# Boundary Control for $2 \times 2$ Elliptic Systems with Conjugation Conditions 

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#### Abstract

In this paper, we consider $2 \times 2$ non-cooperative elliptic system involving Laplace operator defined on bounded, continuous and strictly Lipschitz domain of $R^{n}$. First we prove the existence and uniqueness for the state of the system under conjugation conditions; then we discuss the existence of the optimal control of boundary type with Neumann conditions, and we find the set of equations and inequalities that characterize it.


Keywords: Boundary Control; Elliptic Systems; Conjugation Conditions; Non-Cooperative Systems; Neumann Conditions

## 1. Introduction

So many optimal control problems governed by partial differential equations have been studied as in [1-3].

Systems governed by elliptic, parabolic, and hyperbolic operators have been considered, some of which are of distributed type as in [4-12], while some others are of boundary type as in [13-17].
Boundary control problems for non-cooperative $n \times n$ elliptic systems involving Laplace operator have been discussed in [17].

Here, using the theory of [3], we study the boundary control problem for $2 \times 2$ non-cooperative elliptic systems involving Laplace operator but under conjugation conditions.

Let us consider the following elliptic equations:

$$
\begin{cases}-\Delta y_{1}+y_{1}-y_{2}=f_{1} & x \in \Omega_{1} \cup \Omega_{2},  \tag{1.1}\\ -\Delta y_{2}+y_{1}+y_{2}=f_{2} & x \in \Omega_{1} \cup \Omega_{2},\end{cases}
$$

the heterogeneous boundary Neumann conditions:

$$
\begin{equation*}
\frac{\partial y_{1}}{\partial v_{A}}=g_{1}, \frac{\partial y_{2}}{\partial v_{A}}=g_{2}, x \in \Gamma \tag{1.2}
\end{equation*}
$$

and the conjugation conditions:

$$
\begin{equation*}
\left[\frac{\partial y_{1}}{\partial v_{A}}\right]=0,\left[\frac{\partial y_{2}}{\partial v_{A}}\right]=0, x \in \gamma, \tag{1.3}
\end{equation*}
$$

$$
\begin{equation*}
\left\{\frac{\partial y_{1}}{\partial v_{A}}\right\}^{ \pm}=r\left[y_{1}\right],\left\{\frac{\partial y_{2}}{\partial v_{A}}\right\}^{ \pm}=r\left[y_{2}\right], x \in \gamma \tag{1.4}
\end{equation*}
$$

where we have the following notations:
$\Omega$ is a domain that consists of two open, non-intersecting and strictly Lipschitz domains $\Omega_{1}$ and $\Omega_{2}$ from an n-dimensional real linear space $R^{n}$ i.e. $\Omega_{1}, \Omega_{2} \subset R^{n}$ are bounded, continuous, and strictly Lipschitz domains such that

$$
\Omega=\Omega_{1} \cup \Omega_{2}, \quad \Omega_{1} \cap \Omega_{2}=\phi \text { and } \bar{\Omega}=\bar{\Omega}_{1} \cup \bar{\Omega}_{2}
$$

Furthermore, $\quad \Gamma=\left(\partial \Omega_{1} \cup \partial \Omega_{2}\right) \backslash \gamma \quad$ is a boundary of a domain $\bar{\Omega}, \gamma=\partial \Omega_{1} \cap \partial \Omega_{2} \neq \phi$, and $\partial \Omega_{i}$ is a boundary of a domain $\Omega_{i}, i=1,2$.

In addition,

$$
\begin{gathered}
g_{i} \in L^{2}(\Gamma), f_{i} \in L^{2}(\Omega) \quad(i=1,2), \\
0 \leq r=r(x) \leq r_{1}<\infty, r \in C(\gamma),
\end{gathered}
$$

$r_{1}=$ constant, and $\nu$ is an ort of an outer normal to $\Gamma$.
Finally, $[\varphi]=\varphi^{+}-\varphi^{-}$,

$$
\begin{aligned}
& \varphi^{+}=\{\varphi\}^{+}=\varphi(x) \text { under } x \in \partial \Omega_{2} \cap \gamma, \\
& \varphi^{-}=\{\varphi\}^{-}=\varphi(x) \text { under } x \in \partial \Omega_{1} \cap \gamma .
\end{aligned}
$$

The model of system (1) is given by:

$$
\begin{aligned}
\operatorname{Ay}(x) & =A\left(y_{1}, y_{2}\right) \\
& =\left(-\Delta y_{1}+y_{1}-y_{2},-\Delta y_{2}+y_{1}+y_{2}\right)
\end{aligned}
$$

$$
A:\left(H^{1}(\Omega)\right)^{2} \rightarrow\left(L^{2}(\Omega)\right)^{2}
$$

system (1) is called non-cooperative, since the coefficients took the previous form.

We first prove the existence and uniqueness for the state of system (1), then we formulate the control problem. We also prove the existence and uniqueness of the optimal control of boundary type, and we discuss the necessary and sufficient conditions of the optimality.

## 2. The Existence and Uniqueness for the State of System (1)

Since $H^{1}(\Omega) \subseteq L^{2}(\Omega) \subseteq\left(H^{1}(\Omega)\right)^{\prime}$, then by Cartesian product we have the following chain [1]:

$$
\left(H^{1}(\Omega)\right)^{2} \subseteq\left(L^{2}(\Omega)\right)^{2} \subseteq\left(\left(H^{1}(\Omega)\right)^{\prime}\right)^{2}
$$

On $\left(H^{1}(\Omega)\right)^{2} \times\left(H^{1}(\Omega)\right)^{2}$, we define the following bilinear form:

$$
\begin{align*}
a(y, \psi)= & \int_{\Omega} \nabla y_{1} \nabla \psi_{1} \mathrm{~d} x+\int_{\Omega} \nabla y_{2} \nabla \psi_{2} \mathrm{~d} x \\
& +\int_{\Omega}\left(y_{1} \psi_{1}-y_{2} \psi_{1}+y_{1} \psi_{2}+y_{2} \psi_{2}\right) \mathrm{d} x  \tag{2}\\
& +\int_{\gamma} r\left[y_{1}\right]\left[\psi_{1}\right] \mathrm{d} \gamma+\int_{\gamma} r\left[y_{2}\right]\left[\psi_{2}\right] \mathrm{d} \gamma
\end{align*}
$$

The bilinear form (2) is continuous, since:

$$
\begin{aligned}
& |a(y, \psi)| \leq\left\|y_{1}\right\|_{H^{1}(\Omega)}\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+\left\|y_{2}\right\|_{H^{1}(\Omega)}\left\|\psi_{2}\right\|_{H^{1}(\Omega)} \\
& \quad+\left(\left\|y_{1}\right\|_{L^{2}(\Omega)}\left\|\psi_{1}\right\|_{L^{2}(\Omega)}+\left\|y_{2}\right\|_{L^{2}(\Omega)}\left\|\psi_{2}\right\|_{L^{2}(\Omega)}\right. \\
& \left.\quad+\left\|y_{1}\right\|_{L^{2}(\Omega)}\left\|\psi_{2}\right\|_{L^{2}(\Omega)}+\left\|y_{2}\right\|_{L^{2}(\Omega)}\left\|\psi_{1}\right\|_{L^{2}(\Omega)}\right) \\
& \quad+r\left\|\left[y_{1}\right]\right\|_{L^{2}(\gamma)}\left\|\left[\psi_{1}\right]\right\|_{L^{2}(\gamma)}+r\left\|\left[y_{2}\right]\right\|_{L^{2}(\gamma)}\left\|\left[\psi_{2}\right]\right\|_{L^{2}(\gamma)},
\end{aligned}
$$

since the inequalities $\|y\|_{L^{2}(\Omega)} \leq c\|y\|_{H^{1}(\Omega)}$ and $\|y\|_{L^{2}(\gamma)} \leq c_{1}\|y\|_{H^{1}(\Omega)}$ are true [3]. Then we have:

$$
\begin{aligned}
& |a(y, \psi)| \\
& \leq K\left(\left\|y_{1}\right\|_{H^{1}(\Omega)}\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+\left\|y_{2}\right\|_{H^{1}(\Omega)}\left\|\psi_{2}\right\|_{H^{1}(\Omega)}\right. \\
& \left.+\left\|y_{1}\right\|_{H^{1}(\Omega)}\left\|\psi_{2}\right\|_{H^{1}(\Omega)}+\left\|y_{2}\right\|_{H^{1}(\Omega)}\left\|\psi_{1}\right\|_{H^{1}(\Omega)}\right) \\
& \leq K\left(\left\|y_{1}\right\|_{H^{1}(\Omega)}+\left\|y_{2}\right\|_{H^{1}(\Omega)}\right)\left(\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+\left\|\psi_{2}\right\|_{H^{1}(\Omega)}\right) \\
& \leq K\|y\|_{\left.H^{1}(\Omega)\right)^{2}}\|\psi\|_{\left(H^{1}(\Omega)\right)^{2}}, \quad K \text { is constant. }
\end{aligned}
$$

Now, we have the following lemma:

## Lemma 1:

The bilinear form (2) is coercive on $\left(H^{1}(\Omega)\right)^{2}$, that is, there exists $\lambda \in R$, such that:

$$
a(y, y) \geq K_{1}\|y\|_{\left(H^{1}(\Omega)\right)^{2}}^{2}, \quad K_{1}>0
$$

## Proof:

$$
\begin{aligned}
a(y, y)= & \int_{\Omega}\left(\nabla y_{1}\right)^{2}+y_{1}^{2} \mathrm{~d} x+\int_{\Omega}\left(\nabla y_{2}\right)^{2}+y_{2}^{2} \mathrm{~d} x \\
& +\int_{\gamma} r\left(\left[y_{1}\right]^{2}+\left[y_{2}\right]^{2}\right) \mathrm{d} \gamma \\
\geq & \left\|y_{1}\right\|_{H^{1}(\Omega)}^{2}+\left\|y_{2}\right\|_{H^{1}(\Omega)}^{2}, \quad(\text { since } r \geq 0)
\end{aligned}
$$

hence

$$
\begin{equation*}
a(y, y) \geq\|y\|_{\left(H^{1}(\Omega)\right)^{2}}^{2} \tag{3}
\end{equation*}
$$

which proves the coerciveness condition of the bilinear form (2). Then we have the following theorem:

Theorem 1:
For a given $f=\left(f_{1}, f_{2}\right) \in\left(L^{2}(\Omega)\right)^{2}$, there exists a unique solution $y=\left(y_{1}, y_{2}\right) \in\left(H^{1}(\Omega)\right)^{2}$ for system (1).

## Proof:

Since (3) is hold, then by Lax-Milgram lemma, there exists a unique element

$$
y=\left(y_{1}, y_{2}\right) \quad y=\left(y_{1}, y_{2}\right) \in\left(H^{1}(\Omega)\right)^{2}
$$

such that

$$
\begin{equation*}
a(y, \psi)=L(\psi) \quad \forall \psi=\left(\psi_{1}, \psi_{2}\right) \in\left(H^{1}(\Omega)\right)^{2} \tag{4}
\end{equation*}
$$

where $L(\psi)$ is defined by:

$$
\begin{align*}
& L(\psi)=\int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x+\int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x+\int_{\Gamma} g_{1} \psi_{1} \mathrm{~d} \Gamma+\int_{\Gamma} g_{2} \psi_{2} \mathrm{~d} \Gamma, \\
& \forall \psi=\left(\psi_{1}, \psi_{2}\right) \in\left(H^{1}(\Omega)\right)^{2} . \tag{5}
\end{align*}
$$

The linear form (5) is continuous, since:

$$
\begin{aligned}
|L(\psi)|= & \left\|f_{1}\right\|_{L^{2}(\Omega)}\left\|\psi_{1}\right\|_{L^{2}(\Omega)}+\left\|f_{2}\right\|_{L^{2}(\Omega)}\left\|\psi_{2}\right\|_{L^{2}(\Omega)} \\
& +\left\|g_{1}\right\|_{L^{2}(\Gamma)}\left\|\psi_{1}\right\|_{L^{2}(\Gamma)}+\left\|g_{2}\right\|_{L^{2}(\Gamma)}\left\|\psi_{2}\right\|_{L^{2}(\Gamma)},
\end{aligned}
$$

since the inequalities $\|y\|_{L^{2}(\Omega)} \leq c\|y\|_{H^{1}(\Omega)}$
and $\|y\|_{L^{2}(\Gamma)} \leq c_{2}\|y\|_{H^{1}(\Omega)}$ are true [3], then:

$$
\begin{aligned}
|L(\psi)| & \leq c\left\|f_{1}\right\|_{L^{2}(\Omega)}\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+c\left\|f_{2}\right\|_{L^{2}(\Omega)}\left\|\psi_{2}\right\|_{H^{1}(\Omega)} \\
& +c_{2}\left\|g_{1}\right\|_{L^{2}(\Gamma)}\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+c_{2}\left\|g_{2}\right\|_{L^{2}(\Gamma)}\left\|\psi_{2}\right\|_{H^{1}(\Omega)} \\
& \leq\left(c\left\|f_{1}\right\|_{L^{2}(\Omega)}+c_{2}\left\|g_{1}\right\|_{L^{2}(\Gamma)}\right)\left\|\psi_{1}\right\|_{H^{1}(\Omega)} \\
& +\left(c\left\|f_{2}\right\|_{L^{2}(\Omega)}+c_{2}\left\|g_{2}\right\|_{L^{2}(\Gamma)}\right)\left\|\psi_{2}\right\|_{H^{1}(\Omega)},
\end{aligned}
$$

hence

$$
\begin{aligned}
|L(\psi)| & \leq K_{2}\left(\left\|\psi_{1}\right\|_{H^{1}(\Omega)}+\left\|\psi_{2}\right\|_{H^{1}(\Omega)}\right) \\
& =K_{2}\|\psi\|_{\left(H^{1}(\Omega)\right)^{2}} \cdot \quad\left(K_{2} \text { is constant }\right)
\end{aligned}
$$

Now, let us multiply both sides of first equation of (1.1) by $\psi_{1}(x)$, and the second equation by $\psi_{2}(x)$ then integration over $\Omega$, we have:

$$
\begin{aligned}
& \int_{\Omega}\left(-\Delta y_{1}+y_{1}-y_{2}\right) \psi_{1} \mathrm{~d} x=\int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x, \\
& \int_{\Omega}\left(-\Delta y_{2}+y_{1}+y_{2}\right) \psi_{2} \mathrm{~d} x=\int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x .
\end{aligned}
$$

By applying Green's formula:

$$
\begin{aligned}
& \int_{\Omega} \nabla y_{1} \nabla \psi_{1} \mathrm{~d} x-\int_{\Gamma} \psi_{1} \frac{\partial y_{1}}{\partial v_{A}} \mathrm{~d} \Gamma+\int_{\Omega}\left(y_{1}-y_{2}\right) \psi_{1} \mathrm{~d} x \\
& =\int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x, \\
& \int_{\Omega} \nabla y_{2} \nabla \psi_{2} \mathrm{~d} x-\int_{\Gamma} \psi_{2} \frac{\partial y_{2}}{\partial v_{A}} \mathrm{~d} \Gamma+\int_{\Omega}\left(y_{1}+y_{2}\right) \psi_{2} \mathrm{~d} x \\
& =\int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x,
\end{aligned}
$$

by sum the two equations, then comparing the summation with (2), (4) and (5) we obtain:

$$
\begin{aligned}
& \int_{\Omega} \nabla y_{1} \nabla \psi_{1} \mathrm{~d} x+\int_{\Omega} \nabla y_{2} \nabla \psi_{2} \mathrm{~d} x \\
& +\int_{\Omega}\left(y_{1}-y_{2}\right) \psi_{1} \mathrm{~d} x+\int_{\Omega}\left(y_{1}+y_{2}\right) \psi_{2} \mathrm{~d} x \\
& =\int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x+\int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x \\
& \quad+\int_{\Gamma} \psi_{1} \frac{\partial y_{1}}{\partial v_{A}} \mathrm{~d} \Gamma+\int_{\Gamma} \psi_{2} \frac{\partial y_{2}}{\partial v_{A}} \mathrm{~d} \Gamma
\end{aligned}
$$

then we deduce (1.2), which completes the proof.

## 3. Formulation of the Control Problem

The space $\left(L^{2}(\Gamma)\right)^{2}$ is the space of controls. For a control $u=\left(u_{1}, u_{2}\right) \in\left(L^{2}(\Gamma)\right)^{2}$, the state $y(u)=\left(y_{1}(u), y_{2}(u)\right) \in\left(H^{1}(\Omega)\right)^{2}$ of system (1) is given by the solution of the following systems:

$$
\begin{align*}
& \begin{cases}-\Delta y_{1}(u)+y_{1}(u)-y_{2}(u)=f_{1} & x \in \Omega_{1} \cup \Omega_{2}, \\
-\Delta y_{2}(u)+y_{1}(u)+y_{2}(u)=f_{2} & x \in \Omega_{1} \cup \Omega_{2},\end{cases}  \tag{6.1}\\
& \frac{\partial y_{1}}{\partial v_{A}}=g_{1}+u_{1}, \frac{\partial y_{2}}{\partial v_{A}}=g_{2}+u_{2}, x \in \Gamma, \tag{6.2}
\end{align*}
$$

and the conjugation conditions:

$$
\begin{align*}
& {\left[\frac{\partial y_{1}(u)}{\partial v_{A}}\right]=0,\left[\frac{\partial y_{2}(u)}{\partial v_{A}}\right]=0, x \in \gamma}  \tag{6.3}\\
& \left\{\frac{\partial y_{1}(u)}{\partial v_{A}}\right\}^{ \pm}=r\left[y_{1}\right],\left\{\frac{\partial y_{2}(u)}{\partial v_{A}}\right\}^{ \pm}=r\left[y_{2}\right], x \in \gamma \tag{6.4}
\end{align*}
$$

Since there exists a generalized solution $\mathrm{y}(u) \in\left(H^{1}(\Omega)\right)^{2}$ to the boundary value problem (6), then such solution is reasonable on $\Gamma$ of $\bar{\Omega}$, and

$$
\begin{equation*}
\|y(u)\|_{\left(L^{2}(\Gamma)\right)^{2}} \prec \infty \tag{7}
\end{equation*}
$$

The observation equation is given by:

$$
Z(u)=\left(Z_{1}(u), Z_{2}(u)\right)=C y(u)=C\left(y_{1}(u), y_{2}(u)\right)
$$

where $C \in \mathcal{L}\left(\left(L^{2}(\Gamma)\right)^{2} ;\left(L^{2}(\Gamma)\right)^{2}\right)$, namely:

$$
C y(u)=C\left(y_{1}(u), y_{2}(u)\right)=\left(y_{1}(u), y_{2}(u)\right)
$$

i.e.

$$
\begin{equation*}
Z(u)=\left(Z_{1}(u), Z_{2}(u)\right)=y(u)=\left(y_{1}(u), y_{2}(u)\right) \tag{8}
\end{equation*}
$$

For a given $z_{g}=\left(z_{g 1}, z_{g 2}\right) \in\left(L^{2}(\Gamma)\right)^{2}$, the cost function is given by

$$
\begin{align*}
J(u) & =\left\|y_{1}(u)-z_{g 1}\right\|_{L^{2}(\Gamma)}^{2}+\left\|y_{2}(u)-z_{g 2}\right\|_{L^{2}(\Gamma)}^{2}  \tag{9}\\
& +(N u, u)_{\left(L^{2}(\Gamma)\right)^{2}},
\end{align*}
$$

where $N u=\bar{a}(x) u, 0 \prec a_{0} \leq \bar{a}(x) \leq a_{1} \prec \infty$.
The function $y(u) \in\left(H^{1}(\Omega)\right)^{2}$ is specified on the domain $\bar{\Omega}_{1} \cup \bar{\Omega}_{2}$, minimizes the energy functional:

$$
\begin{align*}
\Phi(\psi)= & \int_{\Omega}\left(\nabla \psi_{1}\right)^{2}+\psi_{1}^{2} \mathrm{~d} x+\int_{\Omega}\left(\nabla \psi_{2}\right)^{2}+\psi_{2}^{2} \mathrm{~d} x \\
& +\int_{\gamma} r\left(\left[\psi_{1}\right]^{2}+\left[\psi_{2}\right]^{2}\right) \mathrm{d} \gamma \\
& -2 \int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x-2 \int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x-2 \int_{\Gamma} g_{1} \psi_{1} \mathrm{~d} \Gamma  \tag{10}\\
& -2 \int_{\Gamma} g_{2} \psi_{2} \mathrm{~d} \Gamma-2 \int_{\Gamma} u_{1} \psi_{1} \mathrm{~d} \Gamma-2 \int_{\Gamma} u_{2} \psi_{2} \mathrm{~d} \Gamma
\end{align*}
$$

on $\left(H^{1}(\Omega)\right)^{2}$, and it is the unique solution in $\left(H^{1}(\Omega)\right)^{2}$ to the weakly stated problem of finding an element $y(u) \in\left(H^{1}(\Omega)\right)^{2}$ that meets the following integral equation:

$$
\begin{align*}
& \int_{\Omega} \nabla y_{1} \nabla \psi_{1} \mathrm{~d} x+\int_{\Omega} \nabla y_{2} \nabla \psi_{2} \mathrm{~d} x \\
& \quad+\int_{\Omega} y_{1} \psi_{1}-y_{2} \psi_{1}+y_{1} \psi_{2}+y_{2} \psi_{2} \mathrm{~d} x \\
& \quad+\int_{\gamma} r\left[y_{1}\right]\left[\psi_{1}\right] \mathrm{d} \gamma+\int_{\gamma} r\left[y_{2}\right]\left[\psi_{2}\right] \mathrm{d} \gamma \\
& =\int_{\Omega} f_{1} \psi_{1} \mathrm{~d} x+\int_{\Omega} f_{2} \psi_{2} \mathrm{~d} x+\int_{\Gamma} g_{1} \psi_{1} \mathrm{~d} \Gamma \\
& \quad+\int_{\Gamma} g_{2} \psi_{2} \mathrm{~d} \Gamma+\int_{\Gamma} u_{1} \psi_{1} \mathrm{~d} \Gamma+\int_{\Gamma} u_{2} \psi_{2} \mathrm{~d} \Gamma \\
& \forall \psi(u)=\left(\psi_{1}(u), \psi_{2}(u)\right) \in\left(H^{1}(\Omega)\right)^{2} \tag{11}
\end{align*}
$$

The control problem then is to find $u=\left(u_{1}, u_{2}\right) \in U_{a d}$ such that $J(u) \leq J(v)$, where $U_{a d}$ is a closed convex subset of $\left(L^{2}(\Gamma)\right)^{2}$.

The cost function (9) can be written as (see [1]):

$$
\begin{align*}
J(u)= & \pi(u, u)-2 L(u)+\left\|y_{1}(0)-z_{g 1}\right\|_{L^{2}(\Gamma)}^{2}  \tag{12}\\
& +\left\|y_{2}(0)-z_{g 2}\right\|_{L^{2}(\Gamma)}^{2} .
\end{align*}
$$

In this case, the bilinear form $\pi(.,$.$) and the linear$ form $L($.$) are expressed as:$

$$
\begin{align*}
\pi(u, v)= & \left(y_{1}(u)-y_{1}(0), y_{1}(v)-y_{1}(0)\right)_{L^{2}(\Gamma)} \\
& +\left(y_{2}(u)-y_{2}(0), y_{2}(v)-y_{2}(0)\right)_{L^{2}(\Gamma)}  \tag{13}\\
& +(\bar{a} u, v)_{\left(L^{2}(\Gamma)\right)^{2}}, \\
L(v)= & \left(z_{g 1}-y_{1}(0), y_{1}(v)-y_{1}(0)\right)_{L^{2}(\Gamma)}  \tag{14}\\
+ & \left(z_{g 2}-y_{2}(0), y_{2}(v)-y_{2}(0)\right)_{L^{2}(\Gamma)}
\end{align*}
$$

Now, we prove the continuity of $\pi(u, v)$ and $L(v)$ on $\left(L^{2}(\Gamma)\right)^{2}$ as follows [3]:

Let $\tilde{y}^{\prime}=\tilde{y}\left(u^{\prime}\right)$ and $\tilde{y}^{\prime \prime}=\tilde{y}\left(u^{\prime \prime}\right)$ be solutions from $H^{1}(\Omega)$ to problem (11) under $f=0$ and $g=0$. Then from the bilinear form $a(.,$.$) which is given by (2), we$ can derive the following inequality:

$$
\begin{aligned}
& \left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{\left(H^{1}(\Omega)\right)^{2}}^{2} \\
& \leq \mu a\left(\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}, \tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right) \\
& \leq \mu\left\|u^{\prime}-u^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}\left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}
\end{aligned}
$$

Since $\|y\|_{L^{2}(\Gamma)} \leq c_{2}\|y\|_{H^{1}(\Omega)}$, then:

$$
\begin{aligned}
& \left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}^{2} \\
& \leq c_{2}\left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{H^{1}(\Omega)^{2}}^{2} \\
& \leq c_{2} \mu\left\|u^{\prime}-u^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}\left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}
\end{aligned}
$$

thus, we have:

$$
\left\|\tilde{y}^{\prime}-\tilde{y}^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}} \leq c_{2} \mu\left\|u^{\prime}-u^{\prime \prime}\right\|_{\left(L^{2}(\Gamma)\right)^{2}}
$$

i.e. the function $\tilde{y}(u)$ is continuously dependent on $u$.

Then the continuity of $\pi(u, v)$ and $L(v)$ on $\left(L^{2}(\Gamma)\right)^{2}$ is proved.

The bilinear form $\pi(u, v)$ is coercive on $\left(L^{2}(\Gamma)\right)^{2}$ since $(\bar{a} u, v)=(\sqrt{\bar{a}} u, \sqrt{\bar{a}} v)$,

Thus:

$$
\begin{aligned}
\pi(u, u)= & \left(y_{1}(u)-y_{1}(0), y_{1}(u)-y_{1}(0)\right)_{L^{2}(\Gamma)} \\
& +\left(y_{2}(u)-y_{2}(0), y_{2}(u)-y_{2}(0)\right)_{L^{2}(\Gamma)} \\
& +(\sqrt{\bar{a}} u, \sqrt{\bar{a}} u)_{\left(L^{2}(\Gamma)\right)^{2}} \\
\geq & a_{0}(u, u)_{\left(L^{2}(\Gamma)\right)^{2}} .
\end{aligned}
$$

Then by Lax Milgram lemma, the following theorem is proved. Moreover, it gives the necessary and sufficient conditions of optimality.

## Theorem 2:

Assume that (3) holds, there exists a unique optimal control $u=\left(u_{1}, u_{2}\right) \in U_{a d}$ that is closed convex subset of $\left(L^{2}(\Gamma)\right)^{2}$ and it is then characterized by the following equations and inequalities:

$$
\begin{align*}
& \left\{\begin{array}{l}
-\Delta p_{1}(u)+p_{1}(u)+p_{2}(u)=0, \\
-\Delta p_{2}(u)-p_{1}(u)+p_{2}(u)=0,
\end{array} \quad x \in \Omega_{1} \cup \Omega_{2},\right.  \tag{15.1}\\
& \frac{\partial p_{1}}{\partial v_{A^{*}}}=y_{1}(u)-z_{g 1}, \frac{\partial p_{2}}{\partial v_{A^{*}}}=y_{2}(u)-z_{g 2},  \tag{15.2}\\
& x \in \Gamma, \\
& {\left[\frac{\partial p_{1}(u)}{\partial v_{A^{*}}}\right]=0,\left[\frac{\partial p_{2}(u)}{\partial v_{A^{*}}}\right]=0, x \in \gamma,}  \tag{15.3}\\
& \left\{\frac{\partial p_{1}(u)}{\partial v}\right\}^{ \pm}=r\left[p_{1}\right],\left\{\frac{\partial p_{2}(u)}{\partial v}\right\}^{ \pm}=r\left[p_{2}\right], \tag{15.4}
\end{align*}
$$

$x \in \gamma$,

$$
\begin{equation*}
(p(u)+\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0 \tag{16}
\end{equation*}
$$

together with (6), where
$p(u)=\left(p_{1}(u), p_{2}(u)\right) \in\left(H^{1}(\Omega)\right)^{2}$ is the adjoint state.

## Proof:

The optimal control $u=\left(u_{1}, u_{2}\right) \in\left(L^{2}(\Gamma)\right)^{2}$ is characterized by (see [1])

$$
\begin{equation*}
\pi(u, v-u) \geq L(v-u) \quad \forall v \in U_{a d} \tag{17}
\end{equation*}
$$

by (13), and (14):

$$
\begin{aligned}
& \pi(u, v-u)-L(v-u) \\
& =\left(y(u)-z_{g}, y(v-u)-y(0)\right)_{\left(L^{2}(\Gamma)\right)^{2}} \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0,
\end{aligned}
$$

thus:

$$
\begin{aligned}
& \left(y(u)-z_{g}, y(v)-y(u)\right)_{\left(L^{2}(\Gamma)\right)^{2}} \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0
\end{aligned}
$$

this inequality can be written as

$$
\begin{align*}
& \left(y_{1}(u)-z_{g 1}, y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Gamma)} \\
& +\left(y_{2}(u)-z_{g_{2}}, y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Gamma)}  \tag{18}\\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0 . \tag{22}
\end{align*}
$$

Now, since: $\left(A^{*} p, y\right)=(p, A y)$, then:

$$
\begin{aligned}
& (p, A y)_{\left(L^{2}(\Omega)\right)^{2}} \\
& =\left(p_{1}(u),-\Delta y_{1}(u)+y_{1}(u)-y_{2}(u)\right)_{L^{2}(\Omega)} \\
& +\left(p_{2}(u),-\Delta y_{2}(u)+y_{1}(u)+y_{2}(u)\right)_{L^{2}(\Omega)}
\end{aligned}
$$

by using Green's formula, we obtain:

$$
\begin{align*}
& (p, A y)_{\left(L^{2}(\Omega)\right)^{2}} \\
& =\left(-\Delta p_{1}(u)+p_{1}(u)+p_{2}(u), y_{1}(u)\right)_{L^{2}(\Omega)} \\
& -\left(p_{1}, \frac{\partial y_{1}}{\partial v_{A}}\right)_{L^{2}(\Gamma)}+\left(y_{1}, \frac{\partial p_{1}}{\partial v_{A^{*}}}\right)_{L^{2}(\Gamma)} \\
& +\left(-\Delta p_{2}(u)-p_{1}(u)+p_{2}(u), y_{2}(u)\right)_{L^{2}(\Omega)} \\
& -\left(p_{2}, \frac{\partial y_{2}}{\partial v_{A}}\right)_{L^{2}(\Gamma)}+\left(y_{2}, \frac{\partial p_{2}}{\partial v_{A^{*}}}\right)_{L^{2}(\Gamma)}  \tag{23}\\
& =\left(A^{*} p, y\right)_{\left(L^{2}(\Omega)\right)^{2},} \tag{0}
\end{align*}
$$

then

$$
\begin{align*}
& A^{*} p=A^{*}\left(p_{1}, p_{2}\right) \\
& =\left(-\Delta p_{1}+p_{1}+p_{2},-\Delta p_{2}-p_{1}+p_{2}\right)  \tag{19}\\
& x \in \Omega_{1} \cup \Omega_{2}
\end{align*}
$$

Since the adjoint system takes the form [3]:

$$
\begin{gather*}
A^{*} p(v)=0, x \in \Omega  \tag{20.1}\\
\frac{\partial p}{\partial v_{A^{*}}}=y(v)-z_{g}, x \in \Gamma  \tag{20.2}\\
{\left[\frac{\partial p}{\partial v_{A^{*}}}\right]=0 \text { and }\left\{\frac{\partial p}{\partial v_{A^{*}}}\right\}^{ \pm}=r[p], x \in \gamma,} \tag{20.3}
\end{gather*}
$$

and by using (19), system (15) is proved .
From Green's formula the following equations are true:

$$
\begin{align*}
& \left(-\Delta p_{1}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)} \\
& =\left(\nabla p_{1}(u), \nabla\left(y_{1}(v)-y_{1}(u)\right)\right)_{L^{2}(\Omega)} \\
& \quad-\left(\frac{\partial p_{1}}{\partial v_{A^{*}}}, y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Gamma)},
\end{align*}
$$

Now, we transform (18) by using (15) as follows:

$$
\begin{aligned}
& \left(\frac{\partial p_{1}}{\partial v_{A^{*}}}, y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Gamma)} \\
& +\left(\frac{\partial p_{2}}{\partial v_{A^{*}}}, y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Gamma)} \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0,
\end{aligned}
$$

by (23) and (24), we have:
and

$$
\begin{align*}
& \left(\frac{\partial p_{2}}{\partial v_{A^{*}}}, y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Gamma)} \\
& =\left(\nabla p_{2}(u), \nabla\left(y_{2}(v)-y_{2}(u)\right)\right)_{L^{2}(\Omega)}  \tag{24}\\
& +\left(-p_{1}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)} \\
& +\left(p_{2}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)} .
\end{align*}
$$

by adding

$$
\begin{aligned}
& \left(p_{1}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)}, \\
& \left(p_{2}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)}
\end{aligned}
$$

to the both sides of Equation (21), and

$$
\begin{aligned}
& \left(-p_{1}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)} \\
& \left(p_{2}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)}
\end{aligned}
$$

to the both sides of Equation (22), then by (15) we obtain:

$$
\begin{aligned}
& \left(\frac{\partial p_{1}}{\partial v_{A^{*}}}, y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Gamma)} \\
& =\left(\nabla p_{1}(u), \nabla\left(y_{1}(v)-y_{1}(u)\right)\right)_{L^{2}(\Omega)} \\
& \quad+\left(p_{1}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)} \\
& \quad+\left(p_{2}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)},
\end{aligned}
$$

$$
\begin{aligned}
& \left(\nabla p_{1}(u), \nabla\left(y_{1}(v)-y_{1}(u)\right)\right)_{L^{2}(\Omega)} \\
& +\left(p_{1}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)} \\
& +\left(p_{2}(u), y_{1}(v)-y_{1}(u)\right)_{L^{2}(\Omega)} \\
& +\left(\nabla p_{2}(u), \nabla\left(y_{2}(v)-y_{2}(u)\right)\right)_{L^{2}(\Omega)} \\
& -\left(p_{1}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)} \\
& +\left(p_{2}(u), y_{2}(v)-y_{2}(u)\right)_{L^{2}(\Omega)} \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0,
\end{aligned}
$$

by using (2):

$$
\begin{aligned}
& a\left(y_{1}(v)-y_{1}(u), p_{1}(u)\right)-\int_{\gamma} r\left[y_{1}(v)-y_{1}(u)\right]\left[p_{1}\right] \mathrm{d} \gamma \\
& +a\left(y_{2}(v)-y_{2}(u), p_{2}(u)\right)-\int_{\gamma} r\left[y_{2}(v)-y_{2}(u)\right]\left[p_{2}\right] \mathrm{d} \gamma \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0,
\end{aligned}
$$

from (2), and using Green's formula:

$$
\begin{aligned}
& \left(p_{1}(u), A\left(y_{1}(v)-y_{1}(u)\right)\right)_{L^{2}(\Omega)} \\
& +\int_{\Gamma} p_{1} \frac{\partial\left(y_{1}(v)-y_{1}(u)\right)}{\partial v_{A}} \mathrm{~d} \Gamma \\
& -\int_{\gamma} r\left[y_{1}(v)-y_{1}(u)\right]\left[p_{1}\right] \mathrm{d} \gamma \\
& +\int_{\gamma} r\left[y_{1}(v)-y_{1}(u)\right]\left[p_{1}\right] \mathrm{d} \gamma \\
& +\left(p_{2}(u), A\left(y_{2}(v)-y_{2}(u)\right)\right)_{L^{2}(\Omega)} \\
& +\int_{\Gamma} p_{2} \frac{\partial\left(y_{2}(v)-y_{2}(u)\right)}{\partial v_{A}} \mathrm{~d} \Gamma \\
& -\int_{\gamma} r\left[y_{2}(v)-y_{2}(u)\right]\left[p_{2}\right] \mathrm{d} \gamma \\
& +\int_{\gamma} r\left[y_{2}(v)-y_{2}(u)\right]\left[p_{1}\right] \mathrm{d} \gamma \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2} \geq 0,}
\end{aligned}
$$

from (6), we obtain:

$$
\begin{aligned}
& \int_{\Gamma} p_{1}\left(v_{1}-u_{1}\right) \mathrm{d} \Gamma+\int_{\Gamma} p_{2}\left(v_{2}-u_{2}\right) \mathrm{d} \Gamma \\
& +(\bar{a} u, v-u)_{\left(L^{2}(\Gamma)\right)^{2}} \geq 0,
\end{aligned}
$$

which proves (16).

## Remark:

If the constraints are absent, i.e. when
$U_{a d}=\left(L^{2}(\Gamma)\right)^{2}$, then the equality:
$p(u)+\bar{a} u=0, x \in \Gamma$, follows from condition (16).
Hence

$$
\begin{equation*}
u_{1}=-\frac{p_{1}}{\bar{a}} \quad \text { and } \quad u_{2}=-\frac{p_{2}}{\bar{a}}, \quad x \in \Gamma . \tag{26}
\end{equation*}
$$

## 4. Conclusions

The main result of the paper contains necessary and sufficient conditions of optimality (of Pontryagin's type) for $2 \times 2$ elliptic systems under Neumann conjugation conditions involving Laplace operator defined on bounded, continuous and strictly Lipschitz domain of $R^{n}$, that give characterization of optimal control.

We can consider boundary control problems for $2 \times 2$ and $n \times n$ elliptic distributed systems with Dirichlet conjugation boundary conditions. Also we can consider boundary control problems for parabolic and hyperbolic distributed systems with Dirichlet and Neumann conjugation boundary conditions. The ideas mentioned above will be developed in forthcoming papers.

Also it is evident that by modifying:

- the boundary conditions,
- the nature of the control (distributed, boundary),
- the nature of the observation,
- the initial differential system,
many of variations on the above problem are possible to study with the help of Lions formalism.


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