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Charge-controlled switchable CO2 capture on boron nitride nanomaterials

Abstract

Increasing concerns about the atmospheric CO2 concentration and its impact on the environment are motivating researchers to discover new materials and technologies for efficient CO2 capture and conversion. Here, we report a study of the adsorption of CO2, CH4, and H2 on boron nitride (BN) nanosheets and nanotubes (NTs) with different charge states. The results show that the process of CO2 capture/release can be simply controlled by switching on/off the charges carried by BN nanomaterials. CO2 molecules form weak interactions with uncharged BN nanomaterials and are weakly adsorbed. When extra electrons are introduced to these nanomaterials (i.e., when they are negatively charged), CO2 molecules become tightly bound and strongly adsorbed. Once the electrons are removed, CO2 molecules spontaneously desorb from BN absorbents. In addition, these negatively charged BN nanosorbents show high selectivity for separating CO2 from its mixtures with CH4 and/or H2. Our study demonstrates that BN nanomaterials are excellent absorbents for controllable, highly selective, and reversible capture and release of CO2. In addition, the charge density applied in this study is of the order of 1013 cm-2 of BN nanomaterials and can be easily realized experimentally. 2013 American Chemical Society.

Keywords

switchable, co2, capture, charge, boron, controlled, nitride, nanomaterials

Disciplines

Engineering | Physical Sciences and Mathematics

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Charge-Controlled Switchable CO₂ Capture on Boron Nitride Nanomaterials

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Supporting Information Placeholder

ABSTRACT: Increasing concerns about the atmospheric CO_2 concentration and its impact on the environment are motivating researchers to discover new materials and technologies for efficient CO_2 capture and conversion. Here, we report a study of the adsorption of CO_2 , CH_4 and H_2 on boron nitride (BN) nanosheets and nanotubes (NTs) with different charge states. The results show that the process of CO_2 capture/release can be simply controlled by switching on/off the charges carried by BN nanomaterials. CO_2 molecules form weak interactions with uncharged BN nanomaterials and are weakly adsorbed. When extra electrons are introduced to these nanomaterials (i.e. when they are negatively charged), CO_2 molecules become tightly bound and strongly adsorbed. Once the electrons are removed, CO_2 molecules spontaneously desorb from BN absorbents. In addition, these negatively charged BN nanomaterials are excellent absorbents for controllable, highly selective, and reversible capture and release of CO_2 . In addition, the charge density applied in this study is of the order of 10^{13} cm⁻² of BN nanomaterials, and can be easily realized experimentally.

INTRODUCTION

Climate change due to greenhouse emissions has become a significant global challenge and methods to efficiently and effectively separate, capture, store and convert greenhouse gases, especially CO₂ have attracted increasing attention.¹⁻³ CO₂ is a dominant contributor to the greenhouse effect which leads to climate change and global warming.^{4,5} The current industrial process for CO₂ capture is based on the use of aqueous amine solutions or chilled ammonia, which suffers from relatively low efficiency and issues such as equipment corrosion, solvent loss, and toxicity.⁶ To overcome such disadvantages, solid materials such as metal-organic frameworks (MOFs), carbon and non-carbon nanotubes have been proposed as attractive adsorbents for CO₂ capture⁶⁻⁹. However a regeneration step in the process of CO₂ adsorption/desorption is required generally requiring high temperature to release CO₂. In addition, high CO₂ selectivity means that regeneration is difficult due to the large adsorption energy. Therefore, a major challenge for CO₂ capture is to search for a highperformance and low-cost material with high selectivity and easy regeneration.

In recent years, boron nitride (BN) nanomaterials, such as born nitride nanotubes (BNNTs) and nano-sheets, have attracted increasing attention due to their unique properties.¹⁰⁻²⁵ They have excellent mechanical properties, high thermal conductivity, and high resistance to oxidation.¹⁰⁻¹⁵ These properties make them very promising in a variety of potential applications such as optoelectronic nanodevices, spintronics, light emission, photocatalysts, thermal rectifiers, functional composites, etc.¹⁶⁻²¹ One potential application is for CO₂ capture and storage.²⁶ However, the interactions between neutral BN nanosheets, BNNTs and CO₂ are very weak due to the vacant p-like orbitals of the boron atoms in these materials. In such electron-deficient nanomaterials formation of strong interactions with CO₂ are not favourable because CO₂ is a Lewis acid and it prefers to accept, rather than donate, electrons during reaction.

Here, for the first time we demonstrate that by modifying the charge state of the BN nanomaterials, adsorption/desorption of CO_2 on BN nanosheets and nanotubes can be controlled and reversed. In contrast to previous reports, the CO_2 capture/release occurs spontaneously without any barriers once the charge is injected into, or dismissed from, BN nanostructures. This is the first report of that BN nanomaterials can effectively capture/release CO₂. We also demonstrate the high-selectivity of charged BN nanomaterials in separation of CO₂ from gas mixtures such as CO₂/CH₄ and CO₂/H₂. Here we note that the modification of the charge state of BN nanomaterials can be easily realized experimentally using electrochemical methods, electrospray, electron beam or by gate voltage control.²⁷⁻²⁹

METHODS

The first-principles density-functional theory plus dispersion (DFT-D) calculations were carried out using the DMol3 module in Materials Studio.^{30,31} The BN sheet and BNNTs are fully optimized in the given symmetry using the generalized gradient approximation³² treated by the Perdew-Burke-Ernzerhof exchange-correlation potential with long range dispersion correction via Grimme's scheme.³³ An all electron double numerical atomic orbital augmented by *d*-polarization functions (DNP) is used as basis set. The method has been used to successfully determine interactions of some gases and BN nanomaterials.³⁴ The self-consistent field (SCF) procedure is used with a convergence threshold of 10^{-6} au on energy and electron density. The direct inversion of the iterative subspace technique developed by Pulay is used with a subspace size 6 to speed up SCF convergence on these systems.³⁵ Starting with all possible configurations of the gases adsorbed on the BN nanomaterials, geometry optimizations were performed with a convergence threshold of 0.002 au/Å on the gradient, 0.005 Å on displacements, and 10^{-5} au on the energy. The real-space global cutoff radius is set to be 4.10 Å. For the BN sheet, unit cells range from 5×5 to 10×10 , chosen with a 15 Å vacuum between sheets in order to avoid interactions between periodic images, and the Brillouin zone is sampled by $6 \times 6 \times 1$ kpoints using the Monkhorst-Pack scheme. For BNNTs with index from (5, 5) to (10, 10), tetragonal supercells with dimension $30 \times 30 \times c$ Å³ are used, where c is 10.068 ~ 10.072 Å and optimized to minimize the energy of the nanotube. The values of lengths of 30 Å were large enough to avoid interactions between periodic images. The Brillouin zones are sampled by $1 \times 1 \times 6$ k-points using the Monkhorst-Pack scheme for six BNNTs.

The adsorption energies of CO_2 , CH_4 and H_2 on BN sheets and BNNTs are calculated from Eq. 1.

$$E_{ads} = (E_{BN} + E_{gas}) - E_{BN-gas}$$
(1)

where E_{BN-gas} is the total energy of the BN nanomaterials with the adsorbed gas, E_{BN} is the energy of isolated BN sheets and BNNTs, and E_{gas} is the energy of an isolated gas molecule, such as CO₂, CH₄ and H₂. In this paper, the adsorption energy of each configuration is calculated within the same state. The electron distribution and transfer mechanism are determined using the Mulliken method.³⁶

The charge densities of CO_2 adsorbed on BN sheets and BNNTs are calculated from Eq. 2.

$$\rho = \frac{Q}{s} \tag{2}$$

where ρ (10¹³ cm⁻²) is the electron density of BN sheets and BNNTs, Q is the total charge per unit cell and S is the surface area of the BN sheets and BNNTs. In addition, the surface area can be calculated from Eq. 3 and Eq. 4.

$$S_{(BN sheet)} = \frac{\sqrt{3}}{2}a^2$$
(3)

 $S_{(BNNT)} = 2\pi rc$

(4)

where a is the side length of the BN sheets and with values from $12.580 \sim 25.160$ Å for BN sheets with the unit cells from 5×5 to 10×10 , r is the radius of the BNNTs, and c is the length of the BNNTs, which is $10.068 \sim 10.072$ Å for BNNTs with index from (5, 5) to (10, 10).

RESULTS AND DISCUSSIONS

First, we performed calculations of CO₂ adsorption on an uncharged BN sheet comprised of a 5×5 unit cell with periodic boundaries and a BNNT (5,5). The calculations were carried out using first-principles density-functional theory³² with long range dispersion correction (DFT-D).³³ The computational results show that CO₂ molecules can only form weak interactions with these BN nanomaterials. In order to enhance the ability of CO₂ to be captured on these BN nanomaterials, the charge distributions and electron densities of the BN sheet and BNNT in different charge states were analyzed. Analysis of the results suggests that adding/removing charge to/from the BN nanomaterials allows the control of CO₂ capture and release on/from these BN nanomaterials. In order to prove the above hypotheses, we studied the adsorption/release of CO₂ on BN sheets and BNNTs carrying different charges. Finally, the separation of CO₂ from gas mixtures of CO₂/CH₄ and CO₂/H₂ using negatively charged BN nanostructures is addressed by comparing the adsorption of CH₄ and H₂ on these nanostructures with that of CO_2 .

CO₂ Adsorption on Uncharged BN Sheets and BNNTs



Ead $= 4.46$ kcal/mol	Ead $= 3.79$ kcal/mol
(a)	(b)

Figure 1. Top and side views of physisorbed CO_2 on a (a) 5×5 BN sheet and (b) BNNT (5, 5) in their uncharged states. Atom color code: blue, nitrogen; pink, boron; gray, carbon; red, oxygen.

The minimum energy configurations of CO₂ absorbed on the uncharged 5×5 BN sheet and BNNT (5, 5) are shown in Figure 1. The configuration of the CO_2 and the neutral BN sheet suggests physisorption and the distance between the carbon atom of CO₂ and boron atom of the BN sheet is 3.196 Å (Figure 1 (a)). On physisorption, the linear CO_2 molecule is parallel to the BN sheet, and the absorbed CO₂ molecule shows little structural change compared to a free CO₂ molecule with the O-C-O angle of the absorbed CO_2 being 179.6 °. Due to the weak interactions, the charge transfer between the BN sheet and the absorbed CO_2 molecule is negligible with a value of 0.008e. These results indicate that the adsorption of CO2 on the neutral BN sheet is very weak and the adsorption energy is only 4.46 kcal/mol, which is consistent with the recent report.²⁶ The weak interaction is mainly attributed to the van der Waals forces between CO₂ molecule and the BN

sorbent. Similarly to the case of the BN sheet, the adsorption of CO₂ on neutral BNNT (5, 5) is also a physical process and the adsorption energy is calculated to be 3.79 kcal/mol. The physical absorption and weak interaction are also attributed to the van der Waals interaction between CO₂ and the BNNT. The configuration of the CO₂ adsorbed on BNNT (5, 5) is shown in Figure 1 (b), and the C...O and B...O bond distances are approximately 3.0-3.4 Å. The linear CO₂ molecule attached to the BNNT is parallel to the B-N bond. The O-C-O angle is very close to 180.0°, which means that the geometry of the physically adsorbed CO_2 is similar to that of free CO_2 . The charge transfer between the BNNT and CO₂ is negligible with a value of only 0.002e⁻. We also investigated CO₂ adsorption on BN sheets comprised of different sized unit cells (from 6×6 to 10×10) and BNNTs with different diameters (from (6, 6) to (10, 10)) and found that CO₂ can only form weak interactions with these BN nanomaterials when they are in their neutral states.

Charge Distributions and Electron Densities of Charged BN Sheets and BNNTs

In order to understand what the influence of changing the charge of the BN nanomaterials might have on the interactions between the BN nanomaterials and CO₂, the electron density distributions of the frontier orbitals (i.e. the lowest unoccupied molecular orbital (LUMO)) for the 5×5 BN sheet and BNNT (5, 5) are considered in Figure 2 (a) and (d), respectively. Figure 2 (a) and (d) clearly show that the LUMOs of the BN sheet and BNNT are predominantly distributed on the boron atoms, which suggest that when an extra electron is introduced to the BN nanomaterials, the electron will fill the p-like orbitals of boron atoms of the BN sheet and the BNNT, which is confirmed by comparison of the differences in electron density distributions of different charge states of the BN nanomaterials. Figure 2 shows the differences of electron density distributions between the (b) 0e and 1e⁻, and (c) 0e and 2e⁻ charged BN sheets, as well as the differences of electron density distributions between (e) 0e and 1e, and (f) 0e and 2e charged BNNTs (5,5). We see from Figure 2 that, when extra electrons are introduced to the BN nanomaterials, the electrons spread across the boron atoms of the BN materials, and then those negatively charged BN adsorbents become more likely to denote electrons to CO₂ (Lewis acid) than their neutral state analogues. The Mulliken charge analysis³⁶ also supports the above statements. Detailed information on the Mulliken atomic charges of the BN nanomaterial with different charges is listed in Table S1 of the supporting information. For the neutral, one electron (1e⁻) and two electron (2e⁻) charged 5×5 BN sheets, the Mulliken atomic charges of the nitrogen atoms are -0.428e, -0.422e and -0.415e, respectively. This means that the charge distributions on the nitrogen atoms in BN sheet are not greatly affected by the change in the total charge of the BN sheet. However, a notable change in boron atomic charge distribution is observed in the negatively charged BN sheet, with the atomic charges of boron atoms in 0e, 1e⁻ and 2e⁻ BN sheet being +0.428e, +0.382e and +0.335e, respectively. The charge populations of boron atoms and nitrogen atoms in BNNT (5, 5) with 0e, 1e⁻ and 2e⁻ states are similar to those in the case of the BN sheet. In summary, from the above frontier orbital analysis, electron density distributions and Mulliken atomic charges analysis³⁶, it is noted that when extra electrons are added to the BN materials, they are distributed on the boron atoms (in preference to the nitrogen atoms). As previously mentioned CO_2 is a Lewis acid and prefers to gain electrons during reaction. The negatively charged BN can donate electrons to CO_2 and their strong interaction is expected. Once the negative charges are released from the BN adsorbents, the negative charge distribution (as shown in Figure 2(b), 2(c), 2(e) and 2(f)) on the boron atoms will vanish simultaneously, and the strong interactions between CO_2 and BN sorbents will disappear. This suggests that switching on/off the charge states of BN nanosheets and nanotubes can control their ability to capture/regenerate CO_2 .



Figure 2. The LUMO of (a) 5×5 BN sheet and (d) BNNT (5, 5). These orbitals are drawn with an isosurface value of $0.03e/Å^3$. The colors of the orbitals show the wavefunction (green - positive, red - negative). The differences of electron density distributions (b) between 0e and 1e⁻ as well as (c) between 0e and 2e⁻ charge carrying states of 5×5 BN sheets. The differences of electron density distributions (e) between 0e and 1e⁻ as well as (f) between 0e and 2e⁻ charge carrying states of BNNTs (5, 5). These densities are drawn with an isosurface value of $0.005e/Å^3$. Atom color code: blue, nitrogen; pink, boron.

CO₂ Adsorption on 1e⁻ and 2e⁻ Charged BN Sheets and BNNTs



$E_{ad} = 18.66 \text{ kcal/mol}$	$E_{ad} = 15.87 \text{ kcal/mol}$
(a)	(b)
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Figure 3. Top and side views of chemisorbed CO_2 on 1e⁻ charge carrying state of a (a) 5×5 BN sheet and (b) BNNT (5, 5). Atom color code: blue, nitrogen; pink, boron; gray, carbon; red oxygen.

In order to prove the above hypotheses, we performed calculations of the adsorption/dissociation of CO_2 on negatively charged BN sheets and BNNTs. After addition of one electron to the 5 × 5 BN sheet, the CO_2 molecule strongly interacts with the BN sheet with an adsorption energy of 18.66 kcal/mol (or 78.37 kJ/mol). Ideally the adsorption energies of CO_2 on high-performance adsorbents should be in a range of 40 - 80kJ/mol.³ According to this criterion, the BN sheet with one negative charge on each 5×5 unit cell renders a good sorbent for CO₂ capture. The high adsorption energy indicates strong chemisorption of CO₂ on the negatively charged BN sheet. In the chemisorbed configuration (Figure 3 (a)), the CO₂ molecule structure is drastically distorted, and the C atom is bonded with one boron atom in the BN sheet. Compared with the configuration of CO₂ physically adsorbed on the neutral form of the BN sheet (Figure 1 (a)), the C-B distance shortened from 3.196 Å to 2.072 Å; the O-C-O angle bent from around 180° to 139.1°; and the two double C=O bonds are elongated from 1.176 Å to 1.236 Å. A Mulliken charge population analysis shows that there is 0.570 e⁻ charge transfer from the BN sheet to the CO₂ molecule.

The variation of thermodynamic properties, such as change in Gibbs free energy (Δ G, kcal/mol), enthalpy (Δ H, kcal/mol) and entropy (Δ S, cal/mol K) with temperature (K) has been calculated (Figure 4) in order to study the entropic and temperature effects on CO₂ adsorption on BN sheet (1e[°]). Figure 4 shows that Δ S decreases with temperature increasing from 25 to 200 K, and Δ S is almost constant when the temperature is above 200 K. The value of Δ H is almost constant over the whole temperatures range (i.e. from 25 to 1000 K). This results in Δ G linearly increasing with an increase in temperature. Moreover, Δ G is negative in the temperature range of approximately 25 to 500K, which indicates the adsorption of CO₂ on 1e[°] charged BN sheet to form a chemisorbed configuration is a spontaneous process when the temperature is below 500K.



Figure 4. Variation of thermodynamic properties with temperatures (K) when an isolated CO_2 molecule is chemisorbed by a 1e⁻ charged BN sheet. Squares, triangles and circles correspond to the change in Gibbs free energy (kcal/mol), change in entropy (cal/mol K) and change in enthalpy (kcal/mol), respectively.

Similarly, the CO_2 molecule can also be chemically adsorbed on the surface of negatively charged BNNTs (Figure 3 (b)). The adsorption energy is calculated to be 15.87 kcal/mol (66.97 kJ/mol) based on the DFT-D level of calculations. The high adsorption energy also indicates the excellent potential of BNNTs for adsorbing CO_2 . A Mulliken charge population analysis shows there is 0.449 e⁻ charge transfer from BNNT to the CO_2 molecule, which also supports the strong interactions between them. In this case (Figure 3 (b)), the CO_2 molecule is distorted and one double-bond is broken. The O-C-O bond angle is 119.3°, and the broken C-O bond is significantly elongated to 1.306 Å on the top of the nanotube. The B sites are also considerably pulled out of the tube, and B-N bonds are increased by 0.20 Å. The B-C and B-O distances are 1.903 Å and 1.621 Å, respectively.

The higher adsorption energy and distortion of configurations reflect the stronger interactions between CO_2 and the negatively charged BN sheet/NT than with the neutral forms. The enhanced interactions can be explained from the frontier orbitals analysis, electron density distributions and Mulliken charge populations of the BN sheet/NT. It can be seen that, with extra electrons, boron atoms in the BN nanomaterials become less positively charged than in the neutral nanomaterials. These boron atoms are therefore likely to donate electrons to the CO_2 (Lewis acid) enabling the formation of stronger bonds.



Figure 5. In (a) and (b) energy changes for reactions involving CO_2 adsorption on a 5 × 5 BN sheet are shown. In (a) the change is from physisorption to chemisorption with the BN sheet in a 1e charge state and in (b) the change is from chemisorption to physisorption with the BN sheet in a 0e charge state. In (c) and (d) energy changes are shown for CO_2 adsorption on BNNT (5, 5) in a (c) 1e charge state and (d) neutral charge state. Atom color code: blue, nitrogen; pink, boron; gray, carbon; red, oxygen.

In Figure 5 (a) and (c) we start with the minimum energy configuration of the neutral BN sheet and BNNT with physisorbed CO_2 . An electron is then added to the BN sheet or NT and we consider the change in energy as it relaxes to the chemisorbed state. In Figure 5 (b) and (d) we start with the minimum energy configuration of the negatively charged BN sheet and BNNT with chemisorbed CO_2 . An electron is re-

moved and then the system is allowed to relax, forming physisorbed CO_2 . When an extra electron is introduced to the BN sheet and BNNT, the interactions between CO₂ and BN nanomaterials drastically increase compared with those with the neutral BN sheet and BNNT, and CO2 molecules are chemisorbed on the BN adsorbents (Figure 5 (a) and (c)). The processes are exothermic by 13.87 and 11.91 kcal/mol for BN sheet and BNNT, respectively, with no barrier. On the other hand, after a negative charge is removed from the systems the adsorption of CO₂ on BN adsorbents spontaneously changes from chemisorption into physisorption without any reaction barrier and those processes are exothermic with values of 49.57 and 69.21 kcal/mol for BN sheet and BNNT, respectively. Again the reactions have no barriers to these changes after the electron is removed (Figure 5 (b) and (d)). Here it should noted that the energy for triggering the adsorpbe tion/desorption totally relies on the energy used for charging/discharging of the BN absorbents. Once the charge states were switched on/off, both adsorption and desorption processes of CO₂ on BN adsorbents seem to be spontaneous. In summary, our calculations indicate the BN sheet and BNNT could be new adsorbent materials for controlled capture and release of CO₂ because they can be easily charged and discharged by electrochemical methods, electrospray, electron beam or gate voltage control.

In order to know whether CO₂ can absorb on negatively charged BN sheets with larger unit cells and BNNTs with bigger diameters, we investigate the CO₂ adsorption on different sized BN sheets and BNNTs. Calculations were carried out on BN sheets with unit cells from 6×6 to 10×10 and BNNTs with indices from (6, 6) to (10, 10). The BN sheet lattices are 12.6 Å ~ 25.2 Å and BNNT diameters are 6.8 Å ~ 14.0 Å. Table S2 and S3 list the important structural and energetic properties of CO₂ adsorption on the BN absorbents with different sizes and structures, including bond distances, bond angles, and electron transfer (e) from BN sheet/NTs to CO₂ as well as their adsorption energies. Figure 6 shows the adsorption energies of CO₂ on different sized BN sheets and different diameter BNNTs as functions of their charge densities. For neutral BN sheets/NTs, CO2 molecules weakly interact with them by physisorption and the adsorption energies are around 4 kcal/mol for all BN adsorbents. The charge transfer between CO₂ and BN sheets/NTs are negligible with the values around 0.005 e. However, for one electron charged BN sheets with unit cells from 5×5 to 8×8 and BNNTs with index from (5, 5) to (7, 7), CO₂ molecules and BN adsorbents chemisorb and the adsorption energies are 18.66, 13.94, 10.35 and 7.46 kcal/mol for BN sheet with unit cell 5 \times 5, 6 \times 6, 7 \times 7 and 8 \times 8, and 15.87, 10.74 and 6.42 kcal/mol for BNNTs index (5, 5), (6, 6) and (7, 7), respectively. We can see that with an increase of the BN sheet unit cell from 5×5 to 8×8 and BN diameters from (5, 5) to (7, 7), the adsorption energies, the charge transfer and the interactions between CO2 and BN sheets/BNT decrease gradually. When the unit cell of BN sheet is 9×9 or 10 \times 10, the interactions between CO₂ and the BN sheet are very weak. Similarly when the diameter of the BNNTs is larger than 11.0 Å (with nanotube index (8, 8)), CO₂ molecules are physically adsorbed on the singly charged BNNTs with an adsorption energy of 3.90 kcal/mol, which is very similar to the adsorption of CO_2 on the neutral state BNNT (8, 8). This can be understood from the viewpoint of charge states. When the size of the BN sheets and BNNTs increases, the positive

charge of the boron atoms is increased, resulting in weaker interactions between boron atoms of BN sheets/NTs and CO₂. The electrons allocated to the boron atoms can significant influence the adsorption characteristics of CO2 on BN sheets/NTs, as demonstrated computationally by considering the results from different sized BN sheets/NTs charged with two electrons. Stronger interactions formed between CO₂ and all BN sheets/NTs when charged with 2e⁻ (Table S2 and S3 in supporting information list the detail information of the charge state). The adsorption energies between them are $69.73 \sim$ 21.13 kcal/mol with the lattice of BN sheet increasing from 12.6 Å to 25.2 Å. Absorption energies of 77.27 ~ 30.84 kcal/mol are obtained when the diameter of the BNNTs increases from 6.8 Å to 14.0 Å, respectively. It means that if two negative charges are located on each unit cell, BN nanomaterials with relative larger size can also efficiently capture CO₂.

In those calculations, there is one CO_2 molecule on each BN adsorbent, so the calculations with higher coverages of CO_2 have been carried out. The Table S4 and Figure S1 in the supporting information lists the structural properties and average adsorption energies (kcal/mol) with different coverages of CO_2 on a BN sheets with a 6×6 unit cell and a 2e⁻ charge state and up to 9 molecules of CO_2 (a coverage of 0.25 molecules per B atom). We find that the average adsorption energies of CO_2 on the BN sheet decrease from 53.31 to 4.72 kcal/mol with increase in CO_2 coverage from 0.03 to 0.25 CO_2 molecule per boron atom. The configurations of the CO_2 adsorbed on the BN sheet are also consistent with a change from chemisorption to physisorption with the increase in CO_2 coverage.



Figure 6. Adsorption energies (kcal/mol) of CO₂ on BN sheets with different unit cell sizes and different of BNNTs diameters are shown as functions of electron densities $(10^{13} \text{ cm}^{-2})$. The charge state (2e⁻) is the same for all the computational results shown in Figure 6.

The adsorption energies of CO_2 on BN sheets with different sized unit cells and BNNTs with different diameters, as a function of the electron densities are shown in Figure 6. We can see that the adsorption energies of CO_2 on BN adsorbents depend almost linearly on their electron densities. The adsorption energy of the BNNTs increases more rapidly than that of the BN sheets because of the curvature of BNNTs, which increases as the diameter becomes smaller. These results further demonstrate the feasibility of capturing/releasing CO_2 by using charged/uncharged BN nanostructures, even for large BN sheets/NTs. As mentioned previously, the charge densities used in this study, which are of the order 10^{13} cm⁻² of BN nanomaterial, can be easily achieved, e.g., by using electrochemical method, electrospray, electron beam or gate voltage control.²⁷⁻²⁹ For example, a charge density of 7.4×10^{13} cm⁻² has been obtained with the most common gate.²⁹





Figure 7. Adsorption energies (kcal/mol) of CO₂, CH₄ and H₂ adsorbed on 5×5 BN sheets and BNNTs (5, 5) with neutral, 1e⁻ and 2e⁻ charge states.

In order to demonstrate the high-selectivity of negatively charged BN nanostructures for CO₂ adsorption, the adsorption energies of CH₄ and H₂ on charged and discharged 5×5 BN sheets and BNNTs (5, 5) are calculated and compared with that of CO₂. The results indicate that the adsorption of CH₄ and H₂ on these BN nanomaterials under all conditions considered are physical rather than chemical. Table S5 displays the important parameters for these physical adsorptions, such as bond distances, adsorption energies and electron transfer from BN nanomaterials to CH₄ and H₂. The C...B distances for CH₄ adsorbed on BN nanostructures are between $3.4 \sim 3.6$ Å, and H...B distances for H₂ absorption are in a range of 2.4 ~ 2.7 Å. Figure 7 compares the adsorption energies of CO_2 , CH₄ and H₂ on BNNT (5, 5) and a BN sheet with neutral, 1e and 2e⁻ charged states. For the neutral 5×5 BN sheet and BNNT (5, 5), CH₄ and H₂ physically adsorb and the adsorption energies are around 2 ~ 3 kcal/mol for CH₄ and 1 ~ 2 kcal/mol for H₂, respectively. The interactions are similar to those of the CO₂ molecules with neutral BN sheets and BNNTs. When BN nanostructures are injected with one electron and two electrons, the adsorption of CH₄ and H₂ on them remains physical and the change of adsorption energies are very small in comparison to neutral BN materials. In contrast, the absorption of CO₂ on one-electron charged BN sheets and BNNTs are much stronger. Their chemisorption energies are 18.66 and 15.87 kcal/mol for the BN sheet and BNNT (5, 5),

respectively. With two extra electrons, the interactions between CO₂ and BN sheet and BNNT become very strong, and the adsorption energies are increased to 69.73 and 77.27 kcal/mol for BN sheet and BNNT (5,5), respectively. The charge transfer from 2e⁻ charged BN sheet and BNNT to the CO₂ molecules also increase significantly, with charge transfer values are 0.717 e and 0.983 e, respectively. For BN nanostructures charged with two electrons, CO₂ molecule binds more tightly with them than those charged with one electron. In the chemisorption process CO₂ molecules undergo significant structural distortion (Table S2 and S3), in which the C-B bond between C atom in CO2 and B atom in BN sheet and BNNT become shorter with a lengths 1.670 Å and 1.666 Å, respectively. The O-C-O angles of CO₂ molecules adsorbed on BN sheet and BNNT are bent to 119.5 ° and 124.9 °, respectively. The above comparisons demonstrate that negatively charged BN sheet and BNNTs have very high selectivity for capturing CO2 from CH4 and H2 mixtures, and these nanostructures can serve as good adsorbents for separation of these gases.

CO₂ and Gas Mixture Adsorption on BN Sheet with an Electric Field



Figure 8. (a) Variation of thermodynamic properties with temperatures (K) when an isolated CO_2 molecule adsorbs on a BN sheet to form chemisorption configuration with a vertical electric field of 0.05 a.u. Squares, triangles and circles correspond to the change in Gibbs free energy (kcal/mol), change in entropy

(cal/mol K) and change in enthalpy (kcal/mol), respectively. (b) The energy change due to a change from the physisorbed to chemisorbed configurations with an electric field of 0.05 a.u on 5×5 BN sheet. (c) The energy change due to a change from chemisorbed to physisorbed configurations without an electric field for CO₂ adsorption on 5×5 BN sheet.

In order to further demonstrate the reaction mechanism of charge controlled adsorption/dissociation of CO₂ on BN nanomaterials, the calculations of CO₂ capture and release, and gas mixture (CO_2 , CH_4 and H_2) separation controlled by switching on and off an electric field (0.05 a.u) of the systems have been carried out. Firstly we discuss CO₂ capture/release on the BN sheet by switching on and off the electric field. The optimized geometrical parameters and adsorption energies of the physisorption (a) and chemisorption (b) configurations for CO_2 adsorption on 5 × 5 BN sheet without and with the electric field are presented in Figure S2 in the supporting information. In its physisorption (Figure S2 (a)) and chemisorption (Figure S2 (b)) configurations, the distances between the boron atom in BN sheet and one oxygen atom in CO₂ molecule are 3.001 and 1.651 Å, and the their adsorption energies are calculated to be 2.41 and 19.26 kcal/mol, respectively, which suggests CO₂ strongly interacts with BN sheet in the presence of the electric field. Figure 8(a) shows the temperature dependence of thermodynamic properties, such as ΔG (kcal/mol), ΔH (kcal/mol) and ΔS (cal/mol K) of the adsorption reaction of CO₂ on BN sheet with the electric field. It clearly shows that the ΔS of the reaction decreases firstly and then reaches to a constant with temperature increasing from 200 to 1000 K. Thus ΔG of this reaction increases with the increase of temperature. Moreover, the negative ΔG below 600K indicates the adsorption reaction process is spontaneous when the temperature is below 600K.

Figure 8 (b) and (c) schematically show two reaction processes. In (b), we start with the minimum energy physisorption configuration that the system would adopt without an electric field (Figure S2 (a)), an electric field is applied and then the system is allowed to relax. In (c) we start with the minimum energy chemisorption configuration that the system would adopt when the electric field is on (Figure S2 (b)), the electric field is then switched off and the system is allowed to relax. When the electric field is applied to the system, the interactions between CO₂ and BN sheet drastically increase compared with the case without electric field, and CO₂ molecule is chemisorbed on the BN sheet (Figure 8 (b)), and the process is exothermic by 17.47 kcal/mol without reaction barrier. On the other hand, when the electric field is switched off, adsorption of CO₂ on BN sheet changes from chemisorption into physisorption without any reaction barrier and the process is exothermic by 75.12 kcal/mol (Figure 8 (c)). Similar to the reaction mechanism of charged controlled CO₂ capture/release on the BN materials, the processes of CO₂ capture/release in the presence/absence of an electric field (0.05 a.u.) are found to be also spontaneous. There is an energy cost to the process since it strongly depends on switching on/off an electric field. The CO₂ capture/release controlled by charging/discharging should be very similar to that controlled by switching on/off electric field and the energy cost mainly comes from adding/removing the charges to/from the adsorbents experimentally.

In addition, a gas mixture (one molecule of CO_2 , H_2 and CH_4) adsorption on BN sheet in the presence/absence of the

electric field has also been calculated to address the gas separation by applying electric field on the material. The physisorption and chemisorption configurations of gas mixture adsorption on BN sheet without and with the electric field, the thermodynamic properties of the reaction from isolated gases to chemisorption, and the reaction energies of switching on and off the electric field for physisorption and chemisorption configurations are demonstrated in Figure S3-S5 in the supporting information, respectively. The adsorption energies of the gas mixture on the BN sheet for the two configurations (Figure S3(a)) without and (Figure S3(b)) with electric field are 8.20 and 29.89 kcal/mol, respectively. From the two optimized configurations (Figure S3(a) and Figure S3(b)) of gas mixture adsorption on BN sheet, we can see that CO₂ molecule changes from far binding to BN sheet to tight binding with the B-O bond distances shortening from 2.995 to 1.651 Å. However, the distances of H₂ and CH₄ to BN sheet are almost unaffected by the switching on and off the electric field, which demonstrate that CO₂ can be efficiently separated from the gas mixture by applying the electric field. The energy profiles of reaction processes (Figure S5 in supporting information) of applying electric field to the physisorption configuration (Figure S5(a)) and switching off the electric field with the chemisorption configuration (Figure S5(b)) indicate these processes have no energy barrier. Moreover, ΔG for adsorption of the gas mixture on the BN sheet by applying electric field is negative around 300K (Figure S4), which indicates the chemisorption of CO_2 in the presence of H_2 and CH_4 with the electric field is spontaneous at room temperature.

CONCLUSIONS

In summary, the adsorption/desorption of CO₂ on charged and discharged BN sheets and BNNTs were theoretically investigated. The results show that CO₂ adsorption on BN nanostructures can be drastically enhanced by introducing electrons to the adsorbent. The absorbed CO₂ can be spontaneously released from BN nano-sorbents without any reaction barrier once the electrons are removed. The results suggest that negatively charged BN nanomaterials are excellent sorbents for CO₂ and they can be used to separate CO₂ from gas mixtures, in processes such as natural gas sweetening (CO₂/CH₄) and post-gasification (CO₂/H₂) capture. These results are also supported by the calculations on the adsorption/desorption of CO₂, and gas mixture separation on BN nanostructure in the presence/absence of an electric field. Our investigation demonstrates a versatile approach to CO₂ capture, regeneration and gas separation by charging/discharging the sorbents which can be easily realized.

SUPPORTING INFORMATION

Mulliken atomic charges of the BN nanomaterial with different charges and important structural properties as well as adsorption energies of the gases adsorption on the different size and different charge states of BN nanomaterials are listed in supporting information. This material is available free of charge via the Internet at http://pubs.acs.org.

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SYNOPSIS TOC

