# $C^{*}$-EXTREME ENTANGLEMENT BREAKING MAPS ON OPERATOR SYSTEMS 

SRIRAM BALASUBRAMANIAN* AND NEHA HOTWANI ${ }^{1}$


#### Abstract

Let $\mathcal{E}$ denote the set of all unital entanglement breaking (UEB) linear maps defined on an operator system $\mathcal{S} \subset M_{d}$ and, mapping into $M_{n}$. As it turns out, the set $\mathcal{E}$ is not only convex in the classical sense but also in a quantum sense, namely it is $C^{*}$-convex. The main objective of this article is to describe the $C^{*}$-extreme points of this set $\mathcal{E}$. By observing that every EB map defined on the operator system $\mathcal{S}$ dilates to a positive map with commutative range and also extends to an EB map on $M_{d}$, We show that the $C^{*}$-extreme points of the set $\mathcal{E}$ are precisely the UEB maps that are maximal in the sense of Arveson ([A] and [A69]) and that they are also exactly the linear extreme points of the set $\mathcal{E}$ with commutative range. We also determine their explicit structure, thereby obtaining operator system generalizations of the analogous structure theorem and the Krein-Milman type theorem given in $[\mathrm{BDMS}]$. As a consequence, we show that $C^{*}$-extreme (UEB) maps in $\mathcal{E}$ extend to $C^{*}$-extreme UEB maps on the full algebra. Finally, we obtain an improved version of the main result in [BDMS], which contains various characterizations of $C^{*}$-extreme UEB maps between the algebras $M_{d}$ and $M_{n}$.


## 1. Introduction

The notions of positivity and convexity are fundamental to Mathematical analysis and in particular, to the theory of $C^{*}$-algebras. Among positive maps between $C^{*}$-algebras, the ones that are completely so, are of considerable interest. The study of completely positive maps was initiated by Stinespring and Arveson in the seminal papers $[\mathrm{S}]$,

[^0][A69] and [A72]. Among various results of significant importance in [A69], of particular interest to us is an abstract characterization of the (linear) extreme points of the convex set of completely positive maps between a $C^{*}$-algebra $\mathcal{A}$ and $B(H)$ for some Hilbert space $H$, in terms of the (minimal) Stinespring dilation. As important as classical convexity is, it still has some limitations in the non-commutative setting. Two "non-commutative" convexity notions that have gathered significant attention recently are matrix-convexity and $C^{*}$-convexity, the former introduced and studied by Webster and Winkler in [W] and [WW] and the latter by Hopenwasser, Loebl, Moore and Paulsen in [LP] and [HMP]. Although these notions appear to be similar, they are vastly different as was pointed out by Farenick in [F]. Farenick and Morenz also obtained a complete characterization of $C^{*}$-extreme unital completely positive maps between $C^{*}$-algebras in [FM93] and [FM97]. Further contributions on this and related topics can be found in [BBK], [BK].

Our main objects of focus in this article are entanglement breaking (EB) maps. These maps have drawn considerable attention recently and are particularly sought after in quantum information theory. In the finite dimensional setting, an important aspect of the set of unital entanglement breaking (UEB) maps is that it is $C^{*}$-convex. This warrants the study of $C^{*}$-extreme UEB maps. A complete description of such maps between matrix algebras was obtained in [BDMS]. The purpose of this article is four fold. Firstly, we obtain the explicit structure of UEB maps defined on a finite-dimensional operator system that are maximal with respect to the dilation order (see [A]) on the set of UEB maps (See Theorem 1.10). As a consequence we show that every UEB map on a finite dimensional operator system dilates to a maximal UEB map (See Theorem 1.11). Secondly, we show that UEB maps on a finite-dimensional operator system that are maximal with respect to the dilation order on the set of UEB maps are precisely the $C^{*}$-extreme UEB maps and that they are also exactly the linear extreme UEB maps with commutative range (See Theorem 1.12). As a consequence, we obtain an operator system generalization of a characterization of $C^{*}$-extreme UEB maps as well as the Krein-Milman type theorem for UEB maps given in [BDMS]. (See Corollary 1.14). It is to be noted that the partial order used for obtaining the abstract characterization of $C^{*}$-extreme UEB maps in [BDMS] is the usual partial order (and not the dilation order) on the set of UEB maps, i.e., for UEB maps $\Phi$ and $\Psi, \Phi \leq \Psi$ if and only if $\Psi-\Phi$ is a UEB map. Thirdly, we show that a $C^{*}$-extreme UEB map defined on a finite dimensional
operator system extends to a $C^{*}$-extreme UEB map on the full algebra (See Corollary 1.13). Finally, we obtain an improved version of the main result Theorem 5.3 in [BDMS] (See Theorem 6.2). To help us prove the above mentioned assertions, we make use of the following key observations namely, an EB map defined on a finite dimensional operator system dilates to a positive map with commutative range (See Theorem 1.7) and also has an EB extension to the full algebra (See Theorem 1.5).

Before we explicitly state our main observations and results, we introduce some notations and definitions. Throughout this article, $L$ and $K$ will denote separable Complex Hilbert spaces and $B(L)$ will denote the $\mathrm{C}^{*}$-algebra of bounded linear maps defined on $L$. An operator system $\mathcal{M} \subset B(L)$ is a self adjoint subspace containing the identity operator $I_{L}$. A linear map $\Phi: \mathcal{M} \rightarrow B(K)$ is said to be unital if $\Phi\left(I_{L}\right)=I_{K}$, positive if $\Phi(A)$ is a positive operator in $B(K)$ (in this case we write $\Phi(A) \succeq 0$ or $\Phi(A) \in B(K)^{+}$), whenever $A$ is a positive operator in $\mathcal{M}$. A unital positive linear functional on $\mathcal{M}$ is called a state. The notation $M_{k}$ will denote the space of all $k \times k$ Complex matrices and $M_{k}(\mathcal{M})$, the operator system of all $k \times k$ matrices with entries from $\mathcal{M}$. For a given linear map $\Phi: \mathcal{M} \rightarrow B(K)$ and $k \in \mathbb{N}$, the $\mathbf{k}^{\text {th }}$-ampliation $\Phi_{k}: M_{k}(\mathcal{M}) \rightarrow M_{k}(B(K))$ is defined by $\Phi_{k}\left(\left[A_{i, j}\right]\right)=\left[\Phi\left(A_{i, j}\right)\right]$. Under the identification of $M_{k}(\mathcal{M})$ with $M_{k} \otimes \mathcal{M}$, one sees that $\Phi_{k}:=\mathrm{i}_{k} \otimes \Phi$, where $\mathrm{i}_{k}$ is the identity operator on $M_{k}$. More specifically, $\Phi_{k}: M_{k} \otimes \mathcal{M} \rightarrow M_{k} \otimes B(K)$ is the linear map determined by $\Phi_{k}(X \otimes A)=X \otimes \Phi(A)$. The linear map $\Phi: \mathcal{M} \rightarrow B(K)$ is said to be completely positive (CP) if $\Phi_{k}$ is positive for all $k \in \mathbb{N}$.

Next we define the entanglement breaking property of linear maps. There are a number of competing definitions for the notion of an entanglement breaking map $\Phi: \mathcal{M} \rightarrow B(K)$ that agree with the usual notion in the case that $\mathcal{M}=B(L)$ and $K$ is a finite dimensional Hilbert space. In any case, they all reference the cone of (separable) matrices,

$$
M_{k}^{+} \otimes B(K)^{+}:=\left\{\sum_{m=1}^{\ell} A_{m} \otimes B_{m}: \ell \in \mathbb{N}, A_{m} \in M_{k}^{+}, B_{m} \in B(K)^{+}\right\}
$$

Since we are mainly interested in the case where $\mathcal{M} \subset B(L)$ is an operator system and, $K$ and $L$ are finite dimensional Hilbert spaces, inspired by the various notions of separability introduced in [CH], we will say that the linear map $\Phi: \mathcal{M} \rightarrow B(K)$ is entanglement breaking (EB) if $\left(\mathrm{i}_{k} \otimes \Phi\right)(X) \in \overline{M_{k}^{+} \otimes B(K)^{+}}$, for all $k \in \mathbb{N}$ and $X \in\left(M_{k} \otimes \mathcal{M}\right)^{+}$, where the closure is with respect to the norm topology on $B\left(\mathbb{C}^{k} \otimes K\right)$. Evidently, when $K$ is finite dimensional, $M_{k}^{+} \otimes B(K)^{+}$is
already norm closed. The abbreviations UCP and UEB will be used for "unital completely positive" and "unital entanglement breaking" respectively. The collection of all UEB maps mapping $\mathcal{M}$ to $B(K)$ will be denoted by $\operatorname{UEB}(\mathcal{M}, B(K))$. Finally, since we work mainly in the finite dimensional setting, we make the following notational conventions.

Convention 1.1. Throughout $E=\mathbb{C}^{d}$ and $H=\mathbb{C}^{n}$ and $\mathcal{S} \subset B(E)$ is an operator system.

Given a linear map $\Phi: \mathcal{S} \rightarrow B(H)$, the dual functional $s_{\Phi}$ : $B(H) \otimes \mathcal{S} \rightarrow \mathbb{C}$ associated to it, is the linear mapping determined by

$$
s_{\Phi}(X \otimes A)=\operatorname{trace}\left(\Phi(A) X^{t}\right)
$$

where $t$ is the transpose operator in $B(H)$ induced by a fixed orthonormal basis of $H$. It is easily seen that this definition of $s_{\Phi}$ is independent of the choice of the orthonormal basis and hence of the transpose operator induced by it and that the correspondence $\Phi \mapsto s_{\Phi}$ is bijective. Please see $[\mathrm{St}],[\mathrm{B}]$ for more details.

The Choi matrix $C_{\Phi}$ associated with the linear map $\Phi: B(E) \rightarrow$ $B(H)$ is defined as $\left[\Phi\left(e_{i} e_{j}^{*}\right)\right]_{i, j=1}^{d}=\sum_{i, j=1}^{d} e_{i} e_{j}^{*} \otimes \Phi\left(e_{i} e_{j}^{*}\right) \in B(E) \otimes B(H)$, where $\left\{e_{1}, \ldots, e_{d}\right\}$ is the standard orthonormal basis of $E$. One of the many significant applications of the Choi-matrix is the following wellknown characterization of CP maps due to Choi.

Theorem 1.2. ([C, Theorem 1], [St, Theorem 4.1.8]) Let $\Phi: B(E) \rightarrow$ $B(H)$ be a linear map. The following statements are equivalent.
(i) $\Phi$ is $C P$.
(ii) $\Phi(X)=\sum_{k=1}^{\ell} V_{k}^{*} X V_{k}$, for some linear maps $V_{k}: H \rightarrow E$.
(iii) $C_{\Phi} \in(B(E) \otimes B(H))^{+}$.

The formula for $\Phi$ given in statement (ii) above is called a ChoiKraus decomposition of $\Phi$ and the matrices coefficients $V_{k}$ are called Choi-Kraus operators/coefficients. The minimum number of ChoiKraus operators required to represent $\Phi$ in the form of a Choi-Kraus decomposition is known as the Choi-rank of $\Phi$. The Schmidt rank of the vector $\xi \in E \otimes H$, denoted by $S R(\xi)$, is the smallest natural number $k$ such that $\xi=\sum_{i=1}^{k} x_{i} \otimes y_{i} \in E \otimes H$. Given a completely positive map $\Phi: B(E) \rightarrow B(H)$, by Theorem 1.2 the Choi matrix $C_{\Phi} \in(B(E) \otimes B(H))^{+}$and hence has a spectral decomposition of the form $C_{\Phi}=\sum_{i=1}^{m} \xi_{i} \xi_{i}^{*}$, where $\xi_{i} \in E \otimes H$. Let $\Pi_{\Phi}$ denote the set of all spectral decompositions of $C_{\Phi}$. The Schmidt number of the Choi matrix $C_{\Phi} \in(B(E) \otimes B(H))^{+}$is denoted by $S N\left(C_{\Phi}\right)$ and is defined
as

$$
S N\left(C_{\Phi}\right):=\min _{\Gamma \in \Pi_{\Phi}}\left\{\max _{1 \leq j \leq m}\left\{S R\left(\xi_{j}\right) \mid \Gamma=\sum_{i=1}^{m} \xi_{i} \xi_{i}^{*} \in \Pi_{\Phi}\right\}\right\} .
$$

See [TH] for more details.
In this article, our main focus will be on EB maps. The following are a few well-known characterizations of EB maps that will be used in the sequel.
Theorem 1.3. ([HSR, Theorem 4], [TH]) Let $\Phi: B(E) \rightarrow B(H)$ be a linear map. The following statements are equivalent.
(i) $\Phi$ is $E B$.
(ii) $\Phi(X)=\sum_{j=1}^{m} \phi_{j}(X) R_{j}$, where the $\phi_{j}$ 's are states defined on $B(E)$ and the $R_{j}$ 's are positive operators in $B(H)$.
(iii) $\Phi(X)=\sum_{k=1}^{\ell} V_{k}^{*} X V_{k}$, where each $V_{k}: H \rightarrow E$ is a linear map of rank one.
(iv) $C_{\Phi} \in B(E)^{+} \otimes B(H)^{+}$. (i.e., $C_{\Phi}$ is separable.)
(v) $S N\left(C_{\Phi}\right)=1$.

The formula for $\Phi$ given in statement (ii) of the above theorem is called a Holevo form of $\Phi$. The EB-rank of $\Phi$ is defined as the minimum number of Choi-Kraus operators of rank one, required to represent $\Phi$ as in statement (iii) of Theorem 1.3. Please see [PPPR] for more details.
"Extremal" UEB maps are of utmost importance to us. Let $\mathcal{S} \subset$ $B(E)$ be an operator system. Here we mainly consider two notions of extreme points of $\operatorname{UEB}(\mathcal{S}, B(H))$. The UEB map $\Phi: \mathcal{S} \rightarrow B(H)$ is said to be a linear extreme point of $\operatorname{UEB}(\mathcal{S}, B(H))$ if

$$
\Phi=\sum_{i=1}^{k} t_{i} \Phi_{i}
$$

for some UEB maps $\Phi_{i}: \mathcal{S} \rightarrow B(H)$ and $t_{i} \in(0,1)$ satisfying $\sum_{i=1}^{k} t_{i}=$ 1 , then $\Phi_{i}=\Phi$ for all $i$. Linear extreme UCP maps betweeen $C^{*}$ algebras were characterized by Arveson in [A69].

The notion of $C^{*}$-convexity and $C^{*}$-extreme points were first introduced and studied in [LP] and [HMP] respectively. $C^{*}$-extreme UCP maps were extensively studied in [FM93], [FM97], [FZ] and [Z]. In [BDMS], the authors consider $C^{*}$-extreme UEB maps between matrix algebras. Along the same lines, we define $C^{*}$-extreme points in $\operatorname{UEB}(\mathcal{S}, B(H))$.

Let $\Phi_{1}, \Phi_{2}, \ldots, \Phi_{k} \in \operatorname{UEB}(\mathcal{S}, B(H))$. A UEB map $\Phi: \mathcal{S} \rightarrow B(H)$ is said to be a $C^{*}$-convex combination of the UEB maps $\Phi_{i}$, if there
exists $T_{1}, \ldots, T_{k} \in B(H)$ such that $\sum_{i=1}^{k} T_{i}^{*} T_{i}=I_{H}$ and $\Phi(X)=$ $\sum_{i=1}^{k} T_{i}^{*} \Phi_{i}(X) T_{i}$ for every $X \in \mathcal{S}$. The $T_{i}$ 's will be referred to as the coefficients of this $C^{*}$-convex combination. If the coefficients, i.e., the $T_{i}$ 's are positive, then this $C^{*}$-convex combination will be called a positive $C^{*}$-convex combination.

A UEB map $\Phi: \mathcal{S} \rightarrow B(H)$ is a proper $C^{*}$-convex combination of the UEB maps $\Phi_{i}, 1 \leq i \leq k$, if there exists invertible operators $T_{i} \in B(H)$ such that $\sum_{i=1}^{k} T_{i}^{*} T_{i}=I_{H}$ and

$$
\begin{equation*}
\Phi(X)=\sum_{i=1}^{k} T_{i}^{*} \Phi_{i}(X) T_{i} \tag{1.1}
\end{equation*}
$$

for all $X \in \mathcal{S}$. This $C^{*}$-convex combination is trivial if each $\Phi_{i}$ is unitarily equivalent to $\Phi$; that is, there exist unitary operators $U_{i}$ such that $\Phi_{i}(X)=U_{i}^{*} \Phi(X) U_{i}$. The UEB map $\Phi: \mathcal{S} \rightarrow B(H)$ is said to be a $C^{*}$-extreme point of $\operatorname{UEB}(\mathcal{S}, B(H))$ if, every representation of $\Phi$ as a proper $C^{*}$-convex combination is trivial.

Remark 1.4. Given our finite dimensionality assumptions, it is a standard observation that $C^{*}$-extreme UEB maps from $\mathcal{S} \subset B(E)$ mapping into $B(H)$ are also linear extreme UEB maps. Please see [Z, Theorem 2.2.2] for a proof of this fact for UCP maps. However, not every linear extreme UEB map is $C^{*}$-extreme, as an example in [HSR] shows. See also example 5.7 in [BDMS].

With the above given definitions and remarks, we proceed to state our main observations and results.
1.1. EB maps - Extension and Structure. A well-known extension theorem for CP maps due to Arveson says that every CP map defined on an operator system in a $C^{*}$-algebra mapping into $B(K)$ for some Hilbert space $K$, has a CP extension (See [P, Theorem 7.5]). Our first main observation in this article is the following analogous extension theorem for EB maps.

Theorem 1.5. Let $\mathcal{S} \subset B(E)$ be an operator system and $\phi: \mathcal{S} \rightarrow$ $B(H)$ be an EB map. There exists an EB map $\Phi: B(E) \rightarrow B(H)$ such that $\left.\Phi\right|_{\mathcal{S}}=\phi$.

A proof of Theorem 1.5 is given in Section 2. As an immediate consequence of the above theorem, one obtains the following operator system version of Theorem 1.3.

Corollary 1.6. The equivalence of statements (i), (ii) and (iii) of Theorem 1.3 holds even for UEB maps defined on operator systems. More
precisely, if $\mathcal{S} \subset B(E)$ is an operator system and $\Phi: \mathcal{S} \rightarrow B(H)$ is a linear map, then the following statements are equivalent.
(i) $\Phi$ is $E B$.
(ii) $\Phi$ can be written in the Holevo form, i.e., $\Phi(X)=\sum_{j=1}^{m} \phi_{j}(X) R_{j}$, where the $\phi_{j}$ 's are states defined on the operator system $\mathcal{S}$ and the $R_{j}$ 's are positive operators in $B(H)$.
(iii) $\Phi$ has a Choi-Kraus decomposition with rank-one Choi-Kraus coefficients, i.e., there exists linear maps $V_{k}: H \rightarrow E$ of rank one such that $\Phi(X)=\sum_{k=1}^{\ell} V_{k}^{*} X V_{k}$, for all $X \in \mathcal{S}$.
It is a well-known result due to Stinespring ([P, Theorem 4.1]) that a UCP map defined on a $\mathrm{C}^{*}$-algebra $\mathcal{A}$ mapping into $B(K)$ for some Hilbert space $K$, dilates to a representation. In this article we point out a similar structure theorem for EB maps (for our finite dimensional setting). The following characterization of an EB map in terms of a dilation, is our second main observation in this article.

Theorem 1.7. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow$ $B(H)$ a linear map. The following statements are equivalent.
(i) $\Phi$ is a (unital) $E B$ map.
(ii) $\Phi(X)=V^{*} \Gamma(X) V$ for all $X \in \mathcal{S}$, where $V: H \rightarrow K$ is an isometry for some finite dimensional Hilbert space $K$ and $\Gamma$ : $\mathcal{S} \rightarrow B(K)$ is a (unital) positive map with commutative range.
(iii) $\Phi$ is the compression of a (unital) $E B$ map with commutative range contained in $B(K)$ for some finite dimensional Hilbert space $K$.

A proof of Theorem 1.7 is given in Section 3. It relies mainly on the observation that an entanglement breaking map between matrix algebras factors via the commutative $\mathrm{C}^{*}$-algebra $\ell_{k}^{\infty}$ for some $k$. See [KMP] and [JKPP] for more details. We anticipate that the above structure theorem for "EB" maps being compressions of "positive" maps with commutative range should hold for a much more general setting than is considered here (for instance, for strongly entanglement breaking maps in the infinite dimensional setting (See [LD])).
Remark 1.8. The above theorem can also be deduced from the wellknown fact that a Positive Operator Valued Measure (POVM) dilates to a Projection Valued Measure (PVM) (see [Y]), which in turn can be deduced from Naimark's dilation Theorem ( $[\mathrm{P}$, Theorem 4.6]).
1.2. Maximal UEB maps. In [A] and [A72], Arveson defined the notion of a maximal UCP dilation of a given UCP map on an operator system. He also showed that maximal UCP dilations always exist and
that such maps are precisely the ones that satisfy a unique extension property. Following [A] and [A72], here we consider the UEB analog of maximal dilations (for our finite dimensional setting.) Let $\mathcal{S} \subset B(E)$ be an operator system and let $\Phi: \mathcal{S} \rightarrow B(H)$ be a UEB map. A linear $\operatorname{map} \Psi: \mathcal{S} \rightarrow B(K)$ is said to be a UEB dilation of $\Phi$ if $\Psi$ is a UEB map, $K$ is a separable Hilbert space, and there exists an isometry $V: H \rightarrow K$ such that

$$
\begin{equation*}
\Phi(X)=V^{*} \Psi(X) V \tag{1.2}
\end{equation*}
$$

for all $X \in \mathcal{S}$. In this case we write $\Phi \leq \Psi$. Note that if $H$ is identified with $V H$, then " $\leq$ " is a partial order on the set of all UEB maps defined on $\mathcal{S}$. The UEB dilation $\Psi$ of $\Phi$ is said to be trivial if

$$
\begin{equation*}
\Psi(X) V(x)=V(\Phi(X)(x)) \tag{1.3}
\end{equation*}
$$

for all $X \in \mathcal{S}$ and $x \in H$. The UEB map $\Phi$ is said to be maximal if every UEB dilation of $\Phi$ is trivial.

Remark 1.9. The following observations are immediate from the above definition of maximality for UEB maps.
(i) Since $V V^{*}$ is the projection onto Range $(V)$, it follows that the UEB dilation $\Psi$ is trivial if and only if $V H$ is an invariant subspace for $\Psi(X)$ for all $X \in \mathcal{S}$. Also, since $\mathcal{S}$ is an operator system and $\Psi\left(X^{*}\right)=\Psi(X)^{*}$ for all $X \in \mathcal{S}$, it follows that the UEB dilation $\Psi$ of $\Phi$ is trivial if and only if $V H$ is a reducing subspace for $\Psi(X)$ for every $X \in \mathcal{S}$.
(ii) Since $H$ is finite dimensional and $V: H \rightarrow K$ is an isometry, there is no loss of generality in assuming $K$ to be finite dimensional. (See Lemma 4.4)

The following characterization of maximal UEB maps is one of our main results in this article.

Theorem 1.10. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow$ $B(H)$ be a UEB map. The following statements are equivalent.
(i) $\Phi$ is maximal.
(ii) $\Phi$ has the form

$$
\begin{equation*}
\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i} \tag{1.4}
\end{equation*}
$$

for all $X \in \mathcal{S}$, where the $\phi_{i}$ 's are distinct linear extremal states on $\mathcal{S}, k \leq n$ and the $P_{i}$ 's are mutually orthogonal projections in $B(H)$ such that $\sum_{i=1}^{k} P_{i}=I_{H}$.

A proof of Theorem 1.10 is given in Section 4. Using Theorems 1.7 and 1.10, we also observe the following.

Theorem 1.11. Let $\mathcal{S} \subset B(E)$ be an operator system. Every UEB map $\Phi: \mathcal{S} \rightarrow B(H)$ dilates to a maximal UEB map $\Psi: \mathcal{S} \rightarrow B(K)$ with $\operatorname{dim}(K)<\infty$.

A proof of Theorem 1.11 can be found in Section 4.
1.3. $\mathbf{C}^{*}$-extreme UEB maps. In [BDMS], the authors obtain various characterizations of $C^{*}$-extreme UEB maps between matrix algebras. A primary objective of ours is to characterize $C^{*}$-extreme UEB maps defined on finite dimensional operator systems. One of our main results in this article along these lines is the following.

Theorem 1.12. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow$ $B(H)$ a UEB map. Then the following statements are equivalent.
(i) $\Phi$ is a maximal UEB map.
(ii) $\Phi$ is a $C^{*}$-extreme point of $\operatorname{UEB}(\mathcal{S}, B(H))$.
(iii) $\Phi$ is a linear extreme point of $\operatorname{UEB}(\mathcal{S}, B(H))$ with commutative range.

A proof of Theorem 1.12 is given in Section 5. Combined with Theorem 1.10, the equivalence $(i) \Longleftrightarrow(i i)$ of Theorem 1.12 is an operator system generalization of the equivalence $(i) \Longleftrightarrow(v)$ of [BDMS, Theorem 5.3], which says that $\Phi \in \operatorname{UEB}(B(E), B(H))$ is $C^{*}$-extreme if and only if it has the form given in equation (1.4).

After establishing a characterization as given above, it is only natural to ask whether $C^{*}$-extreme UEB maps defined on operator systems in $B(E)$ extend to $C^{*}$-extreme UEB maps on the whole algebra. This is still only a partially answered question for $C^{*}$-extreme UCP maps. Please see [Z] for more details and some results on this problem. Here, as an application of Theorems 1.10 and 1.12, we obtain the following extension result.

Corollary 1.13. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow$ $B(H)$ be a $C^{*}$-extreme UEB map. There exists a $C^{*}$-extreme UEB map $\Psi: B(E) \rightarrow B(H)$ such that $\left.\Psi\right|_{\mathcal{S}}=\Phi$.

A proof of Corollary 1.13 is given in Section 5.
In [BDMS], a Krein-Milman type theorem was established for the compact convex set $\operatorname{UEB}(B(E), B(H))$. To be precise, it was shown
that the set $\operatorname{UEB}(B(E), B(H))$ equals the $C^{*}$-convex hull of its $C^{*}$ extreme points. As another application of our main results, we obtain the following operator system analog of this result.

Corollary 1.14. $\operatorname{UEB}(\mathcal{S}, B(H))$ equals the $C^{*}$-convex hull of its $C^{*}$ extreme points.

A proof of Corollary 1.14 is given in Section 5.

## 2. Extensions of EB maps

In this section we give a proof of Theorem 1.5. We first include some necessary definitions adapted to our finite dimensional setting. Please see [St, Chapter 5] for more details.

Let $M$ and $L$ be finite dimensional Hilbert spaces. Let $\mathcal{P}(M)$ denote the cone of all positive linear maps from $B(M)$ to $B(M)$. A closed convex cone $\mathcal{C} \subset \mathcal{P}(M)$ is said to be a mapping cone if for each nonzero $A \in B(M)^{+}$, there exists a $\varphi \in \mathcal{C}$ such that $\varphi(A) \neq 0$ and

$$
\phi \circ \sigma \circ \psi \in \mathcal{C}
$$

for all $\sigma \in \mathcal{C}$ and CP maps $\phi, \psi: B(M) \rightarrow B(M)$. The mapping cone $\mathcal{C}$ is said to be symmetric if $\phi \in \mathcal{C}$ implies both $\phi^{*}$ and $t \circ \phi \circ t$ are in $\mathcal{C}$, where $t$ is the transpose operator in $B(M)$ induced by a fixed orthonormal basis of $M$ and $\phi^{*}: B(M) \rightarrow B(M)$ is the adjoint of $\phi$ with respect to the Hilbert-Schmidt inner product on $B(M)$, i.e., $\phi^{*}$ is determined by

$$
\langle\phi(A), B\rangle=\left\langle A, \phi^{*}(B)\right\rangle
$$

for all $A, B \in B(M)$, where $\langle C, D\rangle:=\operatorname{trace}\left(D^{*} C\right)$.
Remark 2.1. Typical examples of mapping cones that are symmetric are $U C P(B(M))$ and $E B(B(M))$.

Suppose that $\mathcal{S} \subset B(L)$ is an operator system and $\mathcal{C} \subset \mathcal{P}(M)$ is a mapping cone. A linear map $\phi: \mathcal{S} \rightarrow B(M)$ is said to be $\mathcal{C}$-positive if the corresponding dual functional $s_{\phi}: B(M) \otimes \mathcal{S} \rightarrow \mathbb{C}$ takes positive values on the cone
$P(\mathcal{S}, \mathcal{C}):=\left\{Y \in B(M) \otimes \mathcal{S}: Y=Y^{*},(\alpha \otimes \mathrm{i})(Y) \succeq 0\right.$ for all $\alpha$ in $\left.\mathcal{C}\right\}$,
where i is the identity map on $\mathcal{S}$.
Lemma 2.2. Let $\mathcal{S} \subset B(E)$ be an operator system and $\phi: \mathcal{S} \rightarrow B(H)$ be an $E B$ map. If $\mathcal{C}$ denotes the mapping cone $E B(B(H))$, then $\phi$ is $\mathcal{C}$-positive.

Proof. Since $\phi$ is an EB map, it follows from [St09, Proposition 1(ii)] that the dual functional $s_{\phi}$ takes positive values on the cone $K$, where $K:=\{X \in(B(H) \otimes \mathcal{S}):(\omega \otimes \mathrm{i})(X) \succeq 0 \quad$ for all states $\omega$ on $B(H)\}$.
Thus it suffices to show that $P(\mathcal{S}, \mathcal{C}) \subseteq K$. To prove this statement, let $\omega$ be a state defined on $B(H)$. Suppose also that $Y=\sum_{i} A_{i} \otimes B_{i} \in$ $P(\mathcal{S}, \mathcal{C})$ and $Z \in B(H)^{+}$is a fixed matrix of rank one. Define the linear map $\gamma: B(H) \rightarrow B(H)$ by $\gamma(X)=\omega(X) Z$. Observe that $\gamma$ has the Holevo form (See Theorem 1.3) and hence $\gamma \in \mathcal{C}$. Since $(\alpha \otimes i)(Y) \succeq 0$ for all $\alpha \in \mathcal{C}$, it follows that

$$
(\gamma \otimes \mathrm{i})(Y)=\sum_{i} \omega\left(A_{i}\right) Z \otimes B_{i}=Z \otimes\left[\sum_{i} \omega\left(A_{i}\right) B_{i}\right] \succeq 0
$$

and therefore $\sum_{i} \omega\left(A_{i}\right) B_{i}=(\omega \otimes \mathrm{i})(Y) \succeq 0$. Thus $Y \in K$ and the proof is complete.

Lemma 2.3. Let $\mathcal{S} \subset B(E)$ be an operator system and $\phi: \mathcal{S} \rightarrow B(H)$ be an EB map. Given $\epsilon>0$, there exists an $E B$ map $\Phi: B(E) \rightarrow B(H)$ such that $\|\Phi(A)-\phi(A)\|<\epsilon\|A\|$ for all $A \in \mathcal{S}$.
Proof. Let $\mathcal{C}=\mathrm{EB}(B(H))$. Since $\phi: \mathcal{S} \rightarrow B(H)$ is an EB map, it follows from Lemma 2.2 that $\phi$ is $\mathcal{C}$-positive. By observing that $\mathcal{C}$ is a symmetric mapping cone, it follows from [St, Theorem 5.1.13] that there exists a sequence $\phi_{j}:=\sum_{i=1}^{r_{j}} \alpha_{i}^{(j)} \circ \psi_{i}^{(j)} \in \mathcal{C}$ with $\alpha_{i}^{(j)} \in \mathcal{C}$ and CP maps $\psi_{i}^{(j)}: \mathcal{S} \rightarrow B(H)$ such that $\phi_{j} \rightarrow \phi$ in norm, in $B(H)$. Note that norm topology on $B(H)$ coincides with the BW-topology due to the finite dimensionality of $B(H)$. By the Arveson extension theorem [P, Theorem 7.5], there exist CP maps $\Psi_{i}^{(j)}: B(E) \rightarrow B(H)$ such that $\left.\Psi_{i}^{(j)}\right|_{\mathcal{S}}=\psi_{i}^{(j)}$. Define $\Phi_{j}:=\sum_{i=1}^{r_{j}} \alpha_{i}^{(j)} \circ \Psi_{i}^{(j)}$. It follows that $\Phi_{j}: B(E) \rightarrow B(H)$ is an EB map for all $j$. Let $A \in \mathcal{S}$ be arbitrary. Since $\phi_{j} \rightarrow \phi$, it follows that there exists an $j_{0} \in \mathbb{N}$ such that $\left\|\phi_{j}(A)-\phi(A)\right\|<\epsilon\|A\|$ for all $j \geq j_{0}$. Thus for all $j \geq j_{0}$, one obtains

$$
\begin{aligned}
\left\|\Phi_{j}(A)-\phi(A)\right\| & =\left\|\sum_{i=1}^{r_{j}}\left(\alpha_{i}^{(j)} \circ \Psi_{i}^{(j)}\right)(A)-\phi(A)\right\| \\
& =\left\|\sum_{i=1}^{r_{j}}\left(\alpha_{i}^{(j)} \circ \psi_{i}^{(j)}\right)(A)-\phi(A)\right\| \\
& =\left\|\phi_{j}(A)-\phi(A)\right\| \\
& <\epsilon\|A\|
\end{aligned}
$$

Defining $\Phi:=\Phi_{j_{0}}$ completes the proof.

Remark 2.4. Recall that $E=\mathbb{C}^{d}$. Due to the convexity of the mapping cone $\mathcal{C}$ and the finite dimensionality of $B(E)$, Caratheodory's theorem [DD, Theorem 16.1.8] implies that the map $\phi_{j}$ in the above proof can be written as $\phi_{j}=\sum_{i=1}^{r} \alpha_{i}^{(j)} \circ \Psi_{i}^{(j)}$, where $r=(2 d)^{2}+1$ (is independent of $j$ ).

Proof of Theorem 1.5. Let $j \in \mathbb{N}$ be arbitrary. By Lemma 2.3, there exists an EB map $\Phi_{j}: B(E) \rightarrow B(H)$ such that $\left\|\Phi_{j}(A)-\phi(A)\right\|<$ $\frac{1}{j}\|A\|$ for all $A \in \mathcal{S}$. Note that $\left\|\Phi_{j}\right\|=\left\|\Phi_{j}\left(I_{E}\right)\right\|$ and

$$
\begin{aligned}
\left\|\Phi_{j}\left(I_{E}\right)\right\| & \leq\left\|\Phi_{j}\left(I_{E}\right)-\phi\left(I_{E}\right)\right\|+\left\|\phi\left(I_{E}\right)\right\| \\
& <\frac{1}{j}+\left\|\phi\left(I_{E}\right)\right\| .
\end{aligned}
$$

Thus the sequence $\left\{\Phi_{j}\right\}_{j \in \mathbb{N}}$ is bounded. Consider the set

$$
\mathcal{K}:=\left\{\Gamma: B(E) \rightarrow B(H): \Gamma \text { is EB and }\|\Gamma\| \leq 1+\left\|\phi\left(I_{E}\right)\right\|\right\}
$$

Observe that $\mathcal{K}$ is compact and $\left\{\Phi_{j}\right\}_{j \in \mathbb{N}}$ is a sequence in $\mathcal{K}$. By the compactness of $\mathcal{K}$, there exists a subsequence $\left\{\Phi_{j_{k}}\right\}_{k \in \mathbb{N}}$ of $\left\{\Phi_{j}\right\}_{j \in \mathbb{N}}$ and an EB map $\Phi \in \mathcal{K}$ such that $\Phi_{j_{k}} \rightarrow \Phi$ as $k \rightarrow \infty$. The proof is complete by observing that $\left.\Phi\right|_{\mathcal{S}}=\phi$.

Remark 2.5. It is to be noted that Theorem 1.5 does not necessarily follow from the extension theorem [St, Theorem 5.2.3]. It was remarked in [St18] that the conclusion of [St, Theorem 5.2.3] holds only for Real operator systems and that this is due to the fact that Krein's Extension Theorem ([St, Theorem A.3.1]) which is formulated only for Real spaces, was incorrectly applied to Complex spaces, while proving [St, Theorem 5.2.3].

## 3. Dilations of UEB maps

In this section, we prove Theorem 1.7. The proof relies on a crucial observation from [RJP, Lemma 3.1]. We begin with the following remark.

Remark 3.1. Let $\Psi: B(E) \rightarrow B(H)$ be a (not necessarily unital) positive map with commutative range. The argument in the proof of [RJP, Lemma 3.1] can easily be adapted to conclude that $\Psi$ is $E B$.

Lemma 3.2. Let $K$ be a finite dimensional Hilbert space and $\mathcal{S} \subset$ $B(E)$ be an operator system. If $\Gamma: \mathcal{S} \rightarrow B(K)$ is a positive map with commutative range, then $\Gamma$ is $E B$.

Proof. By [HJRW, Corollary 2], there exists a commutative $C^{*}$-algebra $\mathcal{B} \subset B(K)$ containing Range $(\Gamma)$ and a CP map $\Theta: B(E) \rightarrow \mathcal{B} \subset B(K)$ such that $\left.\Theta\right|_{\mathcal{S}}=\Gamma$. It follows from Remark 3.1 that $\Theta$ is EB . Since $\Gamma$ is the restriction of the EB map $\Theta$ to $\mathcal{S}$, it is also EB.
Lemma 3.3. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow B(H)$ be an EB map. There exists a commutative $C^{*}$-algebra $\mathcal{A}$, a finite dimensional Hilbert space $K$, an isometry $V: H \rightarrow K$, a unital *algebra homomorphism $\pi: \mathcal{A} \rightarrow B(K)$ and a positive map $\eta: \mathcal{S} \rightarrow \mathcal{A}$ such that $\Phi(X)=V^{*}(\pi \circ \eta)(X) V$, for all $X \in \mathcal{S}$.

Proof. By Theorem 1.5, there exists an EB map $\Psi: B(E) \rightarrow B(H)$ such that $\left.\Psi\right|_{\mathcal{S}}=\Phi$. It follows from Theorem 1.3 and Corollary 1.6, $\Phi$ can be written in Holevo form, i.e., $\Phi(X)=\sum_{j=1}^{m} \phi_{j}(X) R_{j}$ for all $X \in$ $\mathcal{S}$, where the $\phi_{j}$ 's are states on $\mathcal{S}$ and the $R_{j}$ 's are positive operators in $B(H)$. Without loss of generality, assume that $\|\Phi\| \leq 1$.

Case(i) - $\Phi$ is non-unital: Let $\mathcal{A}=\ell_{m+1}^{\infty}$. Define the linear maps $\gamma: \mathcal{A} \rightarrow B(H)$ and $\eta: \mathcal{S} \rightarrow \mathcal{A}$ by $\gamma\left(x_{1}, \ldots, x_{m+1}\right)=\sum_{j=1}^{m+1} x_{j} R_{j}$ and $\eta(X)=\left(\phi_{1}(X), \ldots, \phi_{m}(X), 0\right)$, where $R_{m+1}=I_{H}-\sum_{j=1}^{m} R_{j}=$ $I_{H}-\Phi\left(I_{E}\right) \in B(H)^{+}$. Note that $\gamma$ and $\eta$ are positive maps with $\gamma$ also being unital. In fact, by [P, Theorem 3.9], $\gamma$ is a UCP map since the domain of $\gamma$ is a commutative $\mathrm{C}^{*}$-algebra. By Stinespring's dilation theorem [ P , Theorem 4.1], there exists a finite dimensional Hilbert space $K$, an isometry $V: H \rightarrow K$ and a unital $*$-algebra homomorphism $\pi: \mathcal{A} \rightarrow B(K)$ such that $V^{*} \pi(\cdot) V=\gamma(\cdot)$.

Case(ii) - $\Phi$ is unital: Let $\mathcal{A}=\ell_{m}^{\infty}$. Define the linear maps $\gamma$ : $\mathcal{A} \rightarrow B(H)$ and $\eta: \mathcal{S} \rightarrow \mathcal{A}$ by $\gamma\left(x_{1}, \ldots, x_{m}\right)=\sum_{j=1}^{m} x_{j} R_{j}$ and $\eta(X)=$ $\left(\phi_{1}(X), \ldots, \phi_{m}(X)\right)$. Note that both $\gamma$ and $\eta$ are unital positive maps. Arguing as above, one obtains a finite dimensional Hilbert space $K$, an isometry $V: H \rightarrow K$ and a unital $*$-algebra homomorphism $\pi: \mathcal{A} \rightarrow$ $B(K)$ such that $V^{*} \pi(\cdot) V=\gamma(\cdot)$.

In both cases above, observe that $\Phi=\gamma \circ \eta$. It follows that $\Phi(X)=$ $V^{*}(\pi \circ \eta)(X) V$, for all $X \in \mathcal{S}$.

Proof of Theorem 1.7. To prove $(i) \Longrightarrow(i i)$, observe that by Lemma 3.3 , there exists a commutative $C^{*}$-algebra $\mathcal{A}$, a finite dimensional Hilbert space $K$, an isometry $V: H \rightarrow K$, a unital $*$-algebra homomorphism $\pi: \mathcal{A} \rightarrow B(K)$ and a positive map $\eta: \mathcal{S} \rightarrow \mathcal{A}$ such that $\Phi(X)=V^{*}(\pi \circ \eta)(X) V$ for all $X \in \mathcal{S}$. Define $\Gamma:=\pi \circ \eta$. Since $\pi$ is a $*$-algebra homomorphism and $\eta$ is a positive map with commutative range, it follows that $\Gamma$ is a positive map with commutative range that dilates $\Phi$. Finally, observe that if $\Phi$ is unital, then $\Gamma$ is too.

The implication $(i i) \Longrightarrow$ (iii) follows directly from Lemma 3.2.
To prove $($ iii $) \Longrightarrow(i)$, let $\Phi(X)=V^{*} \Gamma(X) V$ for all $X \in \mathcal{S}$, where $V: H \rightarrow K$ is an isometry for some finite dimensional Hilbert space $K$ and $\Gamma: \mathcal{S} \rightarrow B(K)$ is an EB map with commutative range. It follows from Theorem 1.3 and Corollary 1.6 that $\Gamma(X)=\sum_{j=1}^{\ell} \psi_{j}(X) R_{j}$ for some states $\psi_{j}$ on $\mathcal{S}$ and positive operators $R_{j}$ in $B(K)$. Observe that for each $X \in \mathcal{S}$,

$$
\Phi(X)=V^{*}\left(\sum_{j=1}^{\ell} \psi_{j}(X) R_{j}\right) V=\sum_{j=1}^{\ell} \psi_{j}(X)\left(V^{*} R_{j} V\right)
$$

which is again a map in the Holevo form and hence is EB. Finally, if $\Gamma$ is unital, then $\Phi$ is too.

## 4. Maximal UEB Dilations

In this section we prove Theorems 1.10 and 1.11. The proof of Theorem 1.10 makes use of the following Lemmas, which contain some key properties of maximal UEB maps defined on operator systems.

Lemma 4.1. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow B(H)$ be a UEB map. If $\Phi$ has commutative range, then

$$
\begin{equation*}
\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i} \tag{4.1}
\end{equation*}
$$

where the $\phi_{i}$ 's are distinct states on $\mathcal{S}$ and the $P_{i}$ 's are mutually orthogonal projections in $B(H)$ such that $\sum_{i=1}^{k} P_{i}=I_{H}$.

Proof. Let $\mathcal{A}=C^{*}($ Range $(\Phi)) \subset B(H)$. There exists a unital $*$-algebra isomorphism $\pi: \mathcal{A} \rightarrow \ell_{k}^{\infty}$ for some $k$. Consider $\pi \circ \Phi: \mathcal{S} \rightarrow \ell_{k}^{\infty}$. For each $X \in \mathcal{S}, \pi \circ \Phi(X)=\sum_{i=1}^{k} \lambda_{i, X} e_{i}$, where $\lambda_{i, X}$ 's are scalars (depending on $X$ ) and $\left\{e_{1}, \ldots, e_{k}\right\}$ is the standard basis of $\ell_{k}^{\infty}$. For $1 \leq i \leq k$, define $\phi_{i}: \mathcal{S} \rightarrow \mathbb{C}$ via $\phi_{i}(X)=\lambda_{i, X}$. Since $\pi \circ \Phi$ is a unital positive map, it follows that $\phi_{i}$ 's are states on $\mathcal{S}$. Indeed

$$
\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i}
$$

where $P_{i}=\pi^{-1}\left(\left\{e_{i}\right\}\right)$. Note that the $P_{i}$ 's are mutually orthogonal projections in $B(H)$ such that $\sum_{i=1}^{k} P_{i}=I_{H}$. Without loss of generality, one can assume that the $\phi_{i}$ 's are distinct states on $\mathcal{S}$ because if two $\phi_{i}$ 's are the same, we can sum the corresponding projections, i.e., the $P_{i}$ 's together, and obtain a single projection.

Lemma 4.2. Let $\mathcal{S} \subset B(E)$ be an operator system. If $\Phi: \mathcal{S} \rightarrow B(H)$ is a maximal UEB map, then Range $(\Phi)$ is commutative.

Proof. Since $\Phi$ is a UEB map, it follows from part (iii) of Theorem 1.7 that there exists a finite dimensional Hilbert space $K$, a UEB map $\Gamma: \mathcal{S} \rightarrow B(K)$ with commutative range and an isometry $V: H \rightarrow K$ such that $V^{*} \Gamma(X) V=\Phi(X)$, for all $X \in \mathcal{S}$. Since $\Phi$ is a maximal UEB map, $\Gamma$ is a trivial UEB dilation of $\Phi$. Thus it follows from Remark 1.9 that the subspace $V H$ is an invariant subspace for $\Gamma(X)$ for all $X \in \mathcal{S}$. Since $V V^{*}$ is the projection of $K$ onto Range $(V)$, it follows that

$$
\begin{aligned}
\Phi(X) \Phi(Y)(a) & =V^{*} \Gamma(X) V V^{*} \Gamma(Y) V(a) \\
& =V^{*} \Gamma(X) \Gamma(Y) V(a) \\
& =V^{*} \Gamma(Y) \Gamma(X) V(a) \\
& =V^{*} \Gamma(Y) V V^{*} \Gamma(X) V(a) \\
& =\Phi(Y) \Phi(X)(a), \quad \text { for all } a \in H .
\end{aligned}
$$

Lemma 4.3. Let $\mathcal{S} \subset B(E)$ be an operator system and let $\Phi: \mathcal{S} \rightarrow$ $B(H)$ be a UEB map defined by $\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i}$, where $\phi_{i}$ are states defined on $\mathcal{S}$ and $P_{i} \in B(H)$ are mutually orthogonal projections satisfying $\sum_{i=1}^{k} P_{i}=I_{H}$. If $\Phi$ is a maximal UEB map, then each $\phi_{i}$ is a linear extremal state.

Proof. Fix $i \in\{1,2, \ldots, k\}$. Suppose that $\sigma, \tau$ are states on $\mathcal{S}$ such that

$$
\phi_{i}=t \sigma+(1-t) \tau
$$

for some $t \in(0,1)$. It suffices to show that $\sigma=\tau$. Let $F=\operatorname{Range}\left(P_{i}\right) \subset$ $H$ and $G=F^{\perp} \subset H$. Let $L:=F \oplus F \oplus G$. Define $V: H \rightarrow L$ by

$$
V(x+y)=(\sqrt{t} x, \sqrt{(1-t)} x, y)
$$

for all $x \in F$ and $y \in G$. It is easily seen that $V$ is an isometry. Let $\Psi: \mathcal{S} \rightarrow B(L)$ be defined by

$$
\Psi(X)(x, y, z)=\left(\sigma(X) x, \tau(X) y, \sum_{j \neq i} \phi_{j}(X) P_{j}(z)\right)
$$

for all $X \in \mathcal{S}, x, y \in F$ and $z \in G$. Observe that $\Psi$ is unital. Define the coordinate projections $Q_{1}, Q_{2}, Q_{3}$ on $L$ by $Q_{1}(x, y, z)=$ $(x, 0,0), Q_{2}(x, y, z)=(0, y, 0)$ and $Q_{3}(x, y, z)=(0,0, z)$. One sees that $\Psi(X)=\sigma(X) Q_{1}+\tau(X) Q_{2}+\left(\sum_{j \neq i} \phi_{j}(X) P_{j}\right) Q_{3}$, which is in Holevo form for $\Psi$. By Corollary 1.6, it follows that $\Psi$ is a UEB map. Observe
that

$$
\begin{aligned}
\left\langle V^{*}(u, v, w),(x+y)\right\rangle & =\langle(u, v, w), V(x+y)\rangle \\
& =\langle(u, v, w),(\sqrt{t} x, \sqrt{(1-t)} x, y)\rangle \\
& =\sqrt{t}\langle u, x\rangle+\sqrt{(1-t)}\langle v, x\rangle+\langle w, y\rangle \\
& =\langle\sqrt{t} u+\sqrt{(1-t)} v, x\rangle+\langle w, y\rangle \\
& =\langle\sqrt{t} u+\sqrt{(1-t)} v+w, x+y\rangle,
\end{aligned}
$$

for all $u, v, x \in F$ and $w, y \in G$. Thus $V^{*}$ is defined by

$$
V^{*}(u, v, w):=\sqrt{t} u+\sqrt{(1-t)} v+w .
$$

For each $X \in \mathcal{S}, x \in F$ and $y \in G$, it follows that

$$
\begin{aligned}
V^{*} \Psi(X) V(x+y) & =V^{*} \Psi(X)(\sqrt{t} x, \sqrt{(1-t)} x, y) \\
& =V^{*}\left(\sqrt{t} \sigma(X) x, \sqrt{(1-t)} \tau(X) x, \sum_{j \neq i} \phi_{j}(X) P_{j}(y)\right) \\
& =t \sigma(X) x+(1-t) \tau(X) x+\sum_{j \neq i} \phi_{j}(X) P_{j}(y) \\
& =\phi_{i}(X) x+\sum_{j \neq i} \phi_{j}(X) P_{j}(y) \\
& =\sum_{j=1}^{k} \phi_{j}(X) P_{j}(x+y) \\
& =\Phi(X)(x+y) .
\end{aligned}
$$

Hence $\Psi$ is a UEB dilation of $\Phi$. By hypothesis, $\Psi(X) V H \subset V H$ for all $X \in \mathcal{S}$. Choosing $0 \neq x \in F$, it follows that

$$
\begin{aligned}
(\sqrt{t} \sigma(X) x, \sqrt{(1-t)} \tau(X) x, 0) & =\Psi(X) V(x) \\
& =V(z+w)=(\sqrt{t} z, \sqrt{(1-t)} z, w)
\end{aligned}
$$

for some $z \in F$ and $w \in G$. Hence $\sigma(X) x=\tau(X) x$ for each $X \in \mathcal{S}$. This in turn implies that $\sigma=\tau$ and the proof is complete.

Lemma 4.4. Let $\mathcal{S} \subset B(E)$ be an operator system and let $\Phi: \mathcal{S} \rightarrow$ $B(H)$ be a UEB map. If every UEB dilation $\Psi: \mathcal{S} \rightarrow B(L)$ of $\Phi$ is trivial whenever $L$ is finite dimensional, then $\Phi$ is a maximal UEB map.

Proof. Let $\Psi: \mathcal{S} \rightarrow B(K)$ be a UEB dilation of $\Phi$ where $K$ is an infinite dimensional separable Hilbert space. Let the isometry $V: H \rightarrow K$
be such that $V^{*} \Psi(X) V=\Phi(X)$ for all $X \in \mathcal{S}$. Let $\left\{e_{1}, \ldots, e_{n}\right\}$ denote an orthonormal basis for $V H$. Extend it to an orthonormal basis $\left\{e_{i}: i \in \mathbb{N}\right\}$ for the Hilbert space $K$. For each $m \geq n$, let $K_{m}:=$ $\operatorname{span}\left\{e_{1}, \ldots, e_{m}\right\}, W_{m}: K_{m} \rightarrow K$ denote the inclusion map and $P_{m}$ denote the orthogonal projection of $K$ onto $K_{m}$. Observe that $P_{m}=W_{m}^{*}$. Define $V_{m}=P_{m} \circ V$ and $\Psi_{m}: \mathcal{S} \rightarrow B\left(K_{m}\right)$ to be the compression of $\Psi$ to $K_{m}$, i.e., $\Psi_{m}(X)=\left.P_{m} \Psi(X)\right|_{K_{m}}=W_{m}^{*} \Psi(X) W_{m}$. Let $k \in \mathbb{N}$ and $Z \in\left(M_{k} \otimes \mathcal{S}\right)^{+}$. Since $\Psi$ is an EB map, $\left(\mathrm{i}_{k} \otimes \Psi\right)(Z) \in \overline{M_{k}^{+} \otimes B\left(K_{m}\right)^{+}}$. It follows that $\left(\mathrm{i}_{k} \otimes \Psi_{m}\right)(Z)=\left(\mathrm{i}_{k} \otimes W_{m}\right)^{*}\left(\left(\mathrm{i}_{k} \otimes \Psi\right)(Z)\right)\left(\mathrm{i}_{k} \otimes W_{m}\right) \in$ $\overline{M_{k}^{+} \otimes B\left(K_{m}\right)^{+}}=M_{k}^{+} \otimes B\left(K_{m}\right)^{+}$. Thus $\Psi_{m}: \mathcal{S} \rightarrow B\left(K_{m}\right)$ is a UEB map. Moreover, $V_{m}$ is an isometry and $V_{m}^{*} \Psi_{m}(X) V_{m}=\Phi(X)$ for all $X \in \mathcal{S}$, i.e., $\Psi_{m}$ is a UEB dilation of $\Phi$. Let $a \in H$ and $X \in \mathcal{S}$. Take $V_{m}=P_{m} \circ V$, where $V$ is defined as above. Since $K_{m}$ is finite dimensional, by hypothesis, it follows that $\Psi_{m}$ is a trivial UEB dilation of $\Phi$. By Remark 1.9, one gets that $\Psi_{m}(X) V_{m} H \subset V_{m} H=V H$. Finally, for $x \in H$,

$$
\begin{aligned}
\Psi(X)(V x) & =\lim _{m \rightarrow \infty} P_{m} \Psi(X) P_{m}(V x)=\lim _{m \rightarrow \infty} W_{m}^{*} \Psi(X) W_{m} P_{m}(V x) \\
& =\lim _{m \rightarrow \infty} \Psi_{m}(X) V_{m}(x) \in V H
\end{aligned}
$$

Remark 4.5. It is to be noted that the definition of an EB map adopted here in terms of separability in the norm-closure sense, is not the usual way it is defined in the literature. If one works with a weak*-continuous $C P$ map (or equivalently, a normal CP map), say $\Phi: B(L) \rightarrow B(M)$, then there is a way of defining an $E B$ map that is along the lines of the usual definition namely: $\Phi: B(L) \rightarrow B(M)$ is $E B$ if for every $k \in \mathbb{N}, \Psi: \mathfrak{T}(M) \rightarrow \mathfrak{T}(L)$ maps positive matrices in $M_{k} \otimes \mathfrak{T}(M)$ to separable matrices in $M_{k} \otimes \mathfrak{T}(L)$, where $\mathfrak{T}(L)$ denotes the space of traceclass operators on the Hilbert space L, and $\Psi$ is the unique map whose (Banach space) adjoint $\Psi^{*}$ equals $\Phi$. This certainly is a generalization of the usual definition of an EB map from the finite to the infinite dimensional setting. But since here we mainly work on unital maps defined only on operator systems, defining the "entanglement breaking" property using this type of duality becomes a challenge.

We would like to emphasize that the above remark is only relevant when either of the Hilbert spaces $L$ or $M$ is infinite dimensional and so does not impact the results here.
Proof of Theorem 1.10. (i) $\Longrightarrow$ (ii): Since $\Phi$ is a maximal UEB map, it follows from Lemma 4.2 that $\Phi$ has commutative range. By applying Lemma 4.1, one gets that $\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i}$, where the $\phi_{i}$ 's are distinct states defined on $\mathcal{S}$ and $P_{i} \in B(H)$ are mutually orthogonal
projections satisfying $\sum_{i=1}^{k} P_{i}=I_{H}$. That the $\phi_{i}$ 's are linear extremal states on $\mathcal{S}$ follows from an application of Lemma 4.3.

To prove the implication (ii) $\Longrightarrow$ (i), let $\Psi$ be an arbitrary UEB dilation of $\Phi$; that is, $\Psi: \mathcal{S} \rightarrow B(K)$ is a UEB map such that there exists an isometry $V: H \rightarrow K$ satisfying $\Phi(X)=V^{*} \Psi(X) V$, for all $X \in \mathcal{S}$. By Lemma 4.4, one can assume that $K$ is finite dimensional. It suffices to show that $\Psi(X) V x \in V H$ for all $x \in H$ and $X \in \mathcal{S}$. If $x=0$, then there is nothing to prove. Suppose that $x \neq 0$. We first consider the case $x \in \operatorname{Range}\left(P_{j}\right)$ for some $j$. For such an $x$, it follows that $\Phi(X) x=\phi_{j}(X) x$ and hence $V^{*} \Psi(X) V x=\phi_{j}(X) x$.

Let a Holevo form (see Theorem 1.3 and Corollary 1.6) for $\Psi$ be given by $\Psi(X)=\sum_{i=1}^{k} \psi_{i}(X) R_{i}$, where the $\psi_{i}$ 's are distinct states and the $R_{i}$ 's are positive matrices with $\sum_{i=1}^{k} R_{i}=I_{K}$. Observe that

$$
\begin{align*}
\phi_{j}(X)\|x\|^{2}=\left\langle\phi_{j}(X) x, x\right\rangle & =\langle\Phi(X) x, x\rangle=\left\langle\left(\sum_{i=1}^{k} V^{*} \psi_{i}(X) R_{i} V\right) x, x\right\rangle \\
= & \sum_{i=1}^{k} \psi_{i}(X)\left\langle R_{i} V x, V x\right\rangle \tag{4.2}
\end{align*}
$$

Thus, $\phi_{j}(X)=\sum_{i=1}^{k} \alpha_{i} \psi_{i}(X)$, where $\alpha_{i}=\frac{1}{\|x\|^{2}}\left\langle R_{i} V x, V x\right\rangle \geq 0$. Note that $\alpha_{i}$ 's are independent of $X$. Also, since $V$ is an isometry and $\sum_{i=1}^{k} R_{i}=I_{K}$, it follows that

$$
\begin{equation*}
\sum_{i=1}^{k} \alpha_{i}=\frac{1}{\|x\|^{2}} \sum_{i=1}^{k}\left\langle R_{i} V x, V x\right\rangle=\frac{1}{\|x\|^{2}}\left\langle\left(\sum_{i=1}^{k} R_{i}\right) V x, V x\right\rangle=1 \tag{4.3}
\end{equation*}
$$

Combining equations (4.2) and (4.3), it follows that $\phi_{j}$ is a convex combination of the $\psi_{i}$ 's. Since the $\psi_{i}$ 's are distinct and $\phi_{j}$ is a linear extremal state, it follows that there exists an index $i_{0}$ such that $\alpha_{i_{0}}=1$, $\phi_{j}=\psi_{i_{0}}$ and $\alpha_{i}=\frac{1}{\|x\|^{2}}\left\langle R_{i} V x, V x\right\rangle=0$ for all $i \neq i_{0}$. Since the $R_{i}$ 's are positive, it follows that $R_{i} V x=0$ for all $i \neq i_{0}$, implying $V x=\sum_{i=1}^{k} R_{i} V x=R_{i_{0}} V x$. Thus

$$
\Psi(X) V x=\psi_{i_{0}}(X) R_{i_{0}} V x=\psi_{i_{0}}(X) V x \in V H
$$

Now using the fact that the $P_{i}$ 's are mutually orthogonal projections satisfying $\sum_{i=1}^{k} P_{i}=I_{H}$ and $\oplus_{i=1}^{k}$ Range $\left(P_{i}\right)=H$, it follows that $\Psi(X) V x \in V(H)$ for all $x \in H$, and the proof is complete.

Our next task is to prove Theorem 1.11, before which we prove the following crucial observation concerning a Holevo form of a UEB map.

Lemma 4.6. Let $\mathcal{S} \subset B(E)$ be an operator system. If $\Phi: \mathcal{S} \rightarrow B(H)$ is a UEB map, then $\Phi(X)=\sum_{j=1}^{k} \phi_{j}(X) S_{j}$ for each $X \in \mathcal{S}$, where $\phi_{j}$ are distinct linear extremal states on $\mathcal{S}$ and $S_{j}$ are positive operators in $B(H)$ satisfying $\sum_{j=1}^{k} S_{j}=I_{H}$.
Proof. Let a Holevo form for $\Phi$ be given by

$$
\Phi(X)=\sum_{i=1}^{r} \phi_{i}(X) R_{i}
$$

where the $\phi_{i}$ 's are states on $\mathcal{S}$ and the $R_{i}$ 's are positive operators in $B(H)$ such that $\sum_{i=1}^{r} R_{i}=I_{H}$. Observe that the set of all states on $\mathcal{S}$ is a compact and convex subset of the dual of $\mathcal{S}$. A well known consequence of Caratheodory's theorem (See [DD, Theorem 16.1.8] and [DD, Corollary 16.1.9]) and the Krein-Milman theorem is that in a finite dimensional topological vector space, a compact convex set $\mathcal{C}$ equals the convex hull of its extreme points. Using this fact here, each $\phi_{i}$ can be written as a convex combination of linear extremal states, i.e.,

$$
\phi_{i}(X)=t_{i, 1} \varphi_{i, 1}(X)+\cdots+t_{i, \ell_{i}} \varphi_{i, \ell_{i}}(X)
$$

where $t_{i, j} \in[0,1]$ with $\sum_{j=1}^{\ell_{i}} t_{i, j}=1$ and the $\varphi_{i, j}$ 's are linear extremal states on $\mathcal{S}$. It follows that

$$
\Phi(X)=\sum_{i=1}^{r} \sum_{j=1}^{\ell_{i}} t_{i, j} \varphi_{i, j}(X) R_{i}=\sum_{i=1}^{r} \sum_{j=1}^{\ell_{i}} \varphi_{i, j}(X)\left(t_{i, j} R_{i}\right)
$$

for all $X \in \mathcal{S}$. Note that by combining suitable terms in the above sum, it can be rewritten in the desired form. This completes the proof.

Proof of Theorem 1.11. Using Theorem 1.5 and the Holevo form (see Theorem 1.3 and Corollary 1.6) for EB maps, we have

$$
\begin{equation*}
\Phi(X)=\sum_{i=1}^{r} \phi_{i}(X) R_{i} \tag{4.4}
\end{equation*}
$$

for all $X \in \mathcal{S}$, where the $\phi_{i}$ 's are distinct states on $\mathcal{S}$ and the $R_{i}$ 's are positive operators in $B(H)$ such that $\sum_{i=1}^{r} R_{i}=I_{H}$. By Lemma 4.6, there is no loss of generality in assuming that the $\phi_{i}$ 's in equation (4.4) are distinct linear extremal states on $\mathcal{S}$. By Lemma 3.3, there exists a finite dimensional Hilbert space $K$, an isometry $V: H \rightarrow K, r \in \mathbb{N}$, a unital $*$-algebra homomorphism $\pi: \ell_{r}^{\infty} \rightarrow B(K)$ and a positive map $\eta: \mathcal{S} \rightarrow \ell_{r}^{\infty}$ such that

$$
\Phi(X)=V^{*}(\pi \circ \eta)(X) V
$$

for all $X \in \mathcal{S}$. As observed in the proof of Lemma 3.3, the unitality of $\Phi$ implies the unitality of $\eta$. Note that $\pi\left(x_{1}, \ldots, x_{r}\right)=\sum_{i=1}^{r} x_{i} \pi\left(e_{i}\right)$, where $\left\{e_{1}, \ldots, e_{r}\right\}$ is the standard basis of $\ell_{r}^{\infty}$. Also $(\pi \circ \eta)(X)=$ $\pi\left(\sum_{i=1}^{r} \phi_{i}(X) e_{i}\right)=\sum_{i=1}^{r} \phi_{i}(X) \pi\left(e_{i}\right)$, for all $X \in \mathcal{S}$. Since $\pi$ is a unital $*$-algebra homomorphism, it follows that the $\pi\left(e_{i}\right)$ 's are mutually orthogonal projections and $\sum_{i=1}^{r} \pi\left(e_{i}\right)=I_{K}$. Define $\Psi:=\pi \circ \eta$. It follows from Theorem 1.10 that $\Psi$ is maximal. This completes the proof.

## 5. $C^{*}$-extreme UEB maps on Operator Systems

In this section we prove Theorem 1.12 and Corollary 1.13. The proofs use techniques from [FM93], [FM97], [FZ] and [Z]. We also need the following Lemmas, the first of which contains an equivalent definition of a $C^{*}$-extreme UEB map and the second one contains a description of $\operatorname{UEB}(\mathcal{S}, B(H))$ in terms of $C^{*}$-convex combinations.

Lemma 5.1. Let $\mathcal{S} \subset B(E)$ be an operator system and $\Phi: \mathcal{S} \rightarrow B(H)$ be a UEB map. The following statements are equivalent.
(i) $\Phi$ is a $C^{*}$-extreme point of $\operatorname{UEB}(\mathcal{S}, B(H))$.
(ii) If $\Phi=\sum_{i=1}^{2} T_{i}^{*} \Phi_{i} T_{i}$, for some invertible operators $T_{1}, T_{2} \in$ $B(H)$ satisfying $T_{1}^{*} T_{1}+T_{2}^{*} T_{2}=I_{H}$, then there exist unitaries $U_{1}, U_{2} \in B(H)$ such that $\Phi_{i}(X)=U_{i}^{*} \Phi(X) U_{i}, i=1,2$.
Lemma 5.1 is essentially the operator system version of [BDMS, Proposition 3.2]. The proof given there works equally well for the operator system setting too.

Lemma 5.2. Let $\mathcal{S} \subset B(E)$ be an operator system, $\mathcal{K}$ denote the set $\left\{\Psi: \mathcal{S} \rightarrow B(H): \Psi(X)=g(X) I_{H}, g\right.$ is a linear extremal state on $\left.\mathcal{S}\right\}$, $\mathcal{E}$ denote the $C^{*}$-convex hull of $\mathcal{K}$, i.e., the set of all $C^{*}$-convex combinations of elements of $\mathcal{K}$ and $\mathcal{E}_{+}$denote the set of all positive $C^{*}$-convex combinations of elements of $\mathcal{K}$. The following statements hold.
(i) $\mathcal{E}_{+}=\mathcal{E}=\operatorname{UEB}(\mathcal{S}, B(H))$.
(ii) If $\Gamma \in \mathcal{E}$ is given by $\Gamma(X)=\sum_{i=1}^{m} T_{i}^{*}\left(g_{i}(X) I_{H}\right) T_{i} \in \mathcal{E}$, then $\Gamma \in \mathcal{E}_{+}$and there exists $B_{i} \in B(H)^{+}$such that $\sum_{i=1}^{m} B_{i}^{2}=I_{H}$ and

$$
\Gamma(X)=\sum_{i=1}^{m} B_{i}\left(g_{i}(X) I_{H}\right) B_{i}
$$

for all $X \in \mathcal{S}$.
Proof. Evidently, $\mathcal{E}_{+} \subset \mathcal{E} \subset \operatorname{UEB}(\mathcal{S}, B(H))$. To complete the proof of part $(i)$, it suffices to show that $\operatorname{UEB}(\mathcal{S}, B(H)) \subset \mathcal{E}_{+}$. To this end, let
$\Phi \in \operatorname{UEB}(\mathcal{S}, B(H))$. By Lemma 4.6, it follows that for each $X \in \mathcal{S}$, $\Phi(X)=\sum_{j=1}^{k} \phi_{j}(X) S_{j}$, where $\phi_{j}$ are distinct linear extremal states on $\mathcal{S}$ and $S_{j}$ are positive operators in $B(H)$ satisfying $\sum_{j=1}^{k} S_{j}=I_{H}$. Rewrite

$$
\Phi(X)=\sum_{j=1}^{k} \phi_{j}(X) S_{j}=\sum_{j=1}^{k} \sqrt{S_{j}}\left(\phi_{j}(X) I_{H}\right) \sqrt{S_{j}} .
$$

Thus $\Phi \in \mathcal{E}_{+}$. To prove part (ii), write $T_{i}$ in its polar decomposition, i.e., $T_{i}=U_{i} B_{i}$, where $U_{i}$ is a unitary operator and $B_{i}=\sqrt{T_{i}^{*} T_{i}} \in$ $B(H)^{+}$. Note that $\sum_{i=1}^{m} B_{i}^{2}=\sum_{i=1}^{m} T_{i}^{*} T_{i}=I_{H}$. By part $(i)$, it follows that $\Gamma \in \mathcal{E}_{+}$. Also, for each $X \in \mathcal{S}$,

$$
\begin{aligned}
\Gamma(X) & =\sum_{i=1}^{m} T_{i}^{*}\left(g_{i}(X) I_{H}\right) T_{i}=\sum_{i=1}^{m} B_{i} U_{i}^{*}\left(g_{i}(X) I_{H}\right) U_{i} B_{i} \\
& =\sum_{i=1}^{m} B_{i}\left(g_{i}(X) I_{H}\right) B_{i} .
\end{aligned}
$$

Proof of Theorem 1.12. $(i) \Rightarrow(i i)$ : Using the alternate definition of a $C^{*}$-extreme UEB map given in Lemma 5.1, let $\Phi(X)=T_{1}^{*} \Phi_{1}(X) T_{1}+$ $T_{2}^{*} \Phi_{2}(X) T_{2}$ for all $X \in \mathcal{S}$, where $\Phi_{1}, \Phi_{2}: \mathcal{S} \rightarrow B(H)$ are UEB maps and $T_{1}, T_{2} \in B(H)$ are invertible operators such that $\sum_{i=1}^{2} T_{i}^{*} T_{i}=I_{H}$. It suffices to show that $\Phi_{1}$ and $\Phi_{2}$ are unitarily equivalent to $\Phi$.

Define an isometry $V: H \rightarrow H \oplus H$ via $V(x)=\left(T_{1}(x), T_{2}(x)\right)$ and a linear map $\Psi: \mathcal{S} \rightarrow B(H \oplus H)$ by

$$
\Psi(X)=\left[\begin{array}{cc}
\Phi_{1}(X) & 0 \\
0 & \Phi_{2}(X)
\end{array}\right]
$$

for all $X \in \mathcal{S}$. Observe that $\Psi$ being a direct sum of UEB maps, is a UEB map. In fact, it is a UEB dilation of $\Phi$ since $V^{*} \Psi(X) V=\Phi(X)$ for all $X \in \mathcal{S}$. By hypothesis, it follows that $\Psi$ is a trivial UEB dilation of $\Phi$, that is, $\Psi(X) V=V \Phi(X)$ for all $X \in \mathcal{S}$. Equivalently,

$$
\left[\begin{array}{cc}
\Phi_{1}(X) & 0 \\
0 & \Phi_{2}(X)
\end{array}\right]\left[\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right]=\left[\begin{array}{l}
T_{1} \\
T_{2}
\end{array}\right] \Phi(X)
$$

for all $X \in \mathcal{S}$. Consequently, we have

$$
\Phi_{1}(X)=T_{1} \Phi(X) T_{1}^{-1} \text { and } \Phi_{2}(X)=T_{2} \Phi(X) T_{2}^{-1}
$$

for all $X \in \mathcal{S}$.
Using the polar decomposition of the invertible operator $T_{i}^{*}, i=1,2$, there exists a unitary operator $W_{i} \in B(H)$ and a positive invertible
operator $P_{i} \in B(H)$ given by $\sqrt{T_{i} T_{i}^{*}}$ such that $T_{i}=P_{i} W_{i}, i=1,2$. Hence

$$
\begin{equation*}
\Phi(X)=W_{i}^{*} P_{i}^{-1} \Phi_{i}(X) P_{i} W_{i} \tag{5.1}
\end{equation*}
$$

for $i=1,2$. Since $\Psi$ is a trivial UEB dilation of $\Phi$, it follows that

$$
\begin{equation*}
\Psi(X) V V^{*}=V V^{*} \Psi(X) \tag{5.2}
\end{equation*}
$$

for all $X \in \mathcal{S}$, (See Remark 1.9). Applying the definition of $\Psi$ and $V V^{*}$ in Equation (5.2), one obtains

$$
\Phi_{i}(X) T_{i} T_{i}^{*}=T_{i} T_{i}^{*} \Phi_{i}(X)
$$

for $i=1,2$. Note that $T_{i} T_{i}^{*}$ is a positive operator and $\Phi_{i}(X)$ commutes with $T_{i} T_{i}^{*}$. Hence

$$
\begin{equation*}
\Phi_{i}(X) P_{i}=P_{i} \Phi_{i}(X) \tag{5.3}
\end{equation*}
$$

for each $X \in \mathcal{S}$ and $i=1,2$.
It follows from equations (5.1) and (5.3) that

$$
\Phi(X)=W_{i}^{*} P_{i}^{-1} P_{i} \Phi_{i}(X) W_{i}=W_{i}^{*} \Phi_{i}(X) W_{i}
$$

for each $X \in \mathcal{S}$ and $i=1,2$. Since $W_{1}$ and $W_{2}$ are unitaries, the proof is complete.

The following proof of the implication $(i i) \Rightarrow(i)$ is adapted from [FM93, Theorem 4.1]. Recall the notations $\mathcal{K}, \mathcal{E}$ and $\mathcal{E}_{+}$from Lemma 5.2. Let $\Phi: \mathcal{S} \rightarrow B(H)$ be a $C^{*}$-extreme UEB map. Using Lemma 5.2, write $\Phi(X)=\sum_{i=1}^{m} T_{i}^{*} \Psi_{i}(X) T_{i}$, where $\Psi_{i}(X)=g_{i}(X) I_{H}$ for all $X \in \mathcal{S}$, the $g_{i}$ 's are linear extremal states on $\mathcal{S}$, and $\sum_{i=1}^{m} T_{i}^{*} T_{i}=I_{H}$. By part (ii) of Lemma 5.2, there is no loss of generality in assuming that the $T_{i}$ 's are positive. Let $m$ be the least number of coefficients required to represent $\Phi$ as a positive $C^{*}$-convex combination, i.e., as a $C^{*}$-convex combination of elements of $\mathcal{K}$ with positive coefficients (or equivalently to represent $\Phi$ as an element of $\mathcal{E}_{+}$.) If $m=1$, then there is nothing to prove, due to Theorem 1.10. So assume $m \geq 2$. Note that there must exist an $i$ such that $\left\|T_{i}\right\|=1$. Otherwise $\left\|T_{i}\right\|<1$ for all $i$ and, by modifying Technique-A in [FM93] to our current setting, one can rewrite $\Phi$ as a proper $C^{*}$-convex combination of some $\Gamma_{j} \in \mathcal{E}$ with the additional property that each $\Gamma_{j}$ is a $C^{*}$-convex combination of less than $m$ of the $\Psi_{i}{ }^{\prime}$ 's. Since $\Phi$ is $C^{*}$ - extreme, each $\Gamma_{j}$ is unitarily equivalent to $\Phi$. An application of part (ii) of Lemma 5.2 implies that $\Phi$ is a positive $C^{*}$-convex combination of fewer than $m$ of the $\Psi_{i}$ 's, which contradicts the minimality of $m$. Thus at least one of the $T_{i}$ 's has unit norm. Without loss of generality, assume that $\left\|T_{1}\right\|=1$. Due to the unitality of $\Phi$, note also that $\sum_{i=1}^{m} T_{i}^{2}=I_{H}$. Since $T_{1} \geq 0$ and
$\left\|T_{1}\right\|=1$, it follows that there exists a unitary $U \in B(H)$ such that

$$
U^{*} T_{1} U=\left(\begin{array}{rr}
I & 0 \\
0 & Y_{1}
\end{array}\right)
$$

where $Y_{1}$ is a diagonal matrix satisfying $Y_{1} \geq 0,\left\|Y_{1}\right\|<1$ and $I$ is the identity operator of suitable size. Since $T_{i} \geq 0$ for all $i \geq 2$ and $\sum_{i=1}^{m} T_{i}^{2}=I_{H}$,

$$
U^{*} T_{i} U=\left(\begin{array}{cc}
0 & \\
& Y_{i}
\end{array}\right), \text { with } Y_{i} \geq 0 \text { for all } i \geq 2
$$

Let $W_{i}=U^{*} T_{i} U$ for each $i \in\{1,2, \ldots, m\}$. Observe that $W_{i} \geq 0$ and

$$
\begin{equation*}
U^{*} \Phi(X) U=\sum_{i=1}^{m} W_{i}\left(U^{*} \Psi_{i}(X) U\right) W_{i}=\sum_{i=1}^{m} W_{i} \Psi_{i}(X) W_{i} \tag{5.4}
\end{equation*}
$$

Since $\sum_{i=1}^{m} W_{i}^{2}=I_{H}$, it follows that $\sum_{i=1}^{m} Y_{i}^{2}=I$. Since $\left\|Y_{1}\right\|<1$, $I-Y_{1}^{2}$ invertible. Also since

$$
\sum_{i \geq 2}\left(\begin{array}{cc}
0 & 0 \\
0 & Y_{i}^{2}
\end{array}\right)=\sum_{i \geq 2} W_{i}^{2}=I_{H}-W_{1}^{2}=\left(\begin{array}{cc}
0 & 0 \\
0 & I-Y_{1}^{2}
\end{array}\right)
$$

it follows that $\sum_{i \geq 2} Y_{i}^{2}$ is invertible. Adapting Technique-C from [FM93] to our setting and applying it here allows us to write $\sum_{i \geq 2} W_{i} \Psi_{i}(X) W_{i}$ as a single term $T_{0}^{*} \Gamma_{0}(X) T_{0}$ for some $\Gamma_{0} \in \mathcal{E}$. Note that $T_{0}$ need not be positive. By using the polar decomposition $T_{0}=U_{0} W_{0}$ where $U_{0}$ is unitary and $W_{0}$ is a positive operator, observe that $T_{0}^{*} \Gamma_{0}(X) T_{0}=$ $W_{0} \Psi_{0}(X) W_{0}$, for some $\Psi_{0} \in \mathcal{E}$. Thus

$$
\begin{equation*}
\sum_{i \geq 2} W_{i} \Psi_{i}(X) W_{i}=W_{0} \Psi_{0}(X) W_{0} \tag{5.5}
\end{equation*}
$$

for some positive operator $W_{0}$ and $\Psi_{0} \in \mathcal{E}$.
Indeed $W_{0}^{2}+W_{1}^{2}=I_{H}$ and $W_{0}=\left(\begin{array}{cc}0 & 0 \\ 0 & Y_{0}\end{array}\right)$ for some positive operator $Y_{0}$. Recall the positive matrix $Y_{1}$ and rewrite it as $Y_{1}=\left(\begin{array}{cc}Z_{1} & 0 \\ 0 & 0\end{array}\right)$ where $Z_{1}$ is a positive invertible matrix with $\left\|Z_{1}\right\|<1$. Since $Y_{0}^{2}+Y_{1}^{2}=I$, $Y_{0}=\left(\begin{array}{cc}Z_{0} & 0 \\ 0 & I\end{array}\right)$ for some matrix $Z_{0}$. Thus

$$
W_{1}=\left(\begin{array}{ccc}
I & 0 & 0 \\
0 & Z_{1} & 0 \\
0 & 0 & 0
\end{array}\right) \quad \text { and } \quad W_{0}=\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & Z_{0} & 0 \\
0 & 0 & I
\end{array}\right)
$$

where $Z_{1}, Z_{0}$ are positive matrices. Since $W_{0}^{2}+W_{1}^{2}=I_{H}$, it must be the case that $Z_{0}^{2}+Z_{1}^{2}=I$, which in turn implies that $Z_{0}$ is invertible, since
$\left\|Z_{1}\right\|<1$. Recall that $Z_{1}$ is also invertible. It follows from equations (5.4) and (5.5) and the fact that $\Psi_{1} \in \mathcal{K}$ that

$$
\begin{align*}
& U^{*} \Phi(X) U=W_{0} \Psi_{0}(X) W_{0}+W_{1} \Psi_{1}(X) W_{1} \\
& =\left(\begin{array}{ccc}
I & 0 & 0 \\
0 & Z_{0} & 0 \\
0 & 0 & I
\end{array}\right) \Theta(X)\left(\begin{array}{ccc}
I & 0 & 0 \\
0 & Z_{0} & 0 \\
0 & 0 & I
\end{array}\right)+\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & Z_{1} & 0 \\
0 & 0 & 0
\end{array}\right) \Psi_{1}(X)\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & Z_{1} & 0 \\
0 & 0 & 0
\end{array}\right), \tag{5.6}
\end{align*}
$$

where

$$
\Theta(X)=\left(\begin{array}{lll}
I & 0 & 0  \tag{5.7}\\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right) \Psi_{1}(X)\left(\begin{array}{lll}
I & 0 & 0 \\
0 & 0 & 0 \\
0 & 0 & 0
\end{array}\right)+\left(\begin{array}{lll}
0 & 0 & 0 \\
0 & I & 0 \\
0 & 0 & I
\end{array}\right) \Psi_{0}(X)\left(\begin{array}{ccc}
0 & 0 & 0 \\
0 & I & 0 \\
0 & 0 & I
\end{array}\right)
$$

$\in \mathcal{E}$.
By the invertibility of $Z_{0}$, it is immediate that $Y_{0}=\left(\begin{array}{cc}Z_{0} & 0 \\ 0 & I\end{array}\right)$ is also invertible. Applying Technique-B from [FM93] allows for writing $U^{*} \Phi U$ as a proper $C^{*}$-convex combination of $\Theta$ and some other $\Gamma \in \mathcal{E}$. By hypothesis, $\Phi$ is a $C^{*}$-extreme UEB map. Hence so is $U^{*} \Phi U$. It follows that $\Phi$ is unitarily equivalent to $\Theta$. Let $Q_{1}$ and $Q_{1}^{\perp}$ denote the projections $\left(\begin{array}{lll}I & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0\end{array}\right)$ and $\left(\begin{array}{ccc}0 & 0 & 0 \\ 0 & I & 0 \\ 0 & 0 & I\end{array}\right)$ respectively. Indeed, with respect to the decomposition $H=\operatorname{Range}\left(Q_{1}\right) \oplus \operatorname{Range}\left(Q_{1}^{\perp}\right)$,

$$
\begin{equation*}
\Theta(X)=\left.\left.Q_{1} \Psi_{1}(X)\right|_{\text {Range }\left(Q_{1}\right)} \oplus Q_{1}^{\perp} \Psi_{0}(X)\right|_{\operatorname{Range}\left(Q_{1}^{\perp}\right)}, \tag{5.8}
\end{equation*}
$$

for all $X \in \mathcal{S}$. Since $\left.Q_{1}^{\perp} \Psi_{0}(X)\right|_{\text {Range }\left(Q_{1}^{\perp}\right)} \in \operatorname{UEB}\left(\mathcal{S}, B\left(\right.\right.$ Range $\left.\left.\left(Q_{1}^{\perp}\right)\right)\right)$, one can write

$$
\left.Q_{1}^{\perp} \Psi_{0}(X)\right|_{\text {Range }\left(Q_{1}^{\perp}\right)}:=\sum_{i=1}^{k} S_{i}^{*} \Gamma_{i}(X) S_{i},
$$

where $\Gamma_{i} \in \operatorname{UEB}\left(\mathcal{S}, B\left(\operatorname{Range}\left(Q_{1}^{\perp}\right)\right)\right)$ are given by $\Gamma_{i}(X)=\gamma_{i}(X) I_{\text {Range }\left(Q_{1}^{\perp}\right)}$, $\gamma_{i}$ are linear extremal states defined on $\mathcal{S}, S_{i} \in B\left(\right.$ Range $\left.\left(Q_{1}^{\perp}\right)\right)$ satisfy $\sum_{i=1}^{k} S_{i}^{*} S_{i}=I_{\operatorname{Range}\left(Q_{1}^{\perp}\right)}$. Since $\Theta \in \operatorname{UEB}(\mathcal{S}, B(H))$ is a $C^{*}$-extreme map, it follows from an argument similar to the one on page 770 in [FM93], that each direct summand of $\Theta$ in equation (5.8) is also a $C^{*}$ extreme UEB map. In particular, the UEB map $\left.Q_{1}^{\perp} \Psi_{0}(X)\right|_{\operatorname{Range}\left(Q_{1}^{\perp}\right)}$ is a $C^{*}$-extreme point of $\operatorname{UEB}\left(\mathcal{S}, B\left(\operatorname{Range}\left(Q_{1}^{\perp}\right)\right)\right)$. As before, there is no loss of generality in assuming that $\left\|S_{1}\right\|=1$. We can now repeat
the arguments from before to the current set up and conclude (by taking appropriate direct sums with the zero operator) that there exists a projection $Q_{2}$ in $B(H)$ such that

$$
\Theta(X)=Q_{1} \Psi_{1}(X) Q_{1}+Q_{2} \Delta_{1}(X) Q_{2}+\text { an EB map },
$$

where $Q_{1}$ and $Q_{2}$ are mutually orthogonal projections in $B(H), \Delta_{1}(X)=$ $\gamma_{1}(X) I_{H}$ and $\gamma_{1}$ is a linear extremal state on $\mathcal{S}$. This process has to end after a finite number of steps, due to our finite dimensionality assumptions. This along with the fact that $\Theta$ is unitarily equivalent to $\Phi$ implies that $\Phi(X)=\sum_{i=1}^{\ell} \phi_{i}(X) P_{i}$, for all $X \in \mathcal{S}$, where the $\phi_{i}$ 's are linear extremal states on $\mathcal{S}$ and the $P_{i}$ 's are mutually orthogonal projections in $B(H)$ satisfying $\sum_{i=1}^{\ell} P_{i}=I_{H}$. An application of Theorem 1.10 implies that $\Phi$ is a maximal UEB map and the proof is complete.

To prove the implication (ii) $\Longrightarrow$ (iii), observe that since $\Phi$ is a $C^{*}$-extreme UEB map, by the (proved) equivalence of statements (i) and (ii) and by an application of Lemma 4.2, $\Phi$ has commutative range. That $\Phi$ is a linear extreme UEB map follows from an easy and direct adaptation of Theorem 2.2.2 in [Z] to UEB maps.

Finally, to prove the implication (iii) $\Longrightarrow$ (i), observe that since $\Phi$ is a UEB map with commutative range, Lemma 4.1 implies that

$$
\Phi(X)=\sum_{i=1}^{k} \phi_{i}(X) P_{i}
$$

for some $k \in \mathbb{N}$, where the $\phi_{i}$ 's are distinct states on $\mathcal{S}$ and the $P_{i}$ 's are mutually orthogonal projections in $B(H)$ satisfying $\sum_{i=1}^{k} P_{i}=I_{H}$. By appealing to Theorem 1.10, it suffices to show that $\phi_{i}$ is a linear extremal state for each $i$. Fix $i \in\{1, \ldots, k\}$. Let $\sigma$ and $\tau$ be states defined on $\mathcal{S}$ such that

$$
\phi_{i}=t \sigma+(1-t) \tau
$$

for some $t \in(0,1)$. It is enough to show that $\sigma=\tau$. Observe that $\Phi(X)=t\left(\sigma(X) P_{i}+\sum_{j \neq i} \phi_{j}(X) P_{j}\right)+(1-t)\left(\tau(X) P_{i}+\sum_{j \neq i} \phi_{j}(X) P_{j}\right)$,
for all $X \in \mathcal{S}$. Define the linear maps $\Psi, \Gamma: \mathcal{S} \rightarrow B(H)$ by

$$
\Psi(X)=\sigma(X) P_{i}+\sum_{j \neq i} \phi_{j}(X) P_{j} \text { and } \Gamma(X)=\tau(X) P_{i}+\sum_{j \neq i} \phi_{j}(X) P_{j}
$$

Observe that $\Psi, \Gamma \in \operatorname{UEB}(\mathcal{S}, B(H))$ and $\Phi(X)=t \Psi(X)+(1-t) \Gamma(X)$, for all $X \in \mathcal{S}$. From the assumption that $\Phi$ is linear extreme, it
follows that $\Psi=\Gamma$. This in turn implies that $\sigma=\tau$ and the proof is complete.

Remark 5.3. The following remarks concern the above proof of Theorem 1.12.
(i) The proof of the implication $(i) \Rightarrow$ (ii) in the above theorem works equally well for UCP maps or even just unital positive maps.
(ii) In the proof of the implication $(i i) \Longrightarrow(i)$, it is assumed that $T_{1}$ is not invertible and hence the appearance of the zero block in the definition of $Y_{1}$. If the operator $T_{1}$ is invertible, even then the proof works just fine by letting $Y_{1}=Z_{1}, Y_{0}=Z_{0}$ and by deleting the last row and column in the coefficient matrices occuring in equations (5.6) and (5.7).

Proof of Corollary 1.13. Since $\Phi$ is a $C^{*}$-extreme point in $\operatorname{UEB}(\mathcal{S}, B(H))$, it follows from Theorems 1.12 and 1.10 that $\Phi(X)=\sum_{i=1}^{\ell} \phi_{i}(X) P_{i}$, where the $\phi_{i}$ 's are linear extremal states on $\mathcal{S}$ and the $P_{i}$ 's are orthogonal projections in $B(H)$ such that $\sum_{i=1}^{\ell} P_{i}=I_{H}$. By [Z, Proposition 1.2.4] there exist linear extremal states $\psi_{i}: B(E) \rightarrow \mathbb{C}$ such that $\left.\psi_{i}\right|_{\mathcal{S}}=$ $\phi_{i}$, for each $i$. Define $\Psi: B(E) \rightarrow B(H)$ by $\Psi(X)=\sum_{i=1}^{\ell} \psi_{i}(X) P_{i}$. Observe that $\left.\Psi\right|_{S}=\Phi$. It follows from Theorems 1.10 and 1.12 that $\Psi$ is a $C^{*}$-extreme UEB extension of $\Phi$.

The following example shows that the converse of Corollary 1.13 is not true in general, i.e., the restriction of a $C^{*}$-extreme UEB map on an operator system need not always be $C^{*}$-extreme. Recall that for states defined on operator systems, $C^{*}$-extremality coincides with linear extremality. Let $\mathcal{S} \subset M_{2}(\mathbb{C})$ denote the operator system

$$
\mathcal{S}:=\left\{\left(\begin{array}{ll}
a & b \\
c & a
\end{array}\right) \in M_{2}(\mathbb{C}): a, b, c \in \mathbb{C}\right\} .
$$

Define $\phi: \mathcal{S} \rightarrow \mathbb{C}$ by $\phi(X)=\operatorname{trace}(X E)$ for all $X \in \mathcal{S}$, where $E=$ $\left(\begin{array}{cc}1 / 2 & 0 \\ 0 & 1 / 2\end{array}\right)$. Then $\phi$ is a state on $\mathcal{S}$. In particular, $\phi$ is a UEB map. Define the (distinct) states $\phi_{1}, \phi_{2}: \mathcal{S} \rightarrow \mathbb{C}$ by $\phi_{i}(X)=\operatorname{trace}\left(X F_{i}\right)$, for all $X \in \mathcal{S}, \mathrm{i}=1,2$, where $F_{1}=\left(\begin{array}{cc}1 / 2 & 1 / 2 \\ 1 / 2 & 1 / 2\end{array}\right)$ and $F_{2}=\left(\begin{array}{cc}1 / 2 & -1 / 2 \\ -1 / 2 & 1 / 2\end{array}\right)$.

It follows that

$$
\begin{aligned}
\phi(X) & =\operatorname{trace}\left(X\left(\begin{array}{cc}
1 / 2 & 0 \\
0 & 1 / 2
\end{array}\right)\right)=\frac{1}{2} \operatorname{trace}\left(X\left(F_{1}+F_{2}\right)\right) \\
& =\frac{1}{2}\left(\phi_{1}(X)+\phi_{2}(X)\right)
\end{aligned}
$$

This in turn implies that $\phi$ is not a linear extremal state and hence not a $C^{*}$-extreme state on $\mathcal{S}$. Let $\psi: M_{2} \rightarrow \mathbb{C}$ be the state defined by $\psi(Y)=\operatorname{trace}\left(Y E_{11}\right)$, where $E_{11}=\left(\begin{array}{ll}1 & 0 \\ 0 & 0\end{array}\right)$. Since $E_{11}$ is a projection of rank one, it is well-known that $\psi$ has to be a linear extremal state (See [St] and [Wa]) or equivalently a $C^{*}$-extremal state on $M_{2}$. Finally, observe that $\left.\psi\right|_{\mathcal{S}}=\phi$.

We end this section with a proof of Corollary 1.14, which is a KreinMilman type theorem for UEB maps.
Proof of Corollary 1.14. Recall the notations $\mathcal{E}$ and $\mathcal{K}$ from Lemma 5.2. It follows from Lemma 5.2 that $\mathcal{E}=\operatorname{UEB}(\mathcal{S}, B(H))$. By Theorems 1.10 and 1.12 , it follows that each element of $\mathcal{K}$ is a $C^{*}$-extreme UEB map in $\operatorname{UEB}(\mathcal{S}, B(H))$. This completes the proof.

## 6. $C^{*}$-extreme UEB maps between Matrix algebras

This section contains an improved version of [BDMS, Theorem 5.3], which includes various characterizations of a $C^{*}$-extreme UEB map between matrix algebras (See Theorem 6.2). Recall the convention 1.1 and the definitions of the Choi-rank, EB rank, Schmidt rank and Schmidt number from Section 1. For a positive matrix $X=\sum_{i=1}^{\ell} A_{i} \otimes$ $B_{i} \in B(E \otimes H)$, the partial trace of $X$ with respect to the first coordinate, is denoted by $\operatorname{trace}_{1}(X)$ and is defined as

$$
\operatorname{trace}_{1}(X)=\sum_{i=1}^{\ell} \operatorname{trace}\left(A_{i}\right) B_{i} \in B(H)
$$

The following lemma is a variant of [HSR, Lemma 8].
Lemma 6.1. Let $\Psi: M_{d} \rightarrow M_{n}$ be an EB map and $C_{\Psi}=\sum_{j=1}^{m} \xi_{j} \xi_{j}^{*}$ with $S R\left(\xi_{j}\right)=1$ for each $j \in\{1,2, \ldots, n\}$. If the Choi-rank of $\Psi=n$ and $\operatorname{trace}_{1}\left(C_{\Psi}\right)=\Psi\left(I_{d}\right) \in M_{n}$ is invertible, then $m \geq n$ and there exists $\gamma_{1}, \ldots, \gamma_{n} \in \mathbb{C}^{d} \otimes \mathbb{C}^{n}$ with $S R\left(\gamma_{k}\right)=1$ such that $C_{\Psi}=\sum_{k=1}^{n} \gamma_{k} \gamma_{k}^{*}$.
Proof. Since the Choi-rank of $\Psi$ equals the rank of the Choi matrix $C_{\Psi}$ (See [Wa] and [St]), it is clear that $m \geq n$. Suppose that $m>n$ and $C_{\Psi}=\sum_{j=1}^{m} \xi_{j} \xi_{j}^{*}$ where $S R\left(\xi_{j}\right)=1$ for all $1 \leq j \leq m$. For each
$1 \leq j \leq m$, let $\xi_{j}=x_{j} \otimes y_{j}$ with $\left\|x_{j}\right\|=1$. By hypothesis, $\Psi\left(I_{d}\right)=$ $\operatorname{trace}_{1}\left(C_{\Psi}\right)=\sum_{j=1}^{m} \operatorname{trace}\left(x_{j} x_{j}^{*}\right) y_{j} y_{j}^{*}=\sum_{j=1}^{m} y_{j} y_{j}^{*}$ is invertible. This implies that $\operatorname{Span}\left\{y_{1}, \ldots, y_{m}\right\}=\mathbb{C}^{n}$. Without loss of generality, assume that $\left\{y_{1}, \ldots, y_{n}\right\}$ is a basis for $\mathbb{C}^{n}=$ Range $\left(\operatorname{trace}_{1}\left(C_{\Psi}\right)\right)$. Note that $B=\left\{\xi_{j}: 1 \leq j \leq n\right\}$ is linearly independent. By the Douglas range inclusion property ( $[\mathrm{D}$, Theorem 1]) and due to the hypothesis that the rank of $C_{\Psi}=n$, it follows that $\left\{\xi_{j}: 1 \leq j \leq m\right\} \subset \operatorname{Range}\left(C_{\Psi}\right)$ and that $B$ is in fact a basis for Range $\left(C_{\Psi}\right)$. Consider the sum $\sum_{j=1}^{n+1} \xi_{j} \xi_{j}^{*}$. There exists $\alpha_{j}$ such that $\xi_{n+1}=\sum_{j=1}^{n} \alpha_{j} \xi_{j}$. It follows that $x_{j}=$ $\lambda_{j} x_{n+1}$, for some scalars $\lambda_{j}$, whenever $\alpha_{j} \neq 0,1 \leq j \leq n$. Thus $\left|\lambda_{j}\right|=1$ (since the $x_{j}$ 's are unit vectors) and

$$
x_{n+1} \otimes y_{n+1}=\xi_{n+1}=\sum_{j: \alpha_{j} \neq 0} \alpha_{j} \xi_{j}=\sum_{j: \alpha_{j} \neq 0} x_{n+1} \otimes\left(\lambda_{j} \alpha_{j}\right) y_{j} .
$$

Consequently, $y_{n+1}=\sum_{j: \alpha_{j} \neq 0}\left(\lambda_{j} \alpha_{j}\right) y_{j}$, and

$$
\begin{align*}
\sum_{j=1}^{n+1} \xi_{j} \xi_{j}^{*} & =x_{n+1} x_{n+1}^{*} \otimes\left(\left(\sum_{j: \alpha_{j} \neq 0} y_{j} y_{j}^{*}\right)+y_{n+1} y_{n+1}^{*}\right) \\
& +\sum_{j: \alpha_{j}=0} x_{j} x_{j}^{*} \otimes y_{j} y_{j}^{*} . \tag{6.1}
\end{align*}
$$

Let $r$ denote the cardinality of the set $\left\{j: \alpha_{j}=0\right\}$. Since $\left\{y_{j}: \alpha_{j} \neq\right.$ $0\}$ is linearly independent and $y_{n+1} \in \operatorname{span}\left\{y_{j}: \alpha_{j} \neq 0\right\}$,

$$
\sum_{j: \alpha_{j} \neq 0} y_{j} y_{j}^{*}+y_{n+1} y_{n+1}^{*}=\sum_{k=1}^{n-r} w_{k} w_{k}^{*},
$$

for some non-zero vectors $w_{k} \in \operatorname{span}\left\{y_{j}: \alpha_{j} \neq 0\right\} \subseteq \mathbb{C}^{n}$. Thus,

$$
\sum_{j=1}^{n+1} \xi_{j} \xi_{j}^{*}=\sum_{k=1}^{n-r} \eta_{k} \eta_{k}^{*}+\sum_{j: \alpha_{j}=0} \xi_{j} \xi_{j}^{*}
$$

where $\eta_{k}=x_{n+1} \otimes w_{k}$. Note that the right hand side is a sum consisting of exactly $n$ summands that are vectors in $\mathbb{C}^{d} \otimes \mathbb{C}^{n}$ with Schmidt rank equal to one. Repeat the above process, by adding the $(n+2)^{t h}$ term of $C_{\Psi}$ to $\sum_{j=1}^{n+1} \xi_{j} \xi_{j}^{*}$. Continuing this way, after a finite number of iterations, one can conclude that $C_{\Psi}$ is of the desired form, i.e., $C_{\Psi}=$ $\sum_{k=1}^{n} \gamma_{k} \gamma_{k}^{*}$ with $\gamma_{k} \in \mathbb{C}^{d} \otimes \mathbb{C}^{n}$ and $S R\left(\gamma_{k}\right)=1$. This completes the proof.

Theorem 6.2. Let $\Phi: M_{d} \rightarrow M_{n}$ be a UEB map. The following statements are equivalent.
(i) $\Phi$ is a $C^{*}$-extreme $U E B$ map.
(ii) $\Phi$ is a linear extreme UEB map with commutative range.
(iii) $\Phi$ is a maximal UEB map.
(iv) The map $\Phi$ has the form $\Phi(X)=\sum_{i=1}^{\ell} \phi_{i}(X) P_{i}$, where the $\phi_{i}$ 's are distinct linear extremal states defined on $M_{d}$ and the $P_{i}$ 's are mutually orthogonal projections in $M_{n}$ satisfying $\sum_{i=1}^{\ell} P_{i}=I_{n}$.
(v) The Choi-rank of $\Phi$ is $n$.
(vi) If $C_{\Phi}=\sum_{j=1}^{m} \xi_{j} \xi_{j}^{*}$ with $S R\left(\xi_{j}\right)=1$ for each $j \in\{1,2, \ldots, m\}$, then $m \geq n$ and there exists $\gamma_{1}, \ldots, \gamma_{n} \in \mathbb{C}^{d} \otimes \mathbb{C}^{n}$ with $S R\left(\gamma_{j}\right)=$ 1 such that $C_{\Phi}=\sum_{j=1}^{n} \gamma_{j} \gamma_{j}^{*}$.
(vii) If $\Phi(X)=\sum_{i=1}^{k} V_{i}^{*} X V_{i}$ and rank of $V_{i}$ is one for each $i \in$ $\{1,2, \ldots, k\}$, then $k \geq n$ and there exists matrices $W_{1}, \ldots, W_{n}$ of rank one such that $\Phi(X)=\sum_{i=1}^{n} W_{i}^{*} X W_{i}$.
(viii) The EB-rank of $\Phi$ is $n$.

Proof. The equivalence of $(i),(i i)$ and (iii) follows by letting $\mathcal{S}=B(E)$ in Theorem 1.12

The equivalence (iii) $\Longleftrightarrow(i v)$ follows by letting $\mathcal{S}=B(E)$ in Theorem 1.10.

To prove $(i v) \Longrightarrow(v)$, rewrite $\Phi(X)=\sum_{k=1}^{n} \psi_{k}(X) w_{k} w_{k}^{*}$, where $\psi_{k} \in\left\{\phi_{i}: 1 \leq i \leq \ell\right\}$ and $\left\{w_{k}: 1 \leq k \leq n\right\}$ is an orthonormal set of vectors satisfying $\sum_{k=1}^{n} w_{k} w_{k}^{*}=I_{n}$. For each $k \in\{1,2, \ldots, n\}$, Observe that $\psi_{k}(X)=\operatorname{trace}\left(X u_{k} u_{k}^{*}\right)=\left\langle X u_{k}, u_{k}\right\rangle$ for some unit vector $u_{k}$. Recall the Choi matrix $C_{\Phi} \in M_{d} \otimes M_{n}$. It follows that

$$
\begin{aligned}
C_{\Phi} & =\sum_{i, j=1}^{d} e_{i} e_{j}^{*} \otimes \Phi\left(e_{i} e_{j}^{*}\right)=\sum_{k=1}^{n} \sum_{i, j=1}^{d} e_{i} e_{j}^{*} \otimes\left(u_{k}^{*} e_{i} e_{j}^{*} u_{k}\right) w_{k} w_{k}^{*} \\
& =\sum_{k=1}^{n}\left(\sum_{i=1}^{d}\left(u_{k}^{*} e_{i}\right) e_{i} \otimes w_{k}\right)\left(\sum_{i=1}^{d}\left(e_{i}^{*} u_{k}\right) e_{i}^{*} \otimes w_{k}^{*}\right)=\sum_{k=1}^{n} z_{k} z_{k}^{*}
\end{aligned}
$$

where $z_{k}=\sum_{i=1}^{d}\left(u_{k}^{*} e_{i}\right) e_{i} \otimes w_{k}$. Since $C_{\Phi}$ is a sum of rank-one operators, the rank of $C_{\Phi}$, i.e., the Choi-rank of $\Phi$ is at most $n$. Also since $C_{\Phi}$ is a sum of the positive rank-one operators $z_{k} z_{k}^{*}$, the Douglas range inclusion property ([D, Theorem 1]) implies that $\left\{z_{k}: 1 \leq k \leq n\right\}$ is contained in Range $\left(C_{\Phi}\right)$. By the orthonormality of $\left\{z_{k}: 1 \leq k \leq\right.$ $n\} \subset \operatorname{Range}\left(C_{\Phi}\right)$ it follows that Choi-rank $(\phi) \geq n$. Thus in fact, the Choi-rank of $\Phi$ equals $n$.

Since $\Phi$ is unital, the implication $(v) \Longrightarrow \quad(v i)$ follows from the observation that $\operatorname{trace}_{1}\left(C_{\Phi}\right)=I_{n}$ and a direct application of Lemma 6.1.

The implication $(v i) \Longrightarrow(v i i)$, is a direct consequence of the well-known fact (See [St, Proposition 4.1.6] and [Bh, Theorem 3.1.1] for instance), that every spectral decomposition of $C_{\Phi}$ as a sum of $\ell$ rank one positive matrices $\xi \xi^{*}$ with $S R(\xi)=1$, yields a Choi-Kraus decomposition of $\Phi$ with exactly $\ell$ Choi-Kraus coefficients, each having rank one (and vice-versa).

The implication $(v i i) \Longrightarrow(v i i i)$ is immediate from the definition of the EB-rank of $\Phi$.

The below proof of the implication (viii) $\Longrightarrow$ (iii) is essentially the same as that given in [BDMS]. We include it here for the sake of completeness. Let $\Phi(X)=\sum_{j=1}^{n}\left(v_{j} u_{j}^{*}\right) X\left(u_{j} v_{j}^{*}\right)$, where $\left\{u_{1}, \ldots, u_{n}\right\} \subset$ $\mathbb{C}^{d}$ are unit vectors and $\left\{v_{1}, \ldots, v_{n}\right\} \subset \mathbb{C}^{n}$. It suffices to show that $\left\{v_{1}, \ldots, v_{n}\right\}$ forms an orthonormal basis for $\mathbb{C}^{n}$ or equivalently, the matrix $W:=\left(v_{1}\left|v_{2}\right| \cdots \mid v_{n}\right) \in M_{n}$ is unitary. This easily follows from the fact that $I_{n}=\Phi\left(I_{d}\right)=\sum_{j=1}^{n} v_{j} v_{j}^{*}=W W^{*}$. Let $\left\{u_{j_{k}}: 1 \leq k \leq \ell\right\}$ be the distinct unit vectors among $u_{1}, \ldots, u_{n}$. For $1 \leq k \leq \ell$, define the states $\phi_{k}$ on $M_{d}$ by $\phi_{k}(X)=\operatorname{trace}\left(X u_{j_{k}}\right)$. Rewrite $\Phi$ such that

$$
\Phi(X)=\sum_{j=1}^{n}\left(v_{j} u_{j}^{*}\right) X\left(u_{j} v_{j}^{*}\right)=\sum_{k=1}^{\ell} \phi_{k}(X) P_{k} .
$$

It is easily seen that the $P_{k}$ 's are mutually orthogonal projections satisfying $\sum_{k=1}^{\ell} P_{k}=I_{n}$. This completes the proof.

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Department of Mathematics, IIT Madras, Chennai - 600036, India.
Email address: bsriram@iitm.ac.in, bsriram80@yahoo.co.in
Department of Mathematics, IIT Madras, Chennai - 600036, India.
Email address: ma18d016@smail.iitm.ac.in


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