448
$$\Delta I_{co}(V \to V + \Delta V, opt) = \Delta I_{el}(V \to V + \Delta V) + \Delta I_{pc} (V + \Delta V)^{\alpha}$$
(S3)

449 Therefore, assuming no change in material properties, the coincident pulse response should be 450 the sum of the measured independent voltage and the optical pulse responses shown in Figure 451 2b of the main text. The factor $I_{el}(V + \Delta V)/I_{el}(V)$ is scaled to the voltage pulse ΔV as the 452 additional photocurrent conventionally increases at higher voltage. However even when the 453 photocurrent is scaled to the voltage pulse level as estimated to be a factor of ~ 2 based on the 454 data in Figure 1b(inset), there remains a large difference in the measured and calculated current. Equation (3) yields $2.5 + 1 \ge 4.5 \Delta I_{co}$ which is considerably less than the measured $\Delta I_{co} =$ 455 456 16 pA. In the main text the difference is also presented as charge or the time integral of the 457 current measured during each pulse to demonstrate the sharp current rise is not a short but large 458 peak. This derivation and numerical calculation supports nonlinear current enhancement in the 459 material under coincident pulse stimuli.

460

461 S3. Persistence of Associative Memory

462

The associative memory enhancement effect demonstrated in the present device is volatile. 463 464 After multiple coincident pulses normal current level responses to voltage pulse stimuli are 465 recovered after a period of 300-400 s. This is demonstrated in Figure S1, which shows a device 466 with 5-20 coincident pulses applied to it prior to holding a bias of 2.75 V for an extended period. 467 The 2.75 V bias (red trace) is applied to the device to reveal the fall in current levels more 468 clearly over time. In all cases it takes approximately 300-400 s after the training coincident 469 pulses are applied before the current levels reach normal levels similar to those without 470 coincident pulses applied. The measured half-life (time taken for the enhanced current 471 magnitude to drop by half the its initial value following the coincident pulses) of the memory 472 association effect is indicated for each data and reveal that the memory persistence decreases 473 as the current magnitude increases. This volatility stems from the recombination of atmospheric

474 oxygen and charge carriers generated during the coincident pulse stimulus. At greater current
475 magnitude and therefore higher charge concentration at the device interface the rate of
476 recombination is accelerated.



Figure S2. Memory association enhancement persistence as a function of the number of applied
coincident pulses. Current enhancement persists longer with lower current magnitudes and
decays faster as the current generated from greater number of coincident pulses increases as
shown in the decrease in half life of the persistent current. THE NUMBER OF PULSES IS
HARD TO NOTICE AT THE TOP OF THESE FIGS





Figure S3. (a) A plot showing the current decay following 5 to 20 coincident pulses. (b) Currents normalised and plotted on a log scale, the divergence from a straight line single exponential decay becomes more apparent for larger numbers of applied coincident pulses.

492

493 The decay in persistent current levels does not fit a simple exponential decay for standard 494 persistent photocurrent recombination. In Figure S2 (a) the decay in the current after the 495 coincident pulses are applied are mapped onto a single plot. The rapid decrease in current at 496 higher current levels seems to support a more complicated bi-exponential decay with fast and 497 slow time constants. In Figure S2 (b) the current decay is normalized and plotted on a 498 logarithmic scale. If the decay was purely exponential these data would follow a straight line, 499 however we observe a clear shift in linearity as the number of coincident pulses applied to the 500 device increases.



- 501 502
- 502 503
- 504
- 505

Figure S4. An optical pulse applied while the device is under vacuum (10⁻⁶ torr) results in a much longer decay than the same measurement under atmosphere (see Figure S1), again implicating surface oxygen reactions. The holding bias voltage was 5 V which accounts for the larger photocurrent than was observed at a 2 V holding bias. NO data on simultaneous pulsing in vacuum???

511