Asymptotic safety of gravity-matter systems

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IMPRS *ptfs*

Outline

Introduction

- Motivation
- Truncation

2 Anomalous dimensions

- Scheme for an approximated momentum dependence
- Bounds on the anomalous dimensions

B Results

- Fixed points of dynamical couplings
- Background couplings

Summary and Outlook

Why gravity-matter systems?

- Pure quantum gravity does not exist in nature
- Matter seems to matter Percacci, Perini (2003)

Donà, Eichhorn, Percacci (2013)

Does the non-trivial UV fixed point persist within inclusion of matter?

Why a vertex expansion?

Features:

- Systematic approximation scheme
- Disentangling background and fluctuation fields
- Newton's coupling from graviton three-point function

Challenges:

- Long algebraic computations due to large tensor structures
- Physical choice of tensor projection
- Technical issues for non-flat backgrounds (so far)

$$\Gamma_{k}[\bar{g},\phi] = \sum_{n=0}^{\infty} \frac{1}{n!} \Gamma_{k}^{(\phi_{1}\dots\phi_{n})}[\bar{g},0]\phi_{1}\cdots\phi_{n} \quad \text{with} \quad \phi = (h,c,\bar{c},\psi,\bar{\psi},\varphi)$$

$$= \Gamma_{k}[\bar{g},0] + \Gamma_{k}^{(h)}[\bar{g},0]h + \frac{1}{2}\Gamma_{k}^{(2h)}[\bar{g},0]h^{2} + \frac{1}{3!}\Gamma_{k}^{(3h)}[\bar{g},0]h^{3} + \dots$$

$$+ \frac{1}{2}\Gamma_{k}^{(\bar{c}c)}[\bar{g},0]\bar{c}c + \frac{1}{2}\Gamma_{k}^{(\bar{\psi}\psi)}[\bar{g},0]\bar{\psi}\psi + \frac{1}{2}\Gamma_{k}^{(\varphi\varphi)}[\bar{g},0]\varphi^{2} + \dots$$

- Computation of all *n*-point correlation functions up to third order
- Flow equation for each correlation function from Wetterich equation

Graviton two-point function:

$$\partial_t (\underbrace{\longrightarrow})^{-1} = -\frac{1}{2} \underbrace{\bigcirc}_{+} + \underbrace{\bigcirc}_{-} - 2 \underbrace{\bigcirc}_{+} + \underbrace{$$

Fermionic two-point function:

$$\partial_t (\longrightarrow)^{-1} = -\frac{1}{2} \xrightarrow{\otimes} + +$$

Generating action:

$$\begin{split} S &= S_{\mathsf{EH}} + S_{\mathsf{gh}} + S_{\mathsf{gf}} + \int \mathrm{d}^4 x \sqrt{g} \bar{\psi}_i \nabla \psi_i + \frac{1}{2} \int \mathrm{d}^4 x \sqrt{g} g_{\mu\nu} \partial^{\mu} \varphi_j \partial^{\nu} \varphi_j \\ \text{with} \qquad i = 1, \dots, N_f \qquad \text{and} \qquad j = 1, \dots, N_s \end{split}$$

Tensor structures for vertices:

$$\mathcal{T}^{(\phi_1\dots\phi_n)}(p_1,\dots,p_n;\Lambda_n) = S^{(\phi_1\dots\phi_n)}(p_1,\dots,p_n;\Lambda\to\Lambda_n,G_N\to1) \quad (1)$$

Spin-connection via spin-base invariance formalism Weldon (2000); Gies, Lippoldt (2013)

RG-invariant ansatz for *n*-point function Christiansen, Knorr, Pawlowski, Rodigast (2014)

$$\Gamma^{(\phi_1...\phi_n)} = \prod_{i=1}^n \sqrt{Z_{\phi_i}(p_i^2)} G_n^{\frac{n}{2}-1} \mathcal{T}^{(\phi_1...\phi_n)}(p_1,\ldots,p_n;\Lambda_n)$$
(2)

Graviton two-point function:

$$\Gamma^{(hh)} = Z_h(p^2) \mathcal{T}^{(hh)} \sim Z_h(p^2)(p^2 - 2\Lambda_2)\Pi_{ ext{TT}} + ext{other tensor structures}$$

Fermionic two-point function:

$$\Gamma^{(\overline{\psi}\psi)} = Z_{\psi}(p^2)i\gamma_{\mu}p^{\mu}$$

RG-invariant ansatz for *n*-point function Christiansen, Knorr, Pawlowski, Rodigast (2014)

$$\Gamma^{(\phi_1...\phi_n)} = \prod_{i=1}^n \sqrt{Z_{\phi_i}(p_i^2)} G_n^{\frac{n}{2}-1} \mathcal{T}^{(\phi_1...\phi_n)}(p_1,\ldots,p_n;\Lambda_n)$$
(2)

Graviton two-point function:

 $\Gamma^{(hh)} = Z_h(p^2) \mathcal{T}^{(hh)} \sim Z_h(p^2)(p^2 + M^2) \Pi_{TT}$ + other tensor structures

 $M^2 = -2\Lambda_2 = \text{graviton mass parameter}$

- Transverse-traceless part of graviton *n*-point functions
- $\Gamma^{(3h)}$ at symmetric momentum configuration
- Bilocal momentum projection for $\Gamma^{(3h)}$ Christiansen, Knorr, Meibohm, Pawlowski, MR
- Momentum dependence of $\Gamma^{(\phi\phi)}$ encoded in anomalous dimensions

Truncation

- Flow of *n*-point functions up to order three
- Flat Euclidean background $g_{\mu
 u} = \delta_{\mu
 u} + h_{\mu
 u}$
- Closure of flow equations (gravity sector)

$$G_n(p_1,\ldots,p_n)\equiv G_3=:G$$

$$\Lambda_5\equiv\Lambda_4\equiv\Lambda_3$$

- Minimally coupled scalars and fermions
- No matter-matter-interactions

Parameters of the truncation

$$G, \Lambda_3, M^2, \eta_h(p^2), \eta_c(p^2), \eta_\psi(p^2), \eta_\varphi(p^2), N_s, N_f$$

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The anomalous dimensions obey coupled Fredholm integral equations, e.g.

$$\eta_h(p^2) \sim \int \mathrm{d}\Omega \mathrm{d}q \; q^3 \left(\dot{r}(q^2) - \eta_i(q^2)r(q^2)\right) f_i(p,q,G,\Lambda_3,M^2)$$

In our setup: Four coupled Fredholm integral equations

 \rightarrow Requires large numerical effort for solution (Liouville-Neumann series)

Approximation

Evaluate anomalous dimensions at the peak of the integral q = k

Quality of approximated momentum dependence $(N_s = N_f = 0)$

Full momentum dependence Christiansen, Knorr, Meibohm, Pawlowski, MR (2015)

$$egin{aligned} & (g^*,\mu^*,\lambda_3^*) = (0.66,-0.59,0.11) \ & (heta_{1/2}, heta_3) = (-1.4\pm 4.1\,i,14) \end{aligned}$$

Approximated momentum dependence

$$(g^*, \mu^*, \lambda_3^*) = (0.62, -0.57, 0.095) (\theta_{1/2}, \theta_3) = (-1.3 \pm 4.1 \, i, 12)$$
 (4)

Error of 6% (15%) on fixed point values (critical exponents)

Key message

Momentum dependence of the anomalous dimensions is important

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The regulator is proportional to $Z_\phi \sim k^{-\eta_\phi}$

$$R_{k}^{\phi}(p^{2}) = \left. \Gamma_{k}^{(\phi\phi)}(p^{2}) r_{k}^{\phi}(p^{2}) \right|_{M^{2}=0}$$
(5)

But the regulator has to satisfy (with i = 1 for fermions and i = 2 else)

$$\lim_{k \to \infty} R_k^{\phi}(p) \sim \lim_{k \to \infty} Z_{\phi} k^i \sim \lim_{k \to \infty} k^{i - \eta_{\phi}} \xrightarrow{!} \infty$$
(6)

Bounds on the anomalous dimensions - from the regulator

$$\eta_h < 2\,, \qquad \eta_c < 2\,, \qquad \eta_{\varphi} < 2\,, \qquad \eta_{\psi} < 1$$

Bounds on the anomalous dimensions - from the flow

Contributions to the flow have the typical shape

diagram $\sim (\# - \eta_{\phi})$

with $\# = \{4, 5, 6, \dots\}$ for analytic approximation

The sign of the diagram is regulator dependent for $\eta_{\phi} > \#$

Bounds on the anomalous dimensions - from the flow

$$\eta_h \lessapprox 4$$
, $\eta_c \lessapprox 10$, $\eta_{\varphi} \lessapprox 8$, $\eta_{\psi} \lessapprox 5$

We take the most conservative bound!

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Questions so far?

Flow equations in analytic approximation

Flow of dynamical couplings (matter contributions)

$$\partial_t g \sim -\frac{43}{570\pi} g^2 N_s - \frac{3599}{11400\pi} g^2 N_f ,$$

$$\partial_t \lambda_2 \sim -\frac{1}{24\pi} g N_s + \frac{4}{9\pi} g N_f , \qquad (7)$$

compared to background couplings in background-field approximation

$$\partial_t \bar{g} \sim + \frac{1}{6\pi} \bar{g}^2 N_s + \frac{1}{3\pi} \bar{g}^2 N_f ,$$

$$\partial_t \bar{\lambda} \sim + \frac{1}{12\pi} (3 + 2\bar{\lambda}) \bar{g} N_s - \frac{1}{3\pi} (3 - \bar{\lambda}) \bar{g} N_f .$$
(8)

Do the different signs indicate an error?

No, it's just a computation of different couplings!

- g from $(\Gamma^{(3h)}(p = k) \Gamma^{(3h)}(p = 0))$; numerical integration necessary $Z_h(p^2)$ models remaining p-dependence Christiansen, Knorr, Meibohm, Pawlowski, MR
- λ_3 from $\Gamma^{(3h)}(p=0)$; analytic equation
- μ from $\Gamma^{(2h)}(p=0)$; analytic equation
- $\eta_{\phi}(p^2)$ from $\Gamma^{(\phi\phi)}(p^2)$ with approximated *p*-dependence
- Remark: $\eta_{\varphi}(p^2) = 0$ due to choice of graviton gauge

Inclusion of scalars



 $\begin{array}{ll} \mbox{Reliability bound on truncation } (\eta_h^*>2) & $N_{s_{\rm trunc}}\approx21.5$\\ \mbox{Loss of an attractive UV fixed point} & $N_{s_{\rm stab}}\approx42.6$\\ \mbox{Complete loss of the UV fixed point} & $N_{s_{\rm max}}\approx66.4$\\ \end{array}$

Inclusion of fermions



Reliable, attractive UV fixed point for all N_f

Mixed scalar-fermion systems



One fermion compensates for $\approx 7.1~\text{scalars}$

- ullet The background couplings \bar{g} and $\bar{\lambda}$ are observables for $k \to 0$
- Their flow depends only on dynamical quantities

$$\partial_t \left(\frac{k^2}{\bar{g}}\right) = F_{R^1}(g, \lambda; N_s, N_f)$$

$$\partial_t \left(\frac{\bar{\lambda}k^4}{\bar{g}}\right) = F_{R^0}(g, \lambda; N_s, N_f)$$
(9)

• Same equations as earlier computations if $(g, \lambda)
ightarrow (ar{g}, ar{\lambda})$

Codello, Percacci, Rahmede (2008); Donà, Eichhorn, Percacci (2013)

Background couplings vs dynamical couplings



Qualitative agreement only for small N_s and N_f

Comparison with previous results



DEP: Donà, Eichhorn, Percacci (2013) Good qualitative agreement for small N_s and N_f

Setup:

- Vertex expansion up to third order
- Genuine Newton's coupling from graviton three-point function
- Minimally coupled scalars and fermions

Results:

- All numbers of fermions are compatible with asymptotic safety
- Need to improve truncation for intermediate and large N_s
- Background and dynamical couplings deviate for large N_s and N_f

Outlook - Road towards better truncations

Improvement in the matter sector (Matter-matter-interactions)

- Yukawa-coupling Held, Eichhorn, Pawlowski; in preparation
- φ^4 -coupling Fischer, Pawlowski; in preparation

Improvement in the gravity/mixed sector

- Graviton four-point function Denz, Pawlowski, MR; in preparation
- $h\varphi\varphi$ -coupling Donà, Eichhorn, Labus, Percacci; arXiv:1512.01589

Eichhorn, Labus, MR; in preparation

• Expansion about non-flat backgrounds