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Authors

Altman, Alison B Pemmaraju, C Das Alayoglu, Selim <u>et al.</u>

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Chemical and Morphological Inhomogeneity of Aluminum Metal and Oxides from Soft X-ray Spectromicroscopy

Alison B. Altman,^{1,2} C. Das Pemmaraju,³ Selim Alayoglu,¹ John Arnold,^{1,2} Corwin H. Booth,¹ Augustin Braun,¹ Christopher E. Bunker,⁴ Alexandre Herve,¹ Stefan G. Minasian,^{1*} David Prendergast,^{3*} David K. Shuh,¹ Tolek Tyliszczak⁵

¹ Chemical Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States

² Department of Chemistry, University of California, Berkeley CA 94720, United States

³ Molecular Foundry, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States

⁴ Air Force Research Laboratory, Propulsion Directorate, Wright-Patterson Air Force Base, OH, 45433, United States

⁵ Advanced Light Source, Lawrence Berkeley National Laboratory, Berkeley, CA 94720, United States

Email: sgminasian@lbl.gov, dgprendergast@lbl.gov

ABSTRACT. Oxygen and aluminum K-edge X-ray absorption spectroscopy (XAS), imaging from a scanning transmission X-ray microscope (STXM), and first principles calculations were used to probe the composition and morphology of bulk aluminum metal, α - and γ -Al₂O₃, and several types of aluminum nanoparticles. The imaging results agreed with earlier transmission electron microscopy studies that showed a 2 to 5 nm thick layer of Al₂O₃ on all the Al surfaces. Spectral interpretations were guided by examination of the calculated transition energies, which agreed well with the spectroscopic measurements. Features observed in the experimental O and Al K-edge XAS were used to determine the chemical structure and phase of the Al₂O₃ on the aluminum surfaces. For unprotected 18 and 100 nm Al nanoparticles, this analysis revealed an oxide layer that was similar to γ -Al₂O₃ and comprised of both tetrahedral and octahedral Al coordination sites. For oleic-acid protected Al nanoparticles, only tetrahedral Al oxide coordination sites were observed. The results were correlated to trends in the reactivity of the different materials, which suggests that the structures of different Al₂O₃ layers have an important role in the accessibility of the underlying Al metal towards further oxidation. Combined, the Al K-edge XAS and STXM results provided detailed chemical information that was not obtained from powder X-ray diffraction or imaging from a transmission electron microscope.

INTRODUCTION

Efforts to utilize aluminum and other oxophilic metals require a precise grasp of surface structure and composition, as well as passivation techniques to control surface reactivity. Aluminum materials do not exhibit ideal, bulk-like surface terminations, which becomes an increasingly dominant characteristic of their chemistry as the surface-area to volume ratio increases. In the case of aluminum, the metal typically has a ~5 nm passivating layer of Al_2O_3 ,¹ such that the weight % of O in nanoparticles increases as particle size decreases. Al_2O_3 is commonly found in a number of different amorphous (*am*) and crystalline phases (*a* and γ , among others). However, relatively little is known about how the coexistence of multiple crystalline structures at an interface can result in heterogeneous oxide systems with unusual porosity, electronic structure, thermodynamic stability, and reactivity. Hence, unraveling the rates and mechanisms of Al metal reactivity with O_2 in fine powders, nanoparticles, thin films, and molecular clusters has been the subject of considerable experimental²⁻⁷ and theoretical effort.⁸⁻¹³ Understanding how changes in the phase of Al_2O_3 can impact reactivity has important consequences for scientific and technological applications for aluminum in ceramics, catalysis, coatings, separations,¹⁴⁻²⁰ and energetic materials.²¹⁻²⁵ To optimize the synthesis of well-defined nanoparticles and control their reactivities, characterization tools are needed that can probe the structure of heterogenous materials over multiple length scales and under real-world conditions.

X-ray absorption spectroscopy and scanning transmission X-ray microscopy (STXM-XAS) have emerged as powerful probes of nano- and microscale physical structure and chemical bonding for aluminum materials.^{26,27} Each also has unique advantages and limitations when compared with the electron-based analogs scanning transmission electron microscopy (STEM) and electron energy-loss spectroscopy (STEM-EELS).²⁸ For example, modern STEM capabilities regularly achieve 0.1 nm spatial resolution and monochromated EELS can provide as fine as 0.2 eV energy resolution. STXM can accommodate thicker samples (1 to 20 μ m depending on the X-ray probe energy, relative to 0.5 μ m for STEM when conducted with a fixed energy electron beam), but the spatial resolution is limited to about 25 nm by the focusing ability of the X-ray optic. Modern STXM instruments easily achieve 0.1 eV energy resolution under normal conditions (resolving power $E/\Delta E > 7500$), and XAS collected in the transmission mode can be obtained from the core-levels of many elements in the Periodic Table, including the light atom K-edges for B, C, N, O, F, Na, Mg, Al, and Si. Previous studies have shown that the spectroscopic accuracy of STXM-XAS^{29.32} provides bulk-like electronic structure information that can be interpreted quantitatively within a band structure or molecular orbital model.^{33.37} Additionally, X-ray based microscopes are ideally-suited for sensitive inorganic materials because compositional changes or other damaging effects that can occur on exposure to vacuum or to an electron beam are more easily controlled.^{38,39}

In this study, images from STXM were used to evaluate the chemical speciation and morphology for several reference materials as well as aluminum nanoparticles prepared under different synthetic conditions. The reference materials included a 100 nm aluminum metal foil, α -Al₂O₃, and γ -Al₂O₃. Aluminum nanoparticles (Al NPs) prepared with three different synthetic methodologies were also explored: core-shell Al NPs with 54 nm average diameters, prepared with a protective organic coating of oleic acid (Al–OA NPs);⁴⁰ bare Al NPs with 100 nm average diameters and no protective coating, prepared by the electrical explosion of wires (similar to ALEX[®]);^{39,41} bare Al NPs with 18 nm average diameters

and no protective coating, prepared by pulsed laser ablation of Al targets in an organic solvent.^{42,43} These analytes were specifically chosen because they provided a range of particle sizes, compositions, surface chemistries, and reactivities.^{44,47} For example, previous studies have shown that protecting agents including organic ligands^{40,48-54} and polymer⁵⁵⁻⁵⁷ or metal coatings^{44,58,59} enhanced reactivity properties and provided chemical control over structural attributes like particle size and monodispersity. The STXM-XAS results were consistent with earlier spectromicroscopy studies,^{26,55,60} and also revealed new differences in the surface and micron-scale speciation of Al nanoparticles.

RESULTS AND DISCUSSION

Scanning Transmission X-ray Microscopy. The 100 nm Al foil and 18 and 100 nm Al NPs were obtained from commercial sources and used as received (see Experimental). The protected, core-shell Al–OA NPs were prepared according to the literature procedure, and a 54 nm average size was determined by transmission electron microscopy (TEM) (Supporting Information).⁴⁰ Powder X-ray diffraction (PXRD) patterns were obtained from all the samples to evaluate the purity of analytes used in subsequent STXM measurements (Supporting Information). In each case, the PXRD patterns were consistent with the presence of predominantly face-centered cubic aluminum metal. For the 100 nm Al NPs, the PXRD measurements also provided evidence for a minor component of hydrated alumina, α -Al₂O₃•(H₂O)₃. No evidence for α -Al₂O₃ was observed in the PXRD patterns from the 18 nm Al NPs and 54 nm Al–OA NPs. However, weak reflections were observed that could not be definitively assigned, and may have been evidence of small amounts of Al₂O₃ present in other polymorphs (e.g., β -, θ -, κ -Al₂O₃).

The STXM at the Molecular Environmental Science beamline 11.0.2 of the ALS was used to collect images, elemental maps, and O and Al K-edge XAS spectra from a 100 nm Al foil and the three different types of Al NPs described above.⁶¹⁻⁶³ To minimize degradation of the air- and water-sensitive materials, samples were encapsulated between Si₃N₄ supports in an argon-filled glovebox. In addition, the STXM was filled with a partial atmosphere of helium during acquisition. The STXM was used to collect single-energy images and elemental contrast maps by raster-scanning the sample and collecting transmitted monochromatic light as a function of sample position. Figure 1 shows representative normal contrast images and elemental maps from micron-scale aggregates of nanometer-sized particles. Lighter regions in the elemental maps corresponded to greater concentration of the absorbing atom, and were obtained by subtraction of two images: one taken at an energy just below the edge and another taken at the absorption maximum (see "X-ray Absorption Spectroscopy" below). For example, elemental mapping at the Al K-edge was used to identify particles containing metallic aluminum by subtracting an image obtained at 1550 eV from another at 1560 eV. Regions containing Al₂O₃ were identified at the Al K-edge with energies of 530 and 540 eV.

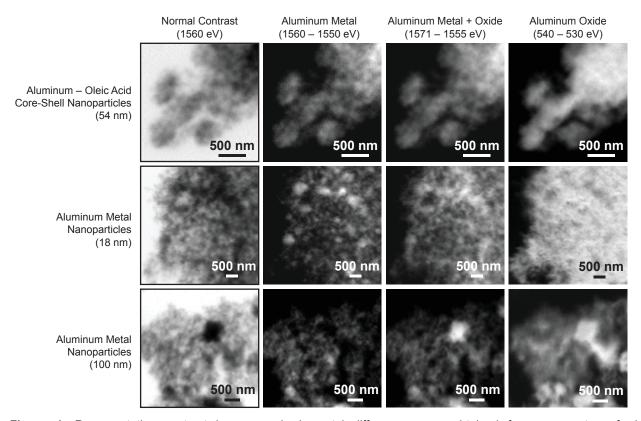


Figure 1. Representative contrast images and elemental difference maps obtained from aggregates of aluminum nanoparticles from which XAS were obtained. Difference maps were produced by obtaining images at the two energies indicated at the top of each column. Subtraction of a background image taken at an energy just below the absorption edge from another image taken at an energy on the absorption maximum generated the difference map, where lighter regions indicate greater concentration of the absorbing atom.

The images in Figure 1 show that all of the NPs had small diameters on the nanometer scale. In powdered form, all the materials formed tight micron-scale aggregates, as had been observed previously for both bare NPs and NPs protected by inorganic and organic surface layers.^{55,60,64} Although many details of the nanoscale structure exceeded the spatial-resolution capability of STXM (40 nm under the experimental conditions used), on a large scale the Al–OA NPs were monodispersed and had a uniform distribution of metallic Al and Al₂O₃. In contrast, the STXM images of the 18 and 100 nm Al NPs revealed inhomogeneities in the distribution of metallic aluminum and Al₂O₃, and showed that particle sizes ranged from the nano to the micrometer scale. These observations agreed well with earlier scanning electron microscopy (SEM) and TEM measurements on ALEX[®] (an ultra-fine aluminum powder produced by the electrical explosion of wires), which showed that while most particle diameters were on the order of 100 nm, particles as small as 20 nm and up to 500 nm in diameter were also observed.³⁹

The STXM images shown in Figure 1 did not provide a direct measure of the thickness of the passivating shells of aluminum oxides. However, relative amounts of metallic aluminum and aluminum oxide in the nanomaterials were evaluated from the images by considering the O and Al K-edge spectroscopic results (*vide infra*) and conducting Singular Value Decomposition (SVD) analysis.⁶⁵ The optical density of a target area in a given image was governed by $OD = \mu \times \rho \times t$, where μ was the mass absorption coefficient, ρ was the density, and t was the sample thickness or path length. In this study, the intensity of the Al K-edge spectrum (OD) at 1560 eV was proportional to the concentration ($\rho \times t$) of metallic

Al in the target area of a given particle. Similarly, the intensity of the O K-edge spectrum (OD) at 540 eV was proportional to the surface concentration of Al_2O_3 in the same target area. Using these concentrations, values for the weight percent of Al and Al_2O_3 were determined and are given in Table 1. To provide an accurate measurement, analyses encompassed a number of individual particles and were conducted on large target areas greater than 10 μ m² and 10,000 total pixels. The chemical composition of surface species was identified from the O and Al K-edge XAS results (*vide infra*). It is worth noting that a limitation of the SVD approach arose from the assumption that surfaces are comprised of a single phase of Al_2O_3 , when a more complex mixture of multiple oxide phases, hydroxides, and hydrated surface species may have been present.

Table 1. Characteristics of AI metal and nanoparticle materials determined from STXM–XAS. For nanoparticle-based samples, estimation of the oxide layer thickness (t_{oxide}) was based on the relative weight percents of each constituent, and assumed an idealized spherical morphology with a homogenous layer of γ -Al₂O₃ (ρ = 3.7 g cm⁻³). For the 100 nm Al foil, estimation of t_{oxide} was based on a homogenous layer of am-Al₂O₃ (ρ = 2.3 g cm⁻³).^a Complete details on the derivation of t_{oxide} are provided in the Supporting Information.

				Weight % of Constituents			
Cmpd	diameter (nm)	passivation	Source	OA	AI_2O_3	metallic Al	t _{oxide} (nm)
Al foil	100	air	Luxel	_	12 ± 1	88 ± 8	3.9 ± 0.4
AI NPs	18	air	US Nano	-	53 ± 5	47 ± 4	1.7 ± 0.2
AI NPs	100	air	US Nano	_	34 ± 3	66 ± 6	5.2 ± 0.5
AI–OA NPs	54	oleic acid	this work	13 ± 1	29 ± 2	58 ± 5	2.3 ± 0.4

^{*a*} Values for the density of *am*-Al₂O₃ vary widely.⁶⁶ Using ρ = 3.6 g cm⁻³ for *am*-Al₂O₃ provides *r*_{oxide} = 2.5 ± 0.3 nm for the 100 nm Al foil.

For the Al foil described above, this analysis showed that the foil had surface concentrations of $3.6 \pm 0.3 \ \mu g \ cm^{-2}$ oxygen and $56 \pm 4.0 \ \mu g \ cm^{-2}$ aluminum. If all the oxygen was associated with a layer of amorphous alumina (*am*-Al₂O₃, $\rho = 2.3 \ g/\ cm^3$),³⁸ then these values correspond to $88 \pm 8 \ wt$. % metallic aluminum and $12 \pm 1 \ wt$. % *am*-Al₂O₃. The thickness of the foil ($104 \pm 11 \ nm$) was calculated by dividing the surface concentration by density of aluminum metal. Likewise, the thicknesses of Al₂O₃ layers on the top and bottom of the Al foil were calculated at $3.9 \pm 0.4 \ nm$ each depending on the value for the density of *am*-Al₂O₃ used. These values are similar to specifications for the Al foil quoted by the manufacturer,⁶⁷ which include thicknesses of 100 and 5 nm and compositions of roughly 86 and 14 wt. % for Al metal and *am*-Al₂O₃, respectively.

Relative amounts of Al metal and oxide for the conventionally-synthesized nanoparticles were determined using SVD analysis in a similar fashion (Table 1). The 18 nm Al NPs had 53 ± 5 wt. % Al₂O₃ and 47 ± 4 wt. % metallic Al, and the 100 nm Al NPs had 34 ± 3 wt. % Al₂O₃ and 66 ± 6 wt. % metallic Al. As described in the introduction, these changes in composition were anticipated given the increased fraction of metallic Al in larger particles. The STXM derived values suggested that there is somewhat less metallic Al in the 100 nm nanoparticles used in this study (66 ± 6 wt. %) than was determined for a different sample of 100 nm ALEX[®] NPs using a volumetric analysis (86% metallic Al) and using energy dispersive spectroscopy (EDS) from TEM (90 ± 5 wt. % total Al and 10 ± 2 wt. % O, which corresponded to 79 ± 5 wt. %

metallic Al and 21 ± 4 wt. % Al₂O₃).^{41,68} As described above for the Al foil, these relative amounts of Al and Al₂O₃ were also used to calculate thicknesses for the Al₂O₃ passivation layer by assuming an ideal spherical morphology and monodisperse 18 and 100 nm particle diameters (Table 1 and Supporting Information). This analysis showed that the conventionallysynthesized 18 nm and 100 nm Al NPs had 1.7 ± 0.2 and 5.2 ± 0.5 nm thick Al₂O₃ layers, respectively.

For the Al–OA core-shell nanoparticles, proportions of 58 ± 5 , 29 ± 2 , 13 ± 1 wt. % were calculated for the metallic core, Al₂O₃, and carbonaceous organic layer, respectively. For comparison, previous ICP-MS and TGA studies have suggested that 40% of the sample mass was metallic aluminum, 25% was Al₂O₃, and 35% was organic.⁴⁰ In both cases, the amount of oxygen relative to carbon was larger than anticipated for oleic acid alone, and clearly indicates that oleic acid passivation did not completely inhibit growth of an aluminum oxide layer. This observation is consistent with earlier studies, which showed that passivation resulted in decomposition of the oleic acid and formation of an organic outer layer on the surface on an intermediate oxide layer.⁴⁰ A thickness of 2.3 ± 0.4 nm was estimated for the Al₂O₃ passivation layer in the Al–OA NPs, which is also in agreement with earlier work showing that organic or inorganic protective coatings limited growth of the Al₂O₃ layer to about 2 nm.^{44,55,69}

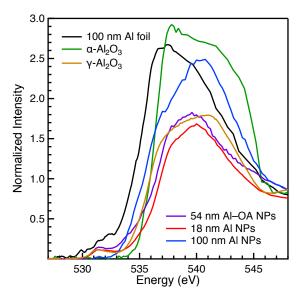


Figure 2. O K-edge XAS data for 54 nm Al–OA nanoparticles (purple), 18 nm Al NPs (red), 100 nm Al NPs (blue), and reference materials including a 100 nm Al metal foil (black), α -Al₂O₃ (green) and γ -Al₂O₃ (brown).

X-ray Absorption Spectroscopy. Figure 2 shows the background-subtracted and normalized oxygen K-edge XAS measured from the bare 18 and 100 nm Al NPs, core-shell Al–OA NPs, and reference materials α -Al₂O₃, γ -Al₂O₃, and a 100 nm Al foil. As previous studies have shown,^{11,30,38} differences in the O K-edge XAS can reflect changes in Al–O bonding found for different phases of aluminum oxides and hydroxides. The O K-edge spectra shown in Figure 2 were similar in many regards; however, some important differences were readily apparent. For example, the spectrum for γ -Al₂O₃ exhibited a small shoulder near 531 – 532 eV that was not observed for α -Al₂O₃. Previous work by Århammar and coworkers suggested that these small features are characteristic of O–O interactions that are present in some phases of

Al₂O₃ but not in crystalline α -Al₂O₃.³⁸ The spectra from the 18 and 100 nm Al NPs and Al–OA NPs each had some of the spectroscopic signatures observed from γ -Al₂O₃. None of the nanoparticle spectra resembled the O K-edge of the 100 nm Al foil (reflecting *am*-Al₂O₃), which was shifted such that the first inflection point was approximately 1 eV lower in energy than was observed for the other materials.

Al K-edge XAS was used to further elaborate on the compositional assignments determined using O K-edge XAS. Figure 3 shows representative background subtracted and normalized Al K-edge spectra obtained from the bare 18 and 100 nm Al NPs, core-shell Al–OA NPs, and reference materials α -Al₂O₃, γ -Al₂O₃, and the 100 nm Al foil. Energies and assignments for the spectral features are summarized in Table 2. Because previous attempts to model and interpret Al Kedge XAS data for the metal and oxide references have been limited,^{31,70,71} this discussion will begin by establishing the bonding descriptions needed to justify the spectral assignments for the nanoparticle materials given in Table 2.

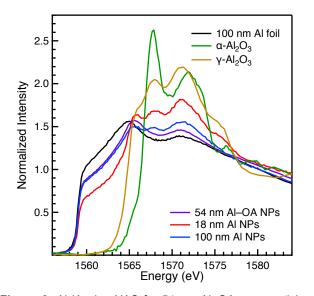


Figure 3. Al K-edge XAS for 54 nm Al–OA nanoparticles (purple), 18 nm Al NPs (red), 100 nm Al NPs (blue), and reference materials including a 100 nm Al metal foil (black), α -Al₂O₃ (green), and γ -Al₂O₃ (brown).

Table 2. AI K-edge feature energies and assignments.

0	0	0
Compound	Energy (eV)	Assignment
α -Al ₂ O ₃	1564.9ª	$1s \rightarrow 2s\text{-}\sigma^{*} \; (a_{1g})$
	1567.8	$1s \to 3p\text{-}\sigma/\pi^*~(t_{\text{1u}})$
	1571.9	$1s \rightarrow 3d\text{-}\sigma^{*} \; (t_{\text{2g}})$
γ -Al ₂ O ₃	1565.4	$1s \rightarrow 3p\text{-}\sigma/\pi^{*} \ (t_{2})$
	1568.0	$1s \to 3p\text{-}\sigma/\pi^{*} \; (t_{\text{\tiny 1u}})$
	1571.2	$1s \rightarrow 3d\text{-}\sigma^{*} \; (t_{\text{2g}})$
100 nm Al foil	1559.0	$1s \to 3p\text{-}\pi^{\text{b}}$
	1564.9	$1s \to 3p\text{-}\sigma^*$
	1570.8	$1s \to 3d\text{-}\sigma^*$
18 nm Al NPs	1559.0	$1s \to 3p\text{-}\pi^{\text{b}}$
	1565.9	$1s \rightarrow 3p\text{-}\sigma/\pi^{*} \ (t_{2})$
	1568.0	$1s \rightarrow 3p\text{-}\sigma/\pi^{**} \ (t_{1u})$
	1571.1	$1s \rightarrow 3d\text{-}\sigma^{*} \; (t_{\text{2g}})$
100 nm Al NPs	1559.0	$1s \rightarrow 3p - \pi^{b}$
	1565.6	$1s \to 3p\text{-}\sigma/\pi^*~(t_2)$
	1567.8	$1s \to 3p\text{-}\sigma/\pi^*~(t_{\text{1u}})$
	1571.1	$1s \rightarrow 3d\text{-}\sigma^{*} \; (t_{\text{2g}})$
54 nm Al–OA NPs	1559.0	$1s \rightarrow 3p - \pi^{b}$
	1565.6	$1s \rightarrow 3p\text{-}\sigma/\pi^{\star}~(t_2)$
	1570.9	$1s \to 3d\text{-}\sigma^*$

^a This transition is weak and dipole-forbidden, but becomes recognizable as a result of some distortion from idealized O_h symmetry which may facilitate some Al 3s and 3p mixing.

Bulk Aluminum Metal. Figure 4 shows the experimental Al K-edge XAS of a 100 nm Al foil and simulations provided by both FEFF 10.0.1^{72,73} and XCH DFT calculations. Both theoretical approaches provided good agreement with the experimental data in both the near-edge and extended energy regions. The onset of the Al K-edge for Al metal, defined by the first inflection point, was fixed at 1559.0 eV and was the calibration point for the other Al K-edge spectra described

below. At higher energies, the edge region was characterized by two very broad features at 1564.9 and 1570.8 eV. Although Al K-edge spectra of Al metal have been reported,⁷⁴⁻⁸¹ to the best of our knowledge no detailed interpretation of the spectral profile has been provided. The discussion below begins by constructing the band structure model for Al metal that was used to understand the Al K-edge XAS features.

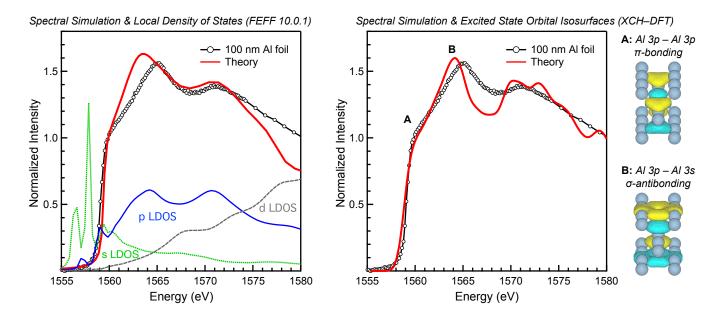


Figure 4. Al K-edge XAS of Al metal (black circles) compared with simulations from FEFF (left) and XCH-DFT (right) calculations. In the left plot, the total simulation (red) is included with the local density of states (LDOS) derived from the FEFF for final states associated with the Al s, p, or d orbitals (green, blue, and gray). In the right plot, the correspondence between transitions and isosurfaces is indicated by the letters A and B. Note that the FEFF calculation is shifted down by approximately 5 eV to better compare to the Al K edge experimental data, which is a typical absolute energy disagreement for FEFF. Similarly, the XCH-DFT simulation was shifted by +1560.65 eV.

From a chemistry perspective, molecular orbital theory can provide a useful starting point for understanding metallic bonding.^{82.84} Figure 5 illustrates this approach for the construction of molecular orbitals in aluminum metal, beginning with the hypothetical aluminum molecule Al₂. To a first approximation, interaction of the two 3s atomic orbitals results in formation of two molecular orbitals of σ -symmetry, which are split in energy as the bonding and antibonding combinations. Similarly, the six 3p orbitals form bonding and antibonding orbitals with both σ - and π -symmetries. Furthermore, because the 3s and 3p σ -bonding orbitals have the same symmetry and similar energies (11.3 and 5.9 eV, respectively),⁸⁴ they can interact via s-p mixing to form new orbitals with different energies. Castleman et al. have previously used vacuum ultraviolet photoelectron imaging to highlight the importance of s-p hybridization in small aluminum clusters formed in the gas phase.⁸⁵ Consequently, Al₂ has a ground-state electronic configuration of $(1\sigma_g^+)^2(1\sigma_u^+)^2(1\pi_u)^2$ where the partially occupied $1\pi_u$ orbitals are bonding, and the empty $2\sigma_g^+$ orbitals are best described as non-bonding resulting from s-p mixing. The orbitals described for Al₂ are split again when an infinite array of Al atoms are brought together in the metal, resulting in a band of tightly packed energy levels corresponding to both bonding and antibonding states.

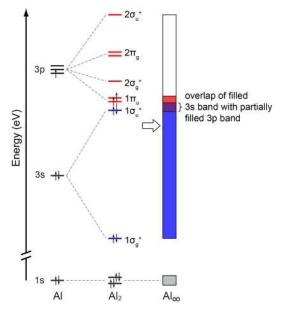


Figure 5. Formation of energy bands in aluminum metal. Orbitals for an isolated Al atom (left) form σ - and π -type interactions in a diatomic molecule (middle) which are split into bonding and antibonding counterparts. Bringing together a large number of atoms in an extended lattice results in a larger number of closely spaced energy levels.

This qualitative picture is consistent with the quantitative Al K-edge XAS simulations obtained using FEFF 10.0.1 and XCH-DFT, shown in Figure 4. As expected, the local density of states provided by the FEFF calculation showed that the occupied bands have predominantly 3s character, while at higher energies the unoccupied bands have primarily Al 3p character. In addition, XCH-DFT calculations were used to generate isosurface plots of the final state electronic orbitals that corresponded to (A) the most dominant transitions near the edge onset, and (B) transitions that were close to the energy of the first main feature at 1564.9 eV (Figure 4). Consistent with the qualitative interpretation provided above, the orbital isosurfaces showed that transitions near the edge onset at 1559 eV involved low-lying empty states with significant π -bonding character between adjacent atoms in the unit cell (Al 1s $\rightarrow \pi^{b}$). The calculations also showed an additional feature at 1563.6 eV (FEFF) and 1564.2 eV (XCH-DFT), which corresponded to the broad peak at 1564.9 eV in the experimental spectrum. The peak was reminiscent of a shape resonance: broad features associated with short-lived final states, given the increasing probability of decay to the continuum at energies above the Al 1s ionization potential.⁸⁶ The orbital isosurfaces in Figure 4 confirmed this interpretation and showed that the states associated with the feature at 1564.9 eV were best described as σ -antibonding between one Al 3p orbital and additional Al 3s orbitals in the plane above and below. Finally, the feature at 1570.8 eV in the experimental and simulated spectra was not easily assigned using the simple band structure provided above; however, the LDOS provided by the FEFF calculations indicated that transitions to states with some Al 3p and Al 3d character may have been present.

Aluminum Oxides. The background-subtracted and normalized Al K-edge XAS of α - and γ -Al₂O₃ are shown in Figure 3. The edge onset for each of the oxides was several eV higher in energy relative to the Al foil. These shifts to higher energy for trivalent oxides resulted from an increase in effective nuclear charge and concomitant stabilization of the Al 1s core levels. Previous Al K-edge analyses of α -Al₂O₃ and other aluminum oxides and minerals have provided frameworks for

assigning transitions in the Al K-edge XAS by constructing models for Al–O bonding in both tetrahedral or octahedral coordination environments.^{27,87-91} A similar approach was applied here to develop interpretations of the Al K-edge spectra of α - and γ -Al₂O₃ that could be related to the nanoparticle systems.

The α -Al₂O₃ spectrum had a weak shoulder at 1564.9 eV and two main features at 1567.8 and 1571.9 eV, and closely resembled previously published spectra.^{27,87-91} In their detailed examination of aluminosilicate minerals, Li and coworkers developed assignments by comparing Al K-edge and Si K-edge spectra and using MO calculations of the AlO₆⁹⁻ cluster. Using their assignments for α -Al₂O₃, the very weak and low energy shoulder at 1564.9 eV was assigned to dipole-forbidden transitions from the Al 1s into Al 3s and O 2p antibonding orbitals of a_{1g} symmetry. The first main feature at 1567.8 eV was assigned to transitions from the Al 1s orbitals into Al 3p and O 2p antibonding orbitals of t_{1u} symmetry. The second main feature at 1571.9 eV was assigned to transitions from the Al 1s orbitals into Al 3p and O 2p antibonding orbitals of t_{1u} symmetry. The second main feature at 1571.9 eV was assigned to transitions from the Al 1s orbitals into Al 3p and O 2p antibonding orbitals of t_{1u} symmetry.

The γ -phase of Al₂O₃ has a spinel-type structure consisting of both octahedral and tetrahedral Al coordination sites; hence, the Al K-edge spectrum of γ -Al₂O₃ was regarded as a superposition of Al K-edge spectra involving the two different symmetries. For example, the two main features at 1568.0 and 1571.2 eV reflected the octahedral coordination environment for Al and were assigned to Al 1s \rightarrow t_{1u} and Al 1s \rightarrow t_{2g} transitions as described above for α -Al₂O₃. Using assignments developed previously for γ -Al₂O₃ and aluminosilicate minerals with strictly tetrahedral Al coordination sites,^{74,90} the low-energy shoulder at 1565.4 eV in the Al K-edge XAS of γ -Al₂O₃ was assigned to transitions from Al 1s orbitals into Al 3p and O 2p antibonding orbitals of t₂ symmetry. The 2.6 eV increase in energy for Al 1s \rightarrow t_{1u} transitions relative to the Al 1s \rightarrow t₂ transitions was attributed to an increase in Al–O overlap expected for six-coordinate octahedral environments.⁹² Finally, the additional high-energy shoulder observed at approximately 1575 eV was attributed to a transition from the Al 1s to final states derived from the Al 3d orbitals of e symmetry.

Aluminum Nanoparticles. The spectra from the unsupported metallic Al foil and nanoparticles each exhibited an edge feature at 1559 eV, which indicated that metallic aluminum was a significant component of the nanoparticle materials. However, the Al K-edge spectra also exhibited rich spectral detail reflecting differences in the surface chemistry. For example, each of the nanoparticle materials had a feature between 1565.6 – 1565.9 eV that was not observed in the spectrum from the Al foil. In addition, α -Al₂O₃, γ -Al₂O₃ and both the 18 and 100 nm Al NPs had features at near 1568 eV that were not observed in the spectrum from either the Al foil or the Al–OA NPs. These differences were associated changes in the nature of the oxide layer, and were rationalized by considering the assignments developed above.

The spectral profiles for the conventionally-synthesized 18 and 100 nm Al NPs and the protected Al–OA NPs were similar in that clearly-resolved features with similar intensities and energies were observed at 1565.9, 1565.6, and 1565.6 eV, respectively. These features were also equivalent in energy to those assigned as Al 1s \rightarrow t₂ transitions observed in the Al K-edge XAS of γ -Al₂O₃ (1565.4 eV). The 18 and 100 nm Al NPs were unique in that they also had features in the Al Kedge XAS at intermediate energies of 1568.0 and 1567.8 eV, and also at higher energies of 1571.1 eV for both, which were recognized as distinct from those observed from the Al foil. These intermediate features were very close to the energy observed for the Al 1s \rightarrow t_{1u} transitions in the Al K-edge spectrum of both α -Al₂O₃ (1567.8 eV) and γ -Al₂O₃ (1568.0 eV). Likewise, the higher energy features were similar to the Al $1s \rightarrow t_{2g}$ transitions observed for γ -Al₂O₃ (1571.2 eV). Taken together, these observations showed that the oxidation of 18 and 100 nm Al NPs results in formation an oxide layer containing both tetrahedral and octahedral Al coordination sides, which was similar to the γ -phase of Al₂O₃. Oxidation of the Al–OA NPs also resulted in formation of an Al₂O₃ layer containing tetrahedral Al coordination sites; however, no signatures of octahedral Al coordination sites were detected.

Conclusion

The results provided in this study have shown that soft X-ray spectromicroscopy can provide unique insight into the complex interfacial chemistry of electropositive, oxophilic metals such as aluminum. Because exposure to vacuum or an electron beam can induce compositional changes in some nanoscale particles and surfaces, spectromicroscopy measurements were conducted at ambient pressure with soft X-rays to ensure that quantitative information was obtained.^{38,39} The normal contrast images and respective elemental maps showed that the thickness of the oxide layer was between 2 and 5 nm for all the aluminum materials studied. The images also showed that the oxide was evenly distributed in oleic-acid protected NPs; however, there was more oxide inhomogeneity in bare, unprotected 18 and 100 nm Al NPs. Features in the Al K-edge XAS of the aluminum nanoparticles were also compared with spectra obtained from reference materials including α -Al₂O₃, γ -Al₂O₃, and a 100 nm aluminum metal film. This comparison revealed that the oleic acid protected 18 and 100 nm Al NPs resulted in formation of both octahedral and tetrahedral coordination sites (similar to those found in γ -Al₂O₃). These observations are consistent with previously published results, and also provided new quantitative insight that was not obtained from analysis of PXRD patterns and TEM images alone. Combined, these results support the viewpoint that advanced synthetic strategies including organic passivation do not eliminate surface oxidation of aluminum metal, but can be used to tune the chemistry of the oxide layer and control its reactivity.⁴⁰

Experimental

General Synthetic Details. All reactions were performed either using standard Schlenk line techniques or in a MBraun inert atmosphere glovebox under a purified nitrogen atmosphere (<0.1 ppm O_2/H_2O). The aluminum foil filter (100 nm thickness) and aluminum nanopowders was purchased from Luxel Corporation and US Research Nanomaterials, Inc, respectively, and used as received. Oleic acid aluminum core-shell nanoparticles were synthesized according to the literature procedure and dried before use.⁴⁰ α and γ -Al₂O₃ were purchased from Aldrich and Strem and dried at 100 °C and 1 x 10⁻³ torr for 24 h before use. Sample preparation and methodology for handling air-sensitive analytes was similar to that described previously.^{29,31,93} Samples were prepared in an argon-filled glovebox by grinding crystals of the analyte into a fine powder with a mortar and pestle, and brushing the powder onto a Si₃N₄ membrane (100 nm, Silson) with a fiber. This method arranged a large number of micron-scale crystals that were suitable for Al K-edge measurements in a relatively compact area. A second membrane was placed over the sample, and the edges were sealed together using Hardman

Double/Bubble[®] 5 minute epoxy.

Beamline Characteristics. STXM methodology was similar to that discussed previously.⁸⁸⁻⁹¹ Single-energy images and Al K-edge XANES spectra were acquired using the STXM instrument at the Molecular Environmental Science (MES) beamline 11.0.2 at the Advanced Light Source (ALS), which was operated in topoff mode at 500 mA, in a ~0.5 atm Hefilled chamber.⁶¹⁻⁶³ The beamline used photons from an elliptically polarizing undulator that delivers photons in the 100 to 2000 eV energy range to a variable-included-angle plane-grating monochromator. For Al K-edge measurements, the photon energy of the high energy diffraction grating (1100 lines/mm) was calibrated to the Al K-edge from the 100 nm aluminum foil filter sample (inflection point = 1559.0 eV). Both of the 4-jaw exit slit widths were set to 30 μ m. The maximum energy resolution $E/\Delta E$ was previously determined at better than 7500, which was consistent with the observed standard deviation for spectral transitions of ± 0.3 eV determined from comparison of spectral features over multiple samples and beam runs. For these measurements, the X-ray beam was focused with a 40 nm zone plate onto the sample, and the transmitted light was detected with a photomultiplier tube. Images at a single energy were obtained by rasterscanning the sample and collecting transmitted monochromatic light as a function of sample position. Spectra at particular regions of interest on the sample image were extracted from the "stack", which is a collection of images recorded at multiple, closely spaced photon energies across the absorption edge. Dwell times used to acquire an image at a single photon energy were 2 ms per pixel. To quantify the absorbance signal, the measured transmitted intensity (I) was converted to optical density using Beer–Lambert's law: $OD = ln(I/I_0) = \mu \rho d$, where I_0 is the incident photon flux intensity, d is the sample thickness, and μ and ρ are the mass absorption coefficient and density of the sample material, respectively. Incident beam intensity was measured through the sample-free region of the Si₃N₄ windows. Regions of particles with an absorption of >1.5 OD were omitted to ensure the spectra were in the linear regime of the Beer-Lambert law. The data were background subtracted using the MBACK algorithm in MATLAB,94 and peak positions were determined from derivative plots generated using the program IGOR 6.0.

FEFF Calculations. FEFF 10.0.1^{72,73} calculations utilized a Hedin-Lundqvist potential and converged using a selfconsistent field approximation within a 6.0 Å cluster of 79 atoms and a full-multiple scattering expansion using a cluster of 55 atoms within a radius 6.6 Å. All other options were chosen to be default.

Al K-edge Spectral Simulations. XAS at the Al K-edge were calculated using the excited electron and corehole approach (XCH).⁹⁴ The XCH approach has been described previously in detail.^{94.96} In this method, the lowest energy X-ray excited state of the system is modeled within an occupation-constrained DFT framework employing a periodic supercell formalism wherein the core-excited atom is described through a core-hole pseudopotential and the screening due to the excited electron is taken into account self-consistently. Higher–lying X-ray excited state energies are approximated through eigenvalue differences obtained from the Kohn-Sham (KS) spectrum of the lowest energy core-excited state. X-ray transition matrix elements are calculated using Fermi's golden rule and typically for light-element K-edges, within a dipole approximation. The XCH method was utilized as implemented in a development version of the Quantum-Espresso package⁹⁷ which provides a plane-wave pseudopotential DFT framework for electronic structure calculations. Ultra-soft pseudopotentials⁹⁸ with the following valence electronic configurations were employed: Al(3s²3p¹), O(2s²2p⁴). These were used together with a plane-wave energy cut-off of 25 Ry. To describe 1s core-excited Al in XANES simulations, a core-hole pseudopotential with the electronic configuration $1s^{1}2s^{2}2p^{6}3s^{2}3p^{2}$ was generated. A 3 x 3 x 3 supercell consisting of 108 atoms was used. The Brillouin zone was sampled through a 3 x 3 x 3 Γ -centered k-point grid during the Kohn-Sham self-consistent field calculation, but the band structure was interpolated over a uniform Γ -centered 5 × 5 × 5 k-point grid using the Shirley interpolation scheme⁹⁹ in order to generate XANES spectra. Within the XCH approach, the relative energy alignment of XAS corresponding to core-excitations on atoms in different chemical environments is carried out through total energy differences.⁹⁵ However, since the method is based on core-hole pseudopotentials and does not explicitly include the energy of core electrons, calibration of the absolute energy position of the calculated spectra with respect to experiment involves a rigid energy shift Δ , which was determined to be 1560.65 eV at the Al K-edge based on the inflection point of the onset of absorption. Orbital isosurfaces were visualized using VESTA-3.¹⁰⁰

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Supporting Information Available: Additional details on the compositional analysis, X-ray absorption spectra, and powder X-ray diffraction data. This material is available free of charge via the Internet at http://pubs.acs.org.

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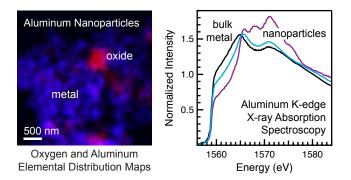
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Synopsis: Composition, structure, and thickness of the Al_2O_3 layers on aluminum metal surfaces were determined using Xray absorption spectroscopy and imaging from a scanning transmission X-ray microscope. The results showed that synthetic strategies including organic passivation may not eliminate surface oxidation, but can be used to tune the chemistry of the oxide layer and control its reactivity. The unique technical approach provided chemical insights that could not be obtained using conventional methods of characterization.