

Charge carrier mobility dependent open-circuit voltage in organic and hybrid solar cells

David Ompong and Jai Singh*

School of Engineering and IT, Faculty of EHSE, Charles Darwin University, Darwin, NT 0909, Australia

Abstract

A better understanding of the open-circuit voltage (V_{oc}) related losses in organic solar cells (OSCs) is desirable in order to assess their photovoltaic performance. We have derived V_{oc} as a function of charge carrier mobilities (μ_e and μ_h) for organic and hybrid solar cells by optimizing the drift-diffusion current density. The V_{oc} thus obtained depends on the energy difference between the highest occupied molecular orbital (HOMO) level and the quasi-Fermi level of holes of the donor material and on the ratio of the electron (μ_e) and hole (μ_h) mobilities in the blend. It is found that the V_{oc} increases with the increase of the mobility ratio μ_e / μ_h . The most loss in V_{oc} is contributed by the energetics of the donor and acceptor materials.

Introduction

Research interest in organic solar cells (OSCs) is currently on the increase mainly because of their cost effectiveness, flexibility, easy fabrication techniques, large scale production and the potential integration of OSCs into a wide variety of devices [1-4]. The development of new materials for photovoltaic applications coupled with device optimization has led to a dramatic increase in OSCs' performance in recent years [5]. A major research focus now lies in finding ways for further optimization of the power conversion efficiency (PCE), guided by a deeper understanding of the fundamental processes that influence the photovoltaic properties of OSCs [6]. The following four processes of OSCs and organic hybrid solar cells (OHSCs) make them remarkably different from their inorganic counterparts: i) photon absorption and exciton generation; ii) diffusion of excitons to the donor acceptor (DA) interface; iii) dissociation and charge separation at the interface; and iv) carrier collection by the electrodes [1,2]. These four processes have to be sufficiently efficient to reduce or eliminate energy losses leading to reduction in the short-circuit current density J_{sc} and open-circuit voltage V_{oc} , and hence, reduction in the power conversion efficiency of OSCs and OHSCs.

The current density J in the drift-diffusion model is a function of both the electrical and chemical potentials gradients, denoted by ∇U and ∇C , respectively. In OSCs, ∇U is negligible because there is no built-in electric field like the one in inorganic solar cells due to the property of p-n junction [7]. Therefore, in OSCs and OHSCs J depends mainly on the gradient of the chemical potential which is a function of V_{oc} as shown below. Thus, J becomes a function of V_{oc} and by optimizing J with respect to V_{oc} one can determine the optimal value of J corresponding to V_{oc} .

It is established that the V_{oc} of OSCs [8-11] depends on the energy difference between the highest occupied molecular orbital (HOMO) of the donor material and lowest unoccupied molecular orbital (LUMO) of the acceptor material or the conduction band of the inorganic nanoparticle in the case of OHSCs [12]. In addition, simulation [5,6] and experimental [13-15] works show that charge transport have effect on PCE of OSCs and a detailed analysis of bulk heterojunction organic solar cells reveals that low V_{oc} is the main factor limiting this efficiency

[9]. This implies that the V_{oc} of an OSC depends on the transport properties of the charge carrier in the material, which has not yet been studied adequately.

In this work, we have derived an analytical expression for V_{oc} by optimizing the drift-diffusion current density J . The V_{oc} thus obtained depends explicitly on the electron and hole mobilities and donor and acceptor HOMO and LUMO energy levels. In a previous study [5], the effective carrier mobility $\mu_{eff} = \sqrt{\mu_e \mu_h}$ is used to define the external voltage applied across an OSC, however in our approach the concept of the effective mobility is not used. Instead, it is found that the V_{oc} depends on the ratio of the electron (μ_e) to hole (μ_h) mobility such that if the ratio ($\mu_e / \mu_h = P$) increases the V_{oc} also increases.

Derivation of Open-Circuit Voltage (V_{oc})

The open-circuit voltage is given by the energy difference between the electron and hole quasi-Fermi levels [7]

$$-qV_{oc} = E_{F,e} - E_{F,h} \quad (1)$$

In OSCs and OHSCs the open-circuit voltage is also related to the HOMO energy level of the donor (E_{HOMO}^A) and the LUMO energy level of the acceptor (E_{LUMO}^A) as [16]:

$$-qV_{oc} = E_{HOMO}^D - E_{LUMO}^A - \Delta E_{loss} \quad (2)$$

where ΔE_{loss} is an empirical value representing energy losses in transporting charge carriers to the electrodes.

According to the drift-diffusion model the total current density J in a semiconductor under bias can be written as the sum of the electron

Correspondence to: Jai Singh, School of Engineering and IT, Faculty of EHSE, Charles Darwin University, Darwin, NT 0909, Australia, Tel: 414-805-9160, Fax: 414-805-9170; E-mail: Jai.Singh@cdu.edu.au

Keywords: charge carrier mobility, donor-acceptor, open-circuit voltage, organic solar cells, quasi-fermi levels

Received: October 29, 2015; **Accepted:** January 11, 2016; **Published:** January 14, 2016

and hole current densities, given by [17]:

$$J = J_n + J_p = \mu_e n \nabla E_{F,e} + \mu_h p \nabla E_{F,h}, \tag{3}$$

where $J_n = \mu_e n \nabla E_{F,e}$ is the electron current density and $J_p = \mu_h p \nabla E_{F,h}$ is the hole current density. Here $n(p)$ is the electron (hole) density, $\mu_e(\mu_h)$ is the electron (hole) mobility, and $\nabla E_{F,e}(\nabla E_{F,h})$ is the gradient of the electron (hole) quasi-Fermi level.

The charge-carrier densities n and p of electrons and holes inside the active layer are, respectively, given as [18]

$$n = N_c \exp[(E_{F,e} - E_{LUMO}^A) / k_B T], \tag{4}$$

and

$$p = N_v \exp[(E_{HOMO}^D - E_{F,h}) / k_B T], \tag{5}$$

where $N_c(N_v)$ is the effective density of states for the LUMO (HOMO) of acceptor (donor) material and $E_{F,e}(E_{F,h})$ is the energy of the corresponding Fermi levels. Using equations (1)-(5), the total current density in (3) can be written as a function of V_{oc} as:

$$J = \mu_e N_c \nabla E_{F,e} \exp[(-2qV_{oc} + E_{F,h} - E_{HOMO}^D + \Delta E_{loss}) / k_B T] + \mu_h N_v \nabla E_{F,h} \exp[(-2qV_{oc} + E_{LUMO}^A - E_{F,e} + \Delta E_{loss}) / k_B T] \tag{6}$$

The total current density J in equation (6) can be optimized with respect to V_{oc} as $\frac{d(J(V_{oc}))}{dV_{oc}} = 0$, which gives:

$$\mu_h N_v \nabla E_{F,h} \exp[(E_{LUMO}^A - E_{F,e}) / k_B T] = -\mu_e N_c \nabla E_{F,e} \exp[(E_{F,h} - E_{HOMO}^D) / k_B T], \tag{7}$$

In OSCs the chemical potential energy gradient ∇C drives the electrons and holes in the opposite direction [7], this explains the significance of the minus sign on the left hand side of equation (7); the minus sign is dropped from here onwards for convenience.

Multiplying both sides of equation (7) by $\exp[(E_{F,h} - E_{HOMO}^D) / k_B T]$ we get:

$$\exp[(-E_g + qV_{oc}) / k_B T] = \mu_e N_c \nabla E_{F,e} / \mu_h N_v \nabla E_{F,h} \exp[2(E_{F,h} - E_{HOMO}^D) / k_B T], \tag{8}$$

where $E_g = |E_{HOMO}^D - E_{LUMO}^A|$ is the effective band gap or the DA interface energy gap. Rearranging equation (8) we obtain V_{oc} as:

$$V_{oc} = \frac{1}{q} \left\{ E_g - 2(E_{F,h} - E_{HOMO}^D) + k_B T \ln \left(\frac{\mu_e N_c \nabla E_{F,e}}{\mu_h N_v \nabla E_{F,h}} \right) \right\}, \tag{9}$$

Following earlier works [5,18] we assume $N_c = N_v$ and $\nabla E_{F,e} = \nabla E_{F,h}$ which gives;

$$qV_{oc} = E_g - \Delta; \text{ where } \Delta = 2(E_{F,h} - E_{HOMO}^D) - k_B T \ln \left(\frac{\mu_e}{\mu_h} \right), \tag{10}$$

Here Δ is the energy loss contributed by the energetic (first term) and charge transport (second term)

Results

We have used equation (10) to calculate V_{oc} in several donor-acceptor (DA) materials listed in Table 1. The input parameters required for each DA in the calculations are also listed in Table 1. In addition, for calculating V_{oc} from equation (10) we need the values of the energy of the donor HOMO (E_{HOMO}^D) and acceptor LUMO (E_{LUMO}^A) which are listed in Table 2. It may be noted that following [18] we have used $(E_{F,h} - E_{HOMO}^D) \approx 0.2$ eV in equation (10) for all DA materials used in Tables 1 and 2. Using these input parameters the calculated values of V_{oc} are listed in Table 2 along with their experimental values obtained for these materials. According to Table 2, the calculated V_{oc} values are in reasonable agreement with those obtained experimentally.

According to equation (10) the V_{oc} increases if the ratio $P > 1$, that means, when the electron mobility is higher than the hole mobility as shown in Figure 1. In a material with equal mobility of electrons and holes, the contribution of the transport term to the V_{oc} vanishes.

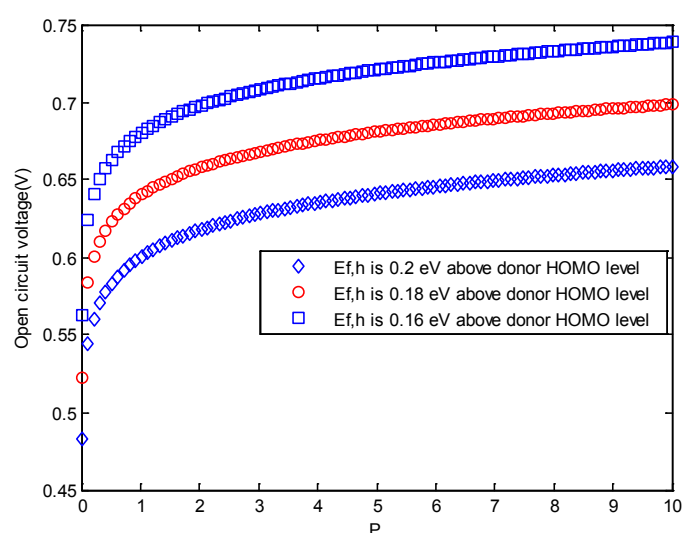
The analytical results of the dependence of V_{oc} on the charge carrier mobilities derived in equation (10) agree with the experimental

Table 1. Input values for calculating V_{oc} with donor –accepter materials forming the active layer, electron mobility μ_e , hole μ_h mobility and mobility ratio ratio P.

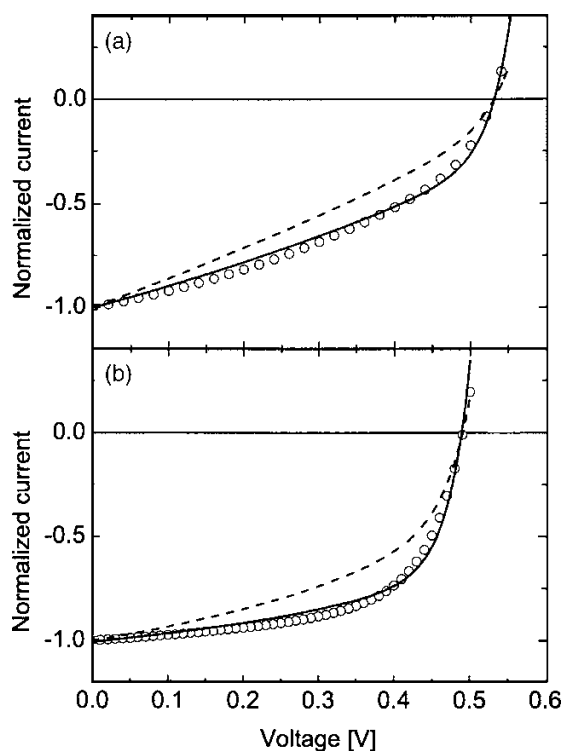
Entry	Active Layer	μ_e (cm ² V ⁻¹ s ⁻¹)	μ_h (cm ² V ⁻¹ s ⁻¹)	$P = \mu_e / \mu_h$	Ref.
OSC	PTB7:PCBM	1.0×10^{-3}	2.0×10^{-4}	5.0	[20]
OSC	PCDTBT:PCBM	2.9×10^{-3}	3.0×10^{-5}	96.7	[21]
OSC	P3HT:PCBM	$\times 10^{-3}$	$\times 10^{-4}$	10.0	[19]
OSC	MDMOPPV: PCBM	$\times 10^{-3}$	$\times 10^{-4}$	10.0	[19]
OSC	PBDTBDD:Bis-PCBM	9.6×10^{-5}	1.3×10^{-4}	0.7	[10]
OSC	PBDTBDD:PCBM	8.8×10^{-4}	1.4×10^{-3}	0.6	[10]
OSC	P3HT: Bis-PCBM	9.6×10^{-5}	1.0×10^{-4}	1.0	[10]
OSC	MEHPPV: PCBM	$\times 10^{-3}$	$\times 10^{-6}$	1000.0	[25]
OSC	Si-PCDTBT:PCBM	2.5×10^{-4}	3.0×10^{-5}	8.3	[22]
OHSC	MDMOPPV:nc-ZnO	2.8×10^{-5}	5.5×10^{-6}	5.1	[23]
OHSC	P3HT: Si-NCs	$\times 10^{-3}$	$\times 10^{-3}$	1.0	[24]

Table 2. Donor-Acceptor materials, Donor HOMO level E_{HOMO}^D , Acceptor LUMO level E_{LUMO}^A , Effective band gap $E_g = |E_{HOMO}^D - E_{LUMO}^A|$, transport loss term $k_B T \ln(\mu_e/\mu_h)$ and V_{oc}^{calc} from equation (10).

Donor material	E_{HOMO}^D (eV)	Acceptor material	E_{LUMO}^A (eV)	E_g (eV)	$k_B T \ln(\mu_e/\mu_h)$ (eV)	V_{oc}^{calc} (V)	V_{oc}^{expt} (V)	Ref.
PTB7	5.15	PCBM	4.06	1.09	0.04	0.73	0.75	[20]
PCDTBT	5.50	PCBM	4.30 [26]	1.20	0.12	0.92	0.85	[21]
P3HT	5.10	PCBM	4.06	1.04	0.06	0.69	0.63	[26]
MDMOPPV	5.36	PCBM	4.06	1.30	0.06	0.96	0.83	[11]
PBDTBDD	5.23	Bis-PCBM	3.80	1.43	-0.01	0.97	1.00	[10]
PBDTBDD	5.23	PCBM	3.94	1.29	-0.01	0.88	0.86	[10]
P3HT	5.10	Bis-PCBM	3.80	1.30	0.00	0.89	0.74	[10]
MEHPPV	5.20	PCBM	3.95	1.00	0.18	0.88	0.74	[13]
Si-PCPDTBT	4.86	PCBM	3.88	0.98	0.05	0.63	0.59	[22]
MDMOPPV	5.20	nc-ZnO	4.20	1.00	0.04	0.64	0.74	[12]
P3HT	5.10	Si-NCs	3.95	1.15	0.00	0.75	0.75	[24]

**Figure 1.** V_{oc} in equation (10) plotted as a function of electron: hole mobility ratio, $P = \mu_e/\mu_h$.

observation as well as with the numerical simulation [23]. In Figures 2a and 2b we have reproduced the J-V characteristics measured on P3HT:PCBM bulk heterojunction organic solar cells (BHJ OSCs) annealed at two different temperatures, 52°C and 70°C, respectively [23]. The measured mobility P3HT:PCBM of electrons and holes is found to be $\mu_e = 2.5 \times 10^{-8}$ ($\text{m}^2\text{V}^{-1}\text{s}^{-1}$), $\mu_h = 3.0 \times 10^{-12}$ ($\text{m}^2\text{V}^{-1}\text{s}^{-1}$) at 52°C (Figure 2a) and $\mu_e = 1.1 \times 10^{-7}$ ($\text{m}^2\text{V}^{-1}\text{s}^{-1}$), $\mu_h = 1.1 \times 10^{-10}$ ($\text{m}^2\text{V}^{-1}\text{s}^{-1}$) at 70°C (Figure 2b) [23]. Using these values, we find that the mobility ratio P decreases from 8.3×10^3 to 1.0×10^3 when one anneals the sample at 52°C and 70°C. According to equation (10), this means that one should get a higher value of V_{oc} at the annealing temperature of 52°C than at 70°C. This result is quite consistent with that shown in Figures 2a and 2b, where the measured and simulated V_{oc} at 52°C is about 0.04 V higher than that at 70°C. Mobility dependent J-V characteristics have also been simulated by assuming $\mu_e = \mu_h$ [5]. The V_{oc} is found to be independent of the charge carrier mobility in the range from 1 to 10^6 cm^2/Vs . According to equation (10) also, the mobility dependent term vanishes for $\mu_e = \mu_h$ and hence V_{oc} becomes constant which is consistent with this result.

**Figure 2.** Measured current-voltage characteristics normalized to the short-circuit current (open circles) of two P3HT/PCBM solar cells annealed at 52°C (a) and 70°C (b). The solid lines denote simulations using slowest carrier recombination constant $\gamma = \frac{\epsilon}{\epsilon} \min(\mu_e, \mu_h)$, while the dashed lines correspond to simulations using average carrier recombination constant $\gamma = \frac{\epsilon}{\epsilon} (\mu_e + \mu_h)$. ϵ is the dielectric constant (Reproduced with permission from (Koster *et al.* [23]. Copyright 2006, AIP Publishing LLC).

Abbreviations

PTB7: (poly[[4,8-bis[(2-ethylhexyl)oxy]benzo[1,2-b:4,5-b']dithiophene-2,6-diyl][3-fluoro-2-(2-ethylhexyl)carbonyl]thieno[3,4-b]thiophenediyl])

PCBM: 1-(3-methoxycarbonyl)-propyl-1-phenyl-(6,6)C

PCDTBT: poly[N-9'-hepta-decanyl-2,7-carbazole-alt-5,5-(4',7'-di-2-thienyl-2',1',3'-benzothiadiazole)]

P3HT: poly(3-hexylthiophene)

MDMOPPV:poly[2-methoxy-5-(3',7'-dimethyloctyloxy)-1-4-phenylene vinylene]

PBDTBDD:poly(((4,8-Bis(5-(2-ethylhexyl)thiophen-2-yl)benzo[1,2-b:4,5-b']dithiophene-2,6-diyl) bis(trimethyl))-co-(5,7-bis(2-ethylhexyl)benzo[1,2-c:4,5-c']dithiophene-4,8-dione))

Bis-PCBM: bisadduct of phenyl-C61-butyric acid methyl ester)

MEHPPV:poly[2-methoxy-5-(2-ethylhexyloxy)-1,4-phenylenevinylene]

Si-PCPDTBT:poly[2,1,3-benzothiadiazole-4,7-diyl[4,4-bis(2-ethylhexyl)-4H-cyclopenta2,1-b:3,4-b']dithiophene-siloe 2,6-diyl]]

nc-ZnO: Zinc oxide nanoparticles

Si NCs: Silicon nanocrystals

Discussions

According to equation (10) the open-circuit voltage becomes equal to the effective band gap energy and hence independent of the charge carrier mobilities when the hole quasi-Fermi level is equal to the HOMO level of the donor molecule and the electron and hole mobilities are equal. It is to be noted that the V_{oc} derived in equation (10), depends on the electron and hole mobilities directly. The material with $\mu_e < \mu_h$ will have greater energy loss Δ and hence lower V_{oc} in comparison with materials with $\mu_e > \mu_h$, which will have lesser Δ and hence higher V_{oc} . From this point of view, one may prefer to use materials with $\mu_e > \mu_h$ for obtaining higher V_{oc} in OSCs.

As stated above, in the calculation of V_{oc} from equation (10), we have assumed a constant value for $(E_{F,h} - E_{HOMO}^D) \approx 0.2$ eV, which is valid only if the charge carrier concentration remains constant and that means the mobilities of charge carriers are not very high or very low. For example, in OSCs based on P3HT:PCBM where a mobility ratio $P=10$ is considered [19], it is found that if both charge carrier mobilities at this ratio are high, then this will lead to the efficient extraction of charge carriers which reduces the charge carrier concentration. This reduction in carrier concentration is expected to draw the hole quasi Fermi level away from the HOMO level of the donor material, which according to equation (10) will reduce the V_{oc} . This will eventually reduce the PCE as found in [19]. Likewise, at low charge carrier mobilities at the same ratio, the recombination will be enhanced which will reduce the short circuit current [6,19], leading to reduction in PCE. In this view, the derived V_{oc} in equation (10) may be regarded to be valid only at moderate electron and hole mobilities leading to high PCE.

For highlighting the role of the charge carrier mobility, it may be desirable to consider the two DA combination materials MDMOPPV:PCBM and P3HT:Bis-PCBM in Table 2. These two combinations have the same effective gap of 1.30 eV but the second term of Δ in equation (10) is 0.06 eV for the first combination and zero for the second (Table 2). As a result the value of Δ is less in the first combination than that in the second, producing higher V_{oc} (0.96 eV) in MDMOPPV:PCBM in comparison with that of 0.89 eV in P3HT:Bis-PCBM. It may be interesting to note that, using $(E_{F,h} - E_{HOMO}^D) \approx 0.2$ eV in equation (10), we get, $\Delta = 0.4 - k_B T \ln(\mu_e/\mu_h)$ which shows that the loss of 0.4 eV due to the energy difference is much bigger than the second term due to the charge transport whose calculated values are listed in column 6 of Table 2.

Conclusion

We have derived a mobility dependant expression for V_{oc} of OSCs

and OHSCs. We have shown that if the difference between the electron and hole mobilities is small, the V_{oc} derived here does not depend on the charge carrier mobilities significantly. According to our model, the V_{oc} of a DA material depends on two terms; the first depends on the energetics and the second on the electron and hole mobility ratio. This may be expected to be useful in predicting the PCE of OSCs and OHSCs prior to their fabrications from a combination of DA materials.

References

- Rita Narayan M, Singh J (2013) Study of the Mechanism and Rate of Exciton Dissociation at the Donor-Acceptor Interface in Bulk-Heterojunction Organic Solar Cells. *Journal of Applied Physics* 114: 735101-735107.
- Omping D, Singh J (2015) Diffusion length and Langevin recombination of singlet and triplet excitons in organic heterojunction solar cells. *Chemphyschem* 16: 1281-1285. [Crossref]
- Narayan MR, Singh J (2012) Roles of Binding Energy and Diffusion Length of Singlet and Triplet Excitons In Organic Heterojunction Solar Cells. *Physica Status Solidi* 9: 2386-2389.
- Narayan MR, Singh J (2015) Excitonic and Photonic Processes In Materials, Singapore, Springer.
- Würfel U, Neher D, Spies A, Albrecht S (2015) Impact of charge transport on current-voltage characteristics and power-conversion efficiency of organic solar cells. *Nat Commun* 6: 6951. [Crossref]
- Deibel C, Wagenpahl A, Dyakonov V (2008) Influence of Charge Carrier Mobility on the Performance of Organic Solar Cells. *Physica Status Solidi (Rrl)- Rapid Research Letters* 2: 175-177.
- Gregg BA (2003) Excitonic Solar Cells. *The Journal of Physical Chemistry B* 107: 4688-4698.
- Brabec CJ, Cravino A, Meissner D, Sariciftci NS, Fromherz T, et al. (2001) Origin of the Open Circuit Voltage of Plastic Solar Cells. *Advanced Functional Materials* 11: 374-380.
- Scharber MC, Mühlbacher D, Koppe M, Denk P, Waldauf C, et al. (2006) Design Rules For Donors In Bulk-Heterojunction Solar Cells-Towards 10% Energy-Conversion Efficiency. *Advanced Materials* 18: 789-794.
- Ye L, Zhang S, Qian D, Wang Q, Hou J (2013) Application of Bis-Pcbm In Polymer Solar Cells With Improved Voltage. *The Journal of Physical Chemistry C* 117: 25360-25366
- Cravino A (2007) Origin of the Open Circuit Voltage Of Donor-Acceptor Solar Cells: Do Polaronic Energy Levels Play A Role? *Applied Physics Letters* 91: 243502.
- Koster LJA, Van Strien WJ, Beek WJE, Blom PWM (2007) Device Operation Of Conjugated Polymer/Zinc Oxide Bulk Heterojunction Solar Cells. *Advanced Functional Materials* 17: 1297-1302.
- Dastoor PC, Meneill CR, Frohne H, Foster CJ, Dean B, et al. (2007) Understanding And Improving Solid-State Polymer/C60-Fullerene Bulk-Heterojunction Solar Cells Using Ternary Porphyrin Blends. *The Journal of Physical Chemistry C* 111: 15415-15426.
- Liu Y, Zhao J, Li Z, Mu C, Ma W, et al. (2014) Aggregation and morphology control enables multiple cases of high-efficiency polymer solar cells. *Scharber* [Crossref]
- Proctor CM, Love JA, Nguyen TQ (2014) Mobility guidelines for high fill factor solution-processed small molecule solar cells. *Adv Mater* 26: 5957-5961. [Crossref]
- Servaites JD, Ratner MA, Marks TJ (2009) Practical Efficiency Limits In Organic Photovoltaic Cells: Functional Dependence of Fill Factor And External Quantum Efficiency. *Applied Physics Letters* 95: 163302.
- Nelson J (2003) *The Physics of Solar Cells*, London, Imperial College Press
- Wagenpahl A, Rauh D, Binder M, Deibel C, Dyakonov V (2010) S-Shaped Current-Voltage Characteristics Of Organic Solar Devices. *Physical Review B* 82: 115306.
- Mandoc MM, Koster LJA, Blom PWM (2007) Optimum Charge Carrier Mobility in Organic Solar Cells. *Applied Physics Letters* 90: 133504.
- Ebenhoch B, Thomson SAJ, Genevicius K, Juška G, Samuel IDW (2015) Charge Carrier Mobility of the Organic Photovoltaic Materials Ptb7 And Pc71bm And Its Influence On Device Performance. *Organic Electronics* 22: 62-68.
- Philippa B, Stolterfoht M, Burn PL, JuÅıka G, Meredith P, et al. (2014) The impact

- of hot charge carrier mobility on photocurrent losses in polymer-based solar cells. *Sci Rep* 4: 5695. [[Crossref](#)]
22. Albrecht S, Vandewal K, Tumbleston JR, Fischer FS, Douglas JD, et al. (2014) On the efficiency of charge transfer state splitting in polymer: fullerene solar cells. *Adv Mater* 26: 2533-2539. [[Crossref](#)]
23. Koster LJA, Mihailetschi VD, Blom PWM (2006) Bimolecular Recombination in Polymer/Fullerene Bulk Heterojunction Solar Cells. *Applied Physics Letters* 88: 052104.
24. Liu CY, Holman ZC, Kortshagen UR (2009) Hybrid solar cells from P3HT and silicon nanocrystals. *Nano Lett* 9: 449-452. [[Crossref](#)]
25. Liu XD, Xu Z, Zhang FJ, Zhao SL, Zhang TH, et al. (2010) Influence of Small-Molecule Material on Performance of Polymer Solar Cells Based On MeH-Ppv: Pcbm Blend. *Chinese Physics B* 19: 118601.
26. Sun Y, Takacs CJ, Cowan SR, Seo JH, Gong X, et al. (2011) Efficient, air-stable bulk heterojunction polymer solar cells using MoO(x) as the anode interfacial layer. *Adv Mater* 23: 2226-2230. [[Crossref](#)]