

Comment on “Proton Spin Structure from Measurable Parton Distributions”

Some time ago, Ji [1–3], using the Belinfante version of the angular momentum operator, derived a beautiful relation between the quark angular momentum and generalized parton distributions. In these Letters, the relation was written as

$$J_q = \frac{1}{2} \int_{-1}^1 dx x [H_q(x, 0, 0) + E_q(x, 0, 0)], \quad (1)$$

and for a decade and a half the quantity J_q has been almost universally interpreted as the expectation value of the longitudinal component of the quark angular momentum in a longitudinally polarized nucleon, i.e., for a nucleon moving along the z direction

$$J_q = \langle\langle J_q^z \rangle\rangle_L. \quad (2)$$

Inspired by the impact-parameter explanation of J_q proposed by Burkardt [4], Ji, Xiong, and Yuan [5] show that a partonic interpretation of the right-hand side of Eq. (1) can be obtained and state that J_q measures the expectation value of the transverse angular momentum of the quarks in a nucleon polarized in the transverse direction i . What they claim to prove is that

$$J_q \propto \langle\langle J_q^{+i} \rangle\rangle_{T_i}, \quad (3)$$

where

$$J_q^{+i} = \int dx^- d^2 x_\perp M_q^{+i}(x) \quad (4)$$

with, in terms of the Belinfante energy-momentum tensor density,

$$M_q^{\mu\rho\sigma}(x) = x^\rho T_q^{\mu\sigma}(x) - x^\sigma T_q^{\mu\rho}(x). \quad (5)$$

Since J^{+i} is a leading-twist operator, it is clear that such a simple partonic interpretation should exist. However, Ji, Xiong, and Yuan misleadingly interpret J^{+i} as the transverse *angular momentum* operator. Indeed, in light-front quantization, the role of time is taken by x^+ , so that J^{+i} is the light-front transverse *boost* operator. In terms of the more conventional instant-form boost (K^i) and rotation (J^i) operators (see, e.g., Refs. [6–8]), the light-front transverse boost operators read

$$J^{+1} = \frac{1}{\sqrt{2}}(K^1 + J^2), \quad J^{+2} = \frac{1}{\sqrt{2}}(K^2 - J^1), \quad (6)$$

while the light-front transverse angular momentum operators are given by

$$J^{-1} = \frac{1}{\sqrt{2}}(K^1 - J^2), \quad J^{-2} = \frac{1}{\sqrt{2}}(K^2 + J^1). \quad (7)$$

The light-front transverse boosts (J^{+i}) are kinematic operators and therefore leading twist, while the light-front transverse angular momenta (J^{-i}) are dynamical operators

and therefore higher twist. It is also easy to see that the quark and gluon spin operators in the $A^+ = 0$ gauge [9]

$$M_{q,\text{spin}}^{\mu\rho\sigma} = \frac{1}{2} \epsilon^{\mu\rho\sigma\nu} \bar{\psi} \gamma_\nu \gamma_5 \psi, \quad (8)$$

$$M_{g,\text{spin}}^{\mu\rho\sigma} = -2 \text{Tr}[F^{\mu\rho} A^\sigma - F^{\mu\sigma} A^\rho] \quad (9)$$

contribute to J^{-i} and not to J^{+i} .

Any genuine transverse angular momentum sum rule is expected to have a frame dependence. The reason is simply the well-known fact that boosts and rotations do not commute. One consequence of this is that special relativity naturally induces spin-orbit correlations. Obviously, there cannot be any spin-orbit correlation with the longitudinal polarization, which is the reason why the longitudinal angular momentum sum rule is frame independent. On the contrary, the transverse polarization is correlated with the momentum, which is at the origin of the frame dependence of the transverse angular momentum sum rule.

In conclusion, the Ji, Xiong, and Yuan partonic interpretation has nothing to do with angular momentum. One cannot simply interpret $(x/2)[H_q(x, 0, 0) + E_q(x, 0, 0)]$ as the density of quark transverse angular momentum in a transversely polarized nucleon. A genuine transverse angular momentum sum rule naturally involves frame dependence, owing to the fact that boosts and rotations do not commute. Since the transverse angular momentum is a dynamical operator in light-front quantization, no simple partonic interpretation is expected. On the contrary, a simple partonic interpretation does exist for the longitudinal component of angular momentum in terms of Wigner distributions [10–16], precisely because it is a kinematic operator.

Finally, one of us (E. L.) [17] recently derived a relation for the instant-form transverse component of the quark angular momentum in a transversely polarized nucleon in terms of the generalized parton distributions H and E , which is frame dependent, and in passing, checked that Eq. (1) is indeed correct for the longitudinal case, with the identification $J_q = \langle\langle J_q^z \rangle\rangle_L$.

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