CERTIFIED REDUCED BASIS METHODS FOR PARAMETRIZED SADDLE POINT PROBLEMS

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To my parents.

Abstract

In this thesis, we present a new reduced basis approach to parametrized saddle point problems. The proposed methods allow us to consider flow problems in parametrized domains with notably greater ease.

Designed for the real-time and many-query context of parameter estimation, optimization, and control, the reduced basis method permits the efficient yet reliable approximation of input-output relationships induced by parametrized partial differential equations. The essential ingredients are: (i) dimension reduction, through Galerkin projection onto a lowdimensional reduced basis space; (ii) certainty, through rigorous *a posteriori* bounds for the errors in the reduced basis approximations; (iii) computational efficiency, through an Offline-Online computational strategy; and (iv) effectiveness, through a greedy sampling approach. Although reduced basis methods are well developed for several classes of partial differential equations, parametrized saddle point problems pose additional difficulties that have not been fully addressed: Parameter-dependent constraints cause complications not only in the choice of stable reduced basis approximation spaces, but also in the construction of rigorous and computationally efficient *a posteriori* error bounds.

Based on an analysis of Brezzi's theory for saddle point problems, this work extends and improves existing reduced basis approaches with respect to *both* the construction of effective reduced basis approximation spaces as well as *a posteriori* error bounds. This is achieved as follows: First, we develop new rigorous *a posteriori* error bounds, which perform better in terms of sharpness as well as computational efficiency. Second, we present and analyze several options to construct the reduced basis approximation spaces. Stability is achieved through enriching the reduced basis approximation space for the primal system variable appropriately; different strategies may be applied, favoring either the approximations for the primal or the Lagrange multiplier variables. Third, both *a posteriori* error bounds and enrichment strategies are employed in a new adaptive sampling procedure for constructing reduced basis approximation spaces that are not only stable but also efficient.

Motivated by applications in the field of microfluidics, the methods are applied to a Stokes flow in a parametrized domain. Providing a direct comparison with former reduced basis techniques, numerical results demonstrate the methods' successful performance in practice.

Zusammenfassung

In der vorliegenden Arbeit präsentieren wir einen neuen Zugang zu Reduzierte-Basis-Methoden für parametrisierte Sattelpunktprobleme. Die daraus resultierenden numerischen Verfahren vereinfachen die Analyse von Strömungsproblemen in parameterabhängigen Gebieten erheblich.

Reduzierte-Basis-Methoden ermöglichen eine effiziente und zuverlässige Approximierung von Input-Output Beziehungen, die durch parametrisierte partielle Differentialgleichungen beschrieben werden. Die wesentlichen Bestandteile sind die Folgenden: (i) Dimensionsreduktion, durch Galerkin Projektion auf einen niedrigdimensionalen Approximationsraum; (ii) Gewissheit, durch *a posteriori* Fehlerschranken; (iii) Effizienz, durch eine Offline-Online Rechenstrategie; und (iv) Effektivität, durch ein Greedy-Sampling-Verfahren. Reduzierte-Basis-Methoden sind gut entwickelt für bestimmte Klassen von partiellen Differentialgleichungen. Parametrisierte Sattelpunktprobleme jedoch stellen uns vor zusätzliche Herausforderungen, die noch nicht vollständig behandelt wurden: Parameterabhängige Nebenbedingungen verursachen Komplikationen nicht nur in der Konstruktion von stabilen Approximationsräumen, sondern auch in der Entwicklung effizienter *a posteriori* Fehlerschätzer.

Diese Arbeit erweitert und verbessert bestehende Ansätze basierend auf einer Analyse der Sattelpunkttheorie von F. Brezzi. Dies erreichen wir wie folgt: Erstens entwickeln wir neue *a posteriori* Fehlerschätzer, die sowohl effizienter als auch schärfer sind als bisherige Schranken. Zweitens präsentieren und untersuchen wir verschiedene Optionen zur Konstruktion des reduzierten Approximationsraums. Stabilität wird erzielt durch eine geeignete Anreicherung des Approximationsraums für die primale Systemvariable; verschiedene Strategien stehen hier zur Auswahl, die sich entweder auf die Approximation der primalen Systemvariable oder die des Lagrange-Multiplikators konzentrieren. Drittens stellen wir ein neues adaptives Sampling-Verfahren vor, das angepasst ist an die Bedürfnisse von Sattelpunktproblemen. Die entwickelten Fehlerschätzer und Stabilisierungstechniken werden hier zielführend eingesetzt zur Konstruktion von Approximationsräumen, die nicht nur stabil, sondern auch effizient sind.

Die hergeleiteten Methoden werden angewendet auf ein Stokes Strömungsproblem, welches motiviert ist durch Problemstellungen im Bereich der Mikrofluidik. Numerische Ergebnisse demonstrieren die Leistungsfähigkeit der entwickelten Verfahren durch einen direkten Vergleich mit derzeit bestehenden Techniken.

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Chapter 1

Introduction

1.1 Motivation

Despite impressive progress in mechanical modeling, numerical analysis, and computer science during the last decades, many engineering problems remain intractable. With the increase in speed and memory capacity of computing devices and the improvement of discretization techniques, also new challenges emerged. Models describing physical problems have to face new requirements in terms of accuracy or real-time applicability, resulting in a numerical complexity that is prohibitive for current technologies.

In many engineering problems, the aim is to optimize, control, or characterize a system whose behavior is governed by a partial differential equation (PDE). Such an analysis often requires the prediction of *outputs* in dependence of *inputs*: Outputs are expressed as functionals evaluated at the solution of the governing PDE, whereas inputs are parameters that define a particular system configuration and enter through the PDE. Outputs of interest may be given by a temperature distribution, structural frequencies, or a channel flowrate; inputs often describe material properties, loads and sources, or geometry. For the efficient numerical simulation of such input-output relationships, classical discretization techniques such as finite differences, finite volume, or finite element methods are generally too expensive. First, the context of process control, parametric modeling, or shape optimization often requires the repeated output prediction for many different input parameter values. Second, numerous branches of science and engineering call for an accurate and reliable response in real time. Third, light computing devices such as smartphones or tablets, which may be employed working in the field but are equipped with less computing power, become appealing alternatives to heavy computing platforms. At the present time, these problems cannot be tackled in a direct, traditional way. The reduced basis method, a model order reduction technique, is one means to address this difficulty.

1.2 Model Order Reduction

Reduced order modeling has received considerable attention as it can often capture the system behavior at significantly less cost. The idea is to replace the original model by an approximate model that is much smaller in terms of dimension but still describes important aspects of the system. This involves a number of challenging issues such as (i) the identi-

fication of crucial system properties that should be preserved by the reduced order model, (ii) the extraction of relevant data from the original model that is needed to generate the reduced order model, (iii) the numerically stable computation of the reduced order model, and (iv) the characterization and quantification of its quality. With its own focus and special needs, every discipline tackles these issues differently; as a result, there exist many model order reduction approaches. Among the most well-known techniques for dynamical systems are proper orthogonal decomposition (POD), balanced truncation, and Krylov subspace methods. For a comprehensive overview of these methods, we refer the reader to, e.g., [3, 8, 21, 30, 43, 90] and the references therein; we shall here only summarize the main ideas that allow us to distinguish the reduced basis method.

One of the central issues of POD is the reduction of data expressing their essential information by means of a few basis vectors (see, e.g., [67, 93]). Assume that we are given a set of data in a high-dimensional space, extracted from the original system by either precomputations or experiments. The idea is then to find a subspace of a fixed dimension rsuch that the error in the projection onto the subspace is minimized. This is achieved by applying the singular value decomposition (SVD) to the data matrix: The optimal subspace is obtained as the span of the r orthonormalized eigenvectors associated with the largest singular values; these vectors are also called the *POD modes*.

Balanced truncation and Krylov subspace methods aim to reduce the order of a dynamical system such that the implied input-output relationship is well approximated. The latter is characterized by an input-output map, the so-called *transfer function*. The objective is then to construct a reduced order model with a transfer function resembling the original one as far as possible.

To this end, balanced truncation (see, e.g., [4, 8, 90]) proceeds as follows: As the name suggests, the system is first *balanced* and then *truncated*. Balancing a dynamical system refers to identifying which states of the system are important to its input-output map. To decide whether it affects the system's input-output map, a state is analyzed with respect to its *controllability* and *observability*: Roughly speaking, a state is controllable if it can be reached through an input; a state is observable if it noticeably contributes to an output. Only the states that are both controllable and observable are of relevance. A useful measure of a state's controllability and observability is given by the *controllability* and *observability* gramians. A dynamical system is called balanced if these two gramians coincide and equal a diagonal matrix containing the so-called *Hankel Singular Values (HSVs)*: In this case, states that are difficult to control are simultaneously difficult to observe; quantifying this behavior, the HSVs provide an order of relevance to the input-output map. We may then truncate the state space dimensions with low HSVs to obtain a lower-dimensional, approximate model.

Krylov subspace methods (see, e.g., [6, 30]) do not rely on the computation of singular values but are based on *moment matching*. Noting that the transfer function of a dynamical system can be expanded into a Laurent series, they assume that essential properties of the input-output relationship are reflected in the associated coefficients; these coefficients are also called *moments*. Given a fixed dimension r, the reduced model of corresponding order is therefore constructed by matching as many moments as possible. Since computing the moments explicitly is inherently numerically unstable, only the observation that moment matching is equivalent to projecting the state space onto suitable Krylov subspaces made the approach feasible in practice.

We emphasize that, even though coming from various fields with different objectives, the above model order reduction techniques share a common basic principle: They all can be interpreted as Galerkin or Petrov-Galerkin type projections of the original space onto a low-dimensional subspace. However, different from the others in its particular choice of this subspace (either given through the span of the POD modes, the range of the controllability and observability gramians, or a Krylov subspace), each method has its own benefits and limitations. Based on a very general concept, POD is applicable to various fields and has been successfully used in, e.g., fluid dynamics, signal analysis, data compression, and control theory (see, e.g., [21, 67, 69] and the references therein); depending on the respective context, it is also known as principal component analysis or the Karhunen-Loève decomposition. The method can be applied to high-complexity linear as well as nonlinear systems (see [3]) and is equipped with useful error estimators (see, e.g., [107] and the references therein). However, POD can yield unpredictable results as it is sensitive to the data being used and important system properties such as stability or passivity are generally not preserved; even though the POD modes are optimal at approximating a given data set, they are not necessarily optimal for describing the dynamics of the system that generates this data (see [93]). Designed for the specific needs in optimal control, balanced truncation preserves system properties such as stability and passivity (see [8, 43]). But, to obtain the required gramians and HSVs, it involves the computationally highly expensive solution of two Lyapunov equations followed by an SVD, and is therefore hardly suitable for large-scale problems (see [4, 93]). The latter can be handled efficiently by Krylov subspace methods (see [6, 30]); however, those often provide a good approximation quality only locally and lack the availability of useful error estimates (see [3, 8]).

To overcome each others limitations, the methods are often combined; see, for example, [93, 108] for approaches combining POD and balanced truncation, and [3] for methods that aim to team the best features of SVN based and Krylov subspace methods.

1.3 The Reduced Basis Method

The reduced basis method is a model order reduction approach that permits the efficient yet reliable approximation of input-output relationships induced by parametrized partial differential equations. In contrast to generic discretization techniques where approximation spaces are not correlated to the physical properties of the system, the method recognizes that the solutions to a parametrized PDE are not arbitrary members of the infinite-dimensional solution space but rather reside or evolve on a much lower-dimensional manifold. Exploitation of this low-dimensionality is the key idea of the reduced basis approach.

Designed for the real-time and many-query context of parameter estimation, optimization, and control, the method provides rapidly convergent and computationally efficient approximations equipped with practicable and rigorous error bounds. In compliance with model order reduction techniques outlined in §1.2, this is achieved via Galerkin projection: Built upon a high-fidelity "truth" finite element discretization, the reduced basis approximation is defined as a Galerkin projection onto a low-dimensional subspace that focuses on the solution manifold induced by the parametrized PDE; following the so-called Lagrange approach, this subspace is constructed as the span of several "truth" solutions of the problem at various parameter values. The error in the reduced basis approximation is then measured relative to the "truth" problem formulation and can be quantified by rigorous *a posteriori* error bounds; the latter are typically derived as relaxations of the error-residual equation or can be obtained through direct application of appropriate *a priori* stability estimates. To the method's key features also belongs an Offline-Online computational strategy that enables the highly efficient (Online) computation of both reduced basis approximations and error bounds for any parameter query at the expense of increased pre-processing (Offline) cost. Finally, reduced basis approximations and error bounds are intimately linked through a greedy sampling approach, in which the (Online-)inexpensive error bounds are used to construct the reduced basis approximation spaces more optimally.

A detailed review of the reduced basis method, including its historical background and further references, can be found in [83, 95]; here, we shall restrict ourselves to the stages that help us embed this thesis in its natural context. The method was introduced in the late 1970s as a successful tool for describing large systems in nonlinear structural mechanics with only a few basis functions (see, e.g., [2, 28, 29, 80, 85, 86]). Thenceforth, literature grew rapidly; it soon extended the approach to a wider class of applications and equations, such as fluid dynamics and the incompressible Navier–Stokes equations (see, e.g., [45, 56, 57, 58, 84]). However, the full potential of the method was not recognized until work in [71, 73, 87] entirely decoupled the reduced basis model generation from the projection stage and introduced the concept of computationally inexpensive a *posteriori* error estimators. Offline-Online decomposition strategies gained computational savings of several orders of magnitude. A posteriori error bounds are clearly crucial to the certification of any particular output prediction, but also play an important role in the effective and efficient model generation: Invoked in a greedy sampling process, they allow us to determine (i) which basis functions might be most effective and (ii) how many basis functions are needed to guarantee a desired accuracy. Much current effort is therefore devoted to the development of computationally inexpensive rigorous a posteriori error bounds [53, 55] and optimal sampling procedures [10, 15, 46].

Relevant to many applications such as heat conduction and linear elasticity in transport and continuum mechanics, work in [41, 73, 87, 95] established the reduced basis method for linear elliptic and parabolic PDEs with an affine parameter dependence. Henceforth, the method has been extended in various directions: Much of recent research focuses on its expansion to noncoercive as well as nonlinear equations in fluid dynamics, including the viscous Burgers' [78, 105, 110], the Stokes [92, 96], and the incompressible Navier–Stokes equations [25, 63, 76, 88, 104]. However, it has also been successfully applied to problems in acoustics [106], electromagnetics [22, 60], and stochastics [11, 49]. An empirical interpolation method (EIM) [7, 19, 26, 38, 40, 94] makes the approach feasible for PDEs with a nonaffine parameter dependence; methods in [54, 70, 74] handle hierarchical systems by combining reduced basis with domain decomposition techniques; and work in [24, 39, 61, 77, 103] provides a reduced basis framework for the efficient solution of parametrized optimal control problems.

For the modeling of time-dependent parametrized PDEs, reduced basis methods are often combined with POD techniques (see §1.2). POD based schemes can be applied within the parametric reduced basis context [16, 44]; however, most approaches apply reduced basis greedy sampling strategies in the (multi-dimensional) parameter domain while using POD in the (one-dimensional) temporal domain (see, e.g., [41, 47, 63]). We note that also balanced truncation and Krylov subspace methods (see §1.2) have recently been extended to parameter-dependent systems; however, many of them in their beginnings, the proposed methods often lack practicable error estimates or can handle only very few parameters (see, e.g., [17, 59] and the references therein).

1.4 Contribution of this Thesis

In this thesis, we present a new reduced basis approach to parametrized saddle point problems. The proposed methods allow us to consider flow problems in parametrized domains with notably greater ease.

1.4.1 Theory — Parametrized Saddle Point Problems

Saddle point problems often arise in practical applications where a certain quantity has to be minimized subject to a set of linear constraints; examples of such applications are mixed finite element methods (see, e.g., [13, 89]) or quadratic programming methods in optimal control (see, e.g., [66] and the references therein). In many cases, not only the primal system variable u but also the Lagrange multiplier p associated with the constraints has a physical interpretation and its computation is also of interest.

Although the reduced basis method is well developed for several classes of partial differential equations (see §1.3), parametrized saddle point problems pose additional difficulties that have not been fully addressed: Parameter-dependent constraints cause complications not only in the choice of stable reduced basis approximation spaces, but also in the construction of rigorous and computationally efficient *a posteriori* error bounds. Currently, stable reduced basis approximations are achieved through an enrichment of the approximation space for *u* by so-called *supremizer* functions [72, 92, 96]; resulting in a reduced basis approximation space for *u* that is *at least* twice as large as for *p*, this is not yet satisfactory. Error estimators are often based on Babuška's theory for noncoercive problems [75, 106]; this provides rigorous upper bounds for the combined error in the reduced basis approximations for *u* and *p*, and requires the expensive computation of lower bounds to the Babuška inf-sup stability constants [53, 55].

In this work, we extend and improve existing reduced basis approaches for parametrized saddle point problems with respect to *both* the construction of effective reduced basis approximation spaces as well as a *posteriori* error bounds. This is achieved as follows: First, we develop new rigorous a *posteriori* error bounds that do not involve the estimation of the Babuška inf-sup stability constants but only much less expensive calculations. Second, separate upper bounds for the errors in the approximations for u and p enable the systematic estimation of engineering outputs depending on either of the two. Third, we introduce a new option to construct stable reduced basis approximation spaces that proves to be very effective. Fourth, recognizing the need for stabilization adaptively, a new sampling procedure assembles reduced basis approximation spaces that are very efficient.

1.4.2 Application — Incompressible Fluid Flow

Our focus shall be on saddle point problems that appear in the modeling of incompressible fluid flow. The incompressible Navier–Stokes equations are the governing equations to describe the motion of a viscous fluid (see, e.g., [36, 89]). They generally consist of two equations, the momentum and the continuity equations, which are derived from the basic principles of conservation of momentum and mass in classical mechanics; the unknowns are the velocity field u of the fluid and the pressure p. In case of a creeping flow, where inertial forces are assumed to be small compared to viscous forces, the nonlinear convection term in the Navier–Stokes equations may be neglected and the equations reduce to the linear Stokes equations.

Quite often, it is crucial to analyze and understand the effects of geometric parameters on a flow, in particular when the objective is to optimize, control, or characterize the system. The reduced basis method allows us to perform such an analysis very efficiently. Earlier work on the Stokes and incompressible Navier–Stokes equations has established the method for nonparametrized domains [63, 104], and there have been several efforts [25, 54, 70, 76, 88, 92, 96] at extending the techniques to parametrized domains. However, due to the difficulties explained above (see §1.4.1), rigorous error bounds are often not treated and the geometric variations considered are very small or applicable only to a very limited set of problems.

Motivated by applications in the field of microfluidics, we apply our new methods to a Stokes flow in a two-dimensional microchannel with a parametrized, rectangular obstacle. Providing a direct comparison with former reduced basis techniques, numerical results demonstrate the proposed methods' successful performance in practice.

1.5 Structure of this Thesis

In Chapter 2, we present the abstract framework of a new reduced basis approach based on Brezzi's theory for saddle point problems. We present new *a posteriori* error estimates that provide rigorous upper bounds for the errors in the approximation of the primal system variable u and the Lagrange multiplier *p separately*; symmetric saddle point problems shall be treated as a special case in which these bounds may be further sharpened. Moreover, we analyze several options to construct the reduced basis approximation spaces with respect to approximation stability as well as their effectiveness. Finally, we present a new adaptive sampling procedure serving the needs of our saddle point context.

In Chapter 3, we apply the reduced basis approach developed in §2 to the Stokes equations in a parametrized domain. We consider a Stokes flow in a two-dimensional microchannel with a parametrized, rectangular obstacle; in the field of microfluidics, such microchannels are used to enhance the mixing between two adjacent laminar streams. Numerical results demonstrate the method's performance in practice through a direct comparison with former techniques: Reduced basis *a posteriori* error bounds are compared with respect to sharpness as well as computational cost; reduced basis approximation spaces are compared with respect to stability, effectiveness, and computational efficiency.

Chapter 2 and Chapter 3 are based on our work in [31, 33].

In Chapter 4, we present a reduced basis method based on a penalty approach. The analysis shows how the introduction of a penalty term affects the development of reduced basis *a posteriori* error bounds as well as the construction of reduced basis approximation spaces. At the expense of a perturbed "truth" approximation, we may develop rigorous *a posteriori* error bounds that are computationally highly efficient as they do not depend on inf-sup stability constants; however, they depend on the penalty parameter, and associated effectivities increase as we approach the nonperturbed problem. Furthermore, the penalty term regularizes the problem such that an additional enrichment of the reduced basis approximation space for u is no longer compulsory to obtain stable approximations.

In Chapter 5, we apply the reduced basis approach developed in §4 to a perturbed Stokes flow in a parametrized domain. Considering the model problem introduced in §3, numerical results demonstrate the effects of the penalty term in practice.

Chapter 4 and Chapter 5 are based on our work in [35].

In Chapter 6, we provide an outlook on how the techniques presented in the previous part of this thesis may be extended to the time-dependent setting. To this end, we consider the instationary Stokes equations. The approach essentially combines techniques developed in §2, §3 and §4, §5 with current reduced basis techniques for parabolic problems. The analysis then shows how time integration affects the development of reduced basis *a posteriori* error bounds as well as the construction of efficient reduced basis approximation spaces. The resulting methods are applied to a Stokes flow in a parametrized microchannel where evolution in time is induced by a time-dependent velocity profile on the inflow boundary. Numerical results focus on differences to the stationary case when demonstrating the method's performance in practice. The chapter is based on our work in [32, 34].

Note that each chapter opens with its own introduction, where we outline the chapter's purpose and content as well as its structure; each chapter closes with its own conclusion, where we emphasize particular benefits and limitations of the approach presented.

Chapter 2

Approximation of Saddle Point Problems

2.1 Introduction

In this chapter, we present the abstract framework of a new reduced basis approach based on Brezzi's theory for saddle point problems. We shall focus on two important aspects: First, we present new *a posteriori* error estimates that, unlike earlier approaches based on Babuška's theory for noncoercive problems, provide upper bounds for the errors in the approximations for the primal system variable and the Lagrange multiplier *separately*. The proposed method is a direct application of techniques in [12, 13] to the reduced basis context (see [111] for initial investigations), which shall be further improved in the special case of a symmetric problem. Second, based on an analysis of Brezzi's theory, we present several options to construct the reduced basis approximation spaces that shall be compared with respect to approximation stability as well as their effectiveness. We achieve stability through enriching the reduced basis approximation space for the primal variable appropriately. Different strategies may be applied, favoring either the approximations for the primal or the Lagrange multiplier variables. Finally, both *a posteriori* error bounds and enrichment strategies are employed in an adaptive sampling procedure for constructing reduced basis approximation spaces that are not only stable but also computationally efficient.

The chapter is organized as follows: In §2.2, we introduce the general formulation of a parametrized saddle point problem and its "truth" approximation upon which we shall build our reduced basis approximation. Section 2.3 defines our reduced basis approximation as the Galerkin projection onto a low-dimensional reduced basis approximation space; we shall recall main properties of the reduced basis system as well as *a priori* convergence results. We develop rigorous reduced basis *a posteriori* error bounds in §2.4, and discuss how to construct the reduced basis approximation space in §2.5; computational efficiency is achieved by the Offline-Online strategy summarized in §2.6. In §2.7, a new adaptive sampling procedure then combines error bounds derived in §2.4 and observations in §2.5 to identify reduced basis approximation spaces that are stable as well as efficient. Finally, in §2.8, we give some concluding remarks.

2.2 General Problem Statement

2.2.1 Formulation

Let X_e and Y_e be two Hilbert spaces with inner products $(\cdot, \cdot)_{X_e}$, $(\cdot, \cdot)_{Y_e}$ and associated norms $\|\cdot\|_{X_e} = \sqrt{(\cdot, \cdot)_{X_e}}$, $\|\cdot\|_{Y_e} = \sqrt{(\cdot, \cdot)_{Y_e}}$, respectively.¹ We define the product space $Z_e \equiv X_e \times Y_e$, with inner product $(\cdot, \cdot)_{Z_e} \equiv (\cdot, \cdot)_{X_e} + (\cdot, \cdot)_{Y_e}$ and norm $\|\cdot\|_{Z_e} = \sqrt{(\cdot, \cdot)_{Z_e}}$. The associated dual spaces are denoted by X'_e , Y'_e , and Z'_e .

Furthermore, let $\mathcal{D} \subset \mathbb{R}^n$ be a prescribed *n*-dimensional, compact parameter set. For any parameter $\mu \in \mathcal{D}$, we then consider the continuous bilinear forms $a(\cdot, \cdot; \mu) : X_e \times X_e \to \mathbb{R}$ and $b(\cdot, \cdot; \mu) : X_e \times Y_e \to \mathbb{R}^2$.

$$\gamma_a^{\mathbf{e}}(\mu) \equiv \sup_{u \in X_{\mathbf{e}}} \sup_{v \in X_{\mathbf{e}}} \frac{a(u, v; \mu)}{\|u\|_{X_{\mathbf{e}}} \|v\|_{X_{\mathbf{e}}}} < \infty, \quad \forall \ \mu \in \mathcal{D},$$
(2.1)

$$\gamma_b^{\mathbf{e}}(\mu) \equiv \sup_{q \in Y_{\mathbf{e}}} \sup_{v \in X_{\mathbf{e}}} \frac{b(v, q; \mu)}{\|q\|_{Y_{\mathbf{e}}} \|v\|_{X_{\mathbf{e}}}} < \infty, \quad \forall \ \mu \in \mathcal{D}.$$

$$(2.2)$$

We moreover assume that $a(\cdot, \cdot; \mu)$ is coercive on X_{e} ,

$$\alpha_a^{\mathrm{e}}(\mu) \equiv \inf_{v \in X_{\mathrm{e}}} \frac{a(v, v; \mu)}{\|v\|_{X_{\mathrm{e}}}^2} > 0, \quad \forall \ \mu \in \mathcal{D},$$

$$(2.3)$$

and that $b(\cdot, \cdot; \mu)$ satisfies the inf-sup condition

$$\beta_{\mathrm{Br}}^{\mathrm{e}}(\mu) \equiv \inf_{q \in Y_{\mathrm{e}}} \sup_{v \in X_{\mathrm{e}}} \frac{b(v, q; \mu)}{\|q\|_{Y_{\mathrm{e}}} \|v\|_{X_{\mathrm{e}}}} > 0, \quad \forall \ \mu \in \mathcal{D}.$$

$$(2.4)$$

By (2.1) and (2.3), the bilinear form $a(\cdot, \cdot; \mu)$ provides with $\|\cdot\|_{X_e,\mu} \equiv \sqrt{a(\cdot, \cdot; \mu)}$ an energy norm on X_e that is equivalent to $\|\cdot\|_{X_e}$ for any $\mu \in \mathcal{D}$. We shall refer to $\beta_{\text{Br}}^e(\mu)$ as the "exact Brezzi" inf-sup constant. We additionally define the "exact Babuška" inf-sup constant

$$\beta_{\mathrm{Ba}}^{\mathrm{e}}(\mu) \equiv \inf_{(u,p)\in Z_{\mathrm{e}}} \sup_{(v,q)\in Z_{\mathrm{e}}} \frac{a(u,v;\mu) + b(v,p;\mu) + b(u,q;\mu)}{\|(u,p)\|_{Z_{\mathrm{e}}}\|(v,q)\|_{Z_{\mathrm{e}}}}, \quad \forall \ \mu \in \mathcal{D};$$
(2.5)

it follows that $\beta_{Ba}^{e}(\mu) > 0$ for all $\mu \in \mathcal{D}$ from the assumptions (2.1)–(2.4) (see [5, 12, 27]).

We now consider the following variational problem: For any given $\mu \in \mathcal{D}$, we find $(u_{e}(\mu), p_{e}(\mu)) \in X_{e} \times Y_{e}$ such that

$$a(u_{e}(\mu), v; \mu) + b(v, p_{e}(\mu); \mu) = f(v; \mu), \quad \forall v \in X_{e}, b(u_{e}(\mu), q; \mu) = g(q; \mu), \quad \forall q \in Y_{e},$$
(2.6)

where $f(\cdot; \mu)$ and $g(\cdot; \mu)$ are bounded linear functionals in X'_{e} and Y'_{e} , respectively. From the results of Brezzi [12] (see, e.g., also [13, 27, 36]), it is well-known that under the assumptions (2.1), (2.2), (2.3), and (2.4), the above problem (2.6) is well-posed and has a unique solution for any $f(\cdot; \mu) \in X'_{e}$, $g(\cdot; \mu) \in Y'_{e}$. We recall that if the bilinear form $a(\cdot, \cdot; \mu)$ is in addition

¹Here and in the following, the subscript e denotes "exact".

 $^{^{2}}$ For clarity of exposition, we suppress the obvious requirement of nonzero elements in the denominators.

symmetric for any $\mu \in \mathcal{D}$, the solution $(u_e(\mu), p_e(\mu))$ to (2.6) corresponds to a saddle point of the Lagrangian functional

$$\mathcal{L}(v,q;\mu) \equiv \frac{1}{2}a(v,v;\mu) + b(v,q;\mu) - f(v;\mu) - g(q;\mu), \quad \forall \ (v,q) \in Z_{\mathbf{e}}.$$

Therefore, even though we in fact consider a more general (not necessarily symmetric) case, we call (2.6) a parametrized saddle point problem.

We pause at this point to briefly comment on the nature of the parameter dependence of our linear and bilinear forms. The efficiency of the reduced basis method relies on an Offline-Online computational decomposition strategy (see §2.6) that requires that all bilinear and linear forms in (2.6) depend affinely on the parameter μ . For instance, we assume that for some $Q_a \in \mathbb{N}$, the bilinear form $a(\cdot, \cdot; \mu)$ can be written as

$$a(u, v; \mu) = \sum_{k=1}^{Q_a} \Theta_a^k(\mu) \, a^k(u, v), \quad \forall \, u, v \in X_e,$$
(2.7)

where for $1 \leq k \leq Q_a$, the parameter-dependent coefficient functions $\Theta_a^k(\mu)$ are continuous over the parameter set \mathcal{D} and the parameter-independent bilinear forms $a^k(\cdot, \cdot)$ are continuous on $X_e \times X_e$. We assume analogous representations for the bilinear form $b(\cdot, \cdot; \mu)$ and the linear functionals $f(\cdot; \mu)$ and $g(\cdot; \mu)$.

2.2.2 Truth Approximation

We now introduce a high-fidelity "truth" approximation upon which our reduced basis approximation will subsequently be built. To this end, let X and Y denote finite-dimensional subspaces of X_e and Y_e , respectively. We define the product space $Z \equiv X \times Y$ and denote by \mathcal{N} the dimension of Z. We emphasize that the dimension \mathcal{N} is typically very large. These "truth" approximation subspaces inherit the inner products and norms of the exact spaces: $(\cdot, \cdot)_X \equiv (\cdot, \cdot)_{X_e}, \|\cdot\|_X \equiv \|\cdot\|_{X_e}, (\cdot, \cdot)_Y \equiv (\cdot, \cdot)_{Y_e}, \|\cdot\|_Y \equiv \|\cdot\|_{Y_e}, \text{ and } (\cdot, \cdot)_Z \equiv (\cdot, \cdot)_{Z_e}, \|\cdot\|_Z \equiv \|\cdot\|_{Z_e}.$

The continuity and coercivity properties (2.1), (2.2), and (2.3) are clearly passed on to the "truth" approximation spaces,

$$\gamma_a(\mu) \equiv \sup_{u \in X} \sup_{v \in X} \frac{a(u, v; \mu)}{\|u\|_X \|v\|_X} < \infty, \quad \forall \ \mu \in \mathcal{D},$$
(2.8)

$$\gamma_b(\mu) \equiv \sup_{q \in Y} \sup_{v \in X} \frac{b(v, q; \mu)}{\|q\|_Y \|v\|_X} < \infty, \quad \forall \ \mu \in \mathcal{D},$$
(2.9)

$$\alpha_a(\mu) \equiv \inf_{v \in X} \frac{a(v, v; \mu)}{\|v\|_X^2} > 0, \quad \forall \ \mu \in \mathcal{D};$$

$$(2.10)$$

thus, $\|\cdot\|_{X,\mu} \equiv \|\cdot\|_{X_{e,\mu}}$ defines a norm on X that is equivalent to $\|\cdot\|_X$ for any $\mu \in \mathcal{D}$. We now further assume that the approximation spaces X and Y are chosen such that they satisfy the Ladyzhenskaya–Babuška–Brezzi (LBB) inf-sup condition (see, e.g., [13])

$$\beta_{\mathrm{Br}}(\mu) \equiv \inf_{q \in Y} \sup_{v \in X} \frac{b(v, q; \mu)}{\|q\|_Y \|v\|_X} \ge \beta_{\mathrm{Br}}^0(\mu) > 0, \quad \forall \ \mu \in \mathcal{D},$$
(2.11)

where $\beta_{Br}^{0}(\mu)$ is a constant independent of the dimension \mathcal{N} . We refer to $\beta_{Br}(\mu)$ as the "truth Brezzi" inf-sup constant. We shall also consider the "truth Babuška" inf-sup constant

$$\beta_{\mathrm{Ba}}(\mu) \equiv \inf_{(u,p)\in Z} \sup_{(v,q)\in Z} \frac{a(u,v;\mu) + b(v,p;\mu) + b(u,q;\mu)}{\|(u,p)\|_Z \|(v,q)\|_Z}, \quad \forall \ \mu \in \mathcal{D};$$
(2.12)

again, it follows from (2.8)–(2.11) that $\beta_{\text{Ba}}(\mu) > 0$ for all $\mu \in \mathcal{D}$. In particular, one can show that $\beta_{\text{Br}}(\mu) \ge \beta_{\text{Ba}}(\mu) \ge C(\alpha_a(\mu), \gamma_a(\mu), \beta_{\text{Br}}(\mu)) > 0$ for all $\mu \in \mathcal{D}$; see, e.g., [109] for an explicit representation of $C(\alpha_a(\mu), \gamma_a(\mu), \beta_{\text{Br}}(\mu))$.

We now define our "truth" approximations to be the Galerkin projections of $u_e(\mu) \in X_e$ and $p_e(\mu) \in Y_e$ onto X and Y, respectively: Given any $\mu \in \mathcal{D}$, we find $(u(\mu), p(\mu)) \in X \times Y$ such that

$$a(u(\mu), v; \mu) + b(v, p(\mu); \mu) = f(v; \mu), \quad \forall v \in X, b(u(\mu), q; \mu) = g(q; \mu), \quad \forall q \in Y.$$
(2.13)

As for the exact problem in §2.2.1, it follows from (2.8), (2.9), (2.10), and (2.11) that the "truth" problem (2.13) has a unique solution for any $f(\cdot; \mu) \in X'_{e}$, $g(\cdot; \mu) \in Y'_{e}$. The bilinear forms $a(\cdot, \cdot; \mu)$ and $b(\cdot, \cdot; \mu)$ define bounded linear operators $A(\mu) : X \to X'$, $B(\mu) : X \to Y'$ and its transpose $B(\mu)^{t} : Y \to X'$ by

$$\begin{split} \langle A(\mu) \, u, v \rangle &= a(u, v; \mu), & \forall \, u, v \in X, \\ \langle B(\mu) \, v, q \rangle &= b(v, q; \mu) = \langle B(\mu)^t q, v \rangle, & \forall \, v \in X, \, \forall \, q \in Y; \end{split}$$

here, $\langle \cdot, \cdot \rangle$ denotes the respective dual pairing. The "truth" system (2.13) can thus be equivalently written as

$$A(\mu) u(\mu) + B(\mu)^{t} p(\mu) = f(\mu) \quad \text{in } X', B(\mu) u(\mu) = g(\mu) \quad \text{in } Y',$$

where $f(\mu) \equiv f(\cdot; \mu)|_X \in X'$ and $g(\mu) \equiv g(\cdot; \mu)|_Y \in Y'$ for all $\mu \in \mathcal{D}$.

2.3 Reduced Basis Approximation

We now turn to the reduced basis method, discussing the approximation procedure, rigorous *a posteriori* error estimators, and the construction of *stable* approximation spaces.

2.3.1 Formulation

Let us suppose for now that we are given a set of nested, low-dimensional reduced basis approximation subspaces $X_N \subset X_{N+1} \subset X$ and $Y_N \subset Y_{N+1} \subset Y$, $N \in \mathbb{N}_{\max} \equiv \{1, \ldots, N_{\max}\}$. We denote by N_X and N_Y the dimensions of X_N and Y_N , respectively, and the total dimension of $Z_N \equiv X_N \times Y_N$ by $N_Z \equiv N_X + N_Y$. The subspaces X_N, Y_N , and Z_N again inherit all inner products and norms of X, Y, and Z, respectively. The reduced basis approximation is then defined as the Galerkin projection onto these low-dimensional subspaces: For any given $\mu \in \mathcal{D}$, we find $u_N(\mu) \in X_N$ and $p_N(\mu) \in Y_N$ such that

$$a(u_{N}(\mu), v_{N}; \mu) + b(v_{N}, p_{N}(\mu); \mu) = f(v_{N}; \mu), \quad \forall v_{N} \in X_{N}, b(u_{N}(\mu), q_{N}; \mu) = g(q_{N}; \mu), \quad \forall q_{N} \in Y_{N}.$$
(2.14)

Written in operator notation, the discrete reduced basis system reads

$$A_N(\mu) u_N(\mu) + B_N(\mu)^t p_N(\mu) = f_N(\mu) \quad \text{in } X'_N, \tag{2.15}$$

$$B_N(\mu) u_N(\mu) = g_N(\mu) \quad \text{in } Y'_N,$$
 (2.16)

where $f_N(\mu) \equiv f(\cdot; \mu)|_{X_N} \in X'_N$, $g_N(\mu) \equiv g(\cdot; \mu)|_{Y_N} \in Y'_N$, and the bounded linear operators $A_N(\mu): X_N \to X'_N$, $B_N(\mu): X_N \to Y'_N$ and its transpose $B_N(\mu)^t: Y_N \to X'_N$ are given by

$$\langle A_N(\mu) \, u_N, v_N \rangle = a(u_N, v_N; \mu), \qquad \forall \, u_N, v_N \in X_N, \langle B_N(\mu) \, v_N, q_N \rangle = b(v_N, q_N; \mu) = \langle B_N(\mu)^t \, q_N, v_N \rangle, \quad \forall \, v_N \in X_N, \, \forall \, q_N \in Y_N.$$

2.3.2 Properties of the Discrete System

We first briefly recall some fundamental properties of the discrete system (2.14). Note that we did not yet impose any special requirements on the reduced basis approximation spaces X_N, Y_N . For a given parameter $\mu \in \mathcal{D}$, the system (2.15), (2.16) is solvable if and only if $g_N(\mu)$ belongs to the range of the operator $B_N(\mu)$. In this case, we can construct $u_N(\mu) \in X_N$ satisfying (2.16) such that $f_N(\mu) - A_N(\mu) u_N(\mu)$ belongs to $\operatorname{im}(B_N(\mu)^t)$ and there exists $p_N(\mu) \in Y_N$ satisfying (2.15). The solution $u_N(\mu)$ is unique, $p_N(\mu)$ is in general determined only up to an element of ker $(B_N(\mu)^t)$. The system (2.15), (2.16) is therefore uniquely solvable if and only if $g_N(\mu)$ belongs to the range of the operator $B_N(\mu)$ and the operator $B_N(\mu)^t$ is injective. In the specific case of finite-dimensional spaces X_N and $Y_N, B_N(\mu)^t$ is injective if and only if $B_N(\mu)$ is surjective. Since this clearly implies that $g_N \in \operatorname{im}(B_N(\mu))$ for any $g_N \in Y'_N$, we obtain the following: For a given parameter $\mu \in \mathcal{D}$, the system (2.15), (2.16) is well-posed if and only if $B_N(\mu)$ is surjective. This may be equivalently expressed by the inf-sup condition

$$\beta_N(\mu) \equiv \inf_{q_N \in Y_N} \sup_{v_N \in X_N} \frac{b(v_N, q_N; \mu)}{\|q_N\|_Y \|v_N\|_X} = \inf_{q_N \in Y_N} \frac{\|B_N(\mu)^t q_N\|_{X'_N}}{\|q_N\|_Y} > 0.$$
(2.17)

(For further details on any of the steps above, we refer the reader to [12, 13] and Appendix B.) In the remainder of this work, a pair of approximation spaces (X_N, Y_N) satisfying (2.17) for any parameter value $\mu \in \mathcal{D}$ is called *stable*.

2.3.3 A Priori Error Estimation

We here consider *a priori* estimates for the errors in the reduced basis approximations. These shall also contribute to the construction of reduced basis approximation spaces in §2.5.

In this section, we assume that the low-dimensional reduced basis spaces X_N, Y_N are constructed such that for any given parameter $\mu \in \mathcal{D}$, a solution $(u_N(\mu), p_N(\mu)) \in X_N \times Y_N$ to (2.14) exists. We then denote the errors in the reduced basis approximations $u_N(\mu) \in X_N$, $p_N(\mu) \in Y_N$, and $(u_N(\mu), p_N(\mu)) \in Z_N$ with respect to the truth approximations by

. . .

$$e_N^u(\mu) \equiv u(\mu) - u_N(\mu) \in X,$$

$$e_N^p(\mu) \equiv p(\mu) - p_N(\mu) \in Y,$$

$$e_N(\mu) \equiv (e_N^u(\mu), e_N^p(\mu)) \in Z.$$
(2.18)

Proposition 2.3.1. For any given $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{max}$, we have

$$\|e_N^u(\mu)\|_X \le \left(1 + \frac{\gamma_a(\mu)}{\alpha_a(\mu)}\right) \inf_{\substack{v_N \in X_N \\ B_N(\mu)v_N = g_N(\mu)}} \|u(\mu) - v_N\|_X + \frac{\gamma_b(\mu)}{\alpha_a(\mu)} \inf_{q_N \in Y_N} \|p(\mu) - q_N\|_Y$$
(2.19)

for the error in the reduced basis approximation $u_N(\mu)$; moreover, if the spaces X_N, Y_N are stable (see §2.3.2), we also obtain

$$\|e_N^p(\mu)\|_Y \le \left(1 + \frac{\gamma_b(\mu)}{\beta_N(\mu)}\right) \inf_{q_N \in Y_N} \|p(\mu) - q_N\|_Y + \frac{\gamma_a(\mu)}{\beta_N(\mu)} \|e_N^u(\mu)\|_X$$
(2.20)

for the error in the reduced basis approximation $p_N(\mu)$, and

$$\inf_{\substack{v_N \in X_N \\ B_N(\mu)v_N = g_N(\mu)}} \|u(\mu) - v_N\|_X \le \left(1 + \frac{\gamma_b(\mu)}{\beta_N(\mu)}\right) \inf_{v_N \in X_N} \|u(\mu) - v_N\|_X.$$
(2.21)

Proof. The statement directly follows from arguments presented in [13] applied to the reduced basis context. \Box

In the special case of a symmetric problem, we may also derive the following result.

Proposition 2.3.2. Let $a(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. Then, for any given $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\max}$, we have

$$\|e_N^u(\mu)\|_X \le 2\sqrt{\frac{\gamma_a(\mu)}{\alpha_a(\mu)}} \inf_{\substack{v_N \in X_N \\ B_N(\mu)v_N = g_N(\mu)}} \|u(\mu) - v_N\|_X + \frac{\gamma_b(\mu)}{\alpha_a(\mu)} \inf_{q_N \in Y_N} \|p(\mu) - q_N\|_Y; \quad (2.22)$$

moreover, if the spaces X_N, Y_N are stable (see §2.3.2), we also obtain

$$\|e_{N}^{p}(\mu)\|_{Y} \leq \left[1 + \frac{\gamma_{b}(\mu)}{\beta_{N}(\mu)} \left(1 + \sqrt{\frac{\gamma_{a}(\mu)}{\alpha_{a}(\mu)}}\right)\right] \inf_{q_{N} \in Y_{N}} \|p(\mu) - q_{N}\|_{Y} + 2\frac{\gamma_{a}(\mu)}{\beta_{N}(\mu)} \inf_{\substack{v_{N} \in X_{N} \\ B_{N}(\mu)v_{N} = g_{N}(\mu)}} \|u(\mu) - v_{N}\|_{X}. \quad (2.23)$$

Proof. We here use techniques very similar to those presented in [13]. Take any parameter $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\text{max}}$. We note that, as a continuous, symmetric, and coercive bilinear form, $a(\cdot, \cdot; \mu)$ defines an inner product on X, and recall that the associated norm is denoted by $\|\cdot\|_{X,\mu} = \sqrt{a(\cdot, \cdot; \mu)}$. By the definition of the reduced basis approximation in §2.3.1 as the Galerkin projection of $(u(\mu), p(\mu))$ onto $X_N \times Y_N$, the errors $e_N^u(\mu)$ and $e_N^p(\mu)$ satisfy

$$a(e_N^u(\mu), v_N; \mu) + b(v_N, e_N^p(\mu); \mu) = 0, \quad \forall v_N \in X_N.$$
(2.24)

First, we prove that (2.22) holds true. For any $v_N \in X_N$ such that $B_N(\mu) v_N = g_N(\mu)$ in Y'_N , we have $v_N - u_N(\mu) \in \ker(B_N(\mu))$ and

$$\begin{split} \|v_N - u_N(\mu)\|_{X,\mu} &= \frac{a(v_N - u_N(\mu), v_N - u_N(\mu); \mu)}{\|v_N - u_N(\mu)\|_{X,\mu}} \\ &\leq \sup_{w_N \in \ker(B_N(\mu))} \frac{a(v_N - u_N(\mu), w_N; \mu)}{\|w_N\|_{X,\mu}} \\ &= \sup_{w_N \in \ker(B_N(\mu))} \frac{a(v_N - u(\mu), w_N; \mu) + a(e_N^u(\mu), w_N; \mu)}{\|w_N\|_{X,\mu}} \\ &= \sup_{w_N \in \ker(B_N(\mu))} \frac{a(v_N - u(\mu), w_N; \mu) - b(w_N, e_N^p(\mu); \mu)}{\|w_N\|_{X,\mu}}, \end{split}$$

where the last equality follows from (2.24). For $w_N \in \ker(B_N(\mu))$, $b(w_N, p_N(\mu); \mu) = b(w_N, q_N; \mu) = 0$ holds for all $q_N \in Y_N$. Inserting this in the inequality above yields, for any $q_N \in Y_N$,

$$\|v_N - u_N(\mu)\|_{X,\mu} \le \sup_{w_N \in \ker(B_N(\mu))} \frac{a(v_N - u(\mu), w_N; \mu) - b(w_N, p(\mu) - q_N; \mu)}{\|w_N\|_{X,\mu}} \le \|v_N - u(\mu)\|_{X,\mu} + \frac{\gamma_b(\mu)}{\sqrt{\alpha_a(\mu)}} \|p(\mu) - q_N\|_Y,$$
(2.25)

where the latter is obtained from the Cauchy–Schwarz inequality for the inner product $a(\cdot, \cdot; \mu)$, (2.9), and (2.10). Using the triangle inequality and (2.25),

$$\begin{aligned} \|e_N^u(\mu)\|_{X,\mu} &\leq \|u(\mu) - v_N\|_{X,\mu} + \|v_N - u_N(\mu)\|_{X,\mu} \\ &\leq 2\|u(\mu) - v_N\|_{X,\mu} + \frac{\gamma_b(\mu)}{\sqrt{\alpha_a(\mu)}}\|p(\mu) - q_N\|_Y, \end{aligned}$$
(2.26)

the *a priori* stability estimate (2.22) then follows from (2.26), (2.8), and (2.10).

We now turn to (2.23). Assuming that X_N and Y_N are stable, the inf-sup condition (2.17) provides

$$\beta_N(\mu) \| q_N - p_N(\mu) \|_Y \le \sup_{v_N \in X_N} \frac{b(v_N, q_N - p_N(\mu); \mu)}{\| v_N \|_X}, \quad \forall q_N \in Y_N.$$
(2.27)

For any $v_N \in X_N$ and $q_N \in Y_N$, we moreover have

$$b(v_N, q_N - p_N(\mu); \mu) = b(v_N, q_N - p(\mu); \mu) + b(v_N, e_N^p(\mu); \mu)$$

= $b(v_N, q_N - p(\mu); \mu) - a(e_N^u(\mu), v_N; \mu),$

where the last equality follows from (2.24). Applying this to (2.27), together with the Cauchy–Schwarz inequality for $a(\cdot, \cdot; \mu)$ and (2.8), we obtain

$$\|q_N - p_N(\mu)\|_Y \le \frac{\gamma_b(\mu)}{\beta_N(\mu)} \|p(\mu) - q_N\|_Y + \frac{\sqrt{\gamma_a(\mu)}}{\beta_N(\mu)} \|e_N^u(\mu)\|_{X,\mu}, \quad \forall \ q_N \in Y_N;$$
(2.28)

the *a priori* error estimate (2.23) thus holds again from the triangle inequality, (2.28), (2.26), and (2.8).

2.4 A Posteriori Error Estimation

We now aim to develop not only efficient reduced order approximations, but also a *posteriori* error estimates that are rigorous, sharp, and computationally inexpensive.

In this section, we again assume that the low-dimensional reduced basis approximation spaces X_N, Y_N are constructed such that for any given parameter $\mu \in \mathcal{D}$, a solution $(u_N(\mu), p_N(\mu)) \in X_N \times Y_N$ to (2.14) exists. We may then derive upper bounds as required for the errors $e_N^u(\mu)$ and $e_N^p(\mu)$ in the reduced basis approximations $u_N(\mu)$ and $p_N(\mu)$, respectively. Symmetric saddle point problems shall be discussed as a special case in which these bounds can be further sharpened.

2.4.1 Ingredients

To formulate rigorous and computationally inexpensive upper bounds for the respective errors defined in (2.18), we first have to introduce further ingredients. The first set of ingredients consists of computationally (Online-)efficient lower and upper bounds to the truth continuity and coercivity constants (2.8) and (2.10),

$$\gamma_{a}^{\mathrm{LB}}(\mu) \leq \gamma_{a}(\mu) \leq \gamma_{a}^{\mathrm{UB}}(\mu), \alpha_{a}^{\mathrm{LB}}(\mu) \leq \alpha_{a}(\mu) \leq \alpha_{a}^{\mathrm{UB}}(\mu), \quad \forall \ \mu \in \mathcal{D},$$

$$(2.29)$$

and to the truth Brezzi and Babuška inf-sup constants (2.11) and (2.12),

$$\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu) \leq \beta_{\mathrm{Br}}(\mu) \leq \beta_{\mathrm{Br}}^{\mathrm{UB}}(\mu), \beta_{\mathrm{Ba}}^{\mathrm{LB}}(\mu) \leq \beta_{\mathrm{Ba}}(\mu) \leq \beta_{\mathrm{Ba}}^{\mathrm{UB}}(\mu), \quad \forall \ \mu \in \mathcal{D}.$$

$$(2.30)$$

The second set of ingredients consists of dual norms of the residuals associated with the reduced basis approximation,

$$\|r_N^1(\cdot;\mu)\|_{X'} = \sup_{v \in X} \frac{r_N^1(v;\mu)}{\|v\|_X}, \quad \|r_N^2(\cdot;\mu)\|_{Y'} = \sup_{q \in Y} \frac{r_N^2(q;\mu)}{\|q\|_Y}, \tag{2.31}$$

where, for all $\mu \in \mathcal{D}$, $r_N^1(\cdot; \mu) \in X'$ and $r_N^2(\cdot; \mu) \in Y'$ are defined as

$$r_N^1(v;\mu) \equiv f(v;\mu) - a(u_N(\mu),v;\mu) - b(v,p_N(\mu);\mu), \quad \forall v \in X,$$
(2.32)

$$r_N^2(q;\mu) \equiv g(q;\mu) - b(u_N(\mu),q;\mu), \qquad \forall q \in Y.$$
(2.33)

The total residual is then given by $r_N((v,q);\mu) \equiv r_N^1(v;\mu) + r_N^2(q;\mu)$ for all $(v,q) \in \mathbb{Z}$, with its dual norm (see Appendix B)

$$\|r_N(\cdot;\mu)\|_{Z'} = \sup_{(v,q)\in\mathbb{Z}} \frac{r_N((v,q);\mu)}{\|(v,q)\|_Z} = \sqrt{\|r_N^1(\cdot;\mu)\|_{X'}^2 + \|r_N^2(\cdot;\mu)\|_{Y'}^2}, \quad \forall \ \mu \in \mathcal{D}.$$
(2.34)

2.4.2 Error Bound Formulation

We may now formulate a posteriori error bounds for the respective errors (2.18) in the reduced basis approximation. The following proposition represents a well-known result for noncoercive problems [72, 76, 79, 106].

Proposition 2.4.1. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, and $\beta_{\text{Ba}}^{\text{LB}}(\mu)$ satisfying (2.30), we define

$$\Delta_N^{\mathrm{Ba}}(\mu) \equiv \frac{\|r_N(\cdot;\mu)\|_{Z'}}{\beta_{\mathrm{Ba}}^{\mathrm{LB}}(\mu)}.$$
(2.35)

Then, $\Delta_N^{\text{Ba}}(\mu)$ is an upper bound for the error $e_N(\mu)$ such that

$$||e_N(\mu)||_Z \le \Delta_N^{\text{Ba}}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\text{max}}.$$
(2.36)

Proof. The above error bound directly follows from the Banach–Nečas–Babuška theorem [5, 27] and (2.30).

Since many engineering outputs of interest such as average vorticity or pressure drop depend on either the primal variable (velocity) $u(\mu)$ or the Lagrange multiplier (pressure) $p(\mu)$, we aim to develop *separate* bounds for the errors in the reduced basis approximations $u_N(\mu)$ and $p_N(\mu)$. Noting that a *posteriori* reduced basis error bounds for coercive and general noncoercive problems can be derived from standard stability results, we apply Brezzi's theory for saddle point problems [12] to our problem setting. Since some aspects of this theory are important also to the construction of reduced basis approximation spaces in §2.5 and §2.7, we briefly summarize the proof below.

Proposition 2.4.2. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, and $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, $\beta_{\text{Br}}^{\text{LB}}(\mu)$ satisfying (2.29) and (2.30), we define

$$\Delta_{N}^{u,\mathrm{Br}}(\mu) \equiv \frac{\|r_{N}^{1}(\cdot;\mu)\|_{X'}}{\alpha_{a}^{\mathrm{LB}}(\mu)} + \left(1 + \frac{\gamma_{a}^{\mathrm{UB}}(\mu)}{\alpha_{a}^{\mathrm{LB}}(\mu)}\right) \frac{\|r_{N}^{2}(\cdot;\mu)\|_{Y'}}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)},\tag{2.37}$$

$$\Delta_N^{p,\mathrm{Br}}(\mu) \equiv \frac{\|r_N^1(\cdot;\mu)\|_{X'}}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)} + \frac{\gamma_a^{\mathrm{UB}}(\mu)}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)} \Delta_N^{u,\mathrm{Br}}(\mu).$$
(2.38)

Then, $\Delta_N^{u,\mathrm{Br}}(\mu)$ and $\Delta_N^{p,\mathrm{Br}}(\mu)$ are upper bounds for the errors $e_N^u(\mu)$ and $e_N^p(\mu)$ such that

$$\|e_N^u(\mu)\|_X \le \Delta_N^{u,\mathrm{Br}}(\mu), \quad \|e_N^p(\mu)\|_Y \le \Delta_N^{p,\mathrm{Br}}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\mathrm{max}}.$$
 (2.39)

Proof. Let $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\text{max}}$. By (2.32), (2.33), and (2.13), the errors $e_N^u(\mu) \in X$ and $e_N^p(\mu) \in Y$ satisfy the equations

$$a(e_N^u(\mu), v; \mu) + b(v, e_N^p(\mu); \mu) = r_N^1(v; \mu), \quad \forall v \in X,$$
(2.40)

$$b(e_N^u(\mu), q; \mu) = r_N^2(q; \mu), \quad \forall q \in Y.$$
 (2.41)

We first consider the error $e_N^u(\mu) \in X$ in the approximation of the primal variable. This may be uniquely decomposed into $e_N^u(\mu) = e_N^0(\mu) + e_N^{\perp}(\mu)$, where $e_N^0(\mu) \in \ker(B(\mu))$ and $e_N^{\perp}(\mu) \in \ker(B(\mu))^{\perp} \equiv \{ v \in X \mid (v, v_0)_X = 0 \text{ for all } v_0 \in \ker(B(\mu)) \}$. From (2.40), $e_N^0(\mu) \in \ker(B(\mu))$ then solves

$$a(e_N^0(\mu), v_0; \mu) = r_N^1(v_0; \mu) - a(e_N^{\perp}(\mu), v_0; \mu), \quad \forall \ v_0 \in \ker(B(\mu)),$$

and it is thus bounded by

$$\|e_N^0(\mu)\|_X \le \frac{1}{\alpha_a(\mu)} \sup_{v_0 \in \ker(B(\mu))} \frac{r_N^1(v_0;\mu) - a(e_N^{\perp}(\mu), v_0;\mu)}{\|v_0\|_X}$$
(2.42)

$$\leq \frac{1}{\alpha_a(\mu)} \Big(\|r_N^1(\cdot;\mu)\|_{X'} + \gamma_a(\mu)\|e_N^{\perp}(\mu)\|_X \Big),$$
(2.43)

using (2.8), (2.10), and the classical Lax–Milgram lemma (see, e.g., [27, 89]). By the LBB inf-sup condition (2.11), it is also true that (see, e.g., [12], [13, §II.1, Proposition 1.2])

$$\beta_{\rm Br}(\mu) = \inf_{v \in \ker(B(\mu))^{\perp}} \sup_{q \in Y} \frac{b(v, q; \mu)}{\|v\|_X \|q\|_Y}.$$
(2.44)

Applied to $e_N^{\perp}(\mu) \in \ker(B(\mu))^{\perp}$, this yields

$$\|e_{N}^{\perp}(\mu)\|_{X} \leq \frac{1}{\beta_{\mathrm{Br}}(\mu)} \sup_{q \in Y} \frac{b(e_{N}^{\perp}(\mu), q; \mu)}{\|q\|_{Y}} = \frac{\|r_{N}^{2}(\cdot; \mu)\|_{Y'}}{\beta_{\mathrm{Br}}(\mu)},$$
(2.45)

where the equality follows from (2.41) as $B(\mu) e_N^{\perp}(\mu) = B(\mu) e_N^u(\mu)$ in Y'. We then obtain the bound (2.37) for $||e_N^u(\mu)||_X$ by combining (2.43), (2.45) and (2.29), (2.30).

We now turn to the error in the approximation of the Lagrange multiplier. From (2.11) and (2.40), we derive that

$$\|e_{N}^{p}(\mu)\|_{Y} \leq \frac{1}{\beta_{\mathrm{Br}}(\mu)} \sup_{v \in X} \frac{b(v, e_{N}^{p}(\mu); \mu)}{\|v\|_{X}} = \frac{1}{\beta_{\mathrm{Br}}(\mu)} \sup_{v \in X} \frac{r_{N}^{1}(v; \mu) - a(e_{N}^{u}(\mu), v; \mu)}{\|v\|_{X}}$$
$$\leq \frac{1}{\beta_{\mathrm{Br}}(\mu)} \Big(\|r_{N}^{1}(\cdot; \mu)\|_{X'} + \gamma_{a}(\mu) \|e_{N}^{u}(\mu)\|_{X} \Big).$$
(2.46)

Together with (2.37) and again (2.29), (2.30), this finally yields (2.38) and (2.39).

Clearly, we can now construct another rigorous upper bound $\Delta_N^{\text{Br}}(\mu)$ for the combined error $e_N(\mu) = (e_N^u(\mu), e_N^p(\mu))$: From (2.39), we obtain

$$\|e_N(\mu)\|_Z \le \sqrt{\left(\Delta_N^{u,\mathrm{Br}}(\mu)\right)^2 + \left(\Delta_N^{p,\mathrm{Br}}(\mu)\right)^2} \equiv \Delta_N^{\mathrm{Br}}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\mathrm{max}}.$$
 (2.47)

As a measure of the quality of the proposed error estimators, we introduce the associated effectivities

$$\eta_N^{\mathrm{Ba}}(\mu) \equiv \frac{\Delta_N^{\mathrm{Ba}}(\mu)}{\|e_N(\mu)\|_Z}, \quad \eta_N^{\mathrm{Br}}(\mu) \equiv \frac{\Delta_N^{\mathrm{Br}}(\mu)}{\|e_N(\mu)\|_Z}, \quad \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\mathrm{max}}.$$
 (2.48)

Effectivities ≥ 1 thus indicate that the error estimate is *rigorous*; effectivities close to unity indicate that the error estimate is *sharp*. The following corollary specifies this behavior for the error bounds derived above.

Corollary 2.4.3. For any given $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\max}$, the effectivities of the a posteriori error bounds (2.35) and (2.47) satisfy

$$1 \le \eta_N^{\mathrm{Ba}}(\mu) \le \frac{\gamma_a(\mu) + \gamma_b(\mu)}{\beta_{\mathrm{Ba}}^{\mathrm{LB}}(\mu)}, \quad 1 \le \eta_N^{\mathrm{Br}}(\mu) \le C_1(\mu) \big(\gamma_a(\mu) + \gamma_b(\mu)\big), \tag{2.49}$$

where $C_1(\mu) > 0$ is a constant depending on $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, and $\beta_{\text{Br}}^{\text{LB}}(\mu)$.

Proof. Let μ be any parameter in \mathcal{D} and $N \in \mathbb{N}_{\max}$. The error bounds are rigorous by (2.36) and (2.47), and it only remains to show the validity of the upper bounds for the effectivities.

First, we consider the effectivities associated with $\Delta_N^{\text{Ba}}(\mu)$. By (2.40), (2.41) and (2.8), (2.9), the total residual satisfies

$$\begin{aligned} r_N((v,q);\mu) &= r_N^1(v;\mu) + r_N^2(q;\mu) = a(e_N^u(\mu),v;\mu) + b(v,e_N^p(\mu);\mu) + b(e_N^u(\mu),q;\mu) \\ &\leq \gamma_a(\mu) \|e_N^u(\mu)\|_X \|v\|_X + \gamma_b(\mu) \|v\|_X \|e_N^p(\mu)\|_Y + \gamma_b(\mu) \|e_N^u(\mu)\|_X \|q\|_Y \\ &= \begin{pmatrix} \|v\|_X \\ \|q\|_Y \end{pmatrix}^t \begin{pmatrix} \gamma_a(\mu) \ \gamma_b(\mu) \\ \gamma_b(\mu) \ 0 \end{pmatrix} \begin{pmatrix} \|e_N^u(\mu)\|_X \\ \|e_N^p(\mu)\|_Y \end{pmatrix}, \quad \forall \ (v,q) \in Z. \end{aligned}$$

This may be estimated in terms of the Euclidean norm $\|\cdot\|_2$, i.e.,

$$r_{N}((v,q);\mu) \leq \|(v,q)\|_{Z} \left\| \begin{pmatrix} \gamma_{a}(\mu) \ \gamma_{b}(\mu) \\ \gamma_{b}(\mu) \ 0 \end{pmatrix} \right\|_{2} \|(e_{N}^{u}(\mu),e_{N}^{p}(\mu))\|_{Z} \\ \leq (\gamma_{a}(\mu) + \gamma_{b}(\mu)) \|e_{N}(\mu)\|_{Z} \|(v,q)\|_{Z}, \quad \forall (v,q) \in Z,$$

where the last inequality follows from the relation $\|\cdot\|_2 \leq \sqrt{\|\cdot\|_1\|\cdot\|_\infty}$ between matrix norms (see Appendix A). Thus, the first part of (2.49) holds from the definition of $\Delta_N^{\text{Ba}}(\mu)$ in (2.35).

We now turn to the effectivities associated with $\Delta_N^{\text{Br}}(\mu)$. Again from (2.40), (2.41), and (2.8), (2.9), it is clearly also true that

$$\|r_N^1(\cdot;\mu)\|_{X'} \le \gamma_a(\mu) \|e_N^u(\mu)\|_X + \gamma_b(\mu) \|e_N^p(\mu)\|_Y,$$
(2.50)

$$\|r_N^2(\cdot;\mu)\|_{Y'} \le \gamma_b(\mu) \|e_N^u(\mu)\|_X.$$
(2.51)

Both residual dual norms (using either the Cauchy–Schwarz inequality or by adding terms) are then particularly bounded by

$$\left(\gamma_a(\mu)^2 + \gamma_b(\mu)^2\right)^{1/2} \|e_N(\mu)\|_Z \le \left(\gamma_a(\mu) + \gamma_b(\mu)\right) \|e_N(\mu)\|_Z,$$

and we obtain for $\Delta_N^{u,\mathrm{Br}}(\mu)$ and $\Delta_N^{p,\mathrm{Br}}(\mu)$ that

$$\Delta_{N}^{u,\mathrm{Br}}(\mu) \leq \left[\frac{1}{\alpha_{a}^{\mathrm{LB}}(\mu)} + \left(1 + \frac{\gamma_{a}^{\mathrm{UB}}(\mu)}{\alpha_{a}^{\mathrm{LB}}(\mu)}\right) \frac{1}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)}\right] \left(\gamma_{a}(\mu) + \gamma_{b}(\mu)\right) \|e_{N}(\mu)\|_{Z},$$

$$\Delta_{N}^{p,\mathrm{Br}}(\mu) \leq \frac{1}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)} \left(1 + \frac{\gamma_{a}^{\mathrm{UB}}(\mu)}{\alpha_{a}^{\mathrm{LB}}(\mu)}\right) \left(1 + \frac{\gamma_{a}^{\mathrm{UB}}(\mu)}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)}\right) \left(\gamma_{a}(\mu) + \gamma_{b}(\mu)\right) \|e_{N}(\mu)\|_{Z};$$

the statement (2.49) now directly follows from the definition of $\Delta_N^{\text{Br}}(\mu)$ in (2.47).

Remark 2.4.4. We here make some additional observations on the *a posteriori* error bounds given in Proposition 2.4.2 that shall be important also to the construction of reduced basis approximation spaces in §2.5.

It is clear from Corollary 2.4.3 that the residual dual norms (2.31) provide an *a posteriori* error estimator for $||e_N(\mu)||_Z$ that is not only reliable but also efficient (in the sense that

the terms are used in the finite element community, i.e., bounding the error up to some constants from above and from below). Now, for any $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\text{max}}$, it follows from (2.42), (2.45), and (2.8) that the error $||e_N^u(\mu)||_X$ is in fact bounded by

$$\|e_N^u(\mu)\|_X \le \frac{1}{\alpha_a(\mu)} \sup_{v_0 \in \ker(B(\mu))} \frac{r_N^1(v_0;\mu)}{\|v_0\|_X} + \left(1 + \frac{\gamma_a(\mu)}{\alpha_a(\mu)}\right) \frac{\|r_N^2(\cdot;\mu)\|_{Y'}}{\beta_{\mathrm{Br}}(\mu)};$$
(2.52)

additionally using (2.51) and

$$\sup_{v_0 \in \ker(B(\mu))} \frac{r_N^1(v_0;\mu)}{\|v_0\|_X} = \sup_{v_0 \in \ker(B(\mu))} \frac{a(e_N^u(\mu), v_0;\mu)}{\|v_0\|_X} \le \gamma_a(\mu) \|e_N^u(\mu)\|_X,$$
(2.53)

we thus obtain a reliable and efficient a posteriori error estimator only for $||e_N^u(\mu)||_X$. \Box

Remark 2.4.5. Following techniques used in [48] to obtain a priori stability estimates for the reduced basis approximations, we can derive an upper bound for the error in $u_N(\mu)$ that is similar to (2.37) but slightly differs in the weighting of the residual dual norms: For any $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{max}$,

$$\frac{1}{2\alpha_{a}^{\text{LB}}(\mu)} \left(\|r_{N}^{1}(\cdot;\mu)\|_{X'} + \frac{\gamma_{a}^{\text{UB}}(\mu)}{\beta_{\text{Br}}^{\text{LB}}(\mu)} \|r_{N}^{2}(\cdot;\mu)\|_{Y'} \right) + \sqrt{\left[\frac{1}{2\alpha_{a}^{\text{LB}}(\mu)} \left(\|r_{N}^{1}(\cdot;\mu)\|_{X'} + \frac{\gamma_{a}^{\text{UB}}(\mu)}{\beta_{\text{Br}}^{\text{LB}}(\mu)} \|r_{N}^{2}(\cdot;\mu)\|_{Y'} \right)\right]^{2} + \frac{\|r_{N}^{2}(\cdot;\mu)\|_{Y'} \|r_{N}^{1}(\cdot;\mu)\|_{X'}}{\alpha_{a}^{\text{LB}}(\mu)\beta_{\text{Br}}^{\text{LB}}(\mu)} \quad (2.54)$$

is a rigorous upper bound for $||e_N^u(\mu)||_X$. The proof follows ideas in [48] and we only outline the main steps. From (2.10), (2.40), and (2.41), we have

$$\begin{aligned} \alpha_a(\mu) \|e_N^u(\mu)\|_X^2 &\leq a(e_N^u(\mu), e_N^u(\mu); \mu) = r_N^1(e_N^u(\mu); \mu) - r_N^2(e_N^p(\mu); \mu) \\ &\leq \|r_N^1(\cdot; \mu)\|_{X'} \|e_N^u(\mu)\|_X + \|r_N^2(\cdot; \mu)\|_{Y'} \|e_N^p(\mu)\|_Y. \end{aligned}$$

Using further (2.46), we obtain the quadratic inequality

$$\begin{aligned} \alpha_{a}(\mu) \|e_{N}^{u}(\mu)\|_{X}^{2} - \left(\|r_{N}^{1}(\cdot;\mu)\|_{X'} + \frac{\gamma_{a}(\mu)}{\beta_{\mathrm{Br}}(\mu)}\|r_{N}^{2}(\cdot;\mu)\|_{Y'}\right) \|e_{N}^{u}(\mu)\|_{X} \\ - \frac{\|r_{N}^{2}(\cdot;\mu)\|_{Y'}\|r_{N}^{1}(\cdot;\mu)\|_{X'}}{\beta_{\mathrm{Br}}(\mu)} \leq 0, \end{aligned}$$

which we may solve for $||e_N^u(\mu)||_X$; this immediately leads to (2.54). Again combined with (2.46), this clearly also yields an associated error bound for $||e_N^p(\mu)||_Y$.

In the special case of a symmetric problem, the *a posteriori* error bounds given in Proposition 2.4.2 can be improved. Now, also the following result holds true.

Proposition 2.4.6. Let $a(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\max}$, and $\alpha_a^{\mathrm{LB}}(\mu)$, $\gamma_a^{\mathrm{UB}}(\mu)$, $\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)$ satisfying (2.29) and (2.30), we define

$$\Delta_N^{u,\text{sym}}(\mu) \equiv \frac{\|r_N^1(\cdot;\mu)\|_{X'}}{\alpha_a^{\text{LB}}(\mu)} + \sqrt{\frac{\gamma_a^{\text{UB}}(\mu)}{\alpha_a^{\text{LB}}(\mu)}} \frac{\|r_N^2(\cdot;\mu)\|_{Y'}}{\beta_{\text{Br}}^{\text{LB}}(\mu)},\tag{2.55}$$

$$\Delta_N^{p,\text{sym}}(\mu) \equiv \left(1 + \sqrt{\frac{\gamma_a^{\text{UB}}(\mu)}{\alpha_a^{\text{LB}}(\mu)}}\right) \frac{\|r_N^1(\cdot;\mu)\|_{X'}}{\beta_{\text{Br}}^{\text{LB}}(\mu)} + \frac{\gamma_a^{\text{UB}}(\mu)}{\beta_{\text{Br}}^{\text{LB}}(\mu)} \frac{\|r_N^2(\cdot;\mu)\|_{Y'}}{\beta_{\text{Br}}^{\text{LB}}(\mu)}.$$
 (2.56)

Then, $\Delta_N^{u,\text{sym}}(\mu)$ and $\Delta_N^{p,\text{sym}}(\mu)$ are upper bounds for the errors $e_N^u(\mu)$ and $e_N^p(\mu)$ such that

$$\|e_N^u(\mu)\|_X \le \Delta_N^{u,\text{sym}}(\mu) < \Delta_N^{u,\text{Br}}(\mu), \quad \|e_N^p(\mu)\|_Y \le \Delta_N^{p,\text{sym}}(\mu) < \Delta_N^{p,\text{Br}}(\mu),$$
(2.57)

for all $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\max}$.

Proof. Let μ be an arbitrary but fixed parameter in \mathcal{D} and $N \in \mathbb{N}_{\max}$. We now proceed as in the proof of Proposition 2.4.2, only that we may now exploit the fact that $a(\cdot, \cdot; \mu)$ defines an inner product on X. We recall that the associated norm is denoted by $\|\cdot\|_{X,\mu} = \sqrt{a(\cdot, \cdot; \mu)}$.

The error $e_N^u(\mu) \in X$ in the approximation of the primal variable may now be uniquely decomposed into $e_N^u(\mu) = \tilde{e}_N^0(\mu) + \tilde{e}_N^{\perp}(\mu)$ where $\tilde{e}_N^0(\mu) \in \ker(B(\mu))$ and $\tilde{e}_N^{\perp}(\mu) \in X$ such that

$$a(\tilde{e}_{N}^{\perp}(\mu), v_{0}; \mu) = 0, \quad \forall v_{0} \in \ker(B(\mu)).$$
 (2.58)

From (2.40) and (2.58), $\tilde{e}_N^0(\mu) \in \ker(B(\mu))$ then solves

$$a(\tilde{e}_{N}^{0}(\mu), v_{0}; \mu) = r_{N}^{1}(v_{0}; \mu) - a(\tilde{e}_{N}^{\perp}(\mu), v_{0}; \mu) = r_{N}^{1}(v_{0}; \mu), \quad \forall v_{0} \in \ker(B(\mu)).$$

Setting here $v_0 = \tilde{e}_N^0(\mu)$, we have

$$\begin{split} \|\tilde{e}_{N}^{0}(\mu)\|_{X,\mu}^{2} &= r_{N}^{1}(\tilde{e}_{N}^{0}(\mu);\mu) \leq \|\tilde{e}_{N}^{0}(\mu)\|_{X} \sup_{v_{0} \in \ker(B(\mu))} \frac{r_{N}^{1}(v_{0};\mu)}{\|v_{0}\|_{X}} \\ &\leq \frac{1}{\sqrt{\alpha_{a}(\mu)}} \|\tilde{e}_{N}^{0}(\mu)\|_{X,\mu} \sup_{v_{0} \in \ker(B(\mu))} \frac{r_{N}^{1}(v_{0};\mu)}{\|v_{0}\|_{X}}, \end{split}$$

where the last inequality follows from (2.10). Hence, $\tilde{e}_N^0(\mu)$ is bounded by

$$\|\tilde{e}_{N}^{0}(\mu)\|_{X,\mu} \leq \frac{1}{\sqrt{\alpha_{a}(\mu)}} \sup_{v_{0} \in \ker(B(\mu))} \frac{r_{N}^{1}(v_{0};\mu)}{\|v_{0}\|_{X}} \leq \frac{\|r_{N}^{1}(\cdot;\mu)\|_{X'}}{\sqrt{\alpha_{a}(\mu)}}.$$
(2.59)

To obtain an upper bound for $\tilde{e}_N^{\perp}(\mu)$, we here consider the inf-sup constant

$$\tilde{\beta}(\mu) \equiv \inf_{q \in Y} \sup_{v \in X} \frac{b(v, q; \mu)}{\|q\|_{Y} \|v\|_{X, \mu}}.$$
(2.60)

From (2.8) and (2.10), we have

$$\frac{\beta_{\rm Br}(\mu)}{\sqrt{\gamma_a(\mu)}} \le \tilde{\beta}(\mu) \le \frac{\beta_{\rm Br}(\mu)}{\sqrt{\alpha_a(\mu)}},\tag{2.61}$$

and thus particularly $\tilde{\beta}(\mu) > 0$ by the LBB inf-sup condition (2.11). Consequently, analogous to (2.44), also

$$\tilde{\beta}(\mu) \|v\|_{X,\mu} \le \sup_{q \in Y} \frac{b(v,q;\mu)}{\|q\|_Y}$$

holds true for any $v \in X$ such that $a(v, v_0; \mu) = 0$ for all $v_0 \in \ker(B(\mu))$. Applied to $\tilde{e}_N^{\perp}(\mu)$ satisfying (2.58), this yields

$$\|\tilde{e}_{N}^{\perp}(\mu)\|_{X,\mu} \leq \frac{1}{\tilde{\beta}(\mu)} \sup_{q \in Y} \frac{b(\tilde{e}_{N}^{\perp}(\mu), q; \mu)}{\|q\|_{Y}} = \frac{\|r_{N}^{2}(\cdot; \mu)\|_{Y'}}{\tilde{\beta}(\mu)},$$
(2.62)

where the equality again follows from (2.41) as $B(\mu) \tilde{e}_N^{\perp}(\mu) = B(\mu) e_N^u(\mu)$ in Y'. Now, using the triangle inequality, we may finally derive that

$$\|e_N^u(\mu)\|_{X,\mu} \le \|\tilde{e}_N^0(\mu)\|_{X,\mu} + \|\tilde{e}_N^{\perp}(\mu)\|_{X,\mu} \le \frac{\|r_N^1(\cdot;\mu)\|_{X'}}{\sqrt{\alpha_a(\mu)}} + \frac{\sqrt{\gamma_a(\mu)}}{\beta_{\mathrm{Br}}(\mu)}\|r_N^2(\cdot;\mu)\|_{Y'}, \quad (2.63)$$

by combining (2.59), (2.62), and (2.61); the bound (2.55) thus follows from (2.63), (2.10), and (2.29), (2.30).

For the error $e_N^p(\mu)$ in the approximation of the Lagrange multiplier, we again obtain from (2.11) and (2.40) that

$$\begin{aligned} \|e_{N}^{p}(\mu)\|_{Y} &\leq \frac{1}{\beta_{\mathrm{Br}}(\mu)} \sup_{v \in X} \frac{b(v, e_{N}^{p}(\mu); \mu)}{\|v\|_{X}} = \frac{1}{\beta_{\mathrm{Br}}(\mu)} \sup_{v \in X} \frac{r_{N}^{1}(v; \mu) - a(e_{N}^{u}(\mu), v; \mu)}{\|v\|_{X}} \\ &\leq \frac{1}{\beta_{\mathrm{Br}}(\mu)} \Big(\|r_{N}^{1}(\cdot; \mu)\|_{X'} + \sqrt{\gamma_{a}(\mu)} \|e_{N}^{u}(\mu)\|_{X,\mu} \Big), \end{aligned}$$
(2.64)

where the last inequality holds by the Cauchy–Schwarz inequality for the inner product

 $a(\cdot, \cdot; \mu)$ and (2.8). Together with (2.63) and again (2.29), (2.30), this leads to (2.56). As it is clearly $\alpha_a(\mu) \leq \gamma_a(\mu)$, we have $\Delta_N^{u, \text{sym}}(\mu) < \Delta_N^{u, \text{Br}}(\mu)$ and $\Delta_N^{p, \text{sym}}(\mu) < \Delta_N^{p, \text{Br}}(\mu)$, which eventually yields (2.57).

Remark 2.4.7. We here comment on the symmetric case in a little more detail.

First, we clearly obtain

$$\Delta_N^{\text{sym}}(\mu) \equiv \sqrt{\left(\Delta_N^{u,\text{sym}}(\mu)\right)^2 + \left(\Delta_N^{p,\text{sym}}(\mu)\right)^2}, \quad \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\text{max}},$$
(2.65)

as a rigorous upper bound for the combined error $e_N(\mu)$ such that

$$||e_N(\mu)||_Z \leq \Delta_N^{\text{sym}}(\mu) < \Delta_N^{\text{Br}}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\text{max}}.$$

Analogous to Corollary 2.4.3, the associated effectivities then satisfy

$$1 \le \eta_N^{\text{sym}}(\mu) \equiv \frac{\Delta_N^{\text{sym}}(\mu)}{\|e_N(\mu)\|_Z} \le C_1^{\text{sym}}(\mu) \big(\gamma_a(\mu) + \gamma_b(\mu)\big), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\text{max}},$$
(2.66)

where $0 < C_1^{\text{sym}}(\mu) < C_1(\mu)$ is a constant depending on $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, and $\beta_{\text{Br}}^{\text{LB}}(\mu)$.

Second, also the error bound (2.52) given in Remark 2.4.4 may be specialized to the symmetric case: Using (2.59), (2.62), (2.61), (2.10), and (2.51), (2.53),

$$\frac{1}{\alpha_a(\mu)} \sup_{v_0 \in \ker(B(\mu))} \frac{r_N^1(v_0;\mu)}{\|v_0\|_X} + \sqrt{\frac{\gamma_a(\mu)}{\alpha_a(\mu)}} \frac{\|r_N^2(\cdot;\mu)\|_{Y'}}{\beta_{\mathrm{Br}}(\mu)}$$

here represents a reliable and efficient a posteriori error estimator for $||e_N^u(\mu)||_X$.

Third, the error $e_N^u(\mu)$ may clearly also be measured in the energy norm $\|\cdot\|_{X,\mu} = \sqrt{a(\cdot,\cdot;\mu)}$. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\max}$, and $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, $\beta_{\text{Br}}^{\text{LB}}(\mu)$ satisfying (2.29) and (2.30), we define

$$\tilde{\Delta}_{N}^{u,\text{sym}}(\mu) \equiv \frac{\|r_{N}^{1}(\cdot;\mu)\|_{X'}}{\sqrt{\alpha_{a}^{\text{LB}}(\mu)}} + \frac{\sqrt{\gamma_{a}^{\text{UB}}(\mu)}}{\beta_{\text{Br}}^{\text{LB}}(\mu)} \|r_{N}^{2}(\cdot;\mu)\|_{Y'}.$$
(2.67)

Then, from (2.63) and (2.29), (2.30), $\tilde{\Delta}_N^{u,\text{sym}}(\mu)$ is an upper bound for $e_N^u(\mu)$ such that

$$\|e_N^u(\mu)\|_{X,\mu} \le \tilde{\Delta}_N^{u,\text{sym}}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\text{max}}.$$
(2.68)

It directly follows from the proof of Proposition 2.4.6 that the effectivities $\tilde{\eta}_N^{u,\text{sym}}(\mu) \equiv \tilde{\Delta}_N^{u,\text{sym}}(\mu)/\|e_N^u(\mu)\|_{X,\mu}$ and $\eta_N^{u,\text{sym}}(\mu) \equiv \Delta_N^{u,\text{sym}}(\mu)/\|e_N^u(\mu)\|_X$ satisfy $\tilde{\eta}_N^{u,\text{sym}}(\mu) \leq \eta_N^{u,\text{sym}}(\mu)$ for all $\mu \in \mathcal{D}, N \in \mathbb{N}_{\text{max}}$.

In the following, we shall not only examine the above *a posteriori* error bounds with respect to sharpness and computational efficiency (see $\S2.6.2$ and $\S3.3.4$), but also emphasize their importance to the construction of efficient reduced basis approximation spaces (see $\S2.7$).

2.5 Construction of Reduced Basis Approximation Spaces

We now turn to the construction of the reduced basis approximation spaces X_N and Y_N . In general, the reduced basis method constructs its low-dimensional approximation spaces by exploiting the parametric dependence of the problem: Solutions to the truth problem (2.13) reside on the subset

$$\mathcal{M} \equiv \{ (u(\mu), p(\mu)) \mid \mu \in \mathcal{D} \} \subset X \times Y,$$

and thus the method typically constructs $X_N \times Y_N$, $N \in \mathbb{N}_{\max}$, by focusing on \mathcal{M} . According to the so-called Lagrange approach, we consider the following option: We assume that we are given a sample of parameter snapshots $\mathcal{D}_N \equiv \{\mu_n \mid 1 \leq n \leq N\} \subset \mathcal{D}, N \in \mathbb{N}_{\max}$; for our present purposes, \mathcal{D}_N may represent any sequence of nested samples in \mathcal{D} , i.e., $\mathcal{D}_1 \subset \mathcal{D}_2 \subset \cdots \subset \mathcal{D}_N$. For $N \in \mathbb{N}_{\max}$, we then define our reduced basis approximation spaces X_N and Y_N as

$$Y_N \equiv \text{span}\{ p(\mu_n) \mid 1 \le n \le N \} = \text{span}\{ \xi_n \mid 1 \le n \le N \},$$
(2.69)

$$X_N^0 \equiv \text{span}\{ u(\mu_n) \mid 1 \le n \le N \},$$
(2.70)

where $\xi_n \in Y$, $1 \leq n \leq N$, denote $(\cdot, \cdot)_Y$ -orthonormal basis functions; we shall refer to this choice as Option 0.

For saddle point problems, building the reduced basis approximation space solely from snapshots $(u(\mu_N), p(\mu_N)) \in \mathcal{M}, N \in \mathbb{N}_{\text{max}}$, is not sufficient. We have already seen in §2.3.2 that the inf-sup condition (2.17) represents an additional requirement for X_N, Y_N to provide a well-posed system (2.14). Now, let μ be an arbitrary but fixed parameter in \mathcal{D} . A priori, it is not known whether (X_N^0, Y_N) satisfies $\beta_N(\mu) > 0$. More precisely, in case of Option 0, we meet either of the following two situations:

- (i) $\beta_N(\mu) = 0$: The system (2.14) is ill-posed as it is not uniquely solvable (see §2.3.2).
- (ii) $\beta_N(\mu) > 0$: The system (2.14) is well-posed (see §2.3.2) and thus uniquely solvable; we obtain approximations $u_N(\mu)$ and $p_N(\mu)$ for $u(\mu)$ and $p(\mu)$, respectively. However, choosing the approximation spaces X_N, Y_N as in (2.69), (2.70) is in fact a special case: Whenever $\beta_N(\mu) > 0$, we in particular have dim $(X_N^0) = \dim(Y_N)$. Consequently, we here encounter a *locking phenomenon*: Positive inf-sup constants $\beta_N(\mu)$ provide a wellposed system (2.14), but the operator $B_N(\mu)$ is not only surjective but bijective; the

space ker $(B_N(\mu))$ is thus trivial and there exists only one solution, $B_N(\mu)^{-1}g_N(\mu)$, satisfying the constraints (2.16). This solution is unlikely to provide an accurate approximation for $u(\mu)$ as components of $u(\mu)$ in ker $(B(\mu))$ determined by the first equation in (2.13) are essentially neglected.

In summary, Option 0 cannot be expected to provide useful approximations.

We now consider further options to construct the reduced basis approximation spaces that aim to avoid the situations above. First, we address the situation in (i) and thus the issue that the inf-sup constants $\beta_N(\mu)$ may approach zero for a given parameter $\mu \in \mathcal{D}$. It is shown in [92, 96] that the space X_N^0 can be enriched such that we obtain a *provably* stable pair (X_N, Y_N) where $\beta_N(\mu) > 0$ holds for all $\mu \in \mathcal{D}$. To explain this in detail, we have to introduce further notations. For $\mu \in \mathcal{D}$ and $1 \leq k \leq Q_b$, let $T_\mu : Y \to X$ and $T^k : Y \to X$ be the Riesz representers given by

$$(T_{\mu}q, v)_X = b(v, q; \mu), \quad (T^k q, v)_X = b^k(v, q), \quad \forall v \in X, q \in Y;$$
 (2.71)

here, $b^k(\cdot, \cdot)$, $1 \le k \le Q_b$, are the parameter-independent bilinear forms in the affine decomposition (2.7) of $b(\cdot, \cdot; \mu)$. Since T_{μ} and T^k satisfy

$$T_{\mu}q = \arg \sup_{v \in X} \frac{b(v, q; \mu)}{\|v\|_{X}}, \quad T^{k}q = \arg \sup_{v \in X} \frac{b^{k}(v, q)}{\|v\|_{X}}, \quad \forall q \in Y,$$
(2.72)

for all $\mu \in \mathcal{D}$, $1 \leq k \leq Q_b$, they are also called *supremizer* functions. By the affine decomposition (2.7) for $b(\cdot, \cdot; \mu)$, we in particular have

$$T_{\mu}q = \sum_{k=1}^{Q_b} \Theta_b^k(\mu) T^k q, \quad \forall q \in Y, \quad \forall \mu \in \mathcal{D}.$$
(2.73)

Now, to any given discrete space $Y_N \subseteq Y$ with basis functions $\xi_n \in Y$, $1 \leq n \leq N_Y$, a sufficient condition for (X_N, Y_N) to form a stable pair is given by

$$T_{\mu}\xi_n \in X_N, \quad \forall \ 1 \le n \le N_Y, \quad \forall \ \mu \in \mathcal{D};$$

$$(2.74)$$

in this case, it can be shown (see [92, 96]) that

$$\beta_N(\mu) \ge \beta_{\mathrm{Br}}(\mu), \quad \forall \ \mu \in \mathcal{D},$$
(2.75)

and (2.17) thus follows from the LBB inf-sup condition (2.11). Due to the affine expansion (2.73), (2.74) is satisfied by the *finite*-dimensional space

span{
$$T^{k}\xi_{n} \mid 1 \le n \le N_{Y}, \ 1 \le k \le Q_{b}$$
} \supseteq { $T_{\mu}\xi_{n} \mid 1 \le n \le N_{Y}, \ \mu \in \mathcal{D}$ }. (2.76)

Associated to Y_N defined as in (2.69), we may therefore choose X_N as

$$X_N^1 \equiv X_N^0 \oplus \text{span}\{ T^k \xi_n \mid 1 \le n \le N, \ 1 \le k \le Q_b \},$$
(2.77)

and obtain by (2.74) and (2.76) pairs of approximation spaces (X_N^1, Y_N) , $N \in \mathbb{N}_{\text{max}}$, which satisfy (2.17) for all parameter values $\mu \in \mathcal{D}$; we shall refer to this choice as Option 1.

However, Option 1 is also rather expensive as it requires $Q_b N$ additional basis functions compared to X_N^0 . For this reason, we consider Option 2 as the following modification (see [88, 96]): We here add only N supremizer functions instead of $Q_b N$ as in (2.77) to X_N^0 ,

$$X_N^2 \equiv X_N^0 \oplus \operatorname{span}\{T_{\mu_n}\xi_n \mid 1 \le n \le N\}.$$
(2.78)

In this case, we can no longer prove stability of (X_N^2, Y_N) a priori but obtain significantly lower space dimensions than in Option 1.

Now, we focus on the situation described in (ii). The primary issue here is not a vanishing inf-sup constant $\beta_N(\mu)$ but rather a trivial space ker $(B_N(\mu))$. In Options 1 and 2, the situation of ker $(B(\mu))$ being trivial has been avoided by an enrichment of X_N^0 with supremizer functions providing particularly large inf-sup constants $\beta_N(\mu)$. For any parameter $\mu \in \mathcal{D}$, it can be expected from (2.52) in Remark 2.4.4 and also the *a priori* error estimates in §2.3.3 that small inf-sup constants $\beta_N(\mu)$ will have more dramatic effects on $\|e_N^p(\mu)\|_Y$ than on $\|e_N^u(\mu)\|_X$ (see also [13, p. 57, Remark 2.8]): By (2.52) and (2.19), (2.22), $\|e_N^u(\mu)\|_X$ clearly profits from a large space ker $(B_N(\mu))$ but does not explicitly depend on the values of $\beta_N(\mu)$. The right-hand side in (2.52) suggests that $\|e_N^u(\mu)\|_X$ first of all benefits from a good testing space ker $(B(\mu)) \cap X_N \subseteq \text{ker}(B_N(\mu))$. For this purpose, we now enrich X_N^0 with additional truth solutions $u(\mu') \in \text{ker}(B(\mu'))$: Given a second sample of parameter snapshots $\mathcal{D}'_N \equiv \{\mu'_n \mid 1 \le n \le N\} \subset \mathcal{D}, \, \mathcal{D}'_N \cap \mathcal{D}_N = \emptyset, N \in \mathbb{N}_{\text{max}}$, we now define Option 3:

$$X_N^3 \equiv X_N^0 \oplus \text{span}\{ u(\mu'_n) \mid 1 \le n \le N \}.$$
(2.79)

Again, we cannot prove a priori that this choice provides stable pairs $(X_N^3, Y_N), N \in \mathbb{N}_{\max}$.

Note that for algebraic stability reasons, we in fact express both Y_N and X_N by orthonormal basis functions: Analogous to (2.69), we also set $X_N = \text{span}\{\phi_m \mid 1 \leq m \leq N_X\}$, where $\phi_m \in X$, $1 \leq m \leq N_X$, are $(\cdot, \cdot)_X$ -orthonormal.

2.6 Offline-Online Computational Procedure

The basic strategy lies in the μ -affine dependence (2.7) of the involved operators. Since much of this machinery is by now standard in reduced basis methods (see, e.g., [87, 95]), we only briefly summarize the Offline-Online procedure and associated computational cost.

2.6.1 Reduced Basis Approximation

Noting that for any $\mu \in \mathcal{D}$, we can expand $u_N(\mu) \in X_N$ and $p_N(\mu) \in Y_N$ into

$$u_N(\mu) = \sum_{m=1}^{N_X} u_{N\,m}(\mu)\phi_m, \quad p_N(\mu) = \sum_{n=1}^{N_Y} p_{N\,n}(\mu)\xi_n, \quad (2.80)$$

the reduced basis system (2.14) may be written as

$$\sum_{n=1}^{N_X} \sum_{k=1}^{Q_a} \Theta_a^k(\mu) A_{N\,im}^k u_{N\,m}(\mu) + \sum_{n=1}^{N_Y} \sum_{k=1}^{Q_b} \Theta_b^k(\mu) B_{N\,ni}^k \, p_{N\,n}(\mu) = \sum_{k=1}^{Q_f} \Theta_f^k(\mu) f_{N\,i}^k, \qquad (2.81)$$

for all $1 \leq i \leq N_X$, and

$$\sum_{m=1}^{N_X} \sum_{k=1}^{Q_b} \Theta_b^k(\mu) B_{Njm}^k u_{Nm}(\mu) = \sum_{k=1}^{Q_g} \Theta_g^k(\mu) g_{Nj}^k, \quad \forall \ 1 \le j \le N_Y;$$
(2.82)

here, the μ -independent quantities are given by

$$\begin{aligned}
A_{N\,ij}^{k} &\equiv a^{k}(\phi_{j},\phi_{i}), \quad 1 \leq i, j \leq N_{X}, \ 1 \leq k \leq Q_{a}, \\
B_{N\,ji}^{k} &\equiv b^{k}(\phi_{i},\xi_{j}), \quad 1 \leq i \leq N_{X}, \ 1 \leq j \leq N_{Y}, \ 1 \leq k \leq Q_{b}, \\
f_{N\,i}^{k} &\equiv f^{k}(\phi_{i}), \quad 1 \leq i \leq N_{X}, \ 1 \leq k \leq Q_{f}, \\
g_{N\,j}^{k} &\equiv g^{k}(\xi_{j}), \quad 1 \leq j \leq N_{Y}, \ 1 \leq k \leq Q_{g}.
\end{aligned}$$
(2.83)

Offline, we first form and store all parameter-independent quantities (2.83) at $O(\mathcal{N}^*)$ computational cost and storage. Online, for any given parameter value $\mu \in \mathcal{D}$, we then perform the required sums in (2.81), (2.82) and solve the resulting $N_Z \times N_Z$ system of linear equations for $(u_N(\mu), p_N(\mu))$ at $O(N_Z^3)$ computational cost. The expensive Offline stage, which is performed only once and whose cost depends on the large truth dimension \mathcal{N} , thus enables the subsequent very inexpensive Online stage that does not depend on \mathcal{N} but only on the considerably smaller dimension N_Z of the reduced basis approximation space.

2.6.2 A Posteriori Error Bounds

It is clear that there are two sets of components to the computation of the *a posteriori* error bounds: the calculation of the residual dual norms (2.31), and the calculation of the required lower and upper bounds (2.29), (2.30) to the coercivity, continuity, and/or inf-sup stability constants. The former is again an application of now standard reduced basis techniques that can be found in, e.g., [95]. The latter is achieved by a successive constraint method (SCM) proposed by Huynh et al. [55]. We here summarize the main parts of the methodology focusing on aspects specific to our saddle point context.

Remark 2.6.1. We note that there exist several approaches for the computation of lower bounds to coercivity and inf-sup stability constants. The applicability of techniques proposed in [71, 87] is unfortunately very restricted (e.g., to the case of positive coefficient functions in (2.7)); for problems involving geometry variations, we rely on more general but also more complicated and Offline-expensive approaches. Compared to earlier proposals [79, 97, 106], the SCM in [55] represents a method that (i) is very generally applicable, (ii) performs better, and (iii) is easier to implement. In addition to accurate lower bounds, it also provides upper bounds that are remarkably sharp and thus constitute valuable estimates for the respective coercivity and inf-sup constants. Nevertheless, for inf-sup constants, it still involves notable Offline computations that depend on a Q^2 -term affine parameter expansion (see (2.99)). This issue has been tackled by more recent techniques based on a "natural norm" [53]. However, though well suited for certain problems, the proposed algorithm is in our situation highly cumbersome and involves eigenvalue problems that are much more difficult to solve; also, no upper bounds are provided.

Dual Norms of the Residuals

We first exploit the μ -affine expansions (2.7) and (2.80) of the involved operators and the reduced basis approximation, respectively, to express the residuals $r_N^1(\cdot;\mu)$ and $r_N^2(\cdot;\mu)$ defined in (2.32) and (2.33) as

$$r_{N}^{1}(v;\mu) = \sum_{k=1}^{Q_{f}} \Theta_{f}^{k}(\mu) f^{k}(v) - \sum_{m=1}^{N_{X}} \sum_{k=1}^{Q_{a}} u_{Nm}(\mu) \Theta_{a}^{k}(\mu) a^{k}(\phi_{m},v) - \sum_{n=1}^{N_{Y}} \sum_{k=1}^{Q_{b}} p_{Nn}(\mu) \Theta_{b}^{k}(\mu) b^{k}(v,\xi_{n}), \quad \forall v \in X, r_{N}^{2}(q;\mu) = \sum_{k=1}^{Q_{g}} \Theta_{g}^{k}(\mu) g^{k}(q) - \sum_{m=1}^{N_{X}} \sum_{k=1}^{Q_{b}} u_{Nm}(\mu) \Theta_{b}^{k}(\mu) b^{k}(\phi_{m},q), \quad \forall q \in Y$$

Setting $Q_1 \equiv Q_f + N_X Q_a + N_Y Q_b$ and $Q_2 \equiv Q_g + N_X Q_b$, this can be written succinctly as

$$r_N^1(v;\mu) = \sum_{k=1}^{Q_1} \Theta_N^{1,k}(\mu) \, r_N^{1,k}(v), \quad \forall \, v \in X,$$
(2.84)

$$r_N^2(q;\mu) = \sum_{k=1}^{Q_2} \Theta_N^{2,k}(\mu) \, r_N^{2,k}(q), \quad \forall \ q \in Y,$$
(2.85)

where the coefficient functions $\Theta_N^{1,k}(\mu)$, $1 \leq k \leq Q_1$, and $\Theta_N^{2,k}(\mu)$, $1 \leq k \leq Q_2$, depend on μ explicitly through the coefficient functions $\Theta_{\mathcal{I}}^k(\mu)$, $1 \leq k \leq Q_{\mathcal{I}}$, $\mathcal{I} = a, b, f, g$, but also implicitly through $u_N(\mu)$ and $p_N(\mu)$; $r_N^{1,k}(\cdot) \in X'$, $1 \leq k \leq Q_1$, and $r_N^{2,k}(\cdot) \in Y'$, $1 \leq k \leq Q_2$, are parameter-independent linear functionals. We now denote by $\hat{e}_N^{1,k} \in X$ and $\hat{e}_N^{2,k} \in Y$ the Riesz representers associated with $r_N^{1,k}(\cdot)$ and $r_N^{2,k}(\cdot)$, respectively, which are obtained as the solutions to the μ -independent linear problems.

problems

$$(\hat{e}_N^{1,k}, v)_X = r_N^{1,k}(v), \quad \forall \ v \in X, \quad \forall \ 1 \le k \le Q_1,$$
(2.86)

$$(\hat{e}_N^{2,k}, q)_Y = r_N^{2,k}(q), \quad \forall \ q \in Y, \quad \forall \ 1 \le k \le Q_2.$$
 (2.87)

By (2.84) and (2.85), the Riesz representers $\hat{e}_N^1(\mu) \in X$ and $\hat{e}_N^2(\mu) \in Y$ associated with the residuals $r_N^1(\cdot;\mu)$ and $r_N^2(\cdot;\mu)$ are then given by

$$\hat{e}_N^1(\mu) = \sum_{k=1}^{Q_1} \Theta_N^{1,k}(\mu) \, \hat{e}_N^{1,k}, \quad \hat{e}_N^2(\mu) = \sum_{k=1}^{Q_2} \Theta_N^{2,k}(\mu) \, \hat{e}_N^{2,k},$$

and it finally follows that

$$\|r_N^1(\cdot;\mu)\|_{X'}^2 = (\hat{e}_N^1(\mu), \hat{e}_N^1(\mu))_X = \sum_{k,l=1}^{Q_1} \Theta_N^{1,k}(\mu) \Theta_N^{1,l}(\mu) (\hat{e}_N^{1,k}, \hat{e}_N^{1,l})_X,$$
(2.88)

$$\|r_N^2(\cdot;\mu)\|_{Y'}^2 = (\hat{e}_N^2(\mu), \hat{e}_N^2(\mu))_Y = \sum_{k,l=1}^{Q_2} \Theta_N^{2,k}(\mu) \Theta_N^{2,l}(\mu) (\hat{e}_N^{2,k}, \hat{e}_N^{2,l})_Y.$$
(2.89)

We may now summarize the Offline-Online decomposition.

In the Offline stage, we find the Riesz representers $\hat{e}_N^{1,k}$, $1 \leq k \leq Q_1$, and $\hat{e}_N^{2,k}$, $1 \leq k \leq Q_2$, and form the associated inner products $(\hat{e}_N^{1,k}, \hat{e}_N^{1,l})_X$, $1 \leq k,l \leq Q_1$, and $(\hat{e}_N^{2,k}, \hat{e}_N^{2,l})_Y$, $1 \leq k,l \leq Q_2$. These cost clearly depend on \mathcal{N} : We have to solve Q_1 and Q_2 linear problems of type (2.86) and (2.87), respectively, and to compute $(Q_1)^2$ and $(Q_2)^2$ inner products in X and Y. In the Online stage, given the reduced basis coefficients $u_N (\mu)$, $1 \leq m \leq N_X$, and $p_{Nn}(\mu)$, $1 \leq n \leq N_Y$, we can then readily compute the coefficient functions $\Theta_N^{1,k}(\mu)$, $1 \leq k \leq Q_1$, and $\Theta_N^{2,k}(\mu)$, $1 \leq k \leq Q_2$. From the stored inner products, we then assemble the sums (2.88) and (2.89) at $O((Q_1)^2) = O((Q_f + N_X Q_a + N_Y Q_b)^2)$ and $O((Q_2)^2) = O((Q_g + N_X Q_b)^2)$ operations, respectively. As desired, the operation count for the Online stage is thus independent of \mathcal{N} .

We now address the second component required for the calculation of the *a posteriori* error bounds.

Lower and Upper Bounds to Coercivity Constants

We first discuss the calculation of (Online-)efficient lower and upper bounds (2.29) to the coercivity constants $\alpha_a(\mu)$.

Following [55], the coercivity constants (2.10) can be written as the minimum of the functional $J : \mathbb{R}^{Q_a} \times \mathcal{D} \to \mathbb{R}, J(y;\mu) \equiv \sum_{k=1}^{Q_a} \Theta_a^k(\mu) y_k$,

$$\alpha_a(\mu) = \min_{y \in \mathcal{Y}} J(y;\mu), \quad \forall \ \mu \in \mathcal{D},$$
(2.90)

over the set $\mathcal{Y} \equiv \{ y \in \mathbb{R}^{Q_a} \mid \exists v_y \in X : y_k = a^k(v_y, v_y) / \|v_y\|_X^2, 1 \leq k \leq Q_a \}$; here, $a^k(\cdot, \cdot)$, $1 \leq k \leq Q_a$, are the parameter-independent bilinear forms in the affine decomposition (2.7). Now, for any $\mu \in \mathcal{D}$, we approximate \mathcal{Y} by sets \mathcal{Y}_{UB} , $\mathcal{Y}_{\text{LB}} \subseteq \mathbb{R}^{Q_a}$ such that

$$\mathcal{Y}_{\rm UB} \subseteq \mathcal{Y} \subseteq \mathcal{Y}_{\rm LB},\tag{2.91}$$

and clearly obtain lower and upper bounds to (2.90) by setting

$$\alpha_a^{\rm LB}(\mu) \equiv \min_{y \in \mathcal{Y}_{\rm LB}} J(y;\mu), \quad \alpha_a^{\rm UB}(\mu) \equiv \min_{y \in \mathcal{Y}_{\rm UB}} J(y;\mu).$$
(2.92)

The sets in (2.91) are constructed as follows: For a given $M_{\alpha} \in \mathbb{N}$ and a sample $C_K \equiv \{\omega_i \mid 1 \leq i \leq K\}$ of parameters in \mathcal{D} , let $\mathcal{P}_{M_{\alpha}}(\mu; C_K)$ denote the set of the M_{α} points closest to μ in the Euclidean norm. For any $\mu \in \mathcal{D}$, we then define

$$\mathcal{Y}_{\text{LB}}(\mu; C_K) \equiv \{ y \in \mathcal{B} \mid J(y; \mu') \ge \alpha_a(\mu'), \forall \mu' \in \mathcal{P}_{M_\alpha}(\mu, C_K) \},$$
(2.93)

$$\mathcal{Y}_{\mathrm{UB}}(C_K) \equiv \{ y^*(\omega) \mid \omega \in C_K \} = \{ y^*(\omega_1), \dots, y^*(\omega_K) \},$$
(2.94)

where $\mathcal{B} \equiv \prod_{k=1}^{Q_a} [b_k^-, b_k^+] \subseteq \mathbb{R}^{Q_a}$ with

$$b_k^- \equiv \inf_{v \in X} \frac{a^k(v, v)}{\|v\|_X^2}, \quad b_k^+ \equiv \sup_{v \in X} \frac{a^k(v, v)}{\|v\|_X^2}, \quad \forall \ 1 \le k \le Q_a,$$
(2.95)

and $y^*(\mu) \in \mathbb{R}^{Q_a}$ is given by

$$y^{*}(\mu)_{k} \equiv \frac{a^{k}(v_{\mu}, v_{\mu})}{\|v_{\mu}\|_{X}^{2}}, \quad \forall \ 1 \le k \le Q_{a}, \quad \text{where } v_{\mu} \equiv \arg \inf_{v \in X} \frac{a(v, v; \mu)}{\|v\|_{X}^{2}}.$$
 (2.96)

These sets then satisfy (2.91) (see [55]) and thus, $\alpha_a^{\text{LB}}(\mu)$ and $\alpha_a^{\text{UB}}(\mu)$ defined by (2.92) provide the required lower and upper bounds to $\alpha_a(\mu)$.

Now, the idea is to construct the parameter set C_K sufficiently large so that for a given $\mu \in \mathcal{D}$, (i) $y^*(\mu)$ will be sufficiently close to a member in $\mathcal{Y}_{\text{UB}}(C_K)$ to provide a good upper bound, and (ii) the constraints in $\mathcal{Y}_{\text{LB}}(\mu; C_K)$ will be sufficiently strict to provide a good lower bound. For this purpose, we build C_K by an Offline greedy algorithm: Starting with an exhaustive sample Ξ of parameters in \mathcal{D} and an initial sample $C_1 \equiv \{\omega_1\}, K = 1$, we enrich C_K by the parameter for which the ratio (as an estimator for the relative error in the lower and upper bounds)

$$\frac{\alpha_a^{\rm UB}(\mu) - \alpha_a^{\rm LB}(\mu)}{\alpha_a^{\rm UB}(\mu)} \tag{2.97}$$

attains its maximum over Ξ . The algorithm stops when (2.97) satisfies a prescribed tolerance $\delta_{\text{tol}}^{\text{SCM}} < 1$; we then set $K_{\text{max}} = K$.

For the Offline stage, the notable computations are: (i) $2Q_a$ eigenproblems over X to form \mathcal{B} (see (2.95)), (ii) K_{\max} eigenproblems over X to obtain $\alpha_a(\omega)$ and v_{ω} (see (2.96)) for all $\omega \in C_{K_{\max}}$, (iii) $K_{\max} Q_a$ inner products over X to form $\mathcal{Y}_{\text{UB}}(C_{K_{\max}})$, and (iv) $|\Xi| K_{\max}$ LP's with $2Q_a + M_\alpha$ constraints to compute $\alpha_a^{\text{LB}}(\mu)$. For any given parameter $\mu \in \mathcal{D}$, we then obtain $\alpha_a^{\text{LB}}(\mu)$, $\alpha_a^{\text{UB}}(\mu)$ in a highly efficient Online stage, where notable computations are only a single LP with $2Q_a + M_\alpha$ constraints associated with $\alpha_a^{\text{LB}}(\mu)$. For the exact choice of $\delta_{\text{tol}}^{\text{SCM}}$ and values for K_{\max} in practice, we refer to §3.3.4, §5.3.3, and §6.8.

Lower and Upper Bounds to Inf-Sup Constants

As suggested in [55], the calculation of (Online-)efficient lower and upper bounds (2.30) to inf-sup constants can be traced back to the calculation of bounds to coercivity constants.

For the Brezzi inf-sup constants (2.11), we obtain from (2.72) that

$$(\beta_{\rm Br}(\mu))^2 = \inf_{q \in Y} \frac{(T_{\mu}q, T_{\mu}q)_X}{\|q\|_Y^2}, \quad \forall \ \mu \in \mathcal{D},$$
(2.98)

and thus $(\beta_{Br}(\mu))^2$ represents the coercivity constant of the continuous, symmetric, and coercive (due to (2.11)) bilinear form $\mathsf{B}(\cdot,\cdot;\mu): Y \times Y \to \mathbb{R}$ with $\mathsf{B}(p,q;\mu) \equiv (T_{\mu}p,T_{\mu}q)_X$ for all $p,q \in Y$. From (2.73), $\mathsf{B}(\cdot,\cdot;\mu)$ can be μ -affinely expanded as

$$\mathsf{B}(p,q;\mu) = (T_{\mu}p, T_{\mu}q)_X = \sum_{\mathsf{K}=1}^{\mathsf{Q}} \Theta_{\mathsf{B}}^{\mathsf{K}}(\mu) \,\mathsf{B}^{\mathsf{K}}(p,q), \quad \forall \ p,q \in Y, \quad \forall \ \mu \in \mathcal{D},$$
(2.99)

where $\mathbf{Q} \equiv Q_b(Q_b + 1)/2$, and $\Theta_{\mathsf{B}}^{\mathsf{K}}(\mu)$, $\mathsf{B}^{\mathsf{K}}(\cdot, \cdot)$ are given by

$$\begin{split} \Theta_{\mathsf{B}}^{\mathsf{K}}(\mu) &= \Theta_{\mathsf{B}}^{(k,l)}(\mu) \equiv \left(\frac{1}{1+\delta_{kl}}\right) \Theta_{b}^{k}(\mu) \,\Theta_{b}^{l}(\mu), \quad \forall \ \mu \in \mathcal{D}, \\ \mathsf{B}^{\mathsf{K}}(p,q) &= \mathsf{B}^{(k,l)}(p,q) \equiv (T^{k}p,T^{l}q)_{X} + (T^{l}p,T^{k}q)_{X}, \quad \forall \ p,q \in Y, \end{split}$$

for all $1 \leq \mathsf{K} \equiv (k, l) \leq \mathsf{Q}, 1 \leq k \leq l \leq Q_b$; here, δ_{kl} denotes the Kronecker delta.

We therefore obtain (Online-)efficient lower and upper bounds to $(\beta_{Br}(\mu))^2$ by the method described for coercivity constants. However, we note that the computational cost for the Offline stage here depend on $\mathbf{Q} = O(Q_b^2) \ge Q_b$.

An analogous procedure can be applied to compute (Online-)efficient lower and upper bounds $\beta_{\text{Ba}}^{\text{LB}}(\mu)$, $\beta_{\text{Ba}}^{\text{UB}}(\mu)$ to the Babuška inf-sup constants (2.12) (see [55]). It only differs in the computational cost for the Offline stage: They here depend on Q^2 where $\max\{Q_a, Q_b\} \leq Q \leq Q_a + Q_b$.

Lower and Upper Bounds to Continuity Constants

For the calculation of (Online-)efficient lower and upper bounds (2.29) to the continuity constants $\gamma_a(\mu)$, we distinguish between the symmetric and the nonsymmetric case: If $a(\cdot, \cdot; \mu)$ is symmetric for all $\mu \in \mathcal{D}$, lower and upper bounds may be computed by an analogous procedure as applied for coercivity constants; the nonsymmetric case is then traced back to the symmetric case by an analogous approach as pursued for inf-sup constants.

The symmetric case: We first consider the case where $a(\cdot, \cdot; \mu)$ is symmetric for all $\mu \in \mathcal{D}$. The continuity constants (2.8) are then given by

$$\gamma_a(\mu) = \sup_{v \in X} \frac{a(v, v; \mu)}{\|v\|_X^2} = \max_{y \in \mathcal{Y}} J(y; \mu), \quad \forall \ \mu \in \mathcal{D},$$

where the functional J and the set \mathcal{Y} are defined as before. Therefore, for any $\mu \in \mathcal{D}$, we here construct sets $\mathcal{Y}_{\text{LB}}(C_K) \subseteq \mathcal{Y} \subseteq \mathcal{Y}_{\text{UB}}(\mu; C_K)$ by

$$\mathcal{Y}_{\mathrm{LB}}(C_K) \equiv \{ y^{**}(\omega) \mid \omega \in C_K \}, \mathcal{Y}_{\mathrm{UB}}(\mu; C_K) \equiv \{ y \in \mathcal{B} \mid J(y; \mu') \le \gamma_a(\mu'), \forall \mu' \in P_{M_\alpha}(\mu; C_K) \},$$

where $y^{**}(\mu) \in \mathbb{R}^{Q_a}$ is defined as

$$y^{**}(\mu)_k \equiv \frac{a^k(w_\mu, w_\mu)}{\|w_\mu\|_X^2}, \quad \forall \ 1 \le k \le Q_a, \quad \text{where } w_\mu \equiv \arg \sup_{v \in X} \frac{a(v, v; \mu)}{\|v\|_X^2};$$

lower and upper bounds to the continuity constants are then obtained by

$$\gamma_a^{\rm LB}(\mu) \equiv \max_{y \in \mathcal{Y}_{\rm LB}(C_K)} J(y;\mu), \quad \gamma_a^{\rm UB}(\mu) \equiv \max_{y \in \mathcal{Y}_{\rm UB}(\mu;C_K)} J(y;\mu).$$

We build the set C_K by the same greedy algorithm as described for coercivity constants, only that we here consider the ratio

$$\frac{\gamma^{\rm UB}_a(\mu)-\gamma^{\rm LB}_a(\mu)}{\gamma^{\rm LB}_a(\mu)}$$

to estimate the error in the bounds; computational cost essentially remain the same as listed for the coercivity constants.

The nonsymmetric case: Similar to (2.98), we obtain for the continuity constants (2.8) that

$$(\gamma_a(\mu))^2 = \sup_{v \in X} \frac{(S_\mu v, S_\mu v)_X}{\|v\|_X^2}, \quad \forall \ \mu \in \mathcal{D},$$

where $S_{\mu}: X \to X$ denotes the Riesz representer associated with $a(\cdot, \cdot; \mu)$ given by

$$(S_{\mu}u, v)_X = a(u, v; \mu), \quad \forall v \in X, \ \forall u \in X, \ \forall \mu \in \mathcal{D}.$$

Thus, $(\gamma_a(\mu))^2$ represents the continuity constant of the continuous (due to (2.8)) and symmetric bilinear form $A(\cdot, \cdot; \mu) : X \times X \to \mathbb{R}$ with $A(u, v; \mu) \equiv (S_{\mu}u, S_{\mu}v)_X$ for all $u, v \in X$, which moreover exhibits a μ -affine expansion analogous to (2.99).

We therefore obtain (Online-)efficient lower and upper bounds to $(\gamma_a(\mu))^2$ by the method described above for the symmetric case. However, as for inf-sup constants, we note that the computational cost for the Offline stage now depend on $Q_a^2 \ge Q_a$.

2.7 Adaptive Sampling Procedures

In this section, all parts of the methodology discussed in §2.4, §2.5, and §2.6 coalesce to tackle the key question of how to construct reduced basis approximation spaces providing rapidly convergent approximations.

Rigorous and computationally efficient reduced basis a posteriori error bounds enable us to invoke a greedy sampling process (see [10, 15] and the references therein) identifying relatively few basis functions that suffice to warrant a desired accuracy for any parameter query: Starting with an exhaustive sample Σ of parameter points in \mathcal{D} and an initial sample μ_1 , it detects the parameter μ_2 for which a chosen error bound $\Delta_N(\mu)$ attains its maximum over Σ and then enriches the current reduced basis approximation spaces with the new basis functions associated with μ_2 (see §2.5); this step is repeated until a prescribed error tolerance is satisfied. In the case of saddle point problems, the inf-sup condition (2.17) represents an additional requirement for the reduced basis approximation spaces. Therefore, according to our observations in §2.5, the above sampling process will in general (e.g., in combination with Option 0 (see (2.70)) not succeed. In earlier work (see, e.g., [75]), it has been applied in combination with either Option 1 (see (2.77)) or Option 2 (see (2.78)), where stability is guaranteed by adding supremizer functions to the reduced basis approximation space X_N for the primal variable *in each step*. For Option 2, the exact procedure is summarized in Algorithm 1.

However, using this procedure, the question arises of whether the resulting reduced basis approximation spaces are indeed computationally efficient and a reduced basis approximation space X_N with a dimension twice as large as for Y_N is required to guarantee stability. Addressing this issue, we here present a new adaptive sampling procedure for saddle point problems: In addition to greedily selecting the parameter snapshots μ_1, \ldots, μ_N (see §2.5), we now also adaptively recognize the need for stabilization and thus identify relatively few basis functions that suffice to warrant *both* a desired accuracy *and* stability for any parameter query. Note that techniques similar to those presented here have also been proposed in [23] in the context of transport dominated problems.

We assume that we are given a stable pair (X_{N-1}, Y_{N-1}) of reduced basis approximation spaces that is based on the parameter snapshots $\mathcal{D}_{N-1} \equiv \{\mu_1, \ldots, \mu_{N-1}\} \subset \Sigma$ (see §2.5). As in the standard greedy procedure (see Algorithm 1), we use our rigorous and (Online-)efficient error bounds (see §2.4 and §2.6) to find $\mu_N \in \Sigma$ and thus the truth solution $(u(\mu_N), p(\mu_N)) \in \mathcal{M}$ that is represented least in the current basis; this solution is then appended to the current basis, and we obtain a subsequent pair (X_N, Y_N) .

From our observations in §2.3.2, §2.3.3, and §2.5, we see that the new pair (X_N, Y_N) will in general provide useful approximations only if the inf-sup constants $\beta_N(\mu)$, $\mu \in \mathcal{D}$, do not become too small. This can be achieved through an additional enrichment of X_N : We may

Algorithm 1 Standard Greedy Algorithm (using Option 2)

1: Choose $\Sigma \subset \mathcal{D}, \, \delta_{\text{tol}} \in (0, 1), \, \text{and} \, \mu_1 \in \Sigma$ 2: Set $N \leftarrow 0$, $\mathcal{D}_N \leftarrow \{\}$, $N_Y \leftarrow 0$, $Y_N \leftarrow \{\}$, $N_X \leftarrow 0$, $X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N+1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $N_Y \leftarrow N_Y + 1, Y_N \leftarrow Y_{N-1} \oplus \operatorname{span}\{ p(\mu_N) \} = Y_{N-1} \oplus \operatorname{span}\{ \xi_N \} (\operatorname{see} (2.69))$ 5: $N_X \leftarrow N_X + 2, X_N \leftarrow X_{N-1} \oplus \operatorname{span}\{u(\mu_N), T_{\mu_N}\xi_N\} \text{ (see (2.78))}$ 6: for all $\mu \in \Sigma$ do 7: Compute $(u_N(\mu), p_N(\mu))$ and $\Delta_N(\mu)$ 8: end for 9: $\mu_{N+1} \equiv \arg \max_{\mu \in \Sigma} \Delta_N(\mu)$ 10:11: **until** $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 12: $N_{\max} \leftarrow N$

Algorithm 2

1: Choose $\Sigma \subset \mathcal{D}, \, \delta_{tol}, \delta_{tol}^{\beta} \in (0, 1), \, and \, \mu_1 \in \Sigma$ 2: Set $N \leftarrow 0$, $\mathcal{D}_N \leftarrow \{\}$, $N_Y \leftarrow 0$, $Y_N \leftarrow \{\}$, $N_X \leftarrow 0$, $X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N+1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $N_Y \leftarrow N_Y + 1, Y_N \leftarrow Y_{N-1} \oplus \operatorname{span}\{p(\mu_N)\} \text{ (see (2.69))}$ 5: $N_X \leftarrow N_X + 1, X_N \leftarrow X_{N-1} \oplus \operatorname{span}\{u(\mu_N)\}$ 6: while (true) do 7: for all $\mu \in \Sigma$ do 8: Compute $\hat{d}_N^{\beta}(\mu) \equiv \max\left\{\frac{\beta_{\mathrm{Br}}^{\mathrm{UB}}(\mu) - \beta_N(\mu)}{\beta_{\mathrm{Dr}}^{\mathrm{UB}}(\mu)}, 0\right\}$ (cf. (2.101)) 9: end for 10: $\mu^* \equiv \arg \max_{\mu \in \Sigma} \hat{d}_N^\beta(\mu)$ 11:if $\hat{d}_N^{\beta}(\mu^*) < \delta_{\text{tol}}^{\beta}$, then 12:break 13:end if 14: $N_X \leftarrow N_X + 1, X_N \leftarrow X_N \oplus \text{span}\{T_{\mu^*}\varrho_N(\mu^*)\} \text{ (see (2.71), (2.100))}$ 15:end while 16:for all $\mu \in \Sigma$ do 17:Compute $(u_N(\mu), p_N(\mu))$ and $\Delta_N(\mu)$ 18:end for 19:20: $\mu_{N+1} \equiv \arg \max_{\mu \in \Sigma} \Delta_N(\mu)$ 21: **until** $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 22: $N_{\max} \leftarrow N$

Algorithm 3

1: Choose $\Sigma \subset \mathcal{D}, \, \delta_{\text{tol}}, \delta_{\text{tol}}^{\beta} \in (0, 1), \text{ and } \mu_1 \in \Sigma$ 2: Set $N \leftarrow 0, \, \mathcal{D}_N \leftarrow \{\}, \, \mathcal{D}' \leftarrow \{\}, \, N_Y \leftarrow 0, \, Y_N \leftarrow \{\}, \, N_X \leftarrow 0, \, X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N+1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $N_Y \leftarrow N_Y + 1, Y_N \leftarrow Y_{N-1} \oplus \operatorname{span}\{p(\mu_N)\} \text{ (see (2.69))}$ 5: if $\mu_N \notin \mathcal{D}'$, then 6: $N_X \leftarrow N_X + 1, X_N \leftarrow X_{N-1} \oplus \operatorname{span}\{u(\mu_N)\}$ 7:end if 8: while (true) do 9: for all $\mu \in \Sigma$ do 10: Compute $(u_N(\mu), p_N(\mu)), \Delta_N(\mu)$, and $\hat{d}_N^{\beta}(\mu) \equiv \max\left\{\frac{\beta_{\rm Br}^{\rm UB}(\mu) - \beta_N(\mu)}{\beta_{\rm Br}^{\rm UB}(\mu)}, 0\right\}$ (cf. (2.101)) 11: 12:end for 13:
$$\begin{split} \mu'_N &\equiv \arg \max_{\mu \in \Sigma} \ \Delta_N(\mu), \ \mu^* \equiv \arg \max_{\mu \in \Sigma} \hat{d}_N^\beta(\mu) \\ \text{if } \hat{d}_N^\beta(\mu^*) < & \delta_{\text{tol}}^\beta, \text{ then} \end{split}$$
14:15: $\mu_{N+1} \equiv \mu'_N$ 16: break 17:end if 18:
$$\begin{split} & N_X \leftarrow N_X + 1 \\ & \text{if } \min_{\mu \in \mathcal{D}' \cup \mathcal{D}_N} \frac{|\mu'_N - \mu|}{|\mu|} \geq 0.1\%, \text{ then } \\ & \mathcal{D}' \leftarrow \mathcal{D}' \cup \{\mu'_N\} \end{split}$$
19:20:21: $X_N \leftarrow X_N \oplus \operatorname{span}\{u(\mu'_N)\} \text{ (see (2.79))}$ 22: 23:else $X_N \leftarrow X_N \oplus \text{span}\{T_{\mu^*} \varrho_N(\mu^*)\} \text{ (see (2.71), (2.100))}$ 24:end if 25:end while 26: 27: until $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 28: $N_{\max} \leftarrow N$

add (i) supremizer functions as in (2.78) or (ii) additional truth solutions as in (2.79). We note from (2.17) that the inf-sup constants $\beta_N(\mu)$ may be computed as the solution to a low-dimensional eigenvalue problem; the associated eigenvector is denoted by $\rho_N(\mu) \in Y_N$,

$$\varrho_N(\mu) \equiv \arg \inf_{q_N \in Y_N} \frac{\|B_N(\mu)^t q_N\|_{X'_N}}{\|q_N\|_Y}, \quad \forall \ \mu \in \mathcal{D}.$$
(2.100)

Whether the pair (X_N, Y_N) needs to be stabilized is now indicated by the relative distance $d_N^\beta(\mu)$ between $\beta_N(\mu)$ and the truth constant $\beta_{\text{Br}}(\mu)$,

$$d_N^{\beta}(\mu) \equiv \max\left\{\frac{\beta_{\mathrm{Br}}(\mu) - \beta_N(\mu)}{\beta_{\mathrm{Br}}(\mu)}, 0\right\}, \quad \forall \ \mu \in \mathcal{D}.$$
 (2.101)

We choose the parameter $\mu^* \in \Sigma$ as the argument that maximizes $d_N^\beta(\mu)$; whenever $d_N^\beta(\mu^*)$ exceeds a prescribed tolerance $\delta_{\text{tol}}^\beta \in (0, 1)$, we additionally enrich the space X_N .

First, in Algorithm 2, we enrich X_N with the supremizer function $T_{\mu^*}\varrho_N(\mu^*) \in X$ (see (2.71)), $X_N^+ \equiv X_N \oplus \text{span}\{T_{\mu^*}\varrho_N(\mu^*)\}$. With $X_N^+ \supset X_N$, the enriched pair (X_N^+, Y_N) then clearly satisfies

$$\beta_N^+(\mu) = \inf_{q_N \in Y_N} \sup_{v_N \in X_N^+} \frac{b(v_N, q_N; \mu)}{\|q_N\|_Y \|v_N\|_X} \ge \inf_{q_N \in Y_N} \sup_{v_N \in X_N} \frac{b(v_N, q_N; \mu)}{\|q_N\|_Y \|v_N\|_X} = \beta_N(\mu),$$

for all $\mu \in \mathcal{D}$. From (2.71), it moreover follows that

$$\frac{\|B_N^+(\mu^*)^t \varrho_N(\mu^*)\|_{(X_N^+)'}}{\|\varrho_N(\mu^*)\|_Y} = \frac{\|B(\mu^*)^t \varrho_N(\mu^*)\|_{X'}}{\|\varrho_N(\mu^*)\|_Y} \ge \beta_{\mathrm{Br}}(\mu^*),$$

where the inequality holds by the definition of $\beta_{\text{Br}}(\mu)$ in (2.11). As long as $d_N^{\beta,+}(\mu^*) > 0$, we therefore obtain $\varrho_N^+(\mu^*) \neq \varrho_N(\mu^*)$ and thus $\beta_N^+(\mu^*) > \beta_N(\mu^*)$ whenever $\beta_N(\mu^*)$ is a simple eigenvalue. Second, in Algorithm 3, we aim to stabilize the pair (X_N, Y_N) by enriching X_N with additional truth solutions $u(\mu'_N)$. Only if this is not possible do we resort to the supremizer functions as in Algorithm 2.

2.8 Concluding Remarks

We present in this chapter the abstract framework of a new reduced basis approach for saddle point problems. Based on Brezzi's theory [12], we not only derive new rigorous *a posteriori* bounds for the errors in the reduced basis approximations, but also present a new option to construct stable reduced basis approximation spaces. A new sampling procedure for saddle point problems then assembles computationally efficient reduced basis approximation spaces by recognizing the need for stabilization adaptively.

The developed rigorous reduced basis *a posteriori* error bounds exhibit significant advantages over existing reduced basis error estimates based on Babuška's theory for general noncoercive problems. First, they do not rely on the highly expensive Offline stage necessary for the efficient Online calculation of lower bounds to the truth Babuška inf-sup constants, but instead on much less expensive calculations associated with the truth continuity, coercivity, and Brezzi inf-sup constants. Second, as separate upper bounds for the errors in the approximations of the primal variable and the Lagrange multiplier, they enable the systematic estimation of engineering outputs depending on either of the two.

Symmetric saddle point problems are discussed as a special case in which the *a posteriori* error bounds may be further sharpened. The analysis presents possible techniques but clearly does not claim to be exhaustive. For example, as current reduced basis techniques allow their exact computation, *a posteriori* error bounds are here exclusively formulated in terms of the residual dual norms (2.31). However, the symmetric case allows us to consider the dual norm of the residual (2.32) with respect to the energy norm $\|\cdot\|_{X,\mu} = \sqrt{a(\cdot,\cdot;\mu)}$,

$$\|r_N^1(\cdot;\mu)\|_{X',\mu} \equiv \sup_{v \in X} \frac{r_N^1(v;\mu)}{\|v\|_{X,\mu}}, \quad \forall \ \mu \in \mathcal{D}.$$

Using the same techniques as presented in the proof of Proposition 2.4.6, the errors $e_N^u(\mu)$ and $e_N^p(\mu)$ in the reduced basis approximations for the primal and Lagrange multiplier variables may then be bounded in terms of $\|r_N^1(\cdot;\mu)\|_{X',\mu}$, $\|r_N^2(\cdot;\mu)\|_{Y'}$, and $\tilde{\beta}(\mu)$ (see (2.60)): For any $\mu \in \mathcal{D}$ and $N \in \mathbb{N}_{\text{max}}$, we can derive that

$$\|e_N^u(\mu)\|_{X,\mu} \le \|r_N^1(\cdot;\mu)\|_{X',\mu} + \frac{1}{\tilde{\beta}(\mu)}\|r_N^2(\cdot;\mu)\|_{Y'},$$
(2.102)

$$\|e_N^p(\mu)\|_Y \le \frac{2}{\tilde{\beta}(\mu)} \|r_N^1(\cdot;\mu)\|_{X',\mu} + \frac{1}{(\tilde{\beta}(\mu))^2} \|r_N^2(\cdot;\mu)\|_{Y'}.$$
(2.103)

For these to provide useful error bounds in the reduced basis context, $||r_N^2(\cdot;\mu)||_{X',\mu}$ and $\tilde{\beta}(\mu)$ need to be estimated Online-efficiently. Through (2.10) and (2.61), this may be done in terms of $||r_N^1(\cdot;\mu)||_{X'}$, $\alpha_a(\mu)$, $\gamma_a(\mu)$, and $\beta_{Br}(\mu)$; in this case, (2.102) leads to (2.55) as well and (2.103) results in an upper bound worse than (2.56).

Even though the reduced basis method is based on the affine decomposition (2.7) of the involved operators, problems with nonaffine parameter dependencies can also be handled efficiently: In this case, techniques as proposed in [7, 40] may be applied to recover the setting of §2.2.

Chapter 3

Approximation of the Stokes Equations

3.1 Introduction

In this chapter, we apply the reduced basis methodology developed in Chapter 2 to a Stokes flow in a parametrized domain.

Though often performed as a step to the more general Navier–Stokes equations, the analysis of the Stokes equations remains relevant in many engineering applications. In the field of microfluidics (see, e.g., [14, 50]), fluids are often geometrically constrained to a submillimeter scale; for example, lab-on-a-chip devices deal with the handling of extremely small fluid volumes down to less than pico liters. At these scales, flows often exhibit Reynolds numbers that are small enough for inertial effects to be irrelevant and they can therefore be described by the Stokes equations (see $\S1.4.2$).

The chapter is organized as follows: In §3.2, we introduce our Stokes model problem. Motivated by applications in the field of microfluidics, we consider a Stokes flow in a twodimensional microchannel with a parametrized rectangular obstacle. We shall state the strong as well as the weak formulation of the problem and describe the mapping procedure that eventually yields the formulation of a parametrized saddle point problem as introduced in §2.2. In §3.3, numerical results then demonstrate: (i) the need to appropriately enrich the reduced basis approximation space for the primal variable; (ii) the significant effects of different enrichment stategies; (iii) the rapid convergence of (stable) reduced basis approximations; and (iv) the advantages of the proposed reduced basis a *posteriori* error bounds with respect to sharpness and computational cost. Finally, in §3.4, we give some concluding remarks.

3.2 Model Problem

Many microfluidic devices with applications in chemistry, biology, or medicine require the mixing and dispersion of reagents (see, e.g., [98, 99]). An important area in microfluidic research is therefore concerned with the design of so-called micromixers [18] enhancing the mixing between adjacent streams in a straight, smooth-walled microchannel. Due to the low value of the Reynolds number, turbulence is uncommon with the consequence that mixing is difficult: In the expected laminar flow, mixing occurs by diffusion alone, which can result in very long mixing times. This is clearly unacceptable whenever complete and rapid mixing

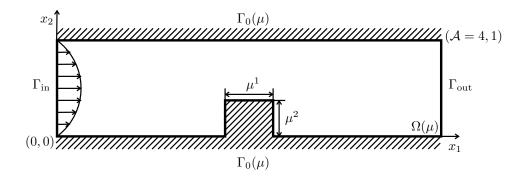


Figure 3.1.: Physical domain $\tilde{\Omega} = \Omega(\mu)$: a parametrized two-dimensional microchannel with an obstacle. The parameters of interest are width μ^1 and height μ^2 of the obstacle.

is desired, as required for the initiation of a chemical reaction. To decrease mixing time as well as mixing length, transverse flows must be generated. In this context, there are two general strategies (see [99]): active methods, where transverse flows are generated by local, oscillatory forcing within the channel; and passive methods, where transverse flows result from the interaction of the externally driven flow with the fixed channel geometry.

Our model problem is motivated by questions appearing in the design of passive micromixers. The challenge here is to generate laminar chaos in an accessible geometry, and a simple design that satisfies this criterion uses obstructions or grooves on one wall of the channel to generate transverse components in steady flows (see [99] and the references therein). Questions of interest then address the exact placing and design of such obstructions. We therefore assume the following simplified setting: We consider a Stokes or creeping flow in a two-dimensional channel with a rectangular obstacle as illustrated in Fig. 3.1. The problem depends on two geometric parameters $\mu \equiv (\mu^1, \mu^2)$ representing width μ^1 and height μ^2 of the obstacle $\mathcal{O}(\mu) \equiv [(\mathcal{A} - \mu^1)/2, (\mathcal{A} + \mu^1)/2] \times [0, \mu^2]$, which are assumed to be varying in the prescribed compact parameter set $\mathcal{D} \equiv [0.1, 0.5]^2 \subset \mathbb{R}^2$. The channel has a fixed aspect ratio, $\mathcal{A} \equiv 4$. The physical domain is thus $\tilde{\Omega} \equiv \Omega(\mu) \equiv ((0, \mathcal{A}) \times (0, 1)) \setminus \mathcal{O}(\mu)$ with its boundary $\tilde{\Gamma} \equiv \Gamma(\mu)$. We further assume fully developed flow conditions with a parabolic velocity profile on the inflow boundary $\Gamma_{\text{in}} \equiv \{0\} \times [0, 1]$, natural outflow conditions on $\Gamma_{\text{out}} \equiv \{4\} \times [0, 1]$, and no-penetration and no-slip velocity conditions on the channel walls and obstacle boundary $\tilde{\Gamma}_0 \equiv \Gamma_0(\mu) \equiv \Gamma(\mu) \setminus (\Gamma_{\text{in}} \cup \Gamma_{\text{out}})$.

3.2.1 Strong Formulation

For the physical domain $\tilde{\Omega}$, $\mu \in \mathcal{D}$, we now seek to find the (inhomogeneous) velocity $\tilde{u}_{e,inh} : \tilde{\Omega} \to \mathbb{R}^2$ and the pressure $\tilde{p}_e : \tilde{\Omega} \to \mathbb{R}$ such that

$$-\tilde{\Delta}\tilde{u}_{\rm e,inh} + \tilde{\nabla}\tilde{p}_{\rm e} = 0 \quad \text{in } \tilde{\Omega}, \tag{3.1}$$

$$\tilde{\nabla} \cdot \tilde{u}_{e,\text{inh}} = 0 \quad \text{in } \tilde{\Omega}, \tag{3.2}$$

with boundary conditions

$$\tilde{u}_{e,inh} = h \quad \text{on } \Gamma_{in}, \quad \tilde{u}_{e,inh} = 0 \quad \text{on } \tilde{\Gamma}_0, \quad \frac{\partial \tilde{u}_{e,inh}}{\partial \tilde{n}} = \tilde{p}_e \tilde{n} \quad \text{on } \tilde{\Gamma}_{out};$$
(3.3)

here, $\tilde{\Delta}$ and $\tilde{\nabla}$ are the Laplacian and gradient operator over the physical domain $\tilde{\Omega}$ (see Appendix A), respectively, \tilde{n} denotes the unit outward normal, and $h : \mathbb{R}^2 \to \mathbb{R}^2$ is given by $h(x) = h(x_1, x_2) \equiv (4x_2(1-x_2), 0)$ for all $x = (x_1, x_2) \in \mathbb{R}^2$.

3.2.2 Weak Formulation

We first introduce the function spaces

$$\begin{split} \tilde{X}_{\mathrm{e,inh}} &\equiv \{ \, \tilde{v} \in (H^1(\tilde{\Omega}))^2 \mid \tilde{v}|_{\Gamma_{\mathrm{in}}} = h, \, \tilde{v}|_{\tilde{\Gamma}_0} = 0 \, \}, \\ \tilde{X}_{\mathrm{e}} &\equiv \{ \, \tilde{v} \in (H^1(\tilde{\Omega}))^2 \mid \tilde{v}|_{\Gamma_{\mathrm{in}} \cup \tilde{\Gamma}_0} = 0 \, \}, \\ \tilde{Y}_{\mathrm{e}} &\equiv L^2(\tilde{\Omega}), \end{split}$$

where $H^1(\tilde{\Omega}) = \{ \tilde{v} \in L^2(\tilde{\Omega}) \mid \tilde{\nabla} \tilde{v} \in (L^2(\tilde{\Omega}))^2 \}$, and $L^2(\tilde{\Omega})$ is the space of square integrable functions over $\tilde{\Omega}$; for more details on Sobolev spaces, we refer to [1]. We then consider the following weak form of the problem (3.1)–(3.3): We seek to find $\tilde{u}_{e,inh} \in \tilde{X}_{e,inh}$ and $\tilde{p}_e \in \tilde{Y}_e$ such that

$$\begin{split} \tilde{a}(\tilde{u}_{\mathrm{e,inh}},\tilde{v}) + b(\tilde{v},\tilde{p}_{\mathrm{e}}) &= 0, \quad \forall \; \tilde{v} \in X_{\mathrm{e}}, \\ \tilde{b}(\tilde{u}_{\mathrm{e,inh}},\tilde{q}) &= 0, \quad \forall \; \tilde{q} \in \tilde{Y}_{\mathrm{e}}, \end{split}$$

where the bilinear forms $\tilde{a} : (H^1(\tilde{\Omega}))^2 \times (H^1(\tilde{\Omega}))^2 \to \mathbb{R}$ and $\tilde{b} : (H^1(\tilde{\Omega}))^2 \times \tilde{Y}_e \to \mathbb{R}$ are defined as

$$\begin{split} \tilde{a}(\tilde{u},\tilde{v}) &\equiv \int_{\tilde{\Omega}} \frac{\partial \tilde{u}_i}{\partial \tilde{x}_j} \frac{\partial \tilde{v}_i}{\partial \tilde{x}_j} d\tilde{x}, \quad \forall \; \tilde{u}, \tilde{v} \in (H^1(\tilde{\Omega}))^2, \\ \tilde{b}(\tilde{v},\tilde{q}) &\equiv -\int_{\tilde{\Omega}} \tilde{q} \frac{\partial \tilde{v}_i}{\partial \tilde{x}_i} d\tilde{x}, \quad \forall \; \tilde{v} \in (H^1(\tilde{\Omega}))^2, \; \tilde{q} \in \tilde{Y}_{\rm e}; \end{split}$$

here, repeated indices imply summation.

3.2.3 Formulation as a Parametrized Saddle Point Problem

We now introduce the lifting function $\tilde{u}_{\rm L} \in X_{\rm e,inh}$ where

$$\tilde{u}_{\rm L}(\tilde{x}_1, \tilde{x}_2) \equiv \begin{cases} (4\tilde{x}_2(1 - \tilde{x}_2)(1 - \tilde{x}_1), 0), & \text{in } \Omega_{\rm L} \equiv [0, 1]^2, \\ (0, 0), & \text{in } \tilde{\Omega} \setminus \Omega_{\rm L}. \end{cases}$$
(3.4)

Note that $\Omega_{\rm L}$ does not include the μ -dependent part of the domain and $\tilde{u}_{\rm L}$ is therefore by construction parameter-independent. The inhomogeneous velocity solution $\tilde{u}_{\rm e,inh}$ may then be decomposed into $\tilde{u}_{\rm e,inh} = \tilde{u}_{\rm L} + \tilde{u}_{\rm e}$, where $\tilde{u}_{\rm e} : \tilde{\Omega} \to \mathbb{R}^2$ satisfies homogeneous Dirichlet boundary conditions on $\Gamma_{\rm in} \cup \tilde{\Gamma}_0$. The governing equations for the (homogeneous) velocity $\tilde{u}_{\rm e}$ and the pressure $\tilde{p}_{\rm e}$ are then: We find $\tilde{u}_{\rm e} \in \tilde{X}_{\rm e}$ and $\tilde{p}_{\rm e} \in \tilde{Y}_{\rm e}$ such that

$$\tilde{a}(\tilde{u}_{e}, \tilde{v}) + \tilde{b}(\tilde{v}, \tilde{p}_{e}) = \tilde{f}(\tilde{v}), \quad \forall \ \tilde{v} \in \tilde{X}_{e}, \\ \tilde{b}(\tilde{u}_{e}, \tilde{q}) = \tilde{g}(\tilde{q}), \quad \forall \ \tilde{q} \in \tilde{Y}_{e},$$

$$(3.5)$$

where the right-hand side is given by $\tilde{f} \equiv -\tilde{a}(\tilde{u}_{\rm L}, \cdot) \in \tilde{X}'_{\rm e}$ and $\tilde{g} \equiv -\tilde{b}(\tilde{u}_{\rm L}, \cdot) \in \tilde{Y}'_{\rm e}$.

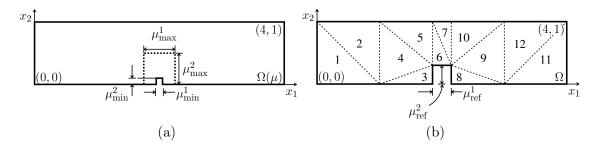


Figure 3.2.: (a) Parametrized domain $\Omega(\mu)$ for $\mu = \mu_{\min} = (0.1, 0.1)$ (solid line) and $\mu = \mu_{\max} = (0.5, 0.5)$ (dotted line); (b) reference domain Ω with $\mu_{\text{ref}} = (\mu_{\text{ref}}^1, \mu_{\text{ref}}^2) = (0.3, 0.3)$, and corresponding decomposition $\overline{\Omega} = \bigcup_{s=1}^{S} \overline{\Omega^s}$, S = 12 (see (3.6)).

To finally obtain a parametrized saddle point problem as defined in §2.2.1, we now transform the problem statement (3.5) to an equivalent problem posed over a parameterindependent reference domain Ω . To this end, $\tilde{\Omega} = \Omega(\mu)$ is traced back to $\Omega \equiv \Omega(\mu_{\text{ref}})$ by a continuous, piecewise affine mapping; for our model problem, we choose $\mu_{\text{ref}} \equiv (0.3, 0.3) \in \mathcal{D}$. The boundary of Ω is then denoted by $\Gamma \equiv \Gamma_0 \cup \Gamma_{\text{in}} \cup \Gamma_{\text{out}}$ where $\Gamma_0 \equiv \Gamma_0(\mu_{\text{ref}})$. For all $\mu \in \mathcal{D}$, the physical domain $\tilde{\Omega} = \Omega(\mu)$ is decomposed into $S \in \mathbb{N}$ subdomains $\tilde{\Omega}^s = \Omega^s(\mu)$, $1 \leq s \leq S$,

$$\overline{\tilde{\Omega}} = \bigcup_{s=1}^{S} \overline{\tilde{\Omega}^s}, \quad \text{such that } \tilde{\Omega}^s \cap \tilde{\Omega}^r = \emptyset, \quad \forall \ 1 \le s < r \le S;$$

for our model problem, this is demonstrated in Fig. 3.2. We in particular obtain

$$\overline{\Omega} = \bigcup_{s=1}^{S} \overline{\Omega^{s}}, \quad \text{where } \Omega^{s} \equiv \Omega^{s}(\mu_{\text{ref}}) \text{ and } \Omega^{s} \cap \Omega^{r} = \emptyset, \quad \forall \ 1 \le s < r \le S,$$
(3.6)

and for any $\mu \in \mathcal{D}$, $1 \leq s \leq S$, the subdomain $\tilde{\Omega}^s$ may be traced back to Ω^s by an affine coordinate transformation $\tilde{\Phi}^s \equiv \Phi^s(\cdot;\mu) : \tilde{\Omega}^s \to \Omega^s$, $\Phi^s(\tilde{x};\mu) \equiv A^s(\mu)\tilde{x} + b^s(\mu)$, where $A^s(\mu) \in \mathbb{R}^{2\times 2}$ nonsingular and $b^s(\mu) \in \mathbb{R}^2$. A more detailed discussion of affine geometry transformations can be found in [95].

We then apply standard techniques to transform the problem statement (3.5) to an equivalent problem statement formulated over the reference domain Ω . We denote the function spaces associated with Ω by X_e and Y_e , respectively. On these spaces, we consider the inner products and norms

$$(u,v)_{X_{e}} \equiv \int_{\Omega} \frac{\partial u_{i}}{\partial x_{j}} \frac{\partial v_{i}}{\partial x_{j}} dx, \quad \forall u,v \in X_{e}, \quad \|\cdot\|_{X_{e}} = \sqrt{(\cdot,\cdot)_{X_{e}}} dx,$$
$$(p,q)_{Y_{e}} \equiv \int_{\Omega} p q dx, \quad \forall p,q \in Y_{e}, \quad \|\cdot\|_{Y_{e}} = \sqrt{(\cdot,\cdot)_{Y_{e}}}.$$

By setting $v(x) \equiv \tilde{v}(\tilde{x})$ and $q(x) \equiv \tilde{q}(\tilde{x})$ for all $x \in \Omega$, we particularly obtain isomorphisms $\tilde{X}_{e} \to X_{e}, \tilde{v} \mapsto v$, and $\tilde{Y}_{e} \to Y_{e}, \tilde{q} \mapsto q$ (see, e.g., [1]). For any $\mu \in \mathcal{D}$, the problem (3.5) is thus equivalent to: Find $u_{e}(\mu) \in X_{e}$ and $p_{e}(\mu) \in Y_{e}$ such that

$$a(u_{e}(\mu), v; \mu) + b(v, p_{e}(\mu); \mu) = f(v; \mu), \quad \forall v \in X_{e}, b(u_{e}(\mu), q; \mu) = g(q; \mu), \quad \forall q \in Y_{e},$$
(3.7)

where the bilinear and linear forms are given by

$$a(u,v;\mu) = \sum_{s=1}^{S} \frac{A_{kj}^{s}(\mu)A_{lj}^{s}(\mu)}{|\det(A^{s}(\mu))|} \int_{\Omega^{s}} \frac{\partial u_{i}}{\partial x_{k}} \frac{\partial v_{i}}{\partial x_{l}} dx, \quad \forall u,v \in X_{e},$$
(3.8)

$$b(v,q;\mu) = -\sum_{s=1}^{S} \frac{A_{ji}^{s}(\mu)}{|\det(A^{s}(\mu))|} \int_{\Omega^{s}} q \frac{\partial v_{i}}{\partial x_{j}} dx, \quad \forall v \in X_{e}, \ q \in Y_{e},$$
(3.9)

$$f(v;\mu) = f(v) = -\int_{\Omega_{\rm L}} \frac{\partial u_{\rm Li}}{\partial x_j} \frac{\partial v_i}{\partial x_j} dx, \quad \forall v \in X_{\rm e},$$
(3.10)

$$g(q;\mu) = g(q) = \int_{\Omega_{\rm L}} q \frac{\partial u_{\rm Li}}{\partial x_i} \, dx, \quad \forall \ q \in Y_{\rm e}.$$
(3.11)

The bilinear forms $a(\cdot, \cdot; \mu)$ and $b(\cdot, \cdot; \mu)$ now satisfy all the assumptions (2.1), (2.2), (2.3), (2.4), and (2.7) (see, e.g., [13, 36, 101]), and we find ourselves exactly in the setting introduced in §2.2.1. We particularly have $Q_a = 10$, $Q_b = 6$, and $Q_f = Q_g = 1$ in the respective μ -affine expansions (2.7).

We choose the truth approximation spaces X and Y in §2.2.2 as the standard conforming \mathbb{P}_2 - \mathbb{P}_1 (quadratic-linear) Taylor-Hood finite element approximation subspaces [100] over a regular triangulation \mathcal{T}_{Ω} of Ω . It is a well-known result (see, e.g., [13, 27, 36, 89]) that in this case, X and Y indeed satisfy the LBB inf-sup condition (2.11).

3.3 Numerical Results

We now apply the reduced basis methodology developed in 2.3-2.7 to our model problem. In this section, all numerical results are attained using the open source software **rb00mit** [65], an implementation of the reduced basis framework within the C++ parallel finite element library **libMesh** [62].

3.3.1 Truth Approximation

Based on a fine mesh with 16,602 elements, the truth system (2.13) has a dimension of $\mathcal{N} = 72,076$. Figure 3.3 shows the velocity streamlines for three different parameter values

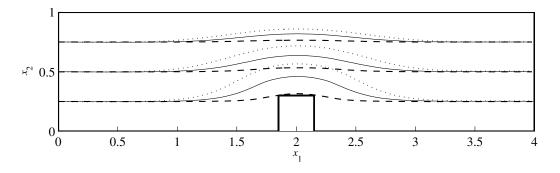


Figure 3.3.: Velocity streamlines for $\mu = (0.1, 0.1)$ (dashed), $\mu = (0.3, 0.3)$ (solid), and $\mu = (0.5, 0.5)$ (dotted) plotted on the reference domain Ω .

 $\mu = (0.1, 0.1), \mu = (0.3, 0.3), \text{ and } \mu = (0.5, 0.5).$ We observe that the μ -dependence over our parameter domain is clearly nontrivial.

3.3.2 Reduced Basis Approximation Spaces

We here present numerical results for the different options introduced in §2.5 to construct the reduced basis approximation spaces X_N and Y_N . For this purpose, the parameter samples \mathcal{D}_N and \mathcal{D}'_N (see §2.5) here have been chosen randomly.

First, we demonstrate in practice that a reduced basis approximation space constructed by using Option 0 (see (2.70)) may indeed not provide useful approximations. For our model problem, we obtain values for the inf-sup constants $\beta_N(\mu)$ that are small but still positive (see Fig. 3.4(a)); the system (2.14) is uniquely solvable, and we obtain approximations $u_N(\mu)$ and $p_N(\mu)$ for $u(\mu)$ and $p(\mu)$, respectively. We thus find ourselves in the situation (ii) described in §2.5 and indeed observe the predicted locking phenomenon: There exists only one solution, $B_N(\mu)^{-1}g_N(\mu)$, satisfying the constraints (2.16) (see Fig. 3.4(b)). Consequently, as we increase N, values of the residual dual norm $\|r_N^2(\cdot;\mu)\|_{Y'}$ associated with the constraints become relatively small and negligible in comparison to large values of $\|r_N^1(\cdot;\mu)\|_{X'}$ (see Fig. 3.5(a)); our approximations $u_N(\mu)$ do not converge satisfactorily, and thus, as suggested by (2.46), neither does $p_N(\mu)$ (see Fig. 3.5(b)).

Constructing the reduced basis approximation spaces X_N, Y_N via Option 1 (see (2.77)), Table 3.1(a) contains values of the inf-sup constants $\beta_N(\mu)$ for several parameter values $\mu \in \mathcal{D}$: We clearly see that they are essentially constant with respect to N and indeed satisfy (2.75). Table 3.1(b) contains the same data when using Option 2 (see (2.78)): We observe that the values of the inf-sup constants $\beta_N(\mu)$ here differ only insignificantly from the values obtained by using Option 1; they are also essentially constant in N and satisfy (2.75). In practice, Option 2 thus provides stable approximation spaces satisfying (2.75) at a fraction of the computational cost of Option 1 (see Fig. 3.6).

1	0	
(a)) Op	tion 1

μ	$\beta_{ m Br}(\mu)$	$eta_5(\mu)$	$\beta_{10}(\mu)$	$\beta_{20}(\mu)$	$eta_{30}(\mu)$	$\beta_{40}(\mu)$
(0.1, 0.1)	$1.073 \cdot 10^{-1}$	$1.138 \cdot 10^{-1}$	$1.137 \cdot 10^{-1}$	$1.135 \cdot 10^{-1}$	$1.135 \cdot 10^{-1}$	$1.135 \cdot 10^{-1}$
(0.5, 0.1)		$1.134 \cdot 10^{-1}$		$1.132 \cdot 10^{-1}$	$1.132 \cdot 10^{-1}$	$1.131 \cdot 10^{-1}$
(0.3, 0.3)		$9.964 \cdot 10^{-2}$			0.000	
(0.1, 0.5)			$8.088 \cdot 10^{-2}$			$8.083 \cdot 10^{-2}$
(0.5, 0.5)	$7.684 \cdot 10^{-2}$	$7.905 \cdot 10^{-2}$	$7.896 \cdot 10^{-2}$	$7.875 \cdot 10^{-2}$	$7.872 \cdot 10^{-2}$	$7.871 \cdot 10^{-2}$

(b) Option 2

μ	$\beta_{ m Br}(\mu)$	$eta_5(\mu)$	$eta_{10}(\mu)$	$\beta_{20}(\mu)$	$eta_{30}(\mu)$	$eta_{40}(\mu)$
(0.1, 0.1)	$1.073 \cdot 10^{-1}$	$1.137 \cdot 10^{-1}$	$1.136 \cdot 10^{-1}$	$1.135 \cdot 10^{-1}$	$1.134 \cdot 10^{-1}$	$1.134 \cdot 10^{-1}$
(0.5, 0.1)	$1.070 \cdot 10^{-1}$		$1.133 \cdot 10^{-1}$	$1.132 \cdot 10^{-1}$		$1.131 \cdot 10^{-1}$
(0.3, 0.3)	$9.583 \cdot 10^{-2}$					$9.957 \cdot 10^{-2}$
(0.1, 0.5)	$7.881 \cdot 10^{-2}$			$8.049 \cdot 10^{-2}$		$8.072 \cdot 10^{-2}$
(0.5, 0.5)	$7.684 \cdot 10^{-2}$	$7.862 \cdot 10^{-2}$	$7.862 \cdot 10^{-2}$	$7.848 \cdot 10^{-2}$	$7.858 \cdot 10^{-2}$	$7.858 \cdot 10^{-2}$

Table 3.1.: Values for the inf-sup constants $\beta_N(\mu)$ (see (2.17)) for several parameter values $\mu \in \mathcal{D} = [0.1, 0.5]^2$ and N evaluated using (a) Option 1 (see (2.77)) and (b) Option 2 (see (2.78)) based on a random sample \mathcal{D}_N of parameter snapshots.

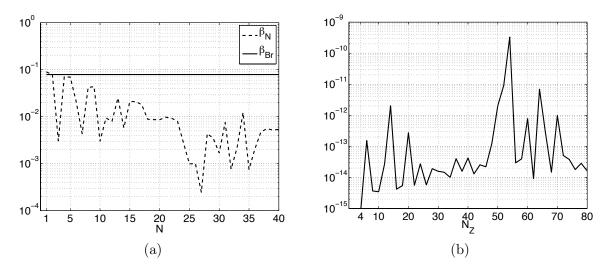


Figure 3.4.: Numerical results for $\mu = (0.1, 0.5)$ using Option 0 (see (2.70)) based on a random sample \mathcal{D}_N of parameter snapshots: (a) the inf-sup constant $\beta_N(\mu)$ (see (2.17)) shown as a function of N together with the truth inf-sup constant $\beta_{\mathrm{Br}}(\mu)$ (see (2.11)); (b) the (absolute) distance $||u_N(\mu) - B_N(\mu)^{-1}g_N(\mu)||_X$ shown as a function of the total dimension N_Z .

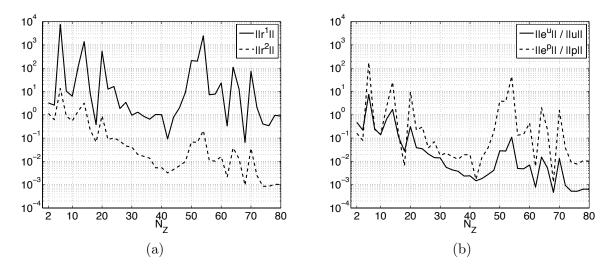


Figure 3.5.: Numerical results for $\mu = (0.1, 0.5)$ using Option 0 (see (2.70)) based on a random sample \mathcal{D}_N of parameter snapshots: (a) the residual dual norms $\|r_N^1(\cdot;\mu)\|_{X'}$, $\|r_N^2(\cdot;\mu)\|_{Y'}$ (see (2.31)) and (b) the relative errors $\|e_N^u(\mu)\|_X/\|u(\mu)\|_X$, $\|e_N^p(\mu)\|_Y/\|p(\mu)\|_Y$ (see (2.18)) shown as functions of the total dimension N_Z .

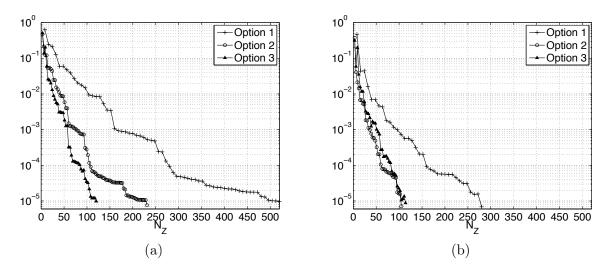


Figure 3.6.: Maximum relative errors (a) $\|e_N^u(\mu)\|_X/\|u(\mu)\|_X$ and (b) $\|e_N^p(\mu)\|_Y/\|p(\mu)\|_Y$ (see (2.18)) shown as functions of the total dimension N_Z computed using Option 1, Option 2, and Option 3 (see Table 3.2) based on random samples \mathcal{D}_N and \mathcal{D}'_N of parameter snapshots; the maximum is taken over 25 parameter values.

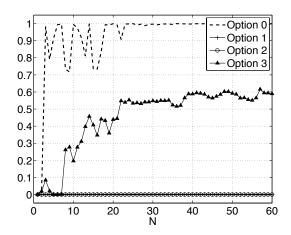


Figure 3.7.: Maximum relative distance $d_N^\beta(\mu)$ (see (2.101)) shown as a function of N computed using Option 0, Option 1, Option 2, and Option 3 (see Table 3.2) based on random samples \mathcal{D}_N and \mathcal{D}'_N of parameter snapshots; the maximum is taken over 25 parameter values.

Figure 3.7 shows the relative distance $d_N^{\beta}(\mu)$ (see (2.101)) between $\beta_N(\mu)$ and the truth constant $\beta_{\text{Br}}(\mu)$. We recall that a relative distance $d_N^{\beta}(\mu) < 1$ implies that $\beta_N(\mu)$ is positive; $d_N^{\beta}(\mu) > 0$ indicates that $\beta_N(\mu) < \beta_{\text{Br}}(\mu)$. In case of Option 1 or Option 2, we thus obtain $d_N^{\beta}(\mu) = 0$ for all $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, due to (2.75). Option 3 (see (2.79)) provides in practice inf-sup constants $\beta_N(\mu)$ that are clearly positive but generally do not satisfy (2.75).

Figure 3.6 then verifies in practice what has been predicted by a priori and a posteriori error estimates (see §2.5): Small inf-sup constants $\beta_N(\mu)$ have more dramatic effects on $\|e_N^p(\mu)\|_Y$ than on $\|e_N^u(\mu)\|_X$. The error $\|e_N^u(\mu)\|_X$ profits from a large space ker $(B_N(\mu))$ but does not explicitly depend on $\beta_N(\mu)$. Therefore, it is not surprising that adding supremizers as in Options 1 and 2, providing through (2.75) comparatively large values for the inf-sup constants $\beta_N(\mu)$, particularly favors the approximations for the Lagrange multiplier. Here, the reduced basis approximations $p_N(\mu)$ for $p(\mu)$ converge much more rapidly than the approximations $u_N(\mu)$ for the primal variable $u(\mu)$ (see Fig. 3.6). Compared to Options 1 and 2, approximation spaces by using Option 3 (see Fig. 3.6(a)). Moreover, in (2.46) and (2.20), the smaller errors $\|e_N^u(\mu)\|_X$ seem to essentially compensate for the effects of the smaller inf-sup constants $\beta_N(\mu)$: Approximations for the Lagrange multiplier are not significantly worse (see Fig. 3.6(b)).

Table 3.2 summarizes all options to construct the reduced basis approximation spaces together with their main respective properties.

Option	X_N	Special properties
0	$X_N^0 \equiv \text{span}\{ u(\mu_n) \mid 1 \le n \le N \}$ (2.70)	$\beta_N(\mu) > 0$ with $d_N^\beta(\mu)$ large $\dim(X_N) = \dim(Y_N)$
1	$X_N^1 \equiv X_N^0 \oplus \text{span}\{ T^k \xi_n \mid 1 \le n \le N, 1 \le k \le Q_b \} $ (2.77)	$\beta_N(\mu) \ge \beta_{\mathrm{Br}}(\mu)$ (in theory and practice)
2	$X_N^2 \equiv X_N^0 \oplus \text{span}\{ T_{\mu_n} \xi_n \mid 1 \le n \le N \} $ (2.78)	$\beta_N(\mu) \ge \beta_{\mathrm{Br}}(\mu) \text{ (in practice)} \ \dim(X_N^2) \ll \dim(X_N^1)$
3	$X_N^3 \equiv X_N^0 \oplus \text{span}\{ u(\mu'_n) \mid 1 \le n \le N \} $ (2.79)	$\beta_N(\mu) > 0$ with $d_N^{\beta}(\mu)$ small good space ker $(B_N(\mu))$

Table 3.2.: Options to construct the reduced basis approximation space X_N for the primal variable, where the reduced basis approximation space Y_N for the Lagrange multiplier is given as in (2.69).

3.3.3 Adaptive Sampling Procedures

We now discuss numerical results for the adaptive sampling procedures developed in §2.7.

First, we note that our Stokes model problem is clearly symmetric and therefore, $\Delta_N^{u,\text{sym}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$, and $\tilde{\Delta}_N^{u,\text{sym}}(\mu)$ represent valid upper bounds for the errors in our reduced basis approximations (see §2.4). All sampling procedures presented in this chapter are now based on an exhaustive random sample $\Sigma \subset \mathcal{D}$ of size $|\Sigma| = 4,900$ and the relative error bound $\Delta_N(\mu) \equiv \tilde{\Delta}_N^{u,\text{sym}}(\mu)/||u_N(\mu)||_X$ (see (2.67)).

Figure 3.8 shows results for Algorithm 2 using different values for $\delta_{tol}^{\beta} \in (0,1)$; a direct

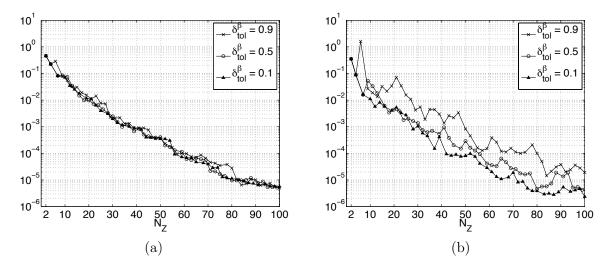


Figure 3.8.: Numerical results for Algorithm 2 (see §2.7) using $\delta_{tol}^{\beta} \in \{0.9, 0.5, 0.1\}$: Maximum relative errors (a) $\|e_N^u(\mu)\|_X/\|u(\mu)\|_X$ and (b) $\|e_N^p(\mu)\|_Y/\|p(\mu)\|_Y$ (see (2.18)) are shown as functions of N_Z ; the maximum is taken over 25 parameter values.

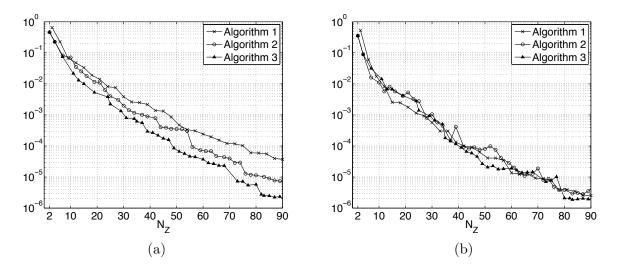


Figure 3.9.: Comparison of Algorithm 2 and Algorithm 3 (both using $\delta_{tol}^{\beta} = 0.1$) to the standard greedy procedure in Algorithm 1 (see §2.7): Maximum relative errors (a) $\|e_N^u(\mu)\|_X/\|u(\mu)\|_X$ and (b) $\|e_N^p(\mu)\|_Y/\|p(\mu)\|_Y$ (see (2.18)) are shown as functions of N_Z ; the maximum is taken over 25 parameter values.

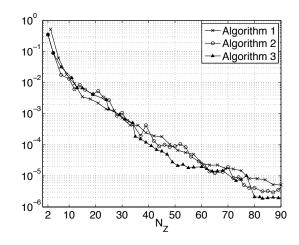


Figure 3.10.: Comparison of Algorithm 2 and Algorithm 3 (both using $\delta_{tol}^{\beta} = 0.1$) to the standard greedy procedure in Algorithm 1 (see §2.7): Maximum relative errors $||e_N(\mu)||_Z/||(u(\mu), p(\mu))||_Z$ (see (2.18)) are shown as functions of N_Z ; the maximum is taken over 25 parameter values.

comparison with Algorithm 3 is then given in Fig. 3.9 and Fig. 3.10. Again, it is verified in practice what a priori and a posteriori error bounds in §2.3.3 and Remark 2.4.4 indicate: Small inf-sup constants $\beta_N(\mu)$ have much stronger effects on $\|e_N^p(\mu)\|_Y$ than on $\|e_N^u(\mu)\|_X$. Even for fairly large tolerances δ_{tol}^{β} , we obtain approximation spaces X_N, Y_N providing very accurate approximations $u_N(\mu)$ that do not necessarily further improve for lower values of δ_{tol}^{β} (see Fig. 3.8(a)). In contrast, accurate approximations for the Lagrange multiplier require a smaller tolerance δ_{tol}^{β} and improve considerably in the dependence of decreasing values of δ_{tol}^{β} (see Fig. 3.8(b)). The errors $\|e_N^u(\mu)\|_X$ benefit primarily from a good testing space ker $(B(\mu)) \cap X_N \subseteq \text{ker}(B_N(\mu))$ (see §2.5) as achieved in Algorithm 3 (see Fig. 3.9(a)).

Figure 3.9 and Figure 3.10 moreover demonstrate that the standard greedy procedure (see Algorithm 1) may indeed generally be inefficient: For our model problem, Algorithm 2 provides accurate approximations for the primal variable at much less computational cost; using Algorithm 3, we obtain even greater computational savings compared to Algorithm 1 of up to 35%.

3.3.4 A Posteriori Error Bounds

We have thus seen how different constructions of the reduced basis approximation spaces affect the convergence of the reduced basis approximations. However, as the associated truth solutions are usually not computed, the errors in our reduced basis approximations are unknown quantities that may only be estimated by *a posteriori* error bounds. We now comment on the sharpness and computational efficiency of the rigorous *a posteriori* error bounds developed in §2.4.

We first compare the proposed *a posteriori* error bounds with respect to sharpness. Figure 3.11 and Figure 3.12 show the maximum errors in the reduced basis approximations for the primal variable and the Lagrange multiplier, respectively, together with the respective

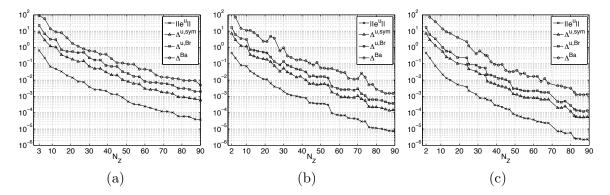


Figure 3.11.: Maximum relative error $\|e_N^u(\mu)\|_X/\|u(\mu)\|_X$ and maximum relative error bounds $\Delta_N^{u, \text{sym}}(\mu)/\|u(\mu)\|_X$, $\Delta_N^{u, \text{Br}}(\mu)/\|u(\mu)\|_X$, and $\Delta_N^{\text{Ba}}(\mu)/\|u(\mu)\|_X$ (see (2.18), (2.55), (2.37), and (2.35)) shown as functions of N_Z for (a) Algorithm 1, (b) Algorithm 2 with $\delta_{\text{tol}}^{\beta} = 0.1$, and (c) Algorithm 3 with $\delta_{\text{tol}}^{\beta} = 0.1$ (see §2.7); the maximum is taken over 25 parameter values; the computation of the error bounds is based on the exact constants $\alpha_a(\mu)$, $\gamma_a(\mu)$, $\beta_{\text{Br}}(\mu)$, and $\beta_{\text{Ba}}(\mu)$.

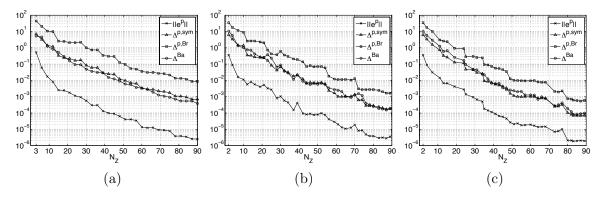


Figure 3.12.: Maximum relative error $\|e_N^p(\mu)\|_Y/\|p(\mu)\|_Y$ and maximum relative error bounds $\Delta_N^{p,\mathrm{sym}}(\mu)/\|p(\mu)\|_Y$, $\Delta_N^{p,\mathrm{Br}}(\mu)/\|p(\mu)\|_Y$, and $\Delta_N^{\mathrm{Ba}}(\mu)/\|p(\mu)\|_Y$ (see (2.18), (2.56), (2.38), and (2.35)) shown as functions of N_Z for (a) Algorithm 1, (b) Algorithm 2 with $\delta_{\mathrm{tol}}^{\beta} = 0.1$, and (c) Algorithm 3 with $\delta_{\mathrm{tol}}^{\beta} = 0.1$ (see §2.7); the maximum is taken over 25 parameter values; the computation of the error bounds is based on the exact constants $\alpha_a(\mu)$, $\gamma_a(\mu)$, $\beta_{\mathrm{Br}}(\mu)$, and $\beta_{\mathrm{Ba}}(\mu)$.

(a)	Algorithm	1
()		

`	/ -										
N	N_Z	$\tilde{\eta}_N^{u,\mathrm{sym}}$	$\eta_N^{u,{\rm sym}}$	$\eta_N^{u,{\rm Br}}$	$\eta_N^{u,{\rm Ba}}$	$\eta_N^{p,{\rm sym}}$	$\eta_N^{p,\mathrm{Br}}$	$\eta_N^{p,\mathrm{Ba}}$	$\eta_N^{\rm sym}$	$\eta_N^{ m Br}$	η_N^{Ba}
5	15	13.21	20.68	58.04	353.6	202.1	1534	296.4	148.0	1238	203.6
10	30	12.34	19.94	63.93	277.5	292.5	2268	326.5	168.6	1559	186.7
15	45	13.22	19.57	57.42	309.9	331.5	3223	417.2	201.2	1598	228.6
20	60	13.29	19.35	58.40	255.0	355.9	3240	415.7	208.1	1580	207.1
25	75	11.97	18.77	57.92	242.9	405.8	3513	476.3	198.4	1560	216.4
30	90	13.55	19.93	54.23	282.3	399.2	3423	453.6	227.0	1678	232.6
35	105	13.62	20.44	53.93	281.7	432.1	2502	519.8	235.2	1535	234.9
40	120	14.28	21.53	57.49	295.4	482.6	3317	495.2	253.0	1526	253.1
(b) Algoi	rithm 2									
\overline{N}	N_Z	$\tilde{\eta}_N^{u,\mathrm{sym}}$	$\eta_N^{u,{\rm sym}}$	$\eta_N^{u,{\rm Br}}$	$\eta_N^{u,\mathrm{Ba}}$	$\eta_N^{p,{\rm sym}}$	$\eta_N^{p,{\rm Br}}$	$\eta_N^{p,{\rm Ba}}$	$\eta_N^{\rm sym}$	η_N^{Br}	η_N^{Ba}
5	12	16.11	23.89	55.18	715.1	98.47	756.3	142.5	88.04	673.4	139.7
10	23	22.45	32.53	62.65	1894	121.8	799.9	234.8	99.72	700.7	233.0
15	34	13.57	25.31	68.15	425.1	152.8	766.6	253.9	130.7	660.6	216.6
20	46	13.17	20.31	61.41	505.3	224.0	1268	307.0	172.7	1081	236.1
25	56	16.44	27.51	63.04	1013	97.71	550.1	258.9	94.87	523.1	250.8
30	68	14.63	23.48	61.14	583.5	237.9	1287	294.2	189.6	1023	233.8
35	78	14.80	22.35	59.63	707.3	126.3	843.6	269.0	118.3	778.2	251.5
40	89	14.13	23.70	61.52	613.1	162.4	800.6	285.4	139.3	685.6	244.3
(c)) Algor	ithm 3									
N	N_Z	$\tilde{\eta}_N^{u,\mathrm{sym}}$	$\eta_N^{u,{\rm sym}}$	$\eta_N^{u,{\rm Br}}$	$\eta_N^{u,\mathrm{Ba}}$	$\eta_N^{p,\mathrm{sym}}$	$\eta_N^{p,{\rm Br}}$	$\eta_N^{p,{\rm Ba}}$	$\eta_N^{\rm sym}$	η_N^{Br}	η_N^{Ba}
5	13	18.36	32.56	68.23	1146	85.38	793.3	162.2	79.09	755.4	157.9
10	29	22.08	33.22	72.73	717.1	183.0	996.8	219.0	152.9	830.9	182.5
15	39	15.56	26.58	63.64	446.0	137.4	721.9	195.0	123.8	649.1	173.1
20	49	15.13	24.88	69.96	699.3	118.1	807.8	174.9	104.2	753.7	162.0
25	60	19.45	33.86	71.13	1034	130.3	577.4	206.5	115.1	509.1	195.4
30	73	20.00	40.18	97.51	976.6	102.5	582.7	187.7	94.86	538.1	179.6
35	83	16.90	28.72	71.93	821.5	136.3	692.0	217.4	124.6	631.3	198.3
40	93	17.87	33.31	82.45	973.4	196.7	1012	301.9	165.4	849.5	288.4
					ı						

Table 3.3.: Maximum effectivities $\tilde{\eta}_N^{u,\text{sym}}(\mu) = \tilde{\Delta}_N^{u,\text{sym}}(\mu)/||e_N^u(\mu)||_{X,\mu}$ (see (2.68)), $\eta_N^{u,\text{sym}}(\mu) = \Delta_N^{u,\text{sym}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{sym}}(\mu) \equiv \Delta_N^{p,\text{sym}}(\mu)/||e_N^p(\mu)||_Y$, $\eta_N^{sym}(\mu) = \Delta_N^{sym}(\mu)/||e_N(\mu)||_Z$ (see (2.57) and (2.66)), and $\eta_N^{u,\text{Br}}(\mu) \equiv \Delta_N^{u,\text{Br}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{Br}}(\mu) \equiv \Delta_N^{p,\text{Br}}(\mu)/||e_N^p(\mu)||_Y$, $\eta_N^{\text{Br}}(\mu) = \Delta_N^{\text{Br}}(\mu)/||e_N(\mu)||_Z$, $\eta_N^{u,\text{Ba}}(\mu) \equiv \Delta_N^{\text{Ba}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{Ba}}(\mu) \equiv \Delta_N^{\text{Ba}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{Ba}}(\mu) \equiv \Delta_N^{\text{Ba}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{Ba}}(\mu) \equiv \Delta_N^{\text{Ba}}(\mu)/||e_N^u(\mu)||_X$, $\eta_N^{p,\text{Ba}}(\mu) \equiv \Delta_N^{\text{Ba}}(\mu)/||e_N^u(\mu)||_X$ (see (2.39), (2.36), and (2.48)) for (a) Algorithm 1, (b) Algorithm 2 with $\delta_{\text{tol}}^{\beta} = 0.1$, and (c) Algorithm 3 with $\delta_{\text{tol}}^{\beta} = 0.1$ (see §2.7); the maximum is taken over 25 parameter values; the computation of the error bounds is based on the exact constants $\alpha_a(\mu)$, $\gamma_a(\mu)$, $\beta_{\text{Br}}(\mu)$, and $\beta_{\text{Ba}}(\mu)$.

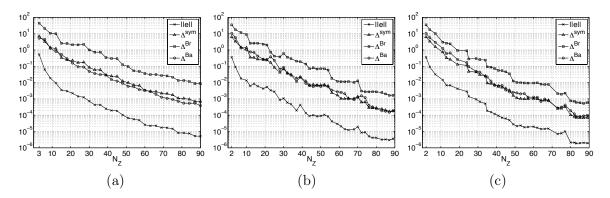


Figure 3.13.: Maximum relative error $||e_N(\mu)||_Z/||(u(\mu), p(\mu))||_Z$ and maximum relative error bounds $\Delta_N^{\text{sym}}(\mu)/||(u(\mu), p(\mu))||_Z$, $\Delta_N^{\text{Br}}(\mu)/||(u(\mu), p(\mu))||_Z$, and $\Delta_N^{\text{Ba}}(\mu)/||(u(\mu), p(\mu))||_Z$ (see (2.18), (2.65), (2.47), and (2.35)) shown as functions of N_Z for (a) Algorithm 1, (b) Algorithm 2 with $\delta_{\text{tol}}^{\beta} = 0.1$, and (c) Algorithm 3 with $\delta_{\text{tol}}^{\beta} = 0.1$ (see §2.7); the maximum is taken over 25 parameter values; the computation of the error bounds is based on the exact constants $\alpha_a(\mu), \gamma_a(\mu), \beta_{\text{Br}}(\mu)$, and $\beta_{\text{Ba}}(\mu)$.

error bounds $\Delta_N^{u,\text{sym}}(\mu)$, $\Delta_N^{u,\text{Br}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$, $\Delta_N^{p,\text{Br}}(\mu)$, and $\Delta_N^{\text{Ba}}(\mu)$ (see §2.4); Figure 3.13 shows the maximum total error in the reduced basis approximation and associated error bounds $\Delta_N^{\text{sym}}(\mu)$, $\Delta_N^{\text{Br}}(\mu)$, and $\Delta_N^{\text{Ba}}(\mu)$ (see §2.4). The maximum is computed over a sample of 25 parameter values. For this sample, to analyze only the effects of the a posteriori error bound formulations and eliminate contributions of the SCM, the exact constants $\alpha_a(\mu)$, $\gamma_a(\mu), \beta_{\rm Br}(\mu), \text{ and } \beta_{\rm Ba}(\mu)$ rather than the lower/upper bounds (2.29), (2.30) are used. Effectivities associated with the error bounds are given in Table 3.3. With maximum effectivities essentially ranging from 50 to 80, $\Delta_N^{u,\text{Br}}(\mu)$ represents a reasonably sharp bound for the error in the reduced basis approximation $u_N(\mu)$ (see Fig. 3.11 and Table 3.3). However, exploiting the symmetry of the problem, $\Delta_N^{u, \text{sym}}(\mu)$ provides a sharper bound not only in theory (see §2.4) but also in practice: Here, effectivities essentially vary between 20 and 40. Furthermore, it is not surprising to observe that the best results are achieved by the bound $\tilde{\Delta}_N^{u,\text{sym}}(\mu)$ that overestimates the error in the reduced basis approximations for the primal variable measured in the energy norm (see (2.68)) only by a factor of approximately 15 (see Table 3.3). In contrast, $\Delta_N^{p,\text{Br}}(\mu)$ and $\Delta_N^{p,\text{sym}}(\mu)$ estimate the error in $p_N(\mu)$ rather pessimistically (see Fig. 3.12). For $\Delta_N^{p,\text{Br}}(\mu)$, exact effectivities strongly depend on how the underlying reduced basis approximation spaces are constructed: Using Algorithm 1, we obtain effectivity values of order $3 \cdot 10^3$ (see Table 3.3(a)); in this case, $\Delta_N^{\text{Ba}}(\mu)$ clearly represents a much sharper bound. Yet, using Algorithm 2 or Algorithm 3, the difference is not as large (see Table 3.3(b) and (c)). For $\Delta_N^{p,\text{sym}}(\mu)$, we still obtain rather large effectivities with average values 362, 161, and 127 in case of Algorithm 1, Algorithm 2, and Algorithm 3, respectively. However, in neither case, do $\Delta_N^{u,\text{sym}}(\mu)$ and $\Delta_N^{p,\text{sym}}(\mu)$ perform worse than $\Delta_N^{\text{Ba}}(\mu)$ but generally provide bounds for the errors in the primal and Lagrange multiplier variables that are much more accurate (see Table 3.3).

We now turn to the computational efficiency of the *a posteriori* error bounds. The SCM

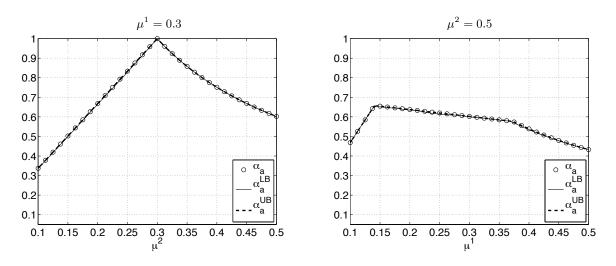


Figure 3.14.: SCM lower and upper bounds $\alpha_a^{\text{LB}}(\mu)$ and $\alpha_a^{\text{UB}}(\mu)$ for the coercivity constants $\alpha_a(\mu)$ (see (2.10), (2.29)) where $|\Xi| = 4,225$, $M_\alpha = \infty$, $\delta_{\text{tol}}^{\text{SCM}} = 0.01$, and $K_{\text{max}} = 35$.

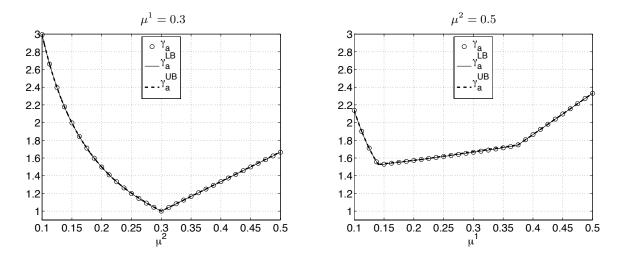


Figure 3.15.: SCM lower and upper bounds $\gamma_a^{\text{LB}}(\mu)$ and $\gamma_a^{\text{UB}}(\mu)$ for the continuity constants $\gamma_a(\mu)$ (see (2.8), (2.29)) where $|\Xi| = 4,225$, $M_\alpha = \infty$, $\delta_{\text{tol}}^{\text{SCM}} = 0.01$, and $K_{\text{max}} = 12$.

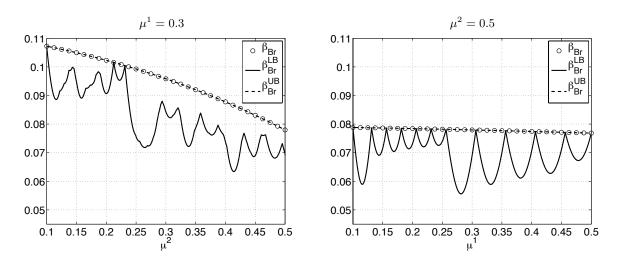


Figure 3.16.: SCM lower and upper bounds $\beta_{\text{Br}}^{\text{LB}}(\mu)$ and $\beta_{\text{Br}}^{\text{UB}}(\mu)$ for the Brezzi inf-sup constants $\beta_{\text{Br}}(\mu)$ (see (2.11), (2.30)) where $|\Xi| = 4,225$, $M_{\alpha} = \infty$, $\delta_{\text{tol}}^{\text{SCM}} = 0.05$, and $K_{\text{max}} = 163$.

(see §2.6.2) enables the (Online-)efficient estimation of the constants $\alpha_a(\mu)$, $\gamma_a(\mu)$, and $\beta_{\rm Br}(\mu)$. We here apply the method for $M_{\alpha} = \infty$ and an exhaustive sample $\Xi \subset \mathcal{D}$ of size $|\Xi| = 4,225$. We set the SCM tolerance $\delta_{\rm tol}^{\rm SCM} = 0.01$ (resp., 0.5) and obtain $K_{\rm max} = 35$ (resp., 6) for the coercivity constants $\alpha_a(\mu)$ and $K_{\rm max} = 12$ (resp., 5) for the continuity constants $\gamma_a(\mu)$. For the Brezzi inf-sup constants $\beta_{\rm Br}(\mu)$, we set $\delta_{\rm tol}^{\rm SCM} = 0.5$ (resp., 0.75) and obtain $K_{\rm max} = 163$ (resp., 111). We receive very accurate (Online-)efficient bounds $\alpha_a^{\rm LB}(\mu)$, $\gamma_a^{\rm UB}(\mu)$, and $\beta_{\rm Br}^{\rm LB}(\mu)$ (see Fig. 3.14, Fig. 3.15, and Fig. 3.16), providing a posteriori error bounds $\Delta_N^{u,{\rm Br}}(\mu)$, $\Delta_N^{u,{\rm sym}}(\mu)$, $\tilde{\Delta}_N^{u,{\rm sym}}(\mu)$ and $\Delta_N^{p,{\rm Br}}(\mu)$, $\Delta_N^{p,{\rm sym}}(\mu)$ that essentially coincide with their values based on the evaluation of the exact constants; associated effectivities thus remain the same as shown in Table 3.3.

Using a 2.66 GHz Intel Core 2 Duo processor, Offline computations necessary for the SCM applied to the Babuška inf-sup constants $\beta_{\text{Ba}}(\mu)$ are infeasible. The difficulty here lies in the computations required for the SCM bounding box \mathcal{B} (see (2.95)): The differential operator associated with the Babuška inf-sup constants exhibits a μ -affine decomposition (2.7) that involves Q terms where max $\{Q_a, Q_b\} \leq Q \leq Q_a + Q_b$; the SCM bounding box then requires the solution of Q(1+Q) generalized eigenvalue problems. (Recall that in our model problem, we have $Q_a = 10$ and $Q_b = 6$; here, Q = 15.) However, even the solution of a single generalized eigenvalue problem associated with \mathcal{B} is a computational challenge: System matrices on the left-hand side are dense, symmetric, not necessarily positive definite, and of a very low rank due to the geometric transformations involved. Consequently, even for our model problem where geometric variations are still relatively simple, we would have to resort to one of the other approaches discussed in Remark 2.6.1 to obtain (Online-)efficient lower bounds to the inf-sup constants $\beta_{\text{Ba}}(\mu)$ (see (2.30)). No matter which of the currently existing methods is used, though, the (Online-)efficient evaluation of $\Delta_N^{\text{Ba}}(\mu)$ requires either Offline costs that are prohibitively expensive compared to the ones needed for $\Delta_N^{u,\text{Br}}(\mu)$, $\Delta_N^{u,\text{sym}}(\mu)$, $\tilde{\Delta}_N^{u,\text{sym}}(\mu)$, $\Delta_N^{p,\text{Br}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$, or a significant loss in terms of accuracy

causing associated effectivities to be much worse than in Table 3.3.

Remark 3.3.1. In case of our model problem, the error bounds introduced in Remark 2.4.5 in fact essentially coincide with $\Delta_N^{u,\text{Br}}(\mu)$ and $\Delta_N^{p,\text{Br}}(\mu)$, respectively, and we therefore do not discuss them separately at this point.

3.3.5 Computation Times

We may now discuss the Online computation times for the proposed methods. For comparison, once the μ -independent parts in (2.7) have been formed, direct computation of the truth approximation $(u(\mu), p(\mu))$ (i.e., assembly and solution of (2.13)) takes on average 6.5 seconds. The rigorous and efficient error bounds $\Delta_N^{u,\text{sym}}(\mu)$ and $\Delta_N^{p,\text{sym}}(\mu)$ allow us to choose the reduced basis system dimension N_Z just large enough to obtain a desired accuracy. In case of Algorithm 1, we need $N_Z = 51$ to achieve a prescribed accuracy of roughly 1% or better in the reduced basis approximations $u_N(\mu)$ (see Fig. 3.11(a)). Once the database has been loaded, the Online calculation of $(u_N(\mu), p_N(\mu))$ (i.e., assembly and solution of (2.14)) and $\Delta_N^{u,\text{sym}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$ for any new value of $\mu \in \mathcal{D}$ takes on average 0.31 and 20.99 milliseconds, respectively, which is in total roughly 300 times faster than direct computation of the truth approximation. In the case of Algorithm 3, the same accuracy is achieved for $N_Z = 35$ (see Fig. 3.11(c)); the Online calculation of $(u_N(\mu), p_N(\mu))$ and $\Delta_N^{u,\text{sym}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$ then takes on average 0.14 and 13.29 milliseconds, respectively, and is thus roughly 480 times faster than direct computation of the truth approximation. Detailed computation times, also for Algorithm 2, are given in Table 3.4.

Method	N_Z	N	$(u_N(\mu),p_N(\mu))$	$\Delta_N^{u,\mathrm{sym}}(\mu), \Delta_N^{p,\mathrm{sym}}(\mu)$	Total		
Algorithm 1	51(78)	17(26)	$0.31 \ (0.79)$	20.99(39.60)	21.20 (40.38)		
Algorithm 2	44 (61)	19(27)	0.21 (0.42)	16.25(24.73)	16.46(25.15)		
Algorithm 3	35(55)	13(23)	0.14(0.32)	13.29(21.59)	13.43(21.92)		
(b) Accuracy of at least 1% (resp., 0.1%) for the reduced basis approximations $p_N(\mu)$							

(a) Accuracy of at least 1% (resp., 0.1%) for the reduced basis approximations $u_N(\mu)$

(a) included of a base $1/6$ (resp., $61/6$) for the reduced scale approximations $p_N(p)$							
Method	N_Z	N	$(u_N(\mu), p_N(\mu))$	$\Delta_N^{u,\mathrm{sym}}(\mu), \Delta_N^{p,\mathrm{sym}}(\mu)$	Total		
Algorithm 1	54 (84)	18 (28)	0.35(0.91)	22.75 (44.70)	23.10(45.61)		
Algorithm 2	44 (72)	19(32)	$0.21 \ (0.58)$	16.25(31.54)	16.46(32.12)		
Algorithm 3	37 (62)	14 (26)	0.16(0.44)	13.94(25.77)	14.09 (26.21)		

Table 3.4.: Average computation times in milliseconds for the Online evaluation of $(u_N(\mu), p_N(\mu))$ (assembly and solution of (2.14)) and the error bounds $\Delta_N^{u,\text{sym}}(\mu)$, $\Delta_N^{p,\text{sym}}(\mu)$ (see (2.55), (2.56)); times are measured using either Algorithm 1, Algorithm 2 with $\delta_{\text{tol}}^{\beta} = 0.1$, or Algorithm 3 with $\delta_{\text{tol}}^{\beta} = 0.1$, with a prescribed accuracy of 1% (resp., 0.1%) for the reduced basis approximations (a) $u_N(\mu)$ and (b) $p_N(\mu)$.

3.4 Concluding Remarks

We present in this chapter numerical results for the reduced basis approach proposed in Chapter 2 applied to a Stokes flow model problem. Compared to earlier approaches, the resulting methods allow us to consider flow problems in parametrized domains with notably greater ease.

The rigorous reduced basis *a posteriori* error bounds developed in §2.4 based on Brezzi's theory for saddle point problems are not only computationally more efficient than earlier error estimates based on Babuška's theory for general noncoercive problems, but numerical results demonstrate their superiority also in terms of sharpness; in particular, the bounds provided for the errors in the approximations for the primal variable are reasonably sharp.

Furthermore, numerical results underline theoretical observations in §2.5 that an enrichment of the reduced basis approximation space with supremizer functions is not necessarily optimal. Adaptively recognizing the need for stabilization, the new sampling procedures for saddle point problems developed in §2.7 construct reduced basis approximation spaces that provide accurate approximations at much less computational cost: Depending on the employed enrichment strategy, we obtain for our Stokes model problem savings compared to the standard greedy approach of up to 31%. This promises even more significant savings when extending the method to nonlinear problems such as the Navier–Stokes equations.

Considering that our Stokes flow model problem belongs to the symmetric case discussed in Chapter 2, future work may be concerned with a more detailed investigation also of the nonsymmetric case. The exact performance of the adaptive sampling procedures as well as the *a posteriori* error bounds (with respect to both computational cost and sharpness) are here neither obvious nor precisely predictable; a clarification through further numerical results may thus be of interest. Also, for simplification, our model problem only involves affine geometry variations. As already mentioned in §2.8, the proposed methods are yet not limited to this case: Techniques as proposed in [7, 40] can be applied, making the proposed approach feasible also, e.g., in the context of nonaffine geometric transformations (see [19, 94]).

Chapter 4

Approximation of Saddle Point Problems by Regularization or Penalty

4.1 Introduction

In the previous Chapters 2 and 3, we discussed the particular difficulties that appear for reduced basis approximations to saddle point problems: Special care must be taken in the construction of stable reduced basis approximation spaces (see §2.5, §2.7, §3.3.2, and §3.3.3), and rigorous reduced basis *a posteriori* error bounds (more specifically, the required lower bounds to inf-sup stability constants) become increasingly expensive and complicated as geometric variations become more complex (see §2.6.2 and §3.3.4); the latter becomes even more prominent in nonlinear problems such as the Navier–Stokes equations.

This chapter now focuses on the development of rapidly convergent reduced basis approximations and associated rigorous *a posteriori* error bounds for saddle point problems that arise from a penalty method applied to a constraint problem such as the Stokes or Navier–Stokes equations (see, e.g., [9, 51, 52] and the references therein). In contrast to reduced basis approaches presented in [19, 63, 70, 76, 88, 92, 94, 104] and the methods discussed in Chapter 2, the penalty formulation not only allows us to approximate the primal system variable as well as the Lagrange multiplier but also enables the development of rigorous error bounds for the reduced basis approximation that do not depend on inf-sup stability constants. The expensive evaluation of (lower bounds to) inf-sup constants is thus avoided and more general parametric variations can therefore be considered with relative ease. However, these benefits do not come without expense: The penalty term introduces an $O(\varepsilon)$ -error in the "truth" approximations for the primal and Lagrange multiplier variables, where ε denotes the penalty parameter.

We shall employ the developed methodology in the setting of a perturbed Stokes flow (see Chapter 5). However, we emphasize that saddle point problems as considered in this chapter may appear also in other applications such as nearly incompressible materials (see, e.g., [9, 13, 36]) where they have a physical interpretation.

The chapter is organized as follows: In §4.2, we introduce the penalty formulation. In §4.3, we define the reduced basis approximation as the Galerkin projection onto a low-dimensional reduced basis approximation space; we shall also state main properties of the reduced basis system including *a priori* convergence results. We then develop rigorous *a posteriori* error

bounds in §4.4, and discuss how to construct the low-dimensional reduced basis approximation space in §4.5; computational efficiency is achieved by the Offline-Online computational strategy summarized in §4.6. A *posteriori* error bounds derived in §4.4 and observations in §4.5 are again combined within an adaptive sampling procedure briefly explained in §4.7. Finally, in §4.8, we provide some concluding remarks.

4.2 General Problem Statement

4.2.1 Formulation

We assume the setting of a parametrized saddle point problem as introduced in §2.2.1. Furthermore, for any parameter $\mu \in \mathcal{D}$, let $c(\cdot, \cdot; \mu) : Y_e \times Y_e \to \mathbb{R}$ be a bilinear form that is continuous and coercive,

$$\gamma_c^{\mathbf{e}}(\mu) \equiv \sup_{p \in Y_{\mathbf{e}}} \sup_{q \in Y_{\mathbf{e}}} \frac{c(p,q;\mu)}{\|p\|_{Y_{\mathbf{e}}} \|q\|_{Y_{\mathbf{e}}}} < \infty, \quad \forall \ \mu \in \mathcal{D},$$
(4.1)

$$\alpha_c^{\mathrm{e}}(\mu) \equiv \inf_{q \in Y_{\mathrm{e}}} \frac{c(q, q; \mu)}{\|q\|_{Y_{\mathrm{e}}}^2} > 0, \quad \forall \ \mu \in \mathcal{D},$$

$$(4.2)$$

and which exhibits a μ -affine representation analogous to (2.7). We note that by (4.1) and (4.2), $c(\cdot, \cdot; \mu)$ provides with $\|\cdot\|_{Y_{e},\mu} \equiv \sqrt{c(\cdot, \cdot; \mu)}$ an energy norm on Y_{e} that is equivalent to $\|\cdot\|_{Y_{e}}$ for any $\mu \in \mathcal{D}$.

For a penalty parameter $\varepsilon > 0$ that will tend to zero, we now replace the problem (2.6) considered in §2.2.1 by a slightly more regular one: Given any $\mu \in \mathcal{D}$, we find $u_{e}^{\varepsilon}(\mu) \in X_{e}$ and $p_{e}^{\varepsilon}(\mu) \in Y_{e}$ such that

$$a(u_{\rm e}^{\varepsilon}(\mu), v; \mu) + b(v, p_{\rm e}^{\varepsilon}(\mu); \mu) = f(v; \mu), \quad \forall v \in X_{\rm e}, b(u_{\rm e}^{\varepsilon}(\mu), q; \mu) - \varepsilon c(p_{\rm e}^{\varepsilon}(\mu), q; \mu) = g(q; \mu), \quad \forall q \in Y_{\rm e}.$$

$$(4.3)$$

By the assumptions (2.1), (2.2), (2.3), (4.1), and (4.2), the problem is well-posed and has a unique solution for any $f(\cdot; \mu) \in X'_e$ and $g(\cdot; \mu) \in Y'_e$ (see, e.g., [13]).

Moreover, as we also assumed (2.4), the penalty solution $(u_e^{\varepsilon}(\mu), p_e^{\varepsilon}(\mu))$ of (4.3) converges to the solution $(u_e(\mu), p_e(\mu))$ of the nonpenalized problem (2.6) as ε approaches zero. For details, we refer to [9, 101] and to the standard texts [13, 36]. For our specific situation, the following proposition states the main result in this context.

Proposition 4.2.1. For $\varepsilon > 0$ and $\mu \in \mathcal{D}$, let $(u_e(\mu), p_e(\mu))$ and $(u_e^{\varepsilon}(\mu), p_e^{\varepsilon}(\mu))$ denote the solutions to (2.6) and (4.3), respectively. Then, for any $\mu \in \mathcal{D}$, we have

$$\|u_{\mathbf{e}}(\mu) - u_{\mathbf{e}}^{\varepsilon}(\mu)\|_{X_{\mathbf{e}}} \le C_{2}^{\mathbf{e}}(\mu)\varepsilon, \quad \|p_{\mathbf{e}}(\mu) - p_{\mathbf{e}}^{\varepsilon}(\mu)\|_{Y_{\mathbf{e}}} \le \frac{\gamma_{a}^{\mathbf{e}}(\mu)}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)}C_{2}^{\mathbf{e}}(\mu)\varepsilon,$$

where the constant $C_2^{\rm e}(\mu) > 0$ depends only on $f(\cdot;\mu)$, $g(\cdot;\mu)$, $\gamma_a^{\rm e}(\mu)$, $\alpha_a^{\rm e}(\mu)$, $\gamma_c^{\rm e}(\mu)$, and $\beta_{\rm Br}^{\rm e}(\mu)$,

$$C_2^{\mathbf{e}}(\mu) \equiv \frac{\gamma_a^{\mathbf{e}}(\mu)\gamma_c^{\mathbf{e}}(\mu)}{\alpha_a^{\mathbf{e}}(\mu)(\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu))^2} \left(1 + \frac{\gamma_a^{\mathbf{e}}(\mu)}{\alpha_a^{\mathbf{e}}(\mu)}\right) \left(\|f(\cdot;\mu)\|_{X_{\mathbf{e}'}} + \frac{\gamma_a^{\mathbf{e}}(\mu)}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)}\|g(\cdot;\mu)\|_{Y_{\mathbf{e}'}}\right).$$

Proof. Take $\varepsilon > 0$ and let μ be any parameter in \mathcal{D} . It follows from (2.6) and (4.3) that

$$a(u_{\mathbf{e}}(\mu) - u_{\mathbf{e}}^{\varepsilon}(\mu), v; \mu) + b(v, p_{\mathbf{e}}(\mu) - p_{\mathbf{e}}^{\varepsilon}(\mu); \mu) = 0, \qquad \forall v \in X_{\mathbf{e}}, \qquad (4.4)$$

$$b(u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu), q; \mu) - \varepsilon c(p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu), q; \mu) = -\varepsilon c(p_{\rm e}(\mu), q; \mu), \quad \forall q \in Y_{\rm e}.$$
(4.5)

First, combining the inf-sup condition (2.4) with (4.4) and (2.1), we obtain that

$$\beta_{\rm Br}^{\rm e}(\mu) \| p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu) \|_{Y_{\rm e}} \leq \sup_{v \in X_{\rm e}} \frac{b(v, p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu); \mu)}{\|v\|_{X_{\rm e}}} = \sup_{v \in X_{\rm e}} \frac{-a(u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu), v; \mu)}{\|v\|_{X_{\rm e}}} \\ \leq \gamma_a^{\rm e}(\mu) \| u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu) \|_{X_{\rm e}}.$$
(4.6)

Second, setting $v = u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu)$, $q = p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu)$ in (4.4), (4.5) and subtracting the second from the first equation yields

$$\|u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu)\|_{X_{\rm e},\mu}^{2} + \varepsilon \|p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu)\|_{Y_{\rm e},\mu}^{2} = \varepsilon c(p_{\rm e}(\mu), p_{\rm e}(\mu) - p_{\rm e}^{\varepsilon}(\mu); \mu),$$
(4.7)

which provides by (2.3), (4.2), and (4.1) that

$$\alpha_a^{\mathrm{e}}(\mu) \| u_{\mathrm{e}}(\mu) - u_{\mathrm{e}}^{\varepsilon}(\mu) \|_{X_{\mathrm{e}}}^2 \leq \varepsilon \gamma_c^{\mathrm{e}}(\mu) \| p_{\mathrm{e}}(\mu) \|_{Y_{\mathrm{e}}} \| p_{\mathrm{e}}(\mu) - p_{\mathrm{e}}^{\varepsilon}(\mu) \|_{Y_{\mathrm{e}}}.$$

We may then apply (4.6) and divide by $\alpha_a^{\rm e}(\mu) \| u_{\rm e}(\mu) - u_{\rm e}^{\varepsilon}(\mu) \|_{X_{\rm e}}$ such that

$$\|u_{\mathbf{e}}(\mu) - u_{\mathbf{e}}^{\varepsilon}(\mu)\|_{X_{\mathbf{e}}} \leq \frac{\gamma_{a}^{\mathbf{e}}(\mu)\gamma_{c}^{\mathbf{e}}(\mu)}{\alpha_{a}^{\mathbf{e}}(\mu)\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)}\|p_{\mathbf{e}}(\mu)\|_{Y_{\mathbf{e}}}\varepsilon.$$
(4.8)

The statement then follows from (4.6), (4.8), and the stability estimate (see [13], cf. (2.38))

$$\|p_{\mathbf{e}}(\boldsymbol{\mu})\|_{Y_{\mathbf{e}}} \leq \frac{1}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\boldsymbol{\mu})} \left(1 + \frac{\gamma_{a}^{\mathbf{e}}(\boldsymbol{\mu})}{\alpha_{a}^{\mathbf{e}}(\boldsymbol{\mu})}\right) \left(\|f(\cdot;\boldsymbol{\mu})\|_{X_{\mathbf{e}'}} + \frac{\gamma_{a}^{\mathbf{e}}(\boldsymbol{\mu})}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\boldsymbol{\mu})}\|g(\cdot;\boldsymbol{\mu})\|_{Y_{\mathbf{e}'}}\right).$$

In the special case that $a(\cdot, \cdot; \mu)$ is symmetric for all $\mu \in \mathcal{D}$, the ε -independent constants given in Proposition 4.2.1 can be sharpened.

Proposition 4.2.2. Let $a(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For $\varepsilon > 0$ and $\mu \in \mathcal{D}$, let $(u_{e}(\mu), p_{e}(\mu))$ and $(u_{e}^{\varepsilon}(\mu), p_{e}^{\varepsilon}(\mu))$ denote the solutions to (2.6) and (4.3), respectively. Then, for any $\mu \in \mathcal{D}$, we have

$$\|u_{\mathbf{e}}(\mu) - u_{\mathbf{e}}^{\varepsilon}(\mu)\|_{X_{\mathbf{e}}} \leq \frac{C_{2}^{\mathbf{e},\mathrm{sym}}(\mu)}{\sqrt{\alpha_{a}^{\mathbf{e}}(\mu)}} \varepsilon, \quad \|p_{\mathbf{e}}(\mu) - p_{\mathbf{e}}^{\varepsilon}(\mu)\|_{Y_{\mathbf{e}}} \leq \frac{\sqrt{\gamma_{a}^{\mathbf{e}}(\mu)}}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)} C_{2}^{\mathbf{e},\mathrm{sym}}(\mu) \varepsilon,$$

where the constant $C_2^{\rm e,sym}(\mu)>0$ may be expressed as

$$C_2^{\mathrm{e,sym}}(\mu) \equiv \frac{\sqrt{\gamma_a^{\mathrm{e}}(\mu)}\gamma_c^{\mathrm{e}}(\mu)}{(\beta_{\mathrm{Br}}^{\mathrm{e}}(\mu))^2} \left[\left(1 + \sqrt{\frac{\gamma_a^{\mathrm{e}}(\mu)}{\alpha_a^{\mathrm{e}}(\mu)}}\right) \|f(\cdot;\mu)\|_{X_{\mathrm{e}'}} + \frac{\gamma_a^{\mathrm{e}}(\mu)}{\beta_{\mathrm{Br}}^{\mathrm{e}}(\mu)} \|g(\cdot;\mu)\|_{Y_{\mathrm{e}'}} \right].$$

Proof. We recall that, as a continuous, symmetric, and coercive bilinear form, $a(\cdot, \cdot; \mu)$ defines an inner product on X_e for any $\mu \in \mathcal{D}$. Therefore, following the steps presented in the previous proof, we now obtain that

$$\beta_{\mathrm{Br}}^{\mathrm{e}}(\mu) \| p_{\mathrm{e}}(\mu) - p_{\mathrm{e}}^{\varepsilon}(\mu) \|_{Y_{\mathrm{e}}} \leq \sqrt{\gamma_{a}^{\mathrm{e}}(\mu)} \| u_{\mathrm{e}}(\mu) - u_{\mathrm{e}}^{\varepsilon}(\mu) \|_{X_{\mathrm{e}},\mu}$$
(4.9)

from the Cauchy-Schwarz inequality for the inner product $a(\cdot, \cdot; \mu)$. Using (4.7), (4.2), (4.1), and (4.9) then yields

$$\|u_{\mathbf{e}}(\mu) - u_{\mathbf{e}}^{\varepsilon}(\mu)\|_{X_{\mathbf{e}},\mu} \le \sqrt{\gamma_{a}^{\mathbf{e}}(\mu)} \frac{\gamma_{c}^{\mathbf{e}}(\mu)}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)} \|p_{\mathbf{e}}(\mu)\|_{Y_{\mathbf{e}}} \varepsilon,$$
(4.10)

and the statement follows from combining (4.10) and the stability estimate (cf. (2.56))

$$\|p_{\mathbf{e}}(\mu)\|_{Y_{\mathbf{e}}} \leq \frac{1}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)} \left[\left(1 + \sqrt{\frac{\gamma_{a}^{\mathbf{e}}(\mu)}{\alpha_{a}^{\mathbf{e}}(\mu)}} \right) \|f(\cdot;\mu)\|_{X_{\mathbf{e}}'} + \frac{\gamma_{a}^{\mathbf{e}}(\mu)}{\beta_{\mathrm{Br}}^{\mathbf{e}}(\mu)} \|g(\cdot;\mu)\|_{Y_{\mathbf{e}}'} \right]$$

r (2.3) or (4.9).

with either (2.3) or (4.9).

4.2.2**Truth Approximation**

As in §2.2.2, we denote by X and Y finite-dimensional subspaces of X_e and Y_e satisfying (2.8), (2.9), (2.10), and (2.11). These in particular also inherit the continuity and coercivity properties of $c(\cdot, \cdot; \mu)$,

$$\gamma_c(\mu) \equiv \sup_{p \in Y} \sup_{q \in Y} \frac{c(p,q;\mu)}{\|p\|_Y \|q\|_Y} < \infty, \quad \forall \ \mu \in \mathcal{D},$$
(4.11)

$$\alpha_c(\mu) \equiv \inf_{q \in Y} \frac{c(q, q; \mu)}{\|q\|_Y^2} > 0, \quad \forall \ \mu \in \mathcal{D};$$

$$(4.12)$$

for any $\mu \in \mathcal{D}$, $\|\cdot\|_{Y,\mu} \equiv \|\cdot\|_{Y_{e},\mu}$ thus defines a norm on Y that is equivalent to $\|\cdot\|_{Y}$. Again, we denote by \mathcal{N} the large dimension of the approximation space $Z = X \times Y$.

We now define the "truth" approximation as follows: For a sufficiently small penalty parameter $\varepsilon > 0$ and any given $\mu \in \mathcal{D}$, we look for $u^{\varepsilon}(\mu) \in X$ and $p^{\varepsilon}(\mu) \in Y$ such that

$$a(u^{\varepsilon}(\mu), v; \mu) + b(v, p^{\varepsilon}(\mu); \mu) = f(v; \mu), \quad \forall v \in X, b(u^{\varepsilon}(\mu), q; \mu) - \varepsilon c(p^{\varepsilon}(\mu), q; \mu) = g(q; \mu), \quad \forall q \in Y.$$

$$(4.13)$$

As for the exact problem in \$4.2.1, it follows from (2.8), (2.9), (2.10), (4.11), and (4.12) that the "truth" problem (4.13) is well-posed and has a unique solution for any $f(\cdot; \mu) \in X'_e$ and $g(\cdot;\mu) \in Y'_{e}$. The system (4.13) may be equivalently written as

$$A(\mu) u^{\varepsilon}(\mu) + B(\mu)^{t} p^{\varepsilon}(\mu) = f(\mu) \quad \text{in } X',$$

$$B(\mu) u^{\varepsilon}(\mu) - \varepsilon C(\mu) p^{\varepsilon}(\mu) = g(\mu) \quad \text{in } Y',$$

where $A(\mu)$, $B(\mu)$, $B(\mu)^t$, $f(\mu)$, and $g(\mu)$ are defined as in §2.2.2, and the bounded linear operator $C(\mu): Y \to Y'$ is given by

$$\langle C(\mu) \, p, q \rangle = c(p,q;\mu), \quad \forall \ p,q \in Y.$$

Since X and Y are moreover assumed to satisfy the LBB inf-sup condition (2.11), the penalty solution $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$ of (4.13) converges to the solution $(u(\mu), p(\mu))$ of the non-penalized problem (2.13) as ε approaches zero; more precisely, corresponding to Proposition 4.2.1 and Proposition 4.2.2, the following results hold true.

Proposition 4.2.3. For $\varepsilon > 0$ and $\mu \in \mathcal{D}$, let $(u(\mu), p(\mu))$ and $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$ denote the solutions to (2.13) and (4.13), respectively. Then, for any $\mu \in \mathcal{D}$, we have

$$\|u(\mu) - u^{\varepsilon}(\mu)\|_X \le C_2(\mu)\varepsilon, \quad \|p(\mu) - p^{\varepsilon}(\mu)\|_Y \le \frac{\gamma_a(\mu)}{\beta_{\mathrm{Br}}(\mu)}C_2(\mu)\varepsilon,$$

where the constant $C_2(\mu) > 0$ depends only on $f(\cdot; \mu)$, $g(\cdot; \mu)$, $\gamma_a(\mu)$, $\alpha_a(\mu)$, $\gamma_c(\mu)$, and $\beta_{Br}(\mu)$,

$$C_2(\mu) \equiv \frac{\gamma_a(\mu)\gamma_c(\mu)}{\alpha_a(\mu)(\beta_{\mathrm{Br}}(\mu))^2} \left(1 + \frac{\gamma_a(\mu)}{\alpha_a(\mu)}\right) \left(\|f(\cdot;\mu)\|_{X'} + \frac{\gamma_a(\mu)}{\beta_{\mathrm{Br}}(\mu)}\|g(\cdot;\mu)\|_{Y'}\right).$$

Proof. Replacing the spaces X_e and Y_e with X and Y, respectively, this is a direct application of Proposition 4.2.1.

Proposition 4.2.4. Let $a(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For $\varepsilon > 0$ and $\mu \in \mathcal{D}$, let $(u(\mu), p(\mu))$ and $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$ denote the solutions to (2.13) and (4.13), respectively. Then, for any $\mu \in \mathcal{D}$, we have

$$\|u(\mu) - u^{\varepsilon}(\mu)\|_{X} \le \frac{C_{2}^{\mathrm{sym}}(\mu)}{\sqrt{\alpha_{a}(\mu)}} \varepsilon, \quad \|p(\mu) - p^{\varepsilon}(\mu)\|_{Y} \le \frac{\sqrt{\gamma_{a}(\mu)}}{\beta_{\mathrm{Br}}(\mu)} C_{2}^{\mathrm{sym}}(\mu) \varepsilon,$$

where the constant $C_2^{\mathrm{sym}}(\mu) > 0$ may be expressed as

$$C_2^{\text{sym}}(\mu) \equiv \frac{\sqrt{\gamma_a(\mu)}\gamma_c(\mu)}{(\beta_{\text{Br}}(\mu))^2} \left[\left(1 + \sqrt{\frac{\gamma_a(\mu)}{\alpha_a(\mu)}} \right) \|f(\cdot;\mu)\|_{X'} + \frac{\gamma_a(\mu)}{\beta_{\text{Br}}(\mu)} \|g(\cdot;\mu)\|_{Y'} \right].$$

Proof. Replacing the spaces X_e and Y_e with X and Y, respectively, this is a direct application of Proposition 4.2.2.

We shall build our reduced basis approximation upon the "truth" discretization (4.13), and we shall measure the error in our reduced basis prediction relative to $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$.

4.3 Reduced Basis Approximation

We now turn to the reduced basis method, discussing the approximation procedure, rigorous a *posteriori* error estimators, and the effects of the penalty parameter ε .

4.3.1 Formulation

Suppose that we are given a set of nested, low-dimensional reduced basis approximation subspaces $X_N \subset X_{N+1} \subset X$ and $Y_N \subset Y_{N+1} \subset Y$, $N \in \mathbb{N}_{\max} \equiv \{1, \ldots, N_{\max}\}$, then the reduced basis approximation is again defined as the Galerkin projection with respect to the truth problem (4.13) onto these low-dimensional subspaces: For $\varepsilon > 0$ sufficiently small and any given $\mu \in \mathcal{D}$, we find $u_N^{\varepsilon}(\mu) \in X_N$ and $p_N^{\varepsilon}(\mu) \in Y_N$ such that

$$a(u_N^{\varepsilon}(\mu), v_N; \mu) + b(v_N, p_N^{\varepsilon}(\mu); \mu) = f(v_N; \mu), \quad \forall v_N \in X_N, b(u_N^{\varepsilon}(\mu), q_N; \mu) - \varepsilon c(p_N^{\varepsilon}(\mu), q_N; \mu) = g(q_N; \mu), \quad \forall q_N \in Y_N.$$

$$(4.14)$$

Written in operator notation, the discrete system reads

$$A_N(\mu) u_N^{\varepsilon}(\mu) + B_N(\mu)^t p_N^{\varepsilon}(\mu) = f_N(\mu) \quad \text{in } X'_N, \tag{4.15}$$

$$B_N(\mu) u_N^{\varepsilon}(\mu) - \varepsilon C_N(\mu) p_N^{\varepsilon}(\mu) = g_N(\mu) \quad \text{in } Y_N', \tag{4.16}$$

where $A_N(\mu)$, $B_N(\mu)$, $B_N(\mu)^t$, $f_N(\mu)$, and $g_N(\mu)$ are defined as in §2.3.1, and the operator $C_N(\mu): Y_N \to Y'_N$ is given by

$$\langle C_N(\mu) p_N, q_N \rangle = c(p_N, q_N; \mu), \quad \forall p_N, q_N \in Y_N.$$

4.3.2 Properties of the Discrete System

The introduction of the penalty term regularizes the reduced basis problem such that no further requirements for the reduced basis approximation spaces X_N and Y_N are needed for (4.14) to be well-posed. Let $\varepsilon > 0$ and $\mu \in \mathcal{D}$. As the bilinear form $c(\cdot, \cdot; \mu)$ satisfies (4.11) and (4.12), the operator $C_N(\mu)$ is nonsingular and (4.16) is equivalent to

$$p_N^{\varepsilon}(\mu) = \frac{1}{\varepsilon} C_N(\mu)^{-1} \Big(B_N(\mu) u_N^{\varepsilon}(\mu) - g_N(\mu) \Big).$$
(4.17)

Inserting this in (4.15), $u_N^{\varepsilon}(\mu)$ is given as the unique solution to the coercive problem

$$\left(A_N(\mu) + \frac{1}{\varepsilon} B_N(\mu)^t C_N(\mu)^{-1} B_N(\mu)\right) u_N^{\varepsilon}(\mu) = f_N(\mu) + \frac{1}{\varepsilon} B_N(\mu)^t C_N(\mu)^{-1} g_N(\mu) \quad \text{in } X'_N,$$

and $p_N^{\varepsilon}(\mu)$ is clearly obtained through (4.17). In contrast to our observations for (2.14) in §2.3.2 where the inf-sup condition (2.17) represents an additional requirement for X_N, Y_N that needs to be satisfied, the system (4.14) is thus uniquely solvable for any parameter $\mu \in \mathcal{D}$ for any choice of X_N, Y_N .

4.3.3 A Priori Error Estimation

We here consider a priori estimates for the errors in our reduced basis approximations.

For any $\varepsilon > 0$ and $\mu \in \mathcal{D}$, we denote the errors in the reduced basis approximations $u_N^{\varepsilon}(\mu) \in X_N, p_N^{\varepsilon}(\mu) \in Y_N$, and $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu)) \in Z_N = X_N \times Y_N$ with respect to the truth penalty approximations by

$$e_N^{u,\varepsilon}(\mu) \equiv u^{\varepsilon}(\mu) - u_N^{\varepsilon}(\mu) \in X,$$

$$e_N^{p,\varepsilon}(\mu) \equiv p^{\varepsilon}(\mu) - p_N^{\varepsilon}(\mu) \in Y,$$

$$e_N^{\varepsilon}(\mu) \equiv (e_N^{u,\varepsilon}(\mu), e_N^{p,\varepsilon}(\mu)) \in Z.$$
(4.18)

The total error $e_N^{\varepsilon}(\mu)$ shall be measured in the ε -dependent energy norm $\|\cdot\|_{Z,\varepsilon,\mu}$ given on the combined space $Z = X \times Y$ by

$$\|(v,q)\|_{Z,\varepsilon,\mu} \equiv \left(a(v,v;\mu) + \varepsilon c(q,q;\mu)\right)^{1/2} = \left(\|v\|_{X,\mu}^2 + \varepsilon \|q\|_{Y,\mu}^2\right)^{1/2}, \quad \forall \ (v,q) \in Z.$$
(4.19)

From (2.8), (2.10), (4.11), and (4.12), this norm is clearly equivalent to $\|\cdot\|_{Z}$; however, we note that this relationship depends on the penalty parameter ε ,

$$\sqrt{\min\{\alpha_a(\mu), \varepsilon \alpha_c(\mu)\}} \, \| \cdot \|_Z \le \| \cdot \|_{Z,\varepsilon,\mu} \le \sqrt{\max\{\gamma_a(\mu), \varepsilon \gamma_c(\mu)\}} \, \| \cdot \|_Z, \quad \forall \ \mu \in \mathcal{D}$$

Proposition 4.3.1. For $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{max}$, we have

$$\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu} \le C\left(\frac{1}{\sqrt{\varepsilon}},\mu\right) \inf_{(v_N,q_N)\in Z_N} \|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu)) - (v_N,q_N)\|_{Z,\varepsilon,\mu},$$
(4.20)

where the constant $C(\frac{1}{\sqrt{\varepsilon}},\mu) > 0$ may be expressed as

$$C\left(\frac{1}{\sqrt{\varepsilon}},\mu\right) \equiv 1 + \max\left\{\frac{\gamma_a(\mu)}{\alpha_a(\mu)},\frac{\gamma_c(\mu)}{\alpha_c(\mu)}\right\} + \frac{\gamma_b(\mu)}{\sqrt{\alpha_a(\mu)\alpha_c(\mu)\varepsilon}}$$

Proof. Let $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\text{max}}$. (We suppress the argument μ in this proof for clarity of exposition.) For any $v_N \in X_N$ and $q_N \in Y_N$, it follows from (4.13) and (4.14) that

$$\begin{aligned} a(v_N - u_N^{\varepsilon}, w_N) + b(w_N, q_N - p_N^{\varepsilon}) &= a(v_N - u^{\varepsilon}, w_N) + b(w_N, q_N - p^{\varepsilon}), \quad \forall w_N \in X_N, \\ b(v_N - u_N^{\varepsilon}, r_N) - \varepsilon \, c(q_N - p_N^{\varepsilon}, r_N) &= b(v_N - u^{\varepsilon}, r_N) - \varepsilon \, c(q_N - p^{\varepsilon}, r_N), \quad \forall r_N \in Y_N. \end{aligned}$$

Setting here $w_N = v_N - u_N^{\varepsilon} \in X_N$, $r_N = q_N - p_N^{\varepsilon} \in Y_N$ and subtracting the second from the first equation yields

$$\begin{aligned} \|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu}^2 &= \|v_N - u_N^{\varepsilon}\|_{X,\mu}^2 + \varepsilon \|q_N - p_N^{\varepsilon}\|_{Y,\mu}^2 \\ &= a(v_N - u^{\varepsilon}, v_N - u_N^{\varepsilon}) + b(v_N - u_N^{\varepsilon}, q_N - p^{\varepsilon}) \\ &- b(v_N - u^{\varepsilon}, q_N - p_N^{\varepsilon}) + \varepsilon c(q_N - p^{\varepsilon}, q_N - p_N^{\varepsilon}), \end{aligned}$$
(4.21)

which, from (2.8), (2.9), (2.10), (4.11), and (4.12), can be estimated by

$$\leq \frac{\gamma_a}{\alpha_a} \|v_N - u^{\varepsilon}\|_{X,\mu} \|v_N - u^{\varepsilon}_N\|_{X,\mu} + \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c}} \|v_N - u^{\varepsilon}_N\|_{X,\mu} \|q_N - p^{\varepsilon}\|_{Y,\mu} \\ + \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c}} \|v_N - u^{\varepsilon}\|_{X,\mu} \|q_N - p^{\varepsilon}_N\|_{Y,\mu} + \varepsilon \frac{\gamma_c}{\alpha_c} \|q_N - p^{\varepsilon}\|_{Y,\mu} \|q_N - p^{\varepsilon}_N\|_{Y,\mu}.$$
(4.22)

Written in matrix-vector notation, this reads

$$\|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu}^2 \le \left(\frac{\|v_N - u_N^{\varepsilon}\|_{X,\mu}}{\sqrt{\varepsilon} \|q_N - p_N^{\varepsilon}\|_{Y,\mu}}\right)^t \left(\frac{\frac{\gamma_a}{\alpha_a}}{\sqrt{\alpha_a\alpha_c\varepsilon}} \frac{\frac{\gamma_b}{\sqrt{\alpha_a\alpha_c\varepsilon}}}{\frac{\gamma_c}{\alpha_c}}\right) \left(\frac{\|v_N - u^{\varepsilon}\|_{X,\mu}}{\sqrt{\varepsilon} \|q_N - p^{\varepsilon}\|_{Y,\mu}}\right),$$

which is further bounded by

$$\leq \left(\max\left\{\frac{\gamma_a}{\alpha_a}, \frac{\gamma_c}{\alpha_c}\right\} + \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c \varepsilon}} \right) \|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu} \|(v_N - u^{\varepsilon}, q_N - p^{\varepsilon})\|_{Z,\varepsilon,\mu}$$

from the same techniques as used in the proof of Corollary 2.4.3. We finally have

$$\|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon}\|_{Z,\varepsilon,\mu} \le \left(\max\left\{\frac{\gamma_a}{\alpha_a}, \frac{\gamma_c}{\alpha_c}\right\} + \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c \varepsilon}}\right)\|(u^{\varepsilon} - v_N, p^{\varepsilon} - q_N)\|_{Z,\varepsilon,\mu}$$
(4.23)

for all $v_N \in X_N$, $q_N \in Y_N$, and (4.20) follows from the triangle inequality

$$\begin{aligned} \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu} &= \|(u^{\varepsilon} - u_N^{\varepsilon}, p^{\varepsilon} - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu} \\ &\leq \|(u^{\varepsilon} - v_N, p^{\varepsilon} - q_N)\|_{Z,\varepsilon,\mu} + \|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu} \end{aligned}$$

and (4.23).

In the case of a symmetric problem (4.3), the constant given in (4.20) can be improved; we may then derive the following *a priori* error estimate.

Proposition 4.3.2. Let $a(\cdot, \cdot; \mu)$ and $c(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\max}$, we then have

$$\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu} \le \left(2 + \frac{\gamma_b(\mu)}{\sqrt{\alpha_a(\mu)\alpha_c(\mu)\varepsilon}}\right) \inf_{(v_N,q_N)\in Z_N} \|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu)) - (v_N,q_N)\|_{Z,\varepsilon,\mu}.$$

Proof. As continuous and symmetric bilinear forms satisfying (2.10) and (4.12), $a(\cdot, \cdot; \mu)$ and $c(\cdot, \cdot; \mu)$ define inner products on X and Y, respectively, for any parameter $\mu \in \mathcal{D}$. Therefore, proceeding as in the previous proof, we may now apply the Cauchy–Schwarz inequality in (4.21) that yields

$$\begin{aligned} \|(v_N - u_N^{\varepsilon}, q_N - p_N^{\varepsilon})\|_{Z,\varepsilon,\mu}^2 \\ &\leq \|v_N - u^{\varepsilon}\|_{X,\mu} \|v_N - u_N^{\varepsilon}\|_{X,\mu} + \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c}} \|v_N - u_N^{\varepsilon}\|_{X,\mu} \|q_N - p^{\varepsilon}\|_{Y,\mu} \\ &+ \frac{\gamma_b}{\sqrt{\alpha_a \alpha_c}} \|v_N - u^{\varepsilon}\|_{X,\mu} \|q_N - p_N^{\varepsilon}\|_{Y,\mu} + \varepsilon \|q_N - p^{\varepsilon}\|_{Y,\mu} \|q_N - p_N^{\varepsilon}\|_{Y,\mu} \end{aligned}$$

instead of (4.22). The statement then follows from analogous steps as presented above. \Box

Through the introduction of the penalty term, we thus obtain a priori error estimates that do not depend on inf-sup constants. However, we note that they depend on the penalty parameter ε : The rate at which the reduced basis approximation $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ converges to the truth solution $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$ behaves as $O(\frac{1}{\sqrt{\varepsilon}})$ and thus may be poor for ε small if no further precautions are taken. Sufficient precautions are, for example, to construct the reduced basis approximation spaces X_N, Y_N such that they are stable as defined in §2.3.2; in this case, there moreover holds the following result.

Proposition 4.3.3. If the spaces X_N, Y_N are stable (see §2.3.2), we have for $\varepsilon \in (0, 1)$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\max}$,

$$\|e_N^{\varepsilon}(\mu)\|_Z \le C(\mu) \inf_{(v_N, q_N)\in Z_N} \|(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu)) - (v_N, q_N)\|_Z,$$

$$(4.24)$$

where the constant $C(\mu) > 0$ depends only on $\gamma_a(\mu)$, $\alpha_a(\mu)$, $\gamma_c(\mu)$, $\alpha_c(\mu)$, $\gamma_b(\mu)$, and $\beta_N(\mu)$.

Proof. We here essentially combine the approach in the proof of Proposition 4.3.1 with techniques used for Proposition 2.3.2. Let $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\text{max}}$. (We suppress all μ -dependence in this proof for clarity of exposition.)

Take any $v_N \in X_N$ and $q_N \in Y_N$. We start as in the proof of Proposition 4.3.1 and obtain (4.21). From (2.8), (2.10), (4.11), (4.12), and (2.9), it then follows that

$$\begin{aligned} \alpha_a \|v_N - u_N^{\varepsilon}\|_X^2 + \varepsilon \alpha_c \|q_N - p_N^{\varepsilon}\|_Y^2 \\ &\leq \gamma_a \|v_N - u^{\varepsilon}\|_X \|v_N - u_N^{\varepsilon}\|_X + \varepsilon \gamma_c \|q_N - p^{\varepsilon}\|_Y \|q_N - p_N^{\varepsilon}\|_Y \\ &+ \gamma_b \|v_N - u_N^{\varepsilon}\|_X \|q_N - p^{\varepsilon}\|_Y + \gamma_b \|v_N - u^{\varepsilon}\|_X \|q_N - p_N^{\varepsilon}\|_Y. \end{aligned}$$
(4.25)

We now proceed as in the proof of Proposition 2.3.2: As the Galerkin projection (4.14) to (4.13) onto the low-dimensional reduced basis approximation spaces X_N, Y_N , the errors $e_N^{u,\varepsilon}$ and $e_N^{p,\varepsilon}$ satisfy

$$a(e_N^{u,\varepsilon}, v_N) + b(v_N, e_N^{p,\varepsilon}) = 0, \quad \forall v_N \in X_N.$$

$$(4.26)$$

Since X_N, Y_N are stable, the inf-sup condition (2.17) and (4.26) provide that

$$\beta_N \|q_N - p_N^{\varepsilon}\|_Y \leq \sup_{v_N \in X_N} \frac{b(v_N, q_N - p_N^{\varepsilon})}{\|v_N\|_X} = \sup_{v_N \in X_N} \frac{b(v_N, q_N - p^{\varepsilon}) + b(v_N, e_N^{p,\varepsilon})}{\|v_N\|_X}$$
$$= \sup_{v_N \in X_N} \frac{b(v_N, q_N - p^{\varepsilon}) - a(e_N^{u,\varepsilon}, v_N)}{\|v_N\|_X}$$
$$\leq \gamma_b \|q_N - p^{\varepsilon}\|_Y + \gamma_a \|e_N^{u,\varepsilon}\|_X,$$
(4.27)

where the last inequality follows from (2.8) and (2.9).

Combining (4.25) and (4.27) now yields

$$\begin{aligned} \alpha_a \|v_N - u_N^{\varepsilon}\|_X^2 &+ \varepsilon \alpha_c \|q_N - p_N^{\varepsilon}\|_Y^2 \\ &\leq \gamma_a \|v_N - u^{\varepsilon}\|_X \|v_N - u_N^{\varepsilon}\|_X + \varepsilon \gamma_c \|q_N - p^{\varepsilon}\|_Y \|q_N - p_N^{\varepsilon}\|_Y \\ &+ \gamma_b \|v_N - u_N^{\varepsilon}\|_X \|q_N - p^{\varepsilon}\|_Y + \frac{\gamma_b}{\beta_N} \|v_N - u^{\varepsilon}\|_X \Big(\gamma_b \|q_N - p^{\varepsilon}\|_Y + \gamma_a \|e_N^{u,\varepsilon}\|_X \Big); \end{aligned}$$

from the triangle inequality $\|e_N^{u,\varepsilon}\|_X \leq \|u^{\varepsilon} - v_N\|_X + \|v_N - u_N^{\varepsilon}\|_X$, the above is further bounded by

$$\leq \left[\gamma_a \left(1 + \frac{\gamma_b}{\beta_N}\right) \|v_N - u^{\varepsilon}\|_X + \gamma_b \|q_N - p^{\varepsilon}\|_Y\right] \|v_N - u^{\varepsilon}_N\|_X \\ + \varepsilon \gamma_c \|q_N - p^{\varepsilon}\|_Y \|q_N - p^{\varepsilon}_N\|_Y + \frac{\gamma_b^2}{\beta_N} \|v_N - u^{\varepsilon}\|_X \|q_N - p^{\varepsilon}\|_Y + \frac{\gamma_a \gamma_b}{\beta_N} \|v_N - u^{\varepsilon}\|_X^2.$$

We apply Young's inequality with $\rho_1, \rho_2 > 0$ such that

$$(2\alpha_a - \rho_1) \|v_N - u_N^{\varepsilon}\|_X^2 + (2\varepsilon\alpha_c - \varepsilon^2\rho_2) \|q_N - p_N^{\varepsilon}\|_Y^2$$

$$\leq \frac{1}{\rho_1} \left[\gamma_a \left(1 + \frac{\gamma_b}{\beta_N} \right) \|v_N - u^{\varepsilon}\|_X + \gamma_b \|q_N - p^{\varepsilon}\|_Y \right]^2 + \frac{\gamma_c^2}{\rho_2} \|q_N - p^{\varepsilon}\|_Y^2$$

$$+ 2\frac{\gamma_b^2}{\beta_N} \|v_N - u^{\varepsilon}\|_X \|q_N - p^{\varepsilon}\|_Y + 2\frac{\gamma_a\gamma_b}{\beta_N} \|v_N - u^{\varepsilon}\|_X^2;$$

choosing $\rho_1 \equiv \alpha_a$ and $\rho_2 \equiv 2\alpha_c$, we obtain $2\varepsilon\alpha_c - \varepsilon^2\rho_2 = 2\varepsilon(1-\varepsilon)\alpha_c > 0$ for $\varepsilon \in (0,1)$ and therefore,

$$\|v_N - u_N^{\varepsilon}\|_X \le C\Big(\|u^{\varepsilon} - v_N\|_X + \|p^{\varepsilon} - q_N\|_Y\Big),\tag{4.28}$$

where C > 0 depends only on γ_a , α_a , γ_c , α_c , γ_b , and β_N . The statement thus follows from the triangle inequality, (4.27), and (4.28).

4.4 A Posteriori Error Estimation

Using the penalty formulation, we may now develop rigorous *a posteriori* error bounds that do not require the computation of lower bounds to inf-sup stability constants.

4.4.1 Ingredients

Before we are able to formulate our *a posteriori* error bounds, we again need to introduce two sets of ingredients. The first set of ingredients consists of lower (and upper bounds) to the truth coercivity constants (2.10) and (4.12),

$$\begin{array}{l}
\alpha_a^{\mathrm{LB}}(\mu) \leq \alpha_a(\mu) \leq \alpha_a^{\mathrm{UB}}(\mu), \\
\alpha_c^{\mathrm{LB}}(\mu) \leq \alpha_c(\mu) \leq \alpha_c^{\mathrm{UB}}(\mu), \\
\end{array} \quad \forall \ \mu \in \mathcal{D}.
\end{array}$$
(4.29)

The second set of ingredients consists of the dual norms of the residuals associated with the reduced basis approximation $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$,

$$\|r_N^{1,\varepsilon}(\cdot;\mu)\|_{X'} = \sup_{v \in X} \frac{r_N^{1,\varepsilon}(v;\mu)}{\|v\|_X}, \quad \|r_N^{2,\varepsilon}(\cdot;\mu)\|_{Y'} = \sup_{q \in Y} \frac{r_N^{2,\varepsilon}(q;\mu)}{\|q\|_Y}, \tag{4.30}$$

where, for all $\mu \in \mathcal{D}$, $r_N^{1,\varepsilon}(\cdot;\mu) \in X'$ and $r_N^{2,\varepsilon}(\cdot;\mu) \in Y'$ are defined as

$$f_N^{1,\varepsilon}(v;\mu) \equiv f(v;\mu) - a(u_N^{\varepsilon}(\mu),v;\mu) - b(v,p_N^{\varepsilon}(\mu);\mu), \quad \forall v \in X,$$
(4.31)

$$r_N^{2,\varepsilon}(q;\mu) \equiv g(q;\mu) - b(u_N^{\varepsilon}(\mu),q;\mu) + \varepsilon c(p_N^{\varepsilon}(\mu),q;\mu), \quad \forall \ q \in Y.$$
(4.32)

4.4.2 Error Bound Formulation

We can now state the following result.

Proposition 4.4.1. For $\varepsilon > 0$, $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, and $\alpha_a^{\text{LB}}(\mu)$, $\alpha_c^{\text{LB}}(\mu)$ satisfying (4.29), we define

$$\Delta_N^{\varepsilon}(\mu) \equiv \left(\frac{\|r_N^{1,\varepsilon}(\cdot;\mu)\|_{X'}^2}{\alpha_a^{\text{LB}}(\mu)} + \frac{\|r_N^{2,\varepsilon}(\cdot;\mu)\|_{Y'}^2}{\varepsilon\alpha_c^{\text{LB}}(\mu)}\right)^{1/2}.$$
(4.33)

Then, $\Delta_N^{\varepsilon}(\mu)$ is an upper bound for the error $e_N^{\varepsilon}(\mu)$ measured in the ε -dependent energy norm (4.19),

$$\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu} \le \Delta_N^{\varepsilon}(\mu), \quad \forall \ \mu \in \mathcal{D}, \ \forall \ N \in \mathbb{N}_{\max}.$$
(4.34)

Proof. We follow the proof of Proposition 4.1 in [35]. Let $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\text{max}}$. From (4.31), (4.32), and (4.13), the errors $e_N^{u,\varepsilon}(\mu)$ and $e_N^{p,\varepsilon}(\mu)$ satisfy the equations

$$a(e_N^{u,\varepsilon}(\mu), v; \mu) + b(v, e_N^{p,\varepsilon}(\mu); \mu) = r_N^{1,\varepsilon}(v; \mu), \quad \forall \ v \in X,$$
(4.35)

$$b(e_N^{u,\varepsilon}(\mu),q;\mu) - \varepsilon c(e_N^{p,\varepsilon}(\mu),q;\mu) = r_N^{2,\varepsilon}(q;\mu), \quad \forall q \in Y.$$
(4.36)

Setting $v = e_N^{u,\varepsilon}(\mu)$, $q = e_N^{p,\varepsilon}(\mu)$ and subtracting the second from the first equation yields

$$\begin{split} \|e_{N}^{u,\varepsilon}(\mu)\|_{X,\mu}^{2} + \varepsilon \,\|e_{N}^{p,\varepsilon}(\mu)\|_{Y,\mu}^{2} &= r_{N}^{1,\varepsilon}(e_{N}^{u,\varepsilon}(\mu);\mu) - r_{N}^{2,\varepsilon}(e_{N}^{p,\varepsilon}(\mu);\mu) \\ &\leq \|r_{N}^{1,\varepsilon}(\cdot;\mu)\|_{X'}\|e_{N}^{u,\varepsilon}(\mu)\|_{X} + \|r_{N}^{2,\varepsilon}(\cdot;\mu)\|_{Y'}\|e_{N}^{p,\varepsilon}(\mu)\|_{Y}; \end{split}$$

from (2.10) and (4.12), this is then further bounded by

$$\leq \frac{\|r_N^{1,\varepsilon}(\cdot;\mu)\|_{X'}}{\sqrt{\alpha_a(\mu)}} \|e_N^{u,\varepsilon}(\mu)\|_{X,\mu} + \frac{\|r_N^{2,\varepsilon}(\cdot;\mu)\|_{Y'}}{\sqrt{\alpha_c(\mu)}} \|e_N^{p,\varepsilon}(\mu)\|_{Y,\mu}$$
$$\leq \left(\|e_N^{u,\varepsilon}(\mu)\|_{X,\mu}^2 + \varepsilon \|e_N^{p,\varepsilon}(\mu)\|_{Y,\mu}^2\right)^{1/2} \left(\frac{\|r_N^{1,\varepsilon}(\cdot;\mu)\|_{X'}^2}{\alpha_a(\mu)} + \frac{\|r_N^{2,\varepsilon}(\cdot;\mu)\|_{Y'}^2}{\varepsilon\alpha_c(\mu)}\right)^{1/2},$$

where the second estimate is obtained from the Cauchy–Schwarz inequality for the Euclidean distance. The statement then directly follows from (4.29).

Again, as a measure of the quality of the proposed error estimators, we consider the effectivities

$$\eta_N^{\varepsilon}(\mu) \equiv \frac{\Delta_N^{\varepsilon}(\mu)}{\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu}}, \quad \varepsilon > 0, \quad \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max}.$$
(4.37)

We recall that effectivities ≥ 1 indicate that the error estimate is rigorous; effectivities close to unity indicate that the error estimate is sharp. The following corollary specifies this behavior for the error bounds derived above.

Corollary 4.4.2. For $\varepsilon \in (0,1]$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\max}$, the effectivity $\eta_N^{\varepsilon}(\mu)$ associated with the error bound (4.33) satisfies

$$1 \le \eta_N^{\varepsilon}(\mu) \le \frac{C_3(\mu)}{\sqrt{\varepsilon}},\tag{4.38}$$

where the constant $C_3(\mu) > 0$ may be expressed as

$$C_3(\mu) \equiv \left(\frac{(\gamma_b(\mu))^2}{\alpha_a^{\mathrm{LB}}(\mu)\alpha_c(\mu)} + \frac{(\gamma_b(\mu))^2}{\alpha_c^{\mathrm{LB}}(\mu)\alpha_a(\mu)} + \frac{(\gamma_a(\mu))^2}{\alpha_a^{\mathrm{LB}}(\mu)\alpha_a(\mu)} + \frac{(\gamma_c(\mu))^2}{\alpha_c^{\mathrm{LB}}(\mu)\alpha_c(\mu)}\right)^{1/2}.$$

Proof. Let $\varepsilon > 0$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\text{max}}$. (Again, we suppress the argument μ for clarity of exposition.) The first inequality in (4.38) immediately follows from (4.34) in Proposition 4.4.1. In order to derive the upper bound for the effectivity, we note from (4.35) and (4.36) that the dual norms of the residuals satisfy

$$\|r_N^{1,\varepsilon}\|_{X'} \leq \gamma_a \|e_N^{u,\varepsilon}\|_X + \gamma_b \|e_N^{p,\varepsilon}\|_Y \leq \frac{\gamma_a}{\sqrt{\alpha_a}} \|e_N^{u,\varepsilon}\|_{X,\mu} + \frac{\gamma_b}{\sqrt{\alpha_c}} \|e_N^{p,\varepsilon}\|_{Y,\mu},$$
$$\|r_N^{2,\varepsilon}\|_{Y'} \leq \gamma_b \|e_N^{u,\varepsilon}\|_X + \varepsilon \gamma_c \|e_N^{p,\varepsilon}\|_Y \leq \frac{\gamma_b}{\sqrt{\alpha_a}} \|e_N^{u,\varepsilon}\|_{X,\mu} + \varepsilon \frac{\gamma_c}{\sqrt{\alpha_c}} \|e_N^{p,\varepsilon}\|_{Y,\mu}$$

by (2.8), (2.9), (2.10), (4.11), and (4.12). Using the Cauchy–Schwarz inequality for the Euclidean distance, we obtain

$$\|r_N^{1,\varepsilon}\|_{X'} \le \left(\frac{\gamma_a^2}{\alpha_a} + \frac{\gamma_b^2}{\varepsilon\alpha_c}\right)^{1/2} \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu}, \quad \|r_N^{2,\varepsilon}\|_{Y'} \le \left(\frac{\gamma_b^2}{\alpha_a} + \varepsilon\frac{\gamma_c^2}{\alpha_c}\right)^{1/2} \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu}.$$
(4.39)

Applying this to (4.33), we have

$$\left(\Delta_N^{\varepsilon}\right)^2 \leq \left[\frac{1}{\alpha_a^{\mathrm{LB}}} \left(\frac{\gamma_a^2}{\alpha_a} + \frac{\gamma_b^2}{\varepsilon\alpha_c}\right) + \frac{1}{\varepsilon\alpha_c^{\mathrm{LB}}} \left(\varepsilon\frac{\gamma_c^2}{\alpha_c} + \frac{\gamma_b^2}{\alpha_a}\right)\right] \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu}^2$$

$$= \left[\frac{1}{\varepsilon} \left(\frac{\gamma_b^2}{\alpha_a^{\mathrm{LB}}\alpha_c} + \frac{\gamma_b^2}{\alpha_c^{\mathrm{LB}}\alpha_a}\right) + \frac{\gamma_a^2}{\alpha_a^{\mathrm{LB}}\alpha_a} + \frac{\gamma_c^2}{\alpha_c^{\mathrm{LB}}\alpha_c}\right] \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu}^2,$$

which finally yields the statement for $\varepsilon \in (0, 1]$.

In the special case of a symmetric problem, the constant given in Corollary 4.4.2 can be sharpened as follows.

Corollary 4.4.3. Let $a(\cdot, \cdot; \mu)$ and $c(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For $\varepsilon \in (0, 1]$, $\mu \in \mathcal{D}$, and $N \in \mathbb{N}_{\max}$, the effectivity $\eta_N^{\varepsilon}(\mu)$ associated with the error bound (4.33) satisfies

$$1 \le \eta_N^{\varepsilon}(\mu) \le \frac{C_3^{\text{sym}}(\mu)}{\sqrt{\varepsilon}},$$

where the constant $C_3^{\text{sym}}(\mu) > 0$ may be expressed as

$$C_{3}^{\text{sym}}(\mu) \equiv \left(\frac{(\gamma_{b}(\mu))^{2}}{\alpha_{a}^{\text{LB}}(\mu)\alpha_{c}(\mu)} + \frac{(\gamma_{b}(\mu))^{2}}{\alpha_{c}^{\text{LB}}(\mu)\alpha_{a}(\mu)} + \frac{\gamma_{a}(\mu)}{\alpha_{a}^{\text{LB}}(\mu)} + \frac{\gamma_{c}(\mu)}{\alpha_{c}^{\text{LB}}(\mu)}\right)^{1/2} \le C_{3}(\mu).$$
(4.40)

Proof. Proceeding as in the previous proof, the dual norms of the residuals may now be estimated by

$$\begin{aligned} \|r_N^{1,\varepsilon}\|_{X'} &\leq \sqrt{\gamma_a} \|e_N^{u,\varepsilon}\|_{X,\mu} + \gamma_b \|e_N^{p,\varepsilon}\|_Y \leq \sqrt{\gamma_a} \|e_N^{u,\varepsilon}\|_{X,\mu} + \frac{\gamma_b}{\sqrt{\alpha_c}} \|e_N^{p,\varepsilon}\|_{Y,\mu}, \\ \|r_N^{2,\varepsilon}\|_{Y'} &\leq \gamma_b \|e_N^{u,\varepsilon}\|_X + \varepsilon \sqrt{\gamma_c} \|e_N^{p,\varepsilon}\|_{Y,\mu} \leq \frac{\gamma_b}{\sqrt{\alpha_a}} \|e_N^{u,\varepsilon}\|_{X,\mu} + \varepsilon \sqrt{\gamma_c} \|e_N^{p,\varepsilon}\|_{Y,\mu}. \end{aligned}$$

from (4.35), (4.36), the Cauchy–Schwarz inequalities for the inner products $a(\cdot; \cdot; \mu)$ and $c(\cdot, \cdot; \mu)$, (2.8), (2.9), (4.11), and (2.10), (4.12). Therefore, (4.39) now becomes

$$\|r_N^{1,\varepsilon}\|_{X'} \le \left(\gamma_a + \frac{\gamma_b^2}{\varepsilon\alpha_c}\right)^{1/2} \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu}, \quad \|r_N^{2,\varepsilon}\|_{Y'} \le \left(\frac{\gamma_b^2}{\alpha_a} + \varepsilon\gamma_c\right)^{1/2} \|e_N^{\varepsilon}\|_{Z,\varepsilon,\mu},$$

and applying this to (4.33) again yields the desired result. Note that the constant C_3^{sym} differs from C_3 in the last two terms; the inequality in (4.40) then follows as we clearly have $\frac{\gamma_a}{\alpha_a}, \frac{\gamma_c}{\alpha_c} \geq 1$ from the definitions in (2.8), (2.10) and (4.11), (4.12).

4.5 Construction of Reduced Basis Approximation Spaces

We introduced several options to construct the low-dimensional reduced basis approximation space in §2.5. Accordingly, the reduced basis approximation space Y_N for the Lagrange multiplier is defined as the span of several truth solutions $p^{\varepsilon}(\mu_n)$, $1 \leq n \leq N$, associated with chosen parameter snapshots $\mu_1, \ldots, \mu_N \in \mathcal{D}$ (see (2.69)),

$$Y_N \equiv \operatorname{span}\{p^{\varepsilon}(\mu_n) \mid 1 \le n \le N\} = \operatorname{span}\{\xi_n \mid 1 \le n \le N\},\tag{4.41}$$

where $\xi_n \in Y$, $1 \leq n \leq N$, are $(\cdot, \cdot)_Y$ -orthonormal basis functions. The reduced basis approximation space X_N for the primal variable may then be constructed by using either Option 0, Option 1, Option 2, or Option 3; for the reader's convenience, we restate their definitions in Table 4.1 but refer to §2.5 for details on their respective properties. Again, for algebraic stability reasons, also X_N is in particular expressed by orthonormal basis functions, i.e., $X_N \equiv \text{span}\{\phi_m \mid 1 \leq m \leq N_X\}$, where $\phi_m \in X$, $1 \leq m \leq N_X$, are $(\cdot, \cdot)_X$ -orthonormal.

Option	X_N
0	$X_N^0 \equiv \operatorname{span}\{ u^{\varepsilon}(\mu_n) \mid 1 \le n \le N \} \text{ (see (2.70))}$
1	$X_N^1 \equiv X_N^0 \oplus \text{span}\{ T^k \xi_n \mid 1 \le n \le N, \ 1 \le k \le Q_b \} \text{ (see (2.77))}$
2	$X_N^2 \equiv X_N^0 \oplus \text{span}\{ T_{\mu_n} \xi_n \mid 1 \le n \le N \} \text{ (see (2.78))}$
3	$X_N^3 \equiv X_N^0 \oplus \operatorname{span}\{ u^{\varepsilon}(\mu'_n) \mid 1 \le n \le N \} \text{ (see (2.79))}$

Table 4.1.: Options to construct of the reduced basis approximation space X_N for the primal variable, where the reduced basis approximation space Y_N for the Lagrange multiplier is given as in (4.41).

We recall (see §4.3.2) that for all options considered in Table 4.1, the problem (4.14) is uniquely solvable and we obtain valid approximations $u_N^{\varepsilon}(\mu)$ and $p_N^{\varepsilon}(\mu)$ for $u^{\varepsilon}(\mu)$ and $p^{\varepsilon}(\mu)$, respectively. In contrast to the situation described in §2.5, we emphasize that also Option 0 no longer disqualifies itself *a priori* but may possibly provide useful results.

4.6 Offline-Online Computational Procedure

Again, we exploit the parameter-affine dependence (2.7) of the problem to invoke an efficient Offline-Online computational strategy.

4.6.1 Reduced Basis Approximation

Noting that for $\varepsilon > 0$ and $\mu \in \mathcal{D}$, we can expand $u_N^{\varepsilon}(\mu) \in X_N$ and $p_N^{\varepsilon}(\mu) \in Y_N$ as

$$u_N^{\varepsilon}(\mu) = \sum_{m=1}^{N_X} u_{N\,m}^{\varepsilon}(\mu)\phi_m, \quad p_N^{\varepsilon}(\mu) = \sum_{n=1}^{N_Y} p_{N\,n}^{\varepsilon}(\mu)\xi_n, \tag{4.42}$$

the reduced basis system (4.14) may be written as

$$\sum_{m=1}^{N_X} \sum_{k=1}^{Q_a} \Theta_a^k(\mu) A_{N\,im}^k u_{N\,m}^{\varepsilon}(\mu) + \sum_{n=1}^{N_Y} \sum_{k=1}^{Q_b} \Theta_b^k(\mu) B_{N\,ni}^k \, p_{N\,n}^{\varepsilon}(\mu) = \sum_{k=1}^{Q_f} \Theta_f^k(\mu) f_{N\,i}^k,$$

for all $1 \leq i \leq N_X$, and

$$\sum_{m=1}^{N_X} \sum_{k=1}^{Q_b} \Theta_b^k(\mu) B_{Njm}^k u_{Nm}^{\varepsilon}(\mu) - \varepsilon \sum_{n=1}^{N_Y} \sum_{k=1}^{Q_c} \Theta_c^k(\mu) C_{Njn}^k p_{Nn}^{\varepsilon}(\mu) = \sum_{k=1}^{Q_g} \Theta_g^k(\mu) g_{Nj}^k,$$

for all $1 \leq j \leq N_Y$; here, the μ -independent quantities A_{Nim}^k , B_{Nni}^k , f_{Ni}^k , and g_{Nj}^k are given as in §2.6.1 and

$$C_{N\,ij}^{k} \equiv c^{k}(\xi_{j},\xi_{i}), \quad 1 \le i, j \le N_{Y}, \ 1 \le k \le Q_{c}.$$

Again, the expensive Offline stage where we form and store all parameter-independent quantities enables the highly efficient Online evaluation of the reduced basis approximation for any parameter query; computational costs essentially remain the same as listed in §2.6.1.

4.6.2 A Posteriori Error Bounds

The computation of the *a posteriori* error bounds $\Delta_N^{\varepsilon}(\mu)$ again consists of two components: the calculation of the residual dual norms (4.30), and the calculation of the required lower (and upper) bounds (4.29) to the coercivity constants $\alpha_a(\mu)$ and $\alpha_c(\mu)$.

As in §2.6.2, we use the μ -affine expansions (2.7) and (4.42) to express the residual $r_N^{1,\varepsilon}(\cdot;\mu)$ defined in (4.31) as

$$\begin{split} r_N^{1,\varepsilon}(v;\mu) &= \sum_{k=1}^{Q_f} \Theta_f^k(\mu) f^k(v) - \sum_{m=1}^{N_X} \sum_{k=1}^{Q_a} u_{N\,m}^{\varepsilon}(\mu) \,\Theta_a^k(\mu) \,a^k(\phi_m,v) \\ &- \sum_{n=1}^{N_Y} \sum_{k=1}^{Q_b} p_{N\,n}^{\varepsilon}(\mu) \,\Theta_b^k(\mu) \,b^k(v,\xi_n), \quad \forall \ v \in X, \end{split}$$

and the residual $r_N^{2,\varepsilon}(\cdot;\mu)$ defined in (4.32) as

$$\begin{split} r_N^{2,\varepsilon}(q;\mu) &= \sum_{k=1}^{Q_g} \Theta_g^k(\mu) g^k(q) - \sum_{m=1}^{N_X} \sum_{k=1}^{Q_b} u_N^{\varepsilon}(\mu) \, \Theta_b^k(\mu) \, b^k(\phi_m,q) \\ &+ \varepsilon \sum_{n=1}^{N_Y} \sum_{k=1}^{Q_c} p_{N n}^{\varepsilon}(\mu) \, \Theta_c^k(\mu) \, c^k(\xi_n,q), \quad \forall \ q \in Y. \end{split}$$

Setting $Q_1 \equiv Q_f + N_X Q_a + N_Y Q_b$ and $Q_2^{\varepsilon} \equiv Q_g + N_X Q_b + N_Y Q_c$, this can be written succinctly as

$$r_N^{1,\varepsilon}(v;\mu) = \sum_{k=1}^{Q_1} \Theta_N^{1,\varepsilon,k}(\mu) r_N^{1,\varepsilon,k}(v), \quad \forall \ v \in X,$$
(4.43)

$$r_N^{2,\varepsilon}(q;\mu) = \sum_{k=1}^{Q_2^{\varepsilon}} \Theta_N^{2,\varepsilon,k}(\mu) \, r_N^{2,\varepsilon,k}(q), \quad \forall \ q \in Y,$$
(4.44)

where the coefficient functions $\Theta_N^{1,\varepsilon,k}(\mu)$, $1 \leq k \leq Q_1$, and $\Theta_N^{2,\varepsilon,k}(\mu)$, $1 \leq k \leq Q_2^{\varepsilon}$, depend on μ explicitly through the coefficient functions $\Theta_{\mathcal{I}}^k(\mu)$, $1 \leq k \leq Q_{\mathcal{I}}$, $\mathcal{I} = a, b, c, f, g$, but also implicitly through $u_N^{\varepsilon}(\mu)$ and $p_N^{\varepsilon}(\mu)$; $r_N^{1,\varepsilon,k}(\cdot) \in X'$, $1 \le k \le Q_1$, and $r_N^{2,\varepsilon,k}(\cdot) \in Y'$, $1 \le k \le Q_2^{\varepsilon}$, are parameter-independent linear functionals. We denote by $\hat{e}_N^{1,\varepsilon,k} \in X$ and $\hat{e}_N^{2,\varepsilon,k} \in Y$ the Riesz representers associated with $r_N^{1,\varepsilon,k}(\cdot)$

We denote by $\hat{e}_N^{1,\varepsilon,k} \in X$ and $\hat{e}_N^{2,\varepsilon,k} \in Y$ the Riesz representers associated with $r_N^{1,\varepsilon,k}(\cdot)$ and $r_N^{2,\varepsilon,k}(\cdot)$, respectively, which are obtained as the solutions to the μ -independent linear problems

$$(\hat{e}_N^{1,\varepsilon,k}, v)_X = r_N^{1,\varepsilon,k}(v), \quad \forall v \in X, \quad \forall 1 \le k \le Q_1,$$

$$(4.45)$$

$$(\hat{e}_N^{2,\varepsilon,k}, q)_Y = r_N^{2,\varepsilon,k}(q), \quad \forall \ q \in Y, \quad \forall \ 1 \le k \le Q_2^{\varepsilon}.$$

$$(4.46)$$

By (4.43) and (4.44), the Riesz representers $\hat{e}_N^{1,\varepsilon}(\mu) \in X$ and $\hat{e}_N^{2,\varepsilon}(\mu) \in Y$ associated with the residuals $r_N^{1,\varepsilon}(\cdot;\mu)$ and $r_N^{2,\varepsilon}(\cdot;\mu)$ are then given by

$$\hat{e}_N^{1,\varepsilon}(\mu) = \sum_{k=1}^{Q_1} \Theta_N^{1,\varepsilon,k}(\mu) \, \hat{e}_N^{1,\varepsilon,k}, \quad \hat{e}_N^{2,\varepsilon}(\mu) = \sum_{k=1}^{Q_2^\varepsilon} \Theta_N^{2,\varepsilon,k}(\mu) \, \hat{e}_N^{2,\varepsilon,k},$$

and it finally follows that

$$\|r_N^{1,\varepsilon}(\cdot;\mu)\|_{X'}^2 = (\hat{e}_N^{1,\varepsilon}(\mu), \hat{e}_N^{1,\varepsilon}(\mu))_X = \sum_{k,l=1}^{Q_1} \Theta_N^{1,\varepsilon,k}(\mu) \Theta_N^{1,\varepsilon,l}(\mu) (\hat{e}_N^{1,\varepsilon,k}, \hat{e}_N^{1,\varepsilon,l})_X,$$
(4.47)

$$\|r_N^{2,\varepsilon}(\cdot;\mu)\|_{Y'}^2 = (\hat{e}_N^{2,\varepsilon}(\mu), \hat{e}_N^{2,\varepsilon}(\mu))_Y = \sum_{k,l=1}^{Q_2^-} \Theta_N^{2,\varepsilon,k}(\mu) \Theta_N^{2,\varepsilon,l}(\mu) (\hat{e}_N^{2,\varepsilon,k}, \hat{e}_N^{2,\varepsilon,l})_Y.$$
(4.48)

We may now summarize the Offline-Online decomposition.

In the Offline stage, we find the Riesz representers $\hat{e}_N^{1,\varepsilon,k}$, $1 \le k \le Q_1$, and $\hat{e}_N^{2,\varepsilon,k}$, $1 \le k \le Q_2^{\varepsilon}$, and form the associated inner products $(\hat{e}_N^{1,\varepsilon,k}, \hat{e}_N^{1,\varepsilon,l})_X$, $1 \le k, l \le Q_1$, and $(\hat{e}_N^{2,\varepsilon,k}, \hat{e}_N^{2,\varepsilon,l})_Y$, $1 \le k, l \le Q_2^{\varepsilon}$. Again (see §2.6.2), these computations clearly depend on \mathcal{N} : We have to solve Q_1 and Q_2^{ε} linear problems of type (4.45) and (4.46), respectively, and to compute $(Q_1)^2$ and $(Q_2^{\varepsilon})^2$ inner products in X and Y. In the Online stage, given the reduced basis coefficients $u_{Nm}^{\varepsilon}(\mu)$, $1 \le m \le N_X$, and $p_{Nn}^{\varepsilon}(\mu)$, $1 \le n \le N_Y$, we then compute the coefficient functions $\Theta_N^{1,\varepsilon,k}(\mu)$, $1 \le k \le Q_1$, and $\Theta_N^{2,\varepsilon,k}(\mu)$, $1 \le k \le Q_2^{\varepsilon}$, and assemble the sums (4.47) and (4.48) at $O((Q_1)^2) = O((Q_f + N_X Q_a + N_Y Q_b)^2)$ and $O((Q_2^{\varepsilon})^2) = O((Q_g + N_X Q_b + N_Y Q_c)^2)$ operations, respectively. As desired, the operation count for the Online stage is thus independent of \mathcal{N} .

We obtain (Online-)efficient lower and upper bounds (4.29) to the coercivity constants $\alpha_a(\mu)$ and $\alpha_c(\mu)$ by the SCM methodology described in §2.6.2.

4.7 Adaptive Sampling Procedures

We discussed adaptive sampling procedures constructing reduced basis approximation spaces (X_N, Y_N) , $N \in \mathbb{N}_{\text{max}}$, that are computationally efficient in §2.7. Again, we consider the standard greedy sampling process combined with Option 2 (see §2.7, Algorithm 1) and the sampling procedures Algorithm 2 and Algorithm 3 where the need for stabilization is recognized adaptively. In addition, according to our observations in §4.5, a standard greedy

approach (see §2.7) based on Option 0 (see Table 4.1) cannot categorically be excluded from further considerations but may provide useful results; the exact sampling procedure is therefore summarized in Algorithm 4. In contrast to the other sampling processes, we here do not pursue additional stabilization techniques but construct X_N and Y_N based solely on the respective truth solutions associated with the greedily selected parameter snapshots. However, we bear in mind that for penalty parameters ε tending to zero, a priori error estimates in §4.3.3 suggest poor convergence for approximation spaces X_N, Y_N that exhibit small inf-sup constants $\beta_N(\mu)$ (see (2.17)).

Algorithm 4 Standard Greedy Algorithm (using Option 0)

1: Choose $\Sigma \subset \mathcal{D}$, $\delta_{tol} \in (0, 1)$, and $\mu_1 \in \Sigma$ 2: Set $N \leftarrow 0$, $\mathcal{D}_N \leftarrow \{\}$, $N_Y \leftarrow 0$, $Y_N \leftarrow \{\}$, $N_X \leftarrow 0$, $X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N+1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $N_Y \leftarrow N_Y + 1, Y_N \leftarrow Y_{N-1} \oplus \operatorname{span} \{ p^{\varepsilon}(\mu_N) \} (\operatorname{see} (4.41))$ 5: $N_X \leftarrow N_X + 1, X_N \leftarrow X_{N-1} \oplus \operatorname{span} \{ u^{\varepsilon}(\mu_N) \} \text{ (see §4.5, Table 4.1)}$ 6: for all $\mu \in \Sigma$ do 7: Compute $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ and $\Delta_N(\mu)$ 8: end for 9: $\mu_{N+1} \equiv \arg \max_{\mu \in \Sigma} \Delta_N(\mu)$ 10: 11: **until** $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 12: $N_{\max} \leftarrow N$

Remark 4.7.1. We here note that Algorithm 2 and Algorithm 3 involve the (Online-)efficient estimation of the truth Brezzi inf-sup constants $\beta_{Br}(\mu)$ and are therefore not optimally suited to the penalty context where we aim to avoid such computations. However, the estimated truth Brezzi inf-sup constants are only used as reference values for the reduced basis inf-sup constants $\beta_N(\mu)$ (see (2.101)). In the penalty context, such a reference could for example also be given by $\beta_{Br}(\mu_{ref})$ for a chosen reference parameter $\mu_{ref} \in \mathcal{D}$.

4.8 Concluding Remarks

We present in this chapter the abstract framework of a new reduced basis method for saddle point problems based on a penalty formulation. At the expense of a less accurate truth approximation, the proposed approach enables the circumvention of the inf-sup condition. The analysis demonstrates how this affects the construction of reduced basis approximation spaces as well as the development of *a posteriori* error bounds.

Through the introduction of the penalty term, fulfillment of the inf-sup condition (2.17) is no longer a compulsory requirement for the reduced basis approximation spaces. In case of small inf-sup constants, *a priori* error estimates predict poor convergence for penalty parameters tending to zero; however, computationally efficient construction techniques where we do not enrich the reduced basis approximation space for the primal variable cannot categorically be excluded from further considerations but may possibly provide useful results.

The developed rigorous a posteriori error bounds are independent of inf-sup stability constants and can therefore be computed very easily. However, they do depend on the penalty parameter ε : As ε decreases and we approach the nonpenalized problem, the error bounds become increasingly pessimistic. Nevertheless, with associated effectivities growing by an order of $O(\frac{1}{\sqrt{\varepsilon}})$, these effects are relatively benign.

Chapter 5

Approximation of the Stokes Equations by Regularization or Penalty

5.1 Introduction

First introduced in the context of constrained optimization, the penalty method gained popularity as a numerical solution method for the Stokes and Navier–Stokes equations in the 1980s and 1990s (see, for example, [9, 20, 52, 81, 82] and the review papers [51, 91]). Combined with reduced integration techniques [112], it allows the elimination of the pressure and circumvention of the Babuška–Brezzi inf-sup condition [9, 13].

In the previous chapter, we presented an abstract framework of how the penalty formulation may be employed in the reduced basis context. With the increase in speed and memory capacity of computing devices, the penalty approach lost popularity in the finite element community after the 1990s; in the parametric reduced basis setting, it still provides a reduction of computational effort that is significant. In this chapter, we apply the method developed in Chapter 4 to a perturbed Stokes flow in a parametrized domain. Considering the model problem introduced in §3.2, numerical results demonstrate the effects of the penalty term in practice.

The chapter is organized as follows: In §5.2, we derive the penalty formulation for the Stokes model problem introduced in §3.2. In §5.3, we then apply the reduced basis approach discussed in §4: Numerical results illustrate the effects of the penalty term on (i) the perturbed truth velocity and pressure approximations, (ii) convergence of the reduced basis approximations, (iii) computational efficiency, and (iv) error bound effectivity.

5.2 Model Problem

We again consider a Stokes flow through a two-dimensional microchannel with a parametrized rectangular obstacle as described in §3.2. We now introduce a penalty term to the continuity equation (3.2),

$$\nabla \cdot \tilde{u}_{e,\text{inh}}^{\varepsilon} = -\varepsilon \, \tilde{p}_{e}^{\varepsilon} \quad \text{in } \tilde{\Omega}, \tag{5.1}$$

for some small penalty parameter $\varepsilon > 0$. We recall that (5.1) essentially approximates the fluid as nearly incompressible. From (3.1), (5.1), and (3.3), the governing equations for the

velocity $\tilde{u}_{\rm e,inh}^{\varepsilon}$ and the pressure $\tilde{p}_{\rm e}^{\varepsilon}$ are then

$$-\tilde{\Delta}\tilde{u}_{e,\text{inh}}^{\varepsilon} + \tilde{\nabla}\tilde{p}_{e}^{\varepsilon} = 0 \quad \text{in } \tilde{\Omega},$$
(5.2)

$$\tilde{\nabla} \cdot \tilde{u}_{\rm e,inh}^{\varepsilon} + \varepsilon \, \tilde{p}_{\rm e}^{\varepsilon} = 0 \quad \text{in } \tilde{\Omega}, \tag{5.3}$$

with boundary conditions

$$\tilde{u}_{e,inh}^{\varepsilon} = h \quad \text{on } \Gamma_{in}, \quad \tilde{u}_{e,inh}^{\varepsilon} = 0 \quad \text{on } \tilde{\Gamma}_{0}, \quad \frac{\partial \tilde{u}_{e,inh}^{\varepsilon}}{\partial \tilde{n}} = \tilde{p}_{e}^{\varepsilon} \tilde{n} \quad \text{on } \tilde{\Gamma}_{out};$$
(5.4)

here, $\tilde{\Delta}$ and $\tilde{\nabla}$ again denote the Laplacian and gradient operator over the physical domain $\tilde{\Omega}$ (see Appendix A), respectively, \tilde{n} is the unit outward normal, and $h : \mathbb{R}^2 \to \mathbb{R}^2$ is given by $h(x) = h(x_1, x_2) \equiv (4x_2(1-x_2), 0)$ for all $x = (x_1, x_2) \in \mathbb{R}^2$.

Following the steps discussed in §3.2.2 and §3.2.3, the problem (5.2)-(5.4) may be written as a parametrized saddle point problem of the form (4.3). Here, for any $\mu \in \mathcal{D}$, the bilinear forms $a(\cdot, \cdot; \mu)$, $b(\cdot, \cdot; \mu)$ and the linear functionals $f(\cdot; \mu)$, $g(\cdot; \mu)$ are given as in (3.8)–(3.11); accordingly, the bilinear form $c(\cdot, \cdot; \mu) : Y_e \times Y_e \to \mathbb{R}$ represents the L^2 -inner product over the physical domain $\tilde{\Omega}$ formulated on the reference domain Ω (see §3.2.3),

$$c(p,q;\mu) = \sum_{s=1}^{S} \frac{1}{|\det(A^{s}(\mu))|} \int_{\Omega^{s}} p \, q \, dx, \quad \forall \, p,q \in Y_{e}.$$
(5.5)

The bilinear forms $a(\cdot, \cdot; \mu)$, $b(\cdot, \cdot; \mu)$, and $c(\cdot, \cdot; \mu)$ then satisfy the assumptions (2.1)–(2.4), (4.1), and (4.2), and we thus find ourselves in the setting introduced in §4.2.1. We particularly have $Q_a = 10$, $Q_b = 6$, $Q_f = Q_g = 1$, and $Q_c = 5$ in the respective μ -affine expansions (2.7). Choosing again the truth approximation spaces X and Y as the standard conforming \mathbb{P}_2 - \mathbb{P}_1 Taylor–Hood finite element approximation subspaces [100] over the regular triangulation \mathcal{T}_{Ω} , we ensure that also (2.11) is satisfied and therefore recover the situation described in §4.2.2.

5.3 Numerical Results

We now apply the reduced basis methodology developed in 4.3-4.7 to our model problem. In this section, all numerical results are attained using the open source software **rb00mit** [65], an implementation of the reduced basis framework within the C++ parallel finite element library **libMesh** [62].

5.3.1 Truth Approximation

First, we compare our perturbed truth approximations $u^{\varepsilon}(\mu)$ and $p^{\varepsilon}(\mu)$ with the solutions $u(\mu)$ and $p(\mu)$ of the nonpenalized Stokes equations ($\varepsilon = 0$).

Again based on a fine mesh with 16,602 elements, the truth system (4.13) has a dimension of $\mathcal{N} = 72,076$. Figure 5.1 shows the velocity streamlines computed using $\varepsilon = 10^{-2}, 10^{-3}, 10^{-4}$, and 10^{-5} ; for comparison, they are also plotted for $\varepsilon = 0$. We observe that our penalized truth approximation is certainly less-than-perfect, but for most engineering applications, the approximation is sufficiently accurate even for reasonably large values

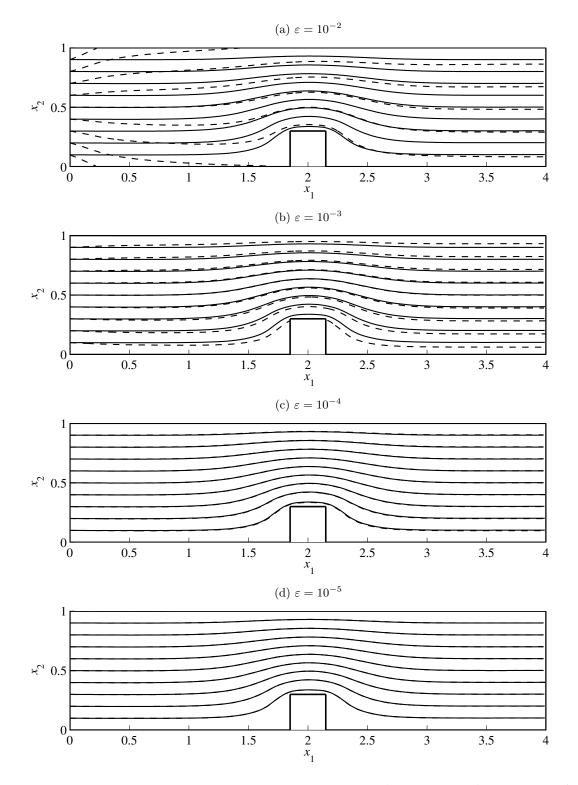


Figure 5.1.: Velocity streamlines computed using (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$; comparison with the solution for $\varepsilon = 0$ (solid).

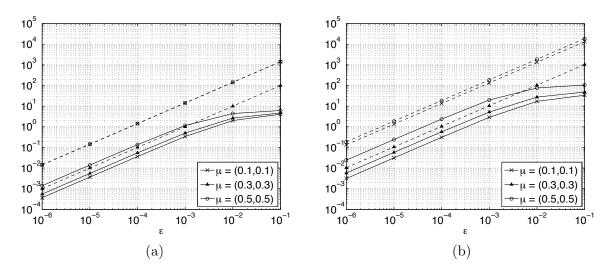


Figure 5.2.: Errors in the penalized truth approximations and associated upper bounds (see Proposition 4.2.4) as functions of the penalty parameter ε for different parameter values μ : (a) $\|u(\mu) - u^{\varepsilon}(\mu)\|_X$ (solid) and $(C_2^{\text{sym}}(\mu)/\sqrt{\alpha_a(\mu)}) \varepsilon$ (dashed), (b) $\|p(\mu) - p^{\varepsilon}(\mu)\|_Y$ (solid) and $(\sqrt{\gamma_a(\mu)}/\beta_{\text{Br}}(\mu))C_2^{\text{sym}}(\mu) \varepsilon$ (dashed).

of the penalty parameter ε . This is further illustrated in Fig. 5.2, which presents the error introduced by the penalty term as a function of ε . Moreover, the results in Fig. 5.2 clearly manifest the $O(\varepsilon)$ -dependence of the error predicted by Proposition 4.2.3 and Proposition 4.2.4.

5.3.2 Reduced Basis Approximation Spaces

We now turn to the reduced basis approximation. We here present numerical results for the different options to construct the reduced basis approximation spaces X_N and Y_N . As discussed in §4.5 and §4.7, we now build X_N and Y_N by using either Algorithm 1, Algorithm 2, Algorithm 3, or Algorithm 4. All sampling procedures are based on an exhaustive random sample $\Sigma \subset \mathcal{D}$ of size $|\Sigma| = 4,900$ and the relative error bound $\Delta_N(\mu) \equiv \Delta_N^{\varepsilon}(\mu)/||e_N^{\varepsilon}(\mu)||_{Z,\varepsilon,\mu}$ (see (4.33)); in Algorithm 2 and Algorithm 3, we further set $\delta_{tol}^{\beta} = 0.1$ (see §2.7).

We first compare the resulting reduced basis approximation spaces X_N, Y_N with respect to stability. We recall that in Algorithm 1, Algorithm 2, and Algorithm 3, the spaces are stabilized through an additional enrichment of the reduced basis approximation space X_N for the primal variable; in Algorithm 4, we abandon any special stabilization techniques. Table 5.1 now provides minimum values of the inf-sup constants $\beta_N(\mu)$ (see (2.17)) for different values of the penalty parameter ε . In case of Algorithm 1, Algorithm 2, and Algorithm 3, the inf-sup constants $\beta_N(\mu)$ are essentially constant in N with values of approximately $8 \cdot 10^{-2}$ (see Table 5.1(a), (b), and (c)); using Algorithm 4, we obtain much lower values that range from $1.5 \cdot 10^{-2}$ to $7 \cdot 10^{-5}$ (see Table 5.1(d)). Furthermore, we note that the constants $\beta_N(\mu)$ show in neither case a noticeable dependence on the penalty parameter ε .

(a) Algorithm 1

ε	$eta_5(\mu)$	$\beta_{10}(\mu)$	$\beta_{20}(\mu)$	$eta_{30}(\mu)$	$\beta_{40}(\mu)$	
		$7.899 \cdot 10^{-2}$		$7.901 \cdot 10^{-2}$		
	$7.850 \cdot 10^{-2}$	$7.853 \cdot 10^{-2}$		$7.856 \cdot 10^{-2}$		
10^{-5}	$7.846 \cdot 10^{-2}$	$7.848 \cdot 10^{-2}$	$7.844 \cdot 10^{-2}$	$7.851 \cdot 10^{-2}$	$7.855 \cdot 10^{-2}$	
(b) Algorithm 2						

(b) Highlinin 2							
ε	$eta_5(\mu)$	$\beta_{10}(\mu)$	$\beta_{20}(\mu)$	$eta_{30}(\mu)$	$eta_{40}(\mu)$		
	$7.846 \cdot 10^{-2}$						
	$7.781 \cdot 10^{-2}$						
	$7.757 \cdot 10^{-2}$						
10^{-5}	$7.754 \cdot 10^{-2}$	$7.745 \cdot 10^{-2}$	$7.735 \cdot 10^{-2}$	$7.752 \cdot 10^{-2}$	$7.761 \cdot 10^{-2}$		

(c) Algorithm 3

ε	$eta_5(\mu)$	$eta_{10}(\mu)$	$eta_{20}(\mu)$	$eta_{30}(\mu)$	$eta_{40}(\mu)$
10^{-2}	$8.089 \cdot 10^{-2}$		$7.960 \cdot 10^{-2}$		
10^{-3}	$7.610 \cdot 10^{-2}$				
10^{-4}			$7.699 \cdot 10^{-2}$		
10^{-5}	$7.682 \cdot 10^{-2}$	$7.767 \cdot 10^{-2}$	$7.742 \cdot 10^{-2}$	$7.739 \cdot 10^{-2}$	$7.725 \cdot 10^{-2}$

(d) Algorithm 4

ε	$eta_5(\mu)$	$\beta_{10}(\mu)$	$\beta_{20}(\mu)$	$eta_{30}(\mu)$	$\beta_{40}(\mu)$
10^{-2}	$2.605 \cdot 10^{-3}$	$2.733 \cdot 10^{-4}$	$9.048\cdot10^{-4}$	$2.456 \cdot 10^{-4}$	
10^{-3}					
10^{-4}		$2.369 \cdot 10^{-2}$	$1.571 \cdot 10^{-3}$		
10^{-5}	$8.031 \cdot 10^{-4}$	$1.478 \cdot 10^{-2}$	$1.076 \cdot 10^{-3}$	$2.061 \cdot 10^{-4}$	$6.661 \cdot 10^{-4}$

Table 5.1.: Minimum inf-sup constants $\beta_N(\mu)$ (see (2.17)) evaluated using (a) Algorithm 1, (b) Algorithm 2, (c) Algorithm 3, and (d) Algorithm 4 (see §2.7 and §4.7) for different values of N and ε ; the minimum is taken over 25 parameter values.

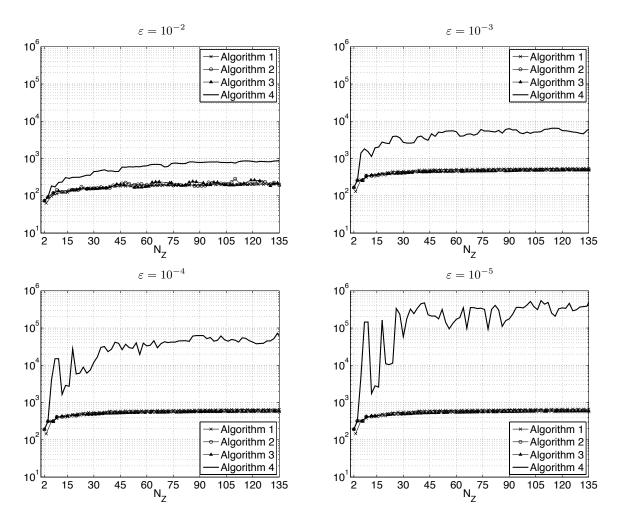


Figure 5.3.: Maximum condition number $\kappa_N^{\varepsilon}(\mu)$ (see (5.6)) of the reduced basis system matrix as a function of N_Z for different values of ε evaluated using Algorithm 1, Algorithm 2, Algorithm 3, and Algorithm 4 (see §2.7 and §4.7); the maximum is taken over 25 parameter values.

In addition to the inf-sup constants $\beta_N(\mu)$, we also consider the condition number

$$\kappa_N^{\varepsilon}(\mu) \equiv \frac{\sigma_N^{\varepsilon,\max}(\mu)}{\sigma_N^{\varepsilon,\min}(\mu)}, \quad \varepsilon > 0, \quad \forall \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max},$$
(5.6)

associated with the reduced basis system (4.15), (4.16); here, $\sigma_N^{\varepsilon,\max}(\mu)$ and $\sigma_N^{\varepsilon,\min}(\mu)$ denote the maximum and minimum singular values of the corresponding system matrix (see also Appendix A). Figure 5.3 presents maximum values of $\kappa_N^{\varepsilon}(\mu)$ as a function of N_Z for different values of ε . We observe that the penalty term has a significant impact on the condition number of the reduced basis system matrix: As ε decreases and we approach the nonpenalized problem, values of $\kappa_N^{\varepsilon}(\mu)$ increase. In case of Algorithm 1, Algorithm 2, and Algorithm 3, this correlation is fairly weak and condition numbers seem to level off at a maximum of roughly 700. In case of Algorithm 4, the situation is different: Here, values of $\kappa_N^{\varepsilon}(\mu)$ explode and the reduced basis system becomes ill-conditioned as the penalty parameter ε tends towards zero.

Accordingly reflected in our reduced basis approximations, this behavior manifests in practice what has been predicted by a priori stability estimates in §4.3.3. We present in Fig. 5.4 the maximum relative errors $||e_N^{u,\varepsilon}(\mu)||_X/||u^{\varepsilon}(\mu)||_X$ and $||e_N^{p,\varepsilon}(\mu)||_Y/||p^{\varepsilon}(\mu)||_Y$ (see (4.18)) in the reduced basis velocity and pressure approximations using Algorithm 4 for different values of the penalty parameter ε ; a direct comparison with the other sampling procedures is then given in Fig. 5.5 and Fig. 5.6. In case of Algorithm 1, Algorithm 2, and Algorithm 3, the velocity approximation spaces are rich enough to ensure stable reduced basis approximations whose convergence is not affected by the penalty parameter ε ; for different values of ε , we obtain approximations $u_N^{\varepsilon}(\mu)$ for $u^{\varepsilon}(\mu)$ (see Fig. 5.5) and $p_N^{\varepsilon}(\mu)$ for $p^{\varepsilon}(\mu)$ (see Fig. 5.6) that are equally accurate and rapidly convergent. In case of Algorithm 4

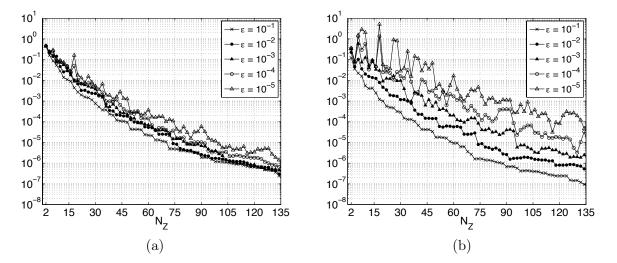


Figure 5.4.: Numerical results for Algorithm 4 (see §4.7) for different values of ε : Maximum relative errors (a) $\|e_N^{u,\varepsilon}(\mu)\|_X/\|u^{\varepsilon}(\mu)\|_X$ and (b) $\|e_N^{p,\varepsilon}(\mu)\|_Y/\|p^{\varepsilon}(\mu)\|_Y$ (see (4.18)) are shown as functions of N_Z ; the maximum is taken over 25 parameter values.

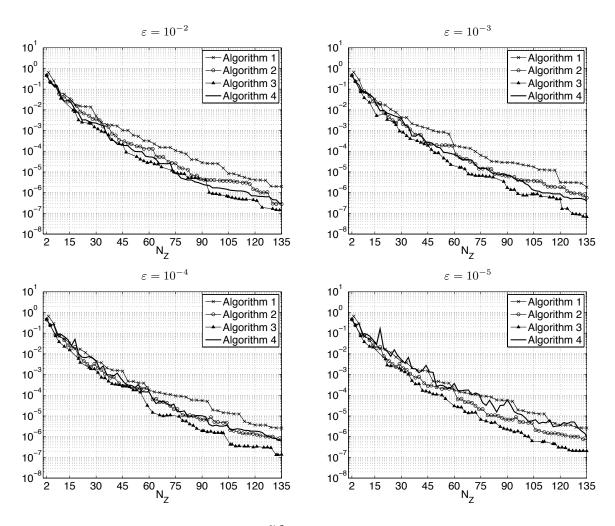


Figure 5.5.: Maximum relative error $||e_N^{u,\varepsilon}(\mu)||_X/||u^{\varepsilon}(\mu)||_X$ (see (4.18)) as a function of N_Z for different values of ε evaluated using Algorithm 1, Algorithm 2, Algorithm 3, and Algorithm 4 (see §2.7 and §4.7); the maximum is taken over 25 parameter values.

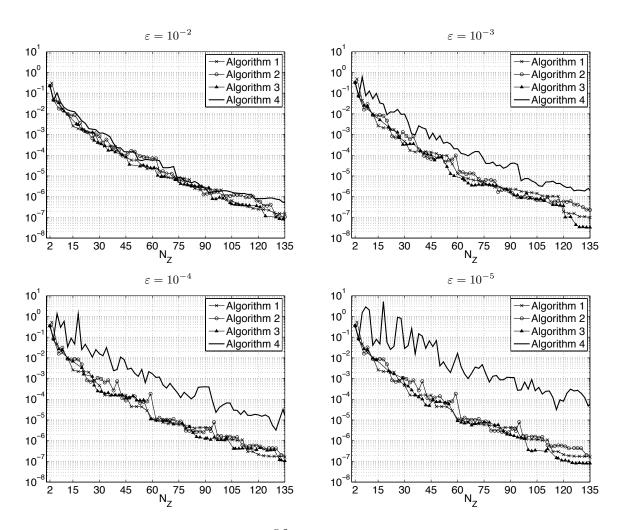


Figure 5.6.: Maximum relative error $||e_N^{p,\varepsilon}(\mu)||_Y/||p^{\varepsilon}(\mu)||_Y$ (see (4.18)) as a function of N_Z for different values of ε evaluated using Algorithm 1, Algorithm 2, Algorithm 3, and Algorithm 4 (see §2.7 and §4.7); the maximum is taken over 25 parameter values.

where the inf-sup constants $\beta_N(\mu)$ associated with X_N and Y_N are comparatively small, convergence of the reduced basis approximations is remarkably good in the case of $\varepsilon \geq 10^{-2}$ but deteriorates for penalty parameters approaching zero (see Fig. 5.4). However, we note that this behavior is in practice much more prominent for the pressure than for the velocity approximations: For $\varepsilon \leq 10^{-3}$, errors in the approximations $p_N^{\varepsilon}(\mu)$ are much larger than in case of the other sampling procedures (see Fig. 5.6) whereas the approximations $u_N^{\varepsilon}(\mu)$ are still very accurate (see Fig. 5.4(a) and Fig. 5.5).

Again, we observe that stabilizing adaptively pays off: Compared to the standard greedy procedure in Algorithm 1, Algorithm 2 and Algorithm 3 achieve reduced basis approximations that are for the pressure equally well (see Fig. 5.6) but for the velocity noticeably better (see Fig. 5.5).

5.3.3 A Posteriori Error Bounds

We now investigate how the error bounds $\Delta_N^{\varepsilon}(\mu)$ developed in §4.4 perform in practice.

The SCM (see §4.6.2 and §2.6.2) enables the (Online-)efficient estimation of the coercivity constants $\alpha_a(\mu)$ and $\alpha_c(\mu)$ (see (2.10) and (4.12)). We here apply the method for $M_{\alpha} = \infty$ and an exhaustive sample $\Xi \subset \mathcal{D}$ of size $|\Xi| = 4,225$. Setting further $\delta_{\text{tol}}^{\text{SCM}} = 0.01$, we obtain $K_{\text{max}} = 35$ for the coercivity constants $\alpha_a(\mu)$, and $K_{\text{max}} = 4$ for the coercivity constants $\alpha_c(\mu)$. We receive very accurate (Online-)efficient lower bounds $\alpha_a^{\text{LB}}(\mu)$ and $\alpha_c^{\text{LB}}(\mu)$ (see Fig. 3.14 in §3.3.4 and Fig. 5.7), providing a posteriori error bounds $\Delta_N^{\varepsilon}(\mu)$ that essentially coincide with their values based on the evaluation of the exact constants $\alpha_a(\mu)$ and $\alpha_c(\mu)$.

In Fig. 5.8, Fig. 5.9, and Fig. 5.10, we now show the error $\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu}$ (see (4.18), (4.19)) and the associated reduced basis *a posteriori* error bound $\Delta_N^{\varepsilon}(\mu)$ (see (4.33)) normalized relative to $\|(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ evaluated using Algorithm 1, Algorithm 2, and Algorithm 3, respectively, for different values of the penalty parameter ε . We see that for $\varepsilon = 10^{-2}$, the error bounds are tight and the effectivities are relatively small; however, the associated truth

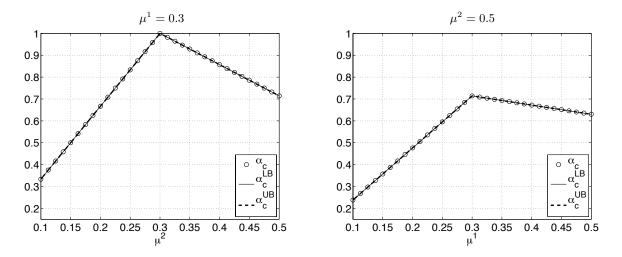


Figure 5.7.: SCM lower and upper bounds $\alpha_c^{\text{LB}}(\mu)$ and $\alpha_c^{\text{UB}}(\mu)$ for the coercivity constants $\alpha_c(\mu)$ (see (4.12), (4.29)) where $|\Xi| = 4,225$, $M_{\alpha} = \infty$, $\delta_{\text{tol}}^{\text{SCM}} = 0.01$, and $K_{\text{max}} = 4$.

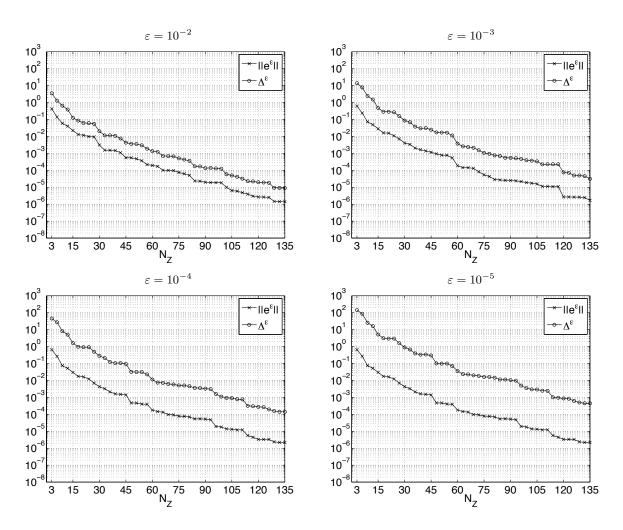


Figure 5.8.: Maximum relative error $\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu}/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.18), (4.19)) and maximum relative error bound $\Delta_N^{\varepsilon}(\mu)/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.33)) as functions of N_Z for different values of ε evaluated using Algorithm 1 (see §2.7); the maximum is taken over 25 parameter values.

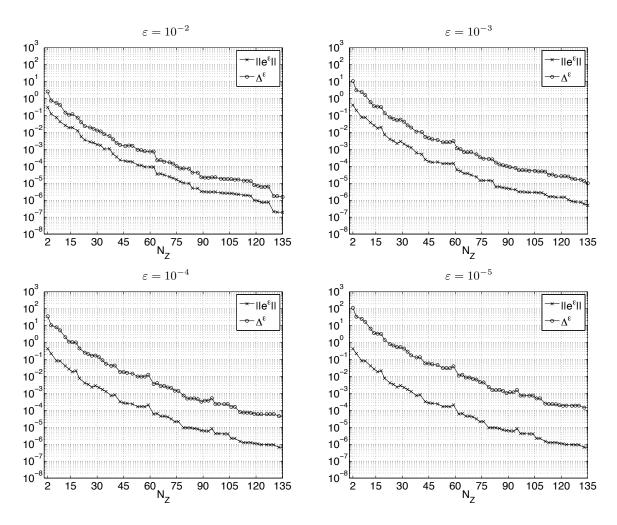


Figure 5.9.: Maximum relative error $\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu}/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.18), (4.19)) and maximum relative error bound $\Delta_N^{\varepsilon}(\mu)/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.33)) as functions of N_Z for different values of ε evaluated using Algorithm 2 (see §2.7); the maximum is taken over 25 parameter values.

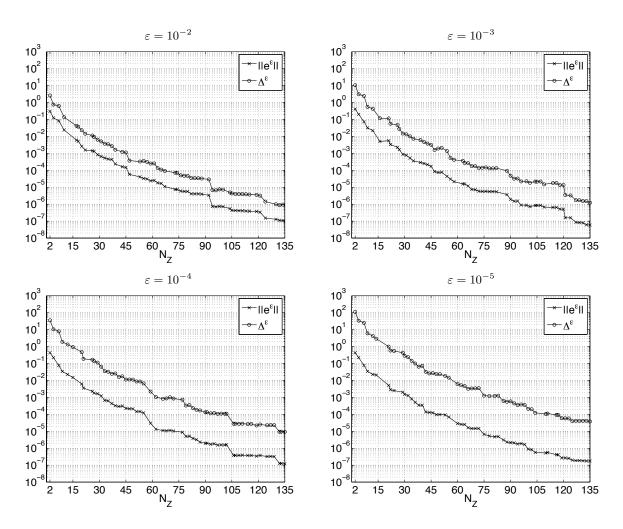


Figure 5.10.: Maximum relative error $\|e_N^{\varepsilon}(\mu)\|_{Z,\varepsilon,\mu}/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.18), (4.19)) and maximum relative error bound $\Delta_N^{\varepsilon}(\mu)/\|(u^{\varepsilon}(\mu),p^{\varepsilon}(\mu))\|_{Z,\varepsilon,\mu}$ (see (4.33)) as functions of N_Z for different values of ε evaluated using Algorithm 3 (see §2.7); the maximum is taken over 25 parameter values.

(a) Algorithm 1						
N	$\varepsilon = 10^{-2}$	$\varepsilon = 10^{-3}$	$\varepsilon = 10^{-4}$	$\varepsilon = 10^{-5}$		
5	$1.163 \cdot 10^1$	$3.357\cdot 10^1$	$1.035 \cdot 10^2$	$3.273 \cdot 10^2$		
10	$1.149\cdot 10^1$	$3.407\cdot 10^1$	$1.065 \cdot 10^2$	$3.368\cdot 10^2$		
15	$1.150 \cdot 10^1$	$3.353 \cdot 10^1$	$1.027\cdot 10^2$	$3.248 \cdot 10^2$		
20	$1.030 \cdot 10^1$	$3.515 \cdot 10^1$	$1.111 \cdot 10^{2}$	$3.511 \cdot 10^2$		
25	$1.300 \cdot 10^1$	$3.224 \cdot 10^1$	$1.108\cdot 10^2$	$3.504 \cdot 10^2$		
30	$1.171 \cdot 10^1$	$3.145 \cdot 10^{1}$	$9.657\cdot 10^1$	$3.053\cdot 10^2$		
35	$1.041 \cdot 10^1$	$3.435 \cdot 10^1$	$1.129 \cdot 10^{2}$	$3.572 \cdot 10^2$		
40	$1.068 \cdot 10^1$	$3.456 \cdot 10^1$	$1.017 \cdot 10^{2}$	$3.217 \cdot 10^{2}$		
(b) Algorithm 2	2				
N	$\varepsilon = 10^{-2}$	$\varepsilon = 10^{-3}$	$\varepsilon = 10^{-4}$	$\varepsilon = 10^{-5}$		
5	$1.167\cdot 10^1$	$3.277 \cdot 10^1$	$1.040 \cdot 10^{2}$	$3.288 \cdot 10^2$		
10	$9.699 \cdot 10^0$	$3.176 \cdot 10^1$	$9.979\cdot 10^1$	$3.153 \cdot 10^2$		
15	$9.903\cdot 10^0$	$3.259 \cdot 10^1$	$1.046 \cdot 10^2$	$3.308\cdot 10^2$		
20	$9.982\cdot 10^0$	$3.076\cdot 10^1$	$9.646 \cdot 10^1$	$3.049 \cdot 10^2$		
25	$9.710\cdot 10^0$	$3.110 \cdot 10^1$	$9.811 \cdot 10^1$	$3.107 \cdot 10^2$		
30	$9.771\cdot 10^0$	$3.727 \cdot 10^1$	$1.260 \cdot 10^{2}$	$3.977\cdot 10^2$		
35	$1.070\cdot 10^1$	$2.899 \cdot 10^1$	$8.831 \cdot 10^1$	$2.790 \cdot 10^{2}$		
40	$1.072 \cdot 10^1$	$2.911 \cdot 10^1$	$9.391 \cdot 10^1$	$2.964\cdot 10^2$		
(c)) Algorithm 3					
N	$\varepsilon = 10^{-2}$	$\varepsilon = 10^{-3}$	$\varepsilon = 10^{-4}$	$\varepsilon = 10^{-5}$		
5	$1.209\cdot 10^1$	$3.486 \cdot 10^1$	$1.051 \cdot 10^2$	$3.313 \cdot 10^2$		
10	$1.249 \cdot 10^1$	$4.018 \cdot 10^1$	$1.215 \cdot 10^2$	$3.155 \cdot 10^2$		
15	$1.191 \cdot 10^1$	$3.548 \cdot 10^1$	$9.780 \cdot 10^1$	$3.327\cdot 10^2$		
20	$1.207\cdot 10^1$	$3.281 \cdot 10^1$	$1.121 \cdot 10^2$	$3.229\cdot 10^2$		
25	$1.083 \cdot 10^1$	$3.473 \cdot 10^1$	$1.094 \cdot 10^2$	$3.539 \cdot 10^2$		
30	$1.160 \cdot 10^1$	$3.322 \cdot 10^1$	$1.102 \cdot 10^2$	$3.623 \cdot 10^2$		
35	$1.146 \cdot 10^1$	$3.073 \cdot 10^1$	$1.061 \cdot 10^2$	$4.799\cdot 10^2$		
40	$1.056 \cdot 10^1$	$3.629 \cdot 10^1$	$9.906 \cdot 10^1$	$3.742 \cdot 10^2$		

Table 5.2.: Maximum effectivities $\eta_N^{\varepsilon}(\mu)$ (see (4.37)) for several values of N and ε evaluated using (a) Algorithm 1, (b) Algorithm 2, and (c) Algorithm 3 (see §2.7); the maximum is taken over 25 parameter values.

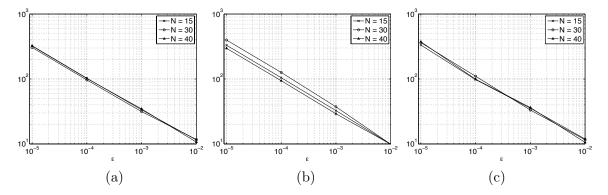


Figure 5.11.: Maximum effectivities $\eta_N^{\varepsilon}(\mu)$ (see (4.37)) as a function of ε for different values of N evaluated using (a) Algorithm 1, (b) Algorithm 2, and (c) Algorithm 3 (see §2.7); the maximum is taken over 25 parameter values.

finite element approximation may not be sufficiently accurate (see Fig. 5.2). As we decrease ε , we obtain a more accurate truth approximation, but the reduced basis error bounds become less sharp. Nevertheless, these results show that even for relatively small values of the penalty parameter, the effectivities remain reasonably small, and the dependence on ε is relatively weak as predicted by Corollary 4.4.2 and Corollary 4.4.3. To further quantify this statement, we present in Table 5.2(a), (b), and (c) the effectivities associated with $\Delta_N^{\varepsilon}(\mu)$ for several values of N and ε evaluated using Algorithm 1, Algorithm 2, and Algorithm 3, respectively. We note that for all three methods, the effectivities remain more or less constant with N. Furthermore, we show in Fig. 5.11 the dependence of the effectivities on the penalty parameter; we clearly observe the $O(\frac{1}{\sqrt{\varepsilon}})$ -dependence predicted by Corollary 4.4.2 and Corollary 4.4.3. The effects of the penalty parameter on the effectivities are thus relatively benign, and we obtain meaningful bounds for reasonably small values of ε . In particular, even at larger effectivities, the error bounds reflect the behavior of the actual error (see Fig. 5.8, Fig. 5.9, and Fig. 5.10), and our greedy sampling processes perform successfully as they are able to choose the parameter snapshots properly (i.e., where the error is largest).

5.3.4 Computation Times

We now discuss the Online computation times. The following timing tests were performed on a 2.66 GHz Intel Core 2 Duo processor. For comparison, once the μ -independent parts in (2.7) have been formed, direct computation of the truth approximation $(u^{\varepsilon}(\mu), p^{\varepsilon}(\mu))$ (i.e., assembly and solution of (2.13)) takes on average 5.7 seconds. Again, the rigorous and efficient error bounds $\Delta_N^{\varepsilon}(\mu)$ allow us to choose the reduced basis system dimension N_Z just large enough to obtain a desired accuracy.

For the penalty parameter $\varepsilon = 10^{-2}$, the error bounds are sharp with effectivities of approximately 11 (see Table 5.2). If we construct the reduced basis approximations spaces by using the standard greedy procedure given in Algorithm 1, we then need $N_Z = 42$ basis functions to achieve a prescribed accuracy of roughly 1% or better in the reduced basis approximations $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ (see Fig. 5.8). Once the database has been loaded, the Online calculation of $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ (i.e., assembly and solution of (4.14)) and $\Delta_N^{\varepsilon}(\mu)$ for any new value of $\mu \in \mathcal{D}$ takes on average 0.21 and 8.94 milliseconds, respectively, which is in total roughly 600 times faster than direct computation of the truth approximation. In case of Algorithm 3, the same accuracy is achieved for $N_Z = 27$ (see Fig. 5.10); the Online calculation of $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ and $\Delta_N^{\varepsilon}(\mu)$ then takes on average 0.1 and 3.87 milliseconds, respectively, and is thus roughly 1,400 times faster than direct computation of the truth approximation.

We now set $\varepsilon = 10^{-5}$. As they become more pessimistic with decreasing values of ε , the error bounds here dictate larger system dimensions at which they guarantee the same order of accuracy. Using Algorithm 1, we need $N_Z = 93$ to achieve an accuracy of roughly 1% or better in the reduced basis approximations $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$; the Online calculation of $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ and $\Delta_N^{\varepsilon}(\mu)$ for any new value of $\mu \in \mathcal{D}$ then takes on average 1.12 and 42.1 milliseconds, respectively, which is in total roughly 130 times faster than direct computation of the truth approximation. Using Algorithm 3, we need $N_Z = 60$; the Online calculation of $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ and $\Delta_N^{\varepsilon}(\mu)$ then takes on average 0.4 and 16.76 milliseconds, respectively, which is in total roughly 330 times faster than direct computation of the truth approximation. Therefore, even for small penalty parameters ε , accurate approximations are guaranteed at significant Online savings.

Detailed computation times for different values of ε are given in Table 5.3.

(a) $\varepsilon = 10^{-2}$						
Method	N_Z	N	$(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$	$\Delta_N^{\varepsilon}(\mu)$	Total	
Algorithm 1	42 (66)	14(22)	0.21 (0.54)	8.94(21.53)	9.14 (22.07)	
Algorithm 2	34(53)	15(23)	0.13(0.30)	5.43(12.67)	5.56 (12.97)	
Algorithm 3	27(47)	10(19)	0.10(0.24)	3.87(10.36)	3.97(10.59)	
(b) $\varepsilon = 10^{-3}$						
Method	N_Z	N	$(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$	$\Delta_N^{\varepsilon}(\mu)$	Total	
Algorithm 1	60(78)	20(26)	0.42(0.78)	17.84(29.90)	18.26(30.69)	
Algorithm 2	43(65)	19(29)	0.19(0.47)	8.39(18.77)	8.58 (19.25)	
Algorithm 3	35(56)	14(23)	0.14(0.35)	5.96(14.39)	6.09 (14.74)	
(c) $\varepsilon = 10^{-4}$						
Method	N_Z	N	$(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$	$\Delta_N^{\varepsilon}(\mu)$	Total	
Algorithm 1	63 (102)	21(34)	0.48(1.38)	19.60(50.55)	20.08 (51.92)	
Algorithm 2	53(79)	23(35)	0.30(0.74)	12.68(27.54)	12.98 (28.28)	
Algorithm 3	51(66)	22(27)	$0.28 \ (0.50)$	11.81(20.04)	12.09(20.54)	
(d) $\varepsilon = 10^{-5}$						
Method	N_Z	N	$(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$	$\Delta_N^{\varepsilon}(\mu)$	Total	
Algorithm 1	93(117)	31(39)	1.12(1.79)	42.10(66.10)	43.22(67.89)	
Algorithm 2	66 (97)	29(44)	0.48(1.12)	19.44(40.69)	19.92(41.80)	
Algorithm 3	60(86)	24(35)	0.40(0.91)	16.76(33.63)	17.16(34.53)	

Table 5.3.: Average computation times in milliseconds for the Online evaluation of $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$ (assembly and solution of (4.14)) and the error bounds $\Delta_N^{\varepsilon}(\mu)$ (see (4.33)): For (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$, times are measured using either Algorithm 1, Algorithm 2, or Algorithm 3 (see §2.7) with a prescribed accuracy of at least 1% (resp., 0.1%) for the reduced basis approximations $(u_N^{\varepsilon}(\mu), p_N^{\varepsilon}(\mu))$.

5.4 Concluding Remarks

We present in this chapter numerical results for the reduced basis approach based on a penalty formulation proposed in Chapter 4. The method is applied to a perturbed Stokes flow in a parametrized microchannel.

As the computation of (Online-)efficient lower bounds to inf-sup stability constants is avoided, we obtain reduced basis approximations and rigorous *a posteriori* error bounds that are computed very easily. Numerical results show that the approximations for both the velocity and the pressure converge rapidly and that the error bounds are reasonably sharp. While the computation of inf-sup constants may be manageable for the Stokes problem, it becomes computationally even more intensive in case of the nonlinear Navier–Stokes equations. The penalty approach enables us to consider flow problems in parametrized domains with significantly less effort and thus provides a convenient way to also handle more complex geometric variations.

Clearly, the method also comes with several disadvantages. First, the introduced penalty term causes an additional error in the truth velocity and pressure approximations. For engineering problems of interest, these errors are generally acceptable but choosing an appropriate value of the penalty parameter ε is a nontrivial issue. Second, the penalty term's regularizing effect is manifested in a significant impact on the condition number of the reduced basis system; yet, stable convergence cannot be guaranteed for reduced basis approximation spaces with small inf-sup constants (2.17). The problem can be resolved by an additional enrichment of the reduced basis approximation space for the primal variable: In this case, we obtain stable approximations whose convergence is not affected by the penalty parameter. Third, our *a posteriori* error bounds become more pessimistic as the penalty parameter ε decreases and our truth approximations become more accurate. However, theoretical and numerical results show that these effects are relatively benign and we obtain meaningful and useful error bounds even for reasonably small values of ε .

Compared to former approaches, sampling procedures that recognize the need for stabilization adaptively again construct reduced basis approximation spaces that are computationally much more efficient. Even though the procedures used are not optimally suited for the penalty context (see Remark 4.7.1), they demonstrate the computational savings that are generally achievable; in the case of our Stokes model problem, we obtain computational savings of up to 40%. Numerical results show that in the penalty context, the reduced basis inf-sup constants $\beta_N(\mu)$ are not necessarily appropriate indicators for an ill-conditioned system; well-suited adaptive sampling procedures should include the condition number $\kappa_N^{\varepsilon}(\mu)$.

Chapter 6

Approximation of the Instationary Stokes Equations

6.1 Introduction

In this chapter, we provide an outlook on how the techniques presented in the previous part of this thesis may be extended to the time-dependent setting. To this end, we consider the instationary Stokes equations. We shall discuss both a method based on the standard formulation as well as a method based on a penalty approach, which combine approaches developed in §2, §3 and §4, §5 with current reduced basis techniques for parabolic problems (see, e.g., [37, 41, 47] and also [63, 64]). The analysis then shows how time integration affects the development of reduced basis a *posteriori* error bounds as well as the construction of computationally efficient reduced basis approximation spaces.

Starting from the standard mixed formulation of the instationary Stokes equations, we develop rigorous a *posteriori* error bounds for the reduced basis velocity approximations. As in the stationary case presented in §2 and §3, they involve the (Online-)estimation of coercivity, continuity, and inf-sup stability constants associated with the diffusion term and incompressibility constraint; in addition, they now also depend on continuity constants associated with the mass term. Employing a penalty formulation, we obtain rigorous upper bounds for the errors in both the velocity and pressure approximations. As in the stationary case presented in §4 and §5, they are computationally very efficient since they do not involve the estimation of inf-sup constants but only depend on coercivity constants associated with the diffusion and penalty terms; however, they again also depend on the penalty parameter such that associated effectivities increase as we approach the nonpenalized problem. To construct efficient reduced basis approximation spaces, we consider a POD greedy procedure (see [37, 46, 47]) that is coupled with adaptive stabilization techniques developed in §2.7. To demonstrate their performance in practice, the methods are then applied to a Stokes flow in a parametrized domain where evolution in time is induced by a time-dependent velocity profile on the inflow boundary.

The chapter is organized as follows: In §6.2, we introduce the general problem formulation and its "truth" approximation. We start from a time-discrete framework already that allows us to directly recover the settings discussed in §2.2 and §4.2; now, we have a saddle point problem associated with each time step. The time discretization scheme is here given by a backward Euler method. Section 6.3 defines our reduced basis approximation as the Galerkin projection onto a low-dimensional reduced basis approximation space. We develop rigorous *a posteriori* error bounds in §6.4. Again, reduced basis approximations and error bounds can be computed Online-efficiently as summarized in §6.5. This enables us to engage adaptive sampling processes for constructing computationally efficient reduced basis approximation spaces, which shall be discussed in §6.6. In §6.7, we introduce our instationary Stokes model problem. Numerical results in §6.8 then illustrate (i) the rapid convergence of reduced basis approximations, (ii) the performance of *a posteriori* error bounds with respect to sharpness, and (iii) computational efficiency. Finally, in §6.9, we give some concluding remarks.

6.2 General Problem Statement

6.2.1 Formulation

We assume the setting as introduced in §2.2.1 and §4.2.1. Furthermore, for any parameter $\mu \in \mathcal{D}$, let $m(\cdot, \cdot; \mu) : X_e \times X_e \to \mathbb{R}$ be a bilinear form that is symmetric, continuous, and positive definite,

$$\gamma_m^{\rm e}(\mu) \equiv \sup_{u \in X_{\rm e}} \sup_{v \in X_{\rm e}} \frac{m(u, v; \mu)}{\|u\|_{X_{\rm e}} \|v\|_{X_{\rm e}}} < \infty, \quad \forall \ \mu \in \mathcal{D}, \tag{6.1}$$

$$m(v, v; \mu) > 0, \quad \forall \ 0 \neq v \in X_{\mathbf{e}}, \quad \forall \ \mu \in \mathcal{D},$$

$$(6.2)$$

and which exhibits a μ -affine representation analogous to (2.7). For any $\mu \in \mathcal{D}$, $m(\cdot, \cdot; \mu)$ then defines an inner product on X_e ; the associated norm shall be denoted by $\|\cdot\|_{\mu} \equiv \sqrt{m(\cdot, \cdot; \mu)}$.

We further assume that we are given a time interval [0,T], T > 0, and linear functionals $f(\cdot;\mu) \in C^0(0,T;X'_e)$ and $g(\cdot;\mu) \in C^0(0,T;Y'_e)$ for all $\mu \in \mathcal{D}$; for a vector space $V, C^0(0,T;V)$ here denotes the space of V-valued functions of class C^0 with respect to $t \in [0,T]$. Throughout this thesis, we directly consider a time-discrete framework: We divide the time interval [0,T] into K subintervals of equal length $\Delta t \equiv T/K$, and define $t^k \equiv k\Delta t$ for all $k = 0, \ldots, K$; for notational convenience, we also introduce $\mathbb{K} \equiv \{1, \ldots, K\}$ and $\mathbb{K}_0 \equiv \mathbb{K} \cup \{0\}$. We then set $f^k(\cdot;\mu) \equiv f(t^k;\mu) \in X'_e$ and $g^k(\cdot;\mu) \equiv g(t^k;\mu) \in Y'_e$ for all $k \in \mathbb{K}_0, \mu \in \mathcal{D}$.

For $\varepsilon \geq 0$, we now consider the following "exact" — more precisely, semi-discrete — problem resulting from a backward Euler method (see, e.g., [27, 42, 89, 102]): For any given $\mu \in \mathcal{D}$, we find $u_e^{\varepsilon,k}(\mu) \in X_e$ and $p_e^{\varepsilon,k}(\mu) \in Y_e$, $k \in \mathbb{K}$, such that $u_e^{\varepsilon,0}(\mu) = 0^1$ and

$$\frac{1}{\Delta t} m(u_{e}^{\varepsilon,k}(\mu) - u_{e}^{\varepsilon,k-1}(\mu), v; \mu)
+ a(u_{e}^{\varepsilon,k}(\mu), v; \mu) + b(v, p_{e}^{\varepsilon,k}(\mu); \mu) = f^{k}(v; \mu), \quad \forall v \in X_{e},
b(u_{e}^{\varepsilon,k}(\mu), q; \mu) - \varepsilon c(p_{e}^{\varepsilon,k}(\mu), q; \mu) = g^{k}(q; \mu), \quad \forall q \in Y_{e},$$
(6.3)

Even though we here use a common notation for simplicity in exposition, we point out that (6.3) states very different problems for $\varepsilon = 0$ and $\varepsilon > 0$, respectively. For $\varepsilon = 0$, we also denote $u_{\rm e}^k(\mu) \equiv u_{\rm e}^{0,k}(\mu)$, $k \in \mathbb{K}_0$, and $p_{\rm e}^k(\mu) \equiv p_{\rm e}^{0,k}(\mu)$, $k \in \mathbb{K}$, for all $\mu \in \mathcal{D}$. For $\varepsilon > 0$,

¹We here assume zero initial conditions for simplicity; note that nonzero initial conditions can be handled as well without much difficulty (see [41]).

corresponding to our discussions in Chapter 4, (6.3) can be considered as a perturbed or regularized version of the problem associated with $\varepsilon = 0$; in this case, we therefore call $(u_e^{\varepsilon,k}(\mu), p_e^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, also the penalty solution. Since these problems differ considerably in their general nature (see §2 and §4), we shall often treat them separately in the following analysis and explicitly distinguish between the two cases $\varepsilon = 0$ and $\varepsilon > 0$. From (6.2) and (2.3), the bilinear form $\frac{1}{\Delta t}m(\cdot, \cdot; \mu) + a(\cdot, \cdot; \mu)$ is coercive on X_e for any $\mu \in \mathcal{D}$. The problem (6.3) is thus uniquely solvable for $(u_e^k(\mu), p_e^k(\mu)), k \in \mathbb{K}$, and $(u_e^{\varepsilon,k}(\mu), p_e^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, as a saddle point problem according to §2.2.1 and §4.2.1, respectively.

6.2.2 Truth Approximation

As in §2.2.2 and §4.2.2, we denote by X and Y finite-dimensional subspaces of X_e and Y_e satisfying the assumptions (2.8)–(2.11), (4.11), and (4.12). These in particular also inherit the inner product $m(\cdot, \cdot; \mu)$ and associated norm $\|\cdot\|_{\mu}$,

$$\gamma_m(\mu) \equiv \sup_{u \in X} \sup_{v \in X} \frac{m(u, v; \mu)}{\|u\|_X \|v\|_X} = \sup_{v \in X} \frac{m(v, v; \mu)}{\|v\|_X^2} < \infty,$$
(6.4)

$$m(v, v; \mu) > 0, \quad \forall \ 0 \neq v \in X, \quad \forall \ \mu \in \mathcal{D}.$$
 (6.5)

Again, we denote by \mathcal{N} the large dimension of the approximation space $Z = X \times Y$.

Our high-fidelity "truth" discretization for (6.3) now reads as follows: For $\varepsilon \geq 0$ and any given $\mu \in \mathcal{D}$, we find $u^{\varepsilon,k}(\mu) \in X$ and $p^{\varepsilon,k}(\mu) \in Y$, $k \in \mathbb{K}$, such that $u^{\varepsilon,0}(\mu) = 0$ and

$$\frac{1}{\Delta t} m(u^{\varepsilon,k}(\mu) - u^{\varepsilon,k-1}(\mu), v; \mu)
+ a(u^{\varepsilon,k}(\mu), v; \mu) + b(v, p^{\varepsilon,k}(\mu); \mu) = f^k(v; \mu), \quad \forall v \in X,
b(u^{\varepsilon,k}(\mu), q; \mu) - \varepsilon c(p^{\varepsilon,k}(\mu), q; \mu) = g^k(q; \mu), \quad \forall q \in Y,$$
(6.6)

In case of $\varepsilon = 0$, we also denote $u^k(\mu) \equiv u^{0,k}(\mu)$, $k \in \mathbb{K}_0$, and $p^k(\mu) \equiv p^{0,k}(\mu)$, $k \in \mathbb{K}$. As the exact problem in §6.2.1, the problem (6.6) is uniquely solvable for $(u^k(\mu), p^k(\mu)), k \in \mathbb{K}$, and $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, according to §2.2.2 and §4.2.2, respectively.

Let $A(\mu)$, $B(\mu)$, $B(\mu)^t$, and $C(\mu)$ be the linear operators associated with the bilinear forms $a(\cdot, \cdot; \mu)$, $b(\cdot, \cdot; \mu)$, and $c(\cdot, \cdot; \mu)$ as introduced in §2.2.2 and §4.2.2; accordingly, we also define the bounded linear operator $M(\mu) : X \to X'$ by

$$\langle M(\mu) u, v \rangle = m(u, v; \mu), \quad \forall u, v \in X.$$

The system (6.6) may then be equivalently written as

$$M(\mu) \frac{u^{\varepsilon,k}(\mu) - u^{\varepsilon,k-1}(\mu)}{\Delta t} + A(\mu) u^{\varepsilon,k}(\mu) + B(\mu)^t p^{\varepsilon,k}(\mu) = f^k(\mu) \quad \text{in } X', \quad k \in \mathbb{K},$$
$$B(\mu) u^{\varepsilon,k}(\mu) - \varepsilon C(\mu) p^{\varepsilon,k}(\mu) = g^k(\mu) \quad \text{in } Y',$$

where $f^k(\mu) \equiv f^k(\cdot;\mu)|_X \in X'$ and $g^k(\mu) \equiv g^k(\cdot;\mu)|_Y \in Y'$ for all $k \in \mathbb{K}_0$.

6.3 Reduced Basis Approximation

We now turn to the reduced basis method, discussing the approximation procedure, rigorous *a posteriori* error estimators, and the construction of stable approximation spaces that capture the causality associated with the parameter dependence *as well as* with evolution in time.

Suppose that we are given a set of nested, low-dimensional reduced basis approximation subspaces $X_N \subset X_{N+1} \subset X$ and $Y_N \subset Y_{N+1} \subset Y$, $N \in \mathbb{N}_{\max} \equiv \{1, \ldots, N_{\max}\}$, then the reduced basis approximation is defined as the Galerkin projection with respect to the truth problem (6.6) onto these low-dimensional subspaces: For $\varepsilon \geq 0$ and any given $\mu \in \mathcal{D}$, we find $u_N^{\varepsilon,k}(\mu) \in X_N$ and $p_N^{\varepsilon,k}(\mu) \in Y_N$, $k \in \mathbb{K}$, such that $u_N^{\varepsilon,0}(\mu) = 0$ and

$$\frac{1}{\Delta t} m(u_N^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k-1}(\mu), v_N; \mu)
+ a(u_N^{\varepsilon,k}(\mu), v_N; \mu) + b(v_N, p_N^{\varepsilon,k}(\mu); \mu) = f^k(v_N; \mu), \quad \forall v_N \in X_N,
b(u_N^{\varepsilon,k}(\mu), q_N; \mu) - \varepsilon c(p_N^{\varepsilon,k}(\mu), q_N; \mu) = g^k(q_N; \mu), \quad \forall q_N \in Y_N,$$
(6.7)

Again, we denote $u_N^k(\mu) \equiv u_N^{0,k}(\mu)$, $k \in \mathbb{K}_0$, and $p_N^k(\mu) \equiv p_N^{0,k}(\mu)$, $k \in \mathbb{K}$. Written in operator notation, the discrete system (6.7) reads

$$\frac{1}{\Delta t} M_N(\mu) \Big(u_N^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k-1}(\mu) \Big)
+ A_N(\mu) u_N^{\varepsilon,k}(\mu) + B_N(\mu)^t p_N^{\varepsilon,k}(\mu) = f_N^k(\mu) \quad \text{in } X'_N,
B_N(\mu) u_N^{\varepsilon,k}(\mu) - \varepsilon C_N(\mu) p_N^{\varepsilon,k}(\mu) = g_N^k(\mu) \quad \text{in } Y'_N, \qquad (6.8)$$

where $A_N(\mu)$, $B_N(\mu)$, $B_N(\mu)^t$, and $C_N(\mu)$ are the linear operators defined in §2.3.1 and §4.3.1, and $M_N(\mu) : X_N \to X'_N$ is given by

$$\langle M_N(\mu)u_N, v_N \rangle = m(u_N, v_N; \mu), \quad \forall u_N, v_N \in X_N, \quad \forall \mu \in \mathcal{D};$$

moreover, $f_N^k(\mu) \equiv f^k(\cdot;\mu)|_{X_N} \in X'_N$ and $g_N^k(\mu) \equiv g^k(\cdot;\mu)|_{Y_N} \in Y'_N$ for all $k \in \mathbb{K}_0$.

The discrete reduced basis system now essentially behaves as in the stationary case: In case of $\varepsilon = 0$, (6.7) is uniquely solvable for $(u_N^k(\mu), p_N^k(\mu)), k \in \mathbb{K}$, if and only if the reduced basis approximation spaces X_N, Y_N are stable (as defined in §2.3.2); in case of $\varepsilon > 0$, (6.7) is uniquely solvable for $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, for any choice of X_N, Y_N (see §4.3.2).

6.4 A Posteriori Error Estimation

We now develop upper bounds for the errors in our reduced basis approximations that are rigorous, sharp, and computationally efficient. As before, the errors shall be measured relative to the respective truth approximations.

In this section, we assume that the low-dimensional reduced basis spaces X_N, Y_N are constructed such that for any given parameter $\mu \in \mathcal{D}$, a solution $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)) \in X_N \times Y_N$, $k \in \mathbb{K}$, to (6.7) exists. For $\mu \in \mathcal{D}$, we then consider the errors

$$e_N^u(\mu) \equiv (e_N^{u,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{u,k}(\mu) \equiv u^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k}(\mu) \in X, \ k \in \mathbb{K},$$

$$e_N^p(\mu) \equiv (e_N^{p,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{p,k}(\mu) \equiv p^{\varepsilon,k}(\mu) - p_N^{\varepsilon,k}(\mu) \in Y, \ k \in \mathbb{K},$$

$$e_N^\varepsilon(\mu) \equiv (e_N^{\varepsilon,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{\varepsilon,k}(\mu) \equiv (e_N^{u,k}(\mu), e_N^{p,k}(\mu)) \in Z, \ k \in \mathbb{K},$$
(6.9)

in the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, with respect to the truth solution $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}$; we note that in particular $e_N^{u,0}(\mu) \equiv u^{\varepsilon,0}(\mu) - u_N^{\varepsilon,0}(\mu) = 0$ from our initial conditions.

To formulate our reduced basis *a posteriori* error bounds, we again rely on the residuals associated with the reduced basis approximation $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$,

$$r_N^{1,k}(\cdot;\mu) \equiv f^k(\mu) - M(\mu) \frac{u_N^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k-1}(\mu)}{\Delta t} - A(\mu) u_N^{\varepsilon,k}(\mu) - B(\mu)^t p_N^{\varepsilon,k}(\mu) \in X', \quad (6.10)$$

$$r_N^{2,k}(\cdot;\mu) \equiv g^k(\mu) - B(\mu) u_N^{\varepsilon,k}(\mu) + \varepsilon C(\mu) p_N^{\varepsilon,k}(\mu) \in Y', \quad (6.11)$$

for $k \in \mathbb{K}$ and $\mu \in \mathcal{D}$.

In the following analysis, we distinguish between the cases $\varepsilon = 0$ and $\varepsilon > 0$.

6.4.1 $\varepsilon = 0$

We here derive rigorous upper bounds for the error $e_N^u(\mu)$ measured in the "spatio-temporal" energy norm

$$\|(v^{j})_{j\in\mathbb{K}}\|_{\ell^{2}(0,k;X)} \equiv \left(\|v^{k}\|_{\mu}^{2} + \Delta t \sum_{j=1}^{k} \|v^{j}\|_{X,\mu}^{2}\right)^{1/2}, \quad (v^{j})_{j\in\mathbb{K}} \subseteq X, \quad k \in \mathbb{K}.$$
(6.12)

As before, our reduced basis a *posteriori* error bounds shall be formulated in terms of the dual norms of the residuals (6.10) and (6.11), and (Online-)efficient lower and upper bounds to the truth continuity, coercivity, and inf-sup stability constants (2.8), (2.10), and (2.11). In addition to (2.29) and (2.30), we also introduce computationally (Online-)efficient (lower and) upper bounds to the truth continuity constants (6.4),

$$\gamma_m^{\rm LB}(\mu) \le \gamma_m(\mu) \le \gamma_m^{\rm UB}(\mu), \quad \forall \ \mu \in \mathcal{D}.$$
(6.13)

We can now state the following result.

Proposition 6.4.1. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, $k \in \mathbb{K}$, and $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, $\beta_{\text{Br}}^{\text{LB}}(\mu)$, $\gamma_m^{\text{UB}}(\mu)$ satisfying (2.29), (2.30), and (6.13), we define

$$\Delta_{N}^{k}(\mu) \equiv \left[\Delta t \sum_{j=1}^{k} \frac{\|r_{N}^{1,j}(\cdot;\mu)\|_{X'}^{2}}{\alpha_{a}^{\mathrm{LB}}(\mu)} + \frac{2}{\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu)} \left(1 + \frac{\gamma_{a}^{\mathrm{UB}}(\mu)}{\alpha_{a}^{\mathrm{LB}}(\mu)}\right) \|r_{N}^{1,j}(\cdot;\mu)\|_{X'} \|r_{N}^{2,j}(\cdot;\mu)\|_{Y'} + \left(\frac{\gamma_{m}^{\mathrm{UB}}(\mu)}{\Delta t} + \frac{(\gamma_{a}^{\mathrm{UB}}(\mu))^{2}}{\alpha_{a}^{\mathrm{LB}}(\mu)}\right) \frac{\|r_{N}^{2,j}(\cdot;\mu)\|_{Y'}}{(\beta_{\mathrm{Br}}^{\mathrm{LB}}(\mu))^{2}}\right]^{1/2}.$$
 (6.14)

Then, $\Delta_N^k(\mu)$ represents an upper bound for the error $e_N^u(\mu)$ measured in the "spatio-temporal" energy norm (6.12),

$$\|e_N^u(\mu)\|_{\ell^2(0,k;X)} \le \Delta_N^k(\mu), \quad \forall k \in \mathbb{K}, \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max}.$$
(6.15)

Proof. Let μ be any parameter in \mathcal{D} , $N \in \mathbb{N}_{\max}$, and $k \in \mathbb{K}$. For clarity of exposition, we suppress the argument μ in this proof.

Take any $1 \leq j \leq k$. From (6.10), (6.11), and (6.6), the errors $e_N^{u,j} \in X$ and $e_N^{p,j} \in Y$ satisfy the equations

$$\frac{1}{\Delta t}m(e_N^{u,j} - e_N^{u,j-1}, v) + a(e_N^{u,j}, v) + b(v, e_N^{p,j}) = r_N^{1,j}(v), \quad \forall v \in X,$$
(6.16)

$$b(e_N^{u,j},q) = r_N^{2,j}(q), \quad \forall \ q \in Y.$$

$$(6.17)$$

By the LBB inf-sup condition (2.11) and (6.16), we have

$$\beta_{\mathrm{Br}} \|e_N^{p,j}\|_Y \le \sup_{v \in X} \frac{b(v, e_N^{p,j})}{\|v\|_X} = \sup_{v \in X} \frac{r_N^{1,j}(v) - a(e_N^{u,j}, v) - \frac{1}{\Delta t}m(e_N^{u,j} - e_N^{u,j-1}, v)}{\|v\|_X} \le \|r_N^{1,j}\|_{X'} + \gamma_a \|e_N^{u,j}\|_X + \frac{\sqrt{\gamma_m}}{\Delta t} \|e_N^{u,j} - e_N^{u,j-1}\|_{\mu},$$
(6.18)

where the last inequality follows from the Cauchy–Schwarz inequality for the inner product $m(\cdot, \cdot)$, (2.8), and (6.4). We then set $v = e_N^{u,j}$, $q = e_N^{p,j}$ in (6.16), (6.17) and subtract the second from the first equation such that

$$\begin{aligned} \frac{1}{\Delta t} m(e_N^{u,j} - e_N^{u,j-1}, e_N^{u,j}) + \|e_N^{u,j}\|_{X,\mu}^2 &= r_N^{1,j}(e_N^{u,j}) - r_N^{2,j}(e_N^{p,j}) \\ &\leq \|r_N^{1,j}\|_{X'} \|e_N^{u,j}\|_X + \|r_N^{2,j}\|_{Y'} \|e_N^{p,j}\|_Y. \end{aligned}$$

Applying now (6.18) and (2.10) yields

$$\begin{split} \frac{1}{\Delta t} \, m(e_N^{u,j} - e_N^{u,j-1}, e_N^{u,j}) + \|e_N^{u,j}\|_{X,\mu}^2 \\ & \leq \frac{1}{\beta_{\mathrm{Br}}} \|r_N^{1,j}\|_{X'} \|r_N^{2,j}\|_{Y'} + \left(\|r_N^{1,j}\|_{X'} + \frac{\gamma_a}{\beta_{\mathrm{Br}}} \|r_N^{2,j}\|_{Y'} \right) \frac{\|e_N^{u,j}\|_{X,\mu}}{\sqrt{\alpha_a}} \\ & + \frac{1}{\Delta t} \frac{\sqrt{\gamma_m}}{\beta_{\mathrm{Br}}} \|r_N^{2,j}\|_{Y'} \|e_N^{u,j} - e_N^{u,j-1}\|_{\mu}, \end{split}$$

which can be further bounded from Young's inequality by

$$\leq \frac{1}{\beta_{\mathrm{Br}}} \|r_N^{1,j}\|_{X'} \|r_N^{2,j}\|_{Y'} + \frac{1}{2\alpha_a} \left(\|r_N^{1,j}\|_{X'} + \frac{\gamma_a}{\beta_{\mathrm{Br}}} \|r_N^{2,j}\|_{Y'} \right)^2 + \frac{1}{2} \|e_N^{u,j}\|_{X,\mu}^2 \\ + \frac{1}{2\Delta t} \frac{\gamma_m}{\beta_{\mathrm{Br}}^2} \|r_N^{2,j}\|_{Y'}^2 + \frac{1}{2\Delta t} \|e_N^{u,j} - e_N^{u,j-1}\|_{\mu}^2.$$

Rearranging terms, the inequality now reads

$$\begin{split} \frac{1}{\Delta t} \left(\|e_N^{u,j}\|_{\mu}^2 - \|e_N^{u,j-1}\|_{\mu}^2 \right) + \|e_N^{u,j}\|_{X,\mu}^2 \\ & \leq \frac{\|r_N^{1,j}\|_{X'}^2}{\alpha_a} + \frac{2}{\beta_{\mathrm{Br}}} \left(1 + \frac{\gamma_a}{\alpha_a} \right) \|r_N^{1,j}\|_{X'} \|r_N^{2,j}\|_{Y'} + \left(\frac{\gamma_m}{\Delta t} + \frac{\gamma_a^2}{\alpha_a} \right) \frac{\|r_N^{2,j}\|_{Y'}^2}{\beta_{\mathrm{Br}}^2}, \end{split}$$

and the result follows from applying the sum $\sum_{j=1}^{k}$, $e_N^{u,0} = 0$, and (2.29), (2.30), (6.13).

Again, in the special case of a symmetric problem, the error bounds given in Proposition 6.4.1 can be improved. We may then derive the following result.

Proposition 6.4.2. Let $a(\cdot, \cdot; \mu)$ be symmetric for all $\mu \in \mathcal{D}$. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\max}$, $k \in \mathbb{K}$, and $\alpha_a^{\text{LB}}(\mu)$, $\gamma_a^{\text{UB}}(\mu)$, $\beta_{\text{Br}}^{\text{LB}}(\mu)$, $\gamma_m^{\text{UB}}(\mu)$ satisfying (2.29), (2.30), and (6.13), we define

$$\Delta_{N}^{\text{sym},k}(\mu) \equiv \left[\Delta t \sum_{j=1}^{k} \frac{\|r_{N}^{1,j}(\cdot;\mu)\|_{X'}^{2}}{\alpha_{a}^{\text{LB}}(\mu)} + \frac{2}{\beta_{\text{Br}}^{\text{LB}}(\mu)} \left(1 + \sqrt{\frac{\gamma_{a}^{\text{UB}}(\mu)}{\alpha_{a}^{\text{LB}}(\mu)}}\right) \|r_{N}^{1,j}(\cdot;\mu)\|_{X'} \|r_{N}^{2,j}(\cdot;\mu)\|_{Y'} + \left(\frac{\gamma_{m}^{\text{UB}}(\mu)}{\Delta t} + \gamma_{a}^{\text{UB}}(\mu)\right) \frac{\|r_{N}^{2,j}(\cdot;\mu)\|_{Y'}^{2}}{(\beta_{\text{Br}}^{\text{LB}}(\mu))^{2}}\right]^{1/2}.$$
 (6.19)

Then, $\Delta_N^{\text{sym},k}(\mu)$ represents an upper bound for the error $e_N^u(\mu)$ measured in the "spatio-temporal" energy norm (6.12),

$$\|e_N^u(\mu)\|_{\ell^2(0,k;X)} \le \Delta_N^{\operatorname{sym},k}(\mu), \quad \forall \ k \in \mathbb{K}, \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max}.$$
(6.20)

Proof. Following the lines of the previous proof, we may now apply the Cauchy–Schwarz inequality for the inner product $a(\cdot, \cdot)$ to obtain

$$\beta_{\rm Br} \|e_N^{p,j}\|_Y \le \|r_N^{1,j}\|_{X'} + \sqrt{\gamma_a} \|e_N^{u,j}\|_{X,\mu} + \frac{\sqrt{\gamma_m}}{\Delta t} \|e_N^{u,j} - e_N^{u,j-1}\|_{\mu},$$

instead of (6.18). Proceeding as before, this yields

$$\begin{split} \frac{1}{\Delta t} \, m(e_N^{u,j} - e_N^{u,j-1}, e_N^{u,j}) + \|e_N^{u,j}\|_{X,\mu}^2 \\ & \leq \frac{1}{\beta_{\mathrm{Br}}} \|r_N^{1,j}\|_{X'} \|r_N^{2,j}\|_{Y'} + \left(\frac{\|r_N^{1,j}\|_{X'}}{\sqrt{\alpha_a}} + \frac{\sqrt{\gamma_a}}{\beta_{\mathrm{Br}}} \|r_N^{2,j}\|_{Y'}\right) \|e_N^{u,j}\|_{X,\mu} \\ & \quad + \frac{1}{\Delta t} \frac{\sqrt{\gamma_m}}{\beta_{\mathrm{Br}}} \|r_N^{2,j}\|_{Y'} \|e_N^{u,j} - e_N^{u,j-1}\|_{\mu}, \end{split}$$

and the statement again follows from applying Young's inequality, the sum $\sum_{j=1}^{k}$, $e_N^{u,0} = 0$, and (2.29), (2.30), (6.13).

6.4.2 $\varepsilon > 0$

We here derive rigorous upper bounds for the error $e_N^{\varepsilon}(\mu)$ measured in the "spatio-temporal" energy norm

$$\|(v^{j},q^{j})_{j\in\mathbb{K}}\|_{\ell^{2}(0,k;Z)} \equiv \left(\|v^{k}\|_{\mu}^{2} + \Delta t \sum_{j=1}^{k} \|v^{j}\|_{X,\mu}^{2} + \varepsilon \|q^{j}\|_{Y,\mu}^{2}\right)^{1/2}, \tag{6.21}$$

where $(v^j, q^j)_{j \in \mathbb{K}} \subseteq Z, k \in \mathbb{K}$.

In addition to the dual norms of the residuals (6.10) and (6.11), we again also rely on (Online-)efficient lower (and upper) bounds (4.29) to the truth coercivity constants (2.10) and (4.12) to formulate our reduced basis *a posteriori* error bounds.

We can then formulate the following result.

Proposition 6.4.3. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, $k \in \mathbb{K}$, and $\alpha_a^{\text{LB}}(\mu)$, $\alpha_c^{\text{LB}}(\mu)$ satisfying (4.29), we define

$$\Delta_{N}^{\varepsilon,k}(\mu) \equiv \left(\Delta t \sum_{j=1}^{k} \frac{\|r_{N}^{1,j}(\cdot;\mu)\|_{X'}^{2}}{\alpha_{a}^{\text{LB}}(\mu)} + \frac{\|r_{N}^{2,j}(\cdot;\mu)\|_{Y'}^{2}}{\varepsilon \alpha_{c}^{\text{LB}}(\mu)}\right)^{1/2}.$$
(6.22)

Then, $\Delta_N^{\varepsilon,k}(\mu)$ represents an upper bound for the error $e_N^{\varepsilon}(\mu)$ measured in the "spatio-temporal" energy norm (6.21),

$$\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)} \le \Delta_N^{\varepsilon,k}(\mu), \quad \forall \ k \in \mathbb{K}, \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max}.$$
(6.23)

Proof. Let μ be any parameter in \mathcal{D} , $N \in \mathbb{N}_{\max}$, and $k \in \mathbb{K}$. For clarity of exposition, we suppress the argument μ in this proof.

Take any $1 \leq j \leq k$. From (6.10), (6.11), and (6.6), the errors $e_N^{u,j} \in X$ and $e_N^{p,j} \in Y$ satisfy the equations

$$\begin{aligned} \frac{1}{\Delta t} \, m(e_N^{u,j} - e_N^{u,j-1}, v) + a(e_N^{u,j}, v) + b(v, e_N^{p,j}) &= r_N^{1,j}(v), \quad \forall \ v \in X, \\ b(e_N^{u,j}, q) - \varepsilon \, c(e_N^{p,j}, q) &= r_N^{2,j}(q), \quad \forall \ q \in Y. \end{aligned}$$

Setting here $v = e_N^{u,j}$, $q = e_N^{p,j}$ and subtracting the second from the first equation, we obtain

$$\frac{1}{\Delta t} m(e_N^{u,j} - e_N^{u,j-1}, e_N^{u,j}) + \|e_N^{u,j}\|_{X,\mu}^2 + \varepsilon \|e_N^{p,j}\|_{Y,\mu}^2 = r_N^{1,j}(e_N^{u,j}) - r_N^{2,j}(e_N^{p,j}) \\
\leq \|r_N^{1,j}\|_{X'} \|e_N^{u,j}\|_X + \|r_N^{2,j}\|_{Y'} \|e_N^{p,j}\|_Y. \quad (6.24)$$

On the right-hand side, we now use (2.10), (4.12), and Young's inequality so that

$$\begin{split} \|r_N^{1,j}\|_{X'}\|e_N^{u,j}\|_X + \|r_N^{2,j}\|_{Y'}\|e_N^{p,j}\|_Y &\leq \frac{\|r_N^{1,j}\|_{X'}}{\sqrt{\alpha_a}}\|e_N^{u,j}\|_{X,\mu} + \frac{\|r_N^{2,j}\|_{Y'}}{\sqrt{\alpha_c}}\|e_N^{p,j}\|_{Y,\mu} \\ &\leq \frac{1}{2}\left(\frac{\|r_N^{1,j}\|_{X'}^2}{\alpha_a} + \|e_N^{u,j}\|_{X,\mu}^2 + \frac{\|r_N^{2,j}\|_{Y'}^2}{\varepsilon\alpha_c} + \varepsilon \|e_N^{p,j}\|_{Y,\mu}^2\right); \end{split}$$

on the left-hand side, we use the Cauchy–Schwarz inequality for the inner product $m(\cdot, \cdot)$ followed by Young's inequality so that

$$m(e_N^{u,j} - e_N^{u,j-1}, e_N^{u,j}) \ge \|e_N^{u,j}\|_{\mu}^2 - \|e_N^{u,j-1}\|_{\mu} \|e_N^{u,j}\|_{\mu} \ge \frac{1}{2} \left(\|e_N^{u,j}\|_{\mu}^2 - \|e_N^{u,j-1}\|_{\mu}^2 \right)$$

Rearranging terms, the inequality (6.24) finally reads

$$\frac{1}{\Delta t} \Big(\|e_N^{u,j}\|_{\mu}^2 - \|e_N^{u,j-1}\|_{\mu}^2 \Big) + \|e_N^{u,j}\|_{X,\mu}^2 + \varepsilon \, \|e_N^{p,j}\|_{Y,\mu}^2 \le \frac{\|r_N^{1,j}\|_{X'}^2}{\alpha_a} + \frac{\|r_N^{2,j}\|_{Y'}^2}{\varepsilon\alpha_c},$$

and the statement follows from applying the sum $\sum_{j=1}^{k}$, $e_N^{u,0} = 0$, and (4.29).

6.5 Offline-Online Computational Procedure

Again, the basic strategy lies in the μ -affine dependence of the involved operators. All μ independent quantities in (2.7) can be formed and stored within a computationally expensive Offline stage, which is performed only once and the cost of which depends on the large finite element dimension \mathcal{N} . For any given parameter $\mu \in \mathcal{D}$, the reduced basis approximation $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, is then computed within a highly efficient Online stage; the cost does not depend on \mathcal{N} but only on the much smaller dimension of the reduced basis approximation space. The computation of the *a posteriori* error bounds again consists of two components: the calculation of the residual dual norms $||r_N^{1,k}(\cdot;\mu)||_{X'}$, $||r_N^{2,k}(\cdot;\mu)||_{Y'}$, $k \in \mathbb{K}$, and the calculation of the required lower and upper bounds (2.29), (2.30), (6.13) and (4.29), respectively, to the involved constants. Except for the additional mass term, the exact computational procedure does not significantly differ from the stationary case; therefore, we do not give further explanations at this point but refer the reader to §2.6, §4.6, and [37].

6.6 Construction of Reduced Basis Approximation Spaces

We now turn to the construction of the reduced basis approximation spaces $X_N, Y_N, N \in \mathbb{N}_{\text{max}}$. According to the reduced basis idea (see §2.5), the low-dimensional spaces X_N, Y_N are again constructed by exploiting the parametric structure of the problem: Basis functions are essentially given by truth solutions associated with several chosen parameter snapshots. To provide accurate approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu))$ for $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, for any parameter query, X_N and Y_N here not only have to appropriately represent the manifold \mathcal{M} induced by the parametric dependence (see §2.5), but also need to capture the causality associated with evolution in time. Again, keeping computational cost to a minimum, we aim to achieve this with as few basis functions as possible.

The POD greedy procedure represents an adaptive sampling process for parabolic problems that properly accounts for temporal and parametric causality: It combines the proper orthogonal decomposition (POD) method in k (see [68, 69]) and also §1.2) with the greedy procedure in μ (see §2.7, [10, 15]). To begin with, we briefly recall the optimality property of the POD as described in [68, 69]. For a given finite set $\mathcal{X}_{\mathcal{I}} \equiv \{\chi_1, \ldots, \chi_{\mathcal{I}}\} \subseteq X$ and $M_X \leq \dim(\operatorname{span}(\mathcal{X}_{\mathcal{I}}))$, the POD basis of rank M_X consists of M_X $(\cdot, \cdot)_X$ -orthonormal basis functions that approximate $\mathcal{X}_{\mathcal{I}}$ best in the sense that

$$\operatorname{span}(\operatorname{POD}_X(\mathcal{X}_{\mathcal{I}}, M_X)) = \arg \inf_{\substack{\mathcal{X} \subseteq \operatorname{span}(\mathcal{X}_{\mathcal{I}}) \\ \dim(\mathcal{X}) = M_X}} \left(\frac{1}{\mathcal{I}} \sum_{i=1}^{\mathcal{I}} \inf_{\chi \in \mathcal{X}} \|\chi_i - \chi\|_X^2 \right)^{1/2};$$

analogously, we denote by $\operatorname{POD}_Y(\mathcal{Y}_{\mathcal{I}}, M_Y)$ the POD basis of rank M_Y for a finite set $\mathcal{Y}_{\mathcal{I}} \subseteq Y$, $M_Y \leq \dim(\operatorname{span}(\mathcal{Y}_{\mathcal{I}}))$. Assuming that we are given a current pair (X_{N-1}, Y_{N-1}) of reduced basis approximation spaces, the POD greedy algorithm now proceeds as follows: In compliance with the greedy approach, it detects the parameter μ_N for which the (Online-)efficient reduced basis error bound attains its maximum over an exhaustive sample $\Sigma \subset \mathcal{D}$. For a prescribed $\Delta N \in \mathbb{K}$, we then compute the POD bases of rank ΔN associated with the truth solutions $u^{\varepsilon,k}(\mu_N)$ and $p^{\varepsilon,k}(\mu_N)$, $k \in \mathbb{K}$; more specifically, we compute $\operatorname{POD}_X(E^u, \Delta N)$ and $\operatorname{POD}_Y(E^p, \Delta N)$ for

$$E^{u} \equiv \{ u^{\varepsilon,k}(\mu_{N}) - \Pi_{X_{N-1}} u^{\varepsilon,k}(\mu_{N}) \mid k \in \mathbb{K} \},\$$

$$E^{p} \equiv \{ p^{\varepsilon,k}(\mu_{N}) - \Pi_{Y_{N-1}} p^{\varepsilon,k}(\mu_{N}) \mid k \in \mathbb{K} \},\$$

where $\Pi_{X_{N-1}}$ and $\Pi_{Y_{N-1}}$ refer to the $(\cdot, \cdot)_X$ - and $(\cdot, \cdot)_Y$ -orthogonal projections on the current reduced basis approximation spaces X_{N-1} and Y_{N-1} , respectively. Finally, the ΔN POD basis functions are appended to X_{N-1} and Y_{N-1} , and we obtain a subsequent pair (X_N, Y_N) . As before, this process is then repeated until a prescribed error tolerance is satisfied. We refer the reader to [37, 46, 47] for a detailed discussion of the POD greedy procedure, and to [63, 64] for an application to the Boussinesq and Fokker–Planck equations.

Algorithm 5 Adaptive Sampling Procedure for $\varepsilon = 0$

1: Choose $\Sigma \subset \mathcal{D}, \, \delta_{\text{tol}}, \delta_{\text{tol}}^{\beta} \in (0, 1), \, \Delta N \in \mathbb{K}, \, \text{and} \, \mu_1 \in \Sigma$ 2: Set $N \leftarrow 0, \, \mathcal{D}_N \leftarrow \{\}, \, \mathcal{D}' \leftarrow \{\}, \, N_Y \leftarrow 0, \, Y_N \leftarrow \{\}, \, N_X \leftarrow 0, \, X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N+1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $E^{p} = \{ p^{k}(\mu_{N}) - \Pi_{Y_{N-1}} p^{k}(\mu_{N}) \mid k \in \mathbb{K} \}$ 5: $N_Y \leftarrow N_Y + \Delta N, Y_N \leftarrow Y_{N-1} \oplus \operatorname{span}(\operatorname{POD}_Y(E^p, \Delta N))$ 6: if $\mu_N \notin \mathcal{D}'$, then 7: $E^{u} = \{ u^{k}(\mu_{N}) - \prod_{X_{N-1}} u^{k}(\mu_{N}) \mid k \in \mathbb{K} \}$ 8: $N_X \leftarrow N_X + \Delta N, X_N \leftarrow X_{N-1} \oplus \operatorname{span}(\operatorname{POD}_X(E^u, \Delta N))$ 9: end if 10: while (true) do 11:for all $\mu \in \Sigma$ do 12:Compute $(u_N^k(\mu), p_N^k(\mu)), k \in \mathbb{K}, \Delta_N(\mu), \text{ and}$ $\hat{d}_N^\beta(\mu) \equiv \max\left\{\frac{\beta_{\mathrm{Br}}^{\mathrm{UB}}(\mu) - \beta_N(\mu)}{\beta_{\mathrm{Br}}^{\mathrm{UB}}(\mu)}, 0\right\} \text{ (cf. (2.101))}$ 13:14: end for 15: $\mu'_N \equiv \arg \max_{\mu \in \Sigma} \Delta_N(\mu), \ \mu^* \equiv \arg \max_{\mu \in \Sigma} \hat{d}_N^\beta(\mu)$ 16:if $\hat{d}_N^{\beta}(\mu^*) < \delta_{\text{tol}}^{\beta}$, then 17: $\mu_{N+1} \equiv \mu'_N$ 18: break 19:end if 20: $\begin{array}{l} \underset{\mu \in \mathcal{D}' \cup \mathcal{D}_N}{\text{if } \min_{\mu \in \mathcal{D}' \cup \mathcal{D}_N} \frac{|\mu'_N - \mu|}{|\mu|}} \geq 0.1\%, \text{ then } \\ \mathcal{D}' \leftarrow \mathcal{D}' \cup \{\mu'_N\} \end{array}$ 21:22: $E^{u} = \{ u^{k}(\mu'_{N}) - \prod_{X_{N}} u^{k}(\mu'_{N}) \mid k \in \mathbb{K} \}$ $N_{X} \leftarrow N_{X} + \Delta N, \ X_{N} \leftarrow X_{N} \oplus \operatorname{span}(\operatorname{POD}_{X}(E^{u}, \Delta N))$ 23: 24:25:else $N_X \leftarrow N_X + 1, X_N \leftarrow X_N \oplus \operatorname{span} \{ T_{\mu^*} \varrho_N(\mu^*) \} \text{ (see (2.71), (2.100))}$ 26:27:end if end while 28:29: until $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 30: $N_{\max} \leftarrow N$

For our saddle point problems, we now couple the above procedure with stabilization techniques developed in §2.7. In §3.3.3 and §5.3.2, best convergence results were achieved by Algorithm 3 where we aim to stabilize through an enrichment of the primal reduced basis approximation space with additional truth solutions. According to these observations, we here apply the sampling procedures presented in Algorithm 5 and Algorithm 6. In case of $\varepsilon = 0$, we again use the distance $d_N^{\beta}(\mu)$ (see (2.101)) of the inf-sup constants $\beta_N(\mu)$ to the truth Brezzi inf-sup constants $\beta_{\text{Br}}(\mu)$ as an indicator whether a current pair of reduced basis approximation spaces needs to be stabilized; the exact procedure is given in Algorithm 5. In case of $\varepsilon > 0$, numerical results in §5.3.2 showed that the inf-sup constants $\beta_N(\mu)$ may not be appropriate indicators for an ill-conditioned system but an adaptive sampling process should be based rather on the condition number $\kappa_N^{\varepsilon}(\mu)$ (see (5.6)); Algorithm 6 now presents a possibility how this could be realized.

Algorithm 6 Adaptive Sampling Procedure for $\varepsilon > 0$ 1: Choose $\Sigma \subset \mathcal{D}, \, \delta_{\text{tol}} \in (0, 1), \, \delta_{\text{tol}}^{\kappa} > 0, \, \Delta N \in \mathbb{K}, \, \text{and} \, \mu_1 \in \Sigma$ 2: Set $N \leftarrow 0$, $\mathcal{D}_N \leftarrow \{\}$, $\mathcal{D}' \leftarrow \{\}$, $N_Y \leftarrow 0$, $Y_N \leftarrow \{\}$, $N_X \leftarrow 0$, $X_N \leftarrow \{\}$ 3: repeat $N \leftarrow N + 1, \mathcal{D}_N \leftarrow \mathcal{D}_{N-1} \cup \{\mu_N\}$ 4: $E^{p} = \{ p^{\varepsilon,k}(\mu_{N}) - \prod_{Y_{N-1}} p^{\varepsilon,k}(\mu_{N}) \mid k \in \mathbb{K} \}$ 5: $N_Y \leftarrow N_Y + \Delta N, \ Y_N \leftarrow Y_{N-1} \oplus \operatorname{span}(\operatorname{POD}_Y(E^p, \Delta N))$ 6: if $\mu_N \notin \mathcal{D}'$, then 7: $E^{u} = \{ u^{\varepsilon,k}(\mu_N) - \prod_{X_{N-1}} u^{\varepsilon,k}(\mu_N) \mid k \in \mathbb{K} \}$ 8: $N_X \leftarrow N_X + \Delta N, X_N \leftarrow X_{N-1} \oplus \operatorname{span}(\operatorname{POD}_X(E^u, \Delta N))$ 9: end if 10:while (true) do 11: for all $\mu \in \Sigma$ do 12:Compute $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}, \Delta_N(\mu), \text{ and } \kappa_N^{\varepsilon}(\mu) \text{ (see (5.6))}$ 13:end for 14: $\mu'_N \equiv \arg \max_{\mu \in \Sigma} \Delta_N(\mu), \ \mu^* \equiv \arg \max_{\mu \in \Sigma} \kappa_N^{\varepsilon}(\mu)$ 15:if $\kappa_N^{\varepsilon}(\mu^*) < \delta_{\text{tol}}^{\kappa}$, then 16:17: $\mu_{N+1} \equiv \mu'_N$ 18:break end if 19: $\begin{array}{l} \underset{\mu \in \mathcal{D}' \cup \mathcal{D}_N}{\text{if } \min_{\mu \in \mathcal{D}' \cup \mathcal{D}_N} \frac{|\mu'_N - \mu|}{|\mu|}} \geq 0.1\%, \text{ then } \\ \mathcal{D}' \leftarrow \mathcal{D}' \cup \{\mu'_N\} \end{array}$ 20: 21: $E^{u} = \{ u^{\varepsilon,k} (\mu'_{N}) - \Pi_{X_{N}} u^{\varepsilon,k} (\mu'_{N}) \mid k \in \mathbb{K} \}$ $N_{X} \leftarrow N_{X} + \Delta N, \ X_{N} \leftarrow X_{N} \oplus \operatorname{span}(\operatorname{POD}_{X}(E^{u}, \Delta N))$ 22: 23:else 24: $N_X \leftarrow N_X + 1, X_N \leftarrow X_N \oplus \operatorname{span}\{T_{\mu^*}\varrho_N(\mu^*)\} \text{ (see (2.71), (2.100))}$ 25:26:end if end while 27:28: **until** $\Delta_N(\mu_{N+1}) < \delta_{\text{tol}}$ 29: $N_{\max} \leftarrow N$

6.7 Model Problem

i

We again consider a Stokes flow in a two-dimensional microchannel with an obstacle as introduced in §3.2; evolution in time is now induced by a time-dependent velocity profile on the inflow boundary.

Let μ be any parameter in \mathcal{D} . For the physical domain $\tilde{\Omega}$ and a given time interval [0, T], T > 0, we now seek to find the (inhomogeneous) velocity $\tilde{u}_{e,inh} : \tilde{\Omega} \times (0, T) \to \mathbb{R}^2$ and the pressure $\tilde{p}_e : \tilde{\Omega} \times (0, T) \to \mathbb{R}$ satisfying

$$\frac{\partial \tilde{u}_{\rm e,inh}}{\partial t} - \tilde{\Delta} \tilde{u}_{\rm e,inh} + \tilde{\nabla} \tilde{p}_{\rm e} = 0 \quad \text{in } \tilde{\Omega} \times (0, T), \tag{6.25}$$

$$\tilde{\nabla} \cdot \tilde{u}_{e,\text{inh}} = 0 \quad \text{in } \tilde{\Omega} \times (0,T),$$
(6.26)

subject to initial conditions $\tilde{u}_{e,inh}(\cdot, 0) = 0$ and with boundary conditions

$$\tilde{u}_{e,inh}(\tilde{x},t) = H(t)h(\tilde{x}) \quad \text{on } \Gamma_{in} \times (0,T), \quad \tilde{u}_{e,inh} = 0 \quad \text{on } \tilde{\Gamma}_0 \times (0,T),$$

$$\frac{\partial \tilde{u}_{e,inh}}{\partial \tilde{n}} = \tilde{p}_{e}\tilde{n} \quad \text{on } \Gamma_{out} \times (0,T);$$
(6.27)

here, $\tilde{\Delta}$ and $\tilde{\nabla}$ again denote the Laplacian and gradient operator over the physical domain $\tilde{\Omega}$ (see Appendix A), \tilde{n} is the unit outward normal, $h : \mathbb{R}^2 \to \mathbb{R}^2$ is given as in §3.2.1 by $h(x) \equiv (4x_2(1-x_2), 0)$ for all $x = (x_1, x_2) \in \mathbb{R}^2$, and we choose $H : [0,T] \to \mathbb{R}$ with $H(t) \equiv t(\sin(2\pi t) + 1)$ for all $t \in [0,T]$. According to the setting introduced in §5.2, we also consider the following perturbation of the problem (6.25)–(6.27): For a sufficiently small $\varepsilon > 0$, we again introduce a penalty term to the continuity equation (6.26) such that

$$\tilde{\nabla} \cdot \tilde{u}_{\mathrm{e,inh}}^{\varepsilon} = -\varepsilon \, p_{\mathrm{e}}^{\varepsilon} \quad \text{in } \tilde{\Omega} \times (0, T).$$
 (6.28)

We now follow the steps discussed in §3.2.2, §3.2.3, and §5.2: We choose the lifting function $\tilde{u}_{\rm L}^{H} \equiv H\tilde{u}_{\rm L}$ where $\tilde{u}_{\rm L}$ is defined as in (3.4), and transform the problem statement for the homogeneous velocity $\tilde{u}_{\rm e}^{\varepsilon} \equiv \tilde{u}_{{\rm e},{\rm inh}}^{\varepsilon} - \tilde{u}_{\rm L}^{H}$ to an equivalent problem posed over the reference domain Ω . Furthermore, as required for the time-discrete setting introduced in §6.2, we divide the time interval [0,T] into K subintervals of equal length $\Delta t \equiv T/K$, and consider a backward Euler method for time integration. The problems (6.25)–(6.27) and (6.25), (6.28), (6.27) may thus be written as a parametrized saddle point problem of the form (6.3). Here, for any $\mu \in \mathcal{D}$, the bilinear forms $a(\cdot, \cdot; \mu)$, $b(\cdot, \cdot; \mu)$, and $c(\cdot, \cdot; \mu)$ are given as in (3.8), (3.9), and (5.5), respectively; accordingly, the bilinear form $m(\cdot, \cdot; \mu) : X_{\rm e} \times X_{\rm e} \to \mathbb{R}$ represents the L^2 -inner product for vector functions over the physical domain $\tilde{\Omega}$ formulated on the reference domain Ω (see §3.2.3),

$$m(u, v; \mu) = \sum_{s=1}^{S} \frac{1}{|\det(A^s(\mu))|} \int_{\Omega^s} u \cdot v \, dx, \quad \forall \, u, v \in X_e,$$
(6.29)

and the linear functionals $f(\cdot; \mu)$ and $g(\cdot; \mu)$ are given by

$$\begin{split} f(v,t;\mu) &= f(v,t) = -H'(t) \int_{\Omega_{\rm L}} u_{\rm L} \cdot v \, dx - H(t) \int_{\Omega_{\rm L}} \frac{\partial u_{\rm Li}}{\partial x_j} \frac{\partial v_i}{\partial x_j} \, dx, \quad \forall \, v \in X_{\rm e}, \, t \in [0,T], \\ g(q,t;\mu) &= g(q,t) = H(t) \int_{\Omega_{\rm L}} q \frac{\partial u_{\rm Li}}{\partial x_i} \, dx, \quad \forall \, q \in Y_{\rm e}, \, t \in [0,T]. \end{split}$$

We recall that the bilinear forms $a(\cdot, \cdot; \mu)$, $b(\cdot, \cdot; \mu)$, and $c(\cdot, \cdot; \mu)$ then satisfy the assumptions (2.1)–(2.4), (4.1), and (4.2). For all $\mu \in \mathcal{D}$, $m(\cdot, \cdot; \mu)$ defines an inner product on X_e such that (6.2) holds true; moreover, there exists a constant $C_4^e(\mu) > 0$ from the Poincaré inequality (see, e.g., [89]) such that

$$m(v, v; \mu) \leq C_4^{\mathrm{e}}(\mu) a(v, v; \mu), \quad \forall v \in X_{\mathrm{e}}, \quad \forall \mu \in \mathcal{D},$$

and thus (6.1) is satisfied with

$$\gamma_m^{\mathrm{e}}(\mu) \equiv \sup_{u \in X_{\mathrm{e}}} \sup_{v \in X_{\mathrm{e}}} \frac{m(u, v; \mu)}{\|u\|_{X_{\mathrm{e}}} \|v\|_{X_{\mathrm{e}}}} \le C_4^{\mathrm{e}}(\mu) \, \gamma_a^{\mathrm{e}}(\mu) < \infty, \quad \forall \ \mu \in \mathcal{D}$$

We particularly have $Q_a = 10$, $Q_b = 6$, $Q_f = Q_g = 1$, and $Q_m = Q_c = 5$ in the respective μ -affine expansions (2.7). Choosing again the truth approximation spaces X and Y as the standard conforming \mathbb{P}_2 - \mathbb{P}_1 Taylor–Hood finite element approximation subspaces [100] over the regular triangulation \mathcal{T}_{Ω} , we ensure that also (2.11) is satisfied and therefore recover the situation described in §6.2.2.

6.8 Numerical Results

We now apply the reduced basis methodology developed in §6.3–§6.6 to our model problem. We set T = 1 and consider a constant time step size Δt corresponding to K = 100 time levels. The truth discretization is based on a fine mesh with a total of $\mathcal{N} = 72,076$ velocity and pressure degrees of freedom. In this section, all numerical results are attained using the open source software rb00mit [65], an implementation of the reduced basis framework within the C++ parallel finite element library libMesh [62].

6.8.1 $\varepsilon = 0$

We first turn to the coercivity, continuity, and inf-sup constants required for our reduced basis procedure. We again obtain (Online-)efficient lower and upper bounds to $\alpha_a(\mu)$, $\gamma_a(\mu)$, and $\beta_{\rm Br}(\mu)$ by using the SCM (see §2.6.2) with the configurations specified in §3.3.4. To estimate the continuity constants $\gamma_m(\mu)$, we apply the method for $M_{\alpha} = \infty$, an exhaustive sample $\Xi \subset \mathcal{D}$ of size $|\Xi| = 4,225$, and $\delta_{\rm tol}^{\rm SCM} = 0.01$. We then obtain $K_{\rm max} = 5$, and Fig. 6.1 presents the resulting accurate (Online-)efficient lower and upper bounds $\gamma_m^{\rm LB}(\mu)$ and $\gamma_m^{\rm UB}(\mu)$.

We now turn to the reduced basis approximation. To build our low-dimensional reduced basis approximation spaces $X_N, Y_N, N \in \mathbb{N}_{\max}$, we apply the POD greedy procedure described in Algorithm 5 (see §6.6). The sampling process is based on an exhaustive random sample $\Sigma \subset \mathcal{D}$ of size $|\Sigma| = 4,900, \Delta N = 2$, and $\delta_{tol}^{\beta} = 0.1$; since our Stokes model problem is clearly symmetric, we here in particular use the relative reduced basis *a posteriori* error bound $\Delta_N(\mu) \equiv \Delta_N^{\text{sym},K}(\mu)/||(u_N^j(\mu))_{j\in\mathbb{K}}||_{\ell^2(0,K;X)}$ (see (6.12), (6.19)). Figure 6.2 now shows the maximum error $||e_N^u(\mu)||_{\ell^2(0,K;X)}$ (see (6.9)) in the reduced basis velocity approximations and associated error bounds $\Delta_N^{\text{sym},K}(\mu)$ and $\Delta_N^K(\mu)$ (see (6.14)) as functions of the dimension N_Z ; Figure 6.3 presents the maximum error $||e_N^u(\mu)||_{\ell^2(0,k;X)}$ and associated error bounds $\Delta_N^{\text{sym},k}(\mu), \Delta_N^k(\mu)$ as functions of $k \in \mathbb{K}$ for several values of N_Z . First, we observe

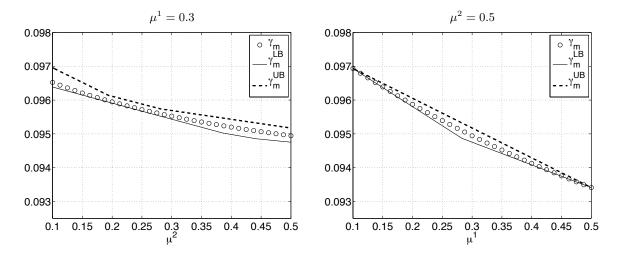


Figure 6.1.: SCM lower and upper bounds $\gamma_m^{\text{LB}}(\mu)$ and $\gamma_m^{\text{UB}}(\mu)$ for the continuity constants $\gamma_m(\mu)$ (see (6.4), (6.13)) where $|\Xi| = 4,225$, $M_\alpha = \infty$, $\delta_{\text{tol}}^{\text{SCM}} = 0.01$, and $K_{\text{max}} = 5$.

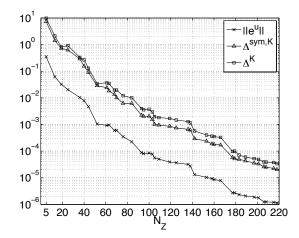


Figure 6.2.: Maximum error $\|e_N^u(\mu)\|_{\ell^2(0,K;X)}$ (see (6.9), (6.12)) and maximum error bounds $\Delta_N^{\operatorname{sym},K}(\mu)$ and $\Delta_N^K(\mu)$ (see (6.19) and (6.14)) normalized with respect to $\|(u^j(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,K;X)}$ shown as functions of N_Z ; the maximum is taken over 25 parameter values.

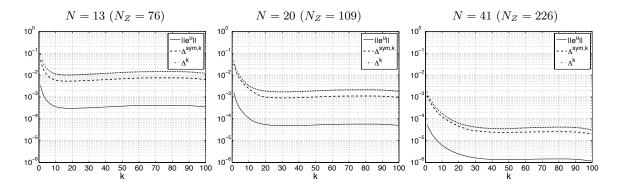


Figure 6.3.: Maximum error $\|e_N^u(\mu)\|_{\ell^2(0,k;X)}$ (see (6.9), (6.12)) and maximum error bounds $\Delta_N^{\operatorname{sym},k}(\mu)$ and $\Delta_N^k(\mu)$ (see (6.19) and (6.14)) normalized with respect to $\|(u^j(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,k;X)}$ shown as functions of $k\in\mathbb{K}$ for several values of N; the maximum is taken over 25 parameter values.

(a) பால		(µ) a	sociated	ΔN	(μ)					
N	N_Z	k = 10	k = 20	k = 30	k = 40	k = 50	k = 60	k = 70	k = 80	k = 90	k = 100
5	36	28.90	30.38	30.96	31.30	31.51	31.61	31.63	31.63	31.52	31.13
10	63	30.41	31.35	31.82	32.20	32.35	32.37	32.29	32.16	32.01	31.91
15	93	25.83	28.05	29.12	29.39	29.48	29.51	29.49	29.44	29.40	29.38
20	109	23.31	24.25	26.10	26.79	27.08	27.23	27.22	27.08	27.00	27.21
25	142	25.29	28.15	29.20	29.54	29.68	29.73	29.72	29.66	29.60	29.64
30	177	26.28	26.05	27.18	28.77	30.06	30.58	30.72	30.70	30.59	30.60
35	201	24.77	24.86	25.70	26.68	27.18	27.36	27.45	27.51	27.50	27.81
40	222	24.18	23.96	24.04	24.19	24.65	25.03	25.37	25.51	25.50	25.54
(b) Effec	tivities η_N^k	$V_{V}(\mu)$ assoc	iated with	n $\Delta^k_N(\mu)$						
N	N_Z	k = 10	k = 20	k = 30	k = 40	k = 50	k = 60	k = 70	k = 80	k = 90	k = 100
5	36	39.58	41.51	42.30	42.76	43.05	43.20	43.22	43.22	43.06	42.53
10	63	39.90	41.66	42.46	42.82	43.08	43.21	43.16	43.00	42.91	42.90
15	93	38.58	43.21	44.23	44.64	44.93	45.13	45.17	45.14	45.18	45.07
20	109	32.59	34.10	35.30	37.11	37.61	37.60	37.34	37.17	36.88	37.19
25	142	35.52	39.26	42.22	42.93	42.87	42.73	42.71	42.63	42.67	43.55
~ ~											
30	177	34.31	34.27	34.61	36.41	37.23	37.31	38.27	39.67	40.42	40.64
$\frac{30}{35}$	177 201	$34.31 \\ 32.86$	$34.27 \\ 33.59$	$34.61 \\ 35.05$	$36.41 \\ 36.52$	$37.23 \\ 37.33$	$37.31 \\ 37.39$	$38.27 \\ 36.98$	$39.67 \\ 36.65$	$40.42 \\ 36.77$	$40.64 \\ 37.54$

(a) Effectivities $\eta_N^{\text{sym},k}(\mu)$ associated with $\Delta_N^{\text{sym},k}(\mu)$

Table 6.1.: Maximum effectivities (a) $\eta_N^{\text{sym},k}(\mu) \equiv \Delta_N^{\text{sym},k}(\mu)/||e_N^u(\mu)||_{\ell^2(0,k;X)}$ (see (6.20)) and (b) $\eta_N^k(\mu) \equiv \Delta_N^k(\mu)/||e_N^u(\mu)||_{\ell^2(0,k;X)}$ (see (6.15)) for several values of $k \in \mathbb{K}$ and N; the maximum is taken over 25 parameter values.

that the reduced basis error and error bounds are roughly uniform in time (see Fig. 6.3) and decrease rapidly as N_Z increases (see Fig. 6.2). We obtain stable, rapidly convergent reduced basis approximations, and rigorous *a posteriori* error bounds that reflect the behavior of the error very accurately. Second, the error bounds are tight. To quantify this statement, we present in Table 6.1 maximum effectivities associated with $\Delta_N^{\text{sym},k}(\mu)$ and $\Delta_N^k(\mu)$ for several values of k and N. We notice that their values remain more or less constant with k. Moreover, as in the stationary case (see §3.3.4), we again benefit from exploiting the symmetry of the problem: Effectivities range from 33 to 45 in case of $\Delta_N^k(\mu)$ (see Table 6.1(b)) and improve in case of $\Delta_N^{\text{sym},k}(\mu)$ by roughly 10 (see Table 6.1(a)). We emphasize at this point that the error bound formulations in (6.14) and (6.19) in fact suggest a growth in time. In practice, this behavior seems rather weak (see Table 6.1) but may be investigated in greater detail within future work.

We now discuss the Online computation times for the proposed method. For comparison, once the μ -independent parts in (2.7) have been formed, direct computation of the truth approximation $(u^k(\mu), p^k(\mu)), k \in \mathbb{K}$, (i.e., assembly and solution of (6.6)) requires roughly 30 seconds on a 2.66 GHz Intel Core 2 Duo processor. We initially take a total reduced basis dimension of $N_Z = 226$. Once the database has been loaded, the Online calculation of $(u_N^k(\mu), p_N^k(\mu)), k \in \mathbb{K}$, (i.e., assembly and solution of (6.7)) and $\Delta_N^{\text{sym},k}(\mu), k \in \mathbb{K}$, for any new value of $\mu \in \mathcal{D}$ takes on average 27.97 and 80.76 milliseconds, respectively, which is in total roughly 270 times faster than direct computation of the truth approximation. Thus, even for this large value of N_Z , we obtain significant Online savings. In practice, however, we quite often need not take such a large value of N_Z — our rigorous and inexpensive error bounds $\Delta_N^{\text{sym},k}(\mu), k \in \mathbb{K}$, allow us to choose the reduced basis dimension just large enough to obtain a desired accuracy. To achieve a prescribed accuracy of at least 1%(resp., 0.1%) in the reduced basis approximations $u_N^k(\mu), k \in \mathbb{K}$, we need $N_Z = 76$ (resp., $N_Z = 109$) (see Fig. 6.2). Again, once the database has been loaded, the Online calculation of $(u_N^k(\mu), p_N^k(\mu)), k \in \mathbb{K}$, and $\Delta_N^{\text{sym},k}(\mu), k \in \mathbb{K}$, for any new value of $\mu \in \mathcal{D}$ takes on average 4.41 (resp., 7.62) and 24.47 (resp., 33.75) milliseconds, respectively, which is in total roughly 1,000 times (resp., 700 times) faster than direct computation of the truth approximation.

6.8.2 $\varepsilon > 0$

Again, the SCM (see §2.6.2) enables the (Online-)efficient estimation of the coercivity constants $\alpha_a(\mu)$ and $\alpha_c(\mu)$; as we here use the same configurations, we refer to §3.3.4 and §5.3.3 for details in this context.

To build our low-dimensional reduced basis approximation spaces $X_N, Y_N, N \in \mathbb{N}_{\text{max}}$, we apply the POD greedy procedure described in Algorithm 6 (see §6.6). The sampling process is based on an exhaustive random sample $\Sigma \subset \mathcal{D}$ of size $|\Sigma| = 4,900$, $\Delta N = 2, \delta_{\text{tol}}^{\kappa} = 10^3$, and the relative reduced basis a posteriori error bound $\Delta_N(\mu) = \Delta_N^{\varepsilon,K}(\mu)/||(u_N^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}||_{\ell^2(0,K;Z)}$ (see (6.21), (6.22)). Figure 6.4 now shows the maximum error $||e_N^{\varepsilon}(\mu)||_{\ell^2(0,K;Z)}$ (see (6.9)) in the reduced basis velocity and pressure approximations together with the associated error bound $\Delta_N^{\varepsilon,K}(\mu)$ as functions of the dimension N_Z for different values of ε . Figure 6.5 then presents the maximum error $||e_N^{\varepsilon}(\mu)||_{\ell^2(0,k;Z)}$ and associated error bound $\Delta_N^{\varepsilon,k}(\mu)$ as functions of $k \in \mathbb{K}$ for several values of N; note that the latter are chosen as the values for which the error bounds $\Delta_N^{\varepsilon,K}(\mu)$ guarantee a prescribed accuracy of at least 1% and

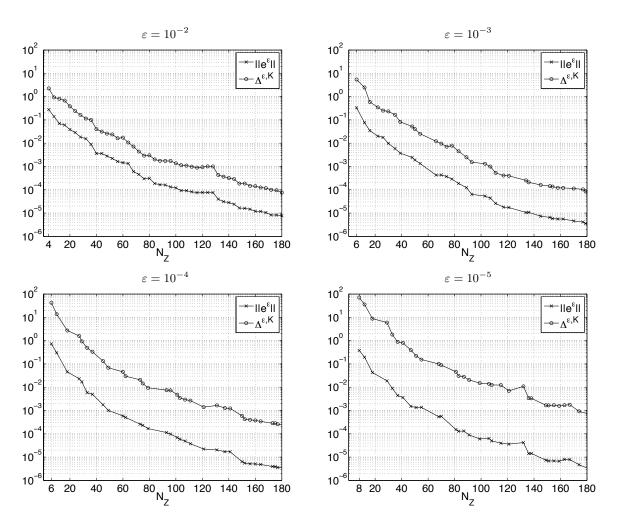


Figure 6.4.: Maximum error $\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,K;Z)}$ (see (6.9), (6.21)) and maximum error bound $\Delta_N^{\varepsilon,K}(\mu)$ (see (6.22)) normalized with respect to $\|(u^{\varepsilon,j}(\mu), p^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,K;Z)}$ shown as functions of N_Z for different values of ε ; the maximum is taken over 25 parameter values.

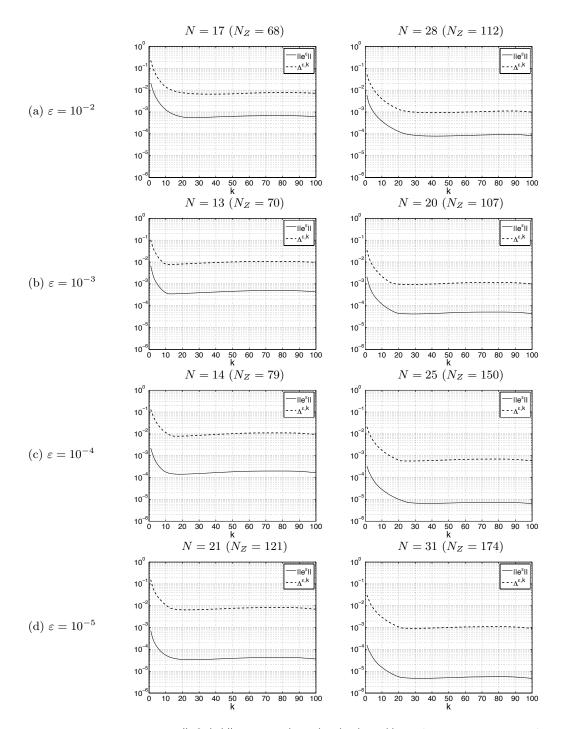


Figure 6.5.: Maximum error $\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)}$ (see (6.9), (6.21)) and maximum error bound $\Delta_N^{\varepsilon,k}(\mu)$ (see (6.22)) normalized with respect to $\|(u^{\varepsilon,j}(\mu), p^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,k;Z)}$ shown as functions of $k \in \mathbb{K}$ for several values of N for (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$; the maximum is taken over 25 parameter values.

(a)) $\varepsilon = 1$	0^{-2}					
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100
5	20	$1.145 \cdot 10^1$	$1.282 \cdot 10^1$	$1.297 \cdot 10^{1}$	$1.296 \cdot 10^{1}$	$1.295 \cdot 10^1$	$1.293 \cdot 10^{1}$
10	40	$1.154\cdot 10^1$	$1.373 \cdot 10^1$	$1.480 \cdot 10^{1}$	$1.479 \cdot 10^1$	$1.476 \cdot 10^1$	$1.466 \cdot 10^{1}$
15	60	$1.202\cdot 10^1$	$1.251 \cdot 10^1$	$1.292 \cdot 10^{1}$	$1.284 \cdot 10^1$	$1.281 \cdot 10^1$	$1.289 \cdot 10^1$
20	80	$1.119\cdot 10^1$	$1.129 \cdot 10^1$	$1.281 \cdot 10^{1}$	$1.314 \cdot 10^1$	$1.268 \cdot 10^1$	$1.268 \cdot 10^{1}$
25	100	$1.154\cdot 10^1$	$1.154\cdot 10^1$	$1.235 \cdot 10^{1}$	$1.248 \cdot 10^1$	$1.239\cdot 10^1$	$1.227 \cdot 10^{1}$
30	120	$1.256 \cdot 10^{1}$	$1.272 \cdot 10^1$	$1.256 \cdot 10^{1}$	$1.269 \cdot 10^1$	$1.268 \cdot 10^1$	$1.254 \cdot 10^{1}$
35	140	$1.132 \cdot 10^1$	$1.126 \cdot 10^{1}$	$1.163 \cdot 10^{1}$	$1.171 \cdot 10^{1}$	$1.159 \cdot 10^1$	$1.159 \cdot 10^{1}$
40	160	$1.252 \cdot 10^1$	$1.240 \cdot 10^{1}$	$1.223 \cdot 10^{1}$	$1.222 \cdot 10^{1}$	$1.202 \cdot 10^1$	$1.212 \cdot 10^{1}$
45	180	$1.241 \cdot 10^{1}$	$1.235 \cdot 10^{1}$	$1.226 \cdot 10^{1}$	$1.206 \cdot 10^{1}$	$1.201 \cdot 10^{1}$	$1.194 \cdot 10^{1}$
(b) $\varepsilon = 1$	0^{-3}					
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100
4	22	$2.691 \cdot 10^1$	$3.139 \cdot 10^1$	$3.273 \cdot 10^1$	$3.293 \cdot 10^1$	$3.275 \cdot 10^1$	$3.227 \cdot 10^1$
8	39	$2.326\cdot 10^1$	$2.499 \cdot 10^1$	$2.545 \cdot 10^{1}$	$2.560 \cdot 10^{1}$	$2.535 \cdot 10^1$	$2.486 \cdot 10^{1}$
12	66	$2.725 \cdot 10^1$	$2.969 \cdot 10^{1}$	$3.071 \cdot 10^1$	$3.166 \cdot 10^1$	$3.155 \cdot 10^1$	$3.101 \cdot 10^{1}$
16	83	$2.176 \cdot 10^1$	$2.759 \cdot 10^{1}$	$2.949 \cdot 10^{1}$	$3.005 \cdot 10^{1}$	$3.007 \cdot 10^1$	$2.991 \cdot 10^{1}$
19	103	$2.358 \cdot 10^1$	$2.367 \cdot 10^{1}$	$2.593 \cdot 10^{1}$	$2.641 \cdot 10^{1}$	$2.645 \cdot 10^{1}$	$2.710 \cdot 10^{1}$
23	121	$2.499 \cdot 10^1$	$2.530 \cdot 10^{1}$	$2.746 \cdot 10^{1}$	$2.814 \cdot 10^{1}$	$2.833 \cdot 10^1$	$2.947 \cdot 10^{1}$
26	145	$2.965 \cdot 10^{1}$	$2.933 \cdot 10^{1}$	$2.945 \cdot 10^{1}$	$2.973 \cdot 10^{1}$	$2.960 \cdot 10^{1}$	$2.983 \cdot 10^{1}$
30	162	$2.663 \cdot 10^{1}$	$2.648 \cdot 10^{1}$	$2.654 \cdot 10^{1}$	$2.759 \cdot 10^{1}$	$2.790 \cdot 10^{1}$	$2.792 \cdot 10^{1}$
34	183	$2.983 \cdot 10^{1}$	$2.902 \cdot 10^1$	$2.903 \cdot 10^{1}$	$2.900 \cdot 10^1$	$2.850 \cdot 10^{1}$	$2.826 \cdot 10^{1}$
(c)	$\varepsilon = 1$	0^{-4}					
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100
3	18	$8.804\cdot 10^1$	$8.882 \cdot 10^1$	$8.986 \cdot 10^1$	$9.150 \cdot 10^1$	$9.159 \cdot 10^1$	$9.113 \cdot 10^1$
7	37	$7.081 \cdot 10^{1}$	$7.651 \cdot 10^{1}$	$7.921 \cdot 10^{1}$	$8.026 \cdot 10^1$	$8.014 \cdot 10^1$	$7.921 \cdot 10^{1}$
10	60	$6.726 \cdot 10^1$	$8.104 \cdot 10^1$	$9.668 \cdot 10^{1}$	$9.737 \cdot 10^1$	$9.639 \cdot 10^1$	$9.608 \cdot 10^{1}$
14	79	$7.970 \cdot 10^{1}$	$9.903 \cdot 10^{1}$	$1.165 \cdot 10^2$	$1.166 \cdot 10^2$	$1.142 \cdot 10^2$	$1.145 \cdot 10^2$
17	101	$9.326 \cdot 10^{1}$	$1.010 \cdot 10^{2}$	$1.074 \cdot 10^{2}$	$1.055 \cdot 10^2$	$1.053 \cdot 10^{2}$	$1.072 \cdot 10^{2}$
21	121	$1.123 \cdot 10^{2}$	$1.117 \cdot 10^{2}$	$1.106 \cdot 10^2$	$1.107 \cdot 10^2$	$1.071 \cdot 10^2$	$1.076 \cdot 10^2$
24	141	$1.058 \cdot 10^2$	$1.055 \cdot 10^2$	$1.043 \cdot 10^2$	$1.033 \cdot 10^2$	$9.825 \cdot 10^{1}$	9.822 $\cdot 10^{1}$
28	160	$9.951 \cdot 10^{1}$	$9.855 \cdot 10^{1}$	$9.974 \cdot 10^{1}$	$9.988 \cdot 10^{1}$	$9.868 \cdot 10^{1}$	$9.903 \cdot 10^{1}$
33	181	$8.304 \cdot 10^{1}$	$8.290 \cdot 10^{1}$	$8.592 \cdot 10^{1}$	$8.635 \cdot 10^{1}$	$8.573 \cdot 10^{1}$	$8.706 \cdot 10^{1}$
(d) $\varepsilon = 1$						
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100
3	18	$2.706\cdot 10^2$	$2.919\cdot 10^2$	$3.311 \cdot 10^{2}$	$3.365 \cdot 10^2$	$3.199\cdot 10^2$	$3.198 \cdot 10^{2}$
7	41	$2.736 \cdot 10^2$	$2.899\cdot 10^2$	$2.965\cdot 10^2$	$2.984\cdot 10^2$	$2.984\cdot 10^2$	$2.952 \cdot 10^{2}$
10	55	$2.240 \cdot 10^{2}$	$2.510\cdot 10^2$	$2.684 \cdot 10^2$	$2.699\cdot 10^2$	$2.698\cdot 10^2$	$2.701 \cdot 10^2$
13	81	$3.073 \cdot 10^{2}$	$3.524 \cdot 10^{2}$	$3.603 \cdot 10^2$	$3.615 \cdot 10^{2}$	$3.609 \cdot 10^{2}$	$3.617 \cdot 10^2$
17	99	$2.696 \cdot 10^2$	$2.860 \cdot 10^{2}$	$3.103 \cdot 10^2$	$3.115 \cdot 10^{2}$	$3.107 \cdot 10^2$	$3.129 \cdot 10^2$
21	121	$3.355 \cdot 10^{2}$	$3.347 \cdot 10^{2}$	$3.391 \cdot 10^2$	$3.415 \cdot 10^2$	$3.396 \cdot 10^2$	$3.371 \cdot 10^2$
24	138	$2.563 \cdot 10^{2}$	$2.950 \cdot 10^{2}$	$3.214 \cdot 10^{2}$	$3.206 \cdot 10^2$	$3.147 \cdot 10^{2}$	$3.210 \cdot 10^2$
28	159	$2.690 \cdot 10^2$	$2.698 \cdot 10^2$	$2.935 \cdot 10^2$	$2.993 \cdot 10^2$	$2.893 \cdot 10^2$	$3.020 \cdot 10^2$
32	183	$2.786 \cdot 10^2$	$2.765 \cdot 10^2$	$3.206 \cdot 10^2$	$3.324\cdot 10^2$	$3.306 \cdot 10^2$	$3.368 \cdot 10^2$

Table 6.2.: Maximum effectivities $\eta_N^{\varepsilon,k}(\mu) \equiv \Delta_N^{\varepsilon,k}(\mu)/\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)}$ (see (6.23)) for several values of $k \in \mathbb{K}$ and N for (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$; the maximum is taken over 25 parameter values.

ε	N_Z	N	$(u_N^{\varepsilon,k}(\mu),p_N^{\varepsilon,k}(\mu)),k\in\mathbb{K}$	$\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$	Total
10^{-2}	68 (112)	17(28)	3.71(7.65)	14.53(26.95)	18.25(34.60)
10^{-3}	70 (107)	13 (20)	3.99(7.43)	17.19(28.25)	21.18(35.67)
10^{-4}	79 (150)	14(25)	4.73(13.59)	20.28(44.70)	25.01 (58.29)
10^{-5}	121 (174)	21 (31)	9.19(17.47)	33.81(54.08)	43.01 (71.55)

Table 6.3.: Average computation times in milliseconds for the Online evaluation of $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, (assembly and solution of (6.7)) and the error bounds $\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$, (see (6.22)) for different values of ε with a prescribed accuracy of at least 1% (resp., 0.1%) for the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$.

0.1% in the reduced basis approximations. First, we again observe that the reduced basis error and error bounds are roughly uniform in time (see Fig. 6.5) and decrease rapidly as N_Z increases (see Fig. 6.4). We obtain stable reduced basis approximations whose rapid convergence is not affected by the penalty parameter, and a posteriori error bounds that are meaningful and rigorous. Second, using the condition numbers $\kappa_N^{\varepsilon}(\mu)$ as an indicator for an ill-conditioned system, Algorithm 6 guarantees stability by properly accounting for the effects of the penalty term: For $\varepsilon = 10^{-2}$, the sampling process recognizes that the reduced basis approximation spaces X_N, Y_N do not have to be stabilized to provide accurate approximations; taking smaller values of ε and thus approaching the nonpenalized problem, an additional enrichment of the reduced basis approximation space for the velocity then becomes more and more necessary. Third, we see that the error bounds are tight for $\varepsilon = 10^{-2}$ but become less sharp as we decrease ε and our perturbed truth approximation becomes more accurate. However, effectivities exhibit a similar $O(\frac{1}{\sqrt{\varepsilon}})$ -dependence on the penalty parameter as observed in the stationary case (see §5.3.3) and remain reasonably small for relatively small values of ε . To further quantify this statement, we present in Table 6.2 the effectivities associated with $\Delta_N^{\varepsilon,k}(\mu)$ for different values of k, N, and ε . We note that their values are fairly constant with k and N and confirm the $O(\frac{1}{\sqrt{\varepsilon}})$ -dependence indicated by Fig. 6.4 and Fig. 6.5. As before, the effects of the penalty parameter on the effectivities are thus relatively being and we obtain useful bounds for reasonably small values of ε .

We close this section by discussing the Online computation times. For comparison, once the μ -independent parts in (2.7) have been formed, direct computation of the truth approximation $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, (i.e., assembly and solution of (6.6)) requires roughly 23 seconds on a 2.66 GHz Intel Core 2 Duo processor. Again, our rigorous and inexpensive reduced basis a posteriori error bounds enable us to choose the reduced basis dimension just large enough to obtain a desired accuracy. Choosing $\varepsilon = 10^{-2}$, the error bounds $\Delta_N^{\varepsilon,k}(\mu)$ are sharp with effectivities of approximately 12 (see Table 6.2(a)) and prescribe a dimension of $N_Z = 68$ to achieve an accuracy of at least 1% in the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$ (see Fig. 6.4). Once the database has been loaded, the Online calculation of $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, (i.e., assembly and solution of (6.7)) and $\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$, for any new value of $\mu \in \mathcal{D}$ then takes on average 3.71 and 14.53 milliseconds, respectively, which is in total roughly 1,200 times faster than direct computation of the truth approximation. Choosing smaller values for ε , the error bounds become more pessimistic and thus dictate a larger system dimension at which they guarantee the same order of accuracy. For $\varepsilon = 10^{-5}$, we need $N_Z = 121$ to achieve a prescribed accuracy of at least 1% in the reduced basis approximations (see Fig. 6.4); the Online calculation of $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, and $\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$, then takes on average 9.19 and 33.81 milliseconds, respectively, which is in total roughly 500 times faster than direct computation of the truth approximation. Thus, even for small penalty parameters ε , accurate approximations are guaranteed at significant Online savings. Detailed computation times for different values of ε are given in Table 6.3.

6.9 Concluding Remarks

In this chapter, we present new reduced basis methods for the instationary Stokes equations.

Combining techniques developed in Chapter 2 with current reduced basis approaches for parabolic problems, we derive new rigorous *a posteriori* bounds for the errors in the reduced basis velocity approximations and a POD greedy procedure that properly accounts for temporal and parametric causality as well as stability. The method provides rapidly convergent reduced basis approximations that are highly efficient and whose accuracy is certified by sharp and inexpensive *a posteriori* error bounds.

An approximation by penalty or regularization again allows for significant Offline savings at the expense of a less accurate truth approximation. Due to the introduced penalty term, an additional enrichment of the reduced basis velocity approximation space is not always necessary to obtain stable approximations; moreover, we obtain *a posteriori* error bounds that do not involve the expensive computation of inf-sup stability constants. As in the stationary case (see Chapter 5), the method provides reduced basis approximations and meaningful *a posteriori* error bounds that are computed very easily; nevertheless, drawbacks such as the disadvantageous dependence of the error bounds on the penalty parameter remain.

Time integration is achieved through a backward Euler method. Clearly, also other time integration schemes may be used. Using a Crank–Nicolson method, often preferred in practice due to its second-order accuracy, we may develop a penalty approach that is very similar to the one presented in this chapter and which is therefore included in Appendix C; in case of $\varepsilon = 0$, useful reduced basis a *posteriori* error bounds could not yet been derived and may therefore be — as well as a *posteriori* error bounds for the reduced basis pressure approximations — part of future work.

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Appendix A

Notations

In this section, we collect some basic notation that is used throughout this work.

Vectors and Matrices

The standard sets of numbers are denoted by \mathbb{N} , $\mathbb{N}_0 \equiv \mathbb{N} \cup \{0\}$, and \mathbb{R} . Let $n, m \in \mathbb{N}$. For a vector $x \in \mathbb{R}^n$, $|x| = \sqrt{x \cdot x}$ denotes the Euclidean norm. For a matrix $A = (A_{ij}) \in \mathbb{R}^{m \times n}$, det(A) refers to the determinant and $\sigma^{\max}(A)$, $\sigma^{\min}(A)$ to the maximum and minimum singular values, respectively. The matrix norms $\|\cdot\|_2$, $\|\cdot\|_1$, and $\|\cdot\|_\infty$ are defined as

$$\|A\|_{2} \equiv \sup_{x \in \mathbb{R}^{n}} \frac{|Ax|}{|x|} = \sigma^{\max}(A), \quad \|A\|_{1} \equiv \max_{j=1,\dots,n} \sum_{i=1}^{m} |A_{ij}|, \quad \|A\|_{\infty} \equiv \max_{i=1,\dots,m} \sum_{j=1}^{n} A_{ij};$$

here and throughout this thesis, we suppress the obvious requirement of nonzero elements in the denominators for clarity of exposition. Note that in particular $||A||_2 \leq \sqrt{||A||_1 ||A||_{\infty}}$ holds true. For a nonsingular matrix, we also consider the condition number $\kappa(A)$,

$$\kappa(A) \equiv ||A||_2 ||A^{-1}||_2 = \frac{\sigma^{\max}(A)}{\sigma^{\min}(A)}$$

Functions

Let $\Omega \subseteq \mathbb{R}^d$, $d \in \mathbb{N}$, be a bounded domain. Then, its boundary is denoted by $\Gamma \equiv \partial \Omega$ and $\overline{\Omega} \equiv \Omega \cup \Gamma$ refers to the closure of Ω . For a scalar function $p : \Omega \to \mathbb{R}$, $\frac{\partial p}{\partial x_i}$ denotes the partial derivative with respect to the variable $x_i, i = 1, \ldots, d$. The gradient, normal derivative, and Laplacian are then given by

$$\nabla p \equiv \left(\frac{\partial p}{\partial x_i}\right)_{i=1}^d, \quad \frac{\partial p}{\partial n} \equiv \nabla p \cdot n, \quad \Delta p \equiv \frac{\partial^2 p}{\partial x_i \partial x_i},$$

respectively, where n denotes the unit outward normal to Γ ; here and throughout this thesis, repeated indices imply summation. For a vector function $u = (u_1, \ldots, u_d) : \Omega \to \mathbb{R}^d$, partial derivatives are defined componentwise, i.e.,

$$\frac{\partial u}{\partial x_i} \equiv \left(\frac{\partial u_j}{\partial x_i}\right)_{j=1}^d, i = 1, \dots, d, \quad \frac{\partial u}{\partial n} \equiv \left(\frac{\partial u_j}{\partial n}\right)_{j=1}^d, \quad \Delta u \equiv (\Delta u_j)_{j=1}^d,$$

and we further consider the divergence operator

$$\operatorname{div}(u) \equiv \nabla \cdot u \equiv \frac{\partial u_i}{\partial x_i}.$$

In this work, $\tilde{\Omega} \equiv \Omega(\mu)$ typically refers to a domain whose exact geometric configuration is specified by a parameter $\mu \in \mathbb{R}^n$, $n \in \mathbb{N}$; its boundary is then denoted by $\tilde{\Gamma} \equiv \partial \tilde{\Omega}$ with unit outward normal \tilde{n} . According to the above notations, $\tilde{\nabla}$, $\frac{\partial}{\partial \tilde{n}}$, and $\tilde{\Delta}$ refer to the gradient, normal derivative, and Laplace operator over $\tilde{\Omega}$,

$$\begin{split} \tilde{\nabla}\tilde{p} &\equiv \left(\frac{\partial\tilde{p}}{\partial\tilde{x}_i}\right)_{i=1}^d, \quad \frac{\partial\tilde{p}}{\partial\tilde{n}} = \tilde{\nabla}\tilde{p}\cdot\tilde{n}, \quad \tilde{\Delta}\tilde{p} \equiv \frac{\partial^2\tilde{p}}{\partial\tilde{x}_i\partial\tilde{x}_i}, \\ \frac{\partial\tilde{u}}{\partial\tilde{n}} &\equiv \left(\frac{\partial\tilde{u}_j}{\partial\tilde{n}}\right)_{j=1}^d, \quad \tilde{\Delta}\tilde{u} = (\tilde{\Delta}\tilde{u}_j)_{j=1}^d, \quad \tilde{\nabla}\cdot\tilde{u} \equiv \frac{\partial\tilde{u}_i}{\partial\tilde{x}_i} \end{split}$$

for scalar and vector functions $\tilde{p}: \tilde{\Omega} \to \mathbb{R}$ and $\tilde{u} = (\tilde{u}_1, \dots, \tilde{u}_d): \tilde{\Omega} \to \mathbb{R}^d$.

Function Spaces

The space of continuous functions on Ω is given by $C(\Omega)$; for $j \in \mathbb{N}$, $C^{j}(\Omega)$ consists of all functions that are *j*-times continuously differentiable on Ω . We denote by $L^{2}(\Omega)$ the standard Lebesgue spaces of square integrable functions over Ω , and $H^{1}(\Omega) \equiv \{v \in L^{2}(\Omega) \mid \nabla v \in (L^{2}(\Omega))^{2}\}$ defines the associated Sobolev space of first order; for details on Sobolev spaces, we refer the reader to [1].

For a Hilbert space V, the associated inner product and norm are denoted by $(\cdot, \cdot)_V$ and $\|\cdot\|_V \equiv \sqrt{(\cdot, \cdot)_V}$, respectively. A second norm $\|\cdot\|$ on V is called equivalent to $\|\cdot\|_V$ if there exist constants c, C > 0 such that

$$c \parallel \cdot \parallel_V \leq \parallel \cdot \parallel \leq C \parallel \cdot \parallel_V.$$

The dual space of V is given by V' and $\langle \cdot, \cdot \rangle \equiv \langle \cdot, \cdot \rangle_{V' \times V}$ denotes by default the associated dual pairing. For an element $f \in V'$, we consider the dual norm

$$||f||_{V'} \equiv \sup_{v \in V} \frac{f(v)}{||v||_V}.$$

For a bounded linear operator $B: V \to V'$, ker(B) and im(B) refer to its kernel and range, respectively,

$$\ker(B) \equiv \{ v \in V \mid Bv = 0 \}, \quad \operatorname{im}(B) \equiv \{ f \in V' \mid \exists v \in V : Bv = f \};$$

moreover, $\ker(B)^{\perp} \equiv \{ v \in V \mid (v, v_0)_V = 0, \forall v_0 \in \ker(B) \}$ denotes the orthogonal complement with respect to the inner product $(\cdot, \cdot)_V$.

Appendix B

Supplementary Comments

In this section, we provide some additional comments on arguments from functional analysis used in this thesis.

On the Reduced Basis System (2.14)

We here comment on the steps outlined in §2.3.2 in a little more detail. To this end, let μ be any parameter in \mathcal{D} .

We first show that the system (2.15), (2.16) is solvable if and only if $g_N(\mu)$ belongs to the range of the operator $B_N(\mu)$. It clearly follows from (2.16) that $g_N(\mu) \in \operatorname{im}(B_N(\mu))$ is a necessary condition for (2.15), (2.16) to be solvable. Furthermore, if we assume that $g_N(\mu) \in$ $\operatorname{im}(B_N(\mu))$, we can without loss of generality find an element $u_N^{\perp}(\mu) \in \ker(B_N(\mu))^{\perp}$ such that $B_N(\mu)u_N^{\perp}(\mu) = g_N(\mu)$ in Y'_N . Due to (2.8) and (2.10), it then follows from the classical Lax–Milgram lemma (see, e.g., [89]) that there exists a unique solution $u_N^0(\mu) \in \ker(B_N(\mu))$ such that

$$a(u_N^0(\mu), v_N; \mu) = f(v_N; \mu) - a(u_N^{\perp}(\mu), v_N; \mu), \quad \forall v_N \in \ker(B_N(\mu)).$$
(B.1)

We now define $u_N(\mu) \equiv u_N^0(\mu) + u_N^{\perp}(\mu) \in X_N$. Then, $u_N(\mu)$ satisfies (2.16) by construction, and it follows from (B.1) that $f_N(\mu) - A_N(\mu)u_N(\mu)$ belongs to the annihilator

$$\ker(B_N(\mu))^0 \equiv \{ f \in X'_N \mid f(v_N) = 0, \forall v_N \in \ker(B_N(\mu)) \}$$

As X_N is a discrete space, ker $(B_N(\mu))^0$ corresponds to the range of the operator $B_N(\mu)^t$ (see, e.g., [13, §II.1, Proposition 1.2]); thus, we have $f_N(\mu) - A_N(\mu)u_N(\mu) \in \operatorname{im}(B_N(\mu)^t)$ and there exists $p_N(\mu) \in Y_N$ such that (2.15) is satisfied.

It remains to show that the solution $u_N(\mu) \in X_N$ is unique. For this purpose, let $(\bar{u}_N(\mu), \bar{p}_N(\mu)) \in X_N \times Y_N$ denote a second solution to (2.15), (2.16). It then follows from (2.15) that

$$a(u_N(\mu) - \bar{u}_N(\mu), v_N; \mu) = 0, \quad \forall v_N \in \ker(B_N(\mu)).$$

Moreover, (2.16) provides that $u_N(\mu) - \bar{u}_N(\mu) \in \ker(B_N(\mu))$; inserting this above yields $||u_N(\mu) - \bar{u}_N(\mu)||_{X,\mu} = 0$ and we obtain the desired result.

On the Residual Dual Norm (2.34)

We here comment on the dual norm of the combined residual $r_N(\cdot; \mu)$ given in (2.34). Let μ be any parameter in \mathcal{D} . As also used in §2.6, the Riesz representation theorem provides unique elements $\hat{e}_N^1(\mu) \in X$ and $\hat{e}_N^2(\mu) \in Y$ such that

$$(\hat{e}_N^1(\mu), v)_X = r_N^1(v; \mu), \quad \forall v \in X, \quad (\hat{e}_N^2(\mu), q)_Y = r_N^2(q; \mu), \quad \forall q \in Y,$$
 (B.2)

for the residuals defined in (2.32) and (2.33); in particular, the respective dual norms can then be represented as

$$\|r_N^1(\cdot;\mu)\|_{X'} = \|\hat{e}_N^1(\mu)\|_X, \quad \|r_N^2(\cdot;\mu)\|_{Y'} = \|\hat{e}_N^2(\mu)\|_Y.$$
(B.3)

Choosing now $\hat{e}_N(\mu) \equiv (\hat{e}_N^1(\mu), \hat{e}_N^2(\mu)) \in \mathbb{Z}$, we obtain from (B.2) that

$$(\hat{e}_N(\mu), (v, q))_Z = (\hat{e}_N^1(\mu), v)_X + (\hat{e}_N^2(\mu), q)_Y = r_N^1(v; \mu) + r_N^2(q; \mu) = r_N((v, q); \mu), \quad \forall (v, q) \in Z,$$

and thus $\hat{e}_N(\mu) \in Z$ defines the Riesz representer associated with the combined residual $r_N(\cdot;\mu) \in Z'$. Consequently,

$$||r_N(\cdot;\mu)||_{Z'} = ||\hat{e}_N(\mu)||_Z = \sqrt{||\hat{e}_N^1(\mu)||_X^2 + ||\hat{e}_N^2(\mu)||_Y^2},$$

and (2.34) follows from (B.3).

Appendix C

Re: Chapter 6 — A Crank–Nicolson Method

For time-integration, we may also employ a Crank–Nicolson method that is often used in practice due to its second-order accuracy (see, e.g., [27, 42, 89, 102]). We here present the corresponding reduced basis methodology for the perturbed problem.

General Problem Statement

Problem Formulation

We assume the setting introduced in §6.2.1. In addition, for a given $\theta \in [0,1]$, we now denote $f_{\theta}^k(\cdot;\mu) \equiv \theta f^k(\cdot;\mu) + (1-\theta)f^{k-1}(\cdot;\mu)$ and $g_{\theta}^k(\cdot;\mu) \equiv \theta g^k(\cdot;\mu) + (1-\theta)g^{k-1}(\cdot;\mu)$ for all $k \in \mathbb{K}$; accordingly, we also set $v_{\theta}^k \equiv \theta v^k + (1-\theta)v^{k-1}$ and $q_{\theta}^k \equiv \theta q^k + (1-\theta)q^{k-1}$ for finite sequences $(v^k)_{k\in\mathbb{K}} \subseteq X_e$ and $(q^k)_{k\in\mathbb{K}} \subseteq Y_e$.

Using a Crank–Nicolson method, our "exact" problem now reads as follows: For $\varepsilon > 0$ and any given $\mu \in \mathcal{D}$, we find $u_{e}^{\varepsilon,k}(\mu) \in X_{e}$ and $p_{e}^{\varepsilon,k}(\mu) \in Y_{e}$, $k \in \mathbb{K}_{0}$, such that $u_{e}^{\varepsilon,0}(\mu) = 0$ and

$$\frac{1}{\Delta t} m(u_{e}^{\varepsilon,k}(\mu) - u_{e}^{\varepsilon,k-1}(\mu), v; \mu)
+ a(u_{e,\theta}^{\varepsilon,k}(\mu), v; \mu) + b(v, p_{e,\theta}^{\varepsilon,k}(\mu); \mu) = f_{\theta}^{k}(v; \mu), \quad \forall v \in X_{e}, k \in \mathbb{K},
b(u_{e}^{\varepsilon,k}(\mu), q; \mu) - \varepsilon c(p_{e}^{\varepsilon,k}(\mu), q; \mu) = g^{k}(q; \mu), \quad \forall q \in Y_{e}, k \in \mathbb{K}_{0},$$
(C.1)

where $\theta = \frac{1}{2}$. Again, as the bilinear form $\frac{1}{\Delta t}m(\cdot, \cdot; \mu) + \theta a(\cdot, \cdot; \mu)$ is coercive on $X_{\rm e}$ for any $\mu \in \mathcal{D}$ from (6.2) and (2.3), the problem (C.1) is uniquely solvable for $(u_{\rm e}^{\varepsilon,k}(\mu), p_{\rm e}^{\varepsilon,k}(\mu))$, $k \in \mathbb{K}_0$, as a saddle point problem according to §4.2.1.

Truth Approximation

Assuming the setting introduced in §6.2.2, our high-fidelity "truth" discretization for (C.1) now reads as follows: For $\varepsilon > 0$ and any given $\mu \in \mathcal{D}$, we find $u^{\varepsilon,k}(\mu) \in X$ and $p^{\varepsilon,k}(\mu) \in Y$, $k \in \mathbb{K}_0$, such that $u^{\varepsilon,0}(\mu) = 0$ and

$$\frac{1}{\Delta t} m(u^{\varepsilon,k}(\mu) - u^{\varepsilon,k-1}(\mu), v; \mu)
+ a(u^{\varepsilon,k}_{\theta}(\mu), v; \mu) + b(v, p^{\varepsilon,k}_{\theta}(\mu); \mu) = f^{k}_{\theta}(v; \mu), \quad \forall v \in X, \ k \in \mathbb{K},
b(u^{\varepsilon,k}(\mu), q; \mu) - \varepsilon c(p^{\varepsilon,k}(\mu), q; \mu) = g^{k}(q; \mu), \quad \forall q \in Y, \ k \in \mathbb{K}_{0},$$
(C.2)

where $\theta = \frac{1}{2}$. Analogous to the exact problem, the problem (C.2) is uniquely solvable for $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}_0$, according to §4.2.2. The system can again be equivalently written in operator notation,

$$M(\mu) \frac{u^{\varepsilon,k}(\mu) - u^{\varepsilon,k-1}(\mu)}{\Delta t} + A(\mu) u^{\varepsilon,k}_{\theta}(\mu) + B(\mu)^t p^{\varepsilon,k}_{\theta}(\mu) = f^k_{\theta}(\mu) \quad \text{in } X', \ k \in \mathbb{K},$$
$$B(\mu) u^{\varepsilon,k}(\mu) - \varepsilon C(\mu) p^{\varepsilon,k}(\mu) = g^k(\mu) \quad \text{in } Y', \ k \in \mathbb{K}_0,$$

where $f_{\theta}^{k}(\mu) \equiv f_{\theta}^{k}(\cdot;\mu)|_{X} \in X'$ for all $k \in \mathbb{K}$.

Reduced Basis Approximation

Suppose that we are given a set of nested, low-dimensional reduced basis approximation subspaces $X_N \subset X_{N+1} \subset X$ and $Y_N \subset Y_{N+1} \subset Y$, $N \in \mathbb{N}_{\max} \equiv \{1, \ldots, N_{\max}\}$, then the reduced basis approximation is defined as the Galerkin projection with respect to the truth problem (C.2) onto these low-dimensional subspaces: For $\varepsilon > 0$ and any given $\mu \in \mathcal{D}$, we find $u_N^{\varepsilon,k}(\mu) \in X_N$ and $p_N^{\varepsilon,k}(\mu) \in Y_N$, $k \in \mathbb{K}_0$, such that $u_N^{\varepsilon,0}(\mu) = 0$ and

$$\frac{1}{\Delta t} m(u_N^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k-1}(\mu), v_N; \mu) + a(u_{N,\theta}^{\varepsilon,k}(\mu), v_N; \mu) + b(v_N, p_{N,\theta}^{\varepsilon,k}(\mu); \mu) = f_{\theta}^k(v_N; \mu), \quad \forall v_N \in X_N, \ k \in \mathbb{K}, b(u_N^{\varepsilon,k}(\mu), q_N; \mu) - \varepsilon c(p_N^{\varepsilon,k}(\mu), q_N; \mu) = g^k(q_N; \mu), \quad \forall q_N \in Y_N, \ k \in \mathbb{K}_0,$$
(C.3)

where $\theta = \frac{1}{2}$. The problem is uniquely solvable for $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}_0$, for any choice of X_N, Y_N (see §4.3.2). Written in operator notation, it also reads

$$\frac{1}{\Delta t} M_N(\mu) \Big(u_N^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k-1}(\mu) \Big)
+ A_N(\mu) u_{N,\theta}^{\varepsilon,k}(\mu) + B_N(\mu)^t p_{N,\theta}^{\varepsilon,k}(\mu) = f_{N,\theta}^k(\mu) \quad \text{in } X'_N, \ k \in \mathbb{K},
B_N(\mu) u_N^{\varepsilon,k}(\mu) - \varepsilon C_N(\mu) p_N^{\varepsilon,k}(\mu) = g_N^k(\mu) \quad \text{in } Y'_N, \ k \in \mathbb{K}_0,$$
(C.4)

where $A_N(\mu)$, $B_N(\mu)$, $B_N(\mu)^t$, $C_N(\mu)$, $M_N(\mu)$, and $g_N^k(\mu)$ are defined as in §6.3 and $f_{N,\theta}^k(\mu) \equiv f_{\theta}^k(\cdot;\mu)|_{X_N} \in X'_N$ for all $k \in \mathbb{K}$.

A Posteriori Error Estimation

We now develop upper bounds for the errors in our reduced basis approximations that are rigorous, sharp, and computationally efficient. As before, the errors shall be measured relative to the respective truth approximations.

For $\mu \in \mathcal{D}$, we consider the errors

$$e_N^u(\mu) \equiv (e_N^{u,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{u,k}(\mu) \equiv u^{\varepsilon,k}(\mu) - u_N^{\varepsilon,k}(\mu) \in X, \ k \in \mathbb{K}_0,$$

$$e_N^p(\mu) \equiv (e_N^{p,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{p,k}(\mu) \equiv p^{\varepsilon,k}(\mu) - p_N^{\varepsilon,k}(\mu) \in Y, \ k \in \mathbb{K}_0,$$

$$e_N^{\varepsilon}(\mu) \equiv (e_N^{\varepsilon,k}(\mu))_{k \in \mathbb{K}}, \text{ where } e_N^{\varepsilon,k}(\mu) \equiv (e_N^{u,k}(\mu), e_N^{p,k}(\mu)) \in Z, \ k \in \mathbb{K}_0,$$

(C.5)

in the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}_0$, with respect to the truth solution $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}_0$; we note that in particular $e_N^{u,0}(\mu) = 0$ from our initial conditions. The error in the reduced basis approximation shall then be measured in the "spatio-temporal" energy norm

$$\|(v^{j}, q^{j})_{j \in \mathbb{K}}\|_{\ell^{2}(0,k;Z)} \equiv \left(\|v^{k}\|_{\mu}^{2} + \Delta t \sum_{j=1}^{k} \|v_{\theta}^{j}\|_{X,\mu}^{2} + \varepsilon \|q_{\theta}^{j}\|_{Y,\mu}^{2}\right)^{1/2},$$
(C.6)

where $(v^j, q^j)_{j \in \mathbb{K}} \subseteq Z, k \in \mathbb{K}$.

To formulate our reduced basis a posteriori error bounds, we again need several ingredients: We rely on (Online-)efficient lower (and upper) bounds (4.29) to the truth coercivity constants (2.10) and (4.12) as well as on the residuals associated with the reduced basis approximation $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}_0$,

$$r_{N,\theta}^{1,k}(\cdot;\mu) \equiv f_{\theta}^{k}(\mu) - M(\mu) \frac{u_{N}^{\varepsilon,k}(\mu) - u_{N}^{\varepsilon,k-1}(\mu)}{\Delta t} - A(\mu) u_{N,\theta}^{\varepsilon,k}(\mu) - B(\mu)^{t} p_{N,\theta}^{\varepsilon,k}(\mu) \in X', \quad (C.7)$$

$$r_{N,\theta}^{2,k}(\cdot;\mu) \equiv g_{\theta}^{k}(\mu) - B(\mu) u_{N,\theta}^{\varepsilon,k}(\mu) + \varepsilon C(\mu) p_{N,\theta}^{\varepsilon,k}(\mu) \in Y', \quad (C.8)$$

where $g_{\theta}^{k}(\mu) \equiv g_{\theta}^{k}(\cdot;\mu)|_{X} \in X'$ for all $k \in \mathbb{K}$ and $\mu \in \mathcal{D}$.

We can now state the following result.

Proposition C.0.1. For any given $\mu \in \mathcal{D}$, $N \in \mathbb{N}_{\text{max}}$, $k \in \mathbb{K}$, and $\alpha_a^{\text{LB}}(\mu)$, $\alpha_c^{\text{LB}}(\mu)$ satisfying (4.29), we define

$$\Delta_N^{\varepsilon,k}(\mu) \equiv \left(\Delta t \sum_{j=1}^k \frac{\|r_{N,\theta}^{1,j}(\cdot;\mu)\|_{X'}^2}{\alpha_a^{\text{LB}}(\mu)} + \frac{\|r_{N,\theta}^{2,j}(\cdot;\mu)\|_{Y'}^2}{\varepsilon \alpha_c^{\text{LB}}(\mu)}\right)^{1/2}.$$
(C.9)

Then, $\Delta_N^{\varepsilon,k}(\mu)$ represents an upper bound for the error $e_N^{\varepsilon}(\mu)$ measured in the "spatio-temporal" energy norm (C.6),

$$\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)} \le \Delta_N^{\varepsilon,k}(\mu), \quad \forall \ k \in \mathbb{K}, \ \mu \in \mathcal{D}, \ N \in \mathbb{N}_{\max}.$$
(C.10)

Proof. Let μ be any parameter in \mathcal{D} , $N \in \mathbb{N}_{\max}$, and $k \in \mathbb{K}$. For clarity of exposition, we suppress the argument μ in this proof.

Take any $1 \le j \le k$. From (C.7), (C.8), and (C.2), the errors satisfy the equations

$$\begin{aligned} \frac{1}{\Delta t} m(e_N^{u,j} - e_N^{u,j-1}, v) + a(e_{N,\theta}^{u,j}, v) + b(v, e_{N,\theta}^{p,j}) &= r_{N,\theta}^{1,j}(v), \quad \forall \ v \in X, \\ b(e_{N,\theta}^{u,j}, q) - \varepsilon \ c(e_{N,\theta}^{p,j}, q) &= r_{N,\theta}^{2,j}(q), \quad \forall \ q \in Y. \end{aligned}$$

Setting here $v = e_{N,\theta}^{u,j}$, $q = e_{N,\theta}^{p,j}$ and subtracting the second from the first equation, we obtain

$$\frac{1}{\Delta t} m(e_N^{u,j} - e_N^{u,j-1}, e_{N,\theta}^{u,j}) + \|e_{N,\theta}^{u,j}\|_{X,\mu}^2 + \varepsilon \|e_{N,\theta}^{p,j}\|_{Y,\mu}^2 = r_{N,\theta}^{1,j}(e_{N,\theta}^{u,j}) - r_{N,\theta}^{2,j}(e_{N,\theta}^{p,j}) \\ \leq \|r_{N,\theta}^{1,j}\|_{X'} \|e_{N,\theta}^{u,j}\|_X + \|r_{N,\theta}^{2,j}\|_{Y'} \|e_{N,\theta}^{p,j}\|_Y.$$
(C.11)

On the right-hand side, we may now apply (2.10), (4.12), and Young's inequality such that

$$\begin{split} \frac{1}{\Delta t} \, m(e_N^{u,j} - e_N^{u,j-1}, e_{N,\theta}^{u,j}) + \|e_{N,\theta}^{u,j}\|_{X,\mu}^2 + \varepsilon \, \|e_{N,\theta}^{p,j}\|_{Y,\mu}^2 \\ & \leq \frac{\|r_{N,\theta}^{1,j}\|_{X'}}{\sqrt{\alpha_a}} \|e_{N,\theta}^{u,j}\|_{X,\mu} + \frac{\|r_{N,\theta}^{2,j}\|_{Y'}}{\sqrt{\alpha_c}} \|e_{N,\theta}^{p,j}\|_{Y,\mu} \\ & \leq \frac{1}{2} \left(\frac{\|r_{N,\theta}^{1,j}\|_{X'}^2}{\alpha_a} + \|e_{N,\theta}^{u,j}\|_{X,\mu}^2 + \frac{\|r_{N,\theta}^{2,j}\|_{Y'}^2}{\varepsilon \alpha_c} + \varepsilon \, \|e_{N,\theta}^{p,j}\|_{Y,\mu}^2 \right); \end{split}$$

rearranging terms then yields

$$\frac{2}{\Delta t} m(e_N^{u,j} - e_N^{u,j-1}, e_{N,\theta}^{u,j}) + \|e_{N,\theta}^{u,j}\|_{X,\mu}^2 + \varepsilon \|e_{N,\theta}^{p,j}\|_{Y,\mu}^2 \le \frac{\|r_{N,\theta}^{1,j}\|_{X'}^2}{\alpha_a} + \frac{\|r_{N,\theta}^{2,j}\|_{Y'}^2}{\varepsilon\alpha_c}.$$

As in particular $\theta = \frac{1}{2}$, the mass term on the left-hand side directly reads

$$\frac{1}{\Delta t} \left(\|e_N^{u,j}\|_{\mu}^2 - \|e_N^{u,j-1}\|_{\mu}^2 \right) + \|e_{N,\theta}^{u,j}\|_{X,\mu}^2 + \varepsilon \, \|e_{N,\theta}^{p,j}\|_{Y,\mu}^2 \le \frac{\|r_{N,\theta}^{1,j}\|_{X'}^2}{\alpha_a} + \frac{\|r_{N,\theta}^{2,j}\|_{Y'}^2}{\varepsilon\alpha_c}, \qquad (C.12)$$

and the statement follows from applying the sum $\sum_{j=1}^{k} e_N^{u,0} = 0$, and (4.29).

Numerical Results

We now apply the developed reduced basis methodology to the model problem described in §6.7. We set T = 1 and consider a constant time step size Δt corresponding to K = 100 time levels. The truth discretization is based on a fine mesh with a total of $\mathcal{N} = 72,076$ velocity and pressure degrees of freedom. In this section, all numerical results are attained using the open source software rb00mit [65], an implementation of the reduced basis framework within the C++ parallel finite element library libMesh [62].

To build our low-dimensional reduced basis approximation spaces $X_N, Y_N, N \in \mathbb{N}_{\max}$, we again apply the POD greedy procedure described in Algorithm 6 (see §6.6). The sampling process is based on an exhaustive random sample $\Sigma \subset \mathcal{D}$ of size $|\Sigma| = 4,900$, $\Delta N = 2, \, \delta_{tol}^{\kappa} = 10^3$, and the relative reduced basis *a posteriori* error bound $\Delta_N(\mu) = \Delta_N^{\varepsilon,K}(\mu)/||(u_N^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}||_{\ell^2(0,K;Z)}$ (see (C.6), (C.9)).

Using a Crank–Nicolson method, numerical results illustrate the same effects as observed from a backward Euler method; we therefore present the following data without further comments at this point and refer the reader to §6.8 and §6.9 for their evaluation.

Computation times for the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, and associated *a posteriori* error bounds $\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$, are given in Table C.1; for comparison, once the μ -independent parts in (2.7) have been formed, direct computation of the truth approximation $(u^{\varepsilon,k}(\mu), p^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, (i.e., assembly and solution of (C.2)) requires roughly 23 seconds on a 2.66 GHz Intel Core 2 Duo processor.

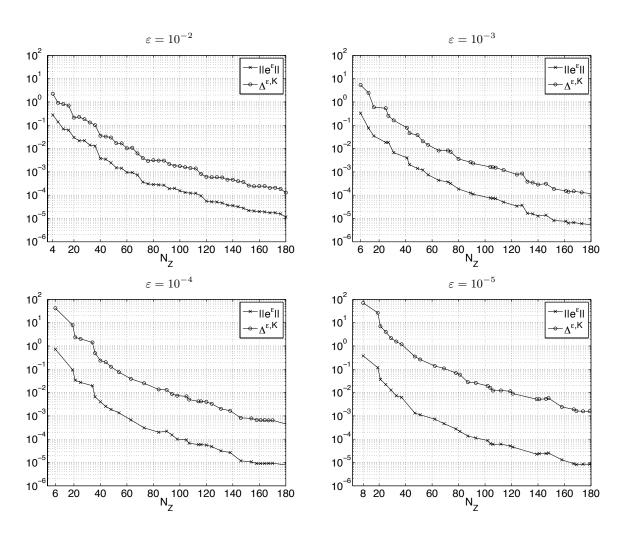


Figure C.1.: Maximum error $\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,K;Z)}$ (see (C.5), (C.6)) and maximum error bound $\Delta_N^{\varepsilon,K}(\mu)$ (see (C.9)) normalized with respect to $\|(u^{\varepsilon,j}(\mu), p^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,K;Z)}$ shown as functions of N_Z for different values of ε ; the maximum is taken over 25 parameter values.

ε	N_Z	N	$(u_N^{\varepsilon,k}(\mu),p_N^{\varepsilon,k}(\mu)),k\in\mathbb{K}$	$\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$	Total
10^{-2}	68 (116)	17(29)	3.71 (8.11)	$14.55\ (28.39)$	18.26(36.50)
10^{-3}	65(124)	12 (22)	3.60(9.46)	15.64(34.74)	19.24 (44.20)
10^{-4}	94 (146)	15(25)	6.30(13.03)	25.06(43.07)	31.37(56.10)
10^{-5}	121(202)	21 (33)	9.19(26.62)	33.94(69.19)	43.13 (95.81)

Table C.1.: Average computation times in milliseconds for the Online evaluation of $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$, (assembly and solution of (C.3)) and the error bounds $\Delta_N^{\varepsilon,k}(\mu), k \in \mathbb{K}$, (see (C.9)) for different values of ε with a prescribed accuracy of at least 1% (resp., 0.1%) for the reduced basis approximations $(u_N^{\varepsilon,k}(\mu), p_N^{\varepsilon,k}(\mu)), k \in \mathbb{K}$.

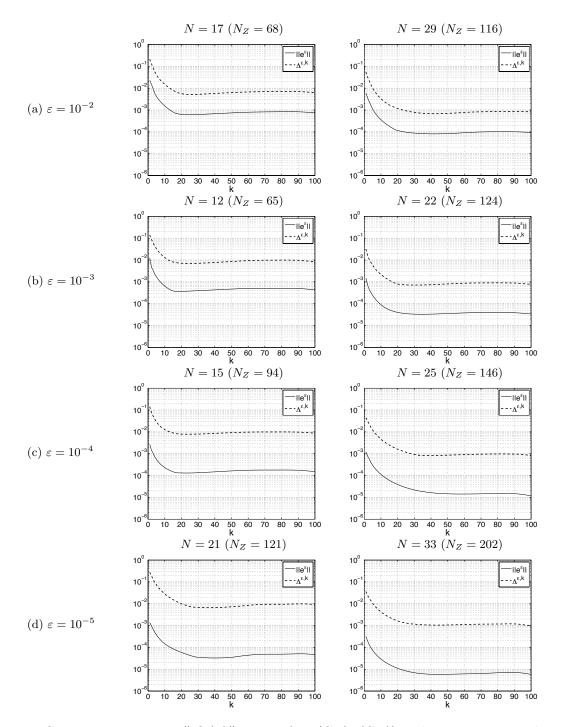


Figure C.2.: Maximum error $\|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)}$ (see (C.5), (C.6)) and maximum error bound $\Delta_N^{\varepsilon,k}(\mu)$ (see (C.9)) normalized with respect to $\|(u^{\varepsilon,j}(\mu), p^{\varepsilon,j}(\mu))_{j\in\mathbb{K}}\|_{\ell^2(0,k;Z)}$ shown as functions of $k \in \mathbb{K}$ for several values of N for (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$; the maximum is taken over 25 parameter values.

(a)	(a) $\varepsilon = 10^{-2}$							
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100	
5	20	$1.131 \cdot 10^1$	$1.211 \cdot 10^{1}$	$1.259 \cdot 10^{1}$	$1.268 \cdot 10^{1}$	$1.267 \cdot 10^1$	$1.252 \cdot 10^{1}$	
10	40	$1.115 \cdot 10^1$	$1.111 \cdot 10^1$	$1.133 \cdot 10^1$	$1.131 \cdot 10^1$	$1.126 \cdot 10^1$	$1.128 \cdot 10^1$	
15	60	$1.123 \cdot 10^1$	$1.168 \cdot 10^1$	$1.247 \cdot 10^1$	$1.253 \cdot 10^1$	$1.250 \cdot 10^1$	$1.248\cdot 10^1$	
20	80	$1.025 \cdot 10^1$	$1.056\cdot 10^1$	$1.152 \cdot 10^{1}$	$1.193 \cdot 10^1$	$1.186\cdot 10^1$	$1.180 \cdot 10^1$	
25	100	$1.260 \cdot 10^1$	$1.224 \cdot 10^1$	$1.154 \cdot 10^1$	$1.160 \cdot 10^1$	$1.162 \cdot 10^1$	$1.149 \cdot 10^1$	
30	120	$1.135 \cdot 10^1$	$1.175 \cdot 10^1$	$1.193 \cdot 10^1$	$1.260 \cdot 10^1$	$1.232\cdot 10^1$	$1.215 \cdot 10^1$	
35	140	$1.139 \cdot 10^1$	$1.193\cdot 10^1$	$1.248 \cdot 10^{1}$	$1.274 \cdot 10^{1}$	$1.283 \cdot 10^1$	$1.294 \cdot 10^1$	
40	160	$1.135 \cdot 10^1$	$1.133\cdot 10^1$	$1.210 \cdot 10^1$	$1.261 \cdot 10^1$	$1.259\cdot 10^1$	$1.244 \cdot 10^{1}$	
45	180	$1.137 \cdot 10^{1}$	$1.145 \cdot 10^{1}$	$1.150 \cdot 10^{1}$	$1.190 \cdot 10^{1}$	$1.203 \cdot 10^1$	$1.224 \cdot 10^{1}$	
(b) $\varepsilon = 1$	0^{-3}						
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100	
3	16	$2.432 \cdot 10^1$	$3.144 \cdot 10^{1}$	$3.393 \cdot 10^1$	$3.479 \cdot 10^1$	$3.442 \cdot 10^1$	$3.365\cdot 10^1$	
7	41	$3.403 \cdot 10^1$	$3.606 \cdot 10^1$	$3.809 \cdot 10^{1}$	$4.015 \cdot 10^1$	$3.835 \cdot 10^1$	$3.970\cdot 10^1$	
11	57	$2.316 \cdot 10^1$	$2.354 \cdot 10^1$	$2.715 \cdot 10^{1}$	$2.868 \cdot 10^1$	$2.881 \cdot 10^1$	$2.810 \cdot 10^1$	
15	80	$2.243 \cdot 10^{1}$	$2.583 \cdot 10^{1}$	$3.045 \cdot 10^1$	$3.177 \cdot 10^1$	$3.209 \cdot 10^1$	$3.186 \cdot 10^{1}$	
18	104	$2.484 \cdot 10^{1}$	$2.543 \cdot 10^{1}$	$2.758 \cdot 10^{1}$	$2.794 \cdot 10^{1}$	$2.741 \cdot 10^{1}$	$2.816 \cdot 10^{1}$	
22	124	$2.409 \cdot 10^{1}$	$2.414 \cdot 10^{1}$	$2.538 \cdot 10^{1}$	$2.791 \cdot 10^{1}$	$2.730 \cdot 10^{1}$	$2.703 \cdot 10^{1}$	
26	140	$2.526 \cdot 10^{1}$	$2.538 \cdot 10^{1}$	$2.589 \cdot 10^{1}$	$2.587 \cdot 10^{1}$	$2.542 \cdot 10^{1}$	$2.591 \cdot 10^{1}$	
29	161	$2.392 \cdot 10^{1}$	$2.412 \cdot 10^{1}$	$2.530 \cdot 10^{1}$	$2.562 \cdot 10^{1}$	$2.548 \cdot 10^{1}$	$2.614 \cdot 10^{1}$	
33	184	$2.574 \cdot 10^{1}$	$2.580\cdot 10^1$	$2.696 \cdot 10^{1}$	$2.790 \cdot 10^{1}$	$2.741 \cdot 10^{1}$	$2.727\cdot 10^1$	
(c)	$\varepsilon = 1$	0^{-4}						
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100	
2	19	$1.112 \cdot 10^{2}$	$1.250 \cdot 10^2$	$1.226 \cdot 10^{2}$	$1.280 \cdot 10^2$	$1.300 \cdot 10^2$	$1.273 \cdot 10^2$	
7	40	$9.345 \cdot 10^{1}$	$9.363 \cdot 10^1$	$1.176 \cdot 10^{2}$	$1.191 \cdot 10^{2}$	$1.167\cdot 10^2$	$1.123 \cdot 10^2$	
11	63	$7.780 \cdot 10^{1}$	$7.809 \cdot 10^1$	$8.898 \cdot 10^{1}$	$8.997 \cdot 10^1$	$9.095 \cdot 10^1$	$9.200 \cdot 10^1$	
13	84	$9.159 \cdot 10^1$	$9.247 \cdot 10^1$	$9.431 \cdot 10^{1}$	$9.236 \cdot 10^1$	$9.097 \cdot 10^1$	$9.132 \cdot 10^{1}$	
16	98	$7.357 \cdot 10^{1}$	$7.571 \cdot 10^1$	$1.034 \cdot 10^{2}$	$1.074 \cdot 10^{2}$	$1.067 \cdot 10^{2}$	$1.044 \cdot 10^{2}$	
21	120	$6.504 \cdot 10^{1}$	$7.179 \cdot 10^1$	$8.352 \cdot 10^{1}$	$8.453 \cdot 10^{1}$	$8.345 \cdot 10^{1}$	$8.272 \cdot 10^{1}$	
24	138	$6.587 \cdot 10^{1}$	$6.658 \cdot 10^{1}$	$7.325 \cdot 10^{1}$	$7.924 \cdot 10^{1}$	$8.110 \cdot 10^{1}$	$8.037 \cdot 10^{1}$	
28	161	$6.715 \cdot 10^{1}$	$6.829 \cdot 10^{1}$	$7.325 \cdot 10^{1}$	$8.367 \cdot 10^{1}$	$8.448 \cdot 10^{1}$	$8.577 \cdot 10^{1}$	
32	181	$7.303 \cdot 10^{1}$	$7.364 \cdot 10^1$	$7.573 \cdot 10^{1}$	$8.103 \cdot 10^{1}$	$8.291 \cdot 10^{1}$	$8.372 \cdot 10^{1}$	
(d) $\varepsilon = 1$							
N	N_Z	k = 10	k = 20	k = 40	k = 60	k = 80	k = 100	
2	19	$3.584 \cdot 10^{2}$	$4.002 \cdot 10^2$	$3.782 \cdot 10^2$	$3.938\cdot 10^2$	$4.090\cdot 10^2$	$4.029\cdot 10^2$	
7	37	$2.557 \cdot 10^2$	$3.281 \cdot 10^2$	$3.393 \cdot 10^2$	$3.369 \cdot 10^2$	$3.348\cdot 10^2$	$3.335 \cdot 10^2$	
10	62	$2.730 \cdot 10^2$	$2.891\cdot 10^2$	$3.212 \cdot 10^2$	$3.096 \cdot 10^2$	$3.087 \cdot 10^2$	$3.165 \cdot 10^2$	
13	81	$2.729 \cdot 10^{2}$	$3.452 \cdot 10^{2}$	$3.700 \cdot 10^2$	$3.714 \cdot 10^{2}$	$3.709 \cdot 10^{2}$	$3.730 \cdot 10^{2}$	
16	102	$2.616 \cdot 10^2$	$3.130 \cdot 10^{2}$	$3.189 \cdot 10^2$	$3.135 \cdot 10^{2}$	$3.124 \cdot 10^{2}$	$3.202 \cdot 10^{2}$	
21	121	$2.092 \cdot 10^2$	$2.249 \cdot 10^2$	$2.791 \cdot 10^2$	$2.969 \cdot 10^{2}$	$2.976 \cdot 10^2$	$3.033 \cdot 10^{2}$	
23	141	$2.572 \cdot 10^{2}$	$2.785 \cdot 10^{2}$	$3.404 \cdot 10^{2}$	$3.487 \cdot 10^{2}$	$3.485 \cdot 10^{2}$	$3.466 \cdot 10^2$	
26	158	$2.168 \cdot 10^2$	$2.240 \cdot 10^{2}$	$2.477 \cdot 10^2$	$2.531 \cdot 10^{2}$	$2.594 \cdot 10^2$	$2.657 \cdot 10^2$	
31	181	$2.479 \cdot 10^2$	$2.476 \cdot 10^2$	$2.641 \cdot 10^2$	$2.757\cdot 10^2$	$2.739\cdot 10^2$	$2.951\cdot 10^2$	

Table C.2.: Maximum effectivities $\eta_N^{\varepsilon,k}(\mu) \equiv \Delta_N^{\varepsilon,k}(\mu) / \|e_N^{\varepsilon}(\mu)\|_{\ell^2(0,k;Z)}$ (see (C.10)) for several values of $k \in \mathbb{K}$ and N for (a) $\varepsilon = 10^{-2}$, (b) $\varepsilon = 10^{-3}$, (c) $\varepsilon = 10^{-4}$, and (d) $\varepsilon = 10^{-5}$; the maximum is taken over 25 parameter values.

Curriculum Vitae

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