

# Ca<sup>2+</sup> Stores Regulate Ryanodine Receptor Ca<sup>2+</sup> Release Channels via Luminal and Cytosolic Ca<sup>2+</sup> Sites

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**ABSTRACT** The free [Ca<sup>2+</sup>] in endoplasmic/sarcoplasmic reticulum Ca<sup>2+</sup> stores regulates excitability of Ca<sup>2+</sup> release by stimulating the Ca<sup>2+</sup> release channels. Just how the stored Ca<sup>2+</sup> regulates activation of these channels is still disputed. One proposal attributes luminal Ca<sup>2+</sup>-activation to luminal facing regulatory sites, whereas another envisages Ca<sup>2+</sup> permeation to cytoplasmic sites. This study develops a unified model for luminal Ca<sup>2+</sup> activation for single cardiac ryanodine receptors (RyR<sub>2</sub>) and RyRs in coupled clusters in artificial lipid bilayers. It is shown that luminal regulation of RyR<sub>2</sub> involves three modes of action associated with Ca<sup>2+</sup> sensors in different parts of the molecule; a luminal activation site (*L*-site, 60 μM affinity), a cytoplasmic activation site (*A*-site, 0.9 μM affinity), and a novel cytoplasmic inactivation site (*I*<sub>2</sub>-site, 1.2 μM affinity). RyR activation by luminal Ca<sup>2+</sup> is demonstrated to occur by a multistep process dubbed luminal-triggered Ca<sup>2+</sup> feedthrough. Ca<sup>2+</sup> binding to the *L*-site initiates brief openings (1 ms duration at 1–10 s<sup>-1</sup>) allowing luminal Ca<sup>2+</sup> to access the *A*-site, producing up to 30-fold prolongation of openings. The model explains a broad data set, reconciles previous conflicting observations and provides a foundation for understanding the action of pharmacological agents, RyR-associated proteins, and RyR<sub>2</sub> mutations on a range of Ca<sup>2+</sup>-mediated physiological and pathological processes.

## INTRODUCTION

In cardiac muscle, contraction is mediated by depolarization and subsequent opening of voltage-dependent, L-type Ca<sup>2+</sup> channels in the cell membrane (sarcolemma). Opening of these channels allows Ca<sup>2+</sup> influx into the cell, which in turn opens Ca<sup>2+</sup> release channels in the sarcoplasmic reticulum (SR) membrane. SR Ca<sup>2+</sup> release can be a small highly localized increase in Ca<sup>2+</sup> concentration called a spark, or can propagate along the cell in a Ca<sup>2+</sup> wave, or can occur simultaneously along the length of the cell resulting in a global increase of Ca<sup>2+</sup> called a transient (1). Ca<sup>2+</sup> sparks are believed to be the elementary SR Ca<sup>2+</sup> release event and Ca<sup>2+</sup> waves and transients are thought to be the summation of spark events (2).

It has long been known that the excitability of Ca<sup>2+</sup> release from muscle SR is substantially increased by its luminal Ca<sup>2+</sup> load ([Ca<sup>2+</sup>]<sub>L</sub>) (3,4). More recently, it has been shown that the excitability of neuronal endoplasmic reticulum (ER) is regulated in the same way, suggesting that modulation of Ca<sup>2+</sup> signaling by [Ca<sup>2+</sup>]<sub>L</sub> is a general phenomenon (5). However, the basis for load-dependent excitability remains controversial. Experiments on muscle cells and isolated SR vesicles have found that the effect of [Ca<sup>2+</sup>]<sub>L</sub> on Ca<sup>2+</sup> release could not be explained entirely by the associated Ca<sup>2+</sup> gradient across the SR and that [Ca<sup>2+</sup>]<sub>L</sub> must somehow control the Ca<sup>2+</sup> permeability of the membrane (6,7). This was confirmed when activity of isolated RyRs in artificial bilayers was found to be modulated by [Ca<sup>2+</sup>]<sub>L</sub> (8). At issue is the mechanism underlying the activation of RyRs by luminal Ca<sup>2+</sup> being attributed to

two quite different processes. The true-luminal hypothesis attributes luminal Ca<sup>2+</sup>-activation and inhibition to Ca<sup>2+</sup> regulatory sites on the luminal side of the RyR (9). The feedthrough hypothesis proposes that luminal Ca<sup>2+</sup> permeates the pore and binds to cytoplasmic activation and inhibition sites (10–12). The latter is supported by the close correlation between the effect of [Ca<sup>2+</sup>]<sub>L</sub> and Ca<sup>2+</sup> flux (lumen to cytoplasm), and by the observation that luminal regulation is dependent on Ca<sup>2+</sup> buffering on the cytoplasmic side of the membrane (11). However, there is an increasing body of evidence suggesting that there are sensing sites for Ca<sup>2+</sup> regulation on the luminal side of the protein. For example, it has been shown that luminal Ca<sup>2+</sup>-activation is abolished by tryptic digestion of the luminal side of cardiac ryanodine receptors (RyR<sub>2</sub>) (13). Although a recent bilayer study has presented evidence that both mechanisms somehow control skeletal ryanodine receptors (RyR<sub>1</sub>) (14), it is not understood how cytoplasmic and luminal Ca<sup>2+</sup> sites explain current observations or reconcile conflicting interpretations in the literature. In this study, measurements of the gating kinetics of isolated RyR<sub>2</sub> in lipid bilayers underlie the proposal of a novel mechanism for activation of RyRs by luminal-triggered Ca<sup>2+</sup> feedthrough. This model predicts that in addition to cytoplasmic sites for Ca<sup>2+</sup>-activation and low affinity Ca<sup>2+</sup>/Mg<sup>2+</sup>-inhibition, RyRs possess a luminal Ca<sup>2+</sup> binding site and a high affinity cytoplasmic Ca<sup>2+</sup> inactivation site (Fig. 1). Ca<sup>2+</sup> binding at the luminal site is sufficient to activate channel openings whereupon Ca<sup>2+</sup> flows from lumen to cytoplasm (feedthrough) either enhances activation or causes inactivation.

Little is known about the mechanisms underlying the activation and termination of Ca<sup>2+</sup> sparks in muscle. The quantal nature of Ca<sup>2+</sup> spark intensities suggests that they

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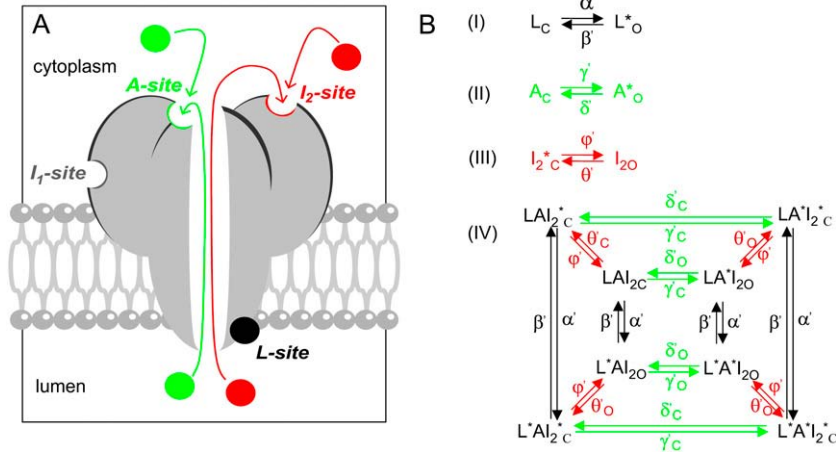


FIGURE 1 The luminal-triggered  $\text{Ca}^{2+}$  feedthrough model for  $[\text{Ca}^{2+}]_L$  regulation of RyRs. (A) Schematic of luminal-triggered  $\text{Ca}^{2+}$  feedthrough.  $\text{Ca}^{2+}$  binding to the  $L$ -site causes channel opening whereupon luminal  $\text{Ca}^{2+}$  has access to the cytoplasmic  $\text{Ca}^{2+}$  activation ( $A$ -site) and inactivation sites ( $I_2$ -site). The  $I_1$ -site that mediates the low affinity  $\text{Ca}^{2+}/\text{Mg}^{2+}$  is not included in the model (see text for details). (B) Kinetic schemes for  $\text{Ca}^{2+}$  binding at the  $L$ -,  $A$ -, and  $I_2$ -sites (Schemes I–III) and the overall scheme (Scheme IV) resulting from the combined action of  $\text{Ca}^{2+}$  at all three sites. Schemes I–III are representations of multistep processes. Hence some reaction rates have complex dependencies on  $[\text{Ca}^{2+}]$  and these are given in the equations listed in Table 1. Asterisks indicate sites occupied with  $\text{Ca}^{2+}$ . Reaction rates that depend on the  $\text{Ca}^{2+}$  current have subscripts  $o$  and  $c$  to indicate rates associated with open and closed channels, respectively. The open and closed status of the channel associated with each kinetic state is indicated by the subscripts  $O$  and  $C$ , respectively.

are produced by the concerted opening of up to 10 RyRs (15). Recent studies of the coupled gating of RyRs in artificial lipid bilayers may provide new insight into the triggering mechanisms for  $\text{Ca}^{2+}$  sparks (14,16,17). These investigations showed that  $\text{Ca}^{2+}$  feedthrough not only regulates the host RyR but also activates adjacent RyRs in a cluster. This study examines the roles of the luminal and cytoplasmic  $\text{Ca}^{2+}$  sites in generating this phenomenon in bilayers and how they might contribute to spark generation in muscle.

## MATERIALS AND METHODS

SR vesicles (containing RyR<sub>2</sub>) were obtained from sheep hearts and were reconstituted into artificial lipid bilayers as previously described (18). Lipid bilayers were formed from phosphatidylethanolamine and phosphatidylcholine (8:2 wt/wt) in *n*-decane (50 mg/ml). During experiments, the *cis* (cytoplasmic) and *trans* (luminal) solutions contained 250 mM  $\text{Cs}^+$  (230 mM  $\text{CsCH}_3\text{O}_3\text{S}$ , 20 mM  $\text{CsCl}$ ) and various concentrations of  $\text{CaCl}_2$ . The *trans* solution was altered by addition of aliquots of stock solutions and the *cis* solution by local perfusion which allowed solution exchange within  $\sim 1$  s (19).

Solutions were pH-buffered with 10 mM *n*-tris[Hydroxymethyl]methyl-2-aminoethanesulfonic acid and solutions were titrated to pH 7.4 using  $\text{CsOH}$ . Free  $[\text{Ca}^{2+}]$  up to 100 nM was estimated using published association constants (20) and the program Bound and Determined (21) and concentrations higher than this were measured using a  $\text{Ca}^{2+}$  electrode (Radiometer, Copenhagen, Denmark).  $[\text{Ca}^{2+}]$  below 10  $\mu\text{M}$  was buffered with 4.5 mM 1,2-bis(*o*-aminophenoxy)ethane-*n,n,n',n'*-tetraacetic acid ( $4\text{K}^+$ ) (BAPTA) and titrated with  $\text{CaCl}_2$ .  $[\text{Ca}^{2+}]$  in the range 10–50  $\mu\text{M}$  was buffered with either sodium citrate (up to 6 mM) or dibromo BAPTA (up to 2 mM).

## Acquisition and analysis of ion channel recordings

Bilayer apparatus and data recording methods have been described previously (14). Electrical potentials are expressed using standard physiological convention (i.e., cytoplasmic side relative to the luminal side at virtual ground). Measurements were carried out at  $23 \pm 2^\circ\text{C}$ . Before analysis, the current signal was digitally filtered at 1 kHz with a Gaussian

filter and sampled at 5 kHz. Single channel properties were measured using Channel2 software (P. W. Gage and M. Smith, Australian National University, Canberra). Open probability ( $P_o$ ) and open and closed durations were calculated from single channel records using a threshold discriminator at 50% of channel amplitude. For experiments in which bilayers contained several RyRs,  $P_o$  was calculated from the time-averaged current divided by the unitary current and the number of channels ( $n$ ). Both methods of calculating  $P_o$  gave similar results. The mean channel open and closed durations ( $\tau_o$  and  $\tau_c$ ) could also be calculated from mean open and closed durations in multichannel recordings ( $\tau_o(n)$  and  $\tau_c(n)$ ), provided that multiple openings were rare:  $\tau_o = \tau_o(n)$  and  $\tau_c = \tau_c(n) \times n$ . The number of channels in each experiment was determined during periods of strong activation, which were usually achieved by turning off the local perfusion and exposing the RyRs to the *cis* bath. The values  $\tau_o$  and  $\tau_c$  were calculated mainly from the recording of 100–1000 opening events. However, under conditions that produced extremely low channel activity, the mean durations were obtained from as few as 40 events covering  $>400$  s of recording (sampling error  $<40^{-1/2}$ , i.e.,  $<16\%$ ). Dwell-time frequency histograms of channel open and closed events were obtained exclusively from single channel recordings and were displayed as probabilities (counts/total number of events). These were plotted using the log-bin method suggested by Sigworth and Sine (22), which displays exponentials as peaked distributions centered at their exponential time constant. Sampling bins were equally spaced on a log scale with 3.5 or 7 bins per decade.

On rare occasions ( $<5\%$  of experiments), RyRs appeared to form coupled clusters in the bilayer, similar to that seen for RyR<sub>1</sub> (14). The experiments on coupled and uncoupled RyRs were analyzed and presented separately. Channel gating in coupled clusters was analyzed using the Hidden Markov Model (23), as described previously (14). The algorithm calculated the transition probability matrix of the underlying Markov process from the raw signal using maximum likelihood criteria. The mean channel opening and closing rates were calculated from the transition probability matrix.

$\text{Ca}^{2+}$ -dependencies of  $P_o$ ,  $\tau_o$ , and opening rate were characterized by fitting these data with Hill curves using the following equations for activation and inactivation (shown here for the case of  $P_o$ , similar equations apply to  $\tau_o$  and opening rate):

$$P_o = P_{\min} + \frac{(P_{\max} - P_{\min})([\text{Ca}^{2+}]/K_a)^{n_a}}{1 + ([\text{Ca}^{2+}]/K_a)^{n_a}},$$

$$P_o = P_{\min} + \frac{(P_{\max} - P_{\min})}{1 + ([\text{Ca}^{2+}]/K_i)^{n_i}}.$$

The expressions  $P_{\min}$  and  $P_{\max}$  are the activities of the minimally and maximally activated channel,  $K_a$  and  $K_i$  are the  $[\text{Ca}^{2+}]$  for half-activation and inhibition, and  $n_a$  and  $n_i$  are the corresponding Hill coefficients. The theory was fitted with the data using the method of least-squares.

## The luminal-triggered $\text{Ca}^{2+}$ feedthrough model: Kinetic schemes and equations

RyRs are known to possess cytoplasmic sites for  $\text{Ca}^{2+}$ -activation ( $A$ -sites  $\sim 1 \mu\text{M}$  affinity) and low-affinity  $\text{Ca}^{2+}/\text{Mg}^{2+}$ -inhibition ( $I_1$ -site,  $\sim 10 \text{ mM}$  affinity, previously named the  $I$ -site). Based on novel data presented below, the model includes a cytoplasmic  $\text{Ca}^{2+}$ -inactivation site (the  $I_2$ -site, which causes partial inactivation with  $\sim 1 \mu\text{M}$  affinity) and a luminal activation site ( $L$ -site,  $\sim 60 \mu\text{M}$  affinity). Fig. 1 *A* shows a schematic representation of the locations of these sites on the RyR protein. It is envisaged that the channel can open if  $\text{Ca}^{2+}$  is bound to either the  $A$ - or  $L$ -sites. Thus luminal  $\text{Ca}^{2+}$  can open the channel by binding to the  $L$ -site, whereupon the flow of  $\text{Ca}^{2+}$  through the pore can increase the cytoplasmic  $[\text{Ca}^{2+}]$  in the vicinity of the pore mouth ( $[\text{Ca}^{2+}]_p$ ) and reinforce channel activation (increasing  $\tau_o$ ) by binding to the  $A$ -site or inactivate the channel by binding to the  $I_2$ -site (decreasing  $\tau_o$ ).

Kinetic schemes for  $\text{Ca}^{2+}$  binding at the  $L$ -,  $A$ -, and  $I_2$ -sites are shown in Fig. 1 *B* (*I–III*) and the overall regulation of RyRs by the three  $\text{Ca}^{2+}$  sites is described by an eight-state model (Scheme *IV*) presenting the amalgamation of Schemes *I–III*. The equations describing the action of luminal and cytoplasmic  $\text{Ca}^{2+}$  on channel activity and the model parameters are listed in Table 1. The  $I_1$ -site was not considered in the model calculations because  $\text{Ca}^{2+}$  feedthrough is unlikely to increase  $[\text{Ca}^{2+}]_p$  to levels where significant  $\text{Ca}^{2+}$  binding will occur at this site.

The first step in developing the luminal-triggered  $\text{Ca}^{2+}$  feedthrough model was to formulate a phenomenological description of RyR<sub>2</sub> regulation by cytoplasmic  $\text{Ca}^{2+}$ . The kinetic schemes used previously to describe cytoplasmic  $\text{Ca}^{2+}$ -activation of RyRs are complex and involve many closed and open states (24,25). To avoid this level of complexity the whole cytoplasmic  $\text{Ca}^{2+}$ -activation process was lumped into two single-step  $\text{Ca}^{2+}$  binding schemes associated with the  $A$ - and  $I_2$ -sites (Schemes *II* and *III* in Fig. 1 *B*). Thus, the parameters  $\gamma$ ,  $\delta$ ,  $\theta$ ,  $\phi$ ,  $\Delta$ , and  $K_i$  in Eqs. 4–10 were adjusted until the model fitted the cytoplasmic  $\text{Ca}^{2+}$  concentration ( $[\text{Ca}^{2+}]_c$ )-dependencies of  $\tau_o$  and  $\tau_c$  (Figs. 5 and 10). The second step was to ascertain the gating properties associated with the  $L$ -site (Scheme *I*). This was done by fitting the voltage/luminal  $\text{Ca}^{2+}$  concentration ( $[\text{Ca}^{2+}]_p$ )-dependence of  $\tau_c$  (Fig. 6) using the parameters  $\alpha$ ,  $\beta$ ,  $K_L$  in Eqs. 2 and 3. The final step was to calculate the effects of  $\text{Ca}^{2+}$  feedthrough on the voltage/ $[\text{Ca}^{2+}]_L$ -depen-

dencies of  $\tau_o$ ,  $\tau_c$ , and their associated dwell-time distributions from theoretical estimates of  $[\text{Ca}^{2+}]_p$  and compare these estimates with the data.

Ion diffusion theory (26) was used to calculate  $[\text{Ca}^{2+}]_p$  and its dependence on  $I_{\text{Ca}}$  (pA) and distance from the pore ( $r$ , nm) under the buffering conditions used here.

$$[\text{Ca}^{2+}]_p = I_{\text{Ca}} \frac{825}{r} \exp\left(\frac{-r}{6}\right) + [\text{Ca}^{2+}]_c. \quad (1)$$

For a given distance from the pore,  $[\text{Ca}^{2+}]_p$  has a linear dependence on  $I_{\text{Ca}}$ , which is encapsulated in the parameters  $X$  in Eq. 7b and  $Y$  in Eq. 10b (see Table 1).  $I_{\text{Ca}}$  was calculated using a rate theory model of RyR conductance (27). This part of the model has only two adjustable parameters,  $X$  and  $Y$ .

Model predictions were made by simulating the time series of channel gating events and using these to generate theoretical open and closed dwell time distributions,  $\tau_o$ ,  $\tau_c$ , and  $P_o$ . For each set of experimental conditions, the model reaction rates were calculated using the equations in Table 1. The sojourn of the simulated channel through its various states was driven by a random number generator in conjunction with the transition probability matrix that was derived from the model reaction rates. Sequential open and closed events were marked by transitions between open and closed states in the model. The effects of limited time resolution in the data were accommodated by amalgamating simulated events shorter than  $<200 \mu\text{s}$ .

## RESULTS

The effects of cytoplasmic and luminal  $\text{Ca}^{2+}$  on RyR<sub>2</sub> in the presence of ATP are presented here first because 1), it is recognized that the effects of luminal  $\text{Ca}^{2+}$  require the presence of cytoplasmic agonists such as ATP, caffeine, or sulmazole (8,28,29); and 2), ATP is the physiological coagonist. Data are presented that define the affinity and gating properties of one luminal and two cytoplasmic  $\text{Ca}^{2+}$  sites on RyR<sub>2</sub>. These properties were used to construct the luminal-triggered  $\text{Ca}^{2+}$  feedthrough model. The results were then fitted by adjusting the model parameters associated with the proximity of the two cytoplasmic  $\text{Ca}^{2+}$  binding sites to the pore mouth. The model is then applied to elucidating the effects of ATP and ion channel coupling on RyR regulation by  $\text{Ca}^{2+}$ .

**TABLE 1** Equations and parameters of the luminal-triggered  $\text{Ca}^{2+}$  feedthrough model

Eq.	Eq.					
<b><i>L</i>-site</b>						
2	$\alpha' = \alpha \times \exp\left(\frac{\Delta eV}{kT}\right) \frac{[\text{Ca}^{2+}]_L^2}{K_L^2 + [\text{Ca}^{2+}]_L^2}$	3	$\beta' = \beta$	Affinity	Opening (0V)	V-dependence
				+ATP $K_L = 60 \mu\text{M}$	$\alpha = 2.7 \text{ s}^{-1}$	$\Delta = 0.4$
				-ATP $K_L = 45 \mu\text{M}$	$\alpha = 0.6 \text{ s}^{-1}$	$\Delta = 0.4$
						Closing
						$\beta = 1000 \text{ s}^{-1}$
						$\beta = 1000 \text{ s}^{-1}$
<b><i>A</i>-site</b>						
4	$\gamma'_c = \gamma [\text{Ca}^{2+}]_c^3$	5	$\delta'_c = \delta [\text{Ca}^{2+}]_c^{-n}$	Affinity	Opening	Feedthrough
				+ATP $K_A = 1 \mu\text{M}$	$\gamma = 43 \text{ s}^{-1} \mu\text{M}^{-3}$	$X = 15 \mu\text{M/pA}$
6	$\gamma'_o = \gamma [\text{Ca}^{2+}]_p^3$	7a,	$\delta'_o = \delta [\text{Ca}^{2+}]_p^{-n}$	-ATP $K_A = 6 \mu\text{M}$	$\gamma = 1.0 \text{ s}^{-1} \mu\text{M}^{-3}$	$X = 15 \mu\text{M/pA}$
		7b	$[\text{Ca}^{2+}]_p = XI_{\text{Ca}} + [\text{Ca}^{2+}]_c$			$\delta = 220 \text{ s}^{-1} \mu\text{M} (n = 1)$
						$\delta = 700 \text{ s}^{-1} (n = 0)$
<b><i>I</i><sub>2</sub>-site</b>						
8	$\theta'_c = \theta \frac{[\text{Ca}^{2+}]_c^2}{K_i^2 + [\text{Ca}^{2+}]_c^2}$	9	$\phi' = \phi$	Affinity	Closing	Feedthrough
				+ATP $K_i = 1.2 \mu\text{M}$	$\theta = 250 \text{ s}^{-1}$	$Y = 0.35 \mu\text{M/pA}$
10a	$\theta'_o = \theta \frac{[\text{Ca}^{2+}]_p^2}{K_i^2 + [\text{Ca}^{2+}]_p^2}$	10b	$[\text{Ca}^{2+}]_p = YI_{\text{Ca}} + [\text{Ca}^{2+}]_c$	-ATP $K_i = 1.2 \mu\text{M}$	$\theta = 800 \text{ s}^{-1}$	$Y = 0.35 \mu\text{M/pA}$
						Opening
						$\phi = 300 \text{ s}^{-1}$
						$\phi = 400 \text{ s}^{-1}$

Equations describe the  $[\text{Ca}^{2+}]$ -dependencies of opening and closing rates associated with each site (refer to Scheme *IV* in Fig. 1 *B*).  $K_A$  was derived from the theoretical half-maximal  $[\text{Ca}^{2+}]_c$ -activation of  $P_o$  and was not used in the fitting of the model. The reaction rates are distinguished from model parameters with an apostrophe (e.g.,  $\alpha'$ ).  $\gamma'$ ,  $\delta'$ , and  $\theta'$  depend on the  $\text{Ca}^{2+}$  current through the pore and therefore take on different values for the open and closed states (indicated by subscripts,  $o$  and  $c$ ). In Eq. 2,  $T$  is the temperature in degrees K,  $k$  is the Boltzmann constant and  $V$  is bilayer voltage.

## Activation of RyR<sub>2</sub> by cytoplasmic Ca<sup>2+</sup> in the presence of ATP

It is shown that the activity of RyR<sub>2</sub> was dependent on both [Ca<sup>2+</sup>]<sub>C</sub> (Fig. 2 A) and [Ca<sup>2+</sup>]<sub>L</sub> (Fig. 2 B). The experiments in this section focus on the properties of the A- and I<sub>2</sub>-sites as revealed by the effects of [Ca<sup>2+</sup>]<sub>C</sub> on RyR<sub>2</sub> activity. Fig. 2 A highlights the gradual transition in the gating properties that occurs as the dominant cytoplasmic agonist switches from ATP to Ca<sup>2+</sup>. RyRs activated by ATP (100 nM [Ca<sup>2+</sup>]<sub>C</sub>, Fig. 2 A, top three traces) had relatively long open and closed events compared to RyRs activated by high [Ca<sup>2+</sup>]<sub>C</sub> (with 2 mM ATP, Fig. 2 A, bottom two traces). To probe the gating mechanisms for this behavior, RyR gating was analyzed by compiling frequency histograms of open and closed times. These could be described by the sum of two exponential decays (these appear as peaked distributions when using the log-bin method in Fig. 3). The time-constants and relative weighting of the exponential components depended on [Ca<sup>2+</sup>]<sub>C</sub> and [Ca<sup>2+</sup>]<sub>L</sub>. Time-constants within each open time distribution had a 3–10-fold separation and lay within the range 1–50 ms (e.g., dashed lines in Fig. 3 A). The closed time distributions varied considerably with [Ca<sup>2+</sup>]<sub>C</sub>. In the presence of 0.1 μM [Ca<sup>2+</sup>]<sub>C</sub>, two widely separated exponentials were clearly resolved (open arrows in Fig. 3 B) in which >80% of the events were associated with the longer time-constant. Raising [Ca<sup>2+</sup>]<sub>C</sub> progressively reduced the longer time constant until both were quite similar at 10 μM (solid arrows in Fig. 3 B). The dwell-time distributions are frequently represented here by their means. The value τ<sub>o</sub> gave a good representation of the average time-constant of the open time distribution and the value for τ<sub>c</sub> lay within 20% of the longer closed time constant.

The effects of [Ca<sup>2+</sup>]<sub>C</sub> on the RyR<sub>2</sub> gating properties, P<sub>o</sub>, τ<sub>o</sub>, and τ<sub>c</sub>, are shown in Fig. 4. Cytoplasmic Ca<sup>2+</sup> activated RyR<sub>2</sub> at μM levels (Fig. 4 A). Luminal Ca<sup>2+</sup> had a marked effect on P<sub>min</sub> at -40 mV (Hill parameters are listed in Table 2), whereas it had a relatively small effect on P<sub>max</sub>. The values for K<sub>a</sub> and hence A-site affinity, showed no significant dependence on [Ca<sup>2+</sup>]<sub>L</sub> or membrane potential (P > 0.06, see Table 2). The [Ca<sup>2+</sup>]<sub>C</sub>-dependencies of τ<sub>o</sub> shown in Fig.

4 B reveal activation and inactivation processes associated with μM [Ca<sup>2+</sup>]<sub>C</sub>. In the virtual absence of luminal Ca<sup>2+</sup> ([Ca<sup>2+</sup>]<sub>L</sub> < 10 μM), τ<sub>o</sub> increased with [Ca<sup>2+</sup>]<sub>C</sub> (channel activation, Fig. 4 B, solid circle), whereas 100 μM [Ca<sup>2+</sup>]<sub>L</sub> unmasked a previously unidentified [Ca<sup>2+</sup>]<sub>C</sub>-dependent inactivation that decreased τ<sub>o</sub> (Fig. 4 B, open circle). Despite this inactivation, raised [Ca<sup>2+</sup>]<sub>C</sub> caused an increase in P<sub>o</sub> via a decrease in τ<sub>c</sub> that overwhelmed the effect of inactivation (P<sub>o</sub> = τ<sub>o</sub>/(τ<sub>c</sub> + τ<sub>o</sub>) with τ<sub>c</sub> ∝ [Ca<sup>2+</sup>]<sub>C</sub><sup>-3</sup> in Fig. 4 C and τ<sub>o</sub> ∝ [Ca<sup>2+</sup>]<sub>C</sub> in Fig. 4 B). The value τ<sub>o</sub> at 1000 μM [Ca<sup>2+</sup>]<sub>L</sub> had values between those at <10 and 100 μM [Ca<sup>2+</sup>]<sub>L</sub> and did not show a significant dependence on [Ca<sup>2+</sup>]<sub>C</sub>. The value τ<sub>c</sub> at 1000 μM [Ca<sup>2+</sup>]<sub>L</sub> was very similar to that seen at 100 μM. The underlying mechanism for the complex effects of [Ca<sup>2+</sup>]<sub>C</sub> and [Ca<sup>2+</sup>]<sub>L</sub> on τ<sub>o</sub> and τ<sub>c</sub> are explained with the description of the model characteristics (see below).

In the model, the activation and inactivation phenomena are attributed to the A- and I<sub>2</sub>-sites, respectively. The properties of these sites are derived from the data in Fig. 4 as follows: RyR gating due to the A- site is described by Scheme II (Fig. 1 B). In the absence of luminal Ca<sup>2+</sup>, the opening and closing rates associated with the A-site are reflected in the channel open and closed times (γ' ≈ 1/τ<sub>c</sub> and δ' ≈ 1/τ<sub>o</sub>). To account for the [Ca<sup>2+</sup>]<sub>C</sub>-dependencies of τ<sub>o</sub> and τ<sub>c</sub> in Fig. 4, δ' ∝ [Ca<sup>2+</sup>]<sub>C</sub> (Fig. 4 B, solid circle, Eqs. 5 and 7) and γ' ∝ [Ca<sup>2+</sup>]<sub>C</sub><sup>3</sup> (Fig. 4 C, solid circle; Eqs. 4 and 6). RyR gating due to the I<sub>2</sub> inactivation site is described by Scheme III (Fig. 1 B) where the equivalent opening and closing rates (φ' and θ') are given by Eqs. 8–10a. Inactivation in response to [Ca<sup>2+</sup>]<sub>C</sub> appears to be only partial (i.e., <40%), because even at 10 μM Ca<sup>2+</sup> where the I<sub>2</sub>-site is saturated, P<sub>o</sub> > 0.6. To account for this the closing rate, θ', is given a sublinear dependence on [Ca<sup>2+</sup>]<sub>P</sub> with an upper limit.

## Bell-shaped dependence of RyR<sub>2</sub> activity on [Ca<sup>2+</sup>]<sub>L</sub> in the presence of ATP

Although several studies have reported a bell-shaped [Ca<sup>2+</sup>]<sub>L</sub>-dependence of RyR<sub>1</sub> P<sub>o</sub> in the presence of ATP, there has

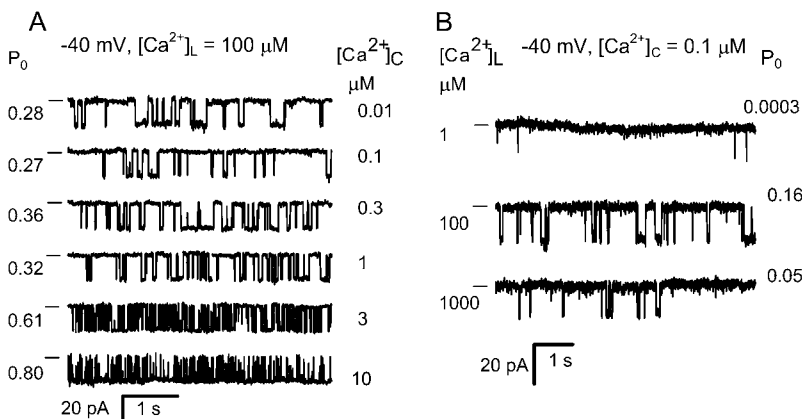


FIGURE 2 The effects of [Ca<sup>2+</sup>]<sub>L</sub> and [Ca<sup>2+</sup>]<sub>C</sub> on the activity of RyR<sub>2</sub> (-40 mV). (A) The effect of [Ca<sup>2+</sup>]<sub>C</sub> on the activity of a RyR in the presence of 100 μM [Ca<sup>2+</sup>]<sub>L</sub>. Increasing [Ca<sup>2+</sup>]<sub>C</sub> increases P<sub>o</sub> in association with a decrease in duration of mean open and closed events. Channel openings are downward current jumps from the baseline (indicated with a dash). (B) The effect of [Ca<sup>2+</sup>]<sub>L</sub> on RyR activity in the presence of 0.1 μM [Ca<sup>2+</sup>]<sub>C</sub> and 2 mM ATP.

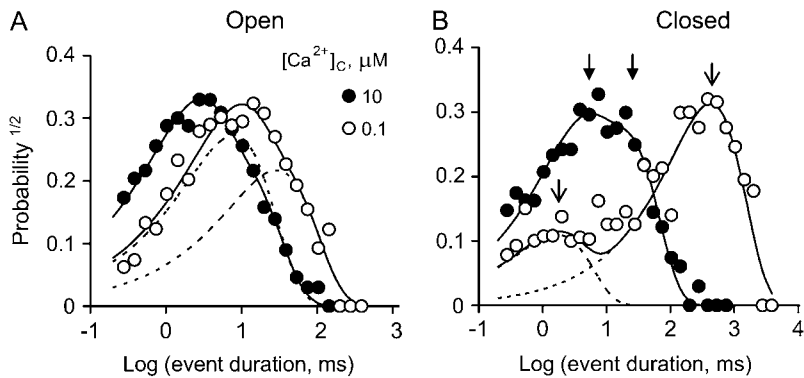


FIGURE 3 Dwell-time probability distributions of channel open (A) and closed (B) events. The data are plotted using the log-bin method of Sigworth and Sine (22). Event duration distributions were compiled from a single RyR, which was activated by 2 mM ATP and the indicated  $[Ca^{2+}]_C$  in the presence of 100  $\mu M [Ca^{2+}]_L$  ( $-40$  mV). An increase in  $[Ca^{2+}]_C$  from 0.1 to 10  $\mu M$  shifted both the open and closed distributions to shorter times. Solid curves show double-exponential fits to the data and dashed curves show individual exponential components for the fits to the open circles. When using the log-bin method the exponential components within these distributions appear as peaks centered at their respective time-constants. The arrows in panel B show the time constants associated with the data at 0.1  $\mu M [Ca^{2+}]_C$  (○, *open arrows*) and 10  $\mu M [Ca^{2+}]_C$  (●, *solid arrows*).

been no equivalent demonstration of this in RyR<sub>2</sub>. In the absence of luminal  $Ca^{2+}$  the activity of RyR<sub>2</sub> was virtually zero ( $P_o = (8 \pm 4) \times 10^{-4}$  at  $[Ca^{2+}]_L \sim 0.1 \mu M$ ,  $n = 6$ ) indicating a near-absolute requirement for luminal  $Ca^{2+}$ . Increasing  $[Ca^{2+}]_L$  (from 1 to 100  $\mu M$  at  $-40$  mV) substantially increased the activity of the RyRs, whereas a further increase to 1 mM decreased channel activity (Fig. 2 B). At  $-60$  mV and  $-40$  mV, voltages favoring the flow of  $Ca^{2+}$  from luminal to cytoplasmic baths, the  $[Ca^{2+}]_L$  for half-activation of  $P_o$ , was 50  $\mu M$  ( $K_a$  for the *left side* of the *bell*) and  $[Ca^{2+}]_L$  for half-inactivation was 1 mM ( $K_i$  for the *right side* of the *bell*, Fig. 5, A and B). Both  $K_a$  and  $K_i$  increased as the bilayer voltage became more positive (Fig. 5, C and D), and at positive voltages the decline in  $P_o$  was beyond the experimental range of  $[Ca^{2+}]_L$  ( $K_i$  could no longer be monitored).

Luminal  $Ca^{2+}$  would not be expected to influence the channel-opening rate ( $1/\tau_o$ ) other than by affecting a luminal site because cytoplasmic sites are inaccessible to luminal ions when the channel is closed. Therefore, analysis of  $\tau_o$  and  $\tau_c$  should provide valuable information of the effect of  $Ca^{2+}$  feedthrough in channel activation. The results show that  $[Ca^{2+}]_L$  had very different effects on  $\tau_o$  and channel

opening rate (Fig. 6). The value  $\tau_o$  had a bell-shaped dependence on  $[Ca^{2+}]_L$ , which is shown at several voltages in Fig. 6 A (the bell is most clearly seen at  $-20$  mV). The  $K_a$  and  $K_i$  values for the  $[Ca^{2+}]_L$ -dependence of  $\tau_o$  over the entire experimental voltage range is shown in Fig. 6 C. The values  $K_a$  and  $K_i$  increased with voltage in a way expected for  $Ca^{2+}$  ions, which must cross the membrane to reach their effector site (*long-dashed line*, Fig. 6 C). These results indicate that luminal  $Ca^{2+}$  activation and inactivation of  $\tau_o$  are mediated by cytoplasmic facing sites. Moreover, the nonadditive effects of luminal and cytoplasmic  $Ca^{2+}$  on channel activation and inhibition indicate that the luminal and cytoplasmic actions are mediated by common binding sites.

Luminal  $Ca^{2+}$  activated RyR<sub>2</sub> by increasing their opening rate and approached a maximum ( $V_{max}$ ) at high concentrations (Fig. 6 B). The voltage-dependence of  $K_a$  and  $V_{max}$  are shown in Fig. 6 D. The  $K_a$  indicates the involvement of a  $Ca^{2+}$  binding site with an affinity ( $K_L$ )  $\sim 60 \mu M$  and the  $V_{max}$  indicates that the  $Ca^{2+}$ -bound site can trigger channel openings at a rate of  $1\text{--}10 \text{ s}^{-1}$ . The  $K_a$  for opening rate did not show a significant voltage-dependence ( $P > 0.15$ , *t*-test) indicating that luminal  $Ca^{2+}$  ions do not move through the *trans*-membrane electric field to reach the site.  $V_{max}$  varied

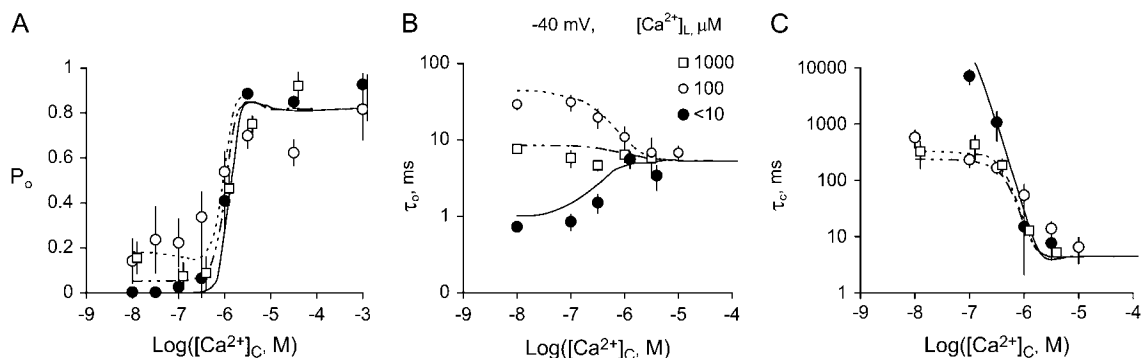


FIGURE 4 The effect of  $[Ca^{2+}]_C$  on the activity of RyR<sub>2</sub> at  $-40$  mV. The values  $P_o$ ,  $\tau_o$ , and  $\tau_c$  were measured in the presence of 2 mM ATP at  $[Ca^{2+}]_L < 10 \mu M$  (●), 100  $\mu M$  (○), or 1000  $\mu M$  (□). (A) Mean open probability  $P_o$  mean  $\pm$  SE. The numbers of experiments and the Hill parameters are listed in Table 2 (Hill fits to the data are not shown). (B, C) The mean  $\pm$  SE of three-to-eight measurements of  $\tau_o$  and  $\tau_c$ . The curves  $[Ca^{2+}]_L < 10 \mu M$  (*solid*), 100  $\mu M$  (*short dashes*), and 1000  $\mu M$  (*short/long dashes*) show the fit to the data of the luminal-triggered  $Ca^{2+}$  feedthrough model using the parameters in Table 1.

**TABLE 2 Hill parameters for cytoplasmic  $\text{Ca}^{2+}$ -activation of RyR<sub>2</sub>  $P_o$** 

No.	Voltage mV	$[\text{Ca}^{2+}]_L$ $\mu\text{M}$	$N$	$P_{\text{max}}$	$P_{\text{min}}$	$K_a$ $\mu\text{M}$	$n_a$
2 mM ATP							
1	-40	<10	11	$0.9 \pm 0.04$	$0.003 \pm 0.04$	$0.9 \pm 0.14$	$2 \pm 1.3$
2	-40	100	16	$0.66 \pm 0.18$	$0.20 \pm 0.07$	$0.5 \pm 0.17$	$2 \pm 2$
3	-40	1000	7	$0.65 \pm 0.06$	$0.07 \pm 0.03$	$1.1 \pm 0.16$	$2.4 \pm 2$
4	40	<10	11	$0.8 \pm 0.03$	$0.01 \pm 0.04$	$0.7 \pm 0.12$	$3.0 \pm 2.4$
5	40	100	15	$0.65 \pm 0.05$	$0.01 \pm 0.05$	$0.6 \pm 0.16$	$1.5 \pm 1.0$
6	40	1000	7	$0.85 \pm 0.03$	$0.02 \pm 0.02$	$1.5 \pm 0.25$	$2.2 \pm 0.8$
0 mM ATP							
7	-40	100	13	$0.54 \pm 0.13$	$0 \pm 0.03$	$5.4 \pm 2.0$	$2 \pm 1.9$
8	40	100	13	$0.50 \pm 0.06$	$0 \pm 0.03$	$7.6 \pm 1.7$	$2 \pm 1.8$

Examples of data are shown in Figs. 4 and 8. Parameters are described in Materials and Methods.  $N$  is the number of experiments averaged for each trace.

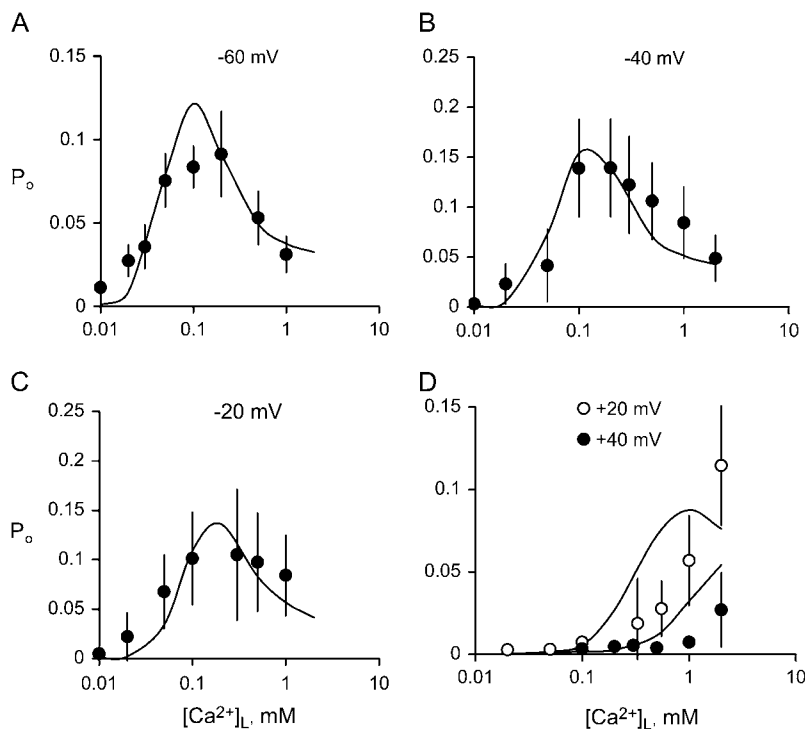
threefold over the experimental voltage range and this could reflect an intrinsic voltage-dependence of the RyR opening mechanism. Taken together, these results indicate that  $[\text{Ca}^{2+}]_L$ -activation of opening rate is mediated by luminal facing sites of action.

The properties of the  $L$ -site are derived from the data in Fig. 6, *B* and *D*, as follows: RyR<sub>2</sub> opening rate associated with the  $L$ -site ( $\alpha'$  in Scheme 1) has an asymptotic dependence on  $[\text{Ca}^{2+}]_L$  (Fig. 6 *B*) that is characteristic of a two-step mechanism [ $(L_C + 2\text{Ca}^{2+}) \leftrightarrow L^*_C \leftrightarrow L^*_O$ ] in which  $\text{Ca}^{2+}$  binding (which is relatively fast) leads to channel opening. In the absence of  $\text{Ca}^{2+}$  feedthrough, the mean duration of these openings ( $\sim 1$  ms) is determined by the closing rate,  $\beta'$ . The Hill coefficient for the  $[\text{Ca}^{2+}]_L$ -dependence of opening rate was 1.6, suggesting a second-order dependence

of  $\alpha'$  on  $[\text{Ca}^{2+}]_L$  in Eq. 2. The voltage-dependence of  $V_{\text{max}}$  is accommodated by imposing a voltage-dependence on  $\alpha'$ .

### Characteristics of the luminal-triggered $\text{Ca}^{2+}$ feedthrough model

The model accounts for the  $[\text{Ca}^{2+}]_L$ ,  $[\text{Ca}^{2+}]_C$ , and voltage-dependencies of  $\tau_o$  and  $\tau_c$  over the range  $-60$  mV to  $+20$  mV (Figs. 4–6, *curves*, the model discrepancy at  $+40$  mV is discussed in the section on model limitations). The model predictions of the  $[\text{Ca}^{2+}]_C$ -dependencies of  $\tau_o$  and  $\tau_c$ , shown in Fig. 4, *B* and *C*, can be understood in terms of the various reaction steps in Scheme 4. In the absence of luminal  $\text{Ca}^{2+}$ ,  $\tau_o$  is primarily determined by the rates  $\delta'_o$  and  $\theta'_o$  of the steps  $LA1_{2C} \leftarrow LA^*I_{2O} \rightarrow LA^*I^*_{2C}$ . These rates have opposite  $[\text{Ca}^{2+}]_C$ -dependencies



**FIGURE 5** The  $[\text{Ca}^{2+}]_L$ -dependence of RyR<sub>2</sub>  $P_o$ . (*A–D*) RyRs were activated at the voltage indicated, by 2 mM ATP and 100 nM  $[\text{Ca}^{2+}]_C$ . Data points show the mean  $\pm$  SE of 3–18 measurements. Solid curves show the fit to the data of the luminal-triggered  $\text{Ca}^{2+}$  feedthrough model using the parameters listed in Table 1.

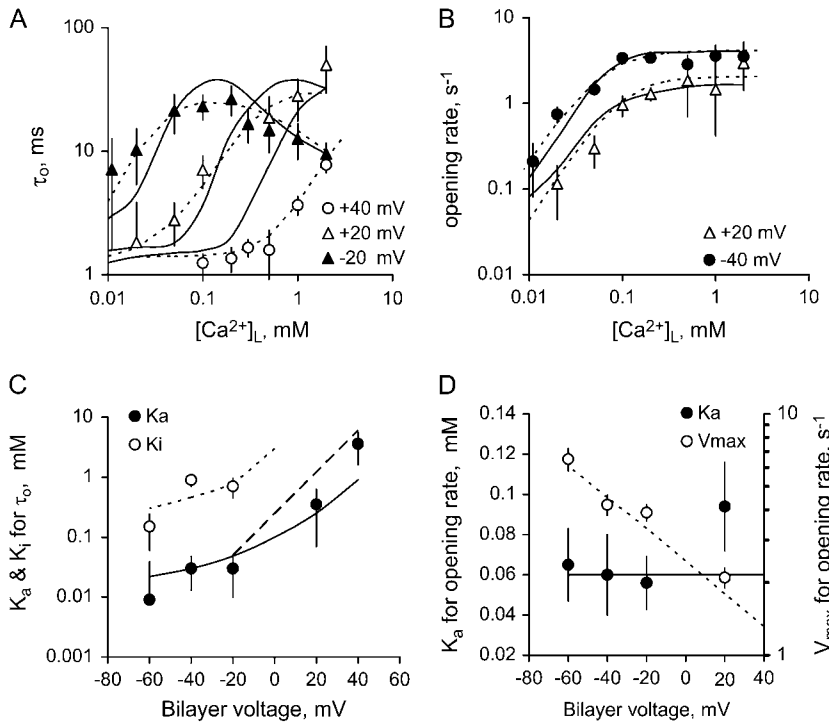


FIGURE 6 The effects of membrane potential on the  $[Ca^{2+}]_L$ -dependent gating of RyR<sub>2</sub>. (A) Mean channel open times at three membrane voltages and (B) channel opening rates at two voltages. Rates were calculated from the inverse of the mean of closed times. The data points show mean  $\pm$  SE of 3–14 measurements. In panels A and B, the dashed curves show Hill fits to the data (parameter values plotted in C and D) and the solid curves show the fits of the luminal-triggered  $Ca^{2+}$  feedthrough model using the parameters listed in Table 1. (C) Half-activating and inactivating  $[Ca^{2+}]_L$ ,  $K_a$  (●), and  $K_i$  (○) for  $\tau_o$ . (D) The  $K_a$  for RyR opening rate (●) and maximum opening rate ( $V_{max}$ , ○). The Hill coefficient for opening rate was  $1.6 \pm 0.3$ . In panels C and D, the solid and dashed curves show the fit to the data of the luminal-triggered  $Ca^{2+}$  feedthrough model. The line (long dashes) shows the voltage-dependence in  $K_a$  expected from translocation of  $Ca^{2+}$  through the membrane voltage (i.e., where  $[Ca^{2+}]_p \propto [Ca^{2+}]_L \exp(-2eV/kT)$ ).

and their effects on the  $[Ca^{2+}]_C$ -dependence of  $\tau_o$  tend to cancel, leaving a small increase in  $\tau_o$  with increasing  $[Ca^{2+}]_C$  (Fig. 4 B, solid curve). When  $[Ca^{2+}]_L = 100 \mu M$  the L-site is mostly bound to  $Ca^{2+}$  so that  $\tau_o$  is primarily determined by the inactivation transitions in  $L^*AI^*_{2C} \leftarrow L^*AI_{2O} \leftrightarrow L^*A^*I_{2O} \rightarrow L^*A^*I^*_{2C}$ . Under these conditions the  $[Ca^{2+}]_C$ -dependent decrease in  $\tau_o$  (Fig. 4 B, short dashes) reflects the  $Ca^{2+}$ -inactivation rate,  $\theta'_o$ . Fig. 4 C shows the model predictions of the  $[Ca^{2+}]_C$ -dependence of  $\tau_c$ . At low  $[Ca^{2+}]_C$   $\tau_c$  is primarily determined by  $\alpha'$  in the reaction step  $LAI_{2C} \rightarrow L^*AI_{2O}$  so that  $\tau_c$  depends on  $[Ca^{2+}]_L$  (compare curves at 10  $\mu M$  and 100  $\mu M$   $[Ca^{2+}]_L$  in Fig. 4 C). Increasing  $[Ca^{2+}]_C$  decreases  $\tau_c$  as  $\gamma'_c$  increasingly determines the channel opening rate via the reaction step  $LAI_{2C} \rightarrow LA^*I_{2O}$ . When  $[Ca^{2+}]_C \geq 3 \mu M$ ,  $\gamma'_c$  becomes very large, then  $\phi'$  in the reactivation step  $LA^*I^*_{2C} \rightarrow LA^*I_{2O}$  determines the lower limit for  $\tau_c$  at high  $[Ca^{2+}]_C$ . When  $[Ca^{2+}]_L = 1000 \mu M$  the  $Ca^{2+}$  feedthrough is large enough to cause saturation of the  $I_2$ -site so that  $[Ca^{2+}]_C$  is unable to modulate inactivation (Fig. 4 B, long/short dashes).

The effects of  $[Ca^{2+}]_L$  and voltage on  $\tau_o$  and  $\tau_c$  are compared with the model in Fig. 6, A and B, respectively. In the absence of  $Ca^{2+}$  feedthrough (low  $[Ca^{2+}]_L$  or large positive voltage),  $\tau_o$  is determined by  $\beta'$  in the reaction  $LAI_{2C} \leftarrow L^*AI_{2O}$ . The contribution of  $Ca^{2+}$  feedthrough is modeled by an extension of the former scheme:  $LAI_{2C} \leftarrow L^*AI_{2O} \leftrightarrow L^*A^*I_{2O} \rightarrow L^*A^*I^*_{2C}$ . Increasing  $Ca^{2+}$  feedthrough biases the state occupancies to the right so that the closing step on the left becomes less frequent. Hence,  $\tau_o$  increases with  $[Ca^{2+}]_L$  until the  $Ca^{2+}$ -dependent closing step on the right

becomes fast enough to decrease  $\tau_o$ . According to the conductance model, the curvature in the voltage-dependencies in the log-plots in Fig. 6 C comes from saturation of  $I_{Ca}$  as the driving force on  $Ca^{2+}$  becomes large. Fig. 6 B shows the corresponding model predictions of channel opening rate at 100 nM  $[Ca^{2+}]_C$ . At these low  $[Ca^{2+}]_C$ , the channel opening rate reflects  $\alpha'$  in the luminal activation step  $LAI_{2C} \rightarrow L^*AI_{2O}$ . The luminal site affinity ( $K_L$ ) determines the  $[Ca^{2+}]_L$  range for the rising phase in the opening rate.

In addition to  $\tau_o$  and  $\tau_c$ , it was found that the model could also account for the effect of  $[Ca^{2+}]_L$  on the shape of the dwell-time distributions in the presence of low  $[Ca^{2+}]_C$ . Fig. 7 shows the distribution of channel open and closed times at three  $[Ca^{2+}]_L$  representing the subactivating (10  $\mu M$ ), maximal-activating (100  $\mu M$ ), and inactivating regions (1 mM) of the bell-shaped dependence of RyR<sub>2</sub> activity on  $[Ca^{2+}]_L$  (−40 mV,  $[Ca^{2+}]_C = 100$  nM). The shape of the closed time distributions varied considerably with  $[Ca^{2+}]_L$  (Fig. 7, A, C, and E). As  $[Ca^{2+}]_L$  increased, the time-constant of the main exponential component decreased from 10 s to 200 ms and the distribution changed from unimodal to bimodal with the appearance of another time constant of  $\sim 1$  ms. Open time distributions (Fig. 7 B, D, and F) were double exponentials as described in Fig. 3. Although the peak of the log-binned open distributions depended on  $[Ca^{2+}]_L$ , the shape of the distributions was similar over the experimental range. The model faithfully reproduced the effect of  $[Ca^{2+}]_L$  on these dwell time distributions (Fig. 7, curves). According to the model, closure of the channel causes termination of  $Ca^{2+}$  feedthrough which, at low  $[Ca^{2+}]_C$ , leads to deactivation of

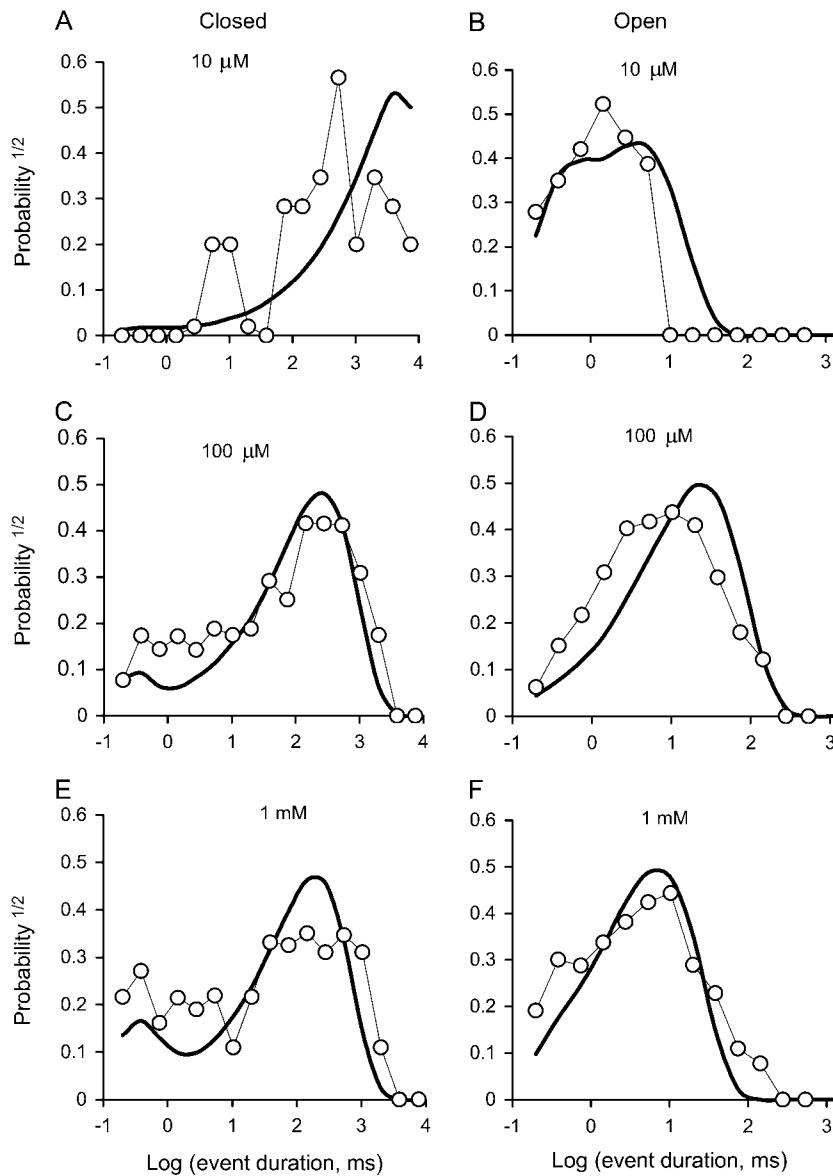


FIGURE 7 The effect of  $[Ca^{2+}]_L$  on closed (A, C, and E) and open (B, D, and F) dwell-time distributions. The data are plotted using the log-bin method of Sigworth and Sine (22). Event duration distributions were compiled from a single RyR, which was activated by cytoplasmic 2 mM ATP (100 nM  $[Ca^{2+}]_C$  and voltage =  $-40$  mV) and the indicated  $[Ca^{2+}]_L$ . The data are plotted using log-bins with 3.5 per decade. Solid curves show simulated distributions generated from the luminal-triggered  $Ca^{2+}$  feedthrough model using the parameters listed in Table 1.

the RyR via  $Ca^{2+}$  dissociation from the A-site, underlying the long closed events. By the same process, channel closures allow  $Ca^{2+}$  dissociation from the  $I_2$ -site, which permits reopening of the channel. The 1-ms time constant in the closed dwell-time distributions is associated with instances of recovery from  $I_2$ -site inactivation that occurred before the channel had time to deactivate.

### The effect of ATP on activation by $[Ca^{2+}]_C$ and $[Ca^{2+}]_L$

The effects of ATP are reexamined here in the light of the new  $Ca^{2+}$  regulation model. Fig. 8 compares the  $[Ca^{2+}]_C$ -dependent properties of RyR<sub>2</sub> in the presence and absence of ATP. ATP increased  $P_o$  over the entire experimental range of  $[Ca^{2+}]_C$  (Fig. 8 A). It increased the maximal activation at

high  $[Ca^{2+}]_C$  ( $P_{max}$ ), decreased the  $[Ca^{2+}]_C$  for half-activation ( $K_a$ ), and altered  $P_o$  at subactivating  $[Ca^{2+}]_C$  ( $P_{min}$ ) depending on  $[Ca^{2+}]_L$  and voltage (Table 2). In accord with previous findings (30), ATP activated RyRs through an increase in  $\tau_o$  and a decrease in  $\tau_c$  (Fig. 8, B and C). In the virtual absence of  $Ca^{2+}$  feedthrough ( $+40$  mV),  $[Ca^{2+}]_C$  caused a much bigger increase in  $\tau_o$  and  $P_o$  in the presence of ATP than in its absence.

Fig. 9 compares the  $[Ca^{2+}]_L$ -dependent properties of RyR<sub>2</sub> in the presence and absence of ATP. Also, as previously reported, ATP markedly amplified the effects of  $[Ca^{2+}]_L$  on  $P_o$  (at low  $[Ca^{2+}]_C$ , Fig. 9 A). In the presence of ATP ( $-40$  mV),  $\tau_o$  showed a bell-shaped dependence on  $[Ca^{2+}]_L$  whereas in the absence of ATP,  $\tau_o$  remained relatively constant (Fig. 9 B). It is interesting to note that ATP had similar effects on the  $[Ca^{2+}]_L$ - and  $[Ca^{2+}]_C$ -dependencies



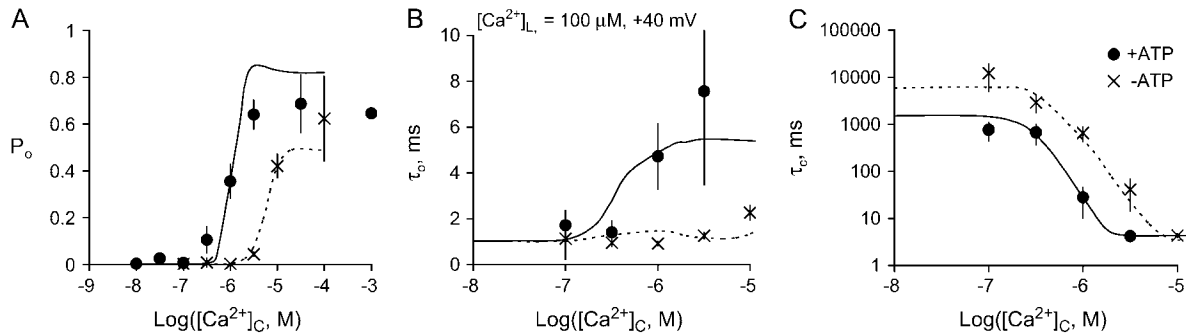


FIGURE 8 The effect of ATP and  $[Ca^{2+}]_C$  on the activity of RyR<sub>2</sub> at +40 mV. The values  $P_o$ ,  $\tau_o$ , and  $\tau_c$  were measured in the presence of 2 mM ATP (●) and in its absence (×). (A) Mean open probability  $P_o$  mean  $\pm$  SE. The numbers of experiments and the Hill parameters are listed in Table 2 (Hill fits to the data are not shown). (B, C) The mean  $\pm$  SE of 3–9 measurements of  $\tau_o$  and  $\tau_c$ . The curves (2 mM ATP, *solid*; and absence of ATP, *dashed*) show the fits of the luminal-triggered  $Ca^{2+}$  feedthrough model using the parameters in Table 1.

of  $\tau_o$  in that, in both cases, an increase in  $Ca^{2+}$  could only increase  $\tau_o$  when ATP was present. This is consistent with common sites of action for cytoplasmic and luminal  $Ca^{2+}$ .  $K_a$  for opening rate was not significantly affected by ATP (Fig. 9 B), indicating that ATP did not modify the affinity of the L-site. Rather, ATP appeared to stabilize the open state of the channel and destabilize the closed state.

In the model, the ability of ATP to increase  $\tau_o$  in response to  $Ca^{2+}$  binding at the A-site is encapsulated in the parameter  $n$  in Eqs. 5 and 7a. In the presence of ATP,  $\tau_o \propto [Ca^{2+}]_C$  (Fig. 8 B, *solid circle*), which is indicative of the channel closing rate,  $\delta \propto [Ca^{2+}]_C^{-1}$  (i.e.,  $n = 1$ ). In the absence of ATP,  $\tau_o$  is nearly independent of  $[Ca^{2+}]_C$  (Fig. 8 B, ×) so that  $n = 0$ . The larger  $\tau_c$  in the absence of ATP (Fig. 8 C, ×) is accommodated in the model by a 13-fold decrease in the opening rate constant associated with the A-site,  $\gamma$  (see Table 1).

The model predictions for the  $[Ca^{2+}]_L$ -dependencies of  $\tau_o$  and  $\tau_c$  in the absence of ATP can be explained along similar lines to those presented above for ATP-activated RyRs. The contribution of  $Ca^{2+}$  feedthrough to  $\tau_o$  depends on the steps  $LAI_{2C} \leftarrow L^*AI_{2O} \leftrightarrow L^*A^*I_{2O} \rightarrow L^*A^*I_{2C}$ . However, in the absence of ATP the states  $L^*A^*I_{2O}$  and  $L^*AI_{2O}$  have

approximately the same average duration. Therefore, although  $Ca^{2+}$  feedthrough biases the state occupancies to the right, in the absence of ATP this causes no significant increase in  $\tau_o$ . Fig. 9 B shows the corresponding model predictions of channel opening rate. The decreased  $[Ca^{2+}]_L$ -dependent opening rate in the absence of ATP was accommodated by a decrease in  $\alpha'$  in the luminal activation step  $LAI_{2C} \rightarrow L^*AI_{2O}$ .

### $Ca^{2+}$ feedthrough couples RyR<sub>2</sub> in lipid bilayers

In six instances, vesicle fusion events incorporated clusters of 3–6 RyR<sub>2</sub> into the bilayer. At subactivating  $[Ca^{2+}]_C$  (100 nM  $Ca^{2+}$  and 2 mM ATP) and when conditions favored  $Ca^{2+}$  feedthrough, the opening of one RyR tended to promote the opening of other RyRs (i.e., channels displayed coupled gating). Fig. 10 shows a recording of six RyRs at positive and negative bilayer potentials. The current trace shows transitions between the current baseline (labeled C for closed) and equally spaced levels corresponding to 1–6 open channels. At –40 mV, channel openings were clearly grouped into bursts and the weighting of current levels in these records markedly deviated from a binomial distribution

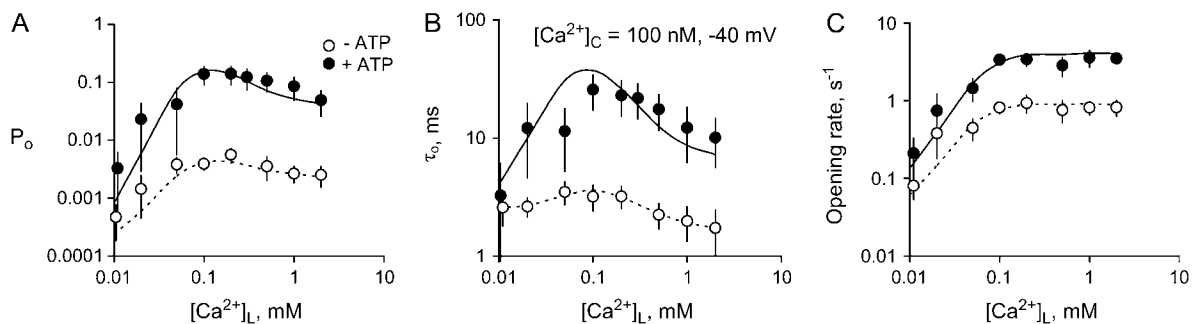


FIGURE 9 The effects of ATP on the  $[Ca^{2+}]_L$ -dependent gating of RyR<sub>2</sub>. (A) The open probability of RyRs (100 nM  $[Ca^{2+}]_C$  and voltage = –40 mV) in the presence of 2 mM ATP (●) and in its absence (○). Also shown are the corresponding  $\tau_o$  (B) and opening rates (C). Solid and dashed curves show the fit to the data of the luminal-triggered  $Ca^{2+}$  feedthrough model using the parameters listed in Table 1. The data points show mean  $\pm$  SE of 3–18 measurements. Hill fits (not shown) to the opening rate reveal that 2 mM ATP increases  $V_{max}$  from  $0.8 \pm 0.1$  to  $4.0 \pm 0.4$  without significantly changing  $K_a$  ( $K_a = 45 \pm 8 \mu M$  and  $60 \pm 20 \mu M$  in the absence and presence of ATP, respectively).

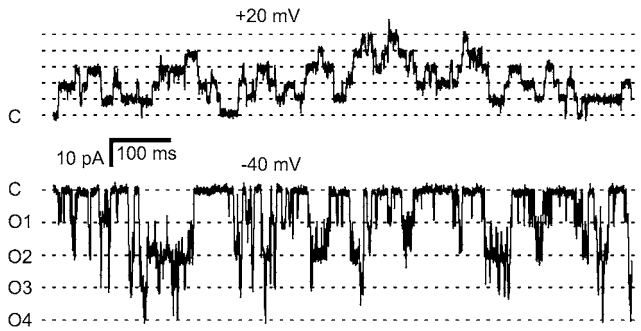


FIGURE 10 Recordings from an experiment with six RyRs in the bilayer showing coupled gating. At +20 mV (*top trace*), the channels appeared to gate independently. The dashed lines indicated the current levels associated with various numbers of open channels. The closed current level is labeled C. At -40 mV (*bottom trace*), the opening of a single channel (transitions from C to O1) was followed closely by openings to higher levels. Only four levels are apparent in this section (labeled O1–O4). The baths contained symmetric 250 mM Cs<sup>+</sup> solutions with 2 mM ATP, [Ca<sup>2+</sup>]<sub>C</sub> = 100 nM, and [Ca<sup>2+</sup>]<sub>L</sub> = 1 mM.

(not shown). Hidden Markov Model analysis of these records (see Materials and Methods) showed that the mean RyR opening rate associated with transitions between the current baseline and level O1 was  $\sim 10\times$  slower than opening rates associated with transitions between higher levels (Fig. 11 A, *solid circle*). Channel closing rates did not depend on the number of open RyRs (Fig. 11 B). Thus, RyR coupling appeared to be mediated by the opening rates. The degree of channel coupling was substantially reduced by decreasing [Ca<sup>2+</sup>]<sub>L</sub> from 1 mM to 0.1 mM or by application of positive bilayer potentials (Fig. 11 A), conditions that oppose Ca<sup>2+</sup> feedthrough. These results are consistent with the notion that the opening of one channel in the bilayer permits Ca<sup>2+</sup> feedthrough that can trigger the opening of neighboring RyRs via their A-sites. The luminal-triggered Ca<sup>2+</sup> feedthrough model was extended to include this possibility by introducing another parameter, Z, which describes the proportionality between  $I_{Ca}$  and its effect on [Ca<sup>2+</sup>]<sub>P</sub> at a neighboring channel. The predicted opening and closing rates for the first opening in a cluster were calculated the same way as for single RyRs. The rates associated with the  $n^{\text{th}}$  channel opening were calculated using a value for

[Ca<sup>2+</sup>]<sub>P</sub> in Eqs. 7b and 10b that incorporated the additional term,  $(n - 1)ZI_{Ca}$ .

The model predictions are compared with the data using a Z value of 0.3. Given the experimental buffering conditions, this corresponds to an interpore separation of  $\sim 3$  nm, which corresponds to the separation of neighboring RyR pores in the triad junction (31). Although the model tends to underestimate the gating rates, it does give a qualitative account of voltage- and [Ca<sup>2+</sup>]<sub>L</sub>-dependent coupling via channel opening rates. In the model, the first channel opening in the cluster is controlled as described for single RyRs (see above). Thus, at 100 nM [Ca<sup>2+</sup>]<sub>C</sub>, 1 mM [Ca<sup>2+</sup>]<sub>L</sub>, and -40 mV, luminal Ca<sup>2+</sup> binding to the L-site is the primary trigger for channel opening whereas Ca<sup>2+</sup> feedthrough serves to reinforce the stability of channel opening. However, the primary trigger for subsequent channel openings is the A-site because feedthrough from the first open channel elevates [Ca<sup>2+</sup>]<sub>P</sub> to 1–3  $\mu$ M. At these [Ca<sup>2+</sup>]<sub>C</sub>, channel opening rates via the A-site are more than 10-fold higher than those attainable via the luminal L-site. In the model, the channel closing rates are slightly increased by the opening of neighboring channels because Ca<sup>2+</sup> feedthrough also contributes to inactivation of neighboring RyRs via I<sub>2</sub>-sites.

### Model limitations

The luminal-triggered Ca<sup>2+</sup> feedthrough model accounts for the Ca<sup>2+</sup> regulation of RyR<sub>2</sub> by the actions of three Ca<sup>2+</sup> binding sites that are linked by Ca<sup>2+</sup> flowing through the pore. However, other Ca<sup>2+</sup> regulation mechanisms could also play a role in controlling Ca<sup>2+</sup> release from the SR. For example, the model does not include the low affinity Ca<sup>2+</sup>/Mg<sup>2+</sup> inhibition mechanism (18). In addition, the [Ca<sup>2+</sup>]<sub>C</sub>-activation mechanism is known to be complex and involves several Ca<sup>2+</sup>-dependent steps (24,25). Simplification of this into A- and I<sub>2</sub>-site gating mechanisms should make the model unsuitable for predicting the complex shape of the dwell-time distributions. With this in mind, it is surprising that the model accounts so well for the [Ca<sup>2+</sup>]<sub>L</sub>-dependent dwell-time distributions at subactivating [Ca<sup>2+</sup>]<sub>C</sub> (Fig. 7). There is also the possibility that the L-site modulates the RyR in other

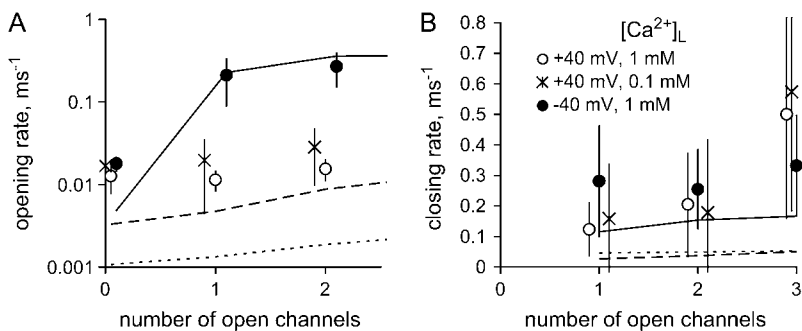


FIGURE 11 The dependence of mean channel opening rate (A) and closing rate (B) on the number of open channels. Rates were measured in the presence of 2 mM ATP and 100 nM [Ca<sup>2+</sup>]<sub>C</sub>. Opening rates were significantly increased ( $P < 0.05$ , *t*-test) by the presence of other open channels under conditions that favored Ca<sup>2+</sup> feedthrough (●, -40 mV [Ca<sup>2+</sup>]<sub>L</sub> = 1 mM, six measurements). When Ca<sup>2+</sup> feedthrough was relatively low, this did not occur (○, +40 mV [Ca<sup>2+</sup>]<sub>L</sub> = 1 mM; ×, -40 mV [Ca<sup>2+</sup>]<sub>L</sub> = 0.1 mM, three measurements each). Closing rates did not significantly depend on the presence of other open channels in the bilayer. The data are compared with predictions of the luminal-triggered Ca<sup>2+</sup> feedthrough model: ●, solid line; ×, long dashes; ○, dashed line.

ways than merely causing channel opening. An example of this has recently been reported in RyR<sub>1</sub> (14) where a luminal Ca<sup>2+</sup> binding site was shown to modulate the A-site affinity for Mg<sup>2+</sup> (preliminary data indicates that this phenomenon does not occur in RyR<sub>2</sub>).

As yet, there has been no direct measure of the Ca<sup>2+</sup> current under experimental conditions. Therefore,  $I_{Ca}$  is estimated from rate theory calculations based on an energy barrier model of the RyR pore (27). A possible failure of the energy barrier model to accurately predict  $I_{Ca}$  under some conditions could underlie the deviation of the model from the data (see the voltage-dependence in  $\tau_o$  that occurred between +20 mV and +40 mV in Fig. 6, A and C). The rate theory model provided a good prediction of the total current in the channel (Cs<sup>+</sup> plus Ca<sup>2+</sup> in this instance, not shown). However, at +40 mV,  $I_{Ca}$  is relatively small and only represents ~1% of the total current. Therefore, it is possible for the model to accurately predict the total current and yet be an order-of-magnitude out in estimating  $I_{Ca}$ . Just a twofold error in  $I_{Ca}$  would account for the deviations between the model and the data at +40 mV. Interestingly,  $\tau_o$  at positive voltages could be better explained by a voltage-dependent equilibrium between Ca<sup>2+</sup> on either end of the pore in which  $[Ca^{2+}]_p \propto [Ca^{2+}]_L \exp(-2eV/kT)$  (Fig. 6 C, long dashes). In any case, the luminal-triggered Ca<sup>2+</sup> feedthrough model provides a good fit to the data between -60 mV and +20 mV, which encompasses the membrane potential of the SR (0 mV, (32)).

The extension of the luminal-triggered Ca<sup>2+</sup> feedthrough model to coupled RyR clusters involved a number of assumptions that are yet to be validated. In the model, coupling occurs between nearest neighbors, although it could also occur between more remote pores. It is not clear that, in the bilayer, RyRs do form the same arrays as they do in muscle. The model assumes that the action of Ca<sup>2+</sup> feedthrough on neighboring channels can be predicted by the steady-state action of raised  $[Ca^{2+}]_C$ . However, neighboring RyRs would be stimulated by a rapid stepwise increase in  $[Ca^{2+}]_C$  and it is known that RyRs can exhibit gating properties that are peculiar to non-steady-state situations (33,34).

## DISCUSSION

The results obtained here support the proposition that RyR activation by Ca<sup>2+</sup> stores is due to Ca<sup>2+</sup> binding sites on both the luminal and cytoplasmic sides of the channel. This study shows that store regulation of RyRs and ER/SR excitability involves three modes of action associated with different parts of the RyR molecule; the L-site, A-site, and I<sub>2</sub>-site (Fig. 1). The L- and I<sub>2</sub>-sites have not been previously identified and this study makes the first measurement of their properties. The effects of  $[Ca^{2+}]_L$  on opening rate in Fig. 6, B and C, highlight the existence of a luminal-facing Ca<sup>2+</sup> activation site (L-site) with an affinity of 60  $\mu$ M. The binding of Ca<sup>2+</sup> to the L-site on its own can activate channel open-

ings of ~1 ms duration at rates of ~1–10 s<sup>-1</sup>. Once the channel is open, the flux of Ca<sup>2+</sup> from the luminal to cytoplasmic sides of the channel increases the cytoplasmic  $[Ca^{2+}]$  near the A-site and produces up to 30-fold prolongation of channel openings. Thus, a large component of the RyR activation by luminal Ca<sup>2+</sup> (up to 97%) is due to the effects of Ca<sup>2+</sup> feedthrough. However, without the luminal site to trigger the initial channel openings, activation by  $[Ca^{2+}]_L$  would not occur. Therefore, the proposed mechanism for store activation of RyRs is luminal-triggered Ca<sup>2+</sup> feedthrough, incorporating both the true luminal and feedthrough hypothesis. This hybrid mechanism would explain the apparently contradictory findings whereby enzyme digestion of luminal RyR domains abolishes luminal Ca<sup>2+</sup> activation (13), but there is a close correlation between RyR activity and Ca<sup>2+</sup> flux through the pore (10).

## Ca<sup>2+</sup>-inactivation

This study makes the first demonstration in single channel recording of a high affinity,  $[Ca^{2+}]_C$ -dependent inactivation in RyRs. It is manifest as a Ca<sup>2+</sup>-dependent reduction in channel  $\tau_o$  (Figs. 2 A and 4 B). The affinity of the I<sub>2</sub>-site is ~1.2  $\mu$ M, which is similar to that of the A-site (0.9  $\mu$ M) so that both activation and inactivation occur over the same range of  $[Ca^{2+}]_C$ .  $[Ca^{2+}]_C$ -activation overrides the effects of the I<sub>2</sub>-site, which causes  $\leq 20\%$  reduction in  $P_o$  in response to cytoplasmic Ca<sup>2+</sup>. However, the I<sub>2</sub>-site causes up to 98% reduction in  $P_o$  in response to feedthrough of luminal Ca<sup>2+</sup> in the presence of physiological (100 nM)  $[Ca^{2+}]_C$ . This is because channel closures block Ca<sup>2+</sup> feedthrough, leading to deactivation via the A-site and substantial lengthening of closed events.

It is well known that cytoplasmic Ca<sup>2+</sup> and Mg<sup>2+</sup> cause identical inhibition of RyRs at high concentrations (~1 mM for RyR<sub>1</sub> and ~10 mM for RyR<sub>2</sub> (35)). This inhibitory action is mediated by a low affinity nonspecific cation site (I<sub>1</sub>-site, previously named the I-site). It was thought that inactivation by high  $[Ca^{2+}]_L$  was mediated by Ca<sup>2+</sup> feedthrough to the I<sub>1</sub>-site (11) because until now, the I<sub>1</sub>-site was the only Ca<sup>2+</sup>-dependent inhibition mechanism that had been clearly identified in RyRs. It is unlikely that the I<sub>1</sub>-site causes inactivation by high  $[Ca^{2+}]_L$  because inactivation has a similar  $[Ca^{2+}]_L$  sensitivity in RyR<sub>1</sub> and RyR<sub>2</sub> (see (11) and (10)) whereas the Ca<sup>2+</sup> sensitivity of the I<sub>1</sub>-site differs 10-fold between the two isoforms. Moreover, the I<sub>1</sub>-site shows no specificity between Ca<sup>2+</sup> and Mg<sup>2+</sup>, whereas the luminal inhibition via the I<sub>2</sub>-site is ~100-fold less sensitive to Mg<sup>2+</sup> than to Ca<sup>2+</sup> (unpublished data).

## Accessibility of A- and I<sub>2</sub>-sites to luminal Ca<sup>2+</sup>

It is noteworthy that even though the A- and I<sub>2</sub>-sites have similar affinity for cytoplasmic Ca<sup>2+</sup>, the effects of luminal

$\text{Ca}^{2+}$  on these sites occur over quite different ranges of  $[\text{Ca}^{2+}]_{\text{L}}$ . The model indicates that the feedthrough effects of  $\text{Ca}^{2+}$  are  $\sim 40$ -fold larger for the *A*-site than for the *I*<sub>2</sub>-site (see *X* and *Y* in Table 1), which suggests a marked difference in the accessibility of these sites to luminal  $\text{Ca}^{2+}$ . This difference is readily explained by the relative proximity of these sites to the  $\text{Ca}^{2+}$  pore (10).  $\text{Ca}^{2+}$  emanating from the pore will diffuse into the cytoplasm and be sequestered by buffering molecules (4.5 mM BAPTA). This leads to a decline in  $[\text{Ca}^{2+}]$  with distance from the pore (see Eq. 1). The luminal-triggered  $\text{Ca}^{2+}$  feedthrough model predicts that each pA of  $\text{Ca}^{2+}$  current through the channel causes a 15  $\mu\text{M}$  increase in  $[\text{Ca}^{2+}]$  at the *A*-site and a 0.35  $\mu\text{M}$  increase at the *I*<sub>2</sub>-site (e.g.,  $I_{\text{Ca}} = 4.9$  pA at  $-40$  mV when  $[\text{Ca}^{2+}]_{\text{L}} = 1$  mM hence  $[\text{Ca}^{2+}]_{\text{p}} = 73$   $\mu\text{M}$  at the *A*-site and 1.7  $\mu\text{M}$  at the *I*<sub>2</sub>-site). This places the *A*- and *I*<sub>2</sub>-sites at 11 nm and 26 nm from the pore, respectively. Given that the furthest point on the RyR from the pore is  $\sim 20$  nm (36), it would seem that the *I*<sub>2</sub>-sites are located at the periphery of the protein or perhaps on an adjacent inhibitory protein. The relative proximity of the *A*- and *I*<sub>2</sub>-sites to the pore is in accord with the observation that  $\text{Ca}^{2+}$  buffering of the cytoplasmic bath affects inactivation much more strongly than activation (10).

### Cytoplasmic agonists and luminal $\text{Ca}^{2+}$ activation

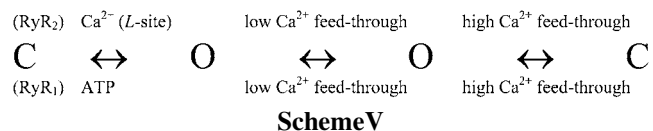
Several studies have shown that  $[\text{Ca}^{2+}]_{\text{L}}$ -dependent activation of RyRs is mainly seen in the presence of agonists such as ATP, caffeine, or sulmazole (8,28,29). Previously it was proposed that these agonists unmask a luminal  $\text{Ca}^{2+}$  sensing site (6). The luminal-triggered  $\text{Ca}^{2+}$  feedthrough model points to a different mechanism in which the *L*-site affinity is not altered by agonists. Rather, they stabilize the channel open conformation, destabilize the closed conformation, and, in conjunction with  $\text{Ca}^{2+}$  feedthrough, markedly increase channel activation in response to  $[\text{Ca}^{2+}]_{\text{L}}$ .

One of these agonists, ATP, is shown here to enhance  $[\text{Ca}^{2+}]_{\text{L}}$ -activation of RyR<sub>2</sub> by three modes of action. Firstly, ATP increases the rate of opening of RyRs that occurs in response to  $\text{Ca}^{2+}$  binding at the *L*-site (Fig. 9 C). Secondly, in accord with previous findings (37), ATP is a cofactor that causes  $\tau_o$  to increase in response to  $\text{Ca}^{2+}$  binding at the *A*-site. Therefore,  $\text{Ca}^{2+}$  feedthrough should increase  $\tau_o$  in the presence of ATP, whereas in the absence of ATP,  $\tau_o$  should not be altered. The results show this to be the case (Fig. 9 B). Thirdly, ATP decreases the rate of inactivation via the *I*<sub>2</sub>-site. In regard to the last two actions, the effects of ATP on the  $[\text{Ca}^{2+}]_{\text{L}}$ -dependence of  $\tau_o$  were entirely consistent with  $\text{Ca}^{2+}$  feedthrough and the observed effects of ATP on the  $[\text{Ca}^{2+}]_{\text{C}}$ -dependence of  $\tau_o$  (Fig. 9, *dashed* and *solid* curves).

Previous studies (30,38–40) have shown that cytoplasmic agonists such as caffeine and ATP have quite different effects on RyR gating than  $\text{Ca}^{2+}$ . Caffeine and ATP produce a much lower  $P_o$  and they generate much longer channel openings and closures than  $\text{Ca}^{2+}$  (Fig. 2 A, see *top trace* showing an RyR

activated by ATP with *bottom trace* of an RyR activated by  $\text{Ca}^{2+}$  and ATP). The different forms of activation by  $\text{Ca}^{2+}$  and ATP can be understood in terms of the luminal-triggered  $\text{Ca}^{2+}$  feedthrough model. Firstly, cytoplasmic  $\text{Ca}^{2+}$  activation via the *A*-site can cause much faster activation rates than ATP, which is limited by the triggering of openings by the *L*-site (hence the longer  $\tau_c$  and the lower  $P_{\text{max}}$  with ATP relative to  $\text{Ca}^{2+}$ ). Secondly, channels activated by cytoplasmic  $\text{Ca}^{2+}$  are more affected by *I*<sub>2</sub>-site mediated inactivation than ATP-activated RyRs (hence  $\tau_o$  is longer ATP than with  $\text{Ca}^{2+}$ ).

Another prediction of the model is that cardiac and skeletal RyR isoforms are differently regulated by luminal  $\text{Ca}^{2+}$ . It is known that there are marked differences in the ways that RyR<sub>1</sub> and RyR<sub>2</sub> are regulated by the cytoplasmic milieu (41). Of particular relevance is that RyR<sub>1</sub> and RyR<sub>2</sub> are differently regulated by cytoplasmic ATP. RyR<sub>1</sub> can be activated by ATP in the absence of cytoplasmic and luminal  $\text{Ca}^{2+}$  (10,14). However, as shown here and elsewhere (28, 42), ATP acts as a cofactor on RyR<sub>2</sub>, which increases its activation in response to  $\text{Ca}^{2+}$  but does not in itself trigger channel openings. This difference in ATP regulation underlies important differences in the way that RyR<sub>1</sub> and RyR<sub>2</sub> respond to luminal  $\text{Ca}^{2+}$  (see Scheme V). It is proposed here that opening of RyR<sub>2</sub> at 100 nM  $[\text{Ca}^{2+}]_{\text{C}}$  is primarily triggered by  $\text{Ca}^{2+}$  binding to the *L*-site whereas  $\text{Ca}^{2+}$  feedthrough serves to reinforce stability of channel opening. However, in RyR<sub>1</sub>, cytoplasmic ATP will trigger channel openings so that the *L*-site is bypassed, resulting in a mechanism that is identical to that previously proposed for RyR<sub>1</sub> (10). It is possible that in the absence of ATP, RyR<sub>1</sub> would default to the RyR<sub>2</sub> mechanism for luminal  $\text{Ca}^{2+}$  regulation provided that RyR<sub>1</sub> also has an *L*-site.



More generally, this model predicts that any cofactor that prolongs channel openings triggered by cytoplasmic  $\text{Ca}^{2+}$  will promote RyR activation by luminal  $\text{Ca}^{2+}$  (the converse will be true for cytoplasmic antagonists). This might be highly relevant to the effects of RyR<sub>2</sub> mutations associated with Sudden Cardiac Death, which are known to enhance activation by luminal  $\text{Ca}^{2+}$  (43) and the ability of polyunsaturated fatty acids, which are RyR antagonists, to protect myocardium against store-overload-induced arrhythmias (44,45). A number of RyR<sub>2</sub> co-proteins have been found to modify  $[\text{Ca}^{2+}]_{\text{L}}$ -regulation of RyRs. The luminal  $\text{Ca}^{2+}$  buffering protein, calsequestrin (CSQ), along with triadin and junctin, are known to bind to RyR<sub>2</sub> and confer luminal  $\text{Ca}^{2+}$  activation, thus being considered as the luminal  $\text{Ca}^{2+}$  sensor (46). However, it is not yet clear if these proteins constitute the *L*-site itself or whether they merely act as a

cofactor that prolongs the channel openings triggered by  $[Ca^{2+}]_C$ . In addition, CSQ dissociation from RyR<sub>2</sub> could be a sensor for Ca<sup>2+</sup> overload of the SR. Exposure of RyR<sub>2</sub> in bilayers to 5 mM  $[Ca^{2+}]_L$  caused a substantial increase in their activity via CSQ dissociation (46). Earlier studies by the same authors found the same treatment caused an increase in the Ca<sup>2+</sup> affinity of the A-site (29). In this study  $[Ca^{2+}]_L$  was kept below 2 mM to minimize the possibility of CSQ dissociation from the RyR. Another cytoplasmic protein, calmodulin (CAM), decreases RyR<sub>2</sub> sensitivity to  $[Ca^{2+}]_C$  via the A-site and appears to modify the action of Ca<sup>2+</sup> feedthrough on RyR<sub>2</sub> activity (47). It is unlikely that CAM contains the A- or I<sub>2</sub>-sites since CAM was not present in this study. CAM is rapidly dissociated from RyRs by the mM  $[Ca^{2+}]$  used to promote SR vesicle fusion with the bilayer.

### Coupled gating of RyR<sub>2</sub> in muscle

The coupled opening of RyRs observed here is very similar to the phenomenon reported previously in RyR<sub>1</sub> and RyR<sub>2</sub> (14,16,17), where it was proposed that coupling occurred when Ca<sup>2+</sup> flow through one channel raised the local  $[Ca^{2+}]_C$  sufficiently to activate neighboring RyRs. This was based on the findings that 1), the coupling only occurred under conditions favoring Ca<sup>2+</sup> feedthrough; 2), coupling did not occur in the presence of a strong cytoplasmic triggering stimulus ( $[Ca^{2+}]_C = 100 \mu\text{M}$ ); and 3) the opening of one channel caused an approximately fivefold reduction in Mg<sup>2+</sup>-inhibition of the other RyRs, suggesting that local  $[Ca^{2+}]_C$  is indeed increased and is competing with Mg<sup>2+</sup> for the A-sites. In this study, RyR<sub>2</sub> coupling was also promoted under conditions favorable to Ca<sup>2+</sup> feedthrough. The coupling of RyR<sub>2</sub> seen here differed slightly from those reported in RyR<sub>1</sub>. There was a much larger relative difference between the first and second opening rates in RyR<sub>2</sub> than RyR<sub>1</sub> (compare 10× for RyR<sub>2</sub> versus 2× for RyR<sub>1</sub>). This probably reflects the different ligand sensitivities of the two isoforms. RyR<sub>1</sub> is more strongly stimulated by ATP than RyR<sub>2</sub> so that the initial opening rate for RyR<sub>1</sub> in a cluster is much higher. However, the rate of subsequent channel openings for RyR<sub>1</sub> and RyR<sub>2</sub> are similar because both isoforms have similar sensitivity to cytoplasmic Ca<sup>2+</sup>.

Ca<sup>2+</sup> sparks are believed to arise from the coupled activation of up to 10 RyRs in the triad junction (15). Freeze fracture electron micrographs show that RyRs within the triad junction are organized into square, two-dimensional arrays (31). The coupling between RyR openings in lipid bilayer experiments suggests that during SR vesicle isolation and reconstitution, fragments of these RyR arrays are retained. The model for luminal-triggered Ca<sup>2+</sup> feedthrough highlights the different mechanisms that might underlie the frequency and morphology of Ca<sup>2+</sup> sparks. Spark frequency is likely to be governed by the opening rate of RyRs in the absence of any open channels, whereas spark morphology would depend on coupling between RyRs. The precise role

of the L-, A-, and I<sub>2</sub>-sites in regulating RyR<sub>2</sub> in the cell is likely to depend on how intracellular Mg<sup>2+</sup> interacts with these sites. The free  $[Mg^{2+}]$  in the cytoplasm and lumen is ~1 mM and this level of Mg<sup>2+</sup> has been shown to lower the apparent Ca<sup>2+</sup> affinity of the A-site to ~10 μM (11) and the L-site to 500 μM (unpublished observations). In the case of the L-site, this places the affinity right in the middle of the physiological range of luminal Ca<sup>2+</sup>.

In vivo, a Ca<sup>2+</sup> spark can occur spontaneously or be initiated by the opening of a single L-type Ca<sup>2+</sup> channel in the sarcolemma (induced sparks) (1). The results here suggest that sparks can be triggered by either binding of luminal Ca<sup>2+</sup> to the L-site (spontaneous sparks) or binding of cytoplasmic Ca<sup>2+</sup> to the A-site (induced and spontaneous sparks). According to the model, spark frequency is governed by the sum of the opening rates associated with L-site and A-sites. Hence spark frequency should increase with SR Ca<sup>2+</sup> load or with increased cytoplasmic Ca<sup>2+</sup> activation. This is consistent with the finding that spark frequency increases with increasing luminal and cytoplasmic stimuli (48,49). In permeabilized cardiac cells, ATP depletion in the presence of 200 nM  $[Ca^{2+}]_C$  was shown to decrease the frequency of spontaneous sparks by >90% even though the Ca<sup>2+</sup> load of the SR was increased (49). This is consistent with bilayer experiments showing that removal of ATP results in a fivefold decrease in RyR opening rate via the L-site (compare  $\alpha \pm$  ATP in Table 1; at 200 nM  $[Ca^{2+}]_C$  the L-site is the main trigger for RyR openings). Interestingly, some RyR activators and inhibitors have only transient effects on spark frequency because any initial change in RyR<sub>2</sub> activity results in an opposing change in SR Ca<sup>2+</sup> load, a process dubbed autoregulation. Why this happens for caffeine and tetracaine (48) and not for ATP (49) is still a mystery. The answer might be found in knowledge of the relative effects of these substances on the L- and A-sites.

According to the model, sparks triggered by either the L- site or A-site should have the same morphology because in either case the subsequent activity of RyRs is governed by the binding of feedthrough Ca<sup>2+</sup> to the A-site and I<sub>2</sub>-sites. The A-sites mediate the coordinated opening of RyRs in a spark and the I<sub>2</sub>-sites govern their termination. Model calculations (not shown) indicate that an I<sub>2</sub>-site inactivation rate of 1000 s<sup>-1</sup> (fourfold higher than measured here) would explain the amplitude distribution and duration of sparks reported by Wang et al. (15). Whether these mechanisms indeed control spark morphology is yet to be determined.

In conclusion, RyRs possess a luminal site for Ca<sup>2+</sup>-activation with an affinity of 60 μM (L-site) and cytoplasmic sites for Ca<sup>2+</sup>-activation and Ca<sup>2+</sup>-inactivation that have ~1 μM affinity (A- and I<sub>2</sub>-sites). Store Ca<sup>2+</sup> regulates RyR activity by binding to the A-, L-, and I<sub>2</sub>-sites and substances that alter channel gating associated with any of these sites will alter the regulation of RyRs by luminal Ca<sup>2+</sup>. A unifying kinetic model is developed that makes the first quantitative predictions of Ca<sup>2+</sup> permeability of the ER/SR. This model provides a framework for understanding the

mode of action of pharmacological agents (e.g., ATP and caffeine), RyR-associated proteins (e.g., calsequestrin), and RyR<sub>2</sub> mutations (e.g., those associated with sudden cardiac death) on a whole range of Ca<sup>2+</sup> mediated physiological and pathological processes.

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