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BOSE-EINSTEIN CONDENSATES

Microscopic magnetic-field imaging

Today's magnetic-field sensors¹ are not capable of making measurements with both high spatial resolution and good field sensitivity. For example, magnetic force microscopy² allows the investigation of magnetic structures with a spatial resolution in the nanometre range, but with low sensitivity, whereas SQUIDS³ and atomic magnetometers⁴ enable extremely sensitive magnetic-field measurements to be made, but at low resolution. Here we use one-dimensional Bose–Einstein condensates in a microscopic field-imaging technique that combines high spatial resolution (within 3 micrometres) with high field sensitivity (300 picotesla).

Trapped cold atoms are ideal magnetic

sensors as they are very sensitive to changes in magnetic-field landscapes, even in the presence of large homogeneous offset fields. Density modulations in trapped thermal atomic clouds have already been used as a measure of magnetic field variation caused by irregular current flow in nearby conductors^{5–8}. We have produced a versatile, high-resolution sensor based on Bose–Einstein condensates (BECs). Its sensitivity is not limited by the temperature T of the cloud, but is rather determined by the chemical potential μ of the condensate, which can be orders of magnitude lower than $k_B T$ (where k_B is Boltzmann's constant).

The principles of the technique are shown in Fig. 1a. A BEC is trapped at the measurement site so that its density profile can be directly imaged. The spatially varying density is a measure of the potential energy and hence of the local magnetic-field variation. To probe spatial magnetic-field variations, we start by confining a one-dimensional BEC (in which μ is smaller than an energy quantum of transverse excitation)⁹ in an elongated magnetic micro-trap with strong transverse and weak longitudinal confinement, created by small conductors mounted on the surface of an atom chip¹⁰.

As a demonstration, we measured the magnetic-field variations above the 100- μm -wide current-carrying wire used to create the trap itself. Scanning the position of the BEC enabled us to reconstruct a full two-dimensional magnetic-field profile (Fig. 1b) near the wire with unprecedented accuracy (sensitivity of 4 nanotesla) at the measurement resolution (3 μm). From this map, we reconstructed the local current flow in the wire¹¹ and found extremely small angular deviations (2×10^{-4} root-mean-square radians) from a straight current path (for details, see supplementary information).

We also investigated an independent field landscape by placing a BEC close (5 μm) to a test wire structure. As long as this structure is grounded and carries no current, the atomic-density profile is homogeneous within the detection sensitivity. This corresponds to an upper bound in potential roughness of less than 10^{-14} eV, corresponding to a temperature of 200 picokelvin (field sensitivity of 300 picotesla). As soon as a small current (about 5 mA) is passed through the wire, a characteristic field profile is imaged.

The technique is applicable not only to magnetic fields, but can also be used to detect variations in electrostatic fields, as can be shown

by charging the probed structure electrically. This measurement is more sensitive than the one discussed above because the trapping parameters can be adjusted independently from the measured potential landscape by using a separate wire for holding the BEC.

The optimal potential sensitivity of a BEC used as a field sensor, $\Delta B = \gamma \Delta N / z_0^3$, is achieved if the trapping parameters are adjusted so that the cloud's transverse size matches the desired spatial resolution z_0 . Here ΔN is the minimal atom-number variation resolved by the imaging system, and γ contains all the atomic-physics parameters of the specific atom ($\gamma = 8.63 \times 10^{-29}$ tesla cubic metres for the ⁸⁷Rb used in our experiment). Currently available CCD (charge-coupled device) cameras allow atom-shot noise-limited detection with a ΔN value of better than 10 atoms per pixel in absorption imaging, so that a sensitivity, ΔB , of 1 nanotesla is possible even at a high spatial resolution of 1 μm (or 1 picotesla at 10 μm). By changing to a different atom with higher mass and/or by reducing the interatomic interaction, a significant increase in sensitivity can be achieved.

A comparison of different magnetic-field measurement techniques^{1–4,12} (see supplementary information) shows that BECs as magnetic sensors could reach unprecedented sensitivity over a large range of spatial resolution. The sample measurements we present here reach higher sensitivities than those obtained with established techniques operating at the same spatial resolution.

Stephan Wildermuth*, Sebastian Hofferberth*, Igor Lesanovsky*, Elmar Haller*, L. Mauritz Andersson*, Sönke Groth*†, Israel Bar-Joseph†, Peter Krüger*, Jörg Schmiedmayer*

*Physikalisches Institut, Universität Heidelberg, Philosophenweg 12, 69120 Heidelberg, Germany e-mail: schmiedmayer@atomchip.org
 †Department of Condensed Matter Physics, The Weizmann Institute of Science, Rehovot 76100, Israel

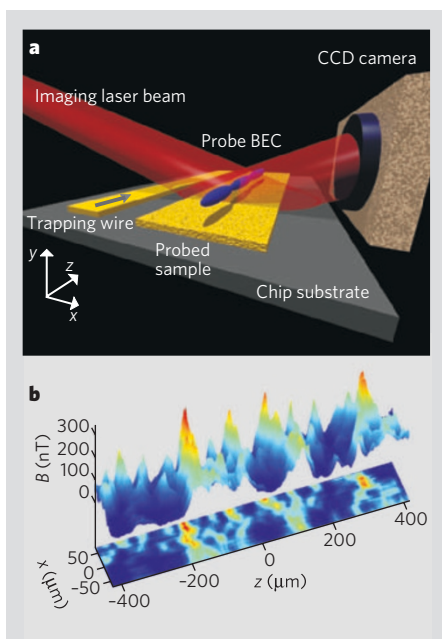


Figure 1 | One-dimensional Bose–Einstein condensate as a magnetic-field sensor.

a, Experimental set-up. A Bose–Einstein condensate is created and trapped by a current-carrying wire mounted on a silicon surface (atom chip) and is positioned above the sample to be probed. **b**, Two-dimensional scan of the magnetic landscape (field component along the wire direction) above a 100- μm -wide and 3.1- μm -tall gold wire. This profile has been reconstructed from 28 equally spaced one-dimensional atomic-density traces measured 10 μm above the current-carrying wire at a homogeneous offset field of 2 millitesla, so that relative field variations of only 4 p.p.m. stemming from slightly irregular current flow could be measured at a spatial resolution of 3 μm .

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