ASYMPTOTIC BEHAVIOR OF SOLUTIONS OF REACTION-DIFFUSION SYSTEMS OF LOTKA-VOLTERRA TYPE

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Dedicated to the memory of Professor P. Hess

1. Introduction. As a mathematical model for the population dynamics of *N*-species in biology, Lotka [12] and Volterra [17] proposed the ordinary differential system of the form:

$$dv_j/dt = (-e_j + b_j^{-1} \sum_{k=1}^{N} a_{jk} v_k) v_j, \quad j = 1, \dots, N,$$
 (LV)

where e_j , $b_j(>0)$, a_{jk} are given constants; and v_j denotes the biomass of the j-species; and investigated the asymptotic behavior of v_1, \ldots, v_N for large time t.

For N=2, there are extensive literatures on (LV) (or (RD) below), e.g., Copell [5], Henry [7], Rothe [16]. However, for $N \ge 3$, little seems to have been known; see Amann [2, 3], Krikorian [11], Fife-Mimura [6], Friedmann-Tzavars [8], Oshime [14] and others.

In the present paper we consider the reaction-diffusion's version of (LV) of the form:

$$\frac{\partial}{\partial t}u_{j} = d_{j}\Delta u_{j} + u_{j}f_{j}(u) \quad (x \in \Omega, \ t > 0)$$

$$\frac{\partial}{\partial \nu}u_{j}\Big|_{\partial\Omega} = 0, \quad (t > 0); \quad u_{j}\Big|_{t=0} = \phi_{j} \quad (j = 1, ..., N),$$
(RD)

where Ω is a bounded domain in \mathbb{R}^n with smooth boundary $\partial\Omega$, d_j is a positive constant, $\partial/\partial\nu$ denotes the outer normal derivative to $\partial\Omega$, and ϕ_j given smooth nonnegative, and not identically zero function satisfying the compatibility condition: $\partial\phi_j/\partial\nu=0$ on $\partial\Omega$. The purpose of the present paper is to study the asymptotic behavior of solutions of (RD) for large t under some assumptions on f_i .

We suppose that f_j , j = 1, ..., N, satisfies the following assumptions.

Received September 1993.

AMS Subject Classification: 35K57.

Assumption 1.1. $f_j(\xi)$ is a smooth function of $\xi \in \mathbb{R}^N_+$ such that there is a positive constant b_j with

$$\sum_{j=1}^{N} b_j \xi_j f_j(\xi) \le 0, \quad \xi = (\xi_1, \dots, \xi_N) \in \mathbb{R}_+^N, \tag{1}$$

$$(\mathbb{R}_+^N = \{ \xi \in \mathbb{R}^N; \ \xi_j \ge 0, \ (j = 1, \dots, N) \}).$$

To state the theorem, we decompose the index set $\Lambda \equiv \{1, ..., N\}$ into three disjoint sets $\Lambda_1, \Lambda_2, \Lambda_3 : \Lambda = \Lambda_1 \cup \Lambda_2 \cup \Lambda_3$ (disjoint). We define the set $\Lambda_{1,s}$ inductively. $\Lambda_{1,1}$ is the set of all j such that

$$\sum_{k=1}^{N} b_k \xi_k f_k(\xi) \le -\delta_1 \xi_j, \quad (\xi \in \mathbb{R}_+^N)$$
 (2)

for some positive constant δ_1 . If $\Lambda_{1,s}$ is defined, then by definition $j \in \Lambda_{1,s+1}$ if and only if either $j \in \Lambda_{1,s}$ or there is an i $(1 \le i \le N)$ such that

$$f_i(\xi) \ge -\delta_2 \xi_j - \rho(|\xi|) \sum_{k \in \Lambda_1} \xi_k, \ (\xi \in \mathbb{R}^N_+), \tag{3}$$

where δ_2 is some positive constant, and $\rho(s)$ some positive increasing function of s: if an index set Λ_0 is empty, we understand $\sum_{j \in \Lambda_0} a_j = 0$. We set

$$\Lambda_1 = \bigcup_s \Lambda_{1,s}$$

 Λ_2 is the set of all $j \in \Lambda \setminus \Lambda_1$ such that

$$|f_j(\xi)| \le \rho(|\xi|) \sum_{k \in \Lambda_1} \xi_k, \quad (\xi \in \mathbb{R}^N_+).$$
 (4)

Here and in what follows, $\rho(s)$ denotes a positive increasing function of s. Finally we set

$$\Lambda_3 = \Lambda - \Lambda_1 - \Lambda_2. \tag{5}$$

Let C be the set of all positive vectors $c = (c_i)_{i \in \Lambda_3}$ (Λ_3 is naturally ordered) such that

$$\left| \sum_{i \in \Lambda_2} c_i f_i(\xi) \right| \le \rho(|\xi|) \sum_{i \in \Lambda_1} \xi_i, \quad (\xi \in \mathbb{R}^N_+). \tag{6}$$

Assumption 1.2. If Λ_3 is non-empty, there is a vector in C.

We denote by r the number of the linearly independent vectors in C, and independent vectors by $c^{(j)} = (c_i^{(j)})_{\Lambda_3}$, $j = 1, \ldots, r$, where a positive vector c means that all the component c_i is positive.

For vector $\gamma = (\gamma_k)_{k \in \Lambda_2} \ge 0$, non-negative real μ_0 , and vector $\mu = (\mu_j)_{j=0}^r$, we define $\Theta(\gamma, \mu)$ by the set of all $\xi \in \mathbb{R}_+^N$ such that

$$\xi_{j} = 0 \ (j \in \Lambda_{1}); \ \xi_{j} = \gamma_{j} \ (\text{for } j \in \Lambda_{2});$$

$$\sum_{j=1}^{N} b_{j} \xi_{j} = \mu_{0}; \quad \text{and} \quad \sum_{k \in \Lambda_{3}} c_{k}^{(j)} \log \xi_{k} = \mu_{j}, \ (j = 1, \dots, r),$$
(7)

where we understand that if Λ_3 is empty set, then $\mu = \mu_0$.

Remark 1.1. Suppose that $\Lambda_3 = \Lambda$, and r = N - 2. Suppose also that

$$\sum_{j=1}^{N} b_j \xi_j f_j(\xi) = 0, \tag{1'}$$

$$\sum_{i=1}^{N} c_j^{(k)} f_j(\xi) = 0, \quad (k = 1, \dots, N-2).$$
 (2')

Then the set

$$\Theta = \{ \xi \in \mathbb{R}_+^N; \sum_{j=1}^N b_j \xi_j = \mu_0, \sum_{j \in \Lambda_3} c_j^{(k)} \log \xi_j = \mu_j, \ (1 \le k \le r) \}$$

is a bounded closed orbit, and so it is a periodic orbit for (RD).

Then our main theorem reads as follows:

Theorem 1.1. Let the above assumptions 1.1 and 1.2 hold. Then any bounded solution u of (RD) converges to a $\Theta(\gamma, \mu)$ as $t \to \infty$, uniformly on $\bar{\Omega}$ for some $\gamma = (\gamma_j)_{j \in \Lambda_2}$ and $\mu = (\mu_j)_{j=0}^r$.

$$\operatorname{dist}(u(x,t),\theta(\gamma,\mu))\to 0$$
 as $t\to\infty$, uniformly on $\bar{\Omega}$,

where

$$\gamma_j > 0$$
 (if Λ_2 is not empty);
 $\mu_0 > 0$ (if $\Lambda_2 \cap \Lambda_3$ is not empty);
 $\mu_j = 0$ (if Λ_3 is empty).

Corollary 1.1. In addition to the assumptions 1.1 and 1.2, assume that $\Lambda_3 = \Lambda$, and r = N - 2, and that (1'), and (2') hold. Then any bounded solution u converges to a periodic orbit for (RD).

The proof is a direct consequence of Theorem 1.1 and Remark 1.1.

We now consider the reaction-diffusion system of the form considered originally by Lotka and Volterra, which is a special case of (RD):

$$\frac{\partial}{\partial t}u_{j} = d_{j}\Delta u_{j} + (-e_{j} + b_{j}^{-1} \sum_{k=1}^{N} a_{jk}u_{k})u_{j}, \quad (x \in \Omega, \ t > 0)$$

$$\frac{\partial}{\partial \nu}u_{j} = 0, \quad (x \in \partial\Omega, \ t > 0); \quad u_{j} = \phi_{j}, \quad (x \in \Omega, \ t = 0).$$
(8)

In this case it is easy to see that assumption 1.1 in Theorem 1.1 implies that $e_j \ge 0$, (j = 1, ..., N) and that the matrix $(a_{jk} + a_{kj})$ is non-positive definite. Let us see the construction of Λ_1 , Λ_2 , Λ_3 more concretely. $j \in \Lambda_{1,1}$ if

$$\sum_{i=1}^{N} b_i (-e_i + b_i^{-1} \sum_{k=1}^{N} a_{ik} \xi_k) \xi_i \le -\delta_1 \xi_j, \quad (\xi \in \mathbb{R}_+^N)$$
 (9)

with some $\delta_1 > 0$. The above condition is equivalent to the one that $e_j > 0$ and the matrix $(a_{jk} + a_{kj})$ is non-positive definite.

If $\Lambda_{1,s}$ is constructed, then $\Lambda_{1,s+1} = \Lambda_{1,s} \cup \Lambda'_{1,s+1}$. Here $j \in \Lambda'_{1,s+1}$ if and only if there is an $i \in \Lambda$ with $a_{ik} \geq 0$ (for $k \in \Lambda - \Lambda_{1,s}$) and with $a_{ij} > 0$. Then we set $\Lambda_1 = \bigcup_s \Lambda_{1,s}$ and $\Lambda_3 = \Lambda - \Lambda_1 - \Lambda_2$. Suppose that there is a positive vector $c = (c_j)_{j \in \Lambda_3}$ with

$$\sum_{i \in \Lambda_3} c_i b_i^{-1} a_{ik} = 0, \quad (k \in \Lambda_2 \cup \Lambda_3). \tag{10}$$

Then the number of the linearly independent vectors in C is the dimension of the kernel of the matrix $(a_{jk})_{j\in\Lambda_3,k\in\Lambda_2\cup\Lambda_3}^*$, (* denotes adjoint matrix). We can now state Theorem 1.1 in a more concrete form.

Theorem 1.2. Assume that $e_j \geq 0$, (j = 1, 2, ..., N) and that the matrix $(a_{jk} + a_{kj})$ is non-positive definite. Assume also that if Λ_3 is non-empty, then there is a c in C satisfying (10). Then any bounded solution of (8) converges to some $\theta(\gamma, \mu)$, uniformly on $\bar{\Omega}$ as $t \to \infty$, uniformly on $\bar{\Omega}$.

Remark 1.2. Let the assumptions in Theorem 1.2 hold. If n=1, (n): space dimension, then it can be shown that any solution is bounded, and so converges to some $\theta(\gamma, \mu)$ as $t \to \infty$, uniformly on $\bar{\Omega}$.

Corollary 1.2. In addition to the assumptions in Theorem 1.2, assume that $a_{jk} = -a_{kj}$, r = N - 2, and $\Lambda_3 = \Lambda \ (\equiv \{1, ..., N\})$. Then any bounded solution of (8) converges to some periodic orbit as $t \to \infty$, uniformly in $\bar{\Omega}$.

Example 1.1. Consider the reaction-diffusion system

$$\frac{\partial}{\partial t} u_j = d_j \partial_x^2 u_j + (u_{j+1} - u_{j-1}) u_j, \ x \in \Omega, \ t > 0$$

$$\frac{\partial}{\partial \nu} u_j = 0 \ (x \in \partial \Omega, \ t > 0); \ u_j(x, 0) = \phi_j(x), \ (x \in \Omega)$$

$$(j = 1, 2, 3; \ u_4 = u_1, \ u_0 = u_3)$$
(11)

 $(\partial_x^2 = \partial^2/\partial x^2)$. It is easy to verify that Assumption 1.1 is satisfied, that $\Lambda_1 = \Lambda_2 =$ empty set and $\Lambda_3 = \{1, 2, 3\}$, and that the kernel of $(a_{jk})^*$ is spanned by the positive vector (1, 1, 1). Thus any bounded solution of (11) tends to the set

$$\Theta = \{ \xi \in \mathbb{R}^3_+; \ \xi_1 + \xi_2 + \xi_3 = \gamma_0, \ \xi_1 \xi_2 \xi_3 = \gamma_1 \}$$

for some positive γ_0 , γ_1 . Clearly Θ is a periodic orbit. Hence any bounded solution of (11) with positive initial data converges to some periodic orbit as $t \to \infty$, uniformly on $\bar{\Omega}$.

Example 1.2. Consider the reaction-diffusion system:

$$\frac{\partial}{\partial t} u_j = d_j \partial_x^2 u_j - e_j u_j + (u_{j+1} - u_{j-1}) u_j,
\frac{\partial}{\partial v} u_j = 0 \ (x \in \partial \Omega, \ t > 0), \ u_j(x, 0) = \phi_j(x) \ (> 0)
(j = 1, ..., N; \ u_0 = u_N, \ u_{N+1} = u_1),$$
(12)

where $e_k = 1$ (k = 1, ..., L); $e_k = (k = L + 1, ..., N)$. Then we can see that

$$\Lambda_1 = \{1, \dots, L, L+2, \dots, 2[\frac{N-L}{2}]\};$$

$$\Lambda_2 = \{L+1, \dots, 2[\frac{N-L+1}{2}] + L-1\}; \quad \Lambda_3 = \text{empty set},$$

where [] denotes Causs symbol; thus any bounded solution u of (12) converges to a constant vector $(\gamma_1, \ldots, \gamma_N)$ where $\gamma_j = 0$, $(j \in \Lambda_1)$ and $\gamma_j > 0$, $(j \in \Lambda_2)$.

2. Some estimates. In this section we give some estimates to be used later. We begin by introducing some notations. $\| \|_p$ denotes the usual L^p -norm over Ω ; we simply write $\| \|$ for $\| \|_2$. $H^{p,2}$ denotes the L^p -Sobolev space of order 2 with the norm $\| \|_{p,2}$. We define the operator P by

$$Pw = \frac{1}{|\Omega|} \int_{\Omega} w(x) \, dx.$$

In what follows, M denotes various constant independent of t; and set

$$g_{j}(\xi) = \xi_{j} f_{j}(\xi)$$

$$K = \sup |u_{j}(x, t)| \quad (x \in \bar{\Omega}, \ t \ge 0, \ j = 1, ..., N)$$

$$K_{1} = \sup_{|\xi| \le K, j} |f_{j}(\xi)|, \quad K_{2} = \sup_{|\xi| \le K, j} |\nabla_{\xi} g_{j}(\xi)|.$$
(13)

Since the initial function is non-negative and not identically zero, it follows from the elementary property of parabolic equations that

$$u_j(x,t) > 0, \quad (x \in \bar{\Omega}, t > 0), \quad j = 1, \dots, N.$$
 (14)

Since we are concerned with behavior of solution for large t, we may assume that (14) hold for $x \in \bar{\Omega}$ and $t \geq 0$.

Lemma 2.1. Let u be a solution of (RD). Then

$$||u_j(t)||_1 \le M, \quad (j \in \Lambda), \tag{15}$$

$$\int_0^t \|u_j(s)\|_1 \, ds \le M, \quad (j \in \Lambda_{1,1}), \tag{16}$$

where M is a constant.

(Here and in what follows we shall simply write $u_i(t)$ for $u_i(x, t)$)

Proof. Integrating the j-th equation in (RD) in x and t over $\Omega \times (0, t)$, multiplying b_j , and taking the summation in j, we get by (2)

$$\int_{\Omega} \langle b, u(x, t) \rangle \, dx + \delta_1 \sum_{j \in \Lambda_{1,1}} \int_0^t \int_{\Omega} b_j u_j(x, s) \, dx \, ds \le \int_{\Omega} \langle b, \phi(x) \rangle \, dx, \tag{17}$$

 (\langle, \rangle) : the inner product in \mathbb{R}^N) from which (15) and (16) follow immediately in view of (14).

Lemma 2.2. Let u be a bounded solution of (RD). Then

$$\int_0^t \|\nabla u_j(s)\|^2 ds \le M, \quad j \in \Lambda; \tag{18}$$

$$\int_{0}^{t} \|u_{j}(s)\|_{1} ds \le M, \quad j \in \Lambda_{1};$$
(19)

and

$$\int_0^t \|g_k(u)\|_1 \, ds \le M, \quad k \in \Lambda_1 \cup \Lambda_2, \tag{20}$$

where M is independent of t.

Proof. By the assumption

$$\sup |u(x,t)| \ (\equiv K) < \infty, \quad (x \in \bar{\Omega}, \ t \ge 0). \tag{21}$$

If $j \in \Lambda_{1,1}$, then (16) implies (19). Let $j \in \Lambda_{1,2} - \Lambda_{1,1}$. Then for some $i \in \Lambda$,

$$f_i(u) \ge \delta_2 u_j - \rho(K) \sum_{k \in \Lambda_{1,1}} u_k.$$

Integration of the *i*-th equation (divided by u_i) in (RD) over $\Omega \times (0, t)$ gives

$$d_{i} \int_{0}^{t} \|(\nabla u_{i}(s))/u_{i}(s)\|^{2} ds + \delta_{2} \int_{0}^{t} \|u_{j}(s)\|_{1} ds$$

$$\leq \int_{\Omega} \log u_{i}(x, t) dx - \int_{\Omega} \log \phi_{i}(x) dx + \rho(K) \sum_{k \in \Lambda_{1,1}} \int_{0}^{t} \|u_{k}(s)\|_{1} ds.$$
(22)

The first term on the right hand side of (22) is, by (21), bounded, since $\log u_i(x, t) \leq \log_+ u_i(x, t) \leq |u_i(x, t)|$. The second term is, by (14), bounded. The third term is, by (16), bounded. Hence the right-side is bounded, and so is the left hand side. Thus (19) holds for $j \in \Lambda_{1,2}$. Inductively we can show (19) holds for $j \in \Lambda_1$. Clearly (20) holds for $j \in \Lambda_1$ in view of (21) and (19). For $j \in \Lambda_2$ we have, by (4) and (21),

$$||g_j(u)||_1 \le \rho(K) \sum_{i \in \Lambda_1} ||u_i||_1,$$

which together with (19) gives (20) with $j \in \Lambda_2$. Taking the inner product of the j-th equation in (RD) with u_j , and then integrating the result in t, we see

$$||u_{j}(t)||^{2} + 2d_{j} \int_{0}^{t} ||\nabla u_{j}(s)||^{2} ds$$

$$= ||\phi_{j}||^{2} + 2 \int_{0}^{t} (g_{j}(u(s)), u_{j}(s)) ds, \quad (j \in \Lambda_{1} \cup \Lambda_{2})$$
(23)

((,): L^2 -inner product). The right hand side is, by (20), bounded, in view of the boundedness of u. This shows that (18) holds for $j \in \Lambda_1 \cup \Lambda_2$. Finally let $j \in \Lambda_3$. Similarly to (22), integrating the i-th equation (multiplied by c_i/u_i), and taking the sum in i one finds that

$$\sum_{i} c_{i} \int_{0}^{t} \|(\nabla u_{i})/u_{i}\|^{2} ds = \sum_{i} c_{i} \left[\int_{\Omega} \log u_{i}(x, t) dx - \int_{\Omega} \log \phi_{i}(x) dx \right] - \sum_{i} c_{i} \int_{0}^{t} \int_{\Omega} f_{i}(u) dx ds.$$
(24)

By Assumption 1.2 the right-hand side of (24) is bounded by

$$(\|u_i(t)\|_1 + \|\log \phi_i\|_1) + \rho(K) \sum_{k \in \Lambda_1} \int_0^t \|u_k\|_1 \, ds,$$

which is, by (15) and (19), bounded. Hence the left hand side of (24) is also bounded. Consequently, by the positivity of c_i ,

$$\int_0^t \|\nabla u_j\|^2 ds \le K^2 \int_0^t \|(\nabla u_j)/u_j\|^2 ds \le M,$$

showing that (18) holds for $j \in \Lambda_3$. This proves Lemma 2.2. \square

To get the L^{∞} -bounds for solutions u of (RD), we introduce an operator $A_{j,p}$ in $L^{p}(\Omega)$:

$$\begin{split} D(A_{j,p}) &= \{ v \in H^{p,2}(\Omega); \ (\partial/\partial \nu) v = 0 \ (\text{on } \partial \Omega) \}; \\ A_{j,p}v &= -d_j \Delta v + \delta_1 v \ (j \in \Lambda_{1,1}); \ = -d_j \Delta v \ (\text{otherwise}). \end{split}$$

We first note that $A_{j,2}$ is a non-negative self-adjoint operator in $L^2(\Omega)$. Let us fix p so that p > n, and write A_j for $A_{j,p}$ for simplicity. Then A_j has the following properties:

- i) the spectral set of A_j consists only of isolated eigenvalues $\{\lambda_j\}$ with $0 \le \lambda_1 < \lambda_2 \le \ldots$, and with finite multiplicities;
- ii) the first eigenvalue λ_1 is positive if and only if $j \in \Lambda_{1,1}$;
- iii) the estimate holds:

$$||v||_{p,2} \le M\{||A_j v||_p + ||v||_p\}, \quad v \in D(A_j);$$
 (25)

- iv) if we define the operator Q by Q = I, (if $\lambda_1 > 0$); Q = I P, (if $\lambda_1 = 0$) (I = identity operator), then $Q e^{-tA_j} = Qe^{-tA_l}Q$;
- v) A_j generates the holomorphic semigroups $\{e^{-tA_j}\}$ in $L^p(\Omega)$ so that

$$||e^{-tA_j}|| \le M; ||Qe^{-tA_j}|| \le Me^{-\beta t}; ||A_je^{-tA_j}|| \le Me^{-\beta t}/t,$$
 (26)

with some positive β ;

vi) the solution u of (RD) can be written as

$$u_j(t) = e^{-tA_j}\phi_j + \int_0^t e^{-(t-s)A_j} g_j^*(u(s)) ds,$$
 (27)

where $g_j^*(u) = g_j(u) + \delta_1 u_j$, $(j \in \Lambda_{1,1})$; $= g_j(u)$ (otherwise) (see Agmon-Douglis-Nirenberg [1], Friedmann [7]).

Lemma 2.3. We have

$$||v||_{\infty} \le M||A_iv||_p + M|Pv|, \quad (v \in D(A_i))$$
 (28)

(| |: absolute value)

Proof. Using the a priori estimate (25) for solutions of elliptic equations, $||Qv||_p \le M||A_jv||_p$, and the Sobolev inequality, we get

$$||v||_{\infty} \leq M(||A_jv||_p + ||Pv||_p),$$

showing (28); note $||Pv||_p \le |Pv||\Omega|^{1/p}$.

Lemma 2.4. Let u be a bounded solution of (RD). Then

$$\int_{0}^{\infty} |Qg_{j}(u(s))|_{p}^{p} ds < \infty, \quad (j = 1, ..., N).$$
 (29)

Proof. If $j \in \Lambda_{1,1}$, then (29) is clear from (16). Suppose $j \in \Lambda - \Lambda_{1,1}$. By integration by parts, $(-\Delta u_j, u_j) = |\nabla u_j|^2$, which is integrable by (18). Since

$$||Qu_j||^2 \le M||A_{i,2}^{1/2}u_j||^2 = M(A_{j,2}u_j, u_j) = Md_j(-\Delta u_j, u_j) = Md_j||\nabla u_j||^2,$$

which is integrable by (18) and since $||Qu_j(t)||_p^p \le 2^{p-2}K^{p-2}||Qu_j(t)||^2$, we see that $||Qu_j(t)||_p^p$ is integrable on $[0, \infty)$. By the mean-value theorem,

$$\|Qg_{j}(u(t))\|_{p} = \|Q[g_{j}(u) - g_{j}(Pu)]\|_{p}$$

$$\leq \sum_{i=1}^{N} \int_{0}^{1} \|Q[\partial_{\xi_{i}}g_{j}(u + s(Pu - u)) \cdot (Pu_{i} - u_{i})]\|_{p} ds$$

$$\leq K_{2}M \sum_{i=1}^{N} \|Pu_{i} - u_{i}\|_{p} \leq K_{2}M \sum_{i=1}^{N} \|Qu_{i}\|_{p},$$
(30)

from which (29) follows. This proves Lemma 2.4.

Lemma 2.5. Let u be a bounded solution of (RD). Then

$$||A_i u_i(t)||_p \to 0 \quad as \quad t \to \infty,$$
 (31)

$$\|Qu_j(t)\|_p \to 0 \quad as \quad t \to \infty,$$
 (32)

$$\|Qu_j(t)\|_{\infty} \to 0 \quad as \quad t \to \infty.$$
 (33)

In particular,

$$\|u_j(t)\|_{\infty} \to 0 \quad as \quad t \to \infty \quad (j \in \Lambda_{1,1}).$$
 (34)

Proof. We first show (32). Applying the Q to both sides of (27), and using (26) and iv), we find that

$$||Qu_j(t)||_p \le Me^{-t\beta}||\phi_j||_p + M\int_0^t e^{-(t-s)\beta}||Qg_j(u(s))||_p ds.$$

Letting $t \to \infty$ in the above inequality, we have, by (29), (32). We next show (31). To this end we express $A_j u_j$ in the form:

$$A_{j}u_{j}(t) = A_{j}e^{-tA_{j}}\phi_{j} + \int_{t/2}^{t} A_{j}e^{-(t-s)A_{j}}(Qg_{j}(u(s)) - Qg_{j}(u(t)) ds$$

$$+ \int_{0}^{t/2} A_{j}e^{-(t-s)A_{j}}Qg_{j}(u(s)) ds + (I - e^{-(t/2)A_{j}})Qg_{j}(u(t))$$

$$(\equiv J_{1} + J_{2} + J_{3} + J_{4}).$$

Clearly, $J_1 \to 0$ as $t \to \infty$ (in L^p). By(26) and (32), $J_4 \to 0$ as $t \to \infty$. From (25) it is easy to see that

$$||J_3||_p \le MK \int_0^{t/2} (t-s)^{-1} e^{-(t-s)\beta} ||Qg_j(u(s))||_p ds$$

from which it follows that $J_3 \to 0$. It remains only to show $J_2 \to 0$. Similarly to (30),

$$\|Qg_{j}(u(t)) - Qg_{j}(u(s))\|_{p} \le MK_{2}\|u(t) - u(s)\|_{p}.$$
(35)

On the other hand, by standard arguments in the theory of (linear) evolution operators, we can show:

$$||u_j(t) - u_j(s)||_p \le MK_1(|t - s| + |t - s|^{1/2}).$$
 (36)

(For the proof see the appendix). Thus it follows from (26), (30), (23) and (36) that

$$||J_2||_p \le \rho(K) \int_{t/2}^t (|t-s|^{-1/2} + |t-s|^{-3/4}) e^{-(t-s)\beta} ds \, \omega(t)^{1/2},$$

where $\omega(t) = \sup_{t/2 \le s \le t} \|Qu(s)\|_p$. Since $\omega(t) \to 0$, it follows that $J_2 \to 0$. This shows (31). (33) and (34) are immediate consequences of (28), (31) and the Sobolev inequality.

Lemma 2.6. Let $j \in \Lambda_3$. Then there is a $\delta > 0$ such that

$$|Pu_i(t)| \ge \delta > 0 \quad (t > 0). \tag{37}$$

Proof. By (24) and (2),

$$\sum_{j \in \Lambda_3} c_j(P(\log u_j(t)) - P(\log \phi_j))$$

$$\geq \int_0^t \int_{\Omega} \sum_{j \in \Lambda_3} c_j f_j(u) \, dx \, ds$$

$$\geq -\rho(K) \sum_{j \in \Lambda_1} \int_0^t \|u_j(s)\|_1 \, ds \qquad [by (6)]$$

$$\geq -M, \quad (> -\infty) \qquad [by (19)].$$

Consequently, by the positivity of c_j ,

$$P\log u_i(t) \ge -M_0, \quad (t > 0) \tag{38}$$

with some constant M_0 . By Jenssen's inequality, (37) follows from (38).

3. Proof of Theorem 1.1. We first show that $u_j(t)$ $(j \in \Lambda_1 \cup \Lambda_2)$ converges to some constant γ_j as $t \to \infty$, uniformly on Ω . We decompose u_j in the form:

$$u_i(t) = Pu_i(t) + Qu_i(t), \quad (= I_1(t) + I_2(t)).$$
 (39)

Then by (33)

$$I_2(t) \to 0$$
, uniformly on Ω . (40)

Applying P to both side of (RD), integrating in x and t and noting $P \Delta u_j = 0$, we see

$$I_1(t) - I_1(s) = \frac{1}{|\Omega|} \int_s^t (\int_{\Omega} g(u) \, dx) \, ds. \tag{41}$$

Since $||g_j(u)||_1$, $(j \in \Lambda_1 \cup \Lambda_2)$ is, by (20), integrable on $[0, \infty)$, it follows that $\{I_1(t)\}$ is a Cauchy sequence. Hence there is a constant γ_j with

$$I_1(t) \rightarrow \gamma_i$$

which together with (40) shows that $u_j(t)$ converges to γ_j as $t \to \infty$, uniformly on $\bar{\Omega}$. Clearly $\gamma_j \geq 0$ since $u_j \geq 0$. Since $||u_j(t)||_1$, $(j \in \Lambda_1)$ is integrable on $(0, \infty)$ by (19), it follows that

$$\gamma_i = 0, \quad (j \in \Lambda_1).$$

Let $j \in \Lambda_2$. Similarly to (24) we get

$$\int_0^t \int_{\Omega} f_j(u) \, dx \, ds + \int_{\Omega} \log \phi_j(x) \, dx \le \int_{\Omega} \log u_j(x, t) \, dx. \tag{42}$$

The first term on the left hand side is bounded from below, and so

$$-M_1 \le$$
 the left hand side of (42),

 M_1 being some positive constant independent of t, since $||f_j(u)||_1$ is integrable on $[0, \infty)$ in view of (4) and (19). Hence letting $t \to \infty$ in (42) we see that the limit of the right hand side is bounded by $-M_1$ from below. Since the limit of the right side is $|\Omega| \log \gamma_j$, it follows that

$$\gamma_j > 0 \quad (j \in \Lambda_2). \tag{43}$$

Set

$$h(t) = \sum_{k=1}^{N} b_k P u_k(t)$$

for simplicity. Then by (1)

$$(d/dt)h(t) = P(\sum_{k=1}^{n} b_k g_k(u)) \le 0$$

since $P \Delta u_j = 0$. Hence h(t) is monotone decreasing in t. h(t) is non-negative, since $u_j(t)$ is non-negative. Thus the limit of h(t) exists, and we denote it by μ_0 . If $\Lambda_1 \neq \Lambda$, then $\mu_0 > 0$ in view of (37), and (43). Finally we show

$$J \equiv \sum_{k \in \Lambda_2} c_k^{(j)} \log u_k(t) \to \mu_j \quad (t \to \infty), \text{ uniformly on } \Omega.$$
 (44)

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Put

$$J_1 = PJ, \quad J_2 = QJ, \quad c_k = c_k^{(j)}.$$

Then

$$J_2 = \sum_k c_k Q[\log(Pu_k(t) + Qu_k(t)) - \log Pu_k(t)]$$

$$= \sum_k c_k Q \log(1 + Qu_k(t)/Pu_k(t))$$

which tends to zero as $t \to \infty$, uniformly on $\bar{\Omega}$, in view of (33) and (37). Simple calculation gives

$$J_1(t) - J_2(s) = \sum_k c_k \int_s^t \left[\int_{\Omega} (d_k |\nabla u_k| / |u_k|^2 + f_k(u)) \, dx \right] ds. \tag{45}$$

By (24), $|(\nabla u_k)/u_k|$ is square integrable on $\Omega \times (0, \infty)$. Also the absolute value of the integral of $\sum_k c_k f_k(u)$ on the right hand side of (44) is, by (6), dominated by the integrable function

$$\rho(K) \sum_{k \in \Lambda_1} \int_s^t \|u_k(\tau)\|_{L^1} d\tau.$$

Hence the right hand side of (45) tends to zero as $s, t \to \infty$. Thus $\{J_1(t)\}$ is a Cauchy sequence, and so is J(t). This shows (44). This completes the proof of Theorem 1.1.

Appendix. Proof of (36). Here we shall give the proof of (36). Set $A_j = A$, $g_j = g$, etc. for simplicity. To show (36), we estimate each term on the right-hand side of the equation

$$u(t) - u(s) = (e^{-tA} - e^{-sA})\phi + \int_{s}^{t} e^{-(t-\tau)A} g(u(\tau)) d\tau + \int_{0}^{s} (e^{-(t-\tau)A} - e^{-(s-\tau)A}) f(u(\tau)) d\tau \ (\equiv J_1 + J_2 + J_3).$$

From the elementary properties in semigroup theory it follows that

$$\begin{aligned} \|(e^{-tA} - e^{-sA})\| &\leq (t - s)\|(e^{-(t - s)A} - I)/((t - s)A)\| \|Ae^{-sA}\| \\ &\leq M(t - s)s^{-1}e^{-s\beta}; \text{ and} \\ \|e^{-tA} - e^{-sA}\| &\leq Me^{-s\beta}. \end{aligned}$$

Hence by the interpolation theorem,

$$||e^{-tA} - e^{-sA}|| \le M(t-s)^{\theta} s^{-\theta} e^{-s\beta} \quad (0 < \theta < 1; \ 0 < s < t).$$
 (A1)

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Similarly,

$$\|(e^{-tA} - e^{-sA})\phi\|_p \le M(t-s)e^{-s\beta}\|A\phi\|_p$$

Thus by (26)

$$||J_1||_p \le M(t-s), ||J_2||_p \le MK_3(t-s),$$

where $K_3 = \sup |g(\xi)|$, $(|\xi| \le K)$. Also by (A1) with $\theta = 1/2$,

$$||J_3|| \leq M(t-s)^{1/2}$$
.

Collecting all the estimates above we get (36).

Acknowledgment: The authors express their gratitude to the referee for useful comments, e.g., on the definition of Λ_1 .

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