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Comparison of $D \rightarrow K_S^0 \pi$ and $D \rightarrow K_L^0 \pi$ Decay Rates

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We present measurements of $D \rightarrow K_S^0 \pi$ and $D \rightarrow K_L^0 \pi$ branching fractions using 281 pb⁻¹ of $\psi(3770)$ data at the CLEO-*c* experiment. We find that $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$ is larger than $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0)$, with an asymmetry of $R(D^0) = 0.108 \pm 0.025 \pm 0.024$. For $\mathcal{B}(D^+ \rightarrow K_S^0 \pi^+)$ and $\mathcal{B}(D^+ \rightarrow K_L^0 \pi^+)$, we observe no measurable difference; the asymmetry is $R(D^+) = 0.022 \pm 0.016 \pm 0.018$. The D^0 asymmetry is consistent with the value based on the U-spin prediction $A(D^0 \rightarrow K^0 \pi^0)/A(D^0 \rightarrow \bar{K}^0 \pi^0) = -\tan^2 \theta_C$, where θ_C is the Cabibbo angle.

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As the dominant decay of the charm quark is to the strange quark, the final states of D^0 and D^+ meson decays typically include K^+ , K^- , K_S^0 , and K_L^0 mesons. While there have been many measurements of D decays to final states

containing K^\pm and K_S^0 , until now there have been no measurements of decays to final states containing a K_L^0 . Typically it has been assumed that the branching fraction for a decay $D \rightarrow K_L^0 X$ will equal that for the corresponding

decay $D \rightarrow K_S^0 X$. However, as pointed out by Bigi and Yamamoto [1], interference between Cabibbo-favored transitions (producing an s quark, and thus a \bar{K}^0) and doubly-Cabibbo-suppressed transitions (producing an \bar{s} quark, and thus a K^0) can lead to a difference in the rates for $D \rightarrow K_L^0 X$ and $D \rightarrow K_S^0 X$. Here we present first measurements of two D decays to K_L^0 , $D^0 \rightarrow K_L^0 \pi^0$, and $D^+ \rightarrow K_L^0 \pi^+$, and we compare the branching fractions with those for $D^0 \rightarrow K_S^0 \pi^0$ and $D^+ \rightarrow K_S^0 \pi^+$. (Throughout, charge-conjugate modes are implied, except where noted.) These comparisons provide information about amplitudes and strong phases in $D \rightarrow K\pi$ decays.

For these measurements we use a 281 pb $^{-1}$ sample of $e^+e^- \rightarrow \psi(3770)$ events, produced by the CESR- c storage ring and recorded with the CLEO- c detector. The CLEO- c detector is a general purpose solenoidal detector which includes a tracking system for measuring momentum and specific ionization (dE/dx) of charged particles, a Ring Imaging Cherenkov detector (RICH) to aid in particle identification, and a CsI calorimeter for detection of electromagnetic showers. The CLEO- c detector is described in detail elsewhere [2–4].

The $\psi(3770)$ resonance is below the threshold for $D\bar{D}\pi$, and so the events of interest, $e^+e^- \rightarrow \psi(3770) \rightarrow D\bar{D}$, have D mesons with energy equal to the beam energy and a unique momentum. Thus, for identifying D^0 and D^+ candidates, we follow Mark III [5] and define the two variables ΔE and beam-constrained mass M_{BC} by:

$$\Delta E \equiv \sum_i E_i - E_{\text{beam}}, \quad M_{\text{BC}} \equiv \sqrt{E_{\text{beam}}^2 - \left| \sum_i \mathbf{p}_i \right|^2},$$

where E_i and \mathbf{p}_i are the energies and momenta of the D decay products. For true D candidates, ΔE will be consistent with zero, and M_{BC} will be consistent with the D mass.

We measure the branching fractions for the decays $D \rightarrow K_S^0 \pi$ by directly reconstructing the final-state particles, where the K_S^0 is reconstructed from $K_S^0 \rightarrow \pi^+ \pi^-$. The decay $D^+ \rightarrow K_S^0 \pi^+$ is measured by a separate CLEO- c analysis [6], but CLEO- c has not previously measured $D^0 \rightarrow K_S^0 \pi^0$. In this Letter, we cite the $D^+ \rightarrow K_S^0 \pi^+$ result and present a measurement of $D^0 \rightarrow K_S^0 \pi^0$.

The decays $D \rightarrow K_L^0 \pi$ have not been previously measured due to the difficulty of K_L^0 reconstruction. Fortunately, the clean $D\bar{D}$ environment allows us to measure these decays without directly detecting the K_L^0 . Instead, we reconstruct all particles in the event except for the K_L^0 —that is, a tag \bar{D} and a π —and infer the presence of a K_L^0 from the missing four-momentum. Our signal is a peak in the missing mass squared distribution at the K_L^0 mass squared.

The situation for D^0 decays has an added complication. When D^0 and \bar{D}^0 are pair-produced through a virtual photon ($J^{\text{PC}} = 1^{--}$), they are in a quantum coherent state. Therefore, the decays of D^0 and \bar{D}^0 are subject to interfer-

ence. This interference has no effect on the overall rate for any particular D^0 or \bar{D}^0 decay, but it does alter how often a particular D^0 decay occurs in combination with a particular \bar{D}^0 decay. Therefore, when D^0 decays are measured with a reconstructed tag \bar{D}^0 , the apparent “branching fractions” of the D^0 will vary according to the decay of the \bar{D}^0 [7]. This effect is especially large for CP eigenstate modes, such as $K_S^0 \pi^0$ and $K_L^0 \pi^0$.

The quantum correlation effects are shown in Table I, where X stands for all modes combined, f stands for a flavored mode, S_+ stands for a CP -even mode, S_- stands for a CP -odd mode, $R_{\text{WS},f}$ is the wrong-sign decay ratio $\mathcal{B}(\bar{D} \rightarrow f)/\mathcal{B}(D \rightarrow f)$, y is the $D^0 - \bar{D}^0$ mixing parameter $y \equiv \Delta\Gamma/2\Gamma$, and $r_f e^{-i\delta_f} \equiv \langle f|\bar{D}^0\rangle/\langle f|D^0\rangle$. An untagged measurement is not altered relative to a measurement using an isolated D^0 . However, measurements of $K_S^0 \pi^0$ (CP -odd) and $K_L^0 \pi^0$ (CP -even), tagged by a flavored \bar{D}^0 decay, are altered by factors of $(1 + R_{\text{WS},f} \mp r_f z_f \mp y)$. These factors depend on the tag mode, and $z_f \equiv 2 \cos\delta_f$ is generally not known. Since $D^0 \rightarrow K_L^0 \pi^0$ must be reconstructed with a tag \bar{D}^0 , we must determine the factor for each tag mode, f . We do this by comparing tagged and untagged $D^0 \rightarrow K_S^0 \pi^0$.

Our procedure is the following: we first measure the branching fraction $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$ by reconstructing this decay without tagging a \bar{D}^0 . Next, we measure the “branching fraction” for $D^0 \rightarrow K_S^0 \pi^0$, with three different flavor tags. Each gives us $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)(1 - C_f)$, where $C_f \equiv (r_f z_f + y)/(1 + R_{\text{WS},f})$. Using $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$ from the untagged measurement, we obtain C_f for each flavor tag. Finally, we measure the “branching fraction” for $D^0 \rightarrow K_L^0 \pi^0$, with the same three flavor tags. Each gives us $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0)(1 + C_f)$. Using the calculated values of C_f , we obtain $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0)$ from each of the three tags. These measurements are then averaged for the final result.

We first measure $D^0 \rightarrow K_S^0 \pi^0$ without searching for a tag \bar{D}^0 . Candidates for $D^0 \rightarrow K_S^0 \pi^0$ are formed by combining a K_S^0 , reconstructed by a pair of charged tracks through the decay $K_S^0 \rightarrow \pi^+ \pi^-$, and a π^0 , reconstructed from a pair of photons detected in the CsI calorimeter. The invariant mass of K_S^0 and π^0 candidates is required to be con-

TABLE I. Untagged (vs X) and tagged (vs f) efficiency-corrected yields for $C = -1$ $D^0\bar{D}^0$ events, to leading order in the mixing parameters. N is the number of $D^0\bar{D}^0$ events, B_i is the branching fraction for mode i for an isolated D^0 , and $z_f \equiv 2 \cos\delta_f$.

	X	f
S_+	$2NB_{S_+}$	$NB_f B_{S_+} (1 + R_{\text{WS},f} + r_f z_f + y)$
S_-	$2NB_{S_-}$	$NB_f B_{S_-} (1 + R_{\text{WS},f} - r_f z_f - y)$
X	-	$2NB_f (1 + R_{\text{WS},f})$

sistent with the known mass, and π^0 candidates are then constrained to the known mass.

Both beam-constrained mass and ΔE are required to be within 3 standard deviations of the nominal value. If there are multiple candidates in one event, we accept only the one whose beam-constrained mass is closest to the nominal D^0 mass. Two sideband subtractions are used to remove background. First, a ΔE sideband subtraction is used to remove the continuum and combinatoric background (a 13% effect). Then, a K_S^0 mass sideband subtraction is used to remove the background from $D^0 \rightarrow \pi^+ \pi^- \pi^0$ events in which $M(\pi^+ \pi^-)$ happens to be within the K_S^0 mass window (a 4% effect). The resulting yield is 7487 ± 101 events. This yield is divided by the detection efficiency, 29.3%, to determine the number of $D^0 \rightarrow K_S^0 \pi^0$ events produced. The efficiency is determined from Monte Carlo simulation, with a correction for π^0 detection efficiency; this correction is determined by comparing π^0 efficiencies measured in data and in our simulation. Finally, we use the total number of $D^0 \bar{D}^0$ events in our sample, 1.031×10^6 (from a separate CLEO-*c* analysis [6]). Dividing the efficiency-corrected yield by twice this number gives the branching fraction.

Systematic uncertainties considered include those from: the ΔE cut ($\pm 0.5\%$), the ΔE sideband subtraction ($\pm 0.8\%$), tracking efficiency ($\pm 0.6\%$), K_S^0 detection efficiency ($\pm 1.8\%$), the K_S^0 sideband subtraction ($\pm 0.3\%$), and the number of $D^0 \bar{D}^0$ events ($\pm 1.4\%$). These total $\pm 2.5\%$. The largest uncertainty is due to π^0 reconstruction efficiency ($\pm 3.8\%$). Although this uncertainty is large, it cancels in the computation of quantum correlation factors and in the comparison of the $D^0 \rightarrow K_S^0 \pi^0$ and $D^0 \rightarrow K_L^0 \pi^0$ branching fractions. Therefore, we keep it separate from the other uncertainties.

We find a branching fraction $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0) = (1.240 \pm 0.017 \pm 0.031 \pm 0.047)\%$, where the last uncertainty is from the π^0 efficiency.

Having determined $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$, we now measure this decay with three different tag modes to obtain the quantum correlation factors $(1 - C_f)$. The three tag modes we use are $\bar{D}^0 \rightarrow K^+ \pi^-$, $\bar{D}^0 \rightarrow K^+ \pi^- \pi^0$, and $\bar{D}^0 \rightarrow K^+ \pi^- \pi^- \pi^+$. The tag \bar{D}^0 is required to be within 3 standard deviations of the nominal values of ΔE and M_{BC} . We select at most one candidate per flavor per tag mode; when multiple candidates pass our requirements, we keep the one with M_{BC} closest to the nominal D^0 mass. We remove fake tag \bar{D}^0 candidates by subtracting the ΔE sideband of the tag.

In the tagged sample, we reconstruct $D^0 \rightarrow K_S^0 \pi^0$ in the same way as in the untagged case. To remove fake $D^0 \rightarrow K_S^0 \pi^0$ candidates, we subtract a K_S^0 mass sideband. (No ΔE sideband subtraction is necessary since, with a tag, the $K_S^0 \pi^0$ signal is essentially free of combinatoric background.)

Although to first order the efficiency of reconstructing the \bar{D}^0 tag cancels in the branching fraction calculation, simulations indicate a slightly larger efficiency for \bar{D}^0 reconstruction when the signal D^0 decays to $K_S^0 \pi^0$. This bias stems from the lower-than-average multiplicity of particles in $D^0 \rightarrow K_S^0 \pi^0$ events. We obtain correction factors for these small biases from Monte Carlo studies.

With the efficiencies from Monte Carlo simulations and the yields in signal and sideband regions, we compute the branching fractions, times quantum correlation factors, in Table II.

The systematic uncertainties are similar to those in the untagged measurement. Track, K_S^0 , and π^0 reconstruction uncertainties are the same, and they will cancel in the ratio of the tagged and untagged results. The only systematic uncertainties from the tag \bar{D}^0 are for the ΔE sideband subtraction and the tag bias correction factor; any other discrepancies in the Monte Carlo simulation would have the same effect on the tag and signal yields.

Finally, we divide these results by $\mathcal{B}(D^0 \rightarrow K_S^0 \pi^0)$, from the untagged measurement, to obtain the three quantum correlation factors $(1 - C_f)$, where f represents the tag mode.

We measure the $D \rightarrow K_L^0 \pi$ branching fractions with a missing mass technique. We reconstruct the tag \bar{D} in 3 \bar{D}^0 modes and 6 D^- modes, and we combine it with a π^0 or π^+ to form missing mass squared: $M_{\text{miss}}^2 \equiv (p_{\text{event}} - p_{\bar{D}} - p_{\pi})^2$. To improve resolution, the tag \bar{D} is constrained to have the expected three-momentum magnitude. The $D \rightarrow K_L^0 \pi$ signal is a peak in M_{miss}^2 at the K_L^0 mass squared ($\sim 0.25 \text{ GeV}^2$).

To remove $D \rightarrow K_S^0 \pi$ events, as well as other backgrounds, we require that the event contain no extra tracks or π^0 's beyond those used in the tag \bar{D} and the π . This veto removes about 90% of $D \rightarrow K_S^0 \pi$ events and a few percent of $D \rightarrow K_L^0 \pi$ events. For $D^0 \rightarrow K_L^0 \pi^0$ only, we also remove an event if it contains an extra $\eta \rightarrow \gamma\gamma$. This removes much of the $D^0 \rightarrow \eta \pi^0$ background. To determine systematic uncertainties from the appearance of fake extra particles in signal events, we compare how often they appear in data and in our simulation, using events in which both D and \bar{D} were fully reconstructed.

TABLE II. Efficiencies, yields, and results for tagged $D^0 \rightarrow K_S^0 \pi^0$ study. No systematic uncertainties are included in the quoted results.

Tag mode	$K^+ \pi^-$	$K^+ \pi^- \pi^0$	$K^+ \pi^- \pi^- \pi^+$
Efficiency	31.74%	31.29%	29.97%
Tag yield—raw	48 095	67 576	75 113
Sideband subtracted	47 440	63 913	71 040
Signal yield—raw	172	248	276
Sideband subtracted	155	203	256
Tag bias correction	1.000	1.014	1.033
$\mathcal{B}(K_S^0 \pi^0)(1 - C_f)$ (%)	1.03 ± 0.09	1.00 ± 0.09	1.16 ± 0.08

As in the tagged $D^0 \rightarrow K_S^0 \pi^0$ study, the tag \bar{D} reconstruction efficiency is higher when the D decays to $K_L^0 \pi$; therefore, we apply correction factors determined from Monte Carlo simulations. The efficiency for observing $D \rightarrow K_L^0 \pi$, given that the tag was found, is also determined in these simulations. It is essentially the efficiency for finding the π without any fake extra particles.

For the $D^0 \rightarrow K_L^0 \pi^0$ branching fraction measurement, the same three \bar{D}^0 decay modes are selected with the same requirements as in the tagged $D^0 \rightarrow K_S^0 \pi^0$ study (except for a minor difference in the order of applying cuts for the $K^+ \pi^- \pi^0$ tag, which results in a slight difference in number of tags). Combining these \bar{D}^0 candidates with π^0 candidates and rejecting events with extra tracks, π^0 's, or η 's, we obtain the M_{miss}^2 plot shown in Fig. 1.

A number of backgrounds slip through our extra track, π^0 , and η vetoes and appear in the M_{miss}^2 plot. The modes $K_S^0 \pi^0$ and $\eta \pi^0$ appear as peaks at essentially the same location as $K_L^0 \pi^0$, $\pi^0 \pi^0$ peaks at $M_{\text{miss}}^2 \approx 0.0$ GeV², and $K^{*0} \pi^0$ peaks at 0.8 GeV². Monte Carlo simulations of these backgrounds are shown in Fig. 1. Other, lesser backgrounds also appear to the right of the $K_L^0 \pi^0$ peak.

To determine the signal and estimate the background, we define a M_{miss}^2 signal region 0.1 to 0.5 GeV², as well as low and high sidebands: -0.1 to 0.1 GeV² and 0.8 to 1.2 GeV². The backgrounds are split into three groups: $D^0 \rightarrow K_S^0 \pi^0$ and $D^0 \rightarrow \eta \pi^0$, $D^0 \rightarrow \pi^0 \pi^0$, and all other backgrounds. For $D^0 \rightarrow K_S^0 \pi^0$ and $D^0 \rightarrow \eta \pi^0$, we use Monte Carlo simulation to determine efficiencies for the

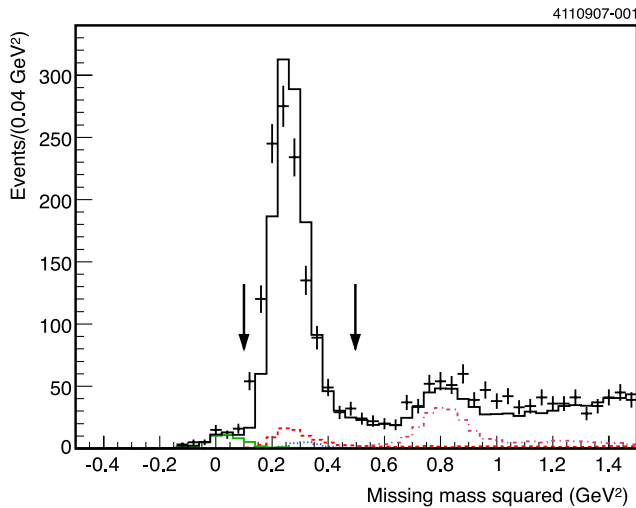


FIG. 1 (color online). Missing mass squared distribution, with all tag modes combined, for $D^0 \rightarrow X \pi^0$, after removing events with extra tracks, π^0 's, or η 's. The points with error bars are data, and the solid line is a Monte Carlo simulation. The dashed, colored lines represent simulations of the peaking backgrounds $D^0 \rightarrow \pi^0 \pi^0$, $K_S^0 \pi^0$, $\eta \pi^0$, and $K^{*0} \pi^0$. The difference in the peak position is due to a minor discrepancy in our calorimeter simulation at large photon energies; the signal region, marked with arrows, encompasses the peak in both distributions.

background subtraction. For $D^0 \rightarrow \pi^0 \pi^0$, we scale the contribution to the signal region according to the yield in the low sideband. For the sum of all other backgrounds, we follow the same procedure with the high sideband. In total, about 10% of the events in the signal region are background, with half coming from $K_S^0 \pi^0$, 1/10 from each of $\eta \pi^0$ and $\pi^0 \pi^0$, and 3/10 from various other decays.

After subtracting all the backgrounds, we obtain the yields and compute branching fractions, times quantum correlation factors, in Table III.

Systematic uncertainties come from the effect on signal efficiency of the veto on extra tracks ($\pm 0.3\%$), the veto on extra π^0 's ($\pm 1.6\%$), the veto on η 's ($\pm 0.5\%$), and the uncertainty in the location and width of the signal peak ($\pm 1.4\%$). Other uncertainties come from the background estimate ($\pm 1.0\%$), ΔE sideband subtraction ($\pm 0.5\%$), and the tag bias correction factor ($\pm 0.2\%$). These total $\pm 2.5\%$. As in $D^0 \rightarrow K_S^0 \pi^0$, π^0 efficiency ($\pm 3.8\%$) is the largest systematic uncertainty; it cancels in the comparison of $D^0 \rightarrow K_S^0 \pi^0$ and $D^0 \rightarrow K_L^0 \pi^0$.

We have determined $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0)(1 + C_f)$ for three different flavor tags f . Using the values of C_f determined from the $D^0 \rightarrow K_S^0 \pi^0$ measurements, we calculate $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0)$ for each tag mode. Finally, we average the results and find $\mathcal{B}(D^0 \rightarrow K_L^0 \pi^0) = (0.998 \pm 0.049 \pm 0.030 \pm 0.038)\%$, where the last uncertainty is from the π^0 efficiency.

The analysis of $D^+ \rightarrow K_L^0 \pi^+$ is similar to $D^0 \rightarrow K_L^0 \pi^0$, though there are a few differences. Since we reconstruct a π^+ instead of a π^0 , the M_{miss}^2 resolution is better. Also, we do not need to correct for quantum correlation. The most significant difference in procedure is that we perform a likelihood fit for the signal and background yields instead of counting events in a signal region.

We reconstruct tag D^- 's in six decay modes: $D^- \rightarrow K^+ \pi^- \pi^-$, $K^+ \pi^- \pi^- \pi^0$, $K_S^0 \pi^-$, $K_S^0 \pi^- \pi^0$, $K_S^0 \pi^- \pi^- \pi^+$, and $K^+ K^- \pi^-$. As before, candidates must have ΔE consistent with zero. We select one candidate per charge per mode based on the best value of ΔE . We fit the M_{BC} distribution for each mode to determine the number of tags, and then pass all candidates with M_{BC} near the peak to be combined with π^+ candidates.

TABLE III. Efficiencies, yields, and results for $D^0 \rightarrow K_L^0 \pi^0$. No systematic uncertainties are included in the quoted results.

Tag mode	$K^+ \pi^-$	$K^+ \pi^- \pi^0$	$K^+ \pi^- \pi^- \pi^+$
Efficiency	55.21%	52.72%	49.88%
Tag yield—raw	48 095	68 000	75 113
Sideband subtracted	47 440	64 280	71 040
Signal yield—raw	367.0	414.5	466.5
Background subtracted	334.8	363.1	418.0
Tag bias correction	1.000	1.037	1.057
$\mathcal{B}(K_L^0 \pi^0)(1 + C_f)$ (%)	1.28 ± 0.08	1.03 ± 0.06	1.12 ± 0.06

The M_{miss}^2 distribution, with all tag modes added together, is shown in Fig. 2. The lines show a fit used to determine the signal yield. The most prominent feature is the signal peak at the K_L^0 mass squared ($\sim 0.25 \text{ GeV}^2$). A number of backgrounds are also present. First, fake D^- candidates produce a background which is estimated from an M_{BC} sideband. All of the other backgrounds come from other D^+ decays. The largest of these are $D^+ \rightarrow K_S^0 \pi^+$ (dashed, green peak under the signal), $\eta \pi^+$ (shoulder on the right side tail of the signal), $\pi^0 \pi^+$ and $\mu^+ \nu_\mu$ (peak on the left of the plot), $\bar{K}^0 \pi^+ \pi^0$, and $\pi^+ \pi^0 \pi^0$. The shapes and efficiencies of these backgrounds are determined from Monte Carlo simulations. The yields of the signal peak and the $\eta \pi^+$, $\pi^0 \pi^+$, and $\mu^+ \nu_\mu$ backgrounds are allowed to vary in the fit; all other yields are fixed based on the efficiencies.

Although Fig. 2 shows all tag modes together, we actually fit each tag mode separately. We calculate a branching fraction from each tag mode using the tag bias correction factor, efficiency, tag D^- yield, and signal $D^+ \rightarrow K_L^0 \pi^+$ yield for that mode. The tag bias correction varies from 1.005 (for $K^+ \pi^- \pi^-$) to 1.047 (for $K_S^0 \pi^- \pi^- \pi^+$). The efficiency averages to 81.6%, and depends little on tag mode. There are a total of 165×10^3 tags, and a total $D^+ \rightarrow K_L^0 \pi^+$ yield of 2023 ± 54 events. The values of the branching fraction calculated from each tag mode are averaged to produce the final result.

Systematic uncertainties include those from: pion reconstruction efficiency ($\pm 0.3\%$) and particle identification ($\pm 0.25\%$), tag bias correction factor ($\pm 0.2\%$), charge of the tag D ($\pm 0.5\%$), extra track and extra π^0 vetoes ($\pm 1.1\%$), signal peak shape ($\pm 0.7\%$), signal peak width

($\pm 1.6\%$), contribution of fake D^- tags ($\pm 0.4\%$), and M_{miss}^2 background yields [$\pm 0.8\%$ from statistical uncertainty in $K_S^0 \pi^+$ background, $\pm 0.3\%$ from $\mathcal{B}(D^+ \rightarrow K_S^0 \pi^+)$, and $\pm 0.5\%$ from all other backgrounds]. The total systematic uncertainty is $\pm 2.4\%$.

We find a branching fraction $\mathcal{B}(D^+ \rightarrow K_L^0 \pi^+) = (1.460 \pm 0.040 \pm 0.035 \pm 0.005)\%$. The final uncertainty is due to the input value of $\mathcal{B}(D^+ \rightarrow K_S^0 \pi^+)$.

To compare $D \rightarrow K_S^0 \pi$ and $D \rightarrow K_L^0 \pi$, we compute the asymmetries

$$R(D) \equiv \frac{\mathcal{B}(D \rightarrow K_S^0 \pi) - \mathcal{B}(D \rightarrow K_L^0 \pi)}{\mathcal{B}(D \rightarrow K_S^0 \pi) + \mathcal{B}(D \rightarrow K_L^0 \pi)}.$$

The D^0 asymmetry (in which the systematic uncertainty for π^0 efficiency cancels) is $R(D^0) = 0.108 \pm 0.025 \pm 0.024$. Using $\mathcal{B}(D^+ \rightarrow K_S^0 \pi^+) = (1.526 \pm 0.022 \pm 0.038)\%$ [6], the D^+ asymmetry is $R(D^+) = 0.022 \pm 0.016 \pm 0.018$.

The asymmetry between $D^0 \rightarrow K_S^0 \pi^0$ and $D^0 \rightarrow K_L^0 \pi^0$ is consistent with SU(3) symmetry, and, in particular, the U-spin subgroup of SU(3). U-spin predicts $A(D^0 \rightarrow K^0 \pi^0)/A(D^0 \rightarrow \bar{K}^0 \pi^0) = -\tan^2 \theta_C$, where θ_C is the Cabibbo angle. This prediction is relatively insensitive to SU(3) breaking [8]. The amplitude ratio can also be predicted from diagrams for these two processes; both have spectator and exchange diagrams which differ only by a factor of $-\tan^2 \theta_C$. However derived, the amplitude ratio implies that the asymmetry is $R(D^0) = 2 \tan^2 \theta_C$. Using $\tan \theta_C = 0.233 \pm 0.001$ [9], we calculate $R(D^0) = 0.109 \pm 0.001$, in good agreement with our measurement.

There is no corresponding U-spin argument for the D^+ decays, so no simple prediction is possible. Diagrams for the Cabibbo-favored and doubly-suppressed decays are different. Both internal and external spectator diagrams contribute to $D^+ \rightarrow \bar{K}^0 \pi^+$, while $D^+ \rightarrow K^0 \pi^+$ has internal spectator and annihilation diagrams. Approximate predictions are, however, possible under certain assumptions. One analysis [10], based on flavor SU(3) with an estimate of symmetry-breaking effects, finds $R(D^+) \approx 0.04$, consistent with our measurement. This analysis also points out that the small asymmetry found for D^+ decays can be interpreted as a large strong phase between two contributing amplitudes in the case of D^+ decays, while the larger asymmetry in the D^0 decays is consistent with a small strong phase.

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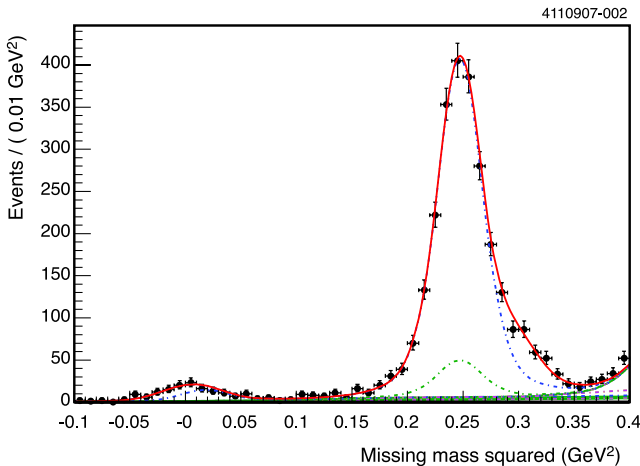


FIG. 2 (color online). Missing mass squared distribution, with all tag modes combined, for $D^+ \rightarrow X \pi^+$, after removing events with extra tracks or π^0 's. The solid line shows a fit for the $D^+ \rightarrow K_L^0 \pi^+$ yield. The many dashed lines represent the various components of the fit, added cumulatively. The small peak under the signal is the contribution of $D^+ \rightarrow K_S^0 \pi^+$ events that are not removed by the extra track and π^0 vetoes.

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