





# Coherence of Eddy Kinetic Energy Variation during Eddy Life Span to Low-Frequency Ageostrophic Energy

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Abstract: The evolution of mesoscale eddies is crucial for understanding the ocean energy cascade. In this study, using global reanalysis sea surface velocity data and a mesoscale eddy trajectory product tracked by satellite altimeters, we aimed to reveal the coherence of eddy kinetic energy (EKE) variation to low-frequency ageostrophic energy during the eddy life span. The variation in EKE throughout the eddy life span was highly coherent to that of the seven-day low-passed ageostrophic kinetic energy, with a correlation coefficient of -0.94. The low-frequency ageostrophic motions supplied 38% of the EKE variation in the growing stage of mesoscale eddies and absorbed 42% in the decaying stage. The evolution rate of the EKE during the eddy life span was consistent with the barotropic conversion rate of the low-frequency ageostrophic motions, further confirming the dominant role of low-frequency ageostrophic motions in eddy growth and decay.

**Keywords:** mesoscale eddy; eddy kinetic energy; eddy life span; low-frequency motion; barotropic conversion rate

# 1. Introduction

Mesoscale eddies are broadly distributed in the global ocean, transporting and regulating the ocean's physical and biogeochemical properties [1–8]. They have a spatial scale of tens to hundreds of kilometers and contain about 90% of the sea surface kinetic energy [9,10]. Mesoscale eddies play crucial roles in the ocean energy cascade. The eddies generally obtain energy from large-scale background currents by barotropic and baroclinic conversion and transfer energy down to submesoscale and smaller scale processes due to ageostrophic instabilities, surface and bottom frictions, and turbulent dissipation [11–18].

Previous studies have revealed the spatiotemporal variations of the EKE in the global ocean using intensive satellite altimeter data and numerical models [18–25]. The variation in kinetic energy throughout the eddy life span has also been derived from sea surface drifters and Argo floats [14]. In climatology, the variation in mesoscale eddy energy is generally divided into three stages over the whole eddy life, i.e., growing, mature, and decaying stages. The EKE rises in the growing stage, sustains a flat plateau in the mature stage, and then falls in the decaying stage [14,26–28]. In addition, the mean EKE of cyclonic eddies (CEs) is generally higher than that of anticyclonic eddies (AEs) [29].

Zhang and Qiu [14] analyzed the variation in submesoscale ageostrophic energy and proposed that submesoscale motions play a critical role in the EKE variation throughout the eddy life span. On the other hand, Huang et al. [30] found that the variation of eddy rotational speed is negatively correlated with that of the current-regulating eddy propagation velocity, indicating the role of background currents in eddy energy variation. Previous



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**Copyright:** © 2022 by the authors. Licensee MDPI, Basel, Switzerland. This article is an open access article distributed under the terms and conditions of the Creative Commons Attribution (CC BY) license (https:// creativecommons.org/licenses/by/ 4.0/). ocean energy analysis has shown that energy conversion can occur between mesoscale eddies and large-scale currents through barotropic and baroclinic instabilities [13,31–35]. However, the quantitative effect of large-scale and low-frequency motions on mesoscale eddy kinetic energy variation throughout the eddy life span remains unclear.

In this study, we aim to quantitatively analyze the contribution of low-frequency ageostrophic motions to the EKE variation throughout the eddy life span in the global ocean, using the eddy trajectory dataset, geostrophic velocity anomaly, and reanalysis current data. The data and methods are described in Section 2. The energy variations of mesoscale eddies and low-frequency ageostrophic motions are compared, and the mechanism of the energy conversion due to barotropic instability (BT) is analyzed in Section 3. Discussion and conclusions are given in Sections 4 and 5, respectively.

# 2. Data and Methods

#### 2.1. Satellite Altimetry Data and Eddy Tracking Dataset

The Archiving Validation and Interpretation of Satellite Oceanographic (AVISO) datasets provided the global gridded geostrophic velocity anomaly data  $U_{ga}$  used in this study, which were accessed from the Copernicus Marine Environment Monitoring Service (CMEMS) website. The datasets were merged products derived from Topography Experiment (TOPEX) Poseidon (T/P), Jason-1/2 (French–US altimeter satellites), ERS-1/2 (European Remote Sensing satellites), and Envisat (European Remote Sensing satellite) altimeters. The gridded data were geophysically/meteorologically corrected (tides, ionosphere, wet, and dry troposphere) and interpolated onto Mercator grids of 0.25° horizontal resolution with global ocean coverage except for high latitude oceans [36]. The temporal resolution of the data is one day, and the time coverage is from 1 January 1993 to 31 December 2019.

The AVISO 27-year mesoscale eddy trajectories atlas product (version 2.0 delayed-time) derived from the satellite altimeter observations is used as a baseline for the study [37]. The daily eddy trajectory data include seven eddy parameters, i.e., tracking number, eddy center position (latitude and longitude), tracking day, polarity information, maximum circum-averaged speed, and eddy radius *R*. The eddy life span  $T_{span}$  is calculated from the track day of eddies with the same tracking number. The data span the period from 1 January 1993 to 31 December 2019 as the geostrophic velocity anomaly  $U_{ga}$ .

The numbers, radius, and lifespan of CEs and AEs in the global ocean are listed in Table 1. One can see that a total of 378,513 eddies are used for energy analysis in this study, of which the number of CEs is slightly higher (3.6%) than that of AEs. In fact, there are generally more CEs than AEs in each ocean, except in the North Atlantic, where the number of CEs is 2.1% lower than AEs. The eddies have a median (mean) lifespan of 51 (73.38) days and a median (mean) radius of 74.20 (82.33) km in total, while the AEs have a larger radius and a longer lifespan than the CEs.

	All	AEs	CEs
Eddy number	378,513	185,937	192,577
Eddy radius (km)	74.20	74.95	73.45
median (mean $\pm$ std)	$(82.33 \pm 35.93)$	$(83.13 \pm 36.23)$	$(81.53 \pm 35.60)$
Eddy lifespan (day)	51.00	51.00	51.00
median (mean $\pm$ std)	$(73.38 \pm 66.40)$	$(74.43 \pm 69.62)$	$(72.60 \pm 63.33)$

Table 1. The numbers, radius, and lifespan of mesoscale eddies in the global ocean.

#### 2.2. Reanalysis Dataset of Total Current Velocity

The total current velocity U, including geostrophic and ageostrophic components, is derived from the global ocean eddy-resolving reanalysis dataset (GLORYS12V1) provided by Copernicus Marine Environment Monitoring Service (CMEMS). The model component is the NEMO (Nucleus for European Modeling of the Ocean) platform, driven at the surface by ECMWF (European Centre for Medium-Range Weather Forecasts) ERA-Interim then ERA5 reanalyzes for recent years. A reduced-order Kalman filter was used to assimilate observations, including along track altimeter data, Satellite Sea Surface Temperature, Sea Ice Concentration, and in situ temperature and salinity (T/S) vertical Profiles. Moreover, a 3D-VAR scheme corrects for temperature and salinity biases that are slowly evolving on a large scale. The spatial and temporal resolutions of the dataset are (1/12)° and one day, respectively. The time coverage is from 1 January 1993 to 31 December 2019, the same as that of  $U_{ga}$ .

We compare the reanalysis current velocity to that from surface six-hourly drifter observations; the derived kinetic energy along eddy tracks is similar to each other (not shown), indicating the reliability of the reanalysis velocity. Furthermore, the reanalysis data have suitable continuity and concurrence with the satellite eddy observations.

## 2.3. Calculation of Geostrophic and Ageostrophic Kinetic Energy

The total current velocities  $\boldsymbol{U}$  include all velocity components induced by different dynamic processes, of which the geostrophic velocity can be provided by the satellite observations. Here, we decompose the total velocity at each grid  $\boldsymbol{U}$  into three components: climatologic mean velocity  $\boldsymbol{U}_0$  from time-averaged  $\boldsymbol{U}$ , geostrophic current velocity anomaly  $\boldsymbol{U}_{ga}$  from satellite observations, and ageostrophic velocity  $\boldsymbol{U}_{ag}$ :

$$\boldsymbol{U}_{ag}(i,j,t) = \boldsymbol{U}(i,j,t) - \boldsymbol{U}_{0}(i,j) - \boldsymbol{U}_{ga}(i,j,t)$$
(1)

where *i* and *j* are the zonal and meridional grids, respectively, and *t* is the time. Thus, eddy kinetic energy (EKE)  $E_{ga}$ , and ageostrophic kinetic energy (AKE)  $E_{ag}$  are calculated as follows:

$$E_{ga}(i,j,t) = \frac{|\mathbf{u}_{ga}(i,j,t)|^2}{2}$$

$$E_{ag}(i,j,t) = \frac{|\mathbf{u}_{ag}(i,j,t)|^2}{2}$$
(2)

We use a moving average filter to obtain the low-frequency part of ageostrophic velocity  $U_{alp}$  from the time series of  $U_{ag}$  at each data grid point; the remaining part of  $U_{ag}$ , subtracting  $U_{alp}$ , is the high-frequency ageostrophic velocity  $U_{ahp}$ , i.e.,  $U_{ahp} = U_{ag} - U_{alp}$ . Considering the periods of the submesoscale motions, a seven-day cutoff frequency is used, as in previous studies [14,38,39]. Since the temporal and spatial scales of ocean motions are generally coherent, the low-frequency ageostrophic motions are generally overlapped by the large-scale motions. Similar to the calculation of  $E_{ag}$ , the high-frequency ageostrophic energy (HAE)  $E_{ahp}$  and the low-frequency ageostrophic energy (LAE)  $E_{alp}$  are calculated from  $U_{ahp}$  and  $U_{alp}$ , respectively.

Figure 1a,b show the global distributions of climatological mean EKE and AKE  $\overline{E_{ga}}$  and  $\overline{E_{ag}}$ , respectively. One can see that the high values of  $\overline{E_{ga}}$  above 600 cm<sup>2</sup>/s<sup>2</sup> are mainly distributed in the western boundary current, the Antarctic circumpolar current (ACC), and the equatorial current regions, where maximum values are even higher than 1000 cm<sup>2</sup>/s<sup>2</sup>. Mesoscale eddies are frequently active in regions with strong currents due to barotropic and baroclinic instabilities of the background currents, while ageostrophic motions could be enhanced due to active nonlinear eddy–eddy interactions and eddy–current interactions. In addition, ageostrophic forces, such as buoyancy fluxes and wind forcing, are generally large in these areas and induce high ageostrophic energy [40–42]. In shallow shelf waters, bottom friction may also contribute to the ageostrophic motions [15,43]. Though high values of  $\overline{E_{ag}}$  and  $\overline{E_{ga}}$  are in close regions, they are dislocated in detail since high ageostrophic motions are prone to appear around mesoscale eddies.  $\overline{E_{ga}}$  have more coverage of high values than  $\overline{E_{ag}}$  indeed.

The difference between  $\overline{E_{ga}}$  and  $\overline{E_{ag}}$  is shown in Figure 1c. The positive values represent stronger geostrophic activities, which are mainly distributed in the western boundary current region and the ACC region, while higher ageostrophic kinetic energy are mainly concentrated in the equatorial region and continental shelf seas, where large winds or bottom friction occurs. The distributions of climatologic LAE  $\overline{E_{alp}}$  and HAE  $\overline{E_{ahp}}$  are sim-

(a`

60°N

30°N

0°

30°S

60°S

60°N

30°N

0°

30°S

60°S

60°N

30°N

0°

30°S

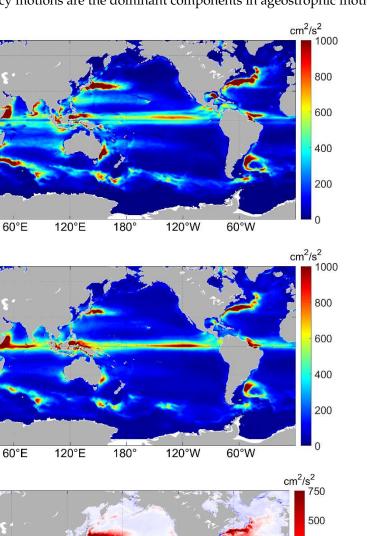
60°S

0°

0°

0°

(b



ilar to that of  $\overline{E_{ag}}$ , but  $\overline{E_{alp}}$  is four times greater than  $\overline{E_{ahp}}$  (not shown). In other words, low-frequency motions are the dominant components in ageostrophic motions.

**Figure 1.** Global distributions of climatological kinetic energy. (a) Eddy geostrophic kinetic energy  $\overline{E_{ga}}$ , (b) ageostrophic kinetic energy  $\overline{E_{ag}}$ , (c) differences between  $\overline{E_{ga}}$  and  $\overline{E_{ag}}$ .

60°W

120°W

250

0

-250

-500

-750

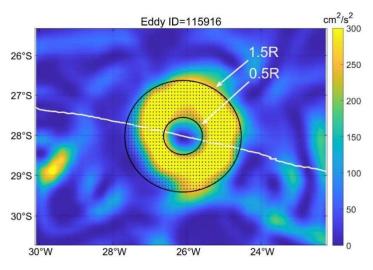
# 2.4. Matching Gridded Data and Normalization

120°E

180°

60°E

The daily center position and corresponding radius *R* of the eddies are derived from the eddy-tracking dataset. The area within 0.5–1.5 *R* around the eddy center is chosen as the main contribution area of the EKE, as in a previous study [14]. The gridded Ega within this area is averaged to represent the EKE of mesoscale eddies at time t, as shown in Figure 2. As a result,  $E_{ga}(eddy,t)$  for a certain eddy can be captured along the eddy track on each day *t* of the eddy life span. In fact, the variation of  $E_{ga}(eddy,t)$  here within 0.5–1.5 *R* is consistent with that derived from the maximum speed of eddies (not shown), although the mean  $E_{ga}$  within 0.5–1.5 *R* is smaller than that from the maximum speed.



**Figure 2.** Matching gridded data along eddy track. The background image shows the EKE field, the white curve represents the eddy trajectory, the black circles give the range of the eddy 0.5–1.5 *R*, and the scatters between the two concentric circles represent the data points used in the energy calculation for each eddy.

The life span of each eddy is then normalized by its total lifespan  $T_{span}$ , i.e.,  $T_e = t/T_{span}$ . Averaging all energy curves along the eddy track in the normalized life span, we could gain the direct mean energy curve  $E_{xx}$  ( $T_e$ ) from all eddies, i.e.,

$$E_{xx}(T_e) = \frac{\sum\limits_{i=1}^{N} E_{xx}(eddy, T_e)}{N}$$
(3)

where *N* is the number of eddy samples, and the subscript *xx* represents the energy categories, such as *ga*, *ag*, *alp* and *ahp*.

In order to clarify the portion of energy variation, we normalize the energies on each day of the life cycle  $E_{xx}(eddy, T_e)$  by its life-mean EKE  $E_{ga}^{ave}(eddy)$ :

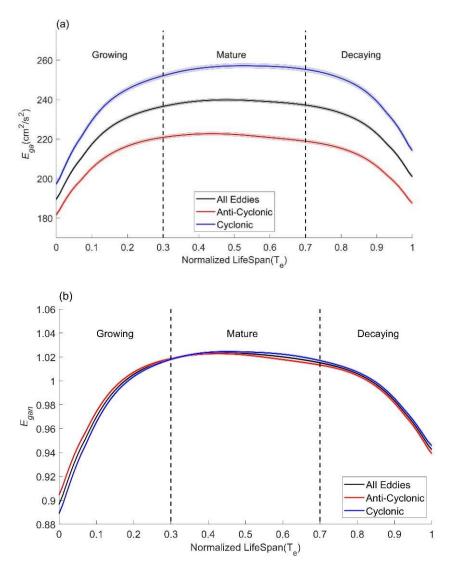
$$E_{xxn}(eddy, T_e) = \frac{E_{xx}(eddy, T_e)}{E_{aa}^{ave}(eddy)}$$
(4)

Finally,  $E_{xxn}(eddy, T_e)$  along the track in the normalized life span for all mesoscale eddies are averaged to obtain the mean curve of  $E_{xxn}(T_e)$  throughout the eddy life cycle, i.e.,  $E_{xxn}(T_e) = \sum_{i=1}^{N} E_{xxn}(eddy, T_e)/N$ . Therefore, the variation curves of EKE  $E_{agn}$ , AKE  $E_{agn}$ , LKE  $E_{alpn}$ , and HKE  $E_{ahpn}$  are derived. The mean standard error,  $std/N^{1/2}$ , is used to measure the deviation of the mean curves, where std is the standard deviation of the data and N is the number of averaged eddy samples.

#### 3. Results

## 3.1. Eddy Kinetic Energy Variation during Eddy Life Span

Figure 3a shows the variation of the direct mean EKE  $E_{ga}$  for all eddies (black curve), AEs (red curve), and CEs (blue curve) along their tracks during their normalized life span. One can see that the three curves show a similar pattern, i.e., increasing at the beginning, maintaining in the middle, and decreasing at the end. Specifically,  $E_{ga}$  for all eddies increases from 190 cm<sup>2</sup>/s<sup>2</sup> at the beginning of the life span to 237 cm<sup>2</sup>/s<sup>2</sup> at 0.3  $T_e$ . Subsequently,  $E_{ga}$  maintains around 240 cm<sup>2</sup>/s<sup>2</sup> until 0.7  $T_e$  and decreases rapidly to 200 cm<sup>2</sup>/s<sup>2</sup> at the end of the life span (black line). The mean  $E_{ga}$  over the life span is 229 cm<sup>2</sup>/s<sup>2</sup>, lower than 245 cm<sup>2</sup>/s<sup>2</sup> for CEs, and higher than 213 cm<sup>2</sup>/s<sup>2</sup> for AEs.



**Figure 3.** Curves of global mean eddy kinetic energy along the eddy track throughout the eddy life span. (a) EKE  $E_{ga}$ , (b) Normalized EKE  $E_{gan}$ . Black curves represent the energy curves for total eddies. Red and blue curves represent the energy curves for AEs and CEs, respectively. Shadows around the energy curves represent the mean standard error.

According to the changing rate of  $E_{ga}(T_e)$  for total eddies with time, we divide the whole eddy life span into three stages with a critical slope of 0.35 cm<sup>2</sup>/s<sup>2</sup> per 0.01  $T_{e}$ : the growing stage from 0 to 0.3  $T_e$ , the mature stage from 0.3 to 0.7  $T_e$ , and the decaying stage from 0.7 to 1  $T_e$ , as shown in Figure 3a. The energy curves of AEs (red line) and CEs (blue line) have similar patterns to the total one (black line), while the EKE values of CEs are systematically higher than those of AEs, as also observed by Chen and Han [29]. Note that the curves of  $E_{ga}(T_e)$  are not symmetric. The energy values at the beginning are 5–15 cm<sup>2</sup>/s<sup>2</sup> smaller than that at the end of the eddy life span. The differences between the curves of CEs and AEs in the growing stage are also smaller than that in the decaying stage.

The global mean of the normalized EKE  $E_{gan}(T_e)$  curve is shown in Figure 3b. One can see that the mean  $E_{gan}(T_e)$  was 1.02 in the mature stage, while the minimum values are 0.90 and 0.94 at the beginning and the end of the eddy life span, respectively. In contrast to the higher  $E_{gan}(T_e)$  for CEs in Figure 3a, the normalized  $E_{gan}(T_e)$  of AEs is generally close to that of CEs, and the  $E_{gan}(T_e)$  of AEs has larger values in the growing stage. The difference may be attributed to more strong CEs than AEs. The statistics reveal 25.7% of CEs have a  $E_{ga}^{ave}(eddy)$  larger than 240 cm<sup>2</sup>/s<sup>2</sup>, while the ratio is only 23.2% for AEs. In the direct mean of  $E_{ga}(T_e)$ , strong eddies contribute more to the mean value, while the normalized mean

curve  $E_{gan}(T_e)$  reveals the real portion of energy variation over the life span for each eddy. On the other hand, the asymmetry of the curve in the beginning and ending are enhanced in the normalized curve since the eddies are weaker in the growing and decaying stages.

Table 2 lists the energy increments of individual components,  $\Delta E_{xx}^G$  in the growing stage and  $\Delta E_{xx}^D$  in the decaying stage for all eddies. One can see that  $\Delta E_{gan}^G = 0.12$  and  $\Delta E_{gan}^D = -0.08$ , i.e.,  $E_{gan}$  increases by 12.1% of  $E_{ga}^{ave}(eddy)$  in the growing stage and decreases by 8% of  $E_{ga}^{ave}(eddy)$  in the decaying stage.

**Table 2.** Increments of  $E_{gn}$ ,  $E_{an}$ ,  $E_{alpn}$ , and  $E_{ahpn}$  and contributions to  $E_{gn}$  in the growing and decaying stages.

	Egn	Ean	E <sub>alpn</sub>	E <sub>ahpn</sub>
Growing stage $\Delta E_{xx}^G (\Delta E_{xx}^G / \Delta E_{gn}^G)$	0.12 (100%)	-0.07 (57%)	-0.05 (38%)	-0.01 (8%)
Decaying stage $\Delta E_{xx}^D (\Delta E_{xx}^D / \Delta E_{gn}^D)$	-0.08 (100%)	0.06 (75%)	0.03 (42%)	0.01 (15%)

# 3.2. Ageostrophic Kinetic Energy Variation during Eddy Life Span

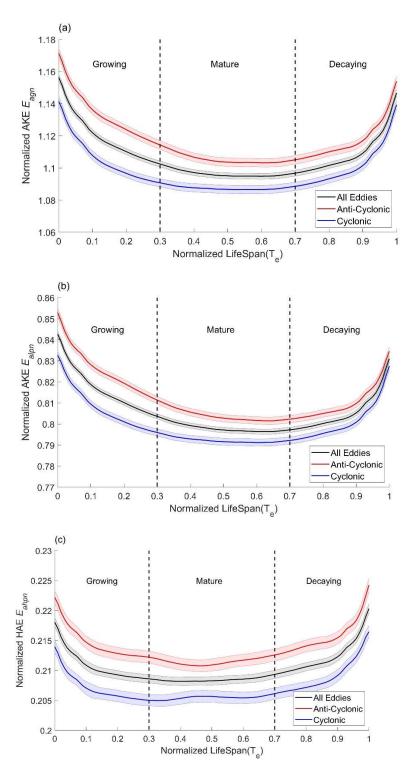
Figure 4a shows the mean curve of the normalized AKE  $E_{agn}(T_e)$  during the eddy life span. One can see that the  $E_{agn}$  of total eddies has a maximum value of 1.16 in the growing stage, a minimum value of 1.09 in the mature stage, and 1.15 in the decaying stage. It has an upside-down variation trend compared to the curve of  $E_{gan}$  shown in Figure 3b. As listed in Table 2,  $E_{agn}$  decreased by 0.07 in the eddy growing stage and increased by 0.06 in the decaying stage, which accounted for 57% and 75% of  $|\Delta E_{gan}|$  in the two stages, respectively. CEs and AEs both have similar variation trends in the normalized AKE, while the values of AEs are higher than that of CEs, probably due to stronger geostrophic strain in AEs [14].

The mean curves of the normalized LAE  $E_{alpn}$  (>7 d) and the normalized HAE  $E_{ahpn}$  (<7 d) during the eddy life span are further shown in Figure 4b,c. One can see that the energy curves of both  $E_{alpn}$  and  $E_{ahpn}$  are similar to that of the total ageostrophic energy  $E_{agn}$  shown in Figure 4a. The maximum and minimum values of  $E_{alpn}$  are 0.84, 0.80, and 0.83 in the three stages, accounting for 62% of the values of total ageostrophic energy  $E_{agn}$ . On the other hand, the extreme values of  $E_{ahpn}$  for all eddies are only 0.22, 0.21, and 0.22 in the three stages. They are four times lower than those of  $E_{alpn}$ , indicating that  $E_{alpn}$  is the dominant component in  $E_{an}$ .

### 3.3. Analysis of Coherence of EKE Variation to Low-Frequency Ageostrophic Motions

As shown in Figure 4, the curves of  $E_{alpn}$  and  $E_{ahpn}$  are highly coherent with that of  $E_{gan}$  during the eddy life span. The correlation coefficient reaches -0.94 and -0.89 with a confidence level of 99% for  $E_{alpn}$  and  $E_{ahpn}$ , respectively. This indicates that both low-frequency and high-frequency ageostrophic motions contribute to EKE variation during the eddy life span, while the low-frequency ageostrophic motions are more closely correlated to the EKE variation.

Table 2 lists the ratios of the contribution. One can see that in the growing stage,  $E_{alpn}$  decreases by 0.05 from 0.84 to 0.79, contributing 38% of  $\Delta E_{gan}^G$ . On the other hand,  $E_{alpn}$  increases by 0.03 (42% of  $\Delta E_{gan}^D$ ) in the decaying stage. On average, the energy provided by the low-frequency ageostrophic motions contributes 40% of the total variation of  $E_{gan}$ . For comparison, the contributions of high-frequency ageostrophic motions to the variation of  $E_{gan}$  are also listed in Table 2. The ratios are only 8% and 15% in the growing and decaying stages, respectively, 3–5 times lower than the ratio of  $E_{alpn}$  to  $E_{gan}$ . The low-frequency ageostrophic motions are a major contributor.

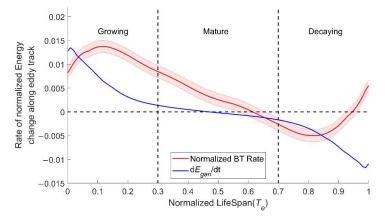


**Figure 4.** Variation curves of global mean ageostrophic kinetic energy along the eddy track throughout the eddy life span. (a) Normalized AKE  $E_{agn}$ , (b) normalized LAE  $E_{alpn}$  (with period > 7 d), and (c) normalized HAE  $E_{ahpn}$  (with period < 7 d). Black curves represent the energy curves for total eddies. Red and blue curves represent the energy curves for AEs and CEs, respectively. Shadows around the energy curves represent the mean standard error.

#### 3.4. Barotropic Conversion Rate

In order to investigate the interactions among eddies and low-frequency ageostrophic motions, we performed energy analysis to quantify the barotropic (BT) conversion rate between the EKE and the LAE. We first integrate the EKE within the range of 0.5-1.5 R

and obtain the integrated-EKE (iEKE) curve for each eddy. Then, the integrated-EKE curve for each eddy is normalized by its life-mean EKE  $E_{ga}^{ave}(eddy)$ , and the global mean of the normalized iEKE curve is obtained, which has similar patterns to the EKE  $E_{ga}$  curve (not shown). We calculate the time derivative  $dE_{gan}/dt$  of the mean normalized iEKE curve in the eddy life span. The resulted  $dE_{gan}/dt$  is shown by the blue curve in Figure 5.



**Figure 5.** Time derivative of normalized iEKE (blue line) and normalized barotropic conversion rate (BTn) (red line) throughout the eddy life span. Shadows around the BTn curves represent the mean standard error.

For the eddy energy variation, the total derivative of eddy kinetic energy without external forcing depends on four dynamic terms, i.e., the shear production of the background current (i.e., BT conversion rate), the buoyant production (i.e., BC conversion rate), the viscosity dissipation, and the energy flux divergence [13,44]. In which the BT conversion rate represents the energy transfer from the background current to the eddy geostrophic current. It is calculated as follows [13]:

$$BT: -\rho_0[\overline{u_{ga}(i,j,t) \cdot \mathbf{U}_{ga}(i,j,t)} \cdot \nabla \left( u_{alp}(i,j,t) + u_0(i,j) \right) + \overline{v_{ga}(i,j,t) \cdot \mathbf{U}_{ga}(i,j,t)} \cdot \nabla \left( v_{alp}(i,j,t) + v_0(i,j) \right)]$$
(5)

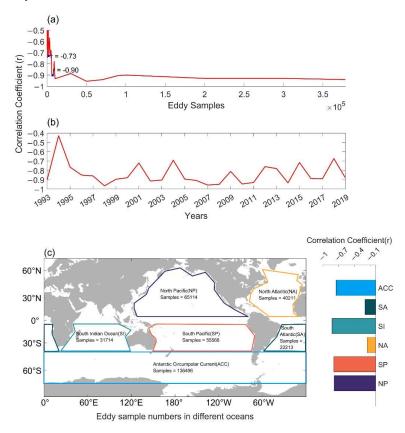
We calculate the BT rate for data grids first and then select values along the eddy track to calculate the integrated BT rate within the range of 0.5–1.5 *R* during the eddy life span. Then, the integrated BT rate curve is also normalized by its life-mean EKE  $E_{ga}^{ave}(eddy)$  for each eddy. Finally, the curves of all global eddies are averaged to obtain the red curve shown in Figure 5. One can see that the normalized BT rate is positive in the eddy growing and mature stage, indicating that the energy transfers from low-frequency motions to eddies, while the BT rate is negative in the decaying stage, where the energy transfers inversely from EKE to LAE.

Comparing the BT rate to the change rate of EKE along the eddy track ( $dE_{gan}/dt$ , blue line in Figure 5), one can see that the two rates have the same decreasing trend during the eddy life span and change the sign of variation rate in the mature stage. Their correlation coefficient reaches 0.73 with 99% significance, indicating that the BT rate is the dominant term in the eddy energy variation. The BT rate is generally larger than the EKE variation rate  $dE_{gan}/dt$ . In the eddy growing stage, EKE increases with a positive dE/dt as LAE transfers to EKE with a positive BT rate; EKE varies little with  $|dE_{gan}/dt|$  less than  $1.5 \times 10^{-3}$  as the BT rate decrease rapidly in the mature stage. Then, EKE decreases with a negative rate in the decaying stage, as the BT rate becomes negative. It is clear that low-frequency ageostrophic motions serve as the energy sources for eddy growth and the energy sink for eddy decay. Furthermore, the BT rate was positive between 0 and 0.62  $T_e$ , while  $dE_{gan}/dt$  becomes negative after 0.47  $T_e$ , implying EKE loss possibly probably due to high-frequency ageostrophic motions, viscous dissipation, wind killing, and energy transfer to potential energy [32,45,46]. The positive anomaly after 0.95  $T_e$  may be caused by eddy shape deforming or eddy-mean flow interaction at the western boundary [47,48].

#### 4. Discussion

Previous investigators found a high correlation between the high-passed ageostrophic kinetic energy and eddy geostrophic kinetic energy during the eddy life span and proposed the importance of submesoscale processes in eddy kinetic energy variation [14]. In this study, we found close coherence in the EKE variation to low-frequency ageostrophic motions such as slowly varying wind-driven flow and planet waves because of their shared frequency band with mesoscale eddies rather than submesoscale processes.

It should be noted that the high coherence of variations of EKE and LAE in this study are a composite result of total eddies in the global ocean. The number of eddy samples N, as well as the locations and occurrence time of eddies, may have an influence on the coherence between EKE and LAE. Figure 6a shows the variation of the correlation coefficient between  $E_{gn}$  and  $E_{alpn}$ , marked as  $r_{j}$  with different numbers of random eddy samples in the global ocean. One can see that r oscillates around 0 as the number of eddy samples N is less than 2000, increases gradually from about -0.73 at N = 2000 to -0.90 as N reaches 7000, and keeps to a little more than -0.92 with N > 10,000. It indicates that the coherence between the curves of EKE and LAE could be established if there were more than 2000 eddies. We further calculate r between  $E_{gn}$  and  $E_{alpn}$  of the eddies in each year. As shown in Figure 6b, r has the largest value of -0.95 in 1998 and then varies between -0.7 and -0.94 with a 2–3 year period in the recent 20 years, which seems to be related to the interannual variability of eddies and background currents influenced by ENSO [23,49,50]. For spatial variability of the coherence shown in Figure 6c, the correlation r between  $E_{gn}$  and  $E_{alpn}$  is larger than -0.8 in most oceans except in the Atlantic, probably due to the narrow and asymmetric ocean basin. Mechanisms about the coherence variation need further pursuing.



**Figure 6.** Variation of correlation coefficient between  $E_{gan}$  and  $E_{alpn}$  curves with (**a**) eddy samples, (**b**) eddy occurrence years, and (**c**) ocean basins.

We find three-stage variations of eddy energy during eddy life for global ocean eddies. In fact, such composite three-stage evolution has been reported for eddies both in the open ocean [2,14,27,29] and in regional seas [26,28,29,51], though the variation rates and the span of each stage may be different in different seas in different years. It should be noted that the three-stage pattern is also the composite result of large amounts of eddies. For a single eddy, the energy curve during its life cycle may be complicated due to various forcing [52,53]. The interaction between eddies, such as merging and splitting, could also cause eddy energy variation during their lives [54,55]. In addition, only eddies with a life cycle of more than 30 days are considered in this study due to data limitations. It will be helpful to comprehensively understand the evolution of mesoscale eddies by separating eddies within different ranges of radius and life span. Furthermore, only sea surface current energy and BT conversion rate are analyzed; the baroclinic instability process and eddy energy conversion to ageostrophic current could be further explored.

## 5. Conclusions

Using satellite altimetry data and reanalysis current velocity data from 1993 to 2019, this study analyzes the role of low-frequency ageostrophic motions in the variation of eddy kinetic energy during eddy life span in global ocean. The quantitative analysis showed that the  $E_{gan}$  curve and the  $E_{alpn}$  curve throughout the eddy life span have a correlation coefficient of -0.94. Low-frequency ageostrophic motion contributes about 38% of the EKE variation in the eddy generation phase and 42% in the decaying phase. Thus, it is the low-frequency ageostrophic motions that are closely coherent with the EKE variation during the eddy life span. In addition, the BT conversion rate analysis confirmed the relationship between the energy transfer between LAE and EKE. Moreover, the EKE decrease in the decaying stage may also be attributed to other processes such as wind work, submesoscale processes, and topography effect.

The new findings in this study may help to improve the understanding of the mechanisms involved in the generation and dissipation of mesoscale eddies and the role of mesoscale eddies in the global ocean energy cascade. Meanwhile, the results in the ocean may also provide a reference for the other geophysical fluid on the Earth, the atmosphere. Different from the same-order geo- and ageostrophic energies in the ocean (Figure 1c), the ageostrophic energy in the atmosphere is weaker than that of geostrophic motion [5,14,56]. Nonetheless, ageostrophic motion is still important in weather-scale processes in the atmosphere and contributes to the enhancement of the geostrophic field [15,57]. The role of ageostrophic motions in the atmospheric energy cascade could be compared to that in the ocean in further studies.

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**Data Availability Statement:** The mesoscale eddy trajectories atlas product META 2.0 DT is distributed by AVISO at: https://www.aviso.altimetry.fr/en/data/products/value-added-products/ global-mesoscale-eddy-trajectory-product.html, (accessed on 22 July 2020). The global gridded current velocity can be accessed at: https://resources.marine.copernicus.eu/product-detail/GLOBAL\_ MULTIYEAR\_PHY\_001\_030/ (accessed on 6 September 2021). The global altimeter satellite gridded geostrophic velocities anomalies are available at: https://resources.marine.copernicus.eu/productdetail/SEALEVEL\_GLO\_PHY\_L4\_MY\_008\_047/ (accessed on 22 July 2020).

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