



Bifurcation Structures in a Bimodal Piecewise Linear Map

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In this paper we present an overview of the results concerning dynamics of a piecewise linear bimodal map. The organizing principles of the bifurcation structures in both regular and chaotic domains of the parameter space of the map are discussed. In addition to the previously reported structures, a family of regions closely related to the so-called U-sequence is described. The boundaries of distinct regions belonging to these structures are obtained analytically using the skew tent map and the map replacement technique.

Keywords: bimodal piecewise linear map, border collision bifurcation, border collision normal form, degenerate bifurcation, homoclinic bifurcation

1. INTRODUCTION

Various bifurcation scenarios have always been in focus of many researchers from different theoretical and applied fields. It is known that bifurcation sequences often differ for *piecewise smooth* (PWS) systems with respect to smooth ones (see, e.g., [1]). In particular, this happens due to *border collision bifurcations* (BCB). This specific notion was initially introduced Nusse and Yorke in [2], but similar phenomena were also investigated earlier in Leonov [3] and Feigin [4] (see also [5, 6] and references therein). Recall that due to a BCB an attracting fixed point may be translated to an attracting cycle of any period or even to cyclic chaotic intervals. The simplest example of such behavior is the *skew tent map*, whose bifurcation scenarios are completely described due to intensive studies [7–10]. Therefore, the skew tent map is often used as a BCB normal form when studying a generic 1D continuous PWS map [5, 11–13]. Another peculiar situation for PWS maps can happen when an eigenvalue (or a conjugate pair of complex eigenvalues) of a cycle crosses the unit circle. Because of certain degeneracy of the map at the bifurcation value, the standard theorems for flip, fold, or Neimark-Sacker bifurcations cannot be applied, and moreover, the result of such a crossing can be atypical. For example, an eigenvalue of a fixed point becomes -1 , but after the bifurcation one observes not a 2-cycle, but a chaotic attractor [14]. Further, the bifurcation structure of the skew tent map also illustrates such a phenomenon as *robustness of chaotic attractors* [15]. That is, in the parameter space there exist *connected* domains that are characterized by chaotic attractors, that never happens in smooth maps.

In the present paper we continue studying the bifurcation structure of the parameter space of a generic 1D continuous piecewise linear bimodal map f . In our previous works basic families of regions related to regular and chaotic dynamics were investigated. In Panchuk et al. [16] we disclosed three basic bifurcation structures related to attracting cycles: the *skew tent map* (STM) structure, the *period adding* (PA) structure, and the *fin* structure. The STM structure concerns cycles whose points are located on the two adjacent branches only (the left and the middle or the middle and the right). The PA structure is composed by regions associated with cycles placed on the two outermost branches (the left and the right). The fin structure is related to periodicity regions having one common boundary with a PA-region.

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In Panchuk et al. [17] we described certain parameter regions corresponding to chaotic attractors (chaoticity regions) and related bifurcation structures. These chaoticity regions neighbor PA and fin regions, and their boundaries are given by *homoclinic* bifurcations of repelling cycles. In general, if a cycle has negative multiplier, then a *merging* bifurcation occurs, and the number of pieces of a chaotic attractor is halved. If a cycle has positive multiplier, an *expansion* bifurcation occurs, and one observes an abrupt increase in size of the attractor with uneven density of points after the bifurcation. For more details see, e.g., [18].

In the current paper we summarize the results of previous works and describe new family of periodicity regions observed in the parameter space of the bimodal map. They emerge from the border that separates the parameter region D_1/D_2 where asymptotic dynamics is associated with two adjacent branches (STM structure) from the region D_0 where all three branches can be involved. For parameters belonging to D_1/D_2 symbolic sequences of the cycles form the so-called *U-sequence* (see, e.g., [19, 20]), among which only basic cycles can be stable while all others are unstable. For parameters in D_0 every cycle with symbolic sequence belonging to U-sequence has a new complementary cycle (whose symbolic sequence differs by one symbol) having a single point on the third branch. If the slope of the third branch is small enough, this new cycle is stable. We refer to this situation as the “stabilization” of a U-sequence cycle.

The paper is organized as follows. In Section 2 we introduce the main definitions and notations used throughout the paper. In Section 3 we recall how the STM structure is formed. Section 4 is devoted to PA and fin regions related to attracting cycles. In Section 5 we explain the phenomena that occur when the invariant absorbing interval involving only one border point collides with the second border point and expands over the third branch. In Section 6 the chaoticity regions surrounding the PA and the fin structures are studied. Section 7 concludes.

All results described are obtained in form of explicit analytic expressions, which can be used for various applied problems. For such a practical usage see, for instance, [21, 22].

2. PRELIMINARIES

2.1. Basic Notations

Let us consider a family of 1D continuous piecewise linear maps $f : \mathbb{R} \rightarrow \mathbb{R}$ such that

$$f : x \mapsto f(x) = \begin{cases} f_{\mathcal{L}}(x) = a_{\mathcal{L}}x + \mu_{\mathcal{L}}, & x \leq d_{\mathcal{L}}, \\ f_{\mathcal{M}}(x) = a_{\mathcal{M}}x + \mu_{\mathcal{M}}, & d_{\mathcal{L}} < x \leq d_{\mathcal{R}}, \\ f_{\mathcal{R}}(x) = a_{\mathcal{R}}x + \mu_{\mathcal{R}}, & x > d_{\mathcal{R}}, \end{cases} \quad (1)$$

$$\text{with } f_{\mathcal{L}}(d_{\mathcal{L}}) = f_{\mathcal{M}}(d_{\mathcal{L}}), \quad f_{\mathcal{M}}(d_{\mathcal{R}}) = f_{\mathcal{R}}(d_{\mathcal{R}}). \quad (2)$$

Here the slopes $a_{\mathcal{L}}, a_{\mathcal{M}}, a_{\mathcal{R}}$, the offsets $\mu_{\mathcal{L}}, \mu_{\mathcal{M}}, \mu_{\mathcal{R}}$, and the border points $d_{\mathcal{L}}, d_{\mathcal{R}}$ are real parameters, and we suppose that $d_{\mathcal{L}} < d_{\mathcal{R}}$. By using the conditions (2) any two of eight parameters can be eliminated, but for the sake of generality, we keep all of them.

In Panchuk et al. [16, 17] the regular and chaotic dynamics of map (1) is studied in the cases (i) $|a_{\mathcal{L}}| < 1, |a_{\mathcal{R}}| < 1$ and (ii)

$0 < a_{\mathcal{L}} < 1, a_{\mathcal{R}} > 1$. In the present work we consider the case $a_{\mathcal{L}} > 0, a_{\mathcal{R}} \geq 0$.

Let us first introduce the notations used throughout the paper.

- We denote by p a point in the parameter space of the map f .
- The values of f at the border points, $\ell := f(d_{\mathcal{L}})$ and $r := f(d_{\mathcal{R}})$, are called (following the tradition of French school on Iteration Theory [23]) *critical points*. Successive images of ℓ and r are denoted as $\ell_i := f^i(\ell)$ and $r_i := f^i(r)$, $i \geq 1$.
- The term *partition* refers to the intervals $I_{\mathcal{L}} = (-\infty, d_{\mathcal{L}}]$, $I_{\mathcal{M}} = (d_{\mathcal{L}}, d_{\mathcal{R}}]$, and $I_{\mathcal{R}} = (d_{\mathcal{R}}, \infty)$. With each partition we associate the corresponding symbol, \mathcal{L}, \mathcal{M} , and \mathcal{R} , respectively. That is, every point $x \in I_s$ is coupled with the symbol s , where $s \in \{\mathcal{L}, \mathcal{M}, \mathcal{R}\}$.
- A cycle $\{x_i\}_{i=0}^{n-1}$ of f is associated with the symbolic sequence $\sigma = s_0 \dots s_i \dots s_{n-1}$, $|\sigma| = n$, $s_i \in \{\mathcal{L}, \mathcal{M}, \mathcal{R}\}$. The cycle is denoted by \mathcal{O}_{σ} .
- Consider a cycle \mathcal{O}_{σ} . Clearly, any cyclic shift of σ : $\sigma_i = s_i \dots s_{n-1} s_0 \dots s_{i-1}$ corresponds to the same cycle \mathcal{O}_{σ} . Resting on the sequence σ_i one can construct a function composition $f^n = f_{s_i} \circ \dots \circ f_{s_{n-1}} \circ f_{s_0} \circ \dots \circ f_{s_{i-1}} := f_{\sigma_i}$. Then every point x_i satisfies $f_{\sigma_i}(x_i) = x_i := x_{\sigma_i}$. For instance, the points of the 3-cycle $\{x_0, x_1, x_2\}$ with $x_0 \in I_{\mathcal{L}}, x_1 \in I_{\mathcal{M}}, x_2 \in I_{\mathcal{R}}$ can be written as $x_0 = x_{\mathcal{L}\mathcal{M}\mathcal{R}}, x_1 = x_{\mathcal{M}\mathcal{R}\mathcal{L}}, x_2 = x_{\mathcal{R}\mathcal{L}\mathcal{M}}$. Clearly, in the main notation for the cycle \mathcal{O}_{σ} any σ_i can be used, and we choose either the shortest one or the most suitable for a particular purpose.
- ν_{σ} denotes the *multiplier* of \mathcal{O}_{σ} , given by $\nu_{\sigma} := \prod_{i=0}^{n-1} a_{s_i}$.
- The *periodicity region* associated with the *stable cycle* \mathcal{O}_{σ} is denoted by \mathcal{P}_{σ} .
- Besides fixed points and cycles the map (1) can also have attracting n -cyclic chaotic intervals denoted as $\mathcal{Q}_n = \{J_i\}_{i=0}^{n-1}$. Hereby every J_i is a closed interval such that (i) f^n is chaotic on J_i , and (ii) $f(J_i) = J_{i+1}$, $i = \overline{0, n-2}$, $f(J_{n-1}) = J_0$. For convenience, J_0 refers to the leftmost interval.
- The symbol \mathcal{C}_n denotes a region in the parameter space related to \mathcal{Q}_n . We also refer to \mathcal{C}_n as a *chaoticity region*.

The regions \mathcal{P}_{σ} and \mathcal{C}_n are confined by boundaries related to certain bifurcations. The boundaries of periodicity regions \mathcal{P}_{σ} are usually associated with one of the following bifurcations: a *border collision bifurcation* (BCB) occurring when a point of a cycle collides with one of the border points; a *degenerate flip bifurcation* (DFB) which occurs when the multiplier of a cycle equals minus one; a *degenerate +1 bifurcation* (DB1) associated with the multiplier being equal plus one. For details on these bifurcations we refer to Sushko and Gardini [14], Nusse and Yorke [2], Banerjee et al. [5], and Zhusubaliyev and Mosekilde [6].

The boundaries of chaoticity regions are often related, besides BCBs and DFBs, to homoclinic bifurcations of a repelling cycle located at the immediate basin boundary of the attractor. For instance, a *merging bifurcation* or an *expansion bifurcation* leads to changing the number of intervals of a chaotic attractor. Whereas a *final bifurcation* (known also as a *boundary crisis*) implies a transformation from a chaotic attractor to a chaotic repeller (see, e.g., [18]). A homoclinic bifurcation is always associated with an unstable cycle \mathcal{O}_{σ} and is defined by the

condition $c_j = x_{\sigma_i}$, where x_{σ_i} is the appropriate point of \mathcal{O}_σ and $c_j = f^j(c)$, $c \in \{\ell, r\}$, is the critical point of the proper rank $j \geq 0$.

For the bifurcation boundaries of periodicity and chaoticity regions we use the following notations:

- $\xi_{\sigma_i}^d$ denotes a BCB of the cycle \mathcal{O}_σ , which occurs when $x_{\sigma_i} = d$ with $d \in \{d_{\mathcal{L}}, d_{\mathcal{R}}\}$.
- η_σ and θ_σ are used for a DFB and a DB1 boundaries, respectively, associated with $\nu_\sigma = -1$ and $\nu_\sigma = +1$.
- $\gamma_{\sigma_i}^{c_j}$, $\zeta_{\sigma_i}^{c_j}$ and $\chi_{\sigma_i}^{c_j}$ refer to homoclinic bifurcations: a merging bifurcation, an expansion bifurcation, and a final bifurcation, respectively.

2.2. Main Characteristic Regions in the Parameter Space

At first we describe the regions characterized by simple dynamics, that is, either an attracting fixed point or divergence to infinity. Then a brief description of regions related to more complex behavior follows (see **Figure 1**).

The function f can have at most three fixed points¹, one in each partition I_s , $s \in \{\mathcal{L}, \mathcal{M}, \mathcal{R}\}$. A fixed point \mathcal{O}_s , $s \in \{\mathcal{L}, \mathcal{R}\}$, given by $x_s = \mu_s / (1 - a_s)$ appears/disappears through collision with the border point d_s . The related bifurcation boundary is

$$\xi_s = \{p : \mu_s = (1 - a_s) d_s\}. \tag{3}$$

Notice that the upper index d_s is omitted in the notation ξ_s because it is clear that at $\xi_{\mathcal{L}}$ the collision occurs with $d_{\mathcal{L}}$ while at $\xi_{\mathcal{R}}$ it occurs with $d_{\mathcal{R}}$. Below, for shortness, we drop upper indices in notations for BCB boundaries whenever it is clear which border point is involved in the bifurcation (namely, when the first symbol of σ is \mathcal{L} or \mathcal{R} the related border point is $d_{\mathcal{L}}$ or $d_{\mathcal{R}}$, respectively).

From $0 \leq a_s < 1$ the stability regions of $\mathcal{O}_{\mathcal{L}}$ and $\mathcal{O}_{\mathcal{R}}$ are obtained:

$$\begin{aligned} \mathcal{P}_{\mathcal{L}} &= \{p : \mu_{\mathcal{L}} < (1 - a_{\mathcal{L}}) d_{\mathcal{L}}, 0 \leq a_{\mathcal{L}} < 1\}, \\ \mathcal{P}_{\mathcal{R}} &= \{p : \mu_{\mathcal{R}} > (1 - a_{\mathcal{R}}) d_{\mathcal{R}}, 0 \leq a_{\mathcal{R}} < 1\}. \end{aligned} \tag{4}$$

The stability region \mathcal{P}_s is separated from the divergence region \mathcal{S}_s by the boundary $a_s = 1$, where

$$\begin{aligned} \mathcal{S}_{\mathcal{L}} &= \{p : \mu_{\mathcal{L}} < (1 - a_{\mathcal{L}}) d_{\mathcal{L}}, \mu_{\mathcal{R}} < (1 - a_{\mathcal{R}}) d_{\mathcal{R}}, a_{\mathcal{L}} > 1\}, \\ \mathcal{S}_{\mathcal{R}} &= \{p : \mu_{\mathcal{R}} > (1 - a_{\mathcal{R}}) d_{\mathcal{R}}, \mu_{\mathcal{L}} > (1 - a_{\mathcal{L}}) d_{\mathcal{L}}, a_{\mathcal{R}} > 1\} \end{aligned} \tag{5}$$

(see **Figure 1** where $\mathcal{P}_{\mathcal{L}}$ and $\mathcal{S}_{\mathcal{R}}$ are shown).

The region

$$\mathcal{P}_{\mathcal{M}} = \{p : \mu_{\mathcal{L}} > (1 - a_{\mathcal{L}}) d_{\mathcal{L}}, \mu_{\mathcal{L}} + (1 + a_{\mathcal{L}}) d_{\mathcal{L}} - (1 + a_{\mathcal{R}}) d_{\mathcal{R}} < \mu_{\mathcal{R}} < (1 - a_{\mathcal{R}}) d_{\mathcal{R}}\}, \tag{6}$$

¹Except the case where the left/the right branch lies on the main diagonal. In this particular situation every point $x \in I_{\mathcal{L}}/x \in I_{\mathcal{R}}$ is a fixed point and has a stable set consisted of all its preimages, if they exist.

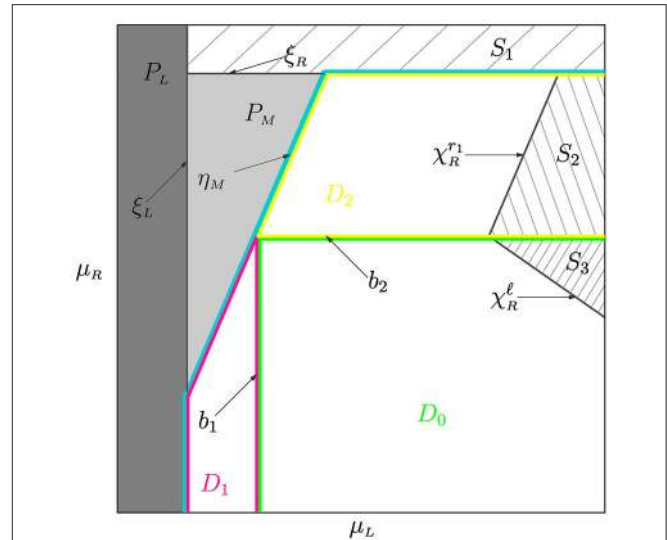


FIGURE 1 | The parameter plane $(\mu_{\mathcal{L}}, \mu_{\mathcal{R}})$ of the map f for $0 < a_{\mathcal{L}} < 1$ and $a_{\mathcal{R}} > 1$. The regions $\mathcal{P}_{\mathcal{L}}, \mathcal{P}_{\mathcal{M}}$ are related to stable fixed points. The regions $\mathcal{S}_1, \mathcal{S}_2, \mathcal{S}_3$ are related to diverging orbits. The regions $\mathcal{D}_1, \mathcal{D}_2$, and \mathcal{D}_0 are associated with asymptotic orbits located in $I_{\mathcal{L}} \cup I_{\mathcal{M}}, I_{\mathcal{M}} \cup I_{\mathcal{R}}$, and $I_{\mathcal{L}} \cup I_{\mathcal{M}} \cup I_{\mathcal{R}}$, respectively.

associated with the attracting fixed point $\mathcal{O}_{\mathcal{M}}$ given by $x_{\mathcal{M}} = \mu_{\mathcal{M}} / (1 - a_{\mathcal{M}})$, is confined by three boundaries. Namely, the BCB boundary $\xi_{\mathcal{L}}$, the BCB boundary $\xi_{\mathcal{R}}$, and the DFB boundary

$$\eta_{\mathcal{M}} = \{p : \mu_{\mathcal{R}} = \mu_{\mathcal{L}} + (1 + a_{\mathcal{L}}) d_{\mathcal{L}} - (1 + a_{\mathcal{R}}) d_{\mathcal{R}}\}, \tag{7}$$

related to the condition $a_{\mathcal{M}} = -1$.

The rest of parameter space is occupied by the region

$$D = \{p : \mu_{\mathcal{L}} > (1 - a_{\mathcal{L}}) d_{\mathcal{L}}, \mu_{\mathcal{R}} < (1 - a_{\mathcal{R}}) d_{\mathcal{R}}, \mu_{\mathcal{R}} < \mu_{\mathcal{L}} + (1 + a_{\mathcal{L}}) d_{\mathcal{L}} - (1 + a_{\mathcal{R}}) d_{\mathcal{R}}\}, \tag{8}$$

which is marked by blue line in **Figure 1**. First, we notice that for certain parameter values asymptotic orbits of f are located only on the two adjacent branches. Namely,

1. If $\ell < d_{\mathcal{R}}$ then only $I_{\mathcal{L}}$ and $I_{\mathcal{M}}$ are relevant for asymptotic dynamics of f . Hence, the surfaces $\xi_{\mathcal{L}}$, $\eta_{\mathcal{M}}$, and b_1 (defined from the condition $\ell = d_{\mathcal{R}}$) confine the region

$$D_1 = \{p : (1 - a_{\mathcal{L}}) d_{\mathcal{L}} < \mu_{\mathcal{L}} < d_{\mathcal{R}} - d_{\mathcal{L}} a_{\mathcal{L}}, \mu_{\mathcal{R}} < \mu_{\mathcal{L}} + (1 + a_{\mathcal{L}}) d_{\mathcal{L}} - (1 + a_{\mathcal{R}}) d_{\mathcal{R}}\} \tag{9}$$

(outlined red in **Figure 1**). The related bifurcation structure—skew tent map (STM for short) structure—is described in Section 3.

2. Similarly, the condition $r > d_{\mathcal{L}}$ holds in the region

$$D_2 = \{p : d_{\mathcal{L}} - d_{\mathcal{R}} a_{\mathcal{R}} < \mu_{\mathcal{R}} < (1 - a_{\mathcal{R}}) d_{\mathcal{R}}, \mu_{\mathcal{R}} < \mu_{\mathcal{L}} + (1 + a_{\mathcal{L}}) d_{\mathcal{L}} - (1 + a_{\mathcal{R}}) d_{\mathcal{R}}\}, \tag{10}$$

(outlined yellow in **Figure 1**). This region, also having the STM structure, is confined by $\xi_{\mathcal{R}}$, $\eta_{\mathcal{M}}$, and the boundary b_2

corresponding to the condition $r = d_{\mathcal{L}}$. For $p \in D_2$ only $I_{\mathcal{M}}$ and $I_{\mathcal{R}}$ are relevant for asymptotic dynamics of f (see Section 3).

Finally, the conditions $\ell > d_{\mathcal{R}}$ and $r < d_{\mathcal{L}}$ define the region

$$D_0 = \{p : \mu_{\mathcal{R}} < d_{\mathcal{L}} - d_{\mathcal{R}}a_{\mathcal{R}}, \mu_{\mathcal{L}} > d_{\mathcal{R}} - d_{\mathcal{L}}a_{\mathcal{L}}\}, \quad (11)$$

delineated green in **Figure 1**). For parameter values belonging to D_0 , orbits of f are located on either $I_{\mathcal{L}}$ and $I_{\mathcal{R}}$, or involve all three partitions. Two basic bifurcation structures observed in the region D_0 are period adding (PA for short) structure and fin structure, described in Section 4.

Figure 2A presents an example of the typical bifurcation structure related to regular dynamics in the $(\mu_{\mathcal{L}}, \mu_{\mathcal{R}})$ parameter plane of the map f (the colors correspond to periods of cycles as indicated in the color bar). In **Figure 2B**, on the contrary, colored zones correspond to chaoticity regions related to chaotic attractors with different number of intervals (colors correspond to the number of intervals), while periodicity regions are painted gray.

In the following sections we characterize different bifurcation structures observed in the parameter space of the map (1).

3. SKEW TENT MAP STRUCTURE

Let $p \in D_1$ or $p \in D_2$, then asymptotic dynamics of the map (1) is associated with two adjacent partitions only. In such a case the map f is locally topologically conjugate to the famous skew tent map, which was extensively studied already three decades ago (see e.g., [7–10]). Below we summarize main facts concerning this bifurcation structure, which we refer to as the *skew tent map (STM) structure*.

Let us consider the skew tent map family $g_{\alpha,\beta,\mu}$

$$g_{\alpha,\beta,\mu} : x \rightarrow g_{\alpha,\beta,\mu}(x) = \begin{cases} g_{\mathcal{A}}(x) = \alpha x + \mu, & x \leq 0, \\ g_{\mathcal{B}}(x) = \beta x + \mu, & x > 0, \end{cases} \quad (12)$$

with $\alpha\beta < 0$. All maps from this family with $\mu > 0$ ($\mu < 0$) are topologically equivalent to $g_{\alpha,\beta,1}$ ($g_{\alpha,\beta,-1}$). Moreover, $g_{\alpha,\beta,-1}$ is topologically equivalent to $g_{\beta,\alpha,1}$. Hence, to describe the STM structure it is enough to fix $\mu = 1$. Non-trivial dynamics in this case is observed for $\alpha > 0, \beta < 0$. Below, for shortness, we drop the indices denoting the parameters and refer to the skew tent map simply as g .

Any point $x < 0$ is associated with the symbol \mathcal{A} , while a point $x > 0$ corresponds to the symbol \mathcal{B} . The critical point is denoted as $c := g(0)$ and its images as $c_i := g^i(c)$.

The stability region of the fixed point $\mathcal{O}_{\mathcal{B}}$ is confined by the DB1 boundary $\theta_{\mathcal{B}} = \{(\alpha, \beta) : \beta = 1\}$ and the DFB boundary $\eta_{\mathcal{B}} = \{(\alpha, \beta) : \beta = -1\}$.

The only possible stable cycles of g are the *basic* n -cycles $\mathcal{O}_{\mathcal{B}\mathcal{A}^{n-1}}$, $n \geq 2$. Each periodicity region $\mathcal{P}_{\mathcal{B}\mathcal{A}^{n-1}}$ is confined by the BCB boundary and the DFB boundary

$$\begin{aligned} \xi_{\mathcal{A}\mathcal{B}\mathcal{A}^{n-2}} &= \left\{ (\alpha, \beta) : \beta = -\frac{1 - \alpha^{n-1}}{(1 - \alpha)\alpha^{n-2}} \right\}, \\ \eta_{\mathcal{B}\mathcal{A}^{n-1}} &= \left\{ (\alpha, \beta) : \beta = -\frac{1}{\alpha^{n-1}} \right\} \end{aligned} \quad (13)$$

(see **Figure 3** where the boundaries $\xi_{\mathcal{A}\mathcal{B}\mathcal{A}}$ and $\eta_{\mathcal{B}\mathcal{A}^2}$ are marked). The BCB boundary $\xi_{\mathcal{A}\mathcal{B}}$ is particular, because it coincides with the DFB boundary $\eta_{\mathcal{B}}$. For every $n \geq 3$ the boundary $\xi_{\mathcal{A}\mathcal{B}\mathcal{A}^{n-2}}$ corresponds to the fold BCB leading to the appearance of the basic cycle $\mathcal{O}_{\mathcal{B}\mathcal{A}^{n-1}}$ and its *complementary* cycle $\mathcal{O}_{\mathcal{B}^2\mathcal{A}^{n-2}}$ which is necessarily unstable. The boundary $\eta_{\mathcal{B}\mathcal{A}^{n-1}}$, $n \geq 3$, is related to the DFB of the cycle $\mathcal{O}_{\mathcal{B}\mathcal{A}^{n-1}}$, which leads to the appearance of $2n$ -cyclic chaotic intervals \mathcal{Q}_{2n} . The DFB of the fixed point $\mathcal{O}_{\mathcal{B}}$ and the 2-cycle $\mathcal{O}_{\mathcal{B}\mathcal{A}}$ are particular as described below.

The chaoticity region \mathcal{C}_{2n} , $n \geq 3$, associated with \mathcal{Q}_{2n} , is the area restricted between $\xi_{\mathcal{A}\mathcal{B}\mathcal{A}^{n-2}}$, $\eta_{\mathcal{B}\mathcal{A}^{n-1}}$, and $\gamma_{\mathcal{B}\mathcal{A}^{n-1}}^{c_{2n}}$. The latter one is related to the *first homoclinic bifurcation* of the cycle $\mathcal{O}_{\mathcal{B}\mathcal{A}^{n-1}}$. This bifurcation occurs when

$$c_{2n} = x_{\mathcal{B}\mathcal{A}^{n-1}} \quad \text{with} \quad c_{2n} = g_{\mathcal{B}\mathcal{A}^{n-1}\mathcal{B}\mathcal{A}^{n-2}\mathcal{B}}(c). \quad (14)$$

The condition (14) holds for

$$\gamma_{\mathcal{B}\mathcal{A}^{n-1}}^{c_{2n}} = \left\{ (\alpha, \beta) : \alpha^{2(n-1)}\beta^3 - \beta + \alpha = 0 \right\}. \quad (15)$$

Crossing the bifurcation boundary $\gamma_{\mathcal{B}\mathcal{A}^{n-1}}^{c_{2n}}$ leads to the transition $\mathcal{Q}_{2n} \Rightarrow \mathcal{Q}_n$, which is also called the *merging bifurcation*. The existence region for n -cyclic chaotic intervals \mathcal{Q}_n is \mathcal{C}_n confined by $\xi_{\mathcal{A}\mathcal{B}\mathcal{A}^{n-2}}$, boundary $\gamma_{\mathcal{B}\mathcal{A}^{n-1}}^{c_{2n}}$, and the boundary $\zeta_{\mathcal{B}\mathcal{A}^{n-2}\mathcal{B}}^{c_n}$ corresponding to the *first homoclinic bifurcation* of the cycle $\mathcal{O}_{\mathcal{B}^2\mathcal{A}^{n-2}}$. This homoclinic bifurcation occurs when

$$c_n = g_{\mathcal{B}\mathcal{A}^{n-1}}(c) = x_{\mathcal{B}\mathcal{A}^{n-2}\mathcal{B}}, \quad (16)$$

which gives

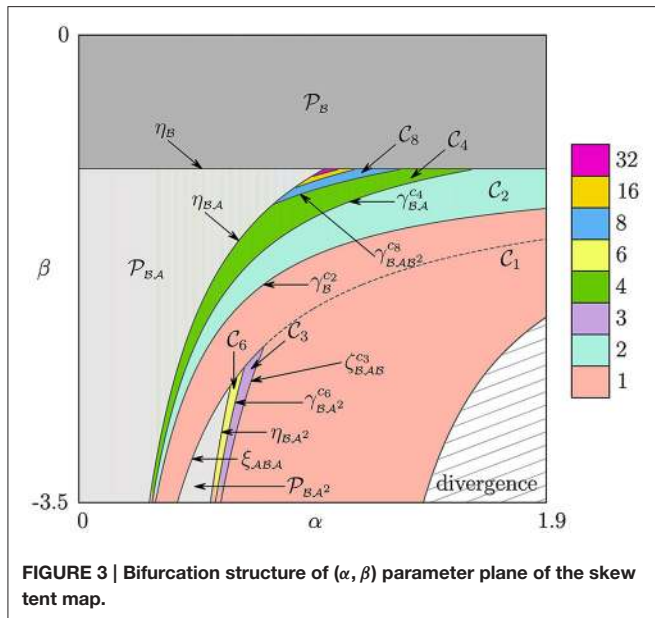
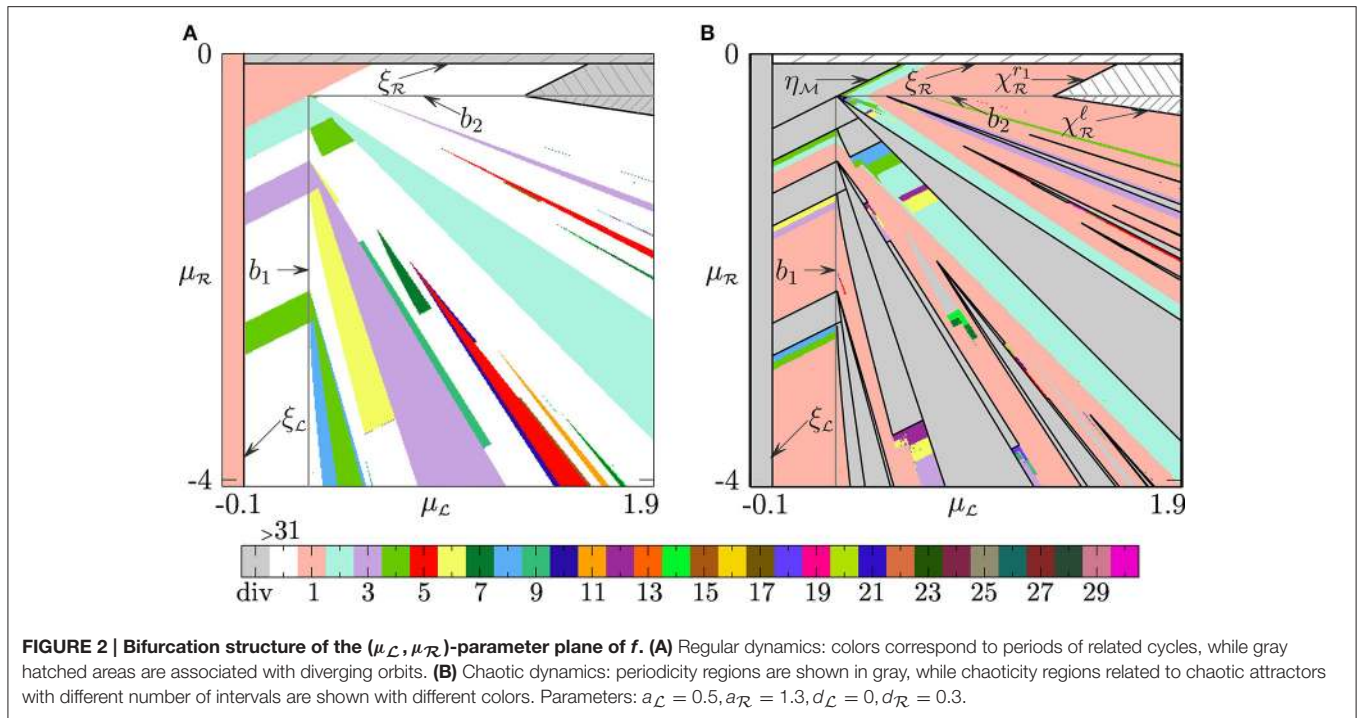
$$\zeta_{\mathcal{B}\mathcal{A}^{n-2}\mathcal{B}}^{c_n} = \left\{ (\alpha, \beta) : \alpha^{n-1}\beta^2 + \beta - \alpha = 0 \right\}. \quad (17)$$

Crossing the boundary $\zeta_{\mathcal{B}\mathcal{A}^{n-2}\mathcal{B}}^{c_n}$ leads to the transition $\mathcal{Q}_n \Rightarrow \mathcal{Q}$, which is also called the *expansion bifurcation*. For more detailed description of merging and expansion bifurcations we refer to Avrutin et al. [18]. In **Figure 3** the boundaries $\gamma_{\mathcal{B}\mathcal{A}^2}^{c_6}$ and $\zeta_{\mathcal{B}\mathcal{A}\mathcal{B}}^{c_3}$ are marked.

Cases $n = 1$ and $n = 2$ differ from the others. The DFB of the fixed point $\mathcal{O}_{\mathcal{B}}$ leads either to an attracting 2-cycle or to 2^m -cyclic chaotic intervals, where $m \geq 1$ depends on particular parameter values. As for the 2-cycle $\mathcal{O}_{\mathcal{A}\mathcal{B}}$, when it loses stability due to DFB, there can appear 2^m -cyclic chaotic intervals \mathcal{Q}_{2^m} , where $m \geq 2$ depends on α, β . Every region \mathcal{C}_{2^m} is adjacent to $\mathcal{C}_{2^{m-1}}$ from one side and to $\mathcal{C}_{2^{m+1}}$ from the other side. The boundary separating \mathcal{C}_{2^m} and $\mathcal{C}_{2^{m+1}}$ is associated with the first homoclinic bifurcation of the 2^m -cycle whose symbolic sequence is the m -th harmonic of \mathcal{B} . Recall that to obtain the k -harmonic ρ_k of \mathcal{B} one has to apply concatenation operator iteratively as follows:

$$\rho_0 := \mathcal{B}, \quad \rho'_0 := \mathcal{A}, \quad \rho_k = \rho_{k-1}\rho'_{k-1}, \quad k \geq 1, \quad (18)$$

where ρ'_{k-1} differs from ρ_{k-1} only by the last symbol. For example, $\rho_1 = \mathcal{B}\mathcal{A}$, $\rho_2 = \mathcal{B}\mathcal{A}\mathcal{B}\mathcal{B}$, $\rho_3 = \mathcal{B}\mathcal{A}\mathcal{B}\mathcal{B}\mathcal{B}\mathcal{A}\mathcal{B}\mathcal{A}$, and so on. Notice that according to the rule (18) it is possible to define harmonics of an arbitrary cycle \mathcal{O}_{σ} provided that x_{σ} is the rightmost point of the cycle. For instance, in the condition



(14) for the first homoclinic bifurcation of $\mathcal{O}_{\mathcal{B}A^{n-1}}$ the symbolic sequence $\mathcal{B}A^{n-1}\mathcal{B}A^{n-2}\mathcal{B}$ is the first harmonic of $\mathcal{B}A^{n-1}$.

The condition for the first homoclinic bifurcation of the 2^m -cycle \mathcal{O}_{ρ_m} is

$$g_{\rho_{m+1}}(c) = x_{\rho_m}, \tag{19}$$

that holds for

$$\gamma_{\rho_m}^{c_{2^{m+1}}} = \left\{ (\alpha, \beta) : \alpha^{2\delta_m} \beta^{2\delta_{m+1}} + \left(\frac{\alpha}{\beta}\right)^{(-1)^{m+1}} - 1 = 0 \right\}, \tag{20}$$

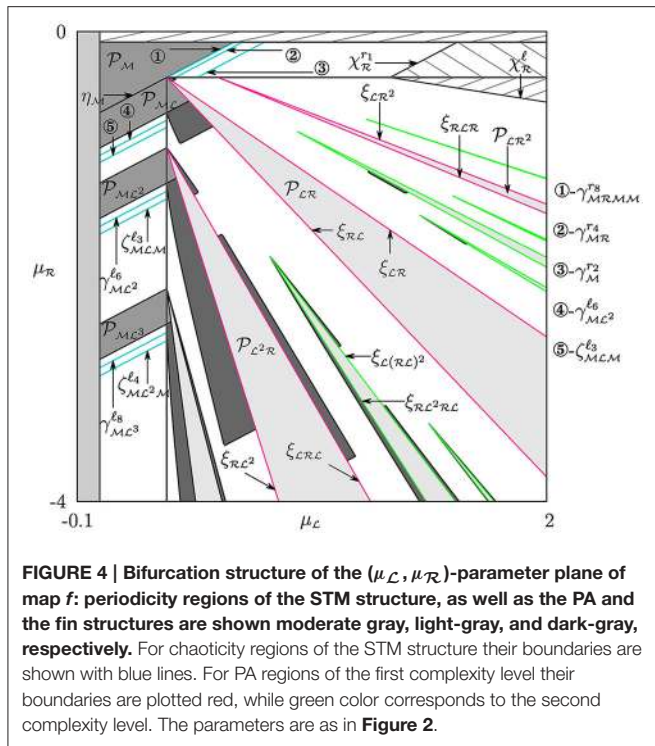
where $\delta_m = (2^m - (-1)^m)/3, m \geq 1$. As mentioned above this bifurcation leads to the transition $\mathcal{Q}_{2^{m+1}} \Rightarrow \mathcal{Q}_{2^m}$. The case $m = 0$ (i.e., the merging bifurcation $\gamma_{\rho_0}^{c_0^2}$) is related to the first homoclinic bifurcation of the fixed point $\mathcal{O}_{\mathcal{B}}$ leading to the transition $\mathcal{Q}_2 \Rightarrow \mathcal{Q}$. For $m \rightarrow \infty$ the curves $\gamma_{\rho_m}^{c_{2^{m+1}}}$ in the (α, β) -plane crowd the point $(\alpha, \beta) = (1, -1)$ (see Figure 3).

The region related to diverging orbits (hatched in Figure 3) is confined by the boundary related to the first homoclinic bifurcation of the fixed point $\mathcal{O}_{\mathcal{A}}$, which occurs when $c_1 = g_{\mathcal{B}}(c) = x_{\mathcal{A}}$:

$$x_{\mathcal{A}}^{c_1} = \left\{ (\alpha, \beta) : \beta = \frac{\alpha}{1 - \alpha} \right\}. \tag{21}$$

At this bifurcation the absorbing interval $J = [c_1, c]$ collides with a boundary of its basin of attraction (given by $\mathcal{O}_{\mathcal{A}}$ and its preimage) and for $\beta < \frac{\alpha}{1 - \alpha}$ the interval J is no longer absorbing causing a typical orbit to diverge. Such bifurcation is also called a *final bifurcation*.

For $p \in D_1$ or $p \in D_2$ the original map f is locally topologically conjugate to g with $\beta = a_{\mathcal{M}}, \alpha = a_{\mathcal{L}}$ or $\alpha = a_{\mathcal{R}}$, respectively, and $\mu = 1$ in the neighborhood of $d_{\mathcal{L}}$ or $d_{\mathcal{R}}$, respectively. Thus, to get the boundaries of STM-regions in the parameter space of f it is enough to substitute \mathcal{L} or \mathcal{R} , respectively, instead of \mathcal{A} and \mathcal{M} instead of \mathcal{B} . In Figure 4 the STM-regions in the $(\mu_{\mathcal{L}}, \mu_{\mathcal{R}})$ parameter plane associated with attracting cycles are colored with moderate gray, while for chaoticity regions only their boundaries are shown (blue lines). The hatched region in D_2 is related to diverging orbits.



4. BIFURCATION STRUCTURES INVOLVING BOTH BORDER POINTS: REGULAR DYNAMICS

Let us briefly recall how other two periodicity structures are formed in the parameter space of f for $p \in D_0$.

4.1. Period Adding Structure

Consider a cycle \mathcal{O}_σ of the map f with $a_L > 0, a_R > 0$ such that all its points belong only to $I_L \cup I_R$. Consequently, σ consists only of the symbols \mathcal{L} and \mathcal{R} . Then the related periodicity region \mathcal{P}_σ is a part of the so-called *period adding structure* (PA structure). Such structures are also known as *Arnold tongues* or *mode-locking tongues* and often appear in the parameter space of a certain class of circle maps, 1D discontinuous maps, two- and higher-dimensional maps near a Neimark-Sacker bifurcation boundary, etc. (see, for example [3, 24–27]). We recall that the regions belonging to PA structure are located in the parameter space according to a specific order based on Farey summation rule. Namely, consider two cycles \mathcal{O}_{σ_1} of period n_1 and \mathcal{O}_{σ_2} of period n_2 with σ_i containing m_i symbols \mathcal{L} . Suppose also that the fractions m_1/n_1 and m_2/n_2 , called also *rotation numbers*, are Farey neighbors, that is, $|m_1n_2 - m_2n_1| = 1$. Then in the parameter space between the two regions related to \mathcal{O}_{σ_1} and \mathcal{O}_{σ_2} there is a region related to a cycle \mathcal{O}_{σ_3} of periodicity $n_1 + n_2$ with σ_3 having $m_1 + m_2$ symbols \mathcal{L} , and thus, the rotation number of \mathcal{O}_{σ_3} is $(m_1 + m_2)/(n_1 + n_2)$. Notice that the cycles $\mathcal{O}_{\sigma_i}, i = 1, 2, 3$, are not necessarily attracting.

In Figure 4 the regions composing the PA structure are shown light-gray.

If $0 < a_L, a_R < 1$ then any PA-cycle is globally attracting inside the region of its existence. Indeed, in this case for $p \in D_0$ there always exists an invariant absorbing interval $J = [r, \ell]$, and hence, there are no divergent orbits. Moreover, we consider an arbitrary PA-cycle \mathcal{O}_σ and immediately get that its multiplier is $\nu_\sigma = a_L^m a_R^{n-m} \in (0, 1)$, where m is the number of \mathcal{L} 's in σ .

However, if one of the slopes is greater than one, while the other is between zero and one, the inequality $\nu_\sigma < 1$ is satisfied only if $n < m(1 - \log_{a_R} a_L)$, so that not all PA-cycles are attracting. There can also be divergent orbits since there is an unstable fixed point on the branch whose slope is greater than one. Furthermore, the interval J is no longer absorbing after the final bifurcation corresponding to the first homoclinic bifurcation of this unstable fixed point. The related bifurcation boundary confines the region in the parameter space in which a typical orbit diverges. For instance, for $0 < a_L < 1, a_R > 1$ (as in Figure 4) the final bifurcation χ_R^ℓ occurs when $x_R = \ell$, and thus, for any $(1 - a_R)(a_L d_L + \mu_L) < \mu_R < d_L - a_R d_R$ all orbits go to infinity.

To simplify calculation of analytic expressions for PA-region boundaries we use the so-called *map replacement technique* [28]. At first we group symbolic sequences of PA-cycles into infinite number of families complying with their *complexity level* (as in [3]). For example, two families related to basic cycles form the *first complexity level*:

$$\Sigma_{1,1} = \{\mathcal{LR}^{n_1}\}_{n_1=1}^\infty, \quad \Sigma_{2,1} = \{\mathcal{RL}^{n_1}\}_{n_1=1}^\infty. \tag{22}$$

Notice that both families contain the common sequence $\mathcal{LR} \equiv \mathcal{RL}$.

To construct the rest of symbolic sequence families we introduce the following *symbolic replacements*:

$$\kappa_m^{\mathcal{L}} := \begin{cases} \mathcal{L} \rightarrow \mathcal{LR}^m \\ \mathcal{R} \rightarrow \mathcal{RL}^m \end{cases}, \quad \kappa_m^{\mathcal{R}} := \begin{cases} \mathcal{L} \rightarrow \mathcal{LR}^m \\ \mathcal{R} \rightarrow \mathcal{RL}^m \end{cases}. \tag{23}$$

Application of a symbolic replacement $\kappa_m^{\mathcal{L}}$ to a symbolic sequence σ means that each symbol \mathcal{L} in σ is replaced by \mathcal{LR}^m and each symbol \mathcal{R} by \mathcal{RL}^m . Similarly, application of $\kappa_m^{\mathcal{R}}$ means replacing all \mathcal{L} 's by \mathcal{LR}^m and all \mathcal{R} 's by \mathcal{RL}^m .

Symbolic sequences of higher complexity levels are obtained by iterative application of $\kappa_m^{\mathcal{L}}$ and $\kappa_m^{\mathcal{R}}$. More precisely, applying $\kappa_{n_2}^{\mathcal{L}}$ and $\kappa_{n_2}^{\mathcal{R}}$ to $\Sigma_{1,1}$ we get the families

$$\begin{aligned} \Sigma_{1,2} &= \{\mathcal{LR}^{n_2} (\mathcal{RL}^m)^{n_1}\}_{n_1, n_2=1}^\infty, \\ \Sigma_{2,2} &= \{\mathcal{LR}^{n_2} (\mathcal{RL}^{n_2})^{n_1}\}_{n_1, n_2=1}^\infty. \end{aligned} \tag{24}$$

And then applying $\kappa_{n_2}^{\mathcal{L}}$ and $\kappa_{n_2}^{\mathcal{R}}$ to $\Sigma_{2,1}$ we obtain

$$\begin{aligned} \Sigma_{3,2} &= \{\mathcal{RL}^{n_2} (\mathcal{LR}^{n_2})^{n_1}\}_{n_1, n_2=1}^\infty, \\ \Sigma_{4,2} &= \{\mathcal{RL}^{n_2} (\mathcal{LR}^{n_2})^{n_1}\}_{n_1, n_2=1}^\infty. \end{aligned} \tag{25}$$

The four families $\Sigma_{i,2}, i = \overline{1,4}$, compose the *second complexity level*. Notice that the families $\Sigma_{1,2}$ and $\Sigma_{3,2}$ contain the common sequence $\mathcal{LRRLR} \equiv \mathcal{RLRLR}$ with $n_1 = n_2 = 1$, while

$\Sigma_{2,2}$ and $\Sigma_{4,2}$ share the sequence $\mathcal{L}\mathcal{R}\mathcal{L}\mathcal{R}\mathcal{L} \equiv \mathcal{R}\mathcal{L}\mathcal{L}\mathcal{R}\mathcal{L}$ with $n_1 = n_2 = 1$.

In short the applied procedure can be written as

$$\begin{aligned} \Sigma_{1,2} &= \kappa_{n_2}^{\mathcal{L}}(\Sigma_{1,1}), \quad \Sigma_{2,2} = \kappa_{n_2}^{\mathcal{R}}(\Sigma_{1,1}), \\ \Sigma_{3,2} &= \kappa_{n_2}^{\mathcal{L}}(\Sigma_{2,1}), \quad \Sigma_{4,2} = \kappa_{n_2}^{\mathcal{R}}(\Sigma_{2,1}). \end{aligned} \tag{26}$$

Further, applying the replacements (23) with $m = n_3$ to four families of the second complexity level, we obtain 2^3 families $\Sigma_{j,3}$, $j = \overline{1, 2^3}$, of the *third complexity level*, and so on. In this way all symbolic sequences of PA-cycles are obtained.

Now let us identify the PA-regions in the parameter space of the map f . Each such region has two boundaries corresponding to BCBs of the related cycle. For the basic cycles $\mathcal{O}_{\mathcal{L}\mathcal{R}^{n_1}}$, $n_1 \geq 1$, these boundaries are obtained directly from the relevant BCB conditions giving

$$\begin{aligned} \xi_{\mathcal{L}\mathcal{R}^{n_1}} &= \{p \in D_0 : \mu_{\mathcal{L}} = \Phi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{L}}, n_1)\}, \\ \xi_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_1-1}} &= \{p \in D_0 : \mu_{\mathcal{L}} = \Psi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{R}}, n_1)\}, \end{aligned} \tag{27}$$

where

$$\Phi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d, n_1) = -\psi(a_{\mathcal{R}}, n_1)\mu_{\mathcal{R}} + \varphi(a_{\mathcal{R}}, a_{\mathcal{L}}, n_1)d, \tag{28}$$

$$\begin{aligned} \Psi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d, n_1) &= -(a_{\mathcal{L}} + \psi(a_{\mathcal{R}}, n_1 - 1))\mu_{\mathcal{R}} \\ &+ a_{\mathcal{R}}\varphi(a_{\mathcal{R}}, a_{\mathcal{L}}, n_1)d \end{aligned} \tag{29}$$

with

$$\varphi(a, b, n) = \frac{1 - a^n b}{a^n}, \quad \psi(a, n) = \frac{1 - a^n}{(1 - a)a^n}. \tag{30}$$

For $\mathcal{O}_{\mathcal{R}\mathcal{L}^{n_1}}$, $n \geq 1$, their boundaries are obtained from (27) by interchanging \mathcal{L} and \mathcal{R} . So that, the basic PA-regions are defined as

$$\mathcal{P}_{\mathcal{L}\mathcal{R}^{n_1}} = \{p \in D_0 : \Psi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{R}}, n_1) < \mu_{\mathcal{L}} < \Phi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{L}}, n_1)\}, \tag{31}$$

$$\mathcal{P}_{\mathcal{R}\mathcal{L}^{n_1}} = \{p \in D_0 : \Psi_{1,1}(a_{\mathcal{R}}, a_{\mathcal{L}}, \mu_{\mathcal{L}}, d_{\mathcal{L}}, n_1) > \mu_{\mathcal{R}} > \Phi_{1,1}(a_{\mathcal{R}}, a_{\mathcal{L}}, \mu_{\mathcal{L}}, d_{\mathcal{R}}, n_1)\}. \tag{32}$$

Further, using the map replacement technique we obtain periodicity regions associated with cycles whose symbolic sequences belong to the families $\Sigma_{1,2}$ and $\Sigma_{2,2}$, given in (24). The trick is as follows. Any symbolic sequence $\tau \in \Sigma_{1,2}$ can be written as $\tau = \kappa_{n_2}^{\mathcal{L}}(\sigma)$ with certain $n_2 \geq 1$ and $\sigma \in \Sigma_{1,1}$ with particular $n_1 \geq 1$. Then to get the BCBs $\xi_{\tau}^{d_{\mathcal{L}}}$ and $\xi_{\tau}^{d_{\mathcal{R}}}$ we take the expressions (27) for $\xi_{\mathcal{L}\mathcal{R}^{n_1}}$ and $\xi_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_1-1}}$ and substitute the parameters $a_{\mathcal{L}}, \mu_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}$ by $a_{\mathcal{L}\mathcal{R}^{n_2}}, \mu_{\mathcal{L}\mathcal{R}^{n_2}}, a_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, \mu_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}$:

$$\begin{aligned} \xi_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{L}\mathcal{R}^{n_1})} &= \{p \in D_0 : \mu_{\mathcal{L}\mathcal{R}^{n_2}} = \Phi_{1,1}(a_{\mathcal{L}\mathcal{R}^{n_2}}, a_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, \\ &\quad \mu_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, d_{\mathcal{L}}, n_1)\}, \\ \xi_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{R}\mathcal{L}\mathcal{R}^{n_1-1})} &= \{p \in D_0 : \mu_{\mathcal{L}\mathcal{R}^{n_2}} = \Psi_{1,1}(a_{\mathcal{L}\mathcal{R}^{n_2}}, a_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, \\ &\quad \mu_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, d_{\mathcal{R}}, n_1)\}. \end{aligned} \tag{33}$$

Recall that $a_{\mathcal{L}\mathcal{R}^{n_2}}, \mu_{\mathcal{L}\mathcal{R}^{n_2}}$ and $a_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, \mu_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}$ denote, respectively, the coefficients of the following composite (auxiliary) functions

$$f_{\mathcal{L}\mathcal{R}^{n_2}} = f_{\mathcal{L}} \circ \underbrace{f_{\mathcal{R}} \circ \dots \circ f_{\mathcal{R}}}_{n_2}, \quad f_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}} = f_{\mathcal{R}} \circ \underbrace{f_{\mathcal{L}} \circ f_{\mathcal{R}} \circ \dots \circ f_{\mathcal{R}}}_{n_2-1}.$$

Schematically, we summarize the applied procedure as follows:

$$\begin{aligned} \xi_{\mathcal{L}\mathcal{R}^{n_1}} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{L}\mathcal{R}^{n_2}]{\mathcal{R} \rightarrow \mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}} \xi_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{L}\mathcal{R}^{n_1})}, \\ \xi_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_1-1}} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{L}\mathcal{R}^{n_2}]{\mathcal{R} \rightarrow \mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}} \xi_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{R}\mathcal{L}\mathcal{R}^{n_1-1})}. \end{aligned} \tag{34}$$

Clearly, the coefficients $a_{\mathcal{L}\mathcal{R}^{n_2}}, \mu_{\mathcal{L}\mathcal{R}^{n_2}}, a_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}, \mu_{\mathcal{R}\mathcal{L}\mathcal{R}^{n_2-1}}$ depend on $a_{\mathcal{L}}, \mu_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}$, and hence, the equalities in (33) can be solved with respect to $\mu_{\mathcal{L}}$, implying the PA-region related to $\mathcal{O}_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{L}\mathcal{R}^{n_1})}$ to be

$$\begin{aligned} \mathcal{P}_{\kappa_{n_2}^{\mathcal{L}}(\mathcal{L}\mathcal{R}^{n_1})} &= \{p \in D_0 : \Psi_{1,2}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{R}}, n_1, n_2) \\ &< \mu_{\mathcal{L}} < \Phi_{1,2}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{L}}, n_1, n_2)\} \end{aligned} \tag{35}$$

where $\Psi_{1,2}, \Phi_{1,2}$ are certain functions (see [16] where expressions for $\Psi_{1,2}, \Phi_{1,2}$ are given in explicit form).

Similarly, any symbolic sequence $\tau \in \Sigma_{2,2}$ can be written as $\tau = \kappa_{n_2}^{\mathcal{R}}(\sigma)$ with $n_2 \geq 1$ and $\sigma \in \Sigma_{1,1}$. To get the expressions for the boundaries of \mathcal{P}_{τ} we perform analogous trick replacing in (27) the parameters $a_{\mathcal{L}}, \mu_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}$ with the coefficients $a_{\mathcal{L}\mathcal{R}\mathcal{L}^{n_2-1}}, \mu_{\mathcal{L}\mathcal{R}\mathcal{L}^{n_2-1}}, a_{\mathcal{R}\mathcal{L}\mathcal{L}^{n_2}}, \mu_{\mathcal{R}\mathcal{L}\mathcal{L}^{n_2}}$ of the related (auxiliary) composite functions. The resulting new equalities can be resolved with respect to $\mu_{\mathcal{L}}$, and the related PA-region is

$$\begin{aligned} \mathcal{P}_{\kappa_{n_2}^{\mathcal{R}}(\mathcal{L}\mathcal{R}^{n_1})} &= \{p \in D_0 : \Psi_{2,2}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{R}}, n_1, n_2) < \mu_{\mathcal{L}} \\ &< \Phi_{2,2}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{L}}, n_1, n_2)\}. \end{aligned} \tag{36}$$

with certain $\Psi_{2,2}, \Phi_{2,2}$.

As for the cycles whose symbolic sequences belong to the families $\Sigma_{3,2}$ and $\Sigma_{4,2}$ (25), boundaries of the related periodicity regions are obtained by interchanging the indices \mathcal{L} and \mathcal{R} and changing the inequality signs to the opposite ones in (36) and (35), respectively.

Similarly, having expressions for the boundaries of periodicity regions of the $(k - 1)$ -th complexity level and using this replacement technique, we can compute the related expressions for periodicity regions of the k -th complexity level (for more detail see [16]).

In **Figure 4** for the regions of the first complexity level their boundaries are colored red, while for those of the second complexity level they are green.

4.2. Fin Structure

In **Figure 2** or in **Figure 4** one can notice that some PA-regions have other (smaller) periodicity regions attached to them, which resemble to certain extent fish fins. Every such fin-like region is associated with a stable kn -cycle, where n is the period of the related PA-cycle and $k \geq 1$. To explain briefly how these regions appear, let us consider the n -th iterate of f .

For the map f^n every point of the PA-cycle, say \mathcal{O}_σ , of the period n is a fixed point. Suppose that the parameter point $p \in \mathcal{P}_\sigma$ moves outside the region crossing transversally, for instance, the boundary $\xi_\sigma^{d\mathcal{L}}$ (see **Figure 5**). Near the BCB the map f^n is locally topologically conjugate in the neighborhood of $d_{\mathcal{L}}$ to a skew tent map $g_{\alpha,\beta,\mu}$ (12) with $\alpha = a_{\mathcal{L}}^l a_{\mathcal{R}}^{n-l}$, $\beta = a_{\mathcal{L}}^{l-1} a_{\mathcal{M}} a_{\mathcal{R}}^{n-l-1}$ (l is the number of symbols \mathcal{L} in σ) and $\mu < 0$ before, $\mu = 0$ at, and $\mu > 0$ after the bifurcation. In other words, the skew tent map is used as the BCB normal form, and the result of such a bifurcation is completely defined by the values of α and β (see, e.g., [5, 10, 11]). In particular, for certain combination of $a_{\mathcal{L}}$, $a_{\mathcal{M}}$ and $a_{\mathcal{R}}$, a stable k -cycle for f^n appears after the BCB. In terms of the original map f , after the BCB of a stable n -cycle \mathcal{O}_σ there occurs a stable kn -cycle, say, \mathcal{O}_τ . Crossing the boundary $\xi_\sigma^{d\mathcal{R}}$ is

As for the trunk region $\mathcal{P}_{\mathcal{R}\mathcal{L}^{n-1}}$ its $k \cdot n$ -fins correspond to cycles whose symbolic sequences are obtained from (37) and (38) by interchanging symbols \mathcal{L} and \mathcal{R} . Whether a trunk region has fins on both sides or none of them is present depends on the parameters.

Every fin region $\mathcal{P}_{\sigma_{kn}^s}$, $s \in \{\mathcal{L}, \mathcal{R}\}$, $k \geq 2$, obviously has a common BCB boundary with the related trunk region. At most $\mathcal{P}_{\sigma_{kn}^s}$ has three more bifurcation boundaries that correspond to two other BCB's and the DFB of the associated cycle. Analytic expressions for the BCB boundaries are obtained by using the skew tent map as border collision normal form. The DFB boundary is given by condition $a_{\mathcal{L}}^l a_{\mathcal{M}} a_{\mathcal{R}}^{kn-l-1} = -1$, where l is the number of symbols \mathcal{L} in σ_{kn}^s . For left fins $\mathcal{P}_{\sigma_{kn}^{\mathcal{L}}}$ the expressions are (see [16])

$$\xi_{\mathcal{L}\mathcal{R}^{n-1}} = \{p \in D_0 : \mu_{\mathcal{L}} = \Phi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{L}}, n - 1)\}, \tag{39}$$

$$\eta_{\sigma_{kn}^{\mathcal{L}}} = \{p \in D_0 : a_{\mathcal{M}} a_{\mathcal{R}}^{k(n-1)} a_{\mathcal{L}}^{k-1} = -1\}, \tag{40}$$

$$\xi_{\sigma_{kn}^{\mathcal{L}}}^{d\mathcal{L}} = \left\{ p \in D_0 : a_{\mathcal{M}} a_{\mathcal{R}}^{n-1} = \frac{a_{\mathcal{L}} a_{\mathcal{R}}^{n-1} ((a_{\mathcal{L}} a_{\mathcal{R}}^{n-1})^{-k+1} - 1)}{a_{\mathcal{L}} a_{\mathcal{R}}^{n-1} - 1} \right\} \tag{41}$$

$$\xi_{\sigma_{kn}^{\mathcal{L}}}^{d\mathcal{R}} = \left\{ p \in D_0 : a_{\mathcal{R}}^{n-1} a_{\mathcal{M}} \frac{(a_{\mathcal{R}}^{n-1} a_{\mathcal{L}})^{k-1} - 1}{a_{\mathcal{R}}^{n-1} a_{\mathcal{L}} - 1} \left(\left(a_{\mathcal{R}}^{n-2} a_{\mathcal{L}} + \frac{a_{\mathcal{R}}^{n-2} - 1}{a_{\mathcal{R}} - 1} \right) \mu_{\mathcal{R}} + a_{\mathcal{R}}^{n-2} \mu_{\mathcal{L}} \right) + a_{\mathcal{R}}^{(n-1)k} a_{\mathcal{M}} a_{\mathcal{L}}^{(k-1)} d_{\mathcal{R}} + \left(a_{\mathcal{R}}^{n-2} a_{\mathcal{M}} + \frac{a_{\mathcal{R}}^{n-2} - 1}{a_{\mathcal{R}} - 1} \right) \mu_{\mathcal{R}} + a_{\mathcal{R}}^{n-2} \mu_{\mathcal{M}} = d_{\mathcal{R}} \right\}, \tag{42}$$

treated likewise with the only difference that $\beta = a_{\mathcal{L}}^l a_{\mathcal{M}} a_{\mathcal{R}}^{n-l-1}$ (for details see [16]).

We refer to $\mathcal{O}_\tau/\mathcal{P}_\tau$ as the $k \cdot n$ -fin cycle/region, and to $\mathcal{O}_\sigma/\mathcal{P}_\sigma$ as the trunk cycle/region. For clarity sake, the region \mathcal{P}_τ is called the left/right $k \cdot n$ -fin if it is attached to the boundary of \mathcal{P}_σ that is related to collision of \mathcal{O}_σ with the left/right border point $d_{\mathcal{L}}/d_{\mathcal{R}}$.

where $\Phi_{1,1}$ is introduced in (28). Notice that the symbolic sequence of the $1 \cdot n$ -fin cycle is $\sigma_n^{\mathcal{L}} \equiv \mathcal{M}\mathcal{R}^{n-1}$, and hence, the related periodicity region is confined by only three boundaries: $\xi_{\mathcal{L}\mathcal{R}^{n-1}}$ (39), $\eta_{\mathcal{M}\mathcal{R}^{n-1}}$ (40), and $\xi_{\mathcal{R}\mathcal{M}\mathcal{R}^{n-2}}$ (42).

Similarly, the bifurcation boundaries confining the right $k \cdot n$ -fin $\mathcal{P}_{\sigma_{kn}^{\mathcal{R}}}$, contiguous to $\mathcal{P}_{\mathcal{L}\mathcal{R}^{n-1}}$, are

$$\xi_{\mathcal{R}\mathcal{L}\mathcal{R}^{n-2}} = \{p \in D_0 : \mu_{\mathcal{L}} = \Psi_{1,1}(a_{\mathcal{L}}, a_{\mathcal{R}}, \mu_{\mathcal{R}}, d_{\mathcal{R}}, n - 1)\} \tag{43}$$

$$\eta_{\sigma_{kn}^{\mathcal{R}}} = \{p \in D_0 : a_{\mathcal{M}} a_{\mathcal{R}}^{k(n-1)-1} a_{\mathcal{L}}^k = -1\}, \tag{44}$$

$$\xi_{\sigma_{kn}^{\mathcal{R}}}^{d\mathcal{R}} = \left\{ p \in D_0 : a_{\mathcal{L}} a_{\mathcal{M}} a_{\mathcal{R}}^{n-2} = \frac{a_{\mathcal{L}} a_{\mathcal{R}}^{n-1} ((a_{\mathcal{L}} a_{\mathcal{R}}^{n-1})^{-k+1} - 1)}{a_{\mathcal{L}} a_{\mathcal{R}}^{n-1} - 1} \right\}, \tag{45}$$

$$\xi_{\sigma_{kn}^{\mathcal{R}}}^{d\mathcal{L}} = \left\{ p \in D_0 : a_{\mathcal{R}}^{n-2} a_{\mathcal{M}} a_{\mathcal{L}} \frac{(a_{\mathcal{R}}^{n-1} a_{\mathcal{L}})^{k-1} - 1}{a_{\mathcal{R}}^{n-1} a_{\mathcal{L}} - 1} \left(\frac{a_{\mathcal{R}}^{n-1} - 1}{a_{\mathcal{R}} - 1} \mu_{\mathcal{R}} + a_{\mathcal{R}}^{n-1} \mu_{\mathcal{L}} \right) + a_{\mathcal{R}}^{(n-1)k-1} a_{\mathcal{M}} a_{\mathcal{L}}^k d_{\mathcal{L}} + a_{\mathcal{M}} \frac{a_{\mathcal{R}}^{n-2} - 1}{a_{\mathcal{R}} - 1} \mu_{\mathcal{R}} + a_{\mathcal{M}} a_{\mathcal{R}}^{n-2} \mu_{\mathcal{L}} + \mu_{\mathcal{M}} = d_{\mathcal{L}} \right\}, \tag{46}$$

Furthermore, the complexity level of the trunk region defines the complexity level of its fins.

Let us consider $\mathcal{P}_{\mathcal{L}\mathcal{R}^{n-1}}$, $n \geq 2$, of the first complexity level. Its left $k \cdot n$ -fin regions, $k \geq 1$, correspond to cycles whose symbolic sequences are

$$\sigma_{kn}^{\mathcal{L}} := (\mathcal{L}\mathcal{R}^{n-1})^{k-1} \mathcal{M}\mathcal{R}^{n-1}, \tag{37}$$

while symbolic sequences of the cycles associated with right $k \cdot n$ -fin regions are

$$\sigma_{kn}^{\mathcal{R}} := (\mathcal{R}\mathcal{L}\mathcal{R}^{n-2})^{k-1} \mathcal{M}\mathcal{L}\mathcal{R}^{n-2}. \tag{38}$$

where $\Psi_{1,1}$ is given in (29). Again the symbolic sequence of the $1 \cdot n$ -fin cycle is $\sigma_n^{\mathcal{R}} \equiv \mathcal{M}\mathcal{L}\mathcal{R}^{n-2}$, and consequently, the periodicity region $\mathcal{P}_{\sigma_{kn}^{\mathcal{R}}}$ is confined by $\xi_{\mathcal{L}\mathcal{R}^{n-1}}$ (39), $\eta_{\mathcal{M}\mathcal{R}^{n-1}}$ (40), and $\xi_{\mathcal{R}\mathcal{M}\mathcal{R}^{n-2}}$ (42).

In **Figure 5** the region $\mathcal{P}_{\mathcal{L}\mathcal{R}^{n-1}}$ and its fins are shown schematically.

The boundaries of fins contiguous to trunk regions $\mathcal{P}_{\mathcal{R}\mathcal{L}^{n-1}}$ can be obtained by interchanging the indices \mathcal{L} and \mathcal{R} in (39–45). The expressions for the boundaries of fin regions attached to trunk regions of higher complexity levels can be obtained using the map replacement

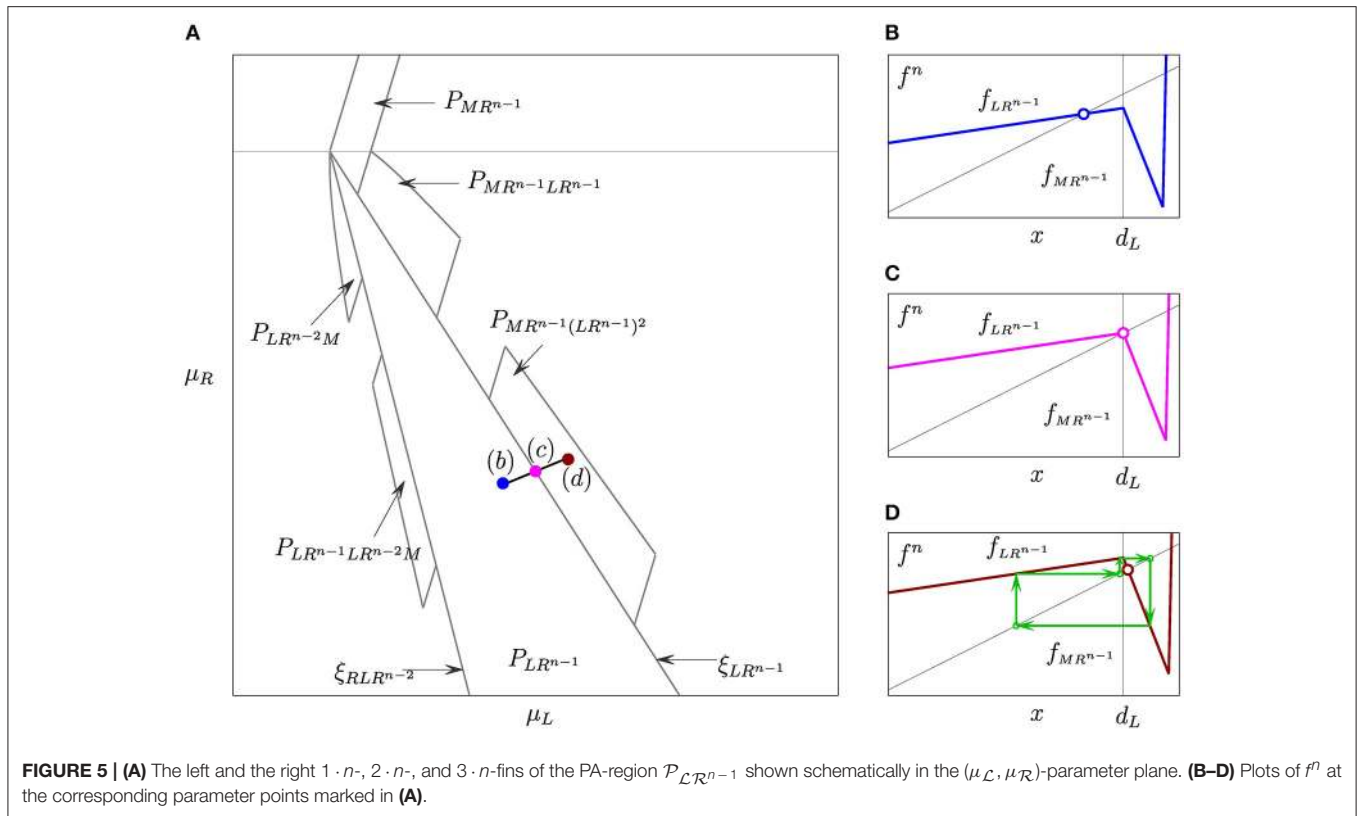


FIGURE 5 | (A) The left and the right $1 \cdot n$ -, $2 \cdot n$ -, and $3 \cdot n$ -fins of the PA-region $\mathcal{P}_{LR^{n-1}}$ shown schematically in the (μ_L, μ_R) -parameter plane. **(B–D)** Plots of f^n at the corresponding parameter points marked in **(A)**.

technique mentioned in Section 4.1 (see [16] for details).

5. EXPANSION OF THE ABSORBING INTERVAL FROM TWO ADJACENT BRANCHES OVER THE THIRD ONE

The present section describes the phenomena that occur when the parameter point p crosses the boundary b_1/b_2 entering D_0 . As mentioned above, for $p \in D_1/D_2$ the invariant absorbing interval J (if it exists) involves only one border point d_L/d_R and is located in two partitions only. At the boundary b_1/b_2 the interval J collides with the second border point and expands into the third partition. This implies modification of bifurcation boundaries confining periodicity and chaoticity regions.

5.1. Prolongation of the Skew Tent Map Structure

Recall that for $p \in D_1$ asymptotic orbits of the map f are located in $I_L \cup I_M$ only. If additionally either $0 \leq a_L \leq 1$, or $a_L > 1$ and $f_M(\ell) > x_L$, then there exists invariant absorbing interval $J = [f_M(d_L), d_L] \subset I_L \cup I_M$. Consequently, all bifurcation boundaries in D_1 engage only the left border point d_L and the critical point ℓ (including its images). Similarly, for $p \in D_2$ the bifurcation conditions are related to the right border point d_R and the critical point r only. Conversely, for $p \in D_0$ the absorbing

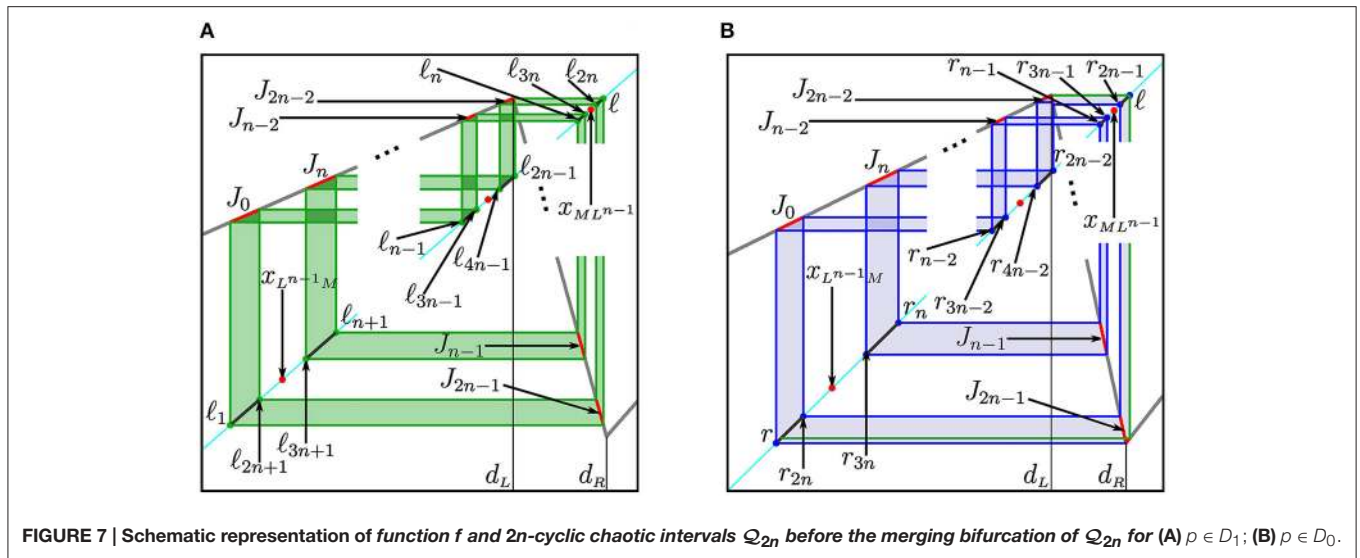
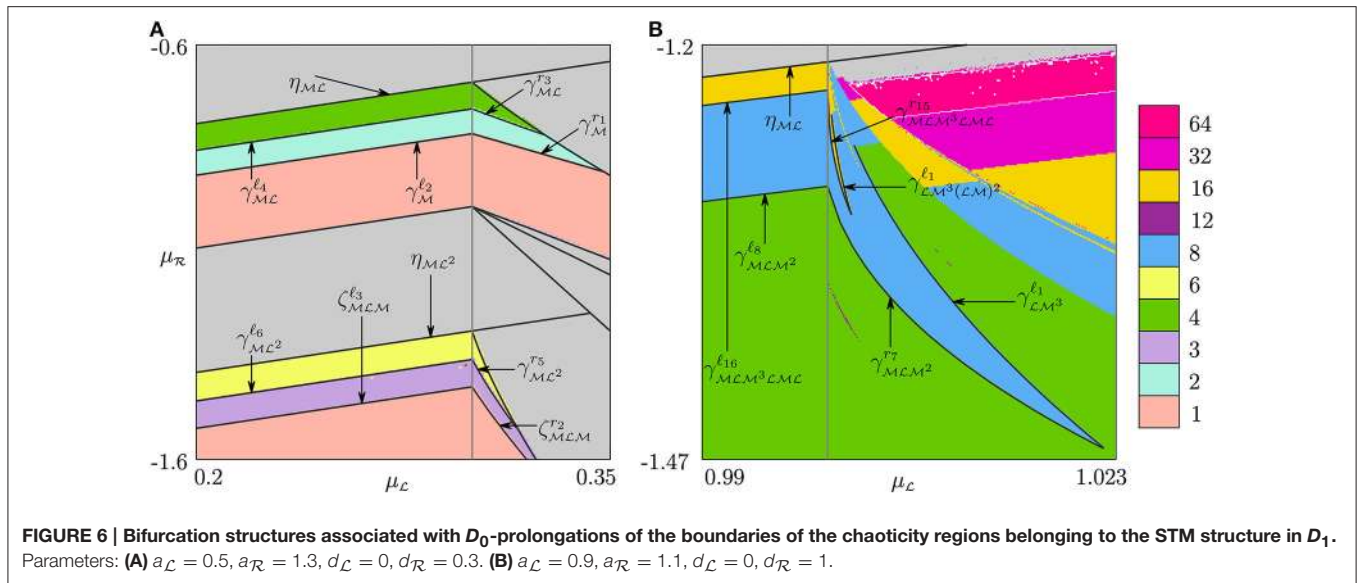
interval J (if existent) extends over all three partitions. Hence, a bifurcation may be associated with any border or critical point.

Let us consider first chaoticity regions C_{2n}, C_n existing after the DFB $\eta_{ML^{n-1}}$ of the attracting cycle $\mathcal{O}_{ML^{n-1}}$, $n \geq 3$ (generic case). We describe how their boundaries change when the parameter point p enters D_0 crossing the border b_1 (see, for instance, the regions C_6 and C_3 shown **Figure 6A**). As described in Section 3, for $p \in D_1$ the DFB of $\mathcal{O}_{ML^{n-1}}$ leads to appearance of $2n$ -cyclic chaotic intervals \mathcal{Q}_{2n} , and thus, the corresponding region C_{2n} is confined by $\eta_{ML^{n-1}}$ and $\gamma_{ML^{n-1}}^{\ell_{2n}}$. However, the part of C_{2n} located in D_0 is confined by the other two boundaries. The first one is the fold BCB boundary $\xi_{LR^{n-1}ML^{n-2}}$, associated with the collision of the fin cycle $\mathcal{O}_{LR^{n-1}ML^{n-1}}$ with d_L . The second boundary is an extension of $\gamma_{ML^{n-1}}^{\ell_{2n}}$ to D_0 , which we refer to as its D_0 -prolongation. This boundary is associated as well with a homoclinic bifurcation of $\mathcal{O}_{ML^{n-1}}$, but its condition differs from the one for $\gamma_{ML^{n-1}}^{\ell_{2n}}$, as we show below.

Figure 7A is a schematic representation of $2n$ -cyclic chaotic intervals $\mathcal{Q}_{2n} = \{J_i\}_{i=0}^{2n-1}$ for $p \in C_{2n} \subset D_1$. It can be shown that the intervals J_i , $i = \overline{0, 2n-2}$, are bounded by ℓ_{i+1} and ℓ_{i+2n+1} and the interval J_{2n-1} by ℓ_{2n} and ℓ . Hence, for $p \in D_1$ the homoclinic bifurcation corresponding to the merging $\mathcal{Q}_{2n} \Rightarrow \mathcal{Q}_n$ can be defined by the following condition:

$$f^{2n}(\ell) \equiv f_{ML^{n-1}ML^{n-2}M}(\ell) = x_{ML^{n-1}} \tag{47}$$

(cf. (14)). For $p \in b_1$, the condition $\ell = d_R$ holds, and hence $\ell_1 = r$, $\ell_{i+1} = r_i$, $i \geq 1$. Therefore, the intervals J_i ,



$i = \overline{0, 2n - 2}$, constituting Q_{2n} are bounded by $l_{i+1} = r_i$ and $l_{i+2n+1} = r_{i+2n}$, while the interval J_{2n-1} is given by $[r_{2n-1}, d_R]$. Hence, the condition (47) becomes

$$\begin{aligned} f^{2n}(\ell) &\equiv f_{ML^{n-1}ML^{n-2}M}(d_R) \equiv f_{L^{n-1}ML^{n-2}M}(r) \\ &\equiv f^{2n-1}(r) = x_{ML^{n-1}}. \end{aligned} \tag{48}$$

As soon as the point p enters the region D_0 by crossing b_1 , the absorbing interval J spreads over all three partitions I_L, I_M , and I_R . For Q_{2n} it means that $J_{2n-1} = [r_{2n-1}, \ell] \subset I_M \cup I_R$ with $d_R \in J_{2n-1}$ (see **Figure 7B**). By definition $f(J_{2n-1}) = J_0$. If p is close to b_1 , the value of ℓ is greater than the value of d_R , but it is still close to d_R , and thus,

$$f(d_R) < f(\ell) < f(r_{2n-1}) \Leftrightarrow r < \ell_1 < r_{2n}.$$

It entails that $J_0 = [f(d_R), f(r_{2n-1})] = [r, r_{2n}]$, while $\ell_1 \in J_0$. As a result, only the right boundary of J_{2n-1} is given by ℓ , and

its left boundary, as well as the boundaries of the intervals $J_i, i = \overline{0, 2n - 2}$, are given by the images of r (of the same rank as for $p \in D_1$). Thus, for $p \in D_0$ the condition for the homoclinic bifurcation leading to the transition $Q_{2n} \Rightarrow Q_n$ coincides with (48). In such a way the D_0 -prolongation of $\gamma_{ML^{n-1}}^{\ell_{2n}}$ is the merging bifurcation boundary

$$\gamma_{ML^{n-1}}^{r_{2n-1}} = \{p \in D_0 : f_{L^{n-1}ML^{n-2}M}(r) = x_{ML^{n-1}}\}. \tag{49}$$

Similarly, we deduce that the chaoticity region C_n consists of two parts, one of which is located in D_1 and the other one in D_0 (see **Figure 6A** for a sample of C_3). The part of C_n located in D_1 is confined by the merging bifurcation boundary $\gamma_{ML^{n-1}}^{\ell_{2n}}$ and the expansion bifurcation boundary $\zeta_{ML^{n-2}M}^{\ell_n}$. The part of C_n located in D_0 is bounded by $\xi_{LR^{n-1}ML^{n-2}}^{\ell_n}$, $\gamma_{ML^{n-1}}^{r_{2n-1}}$ (49) and the D_0 -prolongation of $\zeta_{ML^{n-2}M}^{\ell_n}$, which is associated with

the expansion bifurcation as well and leads to the transition $\mathcal{Q}_n \Rightarrow \mathcal{Q}$. To obtain the condition related to the mentioned D_0 -prolongation we recall that for $p \in b_1$ there holds $\ell = d_{\mathcal{R}}$. There follows that

$$f^n(\ell) \equiv f_{\mathcal{M}\mathcal{L}^{n-1}}(d_{\mathcal{R}}) \equiv f_{\mathcal{L}^{n-1}}(r) \equiv f^{n-1}(r) = x_{\mathcal{M}\mathcal{L}^{n-2}\mathcal{M}}.$$

Hence, for $p \in D_0$ the expansion $\mathcal{Q}_n \Rightarrow \mathcal{Q}$ occurs at

$$\zeta_{\mathcal{M}\mathcal{L}^{n-2}\mathcal{M}}^{r_{n-1}} = \{p \in D_0 : f_{\mathcal{L}^{n-1}}(r) = x_{\mathcal{M}\mathcal{L}^{n-2}\mathcal{M}}\}. \quad (50)$$

Now we turn to chaotic attractors existing after the DFB of $\mathcal{O}_{\mathcal{M}}$ or $\mathcal{O}_{\mathcal{M}\mathcal{L}}$. Although these cases are particular they are treated likewise. The DFB of a stable 2-cycle or a fixed point implies appearance of intervals \mathcal{Q}_{2^m} where $m \geq 1$ depends on parameters. Then there exists a cascade of merging bifurcations ending by the parameter region for that the whole absorbing interval represents a one-piece chaotic attractor. Each boundary $\gamma_{\rho_m}^{\ell_{2^m+1}}$ for $p \in D_1$ has its D_0 -prolongation

$$\gamma_{\rho_m}^{r_{2^m+1-1}} = \{p \in D_0 : f_{\omega_{m+1}}(r) = x_{\rho_m}\} \quad (51)$$

related to the transition $\mathcal{Q}_{2^{m+1}} \Rightarrow \mathcal{Q}_{2^m}$. Here ρ_m is the m -th harmonic of \mathcal{M} (cf. (18) with $\mathcal{A} = \mathcal{L}$ and $\mathcal{B} = \mathcal{M}$), and $\omega_{m+1} = s_2 \dots s_{2^{m+1}}$ such that $\mathcal{M}s_2 \dots s_{2^{m+1}} = \rho_{m+1}$ is $(m + 1)$ -th harmonic of \mathcal{M} .

In **Figure 6B** we observe the D_0 -prolongations $\gamma_{\mathcal{M}\mathcal{L}\mathcal{M}^2(\mathcal{M}\mathcal{L})^2}^{r_{15}}$ and $\gamma_{\mathcal{M}\mathcal{L}\mathcal{M}^2}^{r_7}$, while $\gamma_{\mathcal{M}\mathcal{L}}^{r_3}$ and $\gamma_{\mathcal{M}}^{r_1}$ are visible in **Figure 6A**.

To summarize, we state

Proposition 1 ([17]). *Consider a merging bifurcation boundary $\gamma_{\sigma}^{\ell_{2^n}} \subset D_1$ related to the transition $\mathcal{Q}_{2^n} \Rightarrow \mathcal{Q}_n$, where $|\sigma| = n$, $n \geq 1$, x_{σ} is the rightmost point of the cycle \mathcal{O}_{σ} . Then the D_0 -prolongation $\gamma_{\sigma}^{r_{2^n-1}}$ of $\gamma_{\sigma}^{\ell_{2^n}}$ is given by the homoclinic bifurcation condition $f_{\omega}(r) = x_{\sigma}$ where ω is obtained from the first harmonic of σ by dropping the first symbol \mathcal{M} .*

An expansion bifurcation boundary $\zeta_{\sigma}^{\ell_n} \subset D_1$, $n \geq 3$, and its D_0 -prolongation $\zeta_{\sigma}^{r_{n-1}}$ are related in a similar way.

Further, let us consider an arbitrary chaoticity region \mathcal{C}_{2^m} , $m \geq 2$, which extends from D_1 to D_0 . As it was shown above, when $p \in D_1$ the boundaries of intervals J_i , $i = \overline{0, 2^m - 1}$, constituting the chaotic attractor \mathcal{Q}_{2^m} are defined by the images of ℓ only (see **Figure 7A**). When $p \in D_0$ but still close to b_1 , the rightmost interval J_{2^m-1} is given by $[r_{2^m-1}, \ell]$. The intervals J_i , $i = \overline{0, 2^m - 2}$, are confined by the images of r , in particular, $J_0 = [r, r_{2^m}]$. Moreover, $\ell_1 < r_{2^m}$, and therefore $\ell_1 \in J_0$ (see **Figure 7B**). Then, if p moves away from b_1 , the values ℓ and ℓ_1 increase, and eventually $\ell_1 = r_{2^m}$. Consequently, when p moves further away from b_1 the leftmost interval of \mathcal{Q}_{2^m} becomes $J_0 = [r, \ell_1]$.

Hence, additionally to $\gamma_{\rho_m}^{r_{2^m+1-1}}$ (which is related to r) another homoclinic bifurcation of the 2^m -cycle \mathcal{O}_{ρ_m} related to ℓ may occur. This homoclinic bifurcation defines an additional boundary of \mathcal{C}_{2^m} and occurs when the leftmost point of \mathcal{O}_{ρ_m} coincides with ℓ_1 , that is, at the boundary

$$\gamma_{\rho'_m}^{\ell_1} = \{p \in D_0 : f_{\mathcal{R}}(\ell) = x_{\rho'_m}\}, \quad (52)$$

where $\rho'_m = s_2 \dots s_{2^m}\mathcal{M}$ with $\mathcal{M}s_2 \dots s_{2^m} = \rho_m$ (thus, $x_{\rho'_m}$ is the leftmost point of \mathcal{O}_{ρ_m}). In **Figure 6B** the boundaries $\gamma_{\mathcal{L}\mathcal{M}^3}^{\ell_1}$ and $\gamma_{\mathcal{L}\mathcal{M}^3(\mathcal{L}\mathcal{M})^2}^{\ell_1}$ are shown.

We remark that the homoclinic bifurcation of type (52) can happen only if the corresponding cycle has at least one point in $I_{\mathcal{L}}$. This is not true for the fixed point $\mathcal{O}_{\mathcal{M}}$, and thus, the region \mathcal{C}_2 does not have this additional bifurcation boundary.

In such a way we can formulate

Proposition 2 ([17]). *Suppose that for $p \in D_0$ (in a certain neighborhood of b_1) there exists a D_0 -prolongation $\gamma_{\rho_m}^{r_{2^m+1-1}}$ of $\gamma_{\rho_m}^{\ell_{2^m+1}} \subset D_1$, $m \geq 1$, associated with the transition $\mathcal{Q}_{2^{m+1}} \Rightarrow \mathcal{Q}_{2^m}$. Then there exists a neighborhood $U(\gamma_{\rho_m}^{r_{2^m+1-1}})$ such that for $p \in U(\gamma_{\rho_m}^{r_{2^m+1-1}})$ there exists a boundary $\gamma_{\rho'_m}^{\ell_1}$, where ρ'_m is such that $x_{\rho'_m}$ is the leftmost point of the cycle \mathcal{O}_{ρ_m} . The boundary $\gamma_{\rho'_m}^{\ell_1}$ is associated with the same transition $\mathcal{Q}_{2^{m+1}} \Rightarrow \mathcal{Q}_{2^m}$ and related to a homoclinic bifurcation of the cycle \mathcal{O}_{ρ_m} colliding with an image of the critical point ℓ .*

As for chaoticity regions located in D_2 , to derive expressions for their D_0 -prolongations it is enough to swap \mathcal{L} 's and \mathcal{R} 's, as well as change r to ℓ in (49–51). Similarly, for $p \in D_0$ in a certain neighborhood of b_2 , a chaoticity region \mathcal{C}_{2^m} , $m \geq 1$, has an additional boundary related to a homoclinic bifurcation, whose expression can be derived from (52) by swapping \mathcal{L} 's and \mathcal{R} 's and replacing ℓ_1 by r_1 .

5.2. Periodicity Regions for $a_{\mathcal{R}} \rightarrow 0$: Stabilization of U-Sequence Cycles

In the previous section we described how the boundaries of chaoticity regions forming the STM structure change when the parameter point p enters D_0 by crossing b_1 or b_2 . Every merging or expansion bifurcation boundary existent in D_1/D_2 has its D_0 -prolongation. So does every region related to cyclic chaotic intervals. However, the dynamics of the map f for the parameter values $p \in D_0$ in the neighborhood of b_1/b_2 is not limited to D_0 -prolongations of the STM regions. For instance, in **Figure 8** between the right 2 · 2-fin $\mathcal{P}_{\mathcal{R}\mathcal{L}\mathcal{M}\mathcal{L}}$ of $\mathcal{P}_{\mathcal{R}\mathcal{L}}$ and the left 1 · 3-fin $\mathcal{P}_{\mathcal{M}\mathcal{R}\mathcal{L}}$ of $\mathcal{P}_{\mathcal{R}\mathcal{L}^2}$ one observes several tongues related to n -cycles with small values of n , e.g., with $n = 5, 6, 7, 10$. At the first site the origin of these tongues is unclear. In this section we explain the reason why such regions appear in the neighborhood of b_1 . All stated below can be generalized in the obvious way for b_2 .

To discover the origin of the tongues that are not D_0 -prolongations of the STM structure, recall that for the skew tent map (12) even if the dynamics converges to a chaotic attractor, there exist also many periodic solutions that are unstable. These cycles appear due to BCBs or DFBs and according to a certain universal order with decreasing β . Namely, the symbolic sequences of these cycles form the so-called *U-sequence* where “U” stands for “universal,” since this order is common for a wide class of unimodal maps [19, 20]. Notice that usually U-sequence is described in terms of kneading sequences associated with superstable cycles (in the symbolic representation of which the symbol C related to the extremum of the function is omitted). The skew tent map obviously cannot have superstable cycles, and

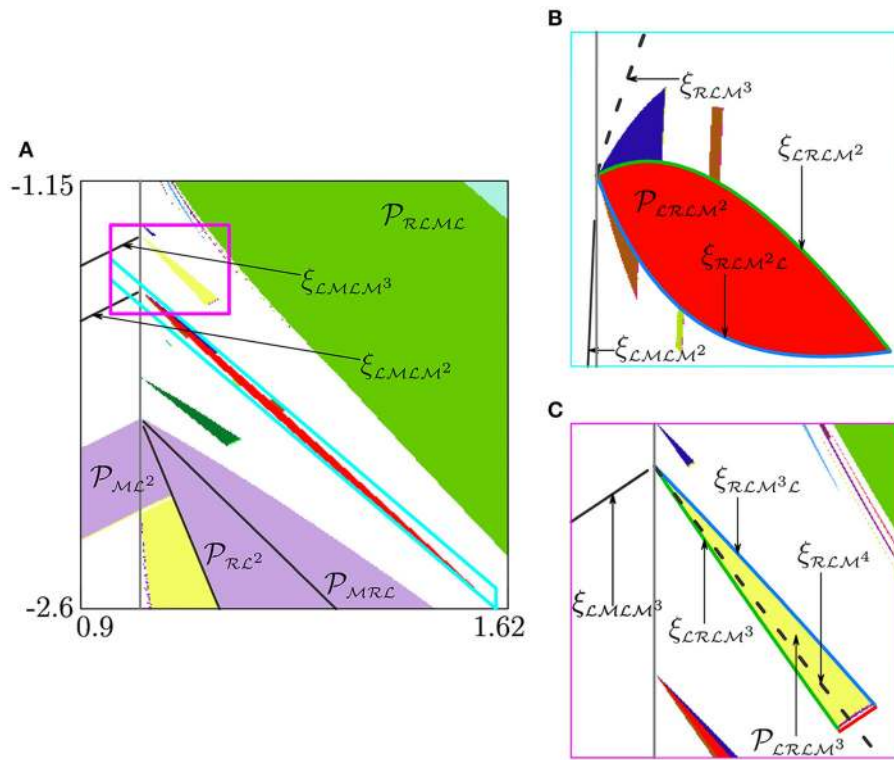


FIGURE 8 | Stabilization of U-sequence cycles. Panels (B,C) are the enlargements of the areas marked in (A) by blue and magenta lines, respectively. The parameters are $a_L = 0.57$, $a_R = 0.2$, $d_L = 0$, $d_R = 1$.

by using the term “U-sequence” we mean the order of appearance of the cycles in terms of their complete symbolic representation.

Another peculiarity of the skew tent map is that some blocks of the U-sequence appear simultaneously. For instance, crossing the DFB boundary η_{BA} (or η_B when $\alpha > 1$) with decreasing β leads to immediate appearance of cycles with symbolic sequences being harmonics of B , defined in (18). Similar statement is true for a basic cycle $\mathcal{O}_{BA^{n-1}}$, $n \geq 3$, that is, all its harmonics appear at the same parameter values (at the boundary $\eta_{BA^{n-1}}$).

In general, the U-sequence is constructed according to a certain iterative procedure and includes also the basic cycles. Detailed description of this construction procedure, as well as explanation of its peculiarities related to the skew tent map, is beyond the scopes of the current paper. Here we only explain how an unstable cycle whose symbolic sequence belongs to U-sequence may be “stabilized” when p enters D_0 crossing b_1 or b_2 . For shortness we refer below to such cycles as *U-sequence cycles*.

Consider, for example, the region related to the 5-cycle observed in **Figure 8A** (see also the enlargement in **Figure 8B**). For $p \in D_1$ the line $\xi_{LM^2M^2}$ denotes the fold BCB, at which two unstable 5-cycles are born: the cycle $\mathcal{O}_{LM^2M^2}$ and its complementary cycle $\mathcal{O}_{M^2LM^2}$. The condition related to this bifurcation is

$$f^5(d_L) = f_{LM^2M^2}(d_L) \equiv f_{M^2LM^2}(d_L) = d_L. \quad (53)$$

In fact, (53) defines two different BCB conditions which are equivalent due to the continuity of f , that is, $f_L(d_L) = f_M(d_L)$. This explains why the two cycles (which coincide at the bifurcation moment) appear simultaneously and why their symbolic sequences after the bifurcation differ by a single symbol.

For $p \in \xi_{LM^2M^2} \cap b_1$ the condition (53) can be rewritten taking into account that $f(d_L) = d_R$:

$$f^5(d_L) = f^4(d_R) = f_{MLM^2}(d_R) \equiv f_{RLM^2}(d_R) = d_L, \quad (54)$$

since $f_M(d_R) = f_R(d_R)$. Applying f one more time to (54), and using again the continuity property of f , we get

$$\begin{aligned} f^5(d_R) &= f_{MLM^2L}(d_R) \equiv f_{RLM^2L}(d_R) \equiv f_{MLM^2M}(d_R) \\ &\equiv f_{RLM^2M}(d_R) = d_R (= f(d_L)). \end{aligned} \quad (55)$$

Similarly to (53), the equalities (55) define four BCB conditions that, nevertheless, are equivalent for $p \in \xi_{LM^2M^2} \cap b_1$. Notice that in (55) the BCB conditions depend on d_R , in contrast to (53) where they depend on d_L . This is similar to how conditions of homoclinic bifurcations change when the parameter point p moves from D_1 to D_0 (as described in Section 5).

Now, for $p \in D_0$ not all the BCB conditions (55) remain equivalent. Moreover, the two cycles $\mathcal{O}_{LM^2M^2}$ and $\mathcal{O}_{M^2LM^2}$ do not appear simultaneously for $p \in D_0$. Instead, each of them creates a complementary pair with one of the two new cycles having a single point in I_R (and, thus, existing only for

$p \in D_0$). Namely, the cycle \mathcal{O}_{LMCM^2} is now complementary to \mathcal{O}_{LRCM^2} , while $\mathcal{O}_{M^2LM^2}$ is complementary to \mathcal{O}_{MRCM^2} . Clearly, the BCB conditions for the two cycles making a pair are equivalent, that is,

$$\begin{aligned} f_{MLCM^2L}(d_R) &\equiv f_{RLCM^2L}(d_R) = d_R, \\ f_{MLCM^2M}(d_R) &\equiv f_{RLCM^2M}(d_R) = d_R. \end{aligned} \quad (56)$$

However, the related bifurcation is not necessarily fold BCB.

Let us describe in more detail the BCBs related to four mentioned 5-cycles. With decreasing μ_R (for the other parameter values as in **Figure 8**) they appear in the following order. First, the two unstable cycles $\mathcal{O}_{M^2LM^2}$ and \mathcal{O}_{MRCM^2} are born due to the fold BCB ξ_{RLCM^3} (dashed curve). Then there occurs the BCB ξ_{LRCM^2} (green curve), at which the point x_{MRCM^2} of \mathcal{O}_{MRCM^2} moves from the middle partition I_M to the left one I_L . Hence, the stable cycle \mathcal{O}_{LRCM^2} appears. Finally, the point of \mathcal{O}_{LRCM^2} located in the right partition I_R moves to the middle partition I_M due to the other BCB ξ_{RLCM^2L} creating the unstable cycle $\mathcal{O}_{M^2LM^2}$. The region \mathcal{P}_{LRCM^2} is then confined by ξ_{RLCM^2L} and ξ_{LRCM^2} . For μ_R being below ξ_{RLCM^2L} the two unstable cycles \mathcal{O}_{LMCM^2} and $\mathcal{O}_{M^2LM^2}$ continue to exist.

We also remark that the region \mathcal{P}_{LRCM^2} has fins similarly to PA-regions. The mechanism of their creation is the same. Indeed, at the BCB boundary of \mathcal{P}_{LRCM^2} related to collision of the cycle with d_L/d_R , the map f^5 is locally (in the neighborhood of d_L/d_R) topologically equivalent to the skew tent map $g_{\alpha,\beta,1}$ with $\alpha = a_L^2 a_M^2 a_R$ and $\beta = a_L a_M^3 a_R / \beta = a_L a_M^3$. Crossing transversally this BCB boundary may lead to appearance of a stable $5k$ -cycle, $k \geq 1$ (cf. Section 4.2). Existence of fins and their number depends on the parameter values.

The scenario due to that the region \mathcal{P}_{LRCM^2} occurs corresponds to what we call the “stabilization” of a U-sequence cycle. That is, for $p \in D_1$ the U-sequence cycle \mathcal{O}_σ is unstable (except for the basic cycles $\mathcal{O}_{MLC^{n-1}}$, $n \geq 2$, which can be stable). When p crosses b_1 entering D_0 a new cycle $\mathcal{O}_{\sigma'}$ (among the others) is born. The symbolic sequence σ' differs from σ by a single symbol. Namely, the symbol M corresponding to the rightmost point of \mathcal{O}_σ is replaced by R in σ' . If the slope a_R is small enough, the cycle $\mathcal{O}_{\sigma'}$ is stable.

The region related to a 6-cycle observed as well in **Figure 8A** has the same origin, though it looks somewhat unlike \mathcal{P}_{LRCM^2} . Let us consider it in more detail and explain why it has different form. For $p \in D_1$ due to the BCB ξ_{LMCM^3} the two unstable 6-cycles \mathcal{O}_{LMCM^3} and $\mathcal{O}_{M^2LM^3}$ appear. Similarly to the case considered above (of the 5-cycle), for $p \in D_0$ there are two more 6-cycles: \mathcal{O}_{LRCM^3} (stable) and \mathcal{O}_{MRCM^3} (unstable). However, in contrast to the previous case, the sequence of bifurcations occurring with decreasing μ_R is different. First, the fold BCB ξ_{RLCM^3L} occurs giving rise to the stable cycle \mathcal{O}_{LRCM^3} and the unstable one \mathcal{O}_{LMCM^3} . Next, the two unstable cycles \mathcal{O}_{MRCM^3} and $\mathcal{O}_{M^2LM^3}$ appear due to the fold BCB ξ_{RLCM^4} . Finally, \mathcal{O}_{LRCM^3} (stable) and \mathcal{O}_{MRCM^3} (unstable) disappear at ξ_{LRCM^3} . Notice that the region \mathcal{P}_{LRCM^3} (differently from \mathcal{P}_{LRCM^2}) has one more boundary η_{LRCM^3} that corresponds to the DFB of the related cycle.

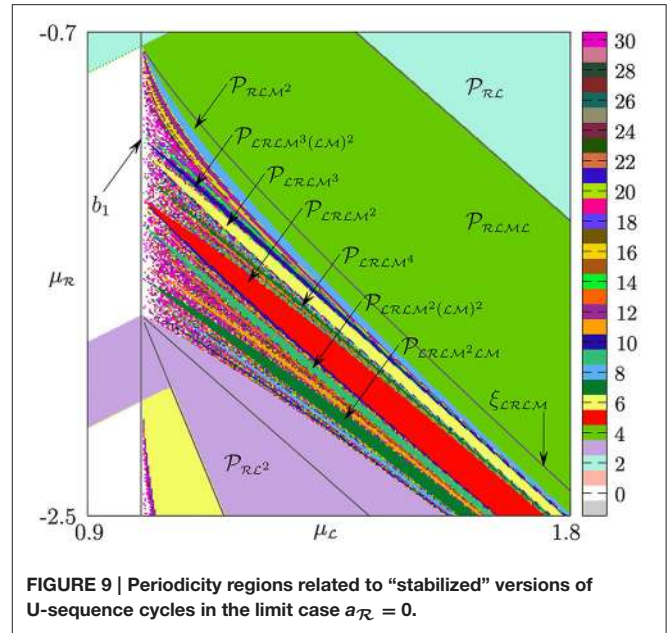


FIGURE 9 | Periodicity regions related to “stabilized” versions of U-sequence cycles in the limit case $a_R = 0$.

The other regions emerging from b_1 and not belonging to PA or fin structures are of the same nature as the two regions described above. Notice that, for a U-sequence cycle $\mathcal{O}_{LM\sigma}$, the higher is its period, the more of its points are located in I_M . Thus, to observe for $p \in D_0$ the region $\mathcal{P}_{LR\sigma}$ corresponding to the “stabilized” version of $\mathcal{O}_{LM\sigma}$, the value a_R should be small enough. In the limit case $a_R = 0$, for every U-sequence cycle there exists the related periodicity region in D_0 (see **Figure 9**), since the multiplier of $\mathcal{O}_{LR\sigma}$ equals zero. Clearly, all these regions have only BCB boundaries and extend to infinite values of μ_L, μ_R .

6. CHAOTICITY REGIONS SURROUNDING THE PERIOD ADDING AND THE FIN STRUCTURES

In Section 4.1 we have recalled how to obtain boundaries of PA-regions. There are chaoticity regions adjacent to them, and now we discuss how their boundaries can be derived using the skew tent map. As an example, **Figure 10** shows the structure of chaoticity regions close to the BCB boundary ξ_{LR} of the PA-region \mathcal{P}_{LR} . One can observe here two hatched domains. The first one, denoted $D_{st}^{L,LR}$, is contiguous to the PA-region \mathcal{P}_{LR} along the boundary ξ_{LR} . For $p \in D_{st}^{L,LR}$ the map f^2 is locally topologically conjugate to the skew tent map. Hence, the bifurcation structure can be described by applying directly the results known for the skew tent map.

In general, in the neighborhood of the boundary $\xi_\sigma^{d_L}/\xi_\sigma^{d_R}$ of every PA-region \mathcal{P}_σ associated with an n -cycle \mathcal{O}_σ , there exist domains denoted $D_{st}^{L,\sigma}/D_{st}^{R,\sigma}$, respectively, such that for related parameter values the iterate f^n is locally (in the neighborhood of d_L/d_R) topologically conjugate to the skew tent map.

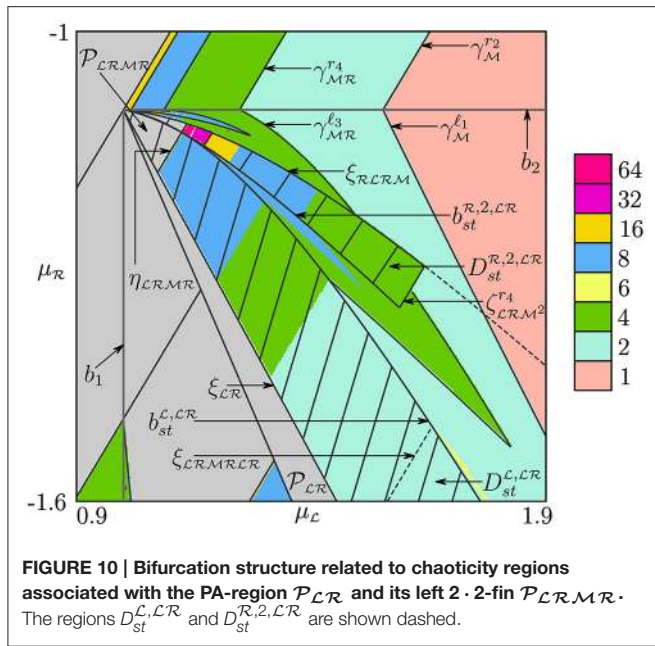


FIGURE 10 | Bifurcation structure related to chaoticity regions associated with the PA-region \mathcal{P}_{LR} and its left 2 · 2-fin \mathcal{P}_{LRMR} . The regions $D_{st}^{L,LR}$ and $D_{st}^{R,2,LR}$ are shown dashed.

The second hatched domain in **Figure 10**, denoted $D_{st}^{R,2,LR}$, is located in the neighborhood of the boundary ξ_{RLRM} of the left 2 · 2-fin region \mathcal{P}_{LRMR} . In general, let us consider an arbitrary PA-region \mathcal{P}_σ , related to the n -cycle \mathcal{O}_σ , and its left/right $k \cdot n$ -fin \mathcal{P}_τ . In the neighborhood of the BCB boundary $\xi_\tau^{d_R}/\xi_\tau^{d_L}$, respectively, there exists a domain $D_{st}^{R,k,\sigma}/D_{st}^{L,k,\sigma}$ such that the restriction of f^{kn} to its absorbing interval is topologically conjugate to the skew tent map.

6.1. Homoclinic Bifurcation Boundaries in $D_{st}^{L,\sigma}$

We consider first the trunk region \mathcal{P}_{LR} and sub-domain $D_{st}^{L,LR}$ contiguous to it along the boundary ξ_{LR} . Clearly, both points of the 2-cycle $\mathcal{O}_{LR} = \{x_{LR}, x_{RL}\}$ are fixed points for the second iterate f^2 , and at $p \in \xi_{LR}$ the point x_{LR} collides with d_L . The resulting dynamics caused by this collision depends on the slopes of f^2 at both sides of d_L at the bifurcation moment. In the neighborhood of the border point the map f^2 is defined as (see **Figure 11**)

$$f^2(x) \equiv \begin{cases} f_{LR}, & f_L^{-1}(d_R) < x \leq d_L, \\ f_{MR}, & d_L < x \leq f_M^{-1}(d_R), \end{cases}$$

where $f_L^{-1}(d_R) = (d_R - \mu_L)/a_L, f_M^{-1}(d_R) = (d_R - \mu_M)/a_M$. If the condition

$$f_{LR}(d_L) < f_M^{-1}(d_R) \tag{57}$$

is satisfied, then f^2 has a local absorbing interval $J = [f_M^{-1}(d_L), f_{MR}(d_L)]$, on which f^2 is topologically conjugate to the skew tent map (12) with $\alpha = a_L a_R$ and $\beta = a_M a_R$. The

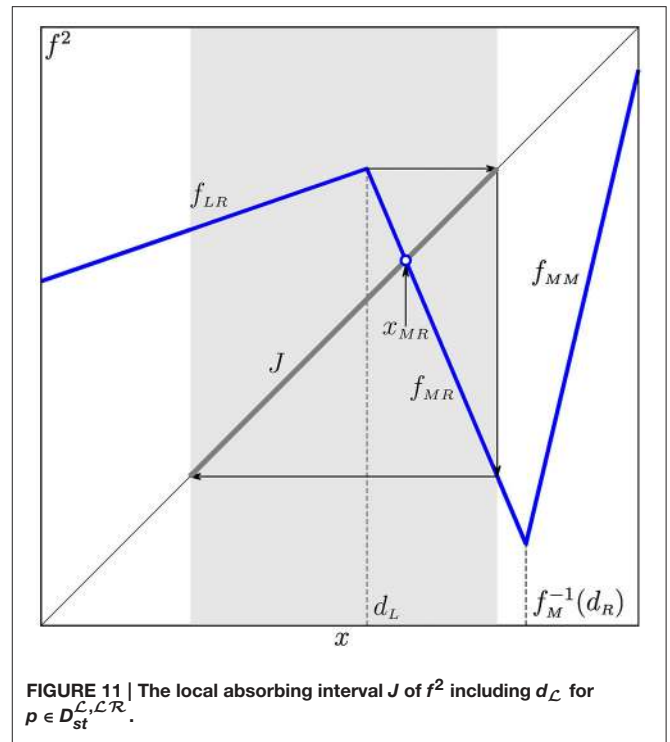


FIGURE 11 | The local absorbing interval J of f^2 including d_L for $p \in D_{st}^{L,LR}$.

parameter region for that the conjugacy takes place is given as

$$D_{st}^{L,LR} = \left\{ p \in D_0 : x_{LR} > d_L, f_{MR}(l) < d_R \right\} = \left\{ p \in D_0 : \mu_L > -\frac{\mu_R}{a_R} + \frac{1 - a_R a_L}{a_R} d_L, a_R a_M l + a_R \mu_M + \mu_R < d_R \right\}. \tag{58}$$

It is located between two boundaries, namely, ξ_{LR} and

$$b_{st}^{L,LR} = \{p \in D_0 : f_{MR}(l) = d_R\}, \tag{59}$$

which follows from the equality $f_{LR}(d_L) = f_M^{-1}(d_R)$.

For $p \in D_{st}^{L,LR}$ expressions for all bifurcation boundaries are derived directly from the related expressions known for the skew tent map (cf. Section 3).

Similarly, for an arbitrary PA-region \mathcal{P}_σ related to an n -cycle $\mathcal{O}_\sigma, n = |\sigma|$, let $\sigma = \mathcal{L}\sigma_1\mathcal{R}\sigma_2$, where the symbolic sequences σ_1 and σ_2 are such that the point $x_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}/x_{\mathcal{R}\sigma_2\mathcal{L}\sigma_1}$ collides with d_L/d_R at the BCB boundary $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}/\xi_{\mathcal{R}\sigma_2\mathcal{L}\sigma_1}$, respectively. For $p \in \mathcal{P}_\sigma$ consider the map f^n , for which every point of \mathcal{O}_σ is a fixed point. Close to the boundary $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$, two branches of f^n joining at the border point d_L are

$$f^n(x) \equiv \begin{cases} f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}, & \check{d} < x \leq d_L, \\ f_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}, & d_L < x \leq \bar{d}. \end{cases}$$

Here \check{d} and \bar{d} are preimages of d_R , such that $f_{\mathcal{L}\sigma_1}(\check{d}) = d_R$ and $f_{\mathcal{M}\sigma_1}(\bar{d}) = d_R$. As before it can be shown that as long as

$$f^n(d_L) \equiv f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}(d_L) < \bar{d} \tag{60}$$

there exists a local absorbing interval

$$J = [f_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}^2(d_{\mathcal{L}}), f_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}(d_{\mathcal{L}})] \quad (61)$$

such that the restriction $f^n|_J$ is topologically conjugate to the skew tent map g with the slopes $\alpha = a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $\beta = a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$. Here $a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$ are the slopes of the composite functions $f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $f_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$, respectively. Such a conjugacy persists until the parameter point p crosses the border $b_{st}^{\mathcal{L},\sigma}$ given by the condition $f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}(d_{\mathcal{L}}) = d$, which is equivalent to $f_{\mathcal{M}\sigma_1}(f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}(d_{\mathcal{L}})) = f_{\mathcal{M}\sigma_1}(d) \equiv d_{\mathcal{R}}$, i.e.,

$$b_{st}^{\mathcal{L},\sigma} = \{p \in D_0 : f_{\sigma_1\mathcal{R}\sigma_2\mathcal{M}\sigma_1}(\ell) = d_{\mathcal{R}}\}. \quad (62)$$

Hence, we can state the following

Proposition 3 ([17]). For $p \in D_{st}^{\mathcal{L},\sigma}$, where the region

$$D_{st}^{\mathcal{L},\sigma} = \{p \in D_0 : x_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2} > d_{\mathcal{L}}, f_{\sigma_1\mathcal{R}\sigma_2\mathcal{M}\sigma_1}(\ell) < d_{\mathcal{R}}\} \quad (63)$$

is confined by $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $b_{st}^{\mathcal{L},\sigma}$ given in (62), the expressions for boundaries of periodicity regions (fins) and chaoticity regions can be obtained from the related expressions known for the skew tent map (62) substituting α by $a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and β by $a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$.

For instance, we consider the left $k \cdot n$ -fin \mathcal{P}_{τ} , $\tau = \mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}$, $k \geq 3$, of the PA-region $\mathcal{P}_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$. This fin region represents the area located between three BCB boundaries: $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$, $\xi_{\tau_i}^{d_{\mathcal{L}}}$ (the point x_{τ_i} of the cycle \mathcal{O}_{τ} collides with $d_{\mathcal{L}}$), $\xi_{\tau_j}^{d_{\mathcal{R}}}$ (the point x_{τ_j} of the cycle \mathcal{O}_{τ} collides with $d_{\mathcal{R}}$), and the DFB boundary η_{τ} . We remark that the boundaries $\xi_{\tau_i}^{d_{\mathcal{L}}}$ and η_{τ} emerge from $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$, as well as the boundaries between chaoticity regions adjacent to $\mathcal{P}_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$.

Analytic expressions for all these boundaries are derived by using the skew tent map. Boundaries of the fin region, $\xi_{\tau_i}^{d_{\mathcal{L}}}$ and η_{τ} , are given in (41) and (40), respectively. Let us consider the merging and the expansion bifurcation boundaries confining chaoticity regions \mathcal{Q}_{2kn} and \mathcal{Q}_{kn} . They can be obtained by applying map replacement technique to expressions for $\gamma_{\mathcal{B}\mathcal{A}^{k-1}}^{c_{2k}}$ (15) and $\zeta_{\mathcal{B}\mathcal{A}^{k-2}\mathcal{B}}^{c_k}$ (17).

Namely, in (15) we substitute all symbols \mathcal{A} with symbolic sequence $\mathcal{L}\sigma_1\mathcal{R}\sigma_2$, all symbols \mathcal{B} with symbolic sequence $\mathcal{M}\sigma_1\mathcal{R}\sigma_2$, critical point c with $f_{\sigma_1\mathcal{R}\sigma_2}(\ell)$, slope α with $a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$, and β with $a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$. Recall that $a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$ are the slopes of the composite functions $f_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $f_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$, respectively. Expressing $a_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ and $a_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2}$ in terms of the original parameters $a_{\mathcal{L}}$, $a_{\mathcal{M}}$, $a_{\mathcal{R}}$ we obtain the merging bifurcation boundary $\gamma_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}}^{\ell_{2kn+n-1}}$ related to transition $\mathcal{Q}_{2kn} \Rightarrow \mathcal{Q}_{kn}$. Likewise, by applying the same replacement to (17) one gets the expansion bifurcation boundary $\zeta_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-2}\mathcal{M}\sigma_1\mathcal{R}\sigma_2}^{\ell_{kn+n-1}}$. In short this procedure can be

represented by the following scheme:

$$\gamma_{\mathcal{B}\mathcal{A}^{k-1}}^{c_{2k}} \xrightarrow[\mathcal{A} \rightarrow \mathcal{L}\sigma_1\mathcal{R}\sigma_2 \quad \mathcal{B} \rightarrow \mathcal{M}\sigma_1\mathcal{R}\sigma_2]{c \rightarrow f_{\sigma_1\mathcal{R}\sigma_2}(\ell)} \gamma_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}}^{\ell_{2kn+n-1}} \quad (64a)$$

$$\zeta_{\mathcal{B}\mathcal{A}^{k-2}\mathcal{B}}^{c_k} \xrightarrow[\mathcal{A} \rightarrow \mathcal{L}\sigma_1\mathcal{R}\sigma_2 \quad \mathcal{B} \rightarrow \mathcal{M}\sigma_1\mathcal{R}\sigma_2]{c \rightarrow f_{\sigma_1\mathcal{R}\sigma_2}(\ell)} \zeta_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-2}\mathcal{M}\sigma_1\mathcal{R}\sigma_2}^{\ell_{kn+n-1}} \quad (64b)$$

For sake of shortness schemes of such kind are used below to describe similar replacement procedures.

To sum up we can state

Proposition 4 ([17]). Let the parameter point p move along the BCB boundary $\xi_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$ starting from nk -fin region $\mathcal{P}_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}}$, $k \geq 3$. Then p first crosses the DFB boundary $\eta_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}}$ (40) leading to $2kn$ -cyclic chaotic intervals \mathcal{Q}_{2kn} , then the merging bifurcation boundary $\gamma_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}}^{\ell_{2kn+n-1}}$ (64a) implying the transition $\mathcal{Q}_{2kn} \Rightarrow \mathcal{Q}_{kn}$, and finally, the expansion bifurcation boundary $\zeta_{\mathcal{M}\sigma_1\mathcal{R}\sigma_2(\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-2}\mathcal{M}\sigma_1\mathcal{R}\sigma_2}^{\ell_{kn+n-1}}$ (64b) leading to $\mathcal{Q}_{kn} \Rightarrow \mathcal{Q}_n$.

The case $k = 2$ is particular. Since a $2n$ -fin cycle corresponds to the 2-cycle of the skew tent map g , its DFB leads to the appearance of $2^m \cdot n$ -cyclic chaotic intervals, where $m \geq 2$ depends on $a_{\mathcal{L}}$, $a_{\mathcal{R}}$. The cascade of the subsequent merging bifurcations is obtained from (20) applying the same replacement as in (64). In Figure 12 an example of this cascade observed inside the domain $D_{st}^{\mathcal{L},\mathcal{L}\mathcal{R}}$ is shown, located close to the left $2 \cdot 2$ -fin region $\mathcal{P}_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}$ of the trunk region $\mathcal{P}_{\mathcal{L}\mathcal{R}}$. One can see three merging bifurcation boundaries, namely, $\gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}(\mathcal{M}\mathcal{R})^2}^{\ell_{17}}$, $\gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}}^{\ell_9}$, and $\gamma_{\mathcal{M}\mathcal{R}}^{\ell_5}$ leading to the transitions $\mathcal{Q}_{2.8} \Rightarrow \mathcal{Q}_{2.4} \Rightarrow \mathcal{Q}_{2.2} \Rightarrow \mathcal{Q}_2$, respectively. To get the expression for $\gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}(\mathcal{M}\mathcal{R})^2}^{\ell_{17}}$ the following replacement

$$\gamma_{\mathcal{B}\mathcal{A}\mathcal{B}^2}^{c_8} \xrightarrow[\mathcal{A} \rightarrow \mathcal{L}\mathcal{R} \quad \mathcal{B} \rightarrow \mathcal{M}\mathcal{R}]{c \rightarrow f_{\mathcal{R}}(\ell)} \gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}(\mathcal{M}\mathcal{R})^2}^{\ell_{17}}$$

can be used. In this way we obtain

$$\gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}(\mathcal{M}\mathcal{R})^2}^{\ell_{17}} = \left\{ p \in D_0 : a_{\mathcal{L}\mathcal{R}}^2 a_{\mathcal{M}\mathcal{R}}^6 + \frac{a_{\mathcal{L}\mathcal{R}}}{a_{\mathcal{M}\mathcal{R}}} - 1 = 0 \right\} \\ = \left\{ p \in D_0 : a_{\mathcal{L}}^2 a_{\mathcal{M}}^6 a_{\mathcal{R}}^8 + \frac{a_{\mathcal{M}}}{a_{\mathcal{L}}} - 1 = 0 \right\}. \quad (65)$$

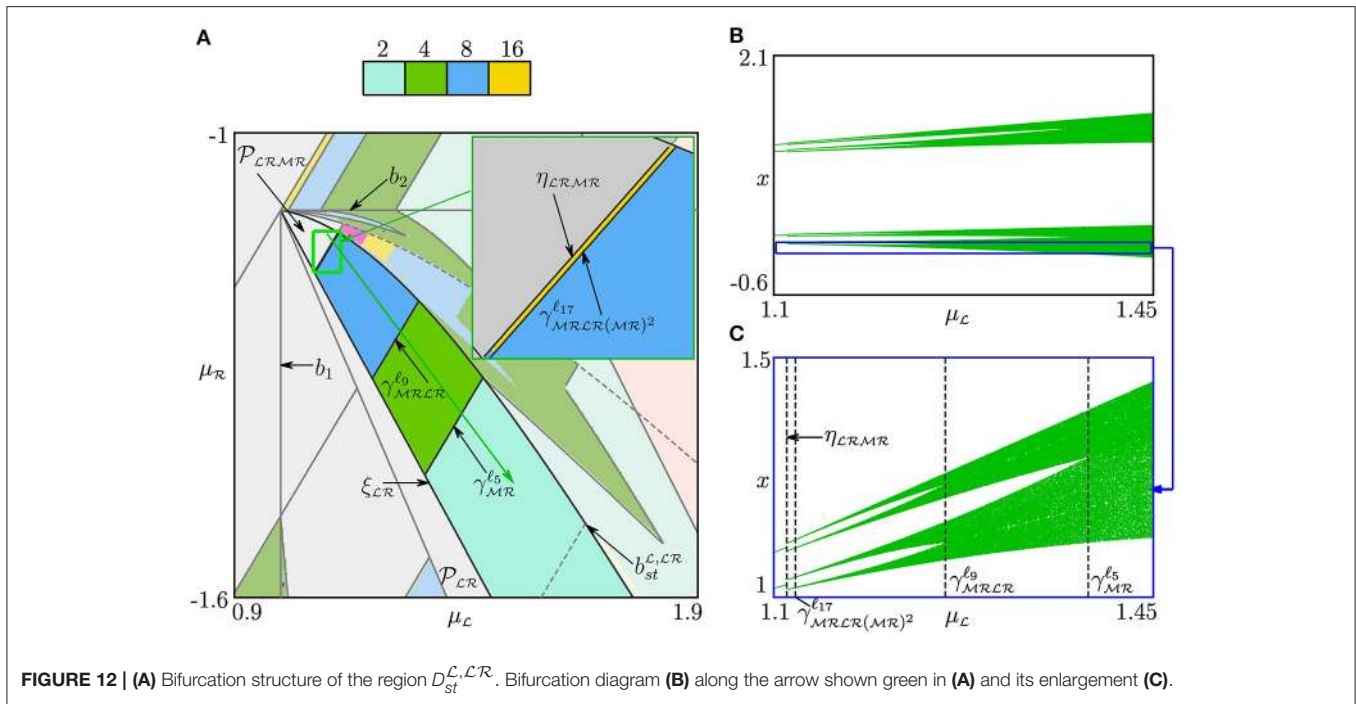
The expressions for the other two boundaries

$$\gamma_{\mathcal{M}\mathcal{R}\mathcal{L}\mathcal{R}}^{\ell_9} = \left\{ p \in D_0 : a_{\mathcal{L}}^2 a_{\mathcal{M}}^2 a_{\mathcal{R}}^4 + \frac{a_{\mathcal{L}}}{a_{\mathcal{M}}} - 1 = 0 \right\}, \quad (66a)$$

$$\gamma_{\mathcal{M}\mathcal{R}}^{\ell_5} = \left\{ p \in D_0 : a_{\mathcal{M}}^2 a_{\mathcal{R}}^2 + \frac{a_{\mathcal{M}}}{a_{\mathcal{L}}} - 1 = 0 \right\} \quad (66b)$$

can be obtained by the same replacement.

We have described above bifurcations of cyclic chaotic intervals and the related boundaries in the neighborhood of $\xi_{\sigma}^{d_{\mathcal{L}}}$ of a PA-region \mathcal{P}_{σ} . To obtain similar bifurcation boundaries in the neighborhood of $\xi_{\sigma}^{d_{\mathcal{R}}}$ one can swap symbols \mathcal{L} and \mathcal{R} as well as interchange the critical points ℓ and r in all analytic expressions.



6.2. Homoclinic Bifurcation Boundaries in $D_{st}^{\mathcal{R},k,\sigma}$

In this section we describe another parameter domain $D_{st}^{\mathcal{R},k,\sigma}$, mentioned above, which is located in the neighborhood of an arbitrary left fin region of a PA-region \mathcal{P}_σ . To describe bifurcation structure of $D_{st}^{\mathcal{R},k,\sigma}$ we again use conjugacy with skew tent map and map replacement technique.

As the first example let us consider the left $2 \cdot 2$ -fin region \mathcal{P}_{LRMR} of the trunk region \mathcal{P}_{LR} (see **Figure 12A**). As explained in Section 4.2 the fin region \mathcal{P}_{LRMR} has the BCB boundary ξ_{RLRM} related to the collision of the cycle \mathcal{O}_{LRMR} with the border $d_{\mathcal{R}}$.

In **Figure 13** we show the plot of f^4 on the interval including the border point $d_{\mathcal{R}}$ for parameter values located in the neighborhood of ξ_{RLRM} . The two linear pieces of f^4 adjoining at $d_{\mathcal{R}}$ are defined as

$$f^4(x) \equiv \begin{cases} f_{MLCRM}, & f_M^{-1}(d_L) \geq x \leq d_{\mathcal{R}}, \\ f_{RLCRM}, & d_{\mathcal{R}} < x \leq f_R^{-1}(d_L), \end{cases}$$

where $f_M^{-1}(d_L) = (d_L - \mu_M)/a_M$, $f_R^{-1}(d_L) = (d_L - \mu_R)/a_R$. If the conditions

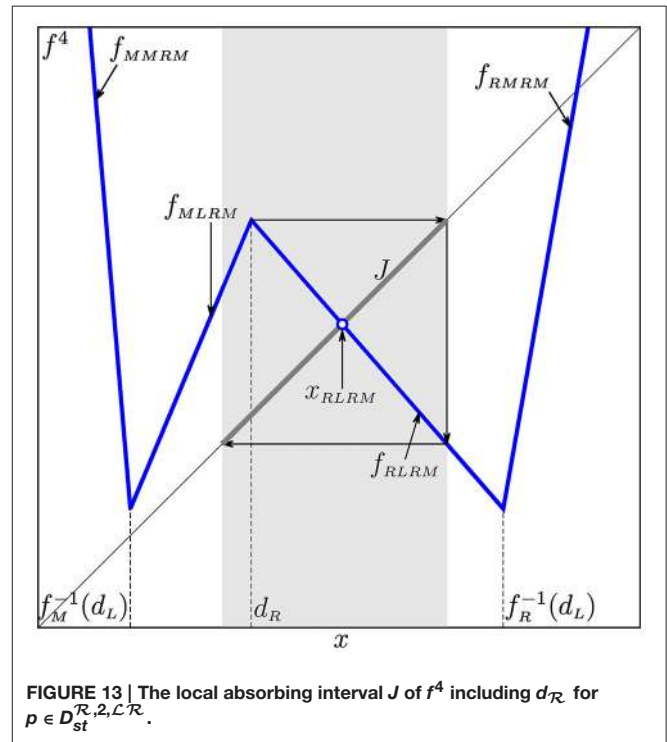
$$f_{RLCRM}(d_{\mathcal{R}}) < f_R^{-1}(d_L), \tag{67a}$$

$$f_{RLCRM}^2(d_{\mathcal{R}}) > x_{MLCRM} \tag{67b}$$

are held, then there exists a local absorbing interval

$$J = [f_{RLCRM}^2(d_{\mathcal{R}}), f_{RLCRM}(d_{\mathcal{R}})] \tag{68}$$

for f^4 . The restriction $f^4|_J$ is topologically conjugate to the skew tent map g with the slopes $\alpha = a_L a_M^2 a_R$ and $\beta = a_L a_M a_R^2$.



Taking into account that $f_{\mathcal{R}}(f_{RLCRM}(d_{\mathcal{R}})) \equiv f_{LRMR}(r)$, $f_{RLCRM}^2(d_{\mathcal{R}}) = f_{LRM}(f_{RLCRM}(r))$, and $x_{MLCRM} = f_{LRM}(x_{LRM}^2)$, the conditions (67) can be written as

$$f_{LRMR}(r) < d_L, \tag{69a}$$

$$f_{LRMR}(r) > x_{LRM}^2. \tag{69b}$$

The conjugacy between $f^4|_J$ and g holds in the parameter region $D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}}$ confined by the BCB boundary $\xi_{\mathcal{R}\mathcal{L}\mathcal{R}\mathcal{M}}$, the boundary

$$b_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}} = \{p \in D_0 : f_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}(r) = d_{\mathcal{L}}\}, \quad (70)$$

and the boundary

$$\zeta_{\mathcal{L}\mathcal{R}\mathcal{M}^2}^{r_4} = \{p \in D_0 : f_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}(r) = x_{\mathcal{L}\mathcal{R}\mathcal{M}^2}\}. \quad (71)$$

Note that for the map f the boundary (71) corresponds to an expansion bifurcation leading from Q_4 to Q . Thus, the region $D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}}$ is defined as

$$D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}} = \{p \in D_0 : x_{\mathcal{R}\mathcal{L}\mathcal{R}\mathcal{M}} > d_{\mathcal{R}}, f_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}(r) < d_{\mathcal{L}}, f_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}(r) > x_{\mathcal{L}\mathcal{R}\mathcal{M}^2}\}. \quad (72)$$

Figure 14A presents the region $D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}}$, where one can observe chaoticity regions Q_{2^m} , $m = \overline{2,6}$. The bifurcation boundaries separating these regions are related to transitions $Q_{2^m} \Rightarrow Q_{2^{m-1}}$. To get expressions for these boundaries we use map replacement technique. Namely, we take the corresponding formulae (20) known for the skew tent map and apply the following replacement:

$$\gamma_{\rho_j}^{c_{j+1}} \xrightarrow[A \rightarrow \mathcal{M}\mathcal{L}\mathcal{R}\mathcal{M} \quad B \rightarrow \mathcal{R}\mathcal{L}\mathcal{R}\mathcal{M}]{\mathcal{C} \rightarrow f_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}(r)} \gamma_{\tau_j}^{r_{2^{j+3}+3}}. \quad (73)$$

Here the symbolic sequence τ_j , $j \geq 0$, is derived from the harmonic ρ_j (18) by substituting $\mathcal{M}\mathcal{L}\mathcal{R}\mathcal{M}$ instead of \mathcal{A} and $\mathcal{R}\mathcal{L}\mathcal{R}\mathcal{M}$ instead of \mathcal{B} .

Summing up we can describe $D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}}$ in the following way. The region $D_{st}^{\mathcal{R},2,\mathcal{L}\mathcal{R}} \setminus \mathcal{P}_{\mathcal{L}\mathcal{R}\mathcal{M}\mathcal{R}}$ consists of $t+1$ chaoticity regions $C_{4,2j}$, $j = \overline{0,t}$, where $t \geq 1$ depends on $a_{\mathcal{L}}$, $a_{\mathcal{R}}$. If the parameter point p crosses the boundary $\gamma_{\tau_j}^{r_{2^{j+3}+3}}$ between the chaoticity regions $C_{2^{j+1}}$ and C_{2^j} , given by

$$\gamma_{\tau_j}^{r_{2^{j+3}+3}} = \left\{ p \in D_{st}^{\mathcal{L}} : a_{\mathcal{L}}^{2(\delta_j+\delta_{j+1})} a_{\mathcal{M}}^{2(2\delta_j+\delta_{j+1})} a_{\mathcal{R}}^{2(\delta_j+2\delta_{j+1})} + \left(\frac{a_{\mathcal{M}}}{a_{\mathcal{R}}}\right)^{(-1)^{j+1}} - 1 = 0 \right\}, j = \overline{0, t-1},$$

then the merging bifurcation $C_{2^{j+1}} \Rightarrow C_{2^j}$ takes place. Here $\delta_j = (2^j - (-1)^j)/3$. In **Figure 14A** the merging bifurcation boundaries $\gamma_{\tau_3}^{r_{67}}$, $\gamma_{\tau_2}^{r_{35}}$, $\gamma_{\tau_1}^{r_{19}}$, and $\gamma_{\tau_0}^{r_{11}}$ are shown. See also 1D bifurcation diagrams in **Figures 14B,C**.

In a similar way we characterize the region $D_{st}^{\mathcal{R},k,\sigma}$, related to the boundary $\xi_{\tau}^{d_{\mathcal{R}}}$ of the left $k \cdot n$ -fin \mathcal{P}_{τ} , $\tau = (\mathcal{L}\sigma_1\mathcal{R}\sigma_2)^{k-1}\mathcal{M}\sigma_1\mathcal{R}\sigma_2$, of the trunk region $\mathcal{P}_{\mathcal{L}\sigma_1\mathcal{R}\sigma_2}$, $k \geq 2$, $n \geq 2$. In order to determine the branches of f^{kn} which are adjacent to one another at the border point, we rewrite the symbolic sequence associated with the fin cycle as follows:

$$\tau \equiv \underline{\mathcal{R}}\sigma_2\mathcal{L}\sigma_1(\mathcal{R}\sigma_2\mathcal{L}\sigma_1)^{k-3}\mathcal{R}\sigma_2\mathcal{L}\sigma_1\mathcal{R}\sigma_2\mathcal{M}\sigma_1. \quad (74)$$

Here the underlined symbols $\underline{\mathcal{R}}$ and $\underline{\mathcal{L}}$ are associated with the points $x_{\underline{\mathcal{R}}\tau_1}$, $x_{\underline{\mathcal{L}}\tau_2}$ such that $x_{\underline{\mathcal{R}}\tau_1}$ collides with $d_{\mathcal{R}}$ at $\xi_{\tau}^{d_{\mathcal{R}}}$, while $x_{\underline{\mathcal{L}}\tau_2}$ collides with $d_{\mathcal{L}}$ at $\xi_{\tau}^{d_{\mathcal{L}}}$. We introduce the abbreviation

$$\omega_1 := \sigma_2\mathcal{L}\sigma_1(\mathcal{R}\sigma_2\mathcal{L}\sigma_1)^{k-3}\mathcal{R}\sigma_2, \quad \omega_2 := \sigma_1\mathcal{R}\sigma_2\mathcal{M}\sigma_1 \quad (75)$$

and rewrite the symbolic sequence (74) as $\mathcal{R}\omega_1\mathcal{L}\omega_2$. Then the points $x_{\underline{\mathcal{R}}\tau_1}$ and $x_{\underline{\mathcal{L}}\tau_2}$ become $x_{\mathcal{R}\omega_1\mathcal{L}\omega_2}$ and $x_{\mathcal{L}\omega_2\mathcal{R}\omega_1}$, respectively. Then in the neighborhood of $d_{\mathcal{R}}$ the iterate f^{kn} is given by

$$f^{kn}(x) \equiv \begin{cases} f_{\mathcal{M}\omega_1\mathcal{L}\omega_2}, & \tilde{d} < x \leq d_{\mathcal{R}}, \\ f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}, & d_{\mathcal{R}} < x < \hat{d} \end{cases}$$

Here \tilde{d} and \hat{d} are the preimages of $d_{\mathcal{L}}$ such that

$$f_{\mathcal{M}\omega_1}(\tilde{d}) = d_{\mathcal{L}}, \quad f_{\mathcal{R}\omega_1}(\hat{d}) = d_{\mathcal{L}}.$$

The map f^{kn} has in the neighborhood of $d_{\mathcal{R}}$ a local absorbing interval

$$J = [f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}^2(d_{\mathcal{R}}), f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}(d_{\mathcal{R}})] \subset [\tilde{d}, \hat{d}], \quad (76)$$

whenever the conditions

$$f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}(d_{\mathcal{R}}) < \hat{d}, \quad f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}^2(d_{\mathcal{R}}) > x_{\mathcal{M}\omega_1\mathcal{L}\omega_2}, \quad (77)$$

are held. The restriction $f^{kn}|_J$ is topologically conjugate to the skew tent map g with

$$\alpha = a_{\mathcal{M}\omega_1\mathcal{L}\omega_2}, \quad \beta = a_{\mathcal{R}\omega_1\mathcal{L}\omega_2}. \quad (78)$$

Taking into account that $f_{\mathcal{R}\omega_1}(f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}(d_{\mathcal{R}})) \equiv f_{\omega_1\mathcal{L}\omega_2\mathcal{R}\omega_1}(r)$, $f_{\mathcal{R}\omega_1\mathcal{L}\omega_2}^2(d_{\mathcal{R}}) = f_{\omega_1\mathcal{L}\omega_2}(f_{\omega_1\mathcal{L}\omega_2\mathcal{R}}(r))$, and $x_{\mathcal{M}\omega_1\mathcal{L}\omega_2} = f_{\omega_1\mathcal{L}\omega_2}(x_{\omega_1\mathcal{L}\omega_2\mathcal{M}})$, the conditions (77) can be rewritten as

$$f_{\omega_1\mathcal{L}\omega_2\mathcal{R}\omega_1}(r) < d_{\mathcal{L}}, \quad f_{\omega_1\mathcal{L}\omega_2\mathcal{R}}(r) > x_{\omega_1\mathcal{L}\omega_2\mathcal{M}}.$$

Therefore, the region $D_{st}^{\mathcal{R},k,\sigma}$ in which the asymptotic dynamics of f can be studied by using the skew tent map g with the slopes given in (78), is confined by the BCB boundary $\xi_{\mathcal{R}\omega_1\mathcal{L}\omega_2}$, the border

$$b_{st}^{\mathcal{R},k,\sigma} = \{p \in D_0 : f_{\omega_1\mathcal{L}\omega_2\mathcal{R}\omega_1}(r) = d_{\mathcal{L}}\}, \quad (79)$$

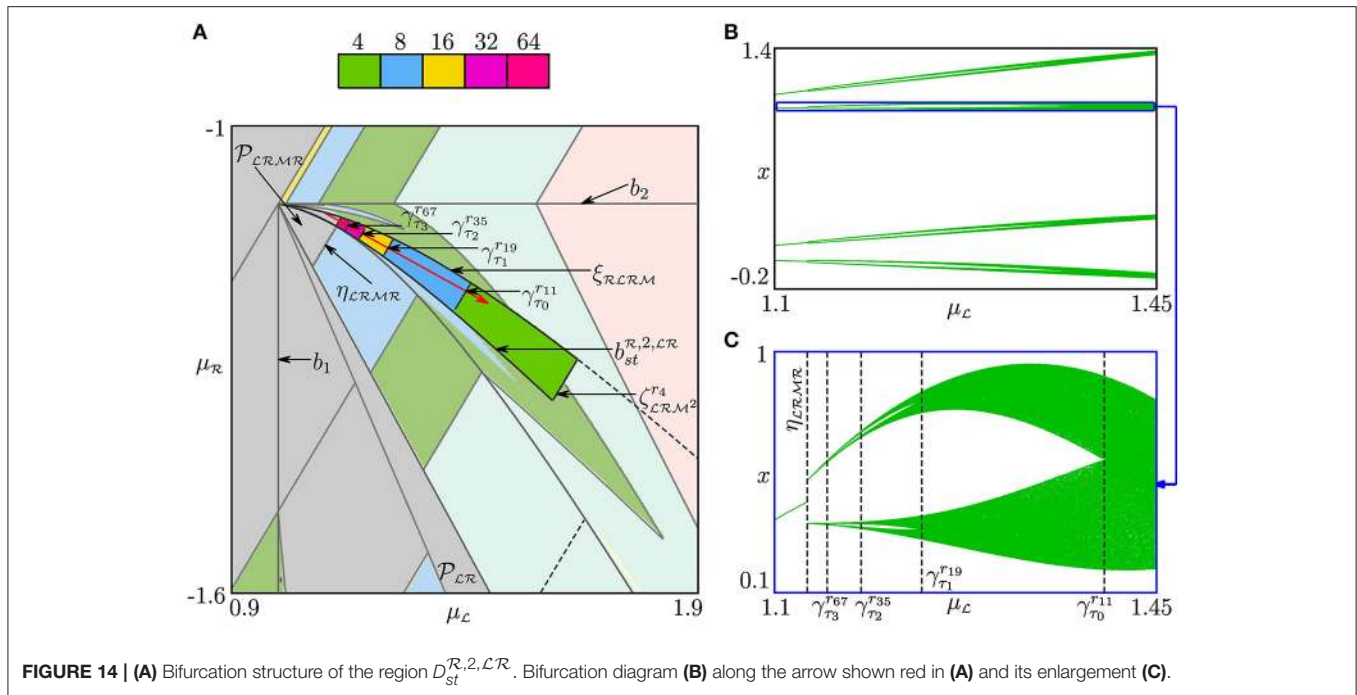
and the expansion bifurcation boundary

$$\zeta_{\omega_1\mathcal{L}\omega_2\mathcal{M}}^{r_{nk}} = \{f_{\omega_1\mathcal{L}\omega_2\mathcal{R}}(r) = x_{\omega_1\mathcal{L}\omega_2\mathcal{M}}\}. \quad (80)$$

Accordingly, the region $D_{st}^{\mathcal{R},k,\sigma}$ is given by

$$D_{st}^{\mathcal{R},k,\sigma} = \{x_{\mathcal{R}\omega_1\mathcal{L}\omega_2} > d_{\mathcal{R}}, f_{\omega_1\mathcal{L}\omega_2\mathcal{R}\omega_1}(r) < d_{\mathcal{L}}, f_{\omega_1\mathcal{L}\omega_2\mathcal{R}}(r) > x_{\omega_1\mathcal{L}\omega_2\mathcal{M}}\}. \quad (81)$$

To sum up we state



Proposition 5 ([17]). The region $D_{st}^{\mathcal{R},k,\sigma} \setminus \mathcal{P}_{\mathcal{R}\omega_1\mathcal{L}\omega_2}$ consists of $t+1$ chaoticity regions $C_{4\cdot 2^j}$, $j = 0, t$, where $t \geq 1$ depends on $a_{\mathcal{L}}$, $a_{\mathcal{R}}$. Each pair of contiguous chaoticity regions $C_{2^{j+1}}$ and C_{2^j} is separated by the merging bifurcation boundary $\gamma_{\tau_j}^{r_{2^{j+1}nk+nk-1}}$ which is derived from (20) applying the following replacement

$$\gamma_{\rho_j}^{c_{2^{j+1}}} \xrightarrow[A \rightarrow \mathcal{M}\omega_1\mathcal{L}\omega_2 \quad B \rightarrow \mathcal{R}\omega_1\mathcal{L}\omega_2]{c \rightarrow \mathcal{f}\omega_1\mathcal{L}\omega_2(r)} \gamma_{\tau_j}^{r_{2^{j+1}nk+nk-1}}, \quad (82)$$

where τ_j is obtained from the harmonic ρ_j (18) by replacing \mathcal{A} with $\mathcal{M}\omega_1\mathcal{L}\omega_2$ and \mathcal{B} with $\mathcal{R}\omega_1\mathcal{L}\omega_2$.

As before, to describe bifurcation structure of the parameter domain $D_{st}^{\mathcal{L},k,\sigma}$, located in the neighborhood of an arbitrary right fin region of a PA-region \mathcal{P}_σ , one can interchange \mathcal{L} and \mathcal{R} as well as the critical points ℓ and r in all analytic expressions.

6.3. Homoclinic Bifurcation Boundaries between $D_{st}^{\mathcal{L},\sigma}$ and $D_{st}^{\mathcal{R},k,\sigma}$

As one can see in Figure 10, between the regions $D_{st}^{\mathcal{L},\mathcal{LR}}$ and $D_{st}^{\mathcal{R},2,\mathcal{LR}}$ there are more bifurcation boundaries related to some homoclinic bifurcations. In contrast to the previous two sections, where chaotic intervals were associated with only one border point, $d_{\mathcal{L}}$ or $d_{\mathcal{R}}$, here we consider the case when chaotic intervals include both border points. This makes impossible direct usage of the results known for the skew tent map.

In Figures 15A,B we show the area between the regions $D_{st}^{\mathcal{L},\mathcal{LR}}$ and $D_{st}^{\mathcal{R},2,\mathcal{LR}}$. As one can observe there are several bifurcation boundaries issuing from the curves $b_{st}^{\mathcal{L},\mathcal{LR}}$ (59) and $b_{st}^{\mathcal{R},2,\mathcal{LR}}$ (70). Right below we briefly describe how these bifurcation boundaries can be derived.

First, recall that $b_{st}^{\mathcal{L},\mathcal{LR}}$ is related to the condition $f_{\mathcal{LR}}(d_{\mathcal{L}}) = f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})$. Until $f_{\mathcal{LR}}(d_{\mathcal{L}}) < f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})$ (that is, inside $D_{st}^{\mathcal{L},\mathcal{LR}}$), the local absorbing interval J of f^2 includes only a single kink point $d_{\mathcal{L}}$. Hence, $f^2|_J$ is defined by only two branches $f_{\mathcal{LR}}$ and $f_{\mathcal{MR}}$. When there holds the opposite inequality $f_{\mathcal{LR}}(d_{\mathcal{L}}) > f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})$, the absorbing interval J spreads over another kink point $f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})$ (to the right of $d_{\mathcal{L}}$) and becomes

$$J = [f_{\mathcal{MR}}(f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})), f_{\mathcal{LR}}(d_{\mathcal{L}})] = [r, f_{\mathcal{R}}(\ell)] \ni f_{\mathcal{M}}^{-1}(d_{\mathcal{R}})$$

(see Figure 15C). Then $f^2|_J$ consists of three branches: $f_{\mathcal{LR}}$, $f_{\mathcal{MR}}$, and $f_{\mathcal{MM}}$. This transformation is similar to the one described in Section 5.1 for the original map f , where the parameter point p crosses the border b_1 . Hence, the bifurcation boundaries emerging from $b_{st}^{\mathcal{L},\mathcal{LR}}$ are of the same sort as D_0 -prolongations. Applying the replacement

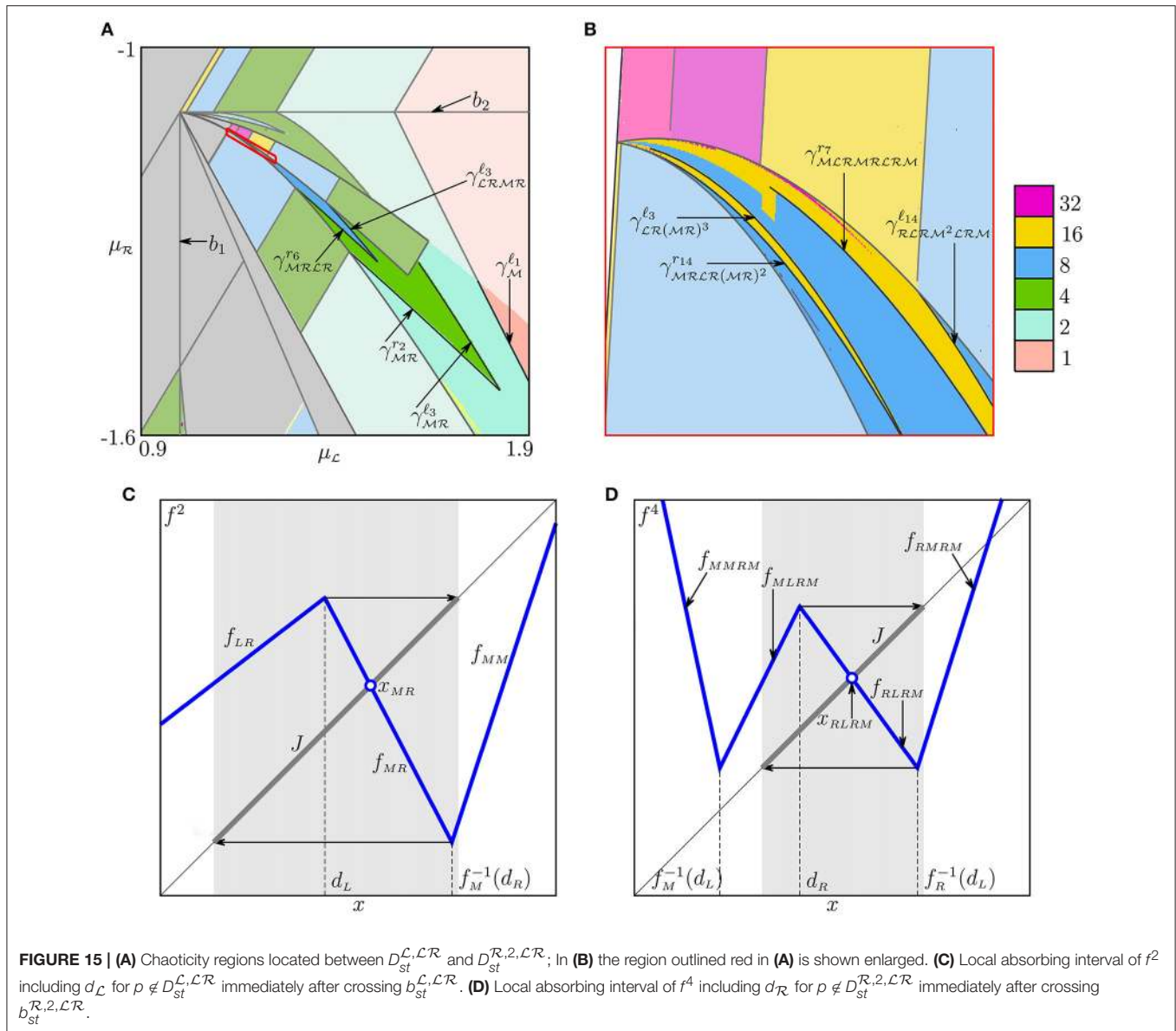
$$\gamma_{\rho_m}^{r_{2^{m+1}-1}} \xrightarrow[\mathcal{L} \rightarrow \mathcal{LR} \quad \mathcal{M} \rightarrow \mathcal{MR}]{r \rightarrow f_{\mathcal{R}}(\ell)} \gamma_{\tau_m}^{r_{2^{2^{m+1}-1}-1}} \quad (83)$$

to (51), one obtains the related analytic expressions for the mentioned boundaries.

Then we apply the replacement

$$\gamma_{\rho'_m}^{\ell_1} \xrightarrow[\mathcal{L} \rightarrow \mathcal{LR} \quad \mathcal{M} \rightarrow \mathcal{MR} \quad \mathcal{R} \rightarrow \mathcal{MM}]{\ell \rightarrow f_{\mathcal{R}}(\ell)} \gamma_{\tau'_m}^{\ell_3} \quad (84)$$

to the additional boundary $\gamma_{\rho'_m}^{\ell_1}$ (52). Here the sequences τ_m and τ'_m are obtained from the sequences ρ_m and ρ'_m , respectively, by replacing \mathcal{L} with \mathcal{LR} and \mathcal{M} with \mathcal{MR} . Recall, that ρ_m consists of the symbols \mathcal{L} and \mathcal{M} and is the m -th harmonic of \mathcal{M} [see (18) with $\mathcal{A} = \mathcal{L}$ and $\mathcal{B} = \mathcal{M}$], while ρ'_m is obtained from



ρ_m by shifting the first symbol \mathcal{M} to the end of the symbolic sequence. Note that in (83), r must be replaced by the value $f_{MR}(f_M^{-1}(d_R)) \equiv f_R(d_R)$ which coincides with r , and that is why the replacement for r is not necessary.

To illustrate the described procedure let us consider the boundary $\gamma_{MR}^{r_2}$ shown in **Figure 15A** which corresponds to the transition $Q_4 \Rightarrow Q_2$. For f^2 this transition is associated with the merging of 2-cyclic chaotic intervals into a single chaotic interval. Thus, we can apply the replacement (83) to the expression for the D_0 -prolongation $\gamma_M^{r_1}$ [i.e., (51) with $m = 0$] obtaining

$$\gamma_{MR}^{r_2} = \{p \in D_0 : f_{LR}(r) = x_{MR}\}. \quad (85)$$

Likewise, we derive the expression for the boundary $\gamma_{MR}^{r_6}$, which is associated with the transition $Q_8 \Rightarrow Q_4$, by using the

expression for the D_0 -prolongation $\gamma_{ML}^{r_3}$ [i.e., (51) with $m = 1$]:

$$\gamma_{MR}^{r_6} = \{p \in D_0 : f_{LR(MR)^2}(r) = x_{MR} \}. \quad (86)$$

The other boundary $\gamma_{LR}^{\ell_3}$ associated with the same transition is derived from the boundary $\gamma_{LM}^{\ell_1}$ (52), $m = 1$, corresponding to another homoclinic bifurcation of O_{ML} . Applying (84) to $\gamma_{LM}^{\ell_1}$ we obtain

$$\gamma_{LR}^{\ell_3} = \{p \in D_0 : f_{RMM}(\ell) = x_{LR} \}. \quad (87)$$

In **Figure 15B** one can also see the boundaries $\gamma_{MR}^{r_{14}}$ and $\gamma_{LR}^{\ell_3}$, which correspond to the transition $Q_{16} \Rightarrow Q_8$. Setting $m = 2$ and applying the replacement (83) to (51) and the

replacement (84) to (52) we get

$$\begin{aligned} \gamma_{MLM^2}^{r_7} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{LR}]{\mathcal{M} \rightarrow \mathcal{MR}} \gamma_{MR\mathcal{LR}(MR)^2}^{r_{14}}, \\ \gamma_{LM^3}^{\ell_1} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{LR}]{\mathcal{M} \rightarrow \mathcal{MR} \quad \mathcal{R} \rightarrow \mathcal{MM}}^{\ell \rightarrow f_{\mathcal{R}}(\ell)} \gamma_{\mathcal{LR}(MR)^3}^{\ell_3}, \end{aligned}$$

which implies

$$\begin{aligned} \gamma_{MR\mathcal{LR}(MR)^2}^{r_{14}} &= \{p \in D_0 : f_{\mathcal{LR}(MR)^2(MR\mathcal{LR})^2}(r) \\ &= x_{MR\mathcal{LR}(MR)^2}\}, \end{aligned} \tag{88}$$

$$\gamma_{\mathcal{LR}(MR)^3}^{\ell_3} = \{p \in D_0 : f_{\mathcal{RMM}}(\ell) = x_{\mathcal{LR}(MR)^3}\}. \tag{89}$$

Analogously we consider the case where the parameter point p leaves the region $D_{st}^{\mathcal{R},2,\mathcal{LR}}$ crossing $b_{st}^{\mathcal{R},2,\mathcal{LR}}$. When $p \in D_{st}^{\mathcal{R},2,\mathcal{LR}}$ the condition $f_{\mathcal{RLCRM}}(d_{\mathcal{R}}) > f_{\mathcal{R}}^{-1}(d_{\mathcal{L}})$ is satisfied and the local absorbing interval J of f^4 in the neighborhood of $d_{\mathcal{R}}$ includes only two linear branches. Otherwise, the interval J spreads over the third branch $f_{\mathcal{RMRM}}$ and becomes

$$\begin{aligned} J &= [f_{\mathcal{RLCRM}}(f_{\mathcal{R}}^{-1}(d_{\mathcal{L}})), f_{\mathcal{RLCRM}}(d_{\mathcal{R}})] \\ &\equiv [f_{\mathcal{RM}}(\ell), f_{\mathcal{LRM}}(r)] \ni f_{\mathcal{R}}^{-1}(d_{\mathcal{L}}). \end{aligned}$$

See **Figure 15D**. Thus, for f^4 crossing the border $b_{st}^{\mathcal{R},2,\mathcal{LR}}$ is similar to crossing b_1 for f . Consequently, the associated boundaries of chaoticity regions for $p \notin D_{st}^{\mathcal{R},2,\mathcal{LR}}$ can be derived as well from (51) and (52).

For example, the transition $\mathcal{Q}_{16} \Rightarrow \mathcal{Q}_8$ occurs at the boundaries $\gamma_{ML}^{r_3}$ and $\gamma_{MLCRM\mathcal{RLCRM}}^{r_7}$. We use the expressions for $\gamma_{ML}^{r_3}$ (51) and $\gamma_{LM}^{\ell_1}$ (52) with $m = 2$ to compute these mentioned boundaries. It is enough to apply the following replacements:

$$\begin{aligned} \gamma_{ML}^{r_3} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{MLCRM}]{r \rightarrow f_{\mathcal{RM}}(\ell)} \gamma_{\mathcal{RLCRM}^2\mathcal{LRM}}^{\ell_{14}}, \\ \gamma_{LM}^{\ell_1} &\xrightarrow[\mathcal{L} \rightarrow \mathcal{MLCRM}]{\ell \rightarrow f_{\mathcal{LRM}}(r)} \gamma_{MLCRM\mathcal{RLCRM}}^{r_7}. \end{aligned} \tag{90}$$

In this way we obtain

$$\begin{aligned} \gamma_{\mathcal{RLCRM}^2\mathcal{LRM}}^{\ell_{14}} &= \{p \in D_0 : f_{\mathcal{RMM}\mathcal{MLCRM}(\mathcal{RLCRM})^2}(\ell) \\ &= x_{\mathcal{RLCRM}^2\mathcal{LRM}}\}, \end{aligned} \tag{91a}$$

$$\begin{aligned} \gamma_{MLCRM\mathcal{RLCRM}}^{r_7} &= \{p \in D_0 : f_{\mathcal{L}(\mathcal{RM})^3}(r) \\ &= x_{MLCRM\mathcal{RLCRM}}\}. \end{aligned} \tag{91b}$$

Regarding the merging bifurcation boundaries $\gamma_{MR}^{\ell_3}$ and $\gamma_{M}^{\ell_1}$ marked in **Figure 15A**, it can be shown that $\gamma_{MR}^{\ell_3}$ is the D_0 -prolongation of $\gamma_{MR}^{r_4} \subset D_2$, and $\gamma_{M}^{\ell_1}$ is the D_0 -prolongation of $\gamma_{M}^{r_2} \subset D_2$.

So far all bifurcation boundaries of the regions C_2, C_4 , and C_8 related to the transitions $\mathcal{Q}_8 \Rightarrow \mathcal{Q}_4 \Rightarrow \mathcal{Q}_2 \Rightarrow \mathcal{Q}$ are

fully described. However, in **Figure 15B** one can see that the region C_{16} has further boundaries different from those described above. We can suggest that the regions $C_{2^m}, m \geq 4$, have much more complicated form and may have other boundaries, whose description falls beyond the scopes of the current paper.

The above example can be generalized. That is, for an arbitrary PA-region \mathcal{P}_{σ} we consider the are located between the regions $D_{st}^{\mathcal{L},\sigma}$ and $D_{st}^{\mathcal{R},k,\sigma}, k \geq 2$. The bifurcation boundaries inside this area are derived in a similar way by using the relevant replacements (for details see [17]).

7. CONCLUSIONS

In this work we carry on studying asymptotic dynamics of a 1D continuous bimodal PWL map f with the two outermost slopes being positive. Previously we have described certain bifurcation structures related to regular as well as to chaotic dynamics of f . Three basic bifurcation structures have been identified: the skew tent map (STM) structure, the period adding (PA) structure, and the fin structure. In case of STM structure the points of a typical orbit are located on two neighbor branches only, so that f is reduced to the skew tent map. In case of PA structure the orbits belong to the outermost branches only, so that the results known for the discontinuous map defined on two partitions with positive slopes can be used. The fin structure is constituted by periodicity regions that have a common boundary with PA-regions. The analytic expressions for bifurcation boundaries defined by BCBs, DFBs, as well as homoclinic bifurcations (such as merging and expansion) have been obtained.

In the current paper we have summarized the earlier results and described another family of tongues associated with regular dynamics observed in the parameter space of f . They emerge from the boundary b_1/b_2 that separates the parameter region D_1/D_2 where f is locally topologically conjugated to the skew tent map (STM structure) from the region D_0 where all three branches can be involved in asymptotic dynamics. These tongues are closely related to the cycles whose symbolic sequences form the U-sequence. For parameters belonging to D_1/D_2 these cycles are located in the two adjacent partitions and are unstable (except for the basic cycles having one point in the middle partition, which can be stable). However, for parameters in D_0 every such cycle has a new complementary cycle with a single point on the third branch. If the slope of the third branch is small enough, this new cycle is stable and one observes the related region in the parameter space of f . In the limit case when the third branch is horizontal (the slope is zero), for every U-sequence cycle from D_1/D_2 there exists the related periodicity region in D_0 . We refer to this situation as the “stabilization” of a U-sequence cycle.

AUTHOR CONTRIBUTIONS

The paper is the result of joint research, the contribution of every author is comparable to the others.

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