

## Confinement of a Large Number of Antiprotons and Production of an Ultraslow Antiproton Beam

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We have used a radio frequency quadrupole decelerator to decelerate antiprotons emerging from the CERN Antiproton Decelerator from MeV- to keV-scale energy, and collected five decelerated pulses in a multiring trap. Some  $5 \times 10^6$  antiprotons were stacked in this way. Cooling of the trapped antiprotons by a simultaneously trapped electron plasma was studied nondestructively via shifts in plasma mode frequencies. We have also demonstrated the first step in extracting a 10–500 eV antiproton beam from the trap.

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The preparation of large samples of antiprotons at extremely low energy is an important step on the road to synthesizing antihydrogen ( $\bar{p}e^+$ ) and protonium ( $\bar{p}p$ ) atoms. Not surprisingly, such exotic atoms can only be efficiently synthesized from component particles at the “chemical” energy scale (eV and lower), and this is far below the mandatory GeV scale of accelerator-produced antiprotons.

Since 2000, this energy gap has been partially bridged by the CERN Antiproton Decelerator (AD), which decelerates antiprotons with  $\sim$ GeV-scale kinetic energy to 5.3 MeV, and reejects them every  $\sim$ 2 min in 90 ns long pulses, containing some  $3 \times 10^7$  particles. This beam energy is, however, still far above the 10 keV range at which antiprotons can readily be captured and cooled in electromagnetic traps [1–3]. Two groups working at the AD [4–7] have nevertheless demonstrated antihydrogen synthesis in such traps. In these experiments,  $\sim$ 70 000  $\mu\text{g}/\text{cm}^2$  degrader foils reduced the antiproton energy further from the MeV scale to the keV scale [1–3]. Many incident antiprotons stop and annihilate within such relatively thick foils, while others emerge with too high an energy to be confined in the potential well of any trap located downstream. Indeed, the largest number of captured particles so far reported for a typical AD pulse, or “shot,” is  $\sim$ 2.5  $\times$  10<sup>4</sup> [8]. We describe here the replacement of these simple degrader foils by a radio frequency quadrupole decelerator (RFQD) [9,10]. This substantial decelerator, 4 m in length, reduced the 5.3 MeV AD beam energy between 10 and 120 keV. Some (5–9)  $\times$  10<sup>6</sup> antiprotons within a single shot fulfilled the phase condition necessary for acceptance and deceleration by the RFQD, the remainder being transmitted without deceleration [9,10]. We expected that the reduced antiproton energy would allow a much larger number of the emerging antiprotons to be captured in a trap placed after the RFQD than is normally the case for

simple degrader foils. For these experiments we used a multiring trap (MRT) comprised of several ring electrodes lying perpendicular to the beam and on the axis of a superconducting solenoid at 10 K, as is shown in Fig. 1.

Since the antiprotons were to spend a long time in this trap, more stringent vacuum conditions were required there to avoid antiproton loss by annihilation on residual gas atoms [3] than in the RFQD, through which they passed only once. Two 90  $\mu\text{g}/\text{cm}^2$  polyethyleneterephthalat (PET) foils located 70 cm upstream of the MRT were therefore used to isolate the trap and RFQD vacuum systems, at  $\leq$  10<sup>-12</sup> mbar and  $\sim$ 10<sup>-9</sup> mbar, respectively. Although these were some 800 times thinner than those used at the entrance of the traps described in Refs. [1–3], antiprotons traversing them unavoidably suffered a finite mean energy loss ( $\sim$  100 keV). The corresponding angular divergence and straggling cause some 30% of the antiprotons to annihilate on the way through, but this is to be compared with similar losses in [1–3] totaling 99.9%, for energy loss 5.3 MeV. The potential of the entire RFQD could be floated by several tens of kilovolts relative to ground, and we tuned the value of this floating voltage so as to maximize the capture efficiency of the (grounded) MRT. This occurred at an RFQD beam output energy of 110 keV.

After passing through the first (upstream) ring, initially biased at ground potential, the antiproton pulse proceeded along the MRT axis to the position of the last one, which being biased at  $-10$  kV, returned it to the first electrode. The trap was closed around the pulse by a switch which changed the bias of this upstream electrode to  $-10$  kV before the first antiprotons in the pulse reached it on their return trip. By this time, most of the 90 ns long pulse had entered the trap.

Once inside the MRT, the antiprotons were to be decelerated and cooled for extraction as an ultraslow antiproton beam for use in future experiments (see, e.g., [11–13]).

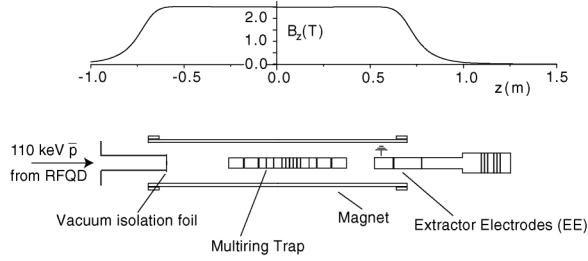


FIG. 1. The MRT and its extractor electrode; the calculated magnetic field distribution  $B_z$  is also shown.

Before the antiprotons arrived, we therefore loaded the trap with a non-neutral plasma of  $\sim 3 \times 10^8$  electrons and allowed this to reach equilibrium, where its shape was spheroidal. This plasma cooled the newly arrived antiprotons continuously, since the thermal energy it received from the antiprotons was dissipated spontaneously by synchrotron radiation in the strong magnetic field with a time constant of about 1 s [3,14]. The MRT design enables a large-volume harmonic potential to be established, within which plasmas containing large numbers of charged particles can be stably confined for long periods [15]. Observing the joint plasma modes then allowed us to monitor the plasma shape, density, and temperature in real time [16].

About  $\sim 1.2 \times 10^6$  among the captured antiprotons could be stored in the MRT for periods of 10 min or more, although they were usually ejected earlier. This figure is  $\sim 50$  times higher than the previous best values obtained with thick degrader foils [8]. With a 30% RFQD transmission efficiency for the decelerated particles [17] and the aforementioned 70% efficiency of transmission through the  $90 \mu\text{g}/\text{cm}^2$  PET foils, the overall antiproton trapping efficiency, defined as the number captured versus the number in the AD shot, was  $\sim 4\%$ . By stacking five AD shots in sequence [8],  $\approx 5 \times 10^6$  antiprotons, the largest number of antiprotons ever stored and cooled, were confined in our MRT.

For diagnostic purposes it was important to know the position and time distribution of antiproton annihilations during the capture, storage, and cooling of the AD pulses. Annihilation detectors consisting of two 2 m long plastic scintillator bars were used for this purpose. They were placed parallel to the trap axis, at distances of 40 and 74 cm from it, and had cross sections of  $4(\text{H}) \times 6(\text{V}) \text{cm}^2$ , thus subtending overlapping solid angles of 0.12 sr seen from the center of the MRT trap axis. Coincident light pulses produced in these scintillators by the passage of charged particles from nearby antiproton annihilations were recorded by photomultiplier tubes at the ends of the bars, and their arrival-time difference permitted the annihilation vertices to be determined axially with a precision of  $\approx 50$  mm. We performed particle-tracking simulations using the GEANT library [18], from which we estimate the overall detection efficiency to be  $\approx 5\%$  at the center of the trapping region. With the MRT residual gas

pressure of  $\approx 10^{-12}$  mbar, most of the annihilations could safely be assumed to occur on the trap cylindrical electrodes (at radius 20 mm) or on other adjacent solid materials, such as the vacuum isolation foil.

We excited axially symmetric electrostatic oscillation ( $l, 0$ ) modes ( $l = 1-4$ ) of the joint plasma by applying white noise to it via one of the trap electrodes. These oscillations induced peaks in the power spectrum at the corresponding mode frequencies, and the shift of these peaks with time was used to monitor the antiproton cooling process. The charge distributions of these modes were axially symmetric and had zero, one, and two nodes along the axial direction, respectively.

Figure 2(a) shows an example of the time evolution of the power spectrum. Resonances corresponding to the (1, 0), (2, 0), and (3, 0) electron plasma modes occur at frequencies 10.6, 17.7, and 22.8 MHz, respectively, when undisturbed by the presence of antiprotons in the trap. After the antiprotons were introduced, the electron plasma received energy from them, as a result of which the (2, 0) and (3, 0) mode frequencies first increased for a few seconds, then decreased with a time constant of a few tens of seconds as thermal equilibrium was reestablished.

To evaluate the temperature variation of the electron plasma from these frequency shifts [19], we consider for each  $l$  the dielectric tensor element along the magnetic field,  $\epsilon_3$ . When the plasma temperature  $T_e$  is sufficiently low to maintain a sharp plasma boundary,  $\epsilon_3$  is approximately given by

$$\epsilon_3 \approx 1 - \frac{\omega_{pe}^2}{\omega_l^2} \left( 1 + \frac{\kappa_B T_e}{m} \frac{3k_l^2}{\omega_l^2} \right), \quad (1)$$

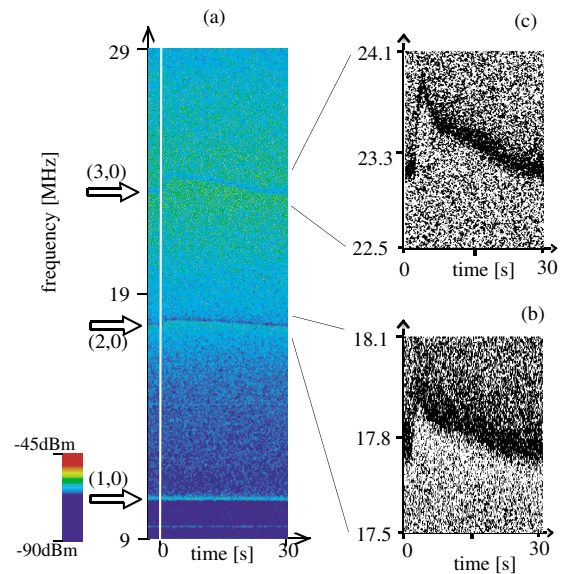


FIG. 2 (color). (a) Time evolution of the power versus frequency spectrum of the electron plasma after antiproton injection at  $t = 0$ . The three arrowed regions correspond to the (1, 0), (2, 0), and (3, 0) modes, respectively; (b) and (c) show the (2, 0) and (3, 0) modes with expanded vertical scales.

where  $k_l \approx \pi(l-1)/(2b)$  is the longitudinal wave number for mode  $(l, 0)$ ,  $2b$  the length of the electron plasma,  $\kappa_B$  the Boltzmann constant,  $m$  the electron mass,  $\omega_l$  the angular frequency of the  $(l, 0)$  mode, and  $\omega_{pe}$  the angular frequency of the electron plasma oscillation. The dispersion relation [20] for a cold spheroidal plasma reduces to [19,21]

$$\epsilon_3 = \left( \frac{\alpha^2 - \epsilon_3}{\alpha^2 - 1} \right)^{1/2} \frac{P_l(\xi_1)Q'_l(\xi_2)}{P'_l(\xi_1)Q_l(\xi_2)}, \quad (2)$$

for a strong magnetic field, where  $\xi_1 = \alpha(\alpha^2 - \epsilon_3)^{-1/2}$ ,  $\xi_2 = \alpha(\alpha^2 - 1)^{-1/2}$ , and  $\alpha$  is the aspect ratio, the ratio of the axial length to the diameter of the plasma.  $P_l$  and  $Q_l$  represent Legendre functions of the first and second kinds, respectively, while  $P'_l$  and  $Q'_l$  denote their derivatives.

Using the observed  $\omega_l$  values ( $l = 2, 3$ , and  $4$ ) before antiproton injection and the electron number previously measured under the same conditions, these equations can be solved to give values  $n_e \approx 4.4 \times 10^7/\text{cm}^3$  and  $\alpha \approx 7.3$  for the plasma density and the aspect ratio. Using  $\alpha$  and  $n_e$  so obtained, the temperature  $T_e$  after antiproton injection was evaluated from  $\omega_l$ , which indicated a maximum plasma temperature raise of  $\approx 0.6$  eV. The variations of the plasma frequencies during the cooling that followed heat transfer from the antiprotons were consistently reproduced by solving rate equations which took into account the synchrotron radiation cooling of electrons with time constant  $\sim 6/(B[\text{T}])^2$  s and the energy transfer between the antiprotons and the electron plasma [19]. Considering that the MRT and the bore tube were cooled to  $\sim 10$  K, the plasma temperature was expected to reach equilibrium at sub eV energies although no explicit measurements were made.

Figure 3 shows typical antiproton annihilation position distributions as a function of the time elapsed since their injection at  $t = 0$ . The high annihilation rate initially seen at the positions of the trap electrodes and the vacuum isolation foil fell rapidly, leaving about half the captured antiprotons in the MRT after a few seconds. The potential depth of the trap was maintained at 50 V during a storage period lasting about 55 s. Once the plasma consisting of electrons and antiprotons had stabilized, few additional annihilations were observed until the antiprotons were released.

The antiproton annihilation distributions in Figs. 3(a) and 3(b) refer respectively to cases in which the electron plasma was ejected or retained after its cooling task had been fulfilled. The antiprotons were extracted continuously from the MRT by ramping the trap potential from 50 to 0 V over a period of some 100 s started at  $t = 55$  s. In the case shown in Fig. 3(a), we ejected the stored electrons before beginning any antiproton extraction procedure by opening and closing the trap for 550 ns several times. This was long enough to allow the electrons to escape, but short enough to retain the antiprotons, these having a much lower velocity than the electrons. As expected, no annihilations

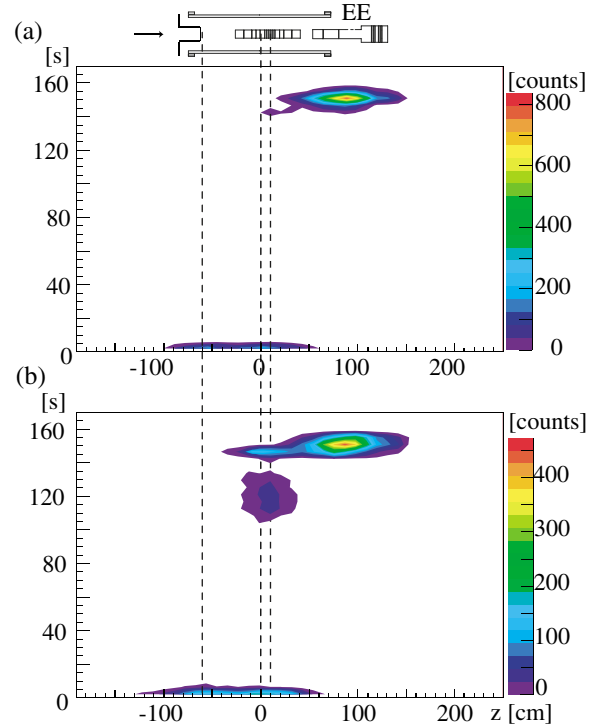


FIG. 3 (color). The  $z$  distribution of antiproton annihilations versus the elapsed time since injection at  $t = 0$ : (a) with and (b) without release of electrons at 40 s (see text). The superposed drawing shows the vacuum isolation foil at  $z = -60$  cm, the MRT center at  $z = 10$  cm, and the first part of the beam line containing the extractor electrodes (EE).

were observed during electron ejection at  $t \approx 40$  s. Furthermore, the fact that none occur between  $t = 55$  s when ramping started and  $t \sim 140$  s confirms that the antiprotons had been cooled into the bottom of the potential well. The strong annihilation peak observed at  $t \approx 150$  s near the downstream extractor electrodes (EE) contained  $\sim 90\%$  of the entire sample of antiprotons trapped in the harmonic potential. Considering that the magnetic field at the EE is 10%–20% of that around the MRT and the inner diameter of the EE is 50 mm, the diameter of the antiproton cloud in the MRT was estimated to be  $\sim 15$ –20 mm.

When [as in Fig. 3(b)] we extracted the antiprotons without first ejecting the electrons, the annihilation pattern changed drastically. The main extraction peak at EE now included only  $\sim 75\%$  of the trapped antiprotons, with the remainder appearing in new peaks at the position of the MRT center, at average times  $t \approx 120$  s and  $t \approx 145$  s. When the first of these new annihilations appeared ( $t \approx 105$  s), the depth of trap potential was  $\sim 25$  V, which is approximately equal to the space charge potential produced by the  $3 \times 10^8$  electrons contained in the electron plasma. The additional peak can thus be well accounted for by annihilations of antiprotons which hit the electrodes due to the  $\mathbf{E} \times \mathbf{B}$  drift caused by instability when electrons overflow from the trap.

A number of new fields of research would be opened up if the extracted antiprotons could be transported along a beam line to experiments located further downstream. These include collision dynamics in the ionization of atoms by antiprotons [22] and antiprotonic atom formation [13], the study of nuclear surface structure via antiprotonic atom formation and annihilation [13,23,24], as well as the formation of antihydrogen atoms [4–7,25].

We conclude from this work that the combination of the RFQD and the MRT results in as many as  $1.2 \times 10^6$  cooled antiprotons being accumulated per AD shot; this is at least 50 times larger than the best values reported with thick-foil energy degraders. The cooling feature of antiprotons was studied nondestructively by monitoring electrostatic plasma modes of electrons. A small number of monoenergetic antiprotons of 10–500 eV was then extracted and delivered to a downstream target area via a beam line with differential pumping capacity of more than 6 orders of magnitude. Further studies along these lines will be reported in a future publication.

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- [1] G. Gabrielse *et al.*, Phys. Rev. Lett. **57**, 2504 (1986).
- [2] M. H. Holzschneider *et al.*, Phys. Lett. A **214**, 279 (1996).
- [3] G. Gabrielse *et al.*, Phys. Rev. Lett. **63**, 1360 (1989).
- [4] M. H. Holzschneider *et al.*, CERN Reports No. CERN/SPSLC 96-47 and No. CERN/SPSC P-302, 1996.
- [5] G. Gabrielse *et al.*, CERN Reports No. CERN/SPSC 97-8 and No. CERN/SPSC P-306, 1997.
- [6] M. Amoretti *et al.*, Nature (London) **419**, 456 (2002).
- [7] ATRAP Collaboration, G. Gabrielse *et al.*, Phys. Rev. Lett. **89**, 213401 (2002); **89**, 233401 (2002).
- [8] G. Gabrielse *et al.*, Phys. Lett. **548B**, 140 (2002).
- [9] W. Pirkel (private communication); Y. Bylinsky, A. M. Lombardi, and W. Pirkel, hep-ex/0008030; CERN Report No. CERN-PS-2000-061-PP; A. M. Lombardi, W. Pirkel, and Y. Bylinsky, in *Proceedings of the 2001 Particle Accelerator Conference, Chicago* (IEEE, Piscataway, 2001), p. 585.
- [10] M. Hori *et al.*, Phys. Rev. Lett. **91**, 123401 (2003).
- [11] J. Eades and F. J. Hartman, Rev. Mod. Phys. **71**, 373 (1999).
- [12] T. Azuma *et al.*, CERN Reports No. CERN/SPSC 97-19 and No. CERN/SPSC P-307, 1997.
- [13] Y. Yamazaki, Nucl. Instrum. Methods Phys. Res., Sect. B **154**, 174 (1999); Hyperfine Interact. **138**, 141 (2001).
- [14] B. R. Beck, J. Faján, and J. Malmberg, Phys. Plasmas **3**, 1250 (1996).
- [15] A. Mohri *et al.*, Jpn. J. Appl. Phys. **37**, 664 (1998).
- [16] H. Higaki and A. Mohri, Phys. Lett. A **235**, 504 (1997).
- [17] A test experiment with protons supplied from a tandem accelerator at Aarhus University showed the deceleration efficiency was about 40%, which indicates that there is still room for a better tuning of the AD and the RFQD.
- [18] CERN Application Software Group, GEANT version 3.21, CERN Program Library Long Writeup W5013, 1994.
- [19] H. Higaki *et al.*, Phys. Rev. E **65**, 046410 (2002).
- [20] D. H. Dubin, Phys. Rev. Lett. **66**, 2076 (1991).
- [21] M. D. Tinkle *et al.*, Phys. Rev. Lett. **72**, 352 (1994).
- [22] H. Knudsen and J. F. Reading, Phys. Rep. **212**, 107 (1992).
- [23] A. Trzińska *et al.*, Phys. Rev. Lett. **87**, 082501 (2001).
- [24] M. Wada *et al.*, in *Non-Neutral Plasma Physics IV*, edited by F. Anderregg *et al.*, AIP Conf. Proc. No. 606 (AIP, New York, 2001), p. 625.
- [25] A. Mohri and Y. Yamazaki, Europhys. Lett. **63**, 207 (2003).