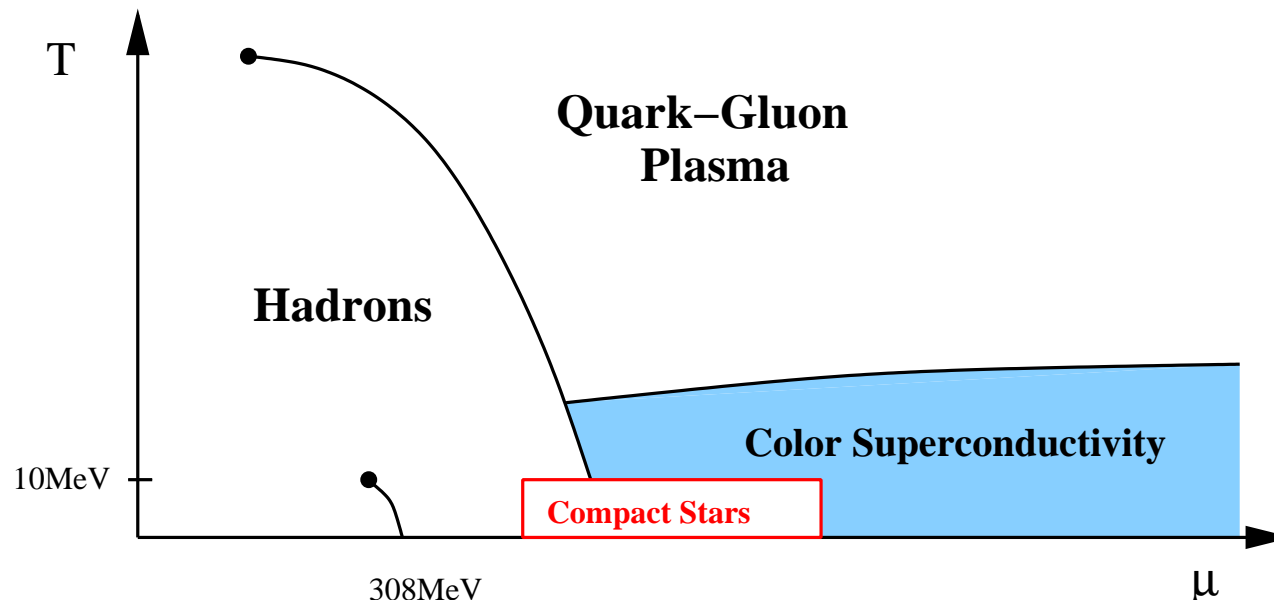


## Bulk viscosity in 2SC quark matter

Mark Alford, Andreas Schmitt, J. Phys. G 34, 67-101 (2007)

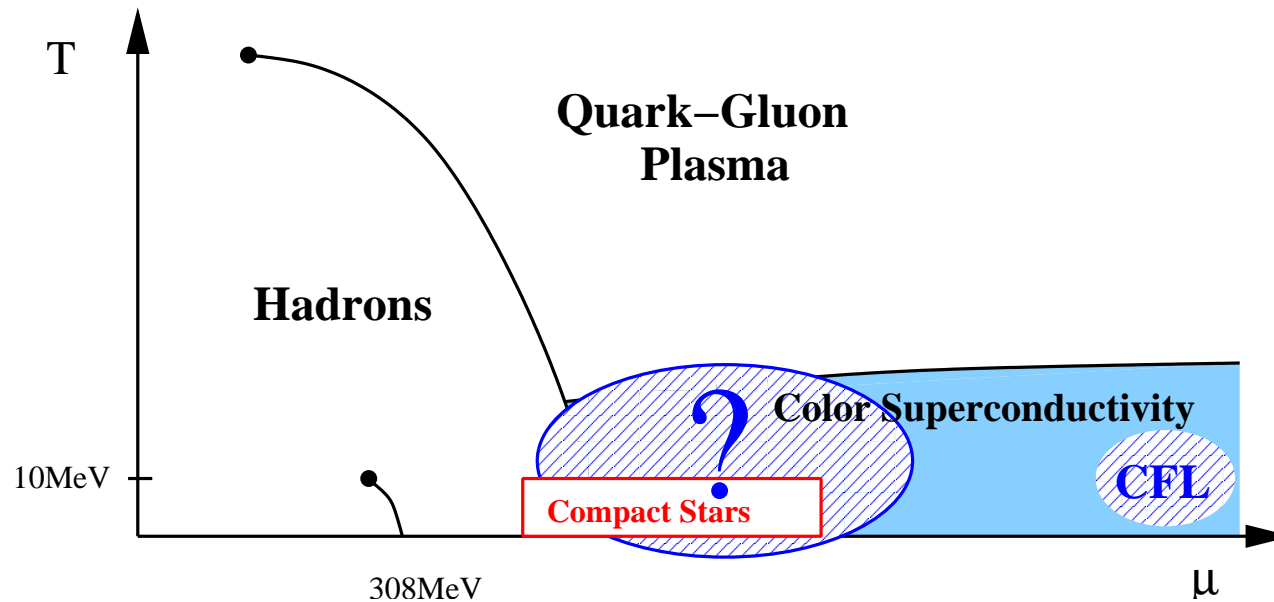
- **Color superconductivity: recent theoretical and astrophysical developments**
- **What is bulk viscosity?**  
**Why is it important for compact stars?**
- **Bulk viscosity in the 2SC phase**

- **QCD phase diagram (1):**  
**What is color superconductivity?**



	Where?	What?	Attractive force	Cooper pairs	Broken gauge group
“usual” superconductor	metals, alloys	ion lattice & electrons	phonons	electrons	$U(1)_{\text{em}}$
<b>color superconductor</b>	<b>neutron stars</b>	<b>quarks &amp; gluons</b>	<b>gluons</b>	<b>quarks</b>	$SU(3)_c$

- **QCD phase diagram (2): Unknown territory**



## Problems at moderate densities:

- perturbative QCD not valid
- strange mass not negligible  
→ neutrality requirements become nontrivial

- **Different approaches**

**Question:**

What is the ground state of deconfined quark matter at moderate densities (in the interior of compact stars)?

1. **Theoretical approach: start from CFL and ask “what is next phase down in density?” (if not hadronic matter)**
2. **Phenomenological approach: “guess” possible phase, compute its properties and compare with astrophysical observations**
3. (Tabletop approach: learn from parallels to cold fermionic atoms in magnetic trap)

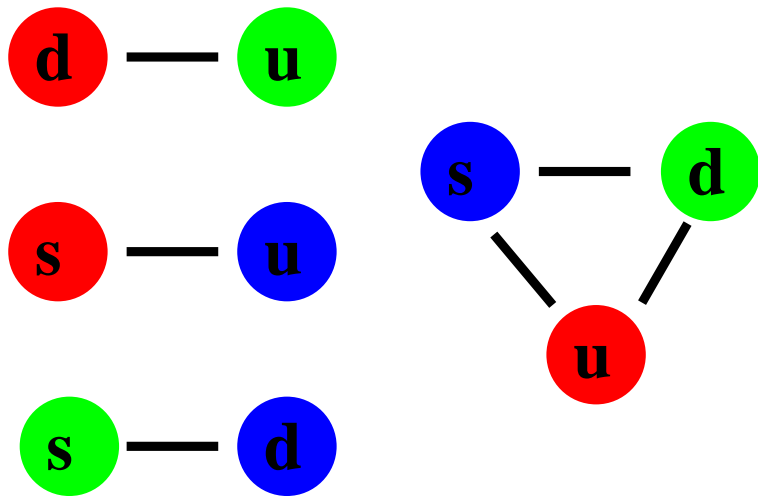
- **On safe grounds: Asymptotically large density**

$$0 \simeq m_s \simeq m_u \simeq m_d \ll \mu \quad \text{all quark masses negligible}$$

“color-flavor locked phase (CFL)”

M. Alford, K. Rajagopal, F. Wilczek, Nucl. Phys. B537, 443 (1999)

$$SU(3)_c \times SU(3)_L \times SU(3)_R \rightarrow SU(3)_{c+L+R}$$

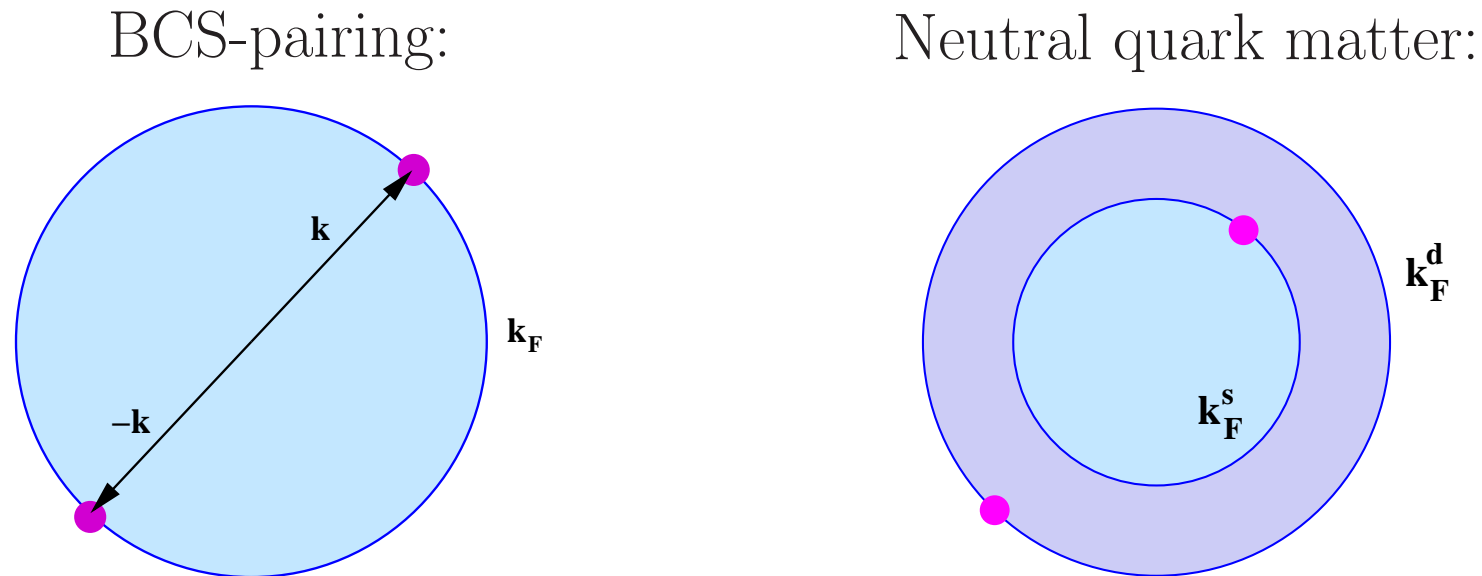


- all quarks form Cooper pairs
- state is automatically color and electrically neutral

- **Large, but not asymptotically large densities**

going down in density  $\Leftrightarrow$  “switching on”  $m_s$  and maintaining neutrality

$\rightarrow$  **mismatch in Fermi momenta of pairing quarks**  
 (“stressed” pairing)

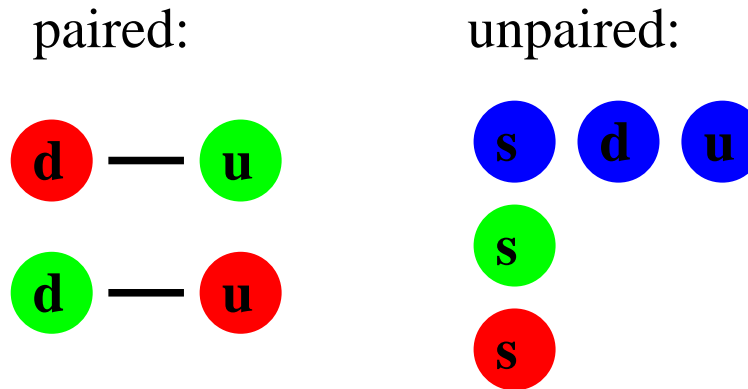


$\rightarrow$  stressed pairing is **unavoidable**

K. Rajagopal, A. Schmitt, PRD 73, 045003 (2006)

- **Less (and less symmetric) pairing**

For instance,  
2SC phase ...



... and many others

- **One-flavor pairing: Color-Spin-Locking, A-phase, ...**  
 T. Schäfer, PRD 62, 094007 (2000)  
 A. Schmitt, Q. Wang and D. H. Rischke, Phys. Rev. D **66**, 114010 (2002)
- **Gapless superconductors: g2SC, gCFL**  
 I. Shovkovy, M. Huang, PLB 564, 205 (2003)  
 M. Alford, C. Kouvaris, K. Rajagopal, PRL **92**, 222001 (2004)
- **Counter-propagating currents: LOFF, meson current**  
 M. Alford, J. Bowers, K. Rajagopal, PRD 63, 074016 (2001)  
 T. Schäfer, PRL 96, 012305 (2006)

- **Astrophysical approach (page 1/2)**
- **neutrino emissivity/cooling of the star**
  - **normal quark matter** N. Iwamoto, PRL 44, 1637 (1980)
  - **CFL** P. Jaikumar, M. Prakash, T. Schäfer, PRD 66, 063003 (2002)
  - **gCFL** M. Alford, P. Jotwani, C. Kouvaris, J. Kundu, K. Rajagopal, PRD, 114011 (2005)
  - **2SC** P. Jaikumar, C.D. Roberts, A. Sedrakian, PRC 73, 042801 (2006)
  - **spin-1** A. Schmitt, I.A. Shovkovy, Q. Wang, PRD 73, 034012 (2006)
  - **LOFF** R. Anglani, G. Nardulli, M. Ruggieri, M. Mannarelli, hep-ph/0607341

### direct Urca processes

$$u + e \rightarrow d + \nu, \quad d \rightarrow u + e + \bar{\nu}$$

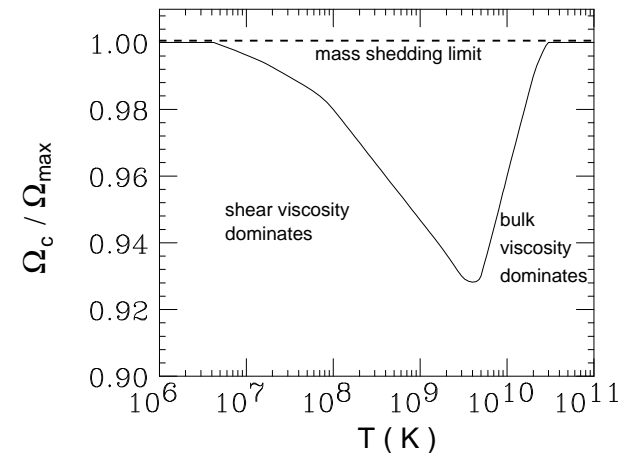
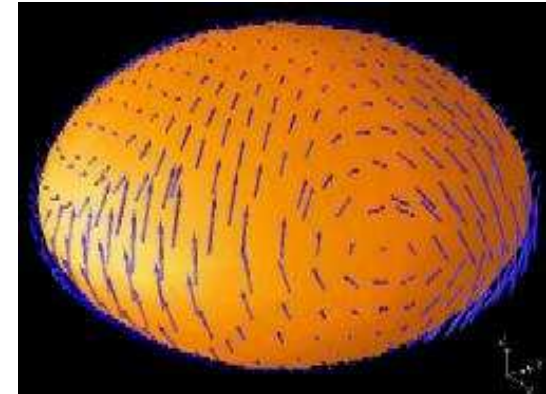
$\nu$ -emissivity sensitive to **magnitude of gap**  
and number and dimensionality of **ungapped modes**

**tendency:** need **fully gapped** phase with **not-too-large gap**  
to explain observed cooling ( $\rightarrow$  **CSL?**)



- **Astrophysical approach (page 2/2)**
- **magnetic fields**
  - **spin-0, e.g. CFL** “rotated electromagnetism”  
→ **no Meissner effect**  
M.G. Alford, J. Berges, K. Rajagopal, NPB 571, 269 (2000)
  - **spin-1 Meissner effect**  
A. Schmitt, Q. Wang, D.H. Rischke, PRL 91, 242301 (2003)
  - **precession of the star** B. Link, PRL 91, 101101 (2003)
  - **crust confinement of magnetic field**  
J.F. Perez-Azorin, J.A. Miralles, U.R.M. Geppert, A&A 451, 1009 (2006)
- **viscosity**
  - **shear, bulk viscosity in CFL**  
C. Manuel, A. Dobado, F.J. Llanes-Estrada, JHEP 0509, 076 (2005)  
M. Alford, M. Braby, S. Reddy, T. Schäfer, in preparation
  - **bulk viscosity in spin-1 phases**  
B.A. Sa’d, I.A. Shovkovy, D.H. Rischke, astro-ph/0607643
  - **bulk viscosity in 2SC phase**  
M.G. Alford, A. Schmitt, J. Phys. G 34, 67-101 (2007)

- **Why compute bulk viscosity?**
  - **r-modes:** non-radial pulsation modes
  - **grow unstable**  
in a **perfect-fluid** rotating star  
→ emission of gravitational waves
  - **spin down** the star drastically and quickly (within days)
- fast rotating stars are observed!  
 $\omega \simeq 1 \text{ms}^{-1}$
- must be some damping mechanism  
→ **bulk/shear viscosity**
- deduce upper limit for  $\omega$  from **viscosity**

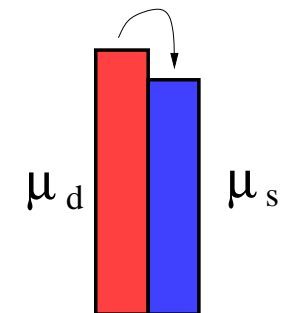
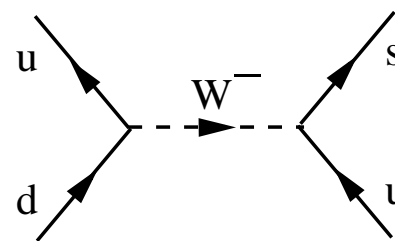
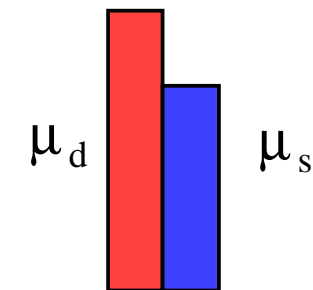
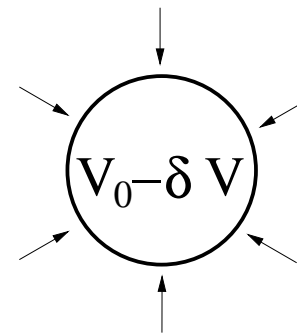
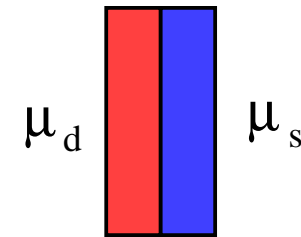
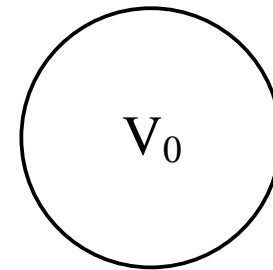
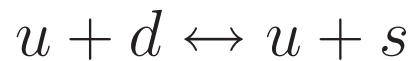


- **What is bulk viscosity?**

- compression/expansion changes chemical composition
- system out of chemical equilibrium

$$\delta\mu \equiv \mu_s - \mu_d \neq 0$$

re-equilibration via

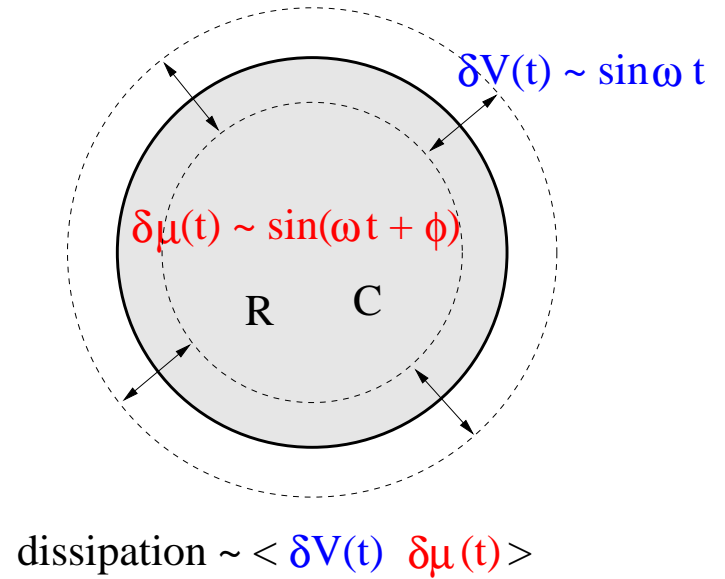
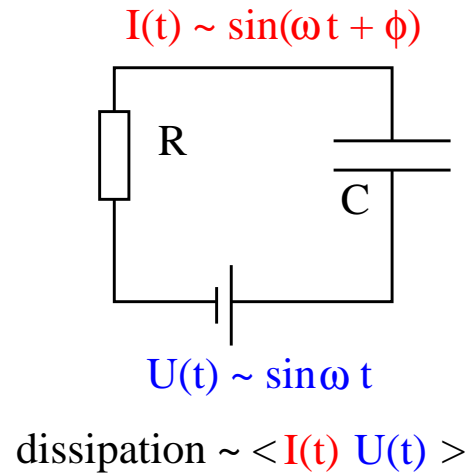


- **2 time scales:**

external oscillation  $\omega$  vs. microscopic rate  $\gamma$

• Bulk viscosity is a resonance phenomenon

Just like an electric circuit!



“capacitance”  $C \leftrightarrow$  inverse microscopic rate  $\gamma^{-1}$   
 (slow process  $\rightarrow$  store large chemical energy)

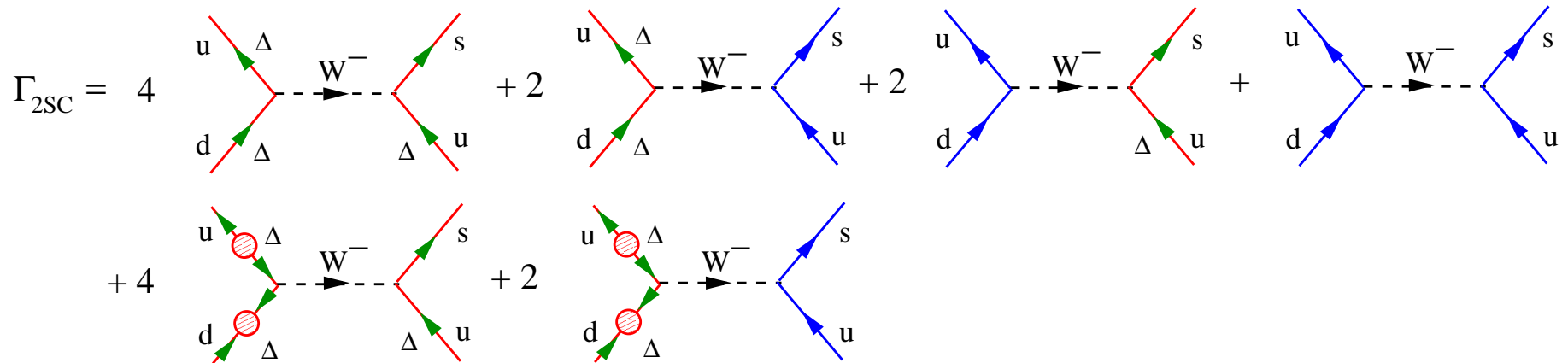
“resistance”  $R \leftrightarrow \left( n_u \frac{\partial \mu_d}{\partial n_u} + n_d \frac{\partial \mu_d}{\partial n_d} - n_s \frac{\partial \mu_s}{\partial n_s} \right)^{-1}$   
 (same dispersion for  $d$  and  $s \rightarrow$  infinite “resistance”  $\rightarrow$  no dissipation)

Bulk viscosity

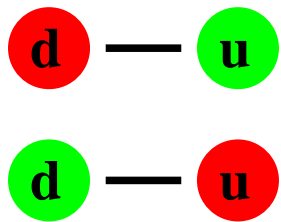
$$\zeta = \alpha \frac{\gamma}{\gamma^2 + \omega^2}$$

$$\alpha \equiv \frac{n_u \frac{\partial \mu_d}{\partial n_u} + n_d \frac{\partial \mu_d}{\partial n_d} - n_s \frac{\partial \mu_s}{\partial n_s}}{\frac{\partial \mu_d}{\partial n_d} + \frac{\partial \mu_s}{\partial n_s}}$$

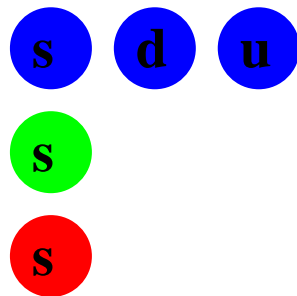
• Compute rate for  $u + d \leftrightarrow u + s$  in 2SC



paired:



unpaired:



small temperatures,

$$T \ll T_c \simeq 30\text{MeV}$$

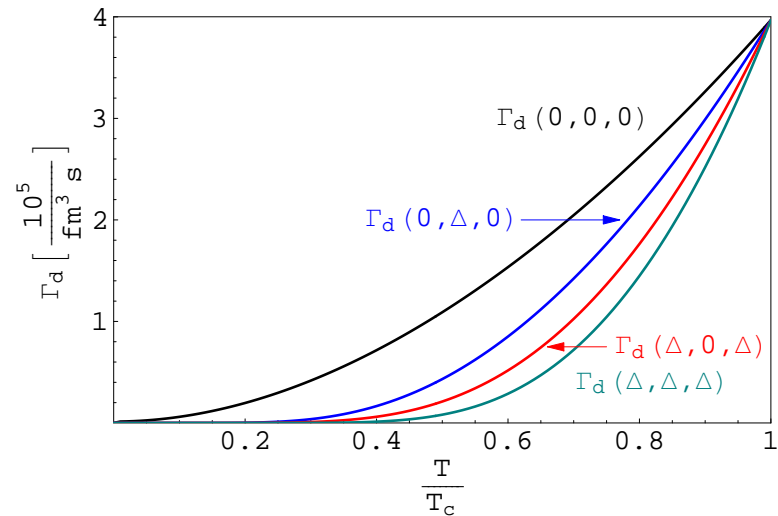
$$\Gamma_{2SC} = \frac{1}{9} \Gamma_{\text{unpaired}}$$

due to **exponential suppression**  
 $\exp(-\Delta/T)$  of gapped modes

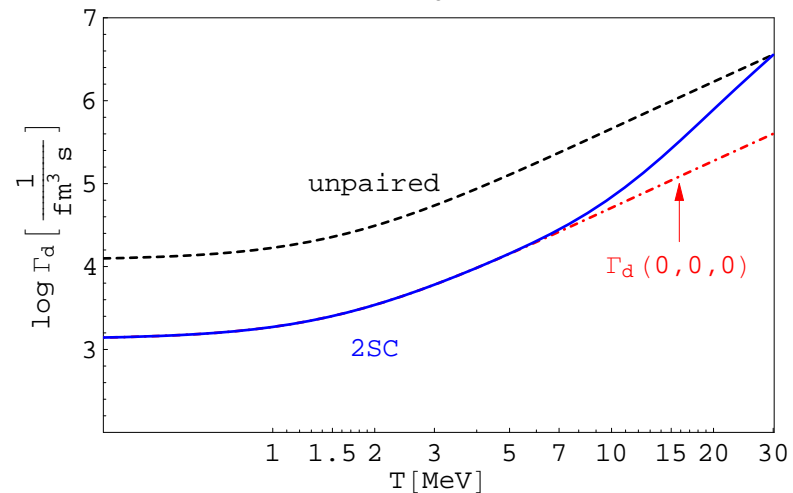
- **Results for all temperatures  $T < T_c$**

(i) fixed  $\delta\mu = \mu_s - \mu_d > 0 \rightarrow$  net production of  $d$  quarks,  $\Gamma_d > 0$

- **contributions of subprocesses**



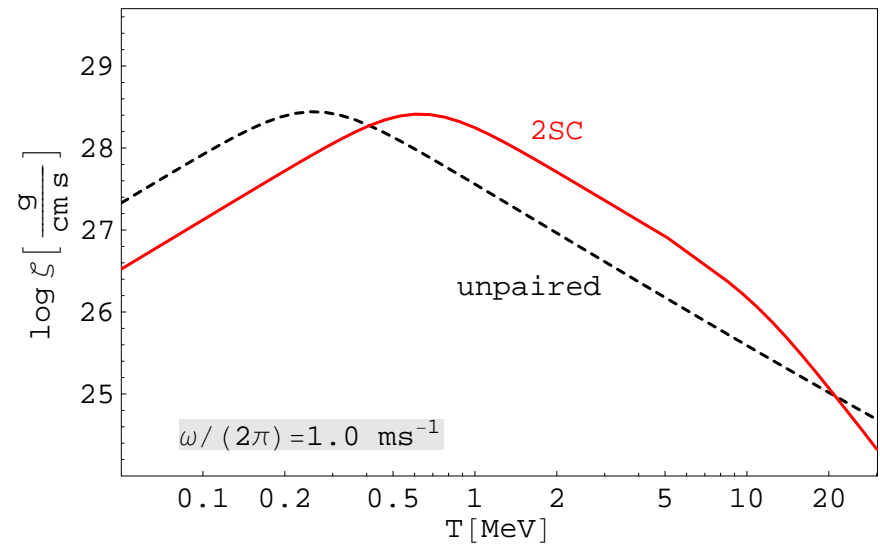
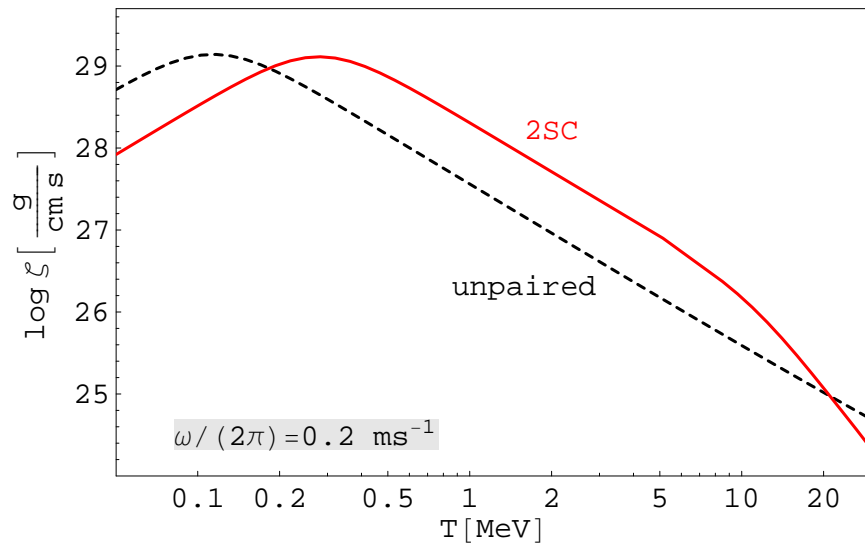
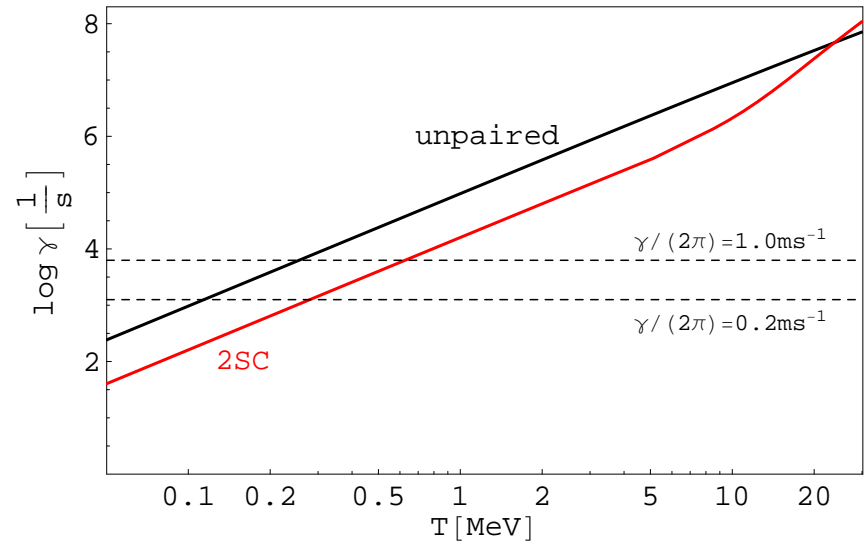
- **total rate compared to unpaired quark matter**



(ii) for bulk viscosity, consider  $\gamma = \frac{\partial \Gamma_d}{\partial \delta\mu}$

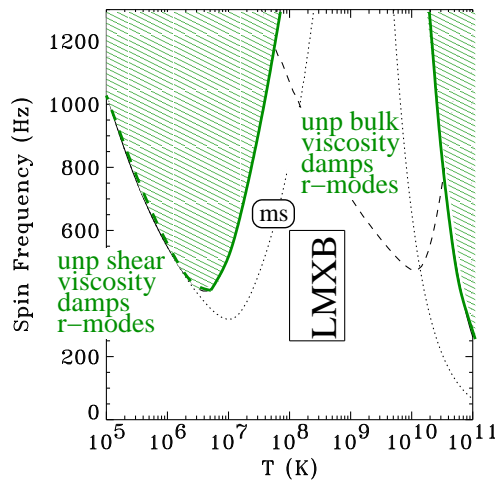
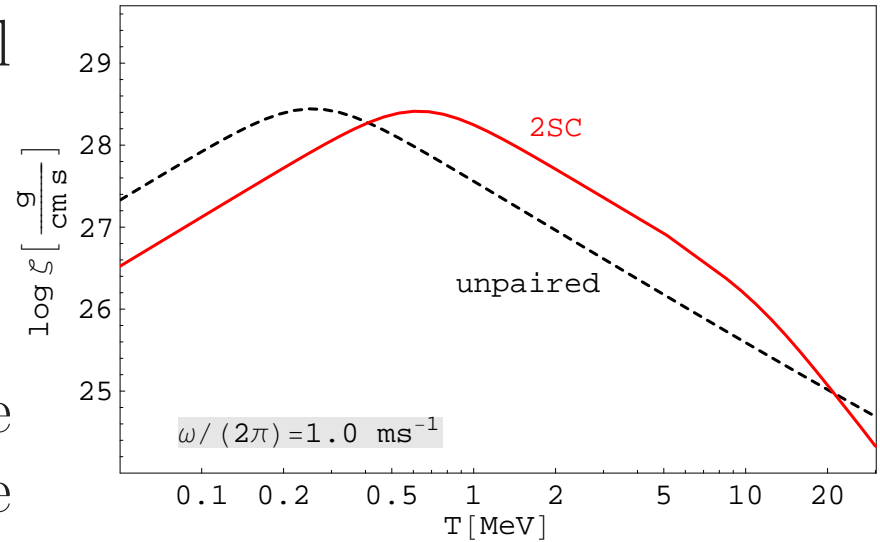
● Results for bulk viscosity

$$\zeta = \alpha \frac{\gamma}{\gamma^2 + \omega^2}$$

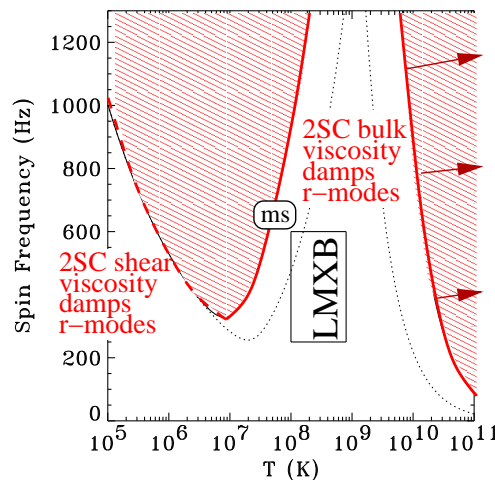


● **Astrophysical implications**

- bulk viscosity in superconductor can be **larger** than in normal phase
- results important for young neutron stars,  $T > 1 \text{ MeV}$
- → first days of neutron star's life (potentially enough for r-mode instabilities to grow)



unpaired



2SC

J. Madsen, PRL 85, 10 (2000)

- ms = millisecond pulsars
- LMXB = low-mass x-ray binaries



- **Conclusions**
- **Bulk viscosity**
  - shear and bulk viscosities damp **r-mode instabilities**
  - **bulk viscosity** of quark matter in a neutron star dominated by **weak processes** (unlike heavy-ion collisions; different **external time scale**)
  - **2SC quark matter** has **larger** bulk viscosity than unpaired quark matter in **very young neutron stars**
- **More general**
  - phase(s) between **CFL** and **hadronic matter** are unknown
  - use astrophysical observations to learn about these phases (**cooling curves, rotation frequencies, magnetic fields ...**)
  - how does quark matter deal with **mismatched Fermi surfaces? exotic superconductors?**