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Blue-LED-excitable NIR-II luminescent lanthanide-doped SrS nanoprobes for ratiometric thermal sensing

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ABSTRACT Lanthanide (Ln3+)-doped near infrared (NIR)-II luminescent nanoprobes have shown great promise in many technological fields, but are currently limited by the low absorption efficiency of Ln3+ due to the forbidden 4f→4f transition. Herein, we report a novel NIR-II luminescent nanoprobe based on efficient energy transfer from Ce³⁺ to Er³⁺ and Nd3+ in sub-10 nm SrS nanocrystals (NCs), which are excitable by using a commercial blue light-emitting diode (LED). Through sensitization by the allowed 4f→5d transition of Ce³⁺, the NCs exhibit strong NIR-II luminescence from Er³⁺ and Nd³⁺ with quantum yields of 2.9% and 2.3%, respectively. Furthermore, by utilizing the intense NIR-II luminescence of Er³⁺ from the thermally coupled Stark sublevels of ⁴I_{13/2}, we demonstrate the application of SrS:Ce3+/Er3+ NCs as blue-LED-excitable NIR-II luminescent nanoprobes for ratiometric thermal sensing. These findings reveal the unique advantages of SrS:Ln3+ NCs in NIR-II luminescence, which may open up a new avenue for exploring novel and versatile luminescent nanoprobes based on Ln3+-doped sulphide NCs.

Keywords: SrS, lanthanide, near-infrared II, nanoprobe, thermal sensing

INTRODUCTION

Lanthanide (Ln3+)-doped luminescent nanocrystals (NCs) emitting in the second near-infrared (NIR-II: 1000-1700 nm) biological window have recently evoked considerable interest, owing to their superior optical properties such as high photochemical stability, sharp emission peaks, long photoluminescence (PL) lifetimes, and large antenna-generated Stokes shift, in parallel with the benefits of minimal background interference and deep tissue penetration of the NIR light [1-9]. These outstanding features make Ln3+-doped NIR-II luminescent NCs ideal candidates as an alternative to traditional fluorescent probes like organic dyes and quantum dots and as a new generation of luminescent nanoprobes in many technological fields, including deep-tissue bioimaging, non-invasive chemical/biological detection, high-seed optical communication, and non-contact thermal sensing [10-22]. Nonetheless, because of the parity-forbidden nature of the $4f{ o}4f$ electronic transitions, Ln3+-doped NIR-II luminescent NCs normally suffer from low absorption and emission efficiencies, and exhibit low-to-medium brightness when compared with organic dyes and quantum dots [23–29]. Therefore, a high-power laser is generally demanded to realize bright NIR-II luminescence in Ln³+-doped NCs, which may limit their widespread applications.

To circumvent the limitation of Ln³+-doped NIR-II luminescent NCs, it is of fundamental importance to introduce an antenna that can effectively harvest the excitation light and sensitize the NIR-II luminescence of Ln³+ emitters [30–33]. In this regard, optical entities of allowed transitions with large absorption cross-sections such as Ce³+, Bi³+, ligand-to-metal charge transfer states, and the host absorption of semiconductors can be effective sensitizers for Ln³+ luminescence [34–37]. The absorption of these species, however, is influenced significantly by the host matrix with respect to the site symmetry, the crystal field (CF) strength, and the covalency [38–41]. Therefore, a judicious screening of host materials for both the sensitizers and Ln³+ emitters is essential to the design of efficient NIR-II luminescent nanoprobes with desired properties.

Hitherto, most of the reported Ln3+-doped NIR-II luminescent nanoprobes have been restricted to fluorides, probably due to the well-established synthetic methods for their upconversion analogues [42-45]. Until recently, the new class of NIR-II luminescent nanoprobes based on Ln3+-doped sulphides (namely, CaS and NaGdS2) have been proposed by our group [46,47]. In comparison with fluorides, sulphide NCs exert a stronger CF and a higher covalency on Ln3+ emitters [48-50]. As a result, the absorption and emission of Ce3+ and Eu2+ ions with allowed 4f→5d transitions locate at lower energies in sulphides (in the visible (vis) region) than in fluorides and oxides (in the ultraviolet (UV) region), which promises Ce3+ and Eu2+ activated sulphides as efficient phosphors in white light-emitting diodes (LEDs) [51-53]. Specifically, the absorption of Ce³⁺ in alkaline-earth sulfides (e.g., CaS and SrS) in the blue region of the spectrum matches well with the emission of the commercial blue-LED chip. This enables the development of blue-LEDexcitable NIR-II luminescent nanoprobes via energy transfer (ET) from Ce³⁺ to NIR-II Ln³⁺ emitters [46]. Moreover, owing to the strong CF level splitting, the luminescence of Ln^{3+} via $4f\rightarrow 4f$ transitions in the sulphide system is characterized by sharp CF emission peaks even at room temperature (RT), which may facilitate the devise of smart thermal sensors based on NIR-II

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luminescence of Ln^{3+} from the thermally coupled CF levels [54.55].

In this work, we develop a unique strategy for the controlled synthesis of Ln³+-doped SrS NCs via a high-temperature coprecipitation method. The optical properties, ET processes, and excited-state dynamics of Ln³+ in Ce³+ singly-doped, Ce³+/Er³+ and Ce³+/Nd³+ co-doped SrS NCs are systematically investigated through concentration- and temperature-dependent steady-state and transient PL spectroscopies. Through sensitization by Ce³+, efficient NIR-II luminescence from Er³+ and Nd³+ is achieved for the first time in sub-10 nm SrS NCs. Furthermore, by virtue of the well-resolved CF transition lines from $^4I_{13/2}$ of Er³+, we show the potential of SrS:Ce³+/Er³+ NCs as blue-LED-excitable NIR-II luminescent nanoprobes for non-contact thermal sensing with a high sensitivity.

EXPERIMENTAL SECTION

Chemicals and materials

The metal acetates of Sr(CH₃COO)₂·0.5H₂O (99.9%), Ce(CH₃-COO)₃·4H₂O (99.99%), Er(CH₃COO)₃·4H₂O (99.99%), and Nd(CH₃COO)₃·4H₂O (99.99%) were bought from Aladdin (China). The organic ligands of oleic acid (OA), oleylamine (OAm), 1-octadecene (ODE), and *N,N'*-diphenylthiourea (DPTU) were purchased from Sigma-Aldrich (China). Ethanol and cyclohexane were bought from Sinopharm (China). All chemicals were used as received without further purification.

Synthesis of SrS:Ln3+ NCs

Monodispersed SrS:Ln3+ NCs were prepared via a co-precipitation method at an elevated temperature. For synthesizing SrS:0.1% Ce³⁺ NCs, the metal acetates including 0.999 mmol of Sr(CH₃COO)₂·H₂O and 0.001 mmol of Ce(CH₃COO)₃·4H₂O (0.01 mol L-1 in aqueous solution, 100 µL) were first loaded into a 100-mL three-necked flask, followed by addition of organic ligands containing 4 mL of OA, 8 mL of OAm, and 8 mL of ODE. Then, the temperature was raised up to 120°C and the mixture was stirred at this temperature for 1h under flowing with N₂ to remove the residual water and oxygen. Thereafter, the temperature was raised up to 180°C and kept for 30 min to dissolve the metal acetates. The solution was then cooled down to RT and 3 mmol of DPTU (in 10 mL of ethanol) was added. After evaporation of ethanol at 80°C for 30 min, the solution was heated to 320°C and stirred for 1 h under N₂ flow to complete the reaction. The resulting mixture was then cooled down to RT and 30 mL of ethanol was added to precipitate the products. The NCs were collected by centrifugation at 8000 ×g for 6 min, and then washed with ethanol and cyclohexane (2:1 in volume) twice. The powder samples were obtained by drying the colloidal NCs at 60°C in a vacuum drying oven for 12 h.

Characterizations

Powder X-ray diffraction (XRD) measurement was performed on the MiniFlex 600 X-ray diffractometer (Rigaku) by using Cu K α 1 (λ = 0.154187 nm) as the radiation source. Transmission electron microscopy (TEM) images were acquired on a TECNAI G2 F20 TEM instrument. Inductively coupled plasma-atomic emission spectroscopy (ICP-AES, Ultima2, Jobin Yvon) and energy dispersive X-ray (EDX) spectrometer were used for the chemical composition analysis. The diffuse reflectance spectra were recorded on a UV/Vis/NIR spectrometer (Lambda 950,

Perkin-Elmer) and BaSO₄ was used as the reference. The FLS980 spectrometer (Edinburgh) was employed for recording the PL excitation and emission spectra and PL decay curves of the samples. A continuous and pulsed xenon lamp (450 W) was used as the excitation source. For absolute PL quantum yield (QY) measurement, the standard barium sulfate-coated integrating sphere (150 nm in diameter, Edinburgh) was mounted on the FLS980 spectrometer and utilized as the sample chamber. The measurements were repeated for three times to produce the average value and standard deviation. For low-temperature PL measurements, a closed cycle cryostat (10-300 K, DE202, Advanced Research Systems) was used as the sample chamber, and the slits of excitation and emission monochromators were set as small as possible to maximize the instrumental resolution. All the spectral data were recorded at RT by using the powder samples unless otherwise noted and corrected for the spectral response of the instruments.

RESULTS AND DISCUSSION

The SrS crystal has a face-centered cubic structure (space group $Fm\overline{3}m$) consisting of edge-sharing SrS₆ octahedra (Fig. 1a) [56]. Ln3+ ions occupy the octahedral Sr2+ site, with charge compensation by S^{2-} deficiency. Monodispersed SrS:Ln³⁺ NCs (Ln = Ce, Nd, and Er) were synthesized through a modified high-temperature co-precipitation method as we previously reported [46]. The as-synthesized NCs were hydrophobic owing to oleic ligands anchored to the surface, and thus can be readily dispersed in nonpolar organic solvents such as cyclohexane, forming a transparent colloidal solution. Upon doping with a trace amount (0.1 mol%) of Ce³⁺, the NCs exhibited bright bluegreen PL under 365-nm UV LED irradiation. XRD patterns of the NCs (Fig. 1b) display intense diffraction peaks which match well with cubic SrS (JCPDS No. 008-0489), indicating high crystallinity and phase purity of the resulting NCs. The diffraction peaks of the NCs turn broader with increasing Ln3+ doping concentration (Fig. S1), suggesting a size reduction of the NCs induced by Ln3+ doping. According to the Debye-Scherrer equation, the mean sizes of the NCs were estimated to be 10.9 \pm 0.4, 10.4 ± 0.7 , 9.6 ± 0.3 , and 9.9 ± 0.6 nm for SrS, SrS:0.1% Ce³⁺, SrS:0.1% Ce³⁺/0.5% Er³⁺, and SrS:0.1% Ce³⁺/0.2% Nd³⁺, respectively. Such a size reduction induced by Ln3+ doping can be ascribed to the charge compensation by extra S²⁻ which introduced transient electric dipoles to the NC surface with negative poles outwards. These transient electric dipoles hindered the diffusion of S²⁻ ions from the solution to the grain surface, thus retarding the growth of SrS NCs [57]. TEM images show that the NCs are monodispersed and roughly spherical with mean sizes of 10.4 \pm 1.1, 10.4 \pm 0.9, 9.5 \pm 0.9, and 9.7 \pm 0.9 nm for SrS, SrS:0.1% Ce^{3+} , SrS:0.1% $Ce^{3+}/0.5\%$ Er^{3+} , and SrS:0.1% $Ce^{3+}/0.2\%$ Nd³⁺ NCs, respectively (Fig. 1c-f and Fig. S2), in good agreement with those determined by XRD measurements. The representative high-resolution TEM (HRTEM) image of the NCs (inset of Fig. 1c) exhibits clear lattice fringes with an observed d spacing of 0.21 nm, corresponding to the (220) plane of cubic SrS. The selected-area electron diffraction (SAED) pattern of the NCs shows intense diffraction rings that can be well indexed into cubic SrS (inset of Fig. 1f). Composition analyses through EDX spectrum and ICP-AES identify the elements of Sr and S and 0.13 mol% of Ce³⁺ dopants in SrS:0.1% Ce³⁺ NCs (Fig. S3). All these results reveal high quality of the as-synthesized Ln³⁺-doped SrS NCs.

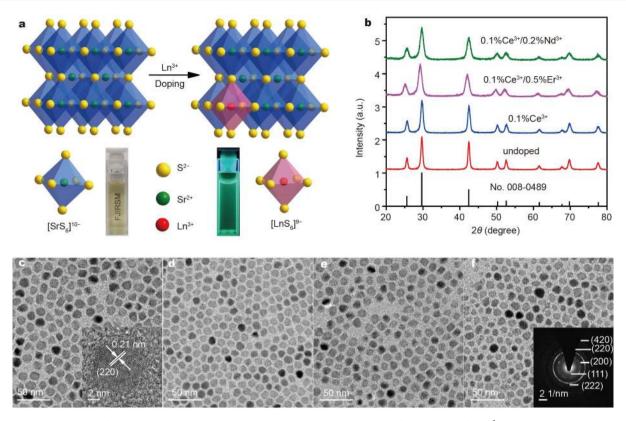


Figure 1 (a) Arrangement of strontium and sulphur atoms in the SrS structure and the crystallographic site for Ln³⁺ dopants. The photographs for the cyclohexane solution (1 mg mL⁻¹) of SrS:0.1% Ce³⁺ NCs under (left) ambient light and (right) 365-nm UV light are presented, showing transparency of the colloidal solution and its bright blue-green PL under UV light. (b) Powder XRD patterns of SrS, SrS:0.1% Ce³⁺, SrS:0.1% Ce³⁺/0.5% Er³⁺, and SrS:0.1% Ce³⁺/0.2% Nd³⁺ NCs. The vertical lines in the bottom represent the standard pattern of cubic SrS (JCPDS No. 008-0489). TEM images of (c) SrS, (d) SrS:0.1% Ce³⁺, (e) SrS:0.1% Ce³⁺/0.5% Er³⁺, and (f) SrS:0.1% Ce³⁺/0.2% Nd³⁺ NCs. The insets show the corresponding HRTEM image and SAED pattern of the NCs.

Optical absorption spectra show that the NCs have large absorbance in the UV region with a maximum at around 252 nm and a long absorption tail extending to 600 nm (Fig. S4). From the Tauc plot of the absorption spectra, the bandgap of the undoped SrS NCs was estimated to be 4.2 eV, which is generally consistent with that (4.3 eV) of bulk SrS [58]. Upon 0.1 mol% of Ce³⁺ doping, an additional absorption band at around 430 nm was explicitly observed, which can be assigned to the 4f→5d transition of octahedral Ce³⁺ in SrS NCs [59]. Under excitation at 410 nm, Ce³⁺ singly-doped NCs exhibited bright blue-green PL with two broad emission bands at 478 and 530 nm (Fig. 2a), corresponding to the electronic transitions of Ce3+ from the lowest ${}^2T_{2g}$ level of the 5d state to the ${}^2F_{5/2}$ and ${}^2F_{7/2}$ levels of the 4f ground state, respectively [59]. By monitoring the Ce³⁺ emission at 530 nm, two intense excitation bands at 263 and 430 nm were observed, ascribing to the host absorption of SrS and the 4f→5d transition of Ce³⁺, respectively [53]. Furthermore, it was observed that the PL intensity of the NCs increased gradually with increasing Ce3+ concentration from 0.01 to 0.1 mol% and then decreased at higher Ce3+ concentrations, indicative of an extremely low quenching concentration (0.1 mol%) of Ce³⁺ in SrS NCs. When the Ce³⁺ concentration exceeded 0.5 mol%, the emission bands of Ce3+ were red-shifted and broadened significantly, due to ET from Ce³⁺ ions emitting at higher energy to slightly perturbed Ce³⁺ ions emitting at lower energies [60], as also evidenced by the decreased PL lifetime of Ce^{3+} from 4.5 ns (0.1 mol%) to 3.9 ns (1 mol%) with increasing Ce³⁺ concentration (Fig. 2b and Table S1). The absolute PLQY

was determined to be 6.1% \pm 0.1% for SrS:0.1% Ce³⁺ NCs, which is higher than that (~2.9%) of 10.2-nm CaS:Ce³⁺ NCs we previously reported [46]. In comparison with CaS, the SrS host exerts weaker covalency and CF strength on Ce³⁺ ions due to the higher metallicity of Sr²⁺ than Ca²⁺. Therefore, the 5d state of Ce³⁺ locates at higher energies in SrS than in CaS, which results in a blueshift in the emission bands of Ce³⁺ (478 and 530 nm) in SrS relative to that (505 and 565 nm) in CaS (Fig. S5). As such, Ce³⁺ ions are less affected by the surface quenching effect and have a higher PLQY in SrS NCs than in CaS NCs with a similar size.

Fig. 2c and Fig. S6 show the temperature-dependent PL emission spectra (10–300 K) of SrS:0.1% Ce³⁺ NCs under excitation at 410 nm. The PL intensity of the NCs decreased gradually with the temperature rises due to the accelerated nonradiative relaxation and enhanced electron-phonon coupling of Ce³⁺ at higher temperatures [61], as further verified by the decreased PL lifetime of Ce³⁺ from 9.1 to 5.0 ns with elevating the temperature from 10 to 300 K (Fig. S7 and Table S2). Based on the Arrhenius plot of the integrated PL intensity of Ce³⁺ *versus* the inverse temperature (Fig. 2d) [62]:

$$I(T) = \frac{I_0}{1 + A\exp(-E_a/k_B T)},$$
(1)

where I(T) and I_0 are the integrated PL intensities at T and 0 K, respectively, E_a is the activation energy, and k_B is the Boltzman constant. The activation energy was determined to be 54.6 meV. The relatively small activation energy of Ce³⁺ in SrS NCs, along

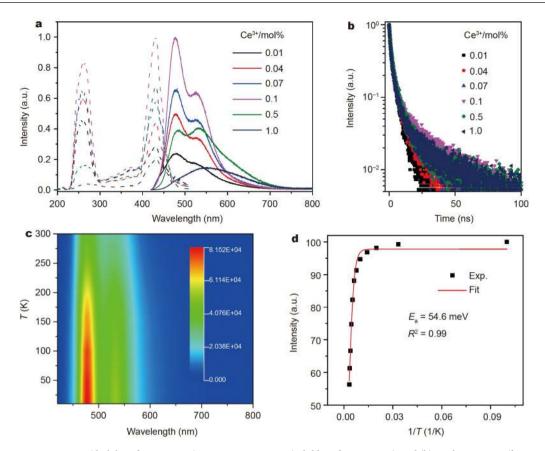


Figure 2 (a) PL excitation spectra (dash line, $\lambda_{em} = 530$ nm), PL emission spectra (solid line, $\lambda_{ex} = 410$ nm), and (b) PL decay curves ($\lambda_{em} = 530$ nm) of SrS:x% Ce³⁺ NCs with different Ce³⁺ concentrations. (c) Wavelength-temperature contour plot of the temperature-dependent PL emission spectra of SrS:0.1% Ce³⁺ NCs under excitation at 410 nm. (d) Integrated PL intensity for the Ce³⁺ emission as a function of temperature, whereby the activation energy (E_a) for the Ce³⁺ emission in SrS:0.1% Ce³⁺ NCs was derived.

with the high covalency and volatility of S^{2-} which result in abundant S^{2-} deficiency in the NCs [63], may account for the low quenching concentration (0.1 mol%) of Ce^{3+} in SrS NCs.

In addition to that of Ce3+ singly-doped SrS NCs, we also investigated the optical properties of Ce3+/Nd3+ and Ce3+/Er3+ co-doped SrS NCs. As shown in Fig. 3a, upon excitation to the Ce³⁺ 5d state at 430 nm, characteristic and sharp emission peaks from the ${}^{4}F_{3/2} \rightarrow {}^{4}I_{I}$ transitions of Nd³⁺ in NIR-I (700–950 nm) at 890 nm (${}^{4}I_{9/2}$) and NIR-II at 1081 nm (${}^{4}I_{11/2}$) and 1354 nm (${}^{4}I_{13/2}$) were detected in Ce³⁺/Nd³⁺ co-doped SrS NCs, along with the blue-green emission of Ce³⁺ at 478 and 530 nm. The PL intensity of Nd³⁺ increased gradually with increasing Nd³⁺ concentrations from 0 to 0.2 mol% and then decreased at higher Nd3+ concentrations owing to the concentration quenching effect [64], with the effective PL lifetime of Nd3+ decreasing from 3.6 ms (0.01 mol%) to 0.2 ms (2 mol%) (Fig. 3b and Table S3). Meanwhile, the PL intensity of Ce3+ decreased significantly with increasing Nd3+ concentration, indicating an efficient ET from Ce³⁺ to Nd³⁺ in the NCs. Such Ce³⁺-to-Nd³⁺ ET was further confirmed by the nearly identical excitation bands for the Ce³⁺ and Nd3+ emissions and the decreased PL lifetime of Ce3+ from 4.5 to 3.1 ns when the Nd^{3+} concentration increased from 0 to 2 mol% (Fig. 3c, d). From the PL lifetime of Ce³⁺, the efficiency of ET from Ce^{3+} to Nd^{3+} was calculated, with the highest value of 31.1% in SrS:0.1% Ce³⁺/2% Nd³⁺ NCs (Table S3). Fig. 3e depicts the ET processes from Ce³⁺ to Nd³⁺ in SrS NCs. Upon excitation at 430 nm, the Ce³⁺ ion is excited to the 5d state, and then ET occurs from the lowest ${}^2T_{2g}$ level of the Ce³⁺ 5d state to the ${}^4G_{5/2}/{}^4G_{7/2}$ levels of Nd³⁺, followed by nonradiative relaxation to ${}^4F_{3/2}$ of Nd³⁺ which results in intense NIR emission of Nd³⁺ through ${}^4F_{3/2} \rightarrow {}^4I_I$ (I = 9/2, 11/2, and 13/2) transitions.

Similarly, upon excitation at 430 nm, both the Ce³⁺ emission at 478 and 530 nm and the typical NIR-II emission from the 4I13/2 \rightarrow ⁴I_{15/2} transition of Er³⁺ at 1539 nm were observed in Ce³⁺/Er³⁺ co-doped SrS NCs (Fig. 4a). The PL intensity of Er³⁺ underwent a gradual increase with increasing Er³⁺ concentration from 0.01 to 0.5 mol% and then decreased at higher Er3+ concentrations at the expense of Ce3+ emission, as a result of efficient ET from Ce3+ to Er3+. Accordingly, the effective PL lifetimes of Er3+ and Ce^{3+} decreased from 13.0 ms (0.01 mol%) and 4.5 ns to 0.3 ms and 2.3 ns, respectively, as the Er3+ concentration increased from 0 to 5 mol% (Fig. 4b, c), due to the concentration quenching effect of Er3+ and Ce3+-to-Er3+ ET. PL excitation spectra of the NCs by monitoring the Er³⁺ emission at 1539 nm and the Ce³⁺ emission at 530 nm exhibited nearly identical excitation bands from the host absorption of SrS and the $4f\rightarrow 5d$ transition of Ce^{3+} (Fig. 4d), confirming the efficient Ce³⁺-to-Er³⁺ ET. From the PL lifetime of Ce3+, the efficiency of ET from Ce3+ to Er3+ was calculated, with the highest value of 48.9% in SrS:0.1% Ce³⁺/5% Er3+ NCs (Table S4). By contrast, negligibly weak PL was detected in Er3+ (or Nd3+) singly-doped SrS NCs. These results demonstrate unambiguously that it is the allowed 4f→5d tran-

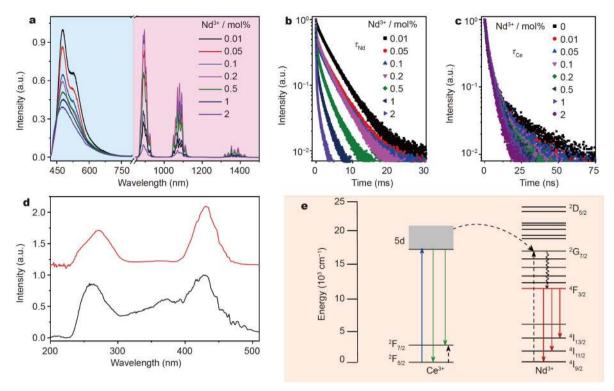


Figure 3 (a) Nd^{3+} concentration-dependent PL emission spectra ($\lambda_{ex} = 430 \text{ nm}$), and PL decay curves from (b) $^4F_{3/2}$ level of Nd^{3+} ($\lambda_{em} = 1081 \text{ nm}$) and (c) 5d state of Ce^{3+} ($\lambda_{em} = 530 \text{ nm}$) in SrS:0.1% Ce^{3+}/x % Nd^{3+} NCs. (d) PL excitation spectra of SrS:0.1% $Ce^{3+}/0.2$ % Nd^{3+} NCs by monitoring the Ce^{3+} (black) and Nd^{3+} (red) emissions at 530 and 1081 nm, respectively, showing the typical excitation bands from the host absorption of SrS and the $4f \rightarrow 5d$ transition of Ce^{3+} for both Ce^{3+} and Nd^{3+} emissions. (e) Simplified energy level schemes for Ce^{3+} and Nd^{3+} in SrS NCs, showing the proposed ET processes from Ce^{3+} to Nd^{3+} . The dashed, curve, and full arrows represent the ET, nonradiative relaxation and radiative transition processes, respectively.

sition of Ce^{3+} which harvests the excitation light and then sensitizes the NIR-II luminescence of Er^{3+} and Nd^{3+} in Ce^{3+}/Er^{3+} and Ce^{3+}/Nd^{3+} co-doped SrS NCs. As illustrated in Fig. 4e, upon excitation to the 5d state of Ce^{3+} , the excitation energy is transferred from the lowest ${}^2T_{2g}$ level of Ce^{3+} to the ${}^2H_{11/2}/{}^4S_{3/2}$ levels of Er^{3+} , followed by nonradiative (or radiative) relaxation to ${}^4I_{13/2}$ of Er^{3+} which leads to intense NIR-II emission of Er^{3+} through ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ transition. The PLQYs for the NIR-II emissions of Er^{3+} and Nd³⁺ sensitized by Ce^{3+} were estimated to be 2.9% \pm 0.1% in SrS:0.1% $Ce^{3+}/0.5\%$ Er^{3+} and 2.3% \pm 0.1% in SrS:0.1% $Ce^{3+}/0.2\%$ Nd³⁺ NCs. It is worth mentioning that, despite its currently low internal PLQYs, the NIR-II PL of Er^{3+} and Nd³⁺ sensitized by the allowed $4f \rightarrow 5d$ transition of Ce^{3+} in SrS NCs is much stronger than that in fluorides such as NaYF4: Yb³⁺/ Er^{3+} NCs via excitation based on the forbidden $4f \rightarrow 4f$ transitions.

To gain deep insights into the ET processes from Ce³⁺ to Er³⁺ and Nd³⁺ in Ce³⁺/Er³⁺ and Ce³⁺/Nd³⁺ co-doped SrS NCs, we carried out temperature-dependent steady-state and transient PL spectroscopic measurements. Fig. 5a shows the temperature-dependent PL emission spectra (10–300 K) of SrS:0.1% Ce³⁺/0.5% Er³⁺ NCs under excitation at 430 nm. It was observed that the integrated PL intensity and PL lifetime for the bluegreen emission of Ce³⁺ decreased gradually with increasing the temperature from 10 to 300 K (Fig. S8 and Table S5), due mainly to larger nonradiative transition probability of Ce³⁺ at higher temperatures, as also observed in Ce³⁺ singly-doped SrS NCs. For comparison, the integrated PL intensity for the NIR-II emission of Er³⁺ exhibited a slight decrease with increasing

temperature from 10 to 100 K and then increased steadily with the temperature rise (Fig. 5b), in parallel with the essentially unchanged PL lifetime of 6.5 ms (Fig. S9). Such a temperature evolution in PL intensity and PL lifetime of Er³⁺ in SrS:Ce³⁺/Er³⁺ NCs was different from those of Nd³⁺ in SrS:Ce³⁺/Nd³⁺ NCs, wherein both the PL intensities and PL lifetimes of Ce³⁺ and Nd³⁺ decreased significantly with the temperature rise (Fig. S10 and Table S6). This implies a phonon-assisted ET process from Ce³⁺ to Er³⁺ in SrS:Ce³⁺/Er³⁺ NCs, in view of the increased occupation number of the phonon state at higher temperatures [65].

Coincidently, we found that the intensities for the ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ CF transition lines of Er3+ experienced drastically distinct temperature dependence. Because of the increased occupation number of the phonon state at higher temperatures, the highlying CF sublevel (R₂) of ⁴I_{13/2} of Er³⁺ will be thermally populated from the low-lying one (R_1) with increasing temperature [54]. As a result, the intensities for the CF transitions of Er3+ from R2 of ⁴I_{13/2} to the higher-lying CF sublevels of ⁴I_{15/2} at 1551 and 1554 nm showed an increase with the temperature rise, whereas those from R₁ of ⁴I_{13/2} to the lower-lying CF sublevels of ⁴I_{15/2} at 1534, 1539, 1541, and 1544 nm decreased (inset of Fig. 5a). It is difficult to accurately identify all the ${}^4I_{13/2} \rightarrow {}^4I_{15/2}$ CF transition lines of Er3+ in SrS NCs, because of the negligibly weak PL upon direct excitation to Er³+ based on the forbidden 4f→4f transitions. Nonetheless, the CF emission peaks of Er3+ at 1539 and 1554 nm, which exhibit opposite temperature dependence in intensity (Fig. 5c), are sharp and well-resolved with little interference from other Stark components at temperatures below

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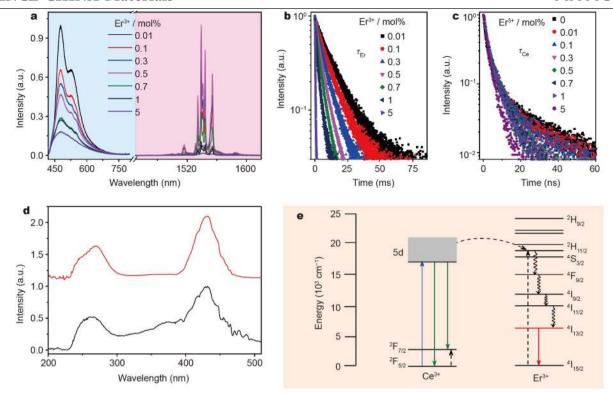


Figure 4 (a) Er^{3+} concentration-dependent PL emission spectra ($\lambda_{ex} = 430 \, \text{nm}$), and PL decay curves from (b) $^4I_{13/2}$ level of Er^{3+} ($\lambda_{em} = 1539 \, \text{nm}$) and (c) 5d state of Ce^{3+} ($\lambda_{em} = 530 \, \text{nm}$) in SrS:0.1% $Ce^{3+}/x\%$ Er^{3+} NCs. (d) PL excitation spectra of SrS:0.1% $Er^{3+}/x\%$ Er³⁺ NCs by monitoring the $Er^{3+}/x\%$ (red) emissions at 530 and 1539 nm, respectively, showing the typical excitation bands from the host absorption of SrS and the $4f \rightarrow 5d$ transition of $Er^{3+}/x\%$ for both $Er^{3+}/x\%$ emissions. (e) Simplified energy level schemes for $Er^{3+}/x\%$ and $Er^{3+}/x\%$ in SrS NCs, showing the proposed ET processes from $Er^{3+}/x\%$ The dashed, curve, and full arrows represent the ET, nonradiative relaxation and radiative transition processes, respectively.

300 K. This feature makes SrS: Ce^{3+}/Er^{3+} NCs an ideal nanoprobe candidate for ratiometric temperature sensing based on the temperature-dependent fluorescent intensity ratio (FIR) of Er^{3+} at 1554 and 1539 nm (I_{1554}/I_{1539}).

As shown in Fig. 5d, the FIR of I_{1554}/I_{1539} increased remarkably from 0.17 to 0.60 with increasing the temperature from 10 to 300 K, as a result of enhanced thermal population of the R_2 sublevel from R_1 of $^4I_{13/2}$ of Er^{3+} at higher temperatures. According to the theory of Boltzmann distribution, the FIR of the thermally coupled Stark sublevels (R_1 and R_2) of $^4I_{13/2}$ can be expressed as [66]

FIR =
$$\frac{I_{1554}}{I_{1539}} = C \exp\left(-\frac{\Delta E}{k_{\rm B}T}\right)$$
, (2)

where C is a constant dictated by the radiative transition rates, degeneracies and the emission cross-sections of R_1 and R_2 sublevels, and ΔE represents the energy gap between R_1 and R_2 . As a consequence, the natural logarithm of the FIR displayed a linear relationship with the inverse temperature (Fig. 5e). By linear fitting to the plot, the energy gap ΔE and the constant C were calculated to be $87.7 \, \mathrm{cm}^{-1}$ and 0.90, respectively. The temperature sensitivity (S) of the nanoprobes, defined as the rate of FIR change *versus* temperature change, is expressed as [67]

$$S(T) = \frac{d(FIR)}{dT} = C \left(\frac{\Delta E}{k_{\rm B} T^2} \right) \exp \left(-\frac{\Delta E}{k_{\rm B} T} \right). \tag{3}$$

By taking the values of C and ΔE into Equation (3), the sensitivity function S(T) versus T was derived and plotted in Fig. 5f. Typically, the maximum S value of 0.33% K^{-1} was achieved at

100 K, which is comparable to that of thermal sensing based on ${\rm Er^{3+}}$ -activated upconversion nanoprobes [68–70]. The high temperature sensitivity, along with the large absorbance of ${\rm Ce^{3+}}$ at 430 nm and the NIR-II emission of ${\rm Er^{3+}}$ in the optical communication window, enables ${\rm SrS:Ce^{3+}/Er^{3+}}$ NCs as blue-LED-excitable NIR-II luminescent nanoprobes for cryogenic temperature sensing in space and energy exploration [71–73].

CONCLUSIONS

In summary, we have systematically investigated Ln³⁺-doped SrS NCs from the controlled synthesis and optical properties to their potential application as a ratiometric temperature sensor. Upon doping with a trace amount (0.1 mol%) of Ce3+, the NCs exhibited bright blue-green PL at 478 and 530 nm with a PLQY of 6.1%. Through sensitization by the allowed 4f→5d transition of Ce³⁺ in the blue region, intense NIR-II luminescence from Er3+ at 1539 nm and Nd3+ at 1081 nm with PLQYs of 2.9% and 2.3%, respectively, was achieved in Ce3+/Er3+ and Ce3+/Nd3+ codoped SrS NCs. By means of temperature-dependent steadystate and transient PL spectroscopies, a phonon-assisted ET from Ce3+ to Er3+ was unveiled. Furthermore, owing to the strong CF level splitting exerted by SrS, the Er3+ luminescence was characterized by sharp and well-resolved CF emission peaks, which enables SrS:Ce3+/Er3+ NCs as blue-LED-excitable NIR-II luminescent nanoprobes for non-contact temperature detection based on the thermally coupled Stark sublevels of ⁴I_{13/2} of Er³⁺. These findings provide fundamental insights into the optical properties and excited-state dynamics of Ln3+ in SrS NCs, thus

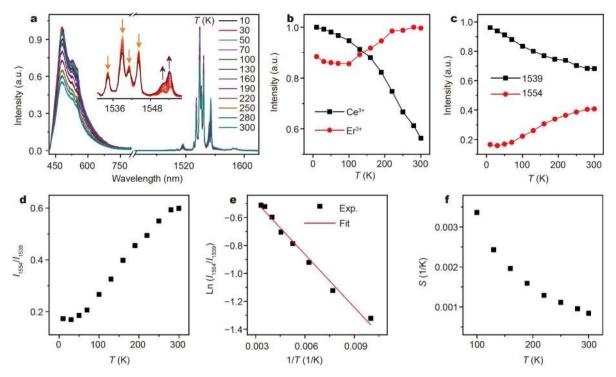


Figure 5 (a) Temperature-dependent PL emission spectra ($\lambda_{ex} = 430 \text{ nm}$) of SrS:0.1% $Ce^{3+}/0.5\%$ Er³⁺ NCs. The inset enlarges the ${}^{4}I_{13/2} \rightarrow {}^{4}I_{15/2}$ CF transitions of Er³⁺, showing drastically different temperature dependence for the CF transition lines of Er³⁺. Integrated PL intensities for (b) the green emission of Ce^{3+} , the NIR-II emission of Er³⁺, and (c) the ${}^{4}I_{13/2} \rightarrow {}^{4}I_{15/2}$ CF transitions of Er³⁺ at 1539 and 1554 nm as a function of temperature. (d) PL intensity ratio between the CF emission peaks of Er³⁺ at 1554 and 1539 nm (I_{1554}/I_{1539}) as a function of temperature. (e) Plot of the FIR of I_{1554}/I_{1539} versus the inverse temperature. (f) The temperature sensitivity (S) of SrS: Ce^{3+}/Er^{3+} nanoprobes as a function of temperature.

laying a foundation for future design of novel NIR-II luminescent nanoprobes based on Ln³⁺-doped sulphides for applications such as deep-tissue temperature detection.

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Author contributions Wei J, Zheng W and Chen X conceived the projects, wrote the paper and were primarily responsible for the experiments. Wei J, Liu Y and Zhang M carried out the synthesis and characterization of the NCs. Wei J, Huang P, Gong Z and Li R measured the PL spectra and analyzed the data. All authors contributed to the general discussion.

Conflict of interest The authors declare that they have no conflict of interest.

Supplementary information Supplementary data are available in the online version of the paper.



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蓝光LED可激发的稀土掺杂SrS近红外二区纳米荧光 探针及温度传感

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摘要 稀土掺杂近红外二区纳米荧光探针在许多技术领域都具有广泛的应用前景,但目前受限于稀土离子的4f→4f禁戒跃迁导致材料的吸收强度弱、发光效率低。本文报道了一种基于SrS纳米晶Ce³+到Er³+和Nd³+高效能量传递的新型蓝光LED可激发的近红外二区纳米荧光探针。通过Ce³+的4f→5d吸收允许跃迁敏化,该纳米晶可实现Er³+和Nd³+的高效近红外二区发光,其荧光量子产率分别为2.9%和2.3%。利用Er³+的⁴ $I_{13/2}$ 热耦合晶体场子能级的近红外二区发光,该纳米晶可作为一种蓝光LED激发的近红外二区纳米荧光探针用于比率式温度传感。这些结果展示了SrS:Ln³+纳米晶在近红外二区波段的发光优点,为多功能稀土硫化物基纳米荧光探针的设计开发提供了新思路。