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# BUCKLING OF A SUBLAMINATE IN A QUASI-ISOTROPIC COMPOSITE LAMINATE

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## SUMMARY

Buckling of a delaminated region can cause high interlaminar stresses which, in turn lead to delamination growth. Hence, buckling strain is an important parameter in assessing the potential for strength loss due to the delamination. The objective of this study was to predict the buckling of an elliptic delamination embedded near the surface of a thick quasi-isotropic laminate. The thickness of the delaminated ply group (the sublaminate) was assumed to be small compared to the total laminate thickness. Finite-element and Rayleigh-Ritz methods were used for the analyses. The Rayleigh-Ritz method was found to be simple, inexpensive, and accurate, except for highly anisotropic delaminated regions. Effects of delamination shape and orientation, material anisotropy, and layup on buckling strains were examined. Results showed that (1) the stress state around the delaminated region is biaxial, which may lead to buckling when the laminate is loaded in tension, (2) buckling strains for multi-directional fiber sublaminates generally are bounded by those for the 0 deg and 90 deg unidirectional sublaminates, and (3) the direction of elongation of the sublaminate that has the lowest buckling strain correlates with the delamination growth direction.

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#### INTRODUCTION

Composite laminates often contain delaminations. Causes of delamination are many and include tool drops, bird strikes, runway debris hits, and manufacturing defects. The shape of a delamination generally resembles an ellipse [1,2]. Delamination can significantly reduce the compressive strength and stiffness of the laminate, primarily because the delaminated region loses flexural stiffness. Even seemingly benign delaminations sometimes initiate localized buckling, thereby creating high interlaminar stresses and extensive delamination growth. Delamination growth may lead to structural instability.

Buckling and postbuckling of a near surface delamination has been addressed by several researchers both analytically and experimentally [1-7]. Analyses have been developed for strip [3,4], rectangular [1,5], circular [2] and elliptical [6,7] delamination shapes. These configurations are shown in figure 1. As indicated in figure 2 the buckled region is referred to herein as the "sublaminate." In refs. 1-7, the laminate was always loaded parallel to one of the principal axes of the sublaminate. The buckling of an elliptical sublaminate of arbitrary orientation and anisotropic material properties has not been addressed.

The objective of this paper is, therefore, to investigate the buckling characteristics of a wide variety of elliptical anisotropic sublaminates delaminated from a quasi-isotropic base laminate. The thickness of the sublaminate is assumed to be small compared to the base laminate. The sublaminate is assumed to be made up of laminated fiber reinforced composite with fibers oriented at various angles relative to the loading axis. A general purpose finite-element program, STAGS [8], was employed to calculate the buckling strain of typical sublaminate configurations. Also, a simple energy method, based on a Rayleigh-Ritz (R-R) formulation, was developed for the

parametric study. The R-R analysis was evaluated by comparing it to the finite element analysis. In this paper results are presented for variety of four ply sublaminates.

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# NOMENCLATURE

a	half-length of an elliptic sublaminate, m
$\left. \left. \begin{array}{c} A_{11}, A_{12}, A_{16} \\ A_{22}, A_{26}, A_{66} \end{array} \right\} \right.$	inplane stiffness coefficients of the sublaminate, N/m
b	half-width of an elliptic sublaminate, m
c <sub>0</sub>	generalized displacement, m
c <sub>1</sub> ,c <sub>2</sub>	generalized displacements, m <sup>-1</sup>
$\left. \begin{array}{c} D_{11}, D_{12}, D_{16} \\ D_{22}, D_{26}, D_{66} \end{array} \right\}$	flexural stiffness coefficients of the sublaminate, N/m
E	Young's modulus, MPa
G	shear modulus, MPa
h	thickness of the sublaminate, m
[K]	bending stiffness matrix of the sublaminate
$[K_1], [K_2], [K_3]$	bending stiffness coefficient matrices
[K <sub>g</sub> ]	geometric stiffness matrix of the sublaminate
N	number of laminae in the sublaminate
N <sub>x'</sub> ,N <sub>y'</sub> ,N <sub>x'y'</sub>	sublaminate stress resultants, N/m
U	strain energy of the deformed sublaminate, J
u	displacement in the x'-direction, m
v	potential energy due to loading, J
v	displacement in y'-direction, m
w	transverse (z) displacement, m
x-y-z	base laminate Cartesian coordinate system
x'-y'-z'	sublaminate Cartesian coordinate system

α	fiber angle of a lamina in a sublaminate, degree
ε	inplane strain
ε x	buckling strain
с Е Х_1	one-term buckling strain
θ	angle between x and x' axes
vlam	Poisson's ratio of the laminate
π	total potential energy of the sublaminate, J
Subscripts	
i	ply number
l,t	along or transverse to fiber direction
x, y	x or y direction
x',y'	x' or y' direction
Definitions	
Base laminate:	group of plies below the sublaminate
Buckling strain:	base laminate strain at which the sublaminate buckles
Delamination:	a separation between plies of a laminate
Sublaminate:	a set of intact plies separated from a thick laminate by a delamination

# DESCRIPTION OF THE PROBLEM

Figures 2(a) and 2(b) show plan and sectional views of a quasi-isotropic composite laminate with a single delamination. As shown, the x-y-z coordinate system is oriented such that the x-axis is along the laminate loading direction. The laminate is loaded to a strain  $\varepsilon_x$  in the x-direction. (A positive value of  $\varepsilon_x$  refers to tensile strain.) The associated strain in the y-direction is  $-v_{lam} \varepsilon_x$ , where  $v_{lam}$  is the base laminate Poisson's ratio.

The plan view of an assumed elliptic delamination is shown in figure 2(a). The set of intact laminae delaminated from the laminate is referred to as the "sublaminate;" the remaining laminate is referred to as the base laminate. The principal axes of the sublaminate are x' and y'; the corresponding semiaxes lengths are a and b. The angle between the axes x' and x is  $\theta$ , which is referred to herein as the sublaminate angle. The sublaminate is assumed to be made up of N laminae;  $\alpha$  is the fiber angle of a lamina measured relative to the x-axis (see figure 2(a)). The sublaminate thickness is h and is assumed to be small compared to the base laminate thickness. Therefore, inplane displacements around the sublaminate boundary can be calculated from the base laminate inplane deformations. Furthermore, the sublaminate lateral dimensions (a and b) are assumed to be relatively large compared to h, and, hence, thin plate linear buckling theory is assumed to be valid.

The buckled shape of a sublaminate is shown in figure 2(b). The transverse deflection w is measured from the sublaminate mid-plane. The transverse displacement and slopes are zero along the sublaminate boundary. The present analyses assume that the sublaminate buckles outward from the base laminate, as shown in figure 2(b). Furthermore, effects of higher modes and inward buckling of the sublaminate are neglected. The in-plane forces acting on the sublaminate due to the laminate strain  $\epsilon_{\rm x}$  were calculated from lamination theory [9] and are shown, schematically, in figure 2(c). Even though the base laminate is under uniaxial stress, the local stress distribution in the sublaminate is generally biaxial due to either a difference in the Poisson's ratio of the sublaminate and the base laminate, or to sublaminate material anisotropy. Note that the buckling initiation problem is reduced to that of an idealized elliptic sublaminate subjected to a set of inplane boundary forces.

#### ANALYSIS

Two analyses were used to calculate the buckling strains of the sublaminate: Finite-element (F-E) and Rayleigh-Ritz (R-R). Both analyses were based on the adjacent equilibrium buckling criterion (see for example ref. 10). The finite-element analysis was performed using the STAGS [8] computer code. The F-E analysis, although versatile and accurate, is a relatively expensive computer analysis, especially for eigen value problems involving a large number of degrees of freedom. Therefore, an additional simplified analysis based on the Rayleigh-Ritz technique was developed. The R-R analysis is known to give accurate results for specially orthotropic sublaminates [11]. The term special orthotropy refers to an orthotropic sublaminate, one of whose principal axes is aligned with the applied load direction. In this study, the R-R method was used for both specially orthotropic and anisotropic sublami-The accuracy of the analysis was assessed using the STAGS finitenates. element analysis. The two analyses are described in the following sections.

# The Finite Element Analysis

The general purpose finite element program STAGS [8] was used in the present study. An eighteen degree-of-freedom triangular plate element having three displacement and two rotational degrees-of-freedom (d.o.f.) at the corner nodes and one rotational d.o.f. at each mid-side node was employed for the analysis. Specially orthotropic cases were analyzed using only a quarter of the sublaminate whereas generally anisotropic cases were analyzed using the full sublaminate idealization. The two idealizations are shown in figure 3. The following boundary conditions were used in the analyses: (1) Quarterplate,  $v = w_{yy'} = 0$  along OA,  $u = w_{yx'} = 0$  along OB, and  $w = w_{yx'} = w_{yy'} = 0$  along AB; (2) Full plate,  $w = w_{yx'} = w_{yy'} = 0$  along the entire outer boundary. The variables u, v, and w are the displacements in the x',y' and

z directions, respectively, and the comma indicates partial differentiation  $(e \cdot g \cdot, w, x' = \frac{\partial w}{\partial x'})$ . For a specified base laminate axial strain  $e_x$ , the sub-laminate boundary forces (see fig. 2(c)) were input as initial stresses in the finite element analysis.

Convergence was studied to determine the degree of mesh refinement required for specially orthotropic and generally anisotropic sublaminates. A quarter plate sublaminate of  $8\times12$  (i.e., eight equal parts along the x'-axis and twelve equal parts along the circumference) mesh is shown in figure 3(a). Calculated buckling strains using  $4\times6$ ,  $8\times12$ , and  $16\times24$  meshes for a unidirectional 0 deg sublaminate indicate that the buckling strain for  $8\times12$  mesh differed by less than 1% from that for  $16\times24$  mesh. Thus, the  $8\times12$  mesh was employed for all isotropic and specially orthotropic sublaminates. A similar convergence study for a unidirectional with  $\alpha = 45$  deg sublaminate (which is anisotropic) indicated that the  $8\times32$  mesh (see fig. 3(b)) predicts buckling strain within 1% from that of the more refined mesh. Thus,  $8\times32$  mesh was employed for all anisotropic cases.

# The Raleigh-Ritz Analysis

Even though the finite element method is versatile it is expensive, especially for eigenvalue analysis of problems with a large number of degrees of freedom. Therefore a simple Raleigh-Ritz analysis, based on the Trefftz criterion [10], is presented to calculate buckling strains of elliptic sublaminates. The procedure consists of three steps: (1) selection of a kinematically admissible transverse displacement function; (2) calculation of the total potential energy; and then (3) application of the Trefftz criterion [10] to yield eigenvalue equations. The eigenvalue equations were solved numerically for different sublaminate configurations and stacking sequences. The governing equations are derived below.

Strain energy of sublaminate. - The strain energy U of the buckled sublaminate is [9]

$$U = \frac{1}{2} \int_{-a}^{a} \int_{-b}^{b} \left\{ D_{11} \left( \frac{\partial^{2} w}{\partial x'^{2}} \right)^{2} + D_{22} \left( \frac{\partial^{2} w}{\partial y'^{2}} \right)^{2} + 2D_{12} \left( \frac{\partial^{2} w}{\partial x'^{2}} \right) \left( \frac{\partial^{2} w}{\partial y'^{2}} \right) + 4D_{66} \left( \frac{\partial^{2} w}{\partial x'^{2} \partial y'} \right)^{2} + \left[ 4 D_{16} \left( \frac{\partial^{2} w}{\partial x'^{2}} \right) \left( \frac{\partial^{2} w}{\partial x'^{2} \partial y'} \right) + D_{26} \left( \frac{\partial^{2} w}{\partial y'^{2}} \right) \left( \frac{\partial^{2} w}{\partial x'^{2} \partial y'} \right) \right] \right] dx' dy'$$
(1)

where the D's are the flexural stiffness constants for an anisotropic sublaminate [9]. The transverse deflection w of the deformed sublaminate is assumed to given by

$$w = \left[1 - \left(\frac{x'}{a}\right)^2 - \left(\frac{y'}{b}\right)^2\right]^2 \left\{C_0 + C_1 x'^2 + C_2 y'^2\right\}$$
(2)

where  $C_0$ ,  $C_1$ , and  $C_2$  are generalized displacements. The w function satisfies zero deflection and slope conditions along the boundary of the sublaminate. Substituting eqn. (2) in eqn. (1) and performing the necessary differentiation and integration, eqn. (1) becomes

$$U = \{C\}^{T} [K]\{C\}$$
 (3)

where  $\{C\}^{T} = \{C_0, C_1, C_2\}$  and the stiffness matrix K is

$$[K] = D_{11}[K_1] + D_{22}[K_2] + 2[K_3](D_{12} + 2D_{66})$$
(4)

The matrices  $K_1$ ,  $K_2$  and  $K_3$  are defined as follows:

$$\begin{bmatrix} K_1 \end{bmatrix} = \frac{\pi b}{a^3} \begin{bmatrix} 4 & 0.5 & 0.5 \\ 1.35 & 0.05 \\ \text{Symmetric} & 0.15 \end{bmatrix}$$
$$\begin{bmatrix} K_2 \end{bmatrix} = \frac{\pi a}{b^3} \begin{bmatrix} 4 & 0.5 & 0.5 \\ 0.15 & 0.05 \\ \text{Symmetric} & 1.35 \end{bmatrix}$$
$$\begin{bmatrix} K_3 \end{bmatrix} = \frac{\pi}{6ab} \begin{bmatrix} 8 & 1 & 1 \\ 1.1 & 0.5 \\ \text{Symmetric} & 1.1 \end{bmatrix}$$

Since the assumed w deflection function is symmetric, coefficients of  $D_{16}$ and  $D_{26}$  vanish in the process of differentiation and integration; hence, these terms do not appear in eqn. (4).

Potential energy of applied loads. - The potential energy V of the stress resultants is given by [9]

$$V = \frac{1}{2} \int_{-a}^{a} \int_{-b}^{b} \left[ N_{x'}(w, y')^{2} + N_{y'}(w, y')^{2} + 2N_{x'y'}(w, y') \right] dx' dy' (5)$$

where the inplane stress resultants  $N_{x'}$ ,  $N_{y'}$ , and  $N_{x'y'}$  are obtained from lamination theory [9] as follows

$$N_{x'} = A_{11} \varepsilon_{x'} + A_{12} \varepsilon_{y'} + A_{1} \varepsilon_{x'y'}$$

$$N_{y'} = A_{12} \varepsilon_{x'} + A_{22} \varepsilon_{y'} + A_{26} \varepsilon_{x'y'}$$

$$N_{x'y'} = A_{16} \varepsilon_{x'} + A_{26} \varepsilon_{y'} + A_{66} \varepsilon_{x'y'}$$
(6)

;

The sublaminate strains  $\varepsilon_{x'}$ ,  $\varepsilon_{y}$ , and  $\varepsilon_{x'y'}$  are expressed in terms of the base laminate strain  $\varepsilon_{x}$ , the Poisson's ratio  $v_{lam}$ , and the sublaminate angle  $\theta$  as follows

$$\varepsilon_{x} = (\cos^{2}\theta - \nu_{lam} \sin^{2}\theta)\varepsilon_{x}$$

$$\varepsilon_{y} = (\sin^{2}\theta - \nu_{lam} \cos^{2}\theta)\varepsilon_{x}$$

$$\varepsilon_{x'y'} = -(1 + \nu_{lam}) \sin 2\theta \varepsilon_{x}$$
(7)

Substituting the expression for w from eqn. (2) into eqn. (5) and performing the necessary differentiation and integration, eqn. (5) simplifies to

$$V = \{C\}^{T} [K_{g}]\{C\} \varepsilon_{x}$$
(8)

The geometric stiffness matrix  $K_g$  is symmetric. The elements in the upper half of the  $[K_g]$  are defined as follows

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$$\begin{split} \kappa_{g_{11}} &= \frac{\pi}{3} \left( \frac{b}{a} \right) \left[ \left( A_{11} - v_{1am} A_{12} \right) + \left( \frac{a}{b} \right)^2 \left( A_{12} - v_{1am} A_{22} \right) \right] \\ \kappa_{g_{12}} &= \frac{\pi}{30ab} \left[ A_{12} - v_{1am} A_{22} \right] \\ \kappa_{g_{13}} &= \frac{\pi}{30ab} \left[ A_{11} - v_{1am} A_{12} \right] \\ \kappa_{g_{22}} &= \frac{\pi}{120} \left( \frac{b}{a} \right) \left[ 3 \left( A_{11} - v_{1am} A_{12} \right) + \left( \frac{a}{b} \right)^2 \left( A_{12} - v_{1am} A_{22} \right) \right] \\ \kappa_{g_{23}} &= 0 \\ \kappa_{g_{33}} &= \frac{\pi}{120} \left( \frac{a}{b} \right) \left[ \left( \frac{b}{a} \right)^2 \left( A_{11} - v_{1am} A_{12} \right) + 3 \left( A_{12} - v_{1am} A_{22} \right) \right] \end{split}$$
 (9)

The shear-extension coupling terms  $A_{16}$  and  $A_{26}$  do not appear in eqn. (9) because the w displacement function was symmetric.

<u>Buckling equation</u>. The total potential energy,  $\Pi$ , of the sublaminate is sum of the strain energy U and the potential energy V.

$$\Pi = U + V \tag{10}$$

Applying the Trefftz criterion [10], i.e.,  $\left| \frac{\partial^2 \Pi}{\partial C_i^2} \right| = 0$ , i = 0, 1, 2 yields the eigenvalue equation.

$$\left| \begin{bmatrix} \mathbf{K} \end{bmatrix} + \varepsilon_{\mathbf{X}} \begin{bmatrix} \mathbf{K} \\ \mathbf{g} \end{bmatrix} \right| = 0 \tag{11}$$

The solution of eqn. (11) results in three eigenvalues which correspond to the three buckling modes. The lowest absolute eigenvalue corresponds to the first buckling mode and is referred to as the buckling strain  $\varepsilon_x$  of the sublaminate.

To illustrate how the material anisotropy and the sublaminate size influence the buckling strain, eqn. (11) is simplified to that for a one term solution. This corresponds to an assumed deflection function  $w = \left[1 - \left(\frac{x'}{a}\right)^2 - \left(\frac{y'}{b}\right)^2\right]^2 C_0$ . Therefore, substituting  $C_1 = C_2 = 0$  in eqns. (3) and (5) and applying the Treffetz criterion yields

$$\varepsilon_{\mathbf{x}_{c1}} = -\frac{12}{b^2} \frac{\left[ \left(\frac{b}{a}\right)^2 D_{11} + \left(\frac{a}{b}\right)^2 D_{22} + \frac{2}{3} \left(D_{12} + 2D_{66}\right) \right]}{A_{11} + \left\{ \left(\frac{a}{b}\right)^2 - \nu_{1am} \right\} A_{12} - \nu_{1am} \left(\frac{a}{b}\right)^2 A_{22}}$$
(12)

where  $\varepsilon_{x}$  is the one term buckling strain. Equation (12) shows that the clsign of  $\varepsilon_{x}$  depends only on the sign of the denominator, which is a func $x_{cl}$  tion of the size and the in-plane stiffness coefficients of the sublaminate, and the laminate Poisson's ratio  $v_{lam}$ . If the denominator is negative, the sign of the buckling strain  $\varepsilon_{xcl}$  is positive. This indicates that the sublaminate may buckle when the base laminate is loaded in tension. This phenomenon may appear unusual, but it is caused by the mismatch of Poisson's ratio between sublaminate and the base laminate. For example, if a quasiisotropic base laminate having a Poisson's ratio of 0.3 is stretched to a strain  $\varepsilon_x$ , it contracts laterally (in the y-direction) by -0.3  $\varepsilon_x$ . A delaminated 90 deg fiber sublaminate, which has a very low Poisson's ratio (say 0.03), contracts only by -0.03  $\varepsilon_x$ . The unequal contraction in the y-direction introduces a net compressive strain of -0.27  $\varepsilon_x$  on the sublaminate. Such compressive strain acting in the fiber direction of the 90 deg fiber sublaminate may cause local buckling, even though the base laminate is loaded in tension.

Equation (12) also shows that the sign of the denominator depends on the aspect ratio b/a of the sublaminate. The effect of sublaminate material properties and aspect ratio on buckling is discussed further in the next section.

#### **RESULTS AND DISCUSSION**

Buckling strains of a surface sublaminate in a quasi-isotropic graphite/ epoxy composite laminate were calculated using the finite-element (F-E) and the Rayleigh-Ritz (R-R) analyses. The R-R analysis was compared with F-E analysis for typical cases. Buckling strains for four-ply sublaminates were calculated for various sublaminate sizes, orientations, fiber angles and layups. The plot of buckling strains versus sublaminate properties is referred to as the buckling curve. Material properties used in the analyses are given in Table 1.

<u>Isotropic sublaminate</u>.- Figure 4 compares buckling strains calculated from the F-E and R-R analyses for an isotropic (material properties for aluminum were used), elliptical sublaminate for different values of b. The F-E results are shown by symbols and the R-R results by a solid line. The two analyses agree very well. The compression buckling strain decreases in

magnitude monotonically as the sublaminate width, b, increases and asymptotically approaches the strip solution (b >> a).

Specially orthotropic sublaminates. - Figure 5 shows buckling curves for 0 deg and 90 deg unidirectional sublaminates. The 0 deg unidirectional sublaminate buckles only by compressive remote strain (as was the case for the isotropic sublaminate). The buckling strain asymptotically approaches a constant value with increased b. The 90 deg unidirectional sublaminate, as already explained with reference to eqn. (12), can buckle under tension and compression remote strain. The buckling caused by tension remote strain is referred to herein as tension buckling and buckling due to compression remote strain is referred to as compression buckling. Results in figure 5 and an examination of buckling modes (not reported here) indicate that the first mode of buckling for a 90 deg unidirectional sublaminate only occurred for remote tension for b < 25 mm and only for remote compression for b > 64 mm. Both compression and tension buckling of a sublaminate are possible for values of b between 25 and 64 mm. Higher modes of buckling involve both positive and negative lateral displacements. The base laminate would interfere with these negative lateral displacements of the sublaminate. To account for this interference requires the solution of a nonlinear contact problem, which is beyond the scope of the present paper. Hence, higher mode buckling strains are not reported here.

The tension buckling strain for b = 12.5 mm is 0.0051, which is higher than the splitting threshold for unidirectional graphite/epoxy laminate (transverse tensile strength = 40 MPa, ref. 12). Therefore, the 90 deg unidirectional sublaminate for b < 19 mm may split before buckling. The splitting would reduce the buckling strain. The two analyses agree very well with each other for both sublaminates.

Anisotropic sublaminates.- The lower half of figure 6 shows the effect of fiber angle  $\alpha$  on buckling strain for unidirectional fiber, circular sublaminates. The compression buckling strain increases with increasing fiber angle. But at fiber angles of 75 deg  $\langle \alpha \langle 90 deg$ , the sublaminate buckles also under tensile laminate strain. The F-E and R-R analyses agree well for  $0 \langle \alpha \langle 55 deg and 80 deg \langle \alpha \langle 90 deg$ . For fiber angles 55 deg  $\langle \alpha \langle 80 deg$ , the R-R analysis overestimates the buckling strain, because the bendingtwisting (D<sub>16</sub> and D<sub>26</sub>) and shear-extension (A<sub>16</sub> and A<sub>26</sub>) coupling terms were not included. However, for other fiber angles the influence of these terms is less significant and hence the R-R results agree with the F-E results.

The upper portion of figure 6 also shows the effect of fiber angle on major Poisson's ratio  $(v_{x'y'})$  of the sublaminate. The Poisson's ratio of the base laminate is also shown in figure 6 for reference. Comparison of the buckling curve and Poisson's ratio curves indicate (1) the buckling strain of the sublaminate is compressive and is relatively small for sublaminate Poisson's ratio greater than the base laminate; (2) buckling strain increases in magnitude as the sublaminate Poisson's ratio decreases and falls below the laminate value; and (3) the first mode of buckling changes from compression to tensile buckling for very low sublaminate Poisson's ratio.

Figure 7 shows the effect of fiber angle on the buckling strain for three unidirectional elliptical sublaminates with different aspect ratios. The length "a" was held constant (a = 25.4 mm); the widths b = a/2, a and 2a. Results for the circular sublaminate in figure 6 are repeated here. Buckling strains for b/a = 0.5 were compressive for  $\alpha < 60$  deg and became tensile for  $\alpha$  close to 90 deg. In contrast, buckling strains of b = 2a sublaminate remained compressive for all fiber angles. Buckling strain magnitude decreased monotonically with increase in sublaminate width, b.

<u>Oblique sublaminate</u>.- The effect of fiber angle on the buckling of oblique sublaminates is shown in figure 8. Buckling curves are shown for unidirectional sublaminate angles  $\theta = 0$ , 15, 30, 45, 60, and 90 degrees. The sublaminate size is a = 25.4, b = 50.8, and h = 0.51 mm. Figure 8 shows that for fiber angles 0 <  $\alpha$  < 38 deg, buckling strain decreases slightly with increasing sublaminate angle. The trend is reversed for  $\alpha > 38$  deg. For  $\theta = 0$  and 15 deg the buckling strain is compressive for all fiber angles; whereas for 30 deg <  $\theta$  < 90 deg the buckling strain changes to tensile for  $\alpha > 70$  deg.

Figure 9 shows the effect of sublaminate angle on buckling strain for  $[0]_4$ ,  $[90]_4$ ,  $[0/90]_8$ ,  $[\pm 45]_8$ , and  $[0/\pm 45/90]$  sublaminates. The figure shows that the buckling strains for  $[0/90]_8$ ,  $[\pm 45]_8$ , and  $[0/\pm 45/90]$  sublaminates are generally bounded by those of the  $[0]_4$  and  $[90]_4$  sublaminates. The buckling strain for a  $[0]_4$  sublaminate decreases with increasing sublaminate angle  $\theta$ ; for a  $[90]_4$  sublaminate the buckling strain increases. Furthermore, buckling strains of  $[0]_4$ ,  $[0/90]_8$ ,  $[\pm 45]_8$ , and  $[0/\pm 45/90]$  sublaminates are quite small (less than 0.0015) and compressive for all sublaminate angles.

<u>Criterion for delamination growth direction</u>.- For the configuration being studied, there are no interlaminar stresses around the delamination boundary until the delaminated region buckles. Whether the delamination grows after buckling depends on many parameters, including sublaminate size, sublaminate in-plane and flexural stiffness, and the interlaminar fracture toughness of the material. If we assume the delamination will grow, the question remains: "What direction will the delamination grow?" The possibility that the direction of initial growth for a circular delamination could be predicted based on changes in the buckling strain was investigated. The delamination was assumed to be initially circular (a = b = 25.4 mm). Then, it was assumed to grow in the load (x) direction by 50% (a = 38.1, and b = 25.4 mm) or transverse to the

load direction by 50% (a = 25.4, and b = 38.1 mm). Buckling strains for the three configurations were calculated for  $[0]_4$ ,  $[0/90]_5$ ,  $[\pm 45]_5$ , and  $[0/\pm 45/90]$  sublaminates. Figure 10 shows the calculated results and the experimentally observed growth directions [1,2,13]. The  $[0]_4$  sublaminate buckling strain is lowest when it is elongated in the load direction. In tests reported in reference 13 the delamination also grew in the load direction, accompanied with splitting. On the other hand the buckling strains for  $[0/90]_5$ ,  $[\pm 45]_5$ , and  $[0/\pm 45/90]$  sublaminates are lower when the delamination is elongated perpendicular to the load direction, which is also the experimentally observed growth direction [1,2]. Apparently, the direction of elongation of the sublaminate that has the lowest buckling strain will be the initial direction of growth of a circular delamination. More tests and analyses are needed to verify this interpretation.

# CONCLUSIONS

Finite-element (F-E) and Rayleigh-Ritz (R-R) analyses were used to predict the buckling strain of a localized delaminated region in a quasiisotropic laminate. The R-R analysis, developed using the Trefftz criterion, is simple and inexpensive compared to the F-E analysis. The accuracy of the R-R analysis was evaluated by comparing with the F-E analysis. Calculated buckling strains from the R-R analysis agreed with the F-E results except for highly anisotropic sublaminates; the neglect of the bending-twisting ( $D_{16}$  and  $D_{26}$ ) and the shear-extension ( $A_{16}$  and  $A_{26}$ ) coupling terms caused significant errors for these sublaminates. However, the region where the R-R method is invalid was small. A parametric study was performed for a four ply sublaminate to assess the influence of material anisotropy, and sublaminate angle shape and layup on buckling strains. The laminate was assumed to be made up of graphite/epoxy composite. This study led to the following conclusions:

- (1) Even though the base laminate is in a uniaxial stress state, the sublaminate is generally in a biaxial stress state either due to a mismatch of Poisson's ratio or to material anisotropy.
- (2) The biaxial stress states in some unidirectional composite sublaminates cause tensile buckling when the sublaminate is elongated in the load direction and fibers are nearly perpendicular to load direction.
- (3) Compressive buckling strain for unidirectional sublaminates increases with the increase in the angle between the load and the fiber directions.
- (4) The buckling strains for  $[0/90]_s$ ,  $[\pm 45]_s$ , and  $[0/\pm 45/90]$  sublaminates are generally bounded by those for  $[0]_4$  and  $[90]_4$  sublaminates.
- (5) The direction of elongation of the sublaminate that has the lowest buckling strain correlates with the initial growth direction of a circular delamination.

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Material	Modu	Poisson ratio		
	EŁ	Et	G <sub>lt</sub>	v <sub>lt</sub>
Aluminum	68.95 (10.0)	68.95 (10.0)	26.32 (3.82)	0.31
Graphite/epoxy	131.0 (19.0)	13.0 (1.89)	6.4 (0.93)	0.34

Table 1.- Material Elastic Properties



Figure 1.- Shapes of delaminations.



Figure 2.- Nomenclature for a buckled sublaminate.



Figure 3.- Finite-element idealizations.



Figure 4.- Aluminum sublaminate: effect of sublaminate width on buckling strain. (a = 25.4 mm, h = 0.51 mm)



Figure 5.- Specially orthotropic sublaminate: effect of sublaminate width (b) on buckling strain. (a = 25.4 mm, h = 0.51 mm, graphite/epoxy)

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Figure 6.- Circular sublaminate: effect of fiber angle on buckling strain. (a = 25.4 mm, h = 0.51 mm, graphite/epoxy)



Figure 7.- Elliptical sublaminate: effect of fiber angle on buckling strain. (a = 25.4 mm, h = 0.51 mm, graphite/epoxy)

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Figure 9.- Oblique sublaminates: effect of sublaminate angle on buckling strain. (a = 25.4 mm, b = 50.8 mm, h = 0.51 mm, graphite/epoxy)

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Figure 10.- Buckling strains for three sublaminate configurations. (h = 0.51 mm, graphite/epoxy)

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