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# Creation and Control of Two-Dimensional Electron Gas Using Al-Based Amorphous Oxides/SrTiO<sub>3</sub> Heterostructures Grown by Atomic Layer Deposition

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### ABSTRACT

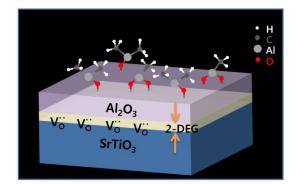
The formation of a two-dimensional electron gas (2-DEG) using SrTiO<sub>3</sub> (STO)-based heterostructures provides promising opportunities in oxide electronics. We realized the formation of 2-DEG using several amorphous layers grown by the atomic layer deposition (ALD) technique at 300 °C which is a process compatible with mass production, and thereby can provide realization of potential applications. We found that the amorphous LaAlO<sub>3</sub> (LAO) layers grown by the ALD process can generate 2-DEG ( $\sim 1 \times 10^{13}$ /cm<sup>2</sup>) with an electron mobility of 4-5 cm<sup>2</sup>/V·s. A much higher electron mobility was observed at lower temperatures. More remarkably, amorphous YAIO<sub>3</sub> (YAO), and Al<sub>2</sub>O<sub>3</sub> layers, which are not polar-perovskitestructured oxides, can create 2-DEG as well. 2-DEG was created by means of the important role of trimethylaluminum, Me<sub>3</sub>Al, as a reducing agent for STO during LAO and YAO ALD as well as the Al<sub>2</sub>O<sub>3</sub> ALD process at 300 °C. The deposited oxide layer also plays an essential role as a catalyst that enables Me<sub>3</sub>Al to reduce the STO. The electrons were localized very near to the STO surface, and the source of carriers was explained based on the oxygen vacancies generated in the STO substrate.

KEYWORDS: 2-D electron gas, amorphous, LaAlO<sub>3</sub>, Al<sub>2</sub>O<sub>3</sub>, SrTiO<sub>3</sub>, atomic layer deposition (ALD), oxygen vacancy

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Creation and Control of Two-Dimensional Electron Gas Using Al-Based Amorphous Oxides/SrTiO<sub>3</sub> Heterostructures Grown by Atomic Layer Deposition



Heterostructures between complex oxide layers are emerging as interesting systems for oxide electronics due to their unique properties.<sup>1</sup> Ohtomo reported the existence of a 2dimensional electron gas (2-DEG) at the heterointerface between the two insulating oxides LaAlO<sub>3</sub> (lanthanum aluminate, LAO) and SrTiO<sub>3</sub> (strontium titanate, STO).<sup>2</sup> The LAO films were grown by pulsed laser deposition (PLD) or molecular beam epitaxy (MBE) on a single crystalline (001) STO substrate.<sup>2-8</sup> Transistor-like devices using these LAO/STO heterostructures have also been demonstrated.3 The most common explanation of the observation of 2-DEG at the LAO/STO interface was based on electronic reconstruction, where the electrons move to the interface to avoid the potential divergence caused by the polar catastrophe.<sup>9</sup> The interface charges are compensated by Ti<sup>3+</sup> ions that provide electrons to the STO conduction band, and these interface electrons are localized near the AlO<sub>2</sub>/LaO/TiO<sub>2</sub> interface, which is called an n-type interface. Basically, this electronic reconstruction can only be realized at a well-prepared abrupt and sharp epitaxial interface between single crystalline LAO and STO. An alternative explanation for the generation of 2-DEG is La inter-diffusion through the interface, doping the surface layer of the STO to be n-type.<sup>10</sup> A third explanation is defect generation such as oxygen vacancies in the STO substrate.<sup>10-12</sup> Recently, it has been discussed that the oxygen vacancies can be generated at the STO substrate during the growth

of the LAO layer as the possible origin of the carrier source.<sup>13,14</sup> In the meantime, the formation of oxygen vacancies has been also supported by theoretical first-principle calculations.<sup>15,16</sup>

Strontium titanate is well-known as a high-*k* dielectric (bulk dielectric constant of 300) with a bandgap of 3.2 eV which has been extensively studied for possible use in electronic devices such as dynamic random access memory (DRAM). Strontium titanate has a cubic perovskite structure with a lattice constant of 0.3905 nm. Lanthanum aluminate is a wide bandgap insulator (5.6 eV) which has a quasi-cubic structure with a lattice constant of 0.3789 nm; thus the lattice mismatch between STO and LAO is small (~3 %).<sup>2</sup>

Up to now, epitaxially grown LAO films have been necessary for the creation of 2-DEG. The epitaxial LAO films were grown on an STO substrate by the PLD process at a relatively high growth temperature (650~850 °C)<sup>2-7,9,11-14,17-20</sup> which is one of the representative physical vapor deposition (PVD) techniques. However, the feasibility of an amorphous layer is rarely reported, except in a recent study using amorphous layers grown by the PLD process at room temperature.<sup>13</sup> Unfortunately, the PLD process has a lower chance of being used in mass production of devices.

ALD is characterized by its unique self-limiting growth mechanism based on alternating saturated surface reactions. The precursor and oxygen source are pulsed into a reactor alternately, and each injection is separated by a purging process using inert gas.<sup>21</sup> This

process provides a lot of advantages in terms of excellent conformality and precise control of thickness as well as large area uniformity. In practice, the ALD process is being widely used for the growth of dielectric films in the mass-production of microelectronic devices.<sup>21</sup>

In this paper, a practical, novel technique is proposed to create and control 2-DEG using several amorphous films on STO substrates in which oxygen vacancies act as the source of the carriers. An amorphous LAO layer grown by an ALD process creates 2-DEG by generating oxygen vacancies in the STO substrate. More remarkably, amorphous Al<sub>2</sub>O<sub>3</sub> as well as YAlO<sub>3</sub> (YAO) grown by ALD on STO substrates, which are not polar-perovskite structured oxides, can create 2-DEG at the interface. The aluminum source for these ALD processes, trimethylaluminum (TMA), creates oxygen vacancies in the STO substrate during the growth of several aluminum-containing layers, including amorphous LAO, YAO and Al<sub>2</sub>O<sub>3</sub>. However, treatment with TMA vapor alone does not create any oxygen vacancies. The presence of a growing oxide layer (LAO, Al<sub>2</sub>O<sub>3</sub>, YAlO<sub>3</sub>, etc.) produced by water vapor as a co-reactant is also needed to catalyze the reduction. 2-DEG density was controlled precisely and reproducibly, which indicates that the formation of oxygen vacancies was well-controlled by the ALD process. Growth of other oxides not containing Al, such as Y<sub>2</sub>O<sub>3</sub>, La<sub>2</sub>O<sub>3</sub>, or LaYO<sub>3</sub>, did not produce 2-DEG.

All films used in this experiment, including the LAO layer, were grown by ALD at 300 °C. The ALD technique enabled the formation of smooth and high quality oxide thin films at a low growth temperature (300 °C) with atomic precision, which is essential to the functionality of oxide heterostructures. The Al<sub>2</sub>O<sub>3</sub> ALD process offers great advantages in terms of mass-production compatibility due to its ideal ALD reaction.<sup>22</sup> The growth of an amorphous layer is easy to achieve by the ALD process with a large process margin; thus it is promising for the realization of potential applications.

### **Results and Discussions**

Figure 1 (a) shows the cross sectional high-resolution transmission electron microscopy (TEM) image of the LAO layer grown by ALD at 300 °C on an STO substrate. The amorphous LAO layer grown on a single crystalline STO substrate is clearly observed. Figure 1 (b) shows the [001] zone axis electron diffraction pattern at the interfacial region of the LAO/STO heterostructure. The diffraction pattern consists of single spots, and those diffraction spots are exactly consistent with those of single crystalline STO<sup>23</sup>, which indicates that the STO substrate used in this experiment is a perfect single crystalline material. However, no diffraction spots were observed from the LAO layer, which confirms the amorphous nature of the LAO layer.

The surface morphologies of the samples were investigated by means of an atomic force microscope (AFM). Figure 2 (a) shows the surface morphology of a TiO<sub>2</sub>-terminated STO substrate.<sup>24,25</sup> The surface consists of regular terrace structures with a high smoothness (rms roughness of 0.15 nm). The height difference between the adjacent terraces is about 0.4 nm which corresponds to one unit-cell of STO as shown in Fig. 2 (b). Figure 2 (c) shows the surface morphology of a 6 nm thick LAO film grown on an STO substrate by the ALD process at 300 °C. The roughness of the LAO surface (about 0.2 nm) is slightly higher than that of the STO substrate. The background terrace structure was preserved although the LAO film was grown on an STO substrate. This result implies that LAO films grown by ALD are very smooth, replicating the underlying terrace structure. Figure 2 (d) shows the sheet resistance  $(R_s)$  and sheet carrier density  $(n_s)$  as a function of the amorphous LAO film thickness grown on the STO substrate (LAO/STO heterostructure). A stoichiometric (51.4 at. % of La) LAO film was grown by an ALD cycle ratio of 2:1 (La:Al) as confirmed by RBS (Fig. S1 (a)), and the thicknesses of the films grown on the STO substrate were estimated by an X-ray reflectivity (XRR) measurement (Fig. S2). Even with amorphous LAO films, the carriers can be created and a critical thickness exists which is required to make the channel to be conducting. The insulating interface of the LAO/STO heterostructures becomes conducting above a LAO thickness of ~3 nm, and the sheet resistance and sheet carrier density remain constant above this critical thickness. Here, the type of carrier is an electron according to the sign of the Hall voltage. This transition behavior of the amorphous LAO/STO heterostructure is similar to the epitaxial LAO/STO heterostructure.<sup>3,6</sup> The critical thickness of ~3 nm is somewhat thicker than the commonly reported value (~4 unit cells) from the epitaxial LAO films grown by PLD.<sup>3</sup> The sheet resistance and sheet carrier density of the amorphous LAO/STO heterostructures were  $\sim 1x10^5 \Omega/sq$ . and  $\sim 1x10^{13}/cm^2$ , respectively, where the mobility was 4~5 cm<sup>2</sup>/V·s as shown in Fig. 2 (e), whereas those of the insulating heterostructures were  $>10^9 \Omega/sq$ . and  $<10^9/cm^2$  which are beyond the sensitivity of the measurement.

Astonishingly, the 2-DEG can be generated by depositing an amorphous Al<sub>2</sub>O<sub>3</sub> layer as well as a YAO layer on the STO substrate as shown in Fig. 3. Figure 3 (a) shows the schematic diagram of several heterostructures using various oxide layers. The sheet resistance and carrier density of the YAO/STO heterostructure were almost identical with those of the LAO/STO heterostructure as shown in Fig. 3 (b). The sheet resistance of the Al<sub>2</sub>O<sub>3</sub>/STO heterostructure is slightly higher than that of the STO/LAO heterostructure by a factor of 3 because of the reduced carrier density (~3x10<sup>12</sup>/cm<sup>2</sup>). However, the 2-DEG was not generated by depositing amorphous La<sub>2</sub>O<sub>3</sub>, Y<sub>2</sub>O<sub>3</sub> or LaYO<sub>3</sub> (LYO) layers on an STO substrate at 300 °C by ALD as shown in Fig. 3 (b). The basic crystal structure of YAO is an orthorhombic, a space group P<sub>nma</sub> (space group number 62), which lattice constants are: a = 0.5330 nm, b = 0.7375 nm, and c = 0.5180 nm.<sup>26</sup> This is a small distortion of the perovskite unit cell. However, it should be noted that  $Al_2O_3$  is not a polar-perovskite structured oxide with cubic or distorted cubic unit cells. Instead, crystalline  $Al_2O_3$  can be described by trigonal or hexagonal unit cells (space group number 167), while ALD  $Al_2O_3$  is amorphous.

Figure 4 (a) shows the sheet resistance and sheet carrier density of an Al<sub>2</sub>O<sub>3</sub>/STO heterostructure as a function of Al<sub>2</sub>O<sub>3</sub> film thickness when Al<sub>2</sub>O<sub>3</sub> is grown on a STO substrate. The sheet resistance decreased sharply, and the sheet carrier density increased abruptly above an Al<sub>2</sub>O<sub>3</sub> thickness of 1.2 nm. Then the sheet resistance and sheet carrier density remained constant above this critical thickness. The sheet resistance and sheet carrier density of the conducting heterostructure are  $4x10^5 \Omega$ /sq. and  $\sim 3x10^{12}$ /cm<sup>2</sup>, respectively. The carrier density of the Al<sub>2</sub>O<sub>3</sub>/STO heterostructure is slightly lower than that of the LAO/STO heterostructure ( $\sim 1x10^{13}$ /cm<sup>2</sup>) as mentioned earlier. The mobility was constant regardless of the Al<sub>2</sub>O<sub>3</sub> film thickness (Fig. S3 (a)).

Figure 4 (b) shows the sheet resistance and sheet carrier density of the  $LAO/Al_2O_3/STO$  heterostructure as a function of  $Al_2O_3$  film thickness when the  $Al_2O_3$  film was grown prior to the growth of the LAO film. Six nm-thick LAO films were grown on top of  $Al_2O_3$  films to fabricate the  $LAO/Al_2O_3/STO$  heterostructures. The sheet resistance and created carrier density of the  $LAO/Al_2O_3/STO$  heterostructure were identical to those of the  $Al_2O_3/STO$ 

heterostructure, which suggests the creation of the 2-DEG was determined by the  $Al_2O_3$  layer, not by the LAO layer. The mobility was constant irrespective of the  $Al_2O_3$  film thickness (Fig. S3 (a)).

Figure 4 (c) shows the sheet resistance and sheet carrier density of the YAO/STO heterostructure as a function of YAO film thickness. The stoichiometric (51.3 at. % of Y) YAO film was grown by an ALD cycle ratio (Y:Al) of 2:1 as confirmed by Rutherford back-scattering spectroscopy (RBS) (Fig. S1 (b)). The transition behaviors of the resistance, carrier density and mobility depending on the film thickness are almost identical with the LAO/STO heterostructure as shown in Fig. 2 (d). However, the insulating interface of the YAO/STO heterostructure becomes conducting above a YAO thickness of ~4 nm. The critical thickness of ~4 nm for YAO is thicker than that of the LAO/STO heterostructure.

All these layers are amorphous as confirmed by TEM and X-ray diffraction (XRD) (Fig. S4), so the polar catastrophe mechanism can be ruled out as the explanation for the formation of 2-DEG in our heterostructures. The possibility of a La ion doping into the STO substrate was ruled out because the 2-DEG was not generated when the single layer of  $La_2O_3$  was grown on the STO substrate. It has been reported that La atoms can replace Sr sites in STO to generate electrons by  $La_{Sr}$  substitution;<sup>27,28</sup> however, this phenomenon was not observed in this work. The possibility of an Al ion doping into the STO substrate can also be excluded

because it is known that Al atoms are prone to replace Ti sites, where Al atoms can act as acceptors.<sup>29-31</sup> The substituted Al atoms in Ti sites reduce the density of electrons, so that the Al doping technique has been used to reduce the carriers in TiO<sub>2</sub>-based dielectric materials. The chance of a Y ion doping into the STO substrate was also excluded, as well as La ion doping, for the same reason.

The likely origin of the free electrons is the creation of oxygen vacancies in the STO during the growth of aluminum-containing oxides such as LAO, YAO and Al<sub>2</sub>O<sub>3</sub>. The possibility of reduction of the STO substrate prior to the deposition of each film can be excluded. It has been reported that the STO substrate is reduced at a growth temperature of 800 °C and low oxygen pressure of 10<sup>-6</sup> mbar without the actual deposition of an LAO film grown by PLD due to high growth temperature and low oxygen pressure.<sup>14</sup> Basically, the STO substrate used in our experiment is insulating ( $R_s > 10^9 \Omega/sq$ .) before the deposition of the oxides, and the STO substrate still keeps its insulating property even in the chamber without actual growth of the film for a few hours. Also a STO substrate keeps its insulating property even after an exposure of the STO substrate to TMA molecules (see Fig. S6 and supporting information).

The observation of the Ti<sup>3+</sup> binding state from the Ti 2p core-level spectrum can provide direct evidence of oxygen vacancies in the STO. X-ray photoelectron spectroscopy

(XPS) measurements were performed in order to determine the valence state of Ti as shown in Fig. 5. Figure 5 (a) shows the normalized Ti 2p spectra of a 2.5 nm-thick Al<sub>2</sub>O<sub>3</sub>/STO heterostructure, and the signal from the bare STO substrate was used for comparison. An increased Ti<sup>3+</sup> signal at ~456.2 eV was observed in the Al<sub>2</sub>O<sub>3</sub>/STO heterostructure compared to the bare STO substrate, which means that the oxygen vacancies were generated in the STO during the growth of the  $Al_2O_3$  layer by ALD. For the estimation of the Ti<sup>3+</sup> depth profile, angle-resolved XPS was measured with various photoelectron take-off angles with a 2.5 nmthick Al<sub>2</sub>O<sub>3</sub>/STO heterostructure as shown in Fig. 5 (b). The effective electron escape depth decreases with decreasing the takeoff angle; thus the XPS spectrum becomes more surfacesensitive (Fig. S5). The signal of Ti<sup>3+</sup> relative to the Ti<sup>4+</sup> main line increases as the takeoff angle decreases. This indicates that the electrons generated by oxygen vacancies are localized at the STO surface region as a 2-DEG (within  $\sim 2$  nm from the interface), which is smaller than the electron escape depth (~5 nm, see Supporting Information).

The mobilities of all conducting heterostructures (~5  $\text{cm}^2/\text{V}\cdot\text{s}$ ) are consistent with generally reported values of reduced STO having many oxygen vacancies.<sup>11</sup> The carriers of all heterostructures disappeared after annealing at 600 °C for 1 hour in an oxygen atmosphere. The carriers disappeared even after annealing at 300 °C for 1 hour in an oxygen atmosphere. These

facts provide additional evidence that the source of the carriers was generated by oxygen vacancies on the STO side.

The transfer of the oxygen between LAO and STO has been experimentally reported during the growth of each oxide grown by PLD using an <sup>18</sup>O isotope by means of dynamic secondary ion mass spectroscopy (D-SIMS).<sup>17</sup> Oxygen transfer was observed at high growth temperature (750 °C) as the thick LAO film (>50 nm) was grown on the STO substrate. However, the oxygen transfer was significantly reduced as the growth temperature decreased below 650 °C.

When it comes to the ALD process, oxygen transfer and reduction of the substrate is also possible provided the thermodynamic condition is satisfied and the reaction kinetics are fast enough. The metal precursor molecules are introduced on the substrate separately without overlapping the oxygen source in the ALD process. Usually, these metal precursor molecules tend to be oxidized by scavenging the nearby oxygen.<sup>32</sup> The substrate material was reduced during the pulse of the metal precursor without introducing an oxygen source, such as an H<sub>2</sub>O or O<sub>3</sub> vapor, by means of oxidation of the metal precursor.<sup>32</sup> However, the kinetics of the redox reaction can be limited by the low growth temperature of ALD and the strong bond of complex ligands attached to the metal precursor molecules in the ALD reaction.

The Gibbs free energy of reaction ( $\Delta G_r$ ) can indicate whether the reduction of the STO substrate by TMA during the Al<sub>2</sub>O<sub>3</sub> ALD process is feasible.<sup>33-35</sup> Possible overall reactions for the reduction of the STO substrate by the TMA at 573 K are the following:

$$2AI(CH_3)_3 + 8TiO_2 \implies AI_2O_3 + 4Ti_2O_3 + 3CH_4(g) + \frac{3}{2}C_2H_4(g) + \frac{1}{2}O_2(g)$$

 $\Delta G_{r, 573K} = -947.156 \text{ kJ/mol}$  (1)

or

$$2Al(CH_3)_3 + 10TiO_2 \Rightarrow Al_2O_3 + 5Ti_2O_3 + 3CH_4(g) + C_2H_4(g) + CO_2(g) + H_2(g)$$

$$\Delta G_{r, 573K} = -1062.726 \text{ kJ/mol}$$
 (2)

The reaction pathways were expected based on the reports that CH<sub>4</sub> gas was mainly produced with small amounts of C<sub>2</sub>H<sub>4</sub> and CO<sub>2</sub> gases as byproducts when TMA dissociates in the presence of oxygen during the Al<sub>2</sub>O<sub>3</sub> ALD process.<sup>36,37</sup> The change in Gibbs free energies ( $\Delta G_r$ ) is negative for reaction pathways (1) and (2) because of the formation of strong Al-O bonds. Thus, at equilibrium, the oxidation of TMA molecules can drive a thermodynamically favorable reduction of STO substrates.<sup>38,39</sup> Of course, other gaseous byproducts, such as CO or H<sub>2</sub>CO, are also possible, but large negative free energy changes also favor these reactions going to completion at equilibrium.

In one experiment, reduction of the STO by the injection of TMA alone did not occur, which means that the thermal energy was not enough to overcome the activation energy due to the slow kinetics at a reaction temperature of 300 °C, even though the thermodynamic condition was satisfied. However, reduction of the STO did happen when LAO, YAO and Al<sub>2</sub>O<sub>3</sub> films were thicker than their critical thicknesses. As a result, the repetitive pulse (55 cycles) of the Al precursor, TMA, which was used for the deposition of the LAO, YAO and  $Al_2O_3$  films, on the STO substrate without introducing an oxygen source (H<sub>2</sub>O) did not induce the reduction of the STO substrate at 300 °C because of the absence of the LAO, YAO and Al<sub>2</sub>O<sub>3</sub> layers. This is also the case when the repetitive pulses (55 cycles) of TMA are applied to the  $Y_2O_3$  layers previously grown on the STO substrate, irrespective of the Y<sub>2</sub>O<sub>3</sub> film thickness (Fig. S6 (a)). It should be noted that the interface of the  $Y_2O_3/STO$  heterostructure is insulating as shown in Fig. 3 (b). In this case, the carriers were generated by means of following the deposition of a 6 nmthick LAO layer on top of the  $Y_2O_3/STO$  heterostructure as shown in Fig. S6 (a). There still exists a required critical thickness of LAO and Al<sub>2</sub>O<sub>3</sub> layers when those layers are grown on top of the Y<sub>2</sub>O<sub>3</sub> layer (4.7 nm-thick) to create carriers, which emphasizes the important role of pregrown LAO and Al<sub>2</sub>O<sub>3</sub> layers (Fig. S6 (b)). LAO and Al<sub>2</sub>O<sub>3</sub> layers that are thicker than the critical thickness are necessary for the redox reaction even on the Y<sub>2</sub>O<sub>3</sub>/STO heterostructure; thus the reduction of the STO can occur during the metal precursor pulse. As discussed above,

the carriers were generated only after growing a certain-thickness LAO (>3 nm), YAO (>4 nm) and Al<sub>2</sub>O<sub>3</sub> (>1.2 nm) layers on the STO substrate at 300 °C. An important feature is that the Al<sub>2</sub>O<sub>3</sub> component is always included in those oxides, such as LAO, YAO and Al<sub>2</sub>O<sub>3</sub>, which generated the carriers in an interaction with the STO. Here, we claim that  $Al_2O_3$  (>1.2 nm) acts as a catalyst for the reaction; thus the redox reaction of the STO substrate can occur when the TMA molecules are introduced. The Al<sub>2</sub>O<sub>3</sub> layer (>1.2 nm) promoted the kinetics of the redox reaction. Therefore, the redox reaction is strongly enhanced at the interface, which results in the oxygen vacancies generated on the STO surface by means of the TMA molecules' oxidation at 300 °C. The minimum required  $Al_2O_3$  amount for the carrier generation was estimated by ~15 cycles of Al<sub>2</sub>O<sub>3</sub> ALD, which corresponds to a 1.35 nm-thick Al<sub>2</sub>O<sub>3</sub> layer. An important fact suggesting that this redox reaction is limited by kinetics is that the carrier density decreased significantly with the decreasing growth temperature (Fig. S7). The carrier density was decreased by a factor of  $10^3$  at a growth temperature of 250 °C; and it could not be measured at 200 °C due to the measurement limit (decreased by a factor of >10<sup>4</sup>).

It is understood that the outward diffusion of oxygen owing to the redox reaction of the STO is limited by grown oxides such as the LAO, YAO and Al<sub>2</sub>O<sub>3</sub> layers due to their diffusivity of oxygen. The outward diffusion of oxygen from the interface of the heterostructure would be limited by those oxides; thus, the carrier density does not increase above the critical thickness.

This is an interesting feature in terms of the self-limiting control of the carrier by blocking the outward diffusion of oxygen in the redox reaction. It has been reported that metallic species such as Al, Ti, Y, Ba and Zr can interact with the STO substrate at room temperature, which indicates the interfacial redox reaction occurs, and the interactions happen at the STO surfaces.<sup>13,35</sup>

Based on the above discussion, we expected that the carriers would be generated even though the composition of the YAO layer is varied far from the stoichiometry, as long as the amount of the  $Al_2O_3$  component in the YAO layer is enough (>~15 cycles). Figure 4 (d) shows the sheet resistance and carrier density depending on the Y-rich YAO layer thickness. A Y-rich YAO (62.6 at% of Y) film was grown by an ALD cycle ratio (Y:Al) of 4:1 as determined by RBS (Fig. S1 (c)). Although the composition of the YAO layer is out of the stoichiometric range, the transition behavior was also observed. This result supports that the creation of 2-DEG does not originate from the perovskite nature of the YAO layer. The insulating interface of the YAO/STO heterostructure becomes conducting above the YAO thickness of ~4.5 nm. The critical thickness of ~4.5 nm is a little bit thicker than the stoichiometric YAO/STO heterostructure (~4 nm). A Y-rich YAO film (4 Y<sub>2</sub>O<sub>3</sub> cycles per Al<sub>2</sub>O<sub>3</sub> cycle) is thicker than a stoichiometric YAO film (2 Y<sub>2</sub>O<sub>3</sub> cycles per Al<sub>2</sub>O<sub>3</sub> cycle) because more Y<sub>2</sub>O<sub>3</sub> ALD cycles are included as the same Al<sub>2</sub>O<sub>3</sub> cycles in Y-rich YAO ALD than in stoichiometric YAO ALD. The minimum required  $Al_2O_3$  cycles in the Y-rich YAO/STO heterostructure for the carrier generation were almost identical (~15 cycles of  $Al_2O_3$  ALD) with the stoichiometric YAO/STO, and also with the LAO/STO and  $Al_2O_3$ /STO heterostructures, which indicates that the amount of  $Al_2O_3$  in the film is the most important factor of the redox reaction. The creation of 2-DEG during the growth of the  $Al_2O_3$  layer on the STO surface was schematically described according to the above model in Fig. 6.

Figures 7 (a)-(c) show the variation of sheet resistance, mobility and sheet carrier density acquired from Hall measurement results depending on the measurement temperature of the several heterostructures. The sheet resistance variations of all heterostructures were decreased with a decreasing temperature because the mobility increases with a decreasing temperature (Fig. 7 (a)). The mobility variations were identical for all heterostructures which indicates the mobilities were governed by the STO, and the increasing mobility with decreasing temperature is a typical property of STO (Fig. 7 (b)).<sup>28,40</sup> The mobility was further increased with decreasing temperature down to 30 K (Fig. S8). The sheet carrier density of the LAO/STO heterostructure was slowly decreased with decreasing temperature, and the sheet carrier density of the Al<sub>2</sub>O<sub>3</sub>/STO heterostructure was decreased slightly faster than the LAO-based heterostructure (Fig. 7 (c)). However, the insertion of the Y<sub>2</sub>O<sub>3</sub> layer between the STO substrates and the LAO or Al<sub>2</sub>O<sub>3</sub> films produced more metallic channel (E<sub>4</sub> < 10 meV) where

the carrier density was not decreased with decreasing temperature. Figure 7 (d) shows the Arrhenius plot of the sheet carrier densities versus temperature, which gives the activation energy from the slope. The activation energies were close to zero in the LAO/Y2O3/STO and Al<sub>2</sub>O<sub>3</sub>/Y<sub>2</sub>O<sub>3</sub>/STO heterostructures, which indicates metallic channels were formed. At room temperature, the sheet resistance was even decreased from  $1.1 \times 10^5$  to  $6.72 \times 10^4 \Omega/sq$ , when the Y<sub>2</sub>O<sub>3</sub> layer was inserted between the LAO layer and the STO substrate. The sheet carrier density was increased from  $1.16 \times 10^{13}$  to  $1.87 \times 10^{13}$ /cm<sup>2</sup>. The role of Y<sub>2</sub>O<sub>3</sub> between LAO and STO is totally different from that of the epitaxially grown STO/Y<sub>2</sub>O<sub>3</sub>/STO heterostructure, where the single atomic layer of Y<sub>2</sub>O<sub>3</sub> induced an insulating interface,<sup>41</sup> because the electronic reconstruction theory does not fit our system. Although the carrier densities decrease slightly with decreasing temperature in the LAO/STO and Al<sub>2</sub>O<sub>3</sub>/STO heterostructures, the activation energies were very small ( $E_a < 30$  meV) and consistent with the oxygen-deficient STO<sup>42</sup> as shown in Fig. 7 (d) which indicates that conducting channels were successfully formed.

#### **IV. Conclusions**

Up to now, a physical vapor deposition method, such as PLD, was used to generate 2-DEG to grow an epitaxial LAO film on an STO substrate. Here, we showed that 2-DEG can be created by growing an amorphous LAO layer as well as YAO and Al<sub>2</sub>O<sub>3</sub> layers grown by the

ALD process on an STO substrate. It was revealed that the key factor for the creation of 2-DEG was an  $Al_2O_3$  ALD process forming oxygen vacancies in the STO by oxidation of the TMA. In practice, the ALD technique has been used for the mass production in the semiconductor field. Among them, the ALD of  $Al_2O_3$  is strongly recommended in the mass production since it is an easy and cheap process with a large processing window and precise atomic level control. And the TMA precursor has a very simple structure which can induce an ideal ALD reaction. Eventually, the formation of oxygen vacancies on the STO substrate is well-controlled by the ALD process. These heterostructures fabricated at a low temperature (~300 °C) by ALD provide promising opportunities in oxide electronics. In addition, it can allow more chance to the scalability of a device if the 2-DEG can be created by ALD on the epitaxially grown STO films on SiO<sub>2</sub>; likewise the observation of 2-DEG on the epitaxially grown STO films on SiO<sub>2</sub> according to the recent report by Park.<sup>6</sup> Thus, the creation and control of 2-DEG by ALD technique using a TMA precursor is of great importance in terms of the realization of device fabrication. Further process optimization and theoretical calculations will aid in the achievement of devices with even higher performance.

**II. Methods.** *Sample Growth.* LAO films were deposited on TiO<sub>2</sub>-terminated (001) STO single crystals purchased from MTI Corporation, and TiO<sub>2</sub>-termination was achieved by chemical etching.<sup>24,25</sup> The LAO films were grown at a growth temperature of 300 °C by ALD in

a horizontal gas flow reactor at a working pressure of 400 mTorr (base pressure 30 mTorr). All films, including an LAO layer, were grown by ALD at 300 °C. Lanthanum tris(N,N'-diisopropylformamidinate) (Dow Chemical Company) was used as the La-precursor, and H<sub>2</sub>O was used as the oxygen source. TMA (Sigma Aldrich) was used as the Al-precursor and H<sub>2</sub>O was used as the oxygen source for the deposition of Al<sub>2</sub>O<sub>3</sub>. One super cycle consisted of 2 sub-cycles of La<sub>2</sub>O<sub>3</sub> and 1 sub-cycle of Al<sub>2</sub>O<sub>3</sub> for the deposition of the stoichiometric LAO films. A Y<sub>2</sub>O<sub>3</sub> thin film was deposited by ALD using yttrium tris(N,N'-diisopropylacetamidinate) (Dow Chemical Company) and H<sub>2</sub>O at a growth temperature of 300 °C. A stoichiometric YAO film (51.3 at% of Y) was grown with a super cycle ratio of 4:1 (Y:Al).

*Measurement of Film Properties.* The thickness of the film grown on an STO substrate was estimated by X-ray reflectivity (XRR, PANAlytical, X'Pert Pro) measurements using a Cu K $\alpha$  X-ray source. The amorphous phase of the film was analyzed by high-resolution TEM (HRTEM, JEOL 2010F) and glancing angle X-ray diffraction (GAXRD). The film surface morphology was investigated by an atomic force microscope (AFM, Asylum, MFP-3D SA). The composition of the film was evaluated by Rutherford back-scattering spectroscopy (RBS) using a 2 MeV He<sup>2+</sup> beam. X-ray photoelectron spectroscopy (XPS, PHI, VersaProbe II) measurements were performed in order to determine the valence state of the Ti. For the

estimation of the Ti<sup>3+</sup> depth profile, the angle-resolved XPS measurement was examined with various photoelectron take-off angles. The effective electron escape depth decreases with the decreasing take-off angle; thus the analysis becomes more surface-sensitive.

*Electrical Measurement.* For the Hall measurement, four 100 nm-thick Au electrodes (diameter of 1 mm with a 10 nm-thick Ti adhesion layer) were deposited at the corners of 1 cm square samples by e-beam evaporation (Denton) using a shadow mask after the deposition of the thin films onto the STO substrate. Sheet resistance ( $R_s$ ) and sheet carrier density ( $n_s$ ) were measured by the Hall measurement system using the Van der Pauw configuration, and the mobility was evaluated from the relationship between  $R_s$  and  $n_s$ . Ohmic contact property was confirmed by I-V measurement between two Au electrodes (Fig. S9). The measuring temperature was controlled by means of vacuum-assisted cooling with the help of liquid nitrogen, which allowed the lowest temperature of 100 K.

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Figure 4 Sheet resistance and sheet carrier density of (a) Al<sub>2</sub>O<sub>3</sub>/STO and (b) LAO/ Al<sub>2</sub>O<sub>3</sub>/STO heterostructures as a function of Al<sub>2</sub>O<sub>3</sub> film thickness. 6 nm-thick LAO films were grown on top of Al<sub>2</sub>O<sub>3</sub> films in Fig 4 (b). The sheet resistance and sheet carrier density of (c) stoichiometric (51.3 at. % of Y) YAO/STO and Y-rich (62.6 at. % of Y) YAO/STO heterostructures as a function of YAO film thickness. YAO films were grown by ALD using cycle ratios of 2:1 (Y:Al) and 4:1, respectively.

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Figure 6 Schematic diagram of 2-DEG creation by reduction of STO by TMA during the growth of an Al<sub>2</sub>O<sub>3</sub> film at 300 °C. The Al<sub>2</sub>O<sub>3</sub> layer (>15 cycles of Al<sub>2</sub>O<sub>3</sub> ALD) promoted the kinetics of the redox reaction; thus the oxygen vacancies could be generated on the STO surface by oxidation of TMA molecules.

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Figure S1 The composition of (a) LAO film (17 nm-thick) and YAO (18 nm) films with cycle ratios (Y:Al) of (b) 2:1 and (c) 4:1, evaluated by Rutherford back-scattering spectroscopy (RBS) using a 2 MeV He<sup>2+</sup> beam. The estimated atomic ratio of La/[La+Al] in the LAO film was 51.4 %, which indicates the grown film is very close to stoichiometry The estimated atomic ratios of Y/[Y+Al] in YAO films were 51.3 and 62.6% with cycle ratios of 2:1 and 4:1, respectively.

Figure S2 The thickness of (a) LAO, (b) Al<sub>2</sub>O<sub>3</sub> and (c) YAO films grown on an STO substrate estimated by X-ray reflectivity measurement. Films were grown on Ru substrates

simultaneously and the thickness was measured by ellipsometry to confirm the linear growth behavior depends on the deposition cycle number. (d) Representative simulation results of LAO,  $Al_2O_3$  and YAO films grown on the STO substrate.

Figure S3 (a) The mobility of Al<sub>2</sub>O<sub>3</sub>/STO and LAO/Al<sub>2</sub>O<sub>3</sub>/STO heterostructures depend on the Al<sub>2</sub>O<sub>3</sub> film thickness. 6 nm-thick LAO films were grown on top of Al<sub>2</sub>O<sub>3</sub> films to fabricate LAO/Al<sub>2</sub>O<sub>3</sub>/STO heterostructure. (b) The mobility of stoichiometric (51.3% of Y) YAO/STO and Y-rich (62.6% of Y) YAO/STO heterostructures as a function of YAO film thickness. YAO films were grown by ALD using cycle ratios of 2:1 (Y:Al) and 4:1, respectively.

Figure S4 Glancing angle X-ray diffraction spectra of LAO/STO, Al<sub>2</sub>O<sub>3</sub>/STO and YAO (51.3% of Y)/STO heterostructures showing that amorphous films were grown on the STO substrates.

Figure S5 The illustration of the angle-resolved XPS analysis with a 2.5 nm-thick  $Al_2O_3/STO$  heterostructure. The analysis sampling depth changes as a function of the photoelectron take-off angle. The photoelectron take-off angle ( $\Theta$ ) is defined as the angle between the sample surface and detector.

Figure S6 (a) Repetitive pulses of TMA (55 cycles) on  $Y_2O_3/STO$  heterostructure depend on the  $Y_2O_3$  film thickness without introducing an oxygen source (H<sub>2</sub>O), which did not generate the carriers. However, the carriers were generated by means of following the deposition of a 6 nm-thick LAO layer on top of the  $Y_2O_3/STO$  heterostructure. (b) A required critical thickness of LAO and  $Al_2O_3$  layers grown on top of the  $Y_2O_3$  layer (4.7 nm-thick) to create 2-DEG, which underlines the important role of LAO and  $Al_2O_3$  layers.

Figure S7 The influence of ALD growth temperature on the conductivity of LAO/STO heterostructure. The density of carriers decreased drastically with decreasing growth temperature.

Figure S8 The mobility of LAO/STO heterostructure at the low temperature (down to 30 K).

Figure S9 Ohmic I-V behavior between two Au electrodes using KEITHLEY 2400 sourcemeter.

#### References

- (1) Mannhart, J.; Schlom, D. G. Science 2010, 327, 1607-1611.
- (2) Ohtomo, A.; Hwang, H. Y. Nature 2004, 427, 423-426.
- (3) Thiel, S.; Hammerl, G.; Schmehl, A.; Schneider, C. W.; Mannhart, J. Science 2006, 313, 1942–1945.
- Xie, Y.; Bell, C.; Yajima, T.; Hikita, Y.; Hwang, H. Y. Nano Letters 2010, 10, 2588-2591.
- (5) Rijnders, G.; Blank, D. H. A. *Nat Mater* **2008**, *7*, 270–271.
- (6) Park, J. W.; Bogorin, D. F.; Cen, C.; Felker, D. A.; Zhang, Y.;
  Nelson, C. T.; Bark, C. W.; Folkman, C. M.; Pan, X. Q.; Rzchowski,
  M. S.; Levy, J.; Eom, C. B. *Nat Commun* 2010, *1*, 94.
- (7) Cen, C.; Thiel, S.; Hammerl, G.; Schneider, C. W.; Andersen, K. E.;
   Hellberg, C. S.; Mannhart, J.; Levy, J. *Nat Mater* 2008, *7*, 298–302.
- (8) Segal, Y.; Ngai, J. H.; Reiner, J. W.; Walker, F. J.; Ahn, C. H. *Physical Review B* 2009, *80*, 241107.
- (9) Nakagawa, N.; Hwang, H. Y.; Muller, D. A. Nat Mater 2006, 5, 204-209.
- (10) Chambers, S. A. Adv. Mater. 2010, 22, 219-248.
- (11) Herranz, G.; Basletić, M.; Bibes, M.; Carrétéro, C.; Tafra, E.; Jacquet, E.; Bouzehouane, K.; Deranlot, C.; Hamzić, A.; Broto, J. M.; Barthélémy, A.; Fert, A. *Physical Review Letters* 2007, *98*, 216803.
- (12) Siemons, W.; Koster, G.; Yamamoto, H.; Harrison, W. A.; Lucovsky, G.; Geballe, T. H.; Blank, D. H. A.; Beasley, M. R. *Physical Review Letters* 2007, *98*, 196802.
- (13) Chen, Y.; Pryds, N.; Kleibeuker, J. e. E.; Koster, G.; Sun, J.;
  Stamate, E.; Shen, B.; Rijnders, G.; Linderoth, S. *Nano Letters* 2011, *11*, 3774–3778.
- (14) Kalabukhov, A.; Gunnarsson, R.; Börjesson, J.; Olsson, E.; Claeson, T.; Winkler, D. *Physical Review B* 2007, *75*, 121404.
- (15) Zhang, L.; Zhou, X.-F.; Wang, H.-T.; Xu, J.-J.; Li, J.; Wang, E. G.;
   Wei, S.-H. *Physical Review B* 2010, *82*, 125412.
- (16) Zhong, Z.; Xu, P. X.; Kelly, P. J. *Physical Review B* 2010, *82*, 165127.

- (17) Schneider, C. W. E., M.; Marozau, I.; Conder, K.; Doebeli, M.; Hu, Yi; Mallepell, M.; Wokaun, A.; Lippert, T. *Appl. Phys. Lett.* 2010, *97*, 192107.
- (18) Huijben, M.; Rijnders, G.; Blank, D. H. A.; Bals, S.; Aert, S. V.; Verbeeck, J.; Tendeloo, G. V.; Brinkman, A.; Hilgenkamp, H. Nat Mater 2006, 5, 556-560.
- Qiao, L.; Droubay, T. C.; Varga, T.; Bowden, M. E.; Shutthanandan, V.; Zhu, Z.; Kaspar, T. C.; Chambers, S. A. *Physical Review B* 2011, *83*, 085408.
- (20) Liao, Y. C.; Kopp, T.; Richter, C.; Rosch, A.; Mannhart, J. *Physical Review B* 2011, *83*, 075402.
- (21) Leskelä, M.; Ritala, M. *Angewandte Chemie International Edition* **2003**, *42*, 5548–5554.
- (22) Puurunen, R. L. J. Appl. Phys. 2005, 97, 121301.
- (23) Jiang, J. C.; Pan, X. Q.; Chen, C. L. Applied Physics Letters 1998, 72, 909-911.
- (24) Kawasaki, M.; Takahashi, K.; Maeda, T.; Tsuchiya, R.; Shinohara, M.; Ishiyama, O.; Yonezawa, T.; Yoshimoto, M.; Koinuma, H. *Science* 1994, *266*, 1540-1542.
- (25) Koster, G.; Kropman, B. L.; Rijnders, G. J. H. M.; Blank, D. H. A.;
   Rogalla, H. *Applied Physics Letters* 1998, *73*, 2920–2922.
- (26) Diehl, R.; Brandt, G. Materials Research Bulletin 1975, 10, 85-90.
- (27) Frederikse, H. P. R.; Hosler, W. R. *Physical Review* 1967, 161, 822–827.
- (28) Tufte, O. N.; Chapman, P. W. Physical Review 1967, 155, 796-802.
- (29) Kim, S. K.; Lee, S. W.; Han, J. H.; Lee, B.; Han, S.; Hwang, C. S. Advanced Functional Materials 2010, 20, 2989–3003.
- (30) Liang, X.; Meng, Z.; Wu, W. Journal of the American Ceramic Society 2004, 87, 2218-2222.
- (31) Chen, T. L., X. M.; Wu, W. B. J. Appl. Phys. 2005, 98, 064109.
- (32) Lee, S. W.; Han, J. H.; Kim, S. K.; Han, S.; Lee, W.; Hwang, C. S. Chemistry of Materials 2011, 23, 976-983.
- (33) Charles T, C. Surface Science Reports 1997, 27, 1-111.
- (34) Ulrike, D. Surface Science Reports 2003, 48, 53-229.
- (35) Fu, Q.; Wagner, T. Surface Science Reports 2007, 62, 431-498.

- (36) Goldstein, D. N.; McCormick, J. A.; George, S. M. The Journal of Physical Chemistry C 2008, 112, 19530-19539.
- (37) Elliott, S. D.; Scarel, G.; Wiemer, C.; Fanciulli, M.; Pavia, G. Chemistry of Materials 2006, 18, 3764-3773.
- (38) HSC Chemistry, 5.11 edition; Outokumpu Research Oy: Pori, Finland.
- (39) George, S. M. Chemical Reviews 2009, 110, 111-131.
- Son, J.; Moetakef, P.; Jalan, B.; Bierwagen, O.; Wright, N. J.;
   Engel-Herbert, R.; Stemmer, S. Nat Mater 2010, 9, 482-484.
- (41) Jang, H. W.; Felker, D. A.; Bark, C. W.; Wang, Y.; Niranjan, M. K.; Nelson, C. T.; Zhang, Y.; Su, D.; Folkman, C. M.; Baek, S. H.; Lee, S.; Janicka, K.; Zhu, Y.; Pan, X. Q.; Fong, D. D.; Tsymbal, E. Y.; Rzchowski, M. S.; Eom, C. B. *Science* **2011**, *331*, 886-889.
- Liu, Z. Q.; Leusink, D. P.; Wang, X.; Lü, W. M.; Gopinadhan, K.;
  Annadi, A.; Zhao, Y. L.; Huang, X. H.; Zeng, S. W.; Huang, Z.;
  Srivastava, A.; Dhar, S.; Venkatesan, T.; Ariando *Physical Review Letters* 2011, *107*, 146802.

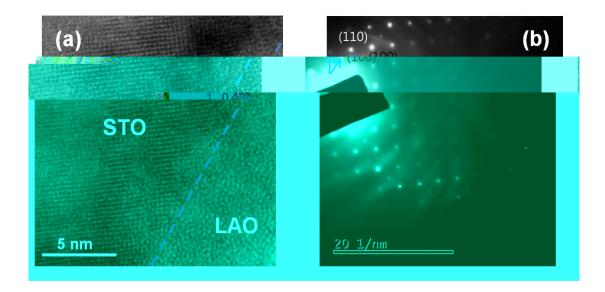


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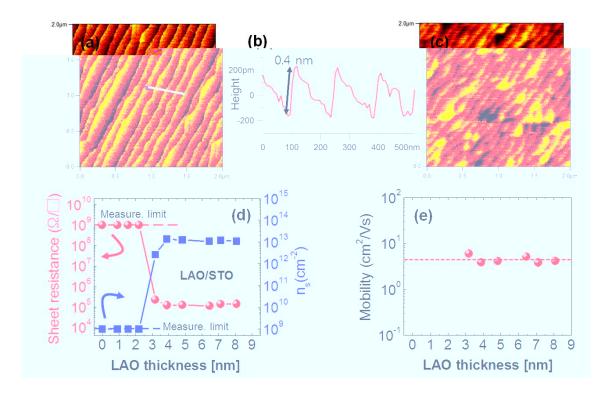


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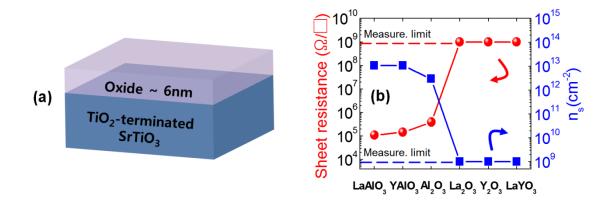


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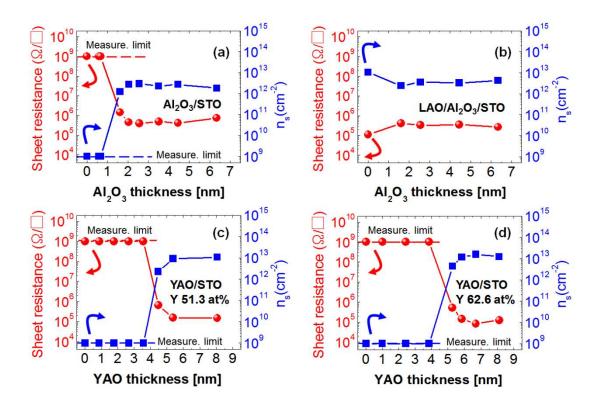


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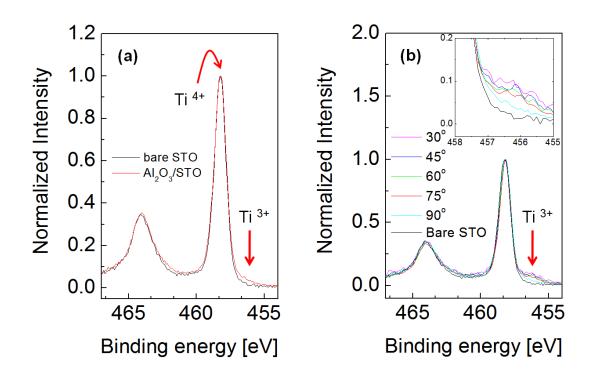


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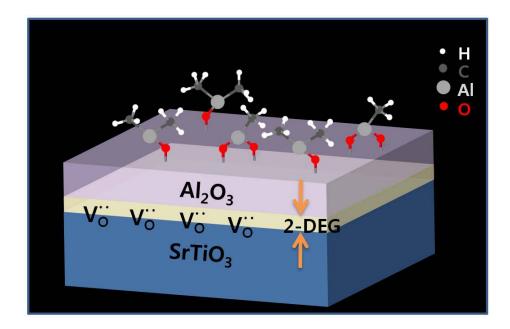


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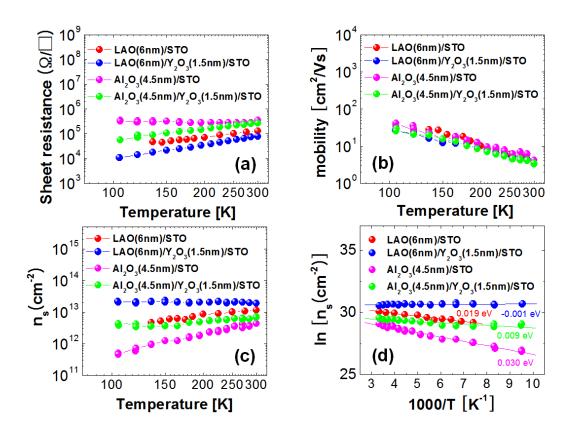


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## Supporting Information for

# Creation and Control of 2-Dimensional Electron Gas Using Al-based Amorphous Oxides/SrTiO<sub>3</sub> Heterostructures Grown by Atomic Layer Deposition

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### **Termination of STO**

The STO crystals were ultrasonically soaked in DI water for 20 min prior to the etching in BHF to make Sr-hydroxide complex on the surface of STO which is used to enhance the etch-selectivity of SrO relative to  $TiO_2$  in a buffered HF solution. And then, the STO crystals were etched in buffered HF solution for 30 s for the  $TiO_2$ -surface termination, followed by the annealed at 950 °C in  $O_2$  ambient for 1 hr in the tube furnace for the surface reconstruction. When the substrate was not treated according to the above procedure, the sheet

resistances of LAO/STO, YAO/STO and  $Al_2O_3$ /STO heterostructures were much higher than those of treated samples by a factor >1000.

#### Thin film growth.

The LAO films were grown on TiO<sub>2</sub>-terminated STO substrate at a growth temperature of 300 °C by atomic layer deposition (ALD) in a horizontal gas flow reactor at a working pressure of 400 mtorr (base pressure 30 mtorr). All films including LAO layer were grown by ALD at 300 °C. Lanthanum tris(N,N'-diisopropylformamidinate) molecules were used as the La-precursor which was kept at 118 °C, and H<sub>2</sub>O was used as the oxygen source which was kept at room temperature. The long enough purge time after each dose was set to be 40 and 80 s, respectively, for the complete removal of excess molecules and preventing of any gas phase reaction. TMA (Sigma Aldrich) was used as the Al-precursor and H<sub>2</sub>O was used as the oxygen source for the deposition of Al<sub>2</sub>O<sub>3</sub>. TMA was kept at room temperature, vaporized into a trap volume ( $\sim 3 \text{ mL}$ ) made of stainless steel and subsequently delivered to the reaction chamber by using pneumatic valves. The purge time after each dose was set to be 40 and 60 s, respectively. One super cycle consists of 2 sub-cycles of La<sub>2</sub>O<sub>3</sub> and 1 sub-cycle of Al<sub>2</sub>O<sub>3</sub> for the deposition of stoichiometric LAO films. Y2O3 thin film was deposited by ALD using yttrium tris-(N,N'diisopropylacetamidinate), Y(iPr<sub>2</sub>amd)<sub>3</sub>, and H<sub>2</sub>O at a growth temperature of 300 °C.

 $Y(iPr_2amd)_3$  molecules were kept at 130 °C. Stoichiometric YAO film (51.3 at% of Y) was grown with a super cycle ratio of 2:1 (Y:Al), and Y-rich YAO film (62.6 at% of Y) was deposited with a super cycle ratio of 4:1 (Y:Al).

#### **Composition of amorphous films (Figures S1)**

The composition of the grown LAO (17 nm) and YAO (18 nm) films were evaluated by Rutherford back-scattering spectroscopy (RBS) using 2 MeV He<sup>2+</sup> beam as shown in Figs. S1. Glassy carbon (Alfa Aesar) was used as a substrate for RBS analysis, which enables the detection of Al peak separated from the substrate background. The simulation error for atomic density was less than 3 %, which gives less than 1.5 at% variation in calculated Al concentrations. The carbon substrate was cleaned with 10 % aqueous HF (5 s), deionized water (30 s), and isopropanol (10 s) prior to drying and ultraviolet-ozone cleaning for 3 min to facilitate the initiation of the film growth. The estimated atomic ratio of La/[La+Al] in LAO film was 51.4 %, which indicates the grown film is very close to the stoichiometry. The estimated atomic ratios of Y/[Y+Al] in YAO film were 51.3 and 62.6 at% with cycle ratios (Y:Al) of 2:1 and 4:1, respectively.

#### **Estimation of amorphous film thickness (Figure S2)**

The thickness of the film grown on STO substrate was estimated by X-ray reflectivity measurement. The X-ray reflectivity data were used to evaluate the film thickness, density, and roughness by the least-squares fitting method using the analytical software, X'Pert Reflectivity 1.1. LAO, Al<sub>2</sub>O<sub>3</sub> and YAO films were grown simultaneously on Ru substrate to confirm the linear growth behavior depending on the deposition cycle number for the thinner regime which was difficult to analyze by XRR due to the weak signal. The estimated thicknesses of (a) LAO, (b) Al<sub>2</sub>O<sub>3</sub> and (c) YAO films were compared with those measured by ellipsometry, which were grown on Ru substrates simultaneously. The simulation results fit well to the raw data from LAO/STO, Al<sub>2</sub>O<sub>3</sub>/STO and YAO/STO heterostructures as shown in Fig. S2 (d). The densities were estimated to 4.85, 4.2, 4.0 g/cm<sup>3</sup> for LAO, YAO and Al<sub>2</sub>O<sub>3</sub> films, respectively. The growth rates of LAO film obtained from both ellipsometry (0.328 nm/super cycle) and XRR (0.326 nm/super cycle) were almost identical which were acquired by the slope of linear fits in Fig. S2 (a). Likewise, the growth rate of  $Al_2O_3$  films obtained from ellipsometry (0.087 nm/cycle) wase consistent with that from XRR (0.086 nm/cycle). The growth behavior of LAO and Al<sub>2</sub>O<sub>3</sub> films grown on Ru substrates were quite linear, which suggests that typical linear growth behavior occurs in these ALD processes. The growth rates of YAO films obtained from XRR (0.180 and 0.285 nm/super cycle for the cycle ratios (Y:Al) of 2:1 and 4:1, respectively) were slightly lower than those measured by ellipsometry (0.191 and 0.318 nm/super cycle, respectively), which were grown on Ru substrates.

#### **Extraction of the mobility (Figure S3)**

The electron mobility was extracted by Hall measurement using the relation between the conductance, carrier density and mobility. Figure S3 (a) shows the electron mobility of Al<sub>2</sub>O<sub>3</sub>/STO and LAO/Al<sub>2</sub>O<sub>3</sub>/STO heterostructures as a function of Al<sub>2</sub>O<sub>3</sub> film thickness, which shows constant electron mobility irrespective of the Al<sub>2</sub>O<sub>3</sub> film thickness. Figure S3 (b) shows the electron mobility of YAO/STO heterostructures as a function of YAO film thickness. The electron mobility was constant regardless of the Y composition and thickness of YAO film.

#### **Confirmation of the amorphous phase grown on STO substrate (Figure S4)**

The crystal structure of the film was examined by X-ray diffraction in glancing angle mode using a Cu K $\alpha$  X-ray source (GAXRD, PANalytical, X'pert Pro). The incidence angle, scan step size and time per step during the GAXRD measurement were 0.4 °, 0.01 ° and 0.5 s, respectively. LAO (16.5 nm), Al<sub>2</sub>O<sub>3</sub> (13.4 nm) and stoichiometric-YAO (12.6 nm) films grown on STO substrate were investigated for the analysis of the phase. No peaks were observed from all films grown on STO substrate as shown in Fig. S4, which indicates that the grown films are amorphous.

#### **Angle-resolved X-ray photoelectron spectroscopy (Figure S5)**

The illustration of the angle-resolved XPS analysis with 2.5 nm-thick Al<sub>2</sub>O<sub>3</sub>/STO heterostructure was described in Fig. S5. The analysis sampling depth changes as a function of the photoelectron take-off angle. The photoelectron take-off angle  $(\Theta)$  is defined as the angle between the sample surface and detector, thus the analysis becomes more surface sensitive with decreasing take-off angle. The normal signal ( $\Theta = 90^{\circ}$ ) from the bare STO substrate was used for a clear comparison in Fig. 5. The signal of Ti<sup>3+</sup> relative to the Ti<sup>4+</sup> main line increases as the photoelectron take-off angle decreases in Al<sub>2</sub>O<sub>3</sub>/STO heterostructure. This indicates that the electrons generated by oxygen vacancies are localized at the STO surface region as a 2-DEG (within  $\sim 2$  nm from the interface, not surface of the 2.5 nm-thick Al<sub>2</sub>O<sub>3</sub>), which is smaller than the electron escape depth (~5 nm). The electron escape depth (~5 nm) was estimated experimentally. The Ti 3d peak was not detected with a 6-nm thick film grown on a STO substrate, however, very weak Ti 3d peak was appeared after an etching of the oxide surface with 1 nm-depth, thus the maximum electron escape depth was expected to be  $\sim 5$  nm. The intensity of Ti<sup>3+</sup> is very weak because the carrier density of Al<sub>2</sub>O<sub>3</sub>/STO heterosture  $(3x10^{12}/cm^2)$ is lower over 1-order of magnitude than any other reports using PLD ( $\sim 10^{14}$ /cm<sup>2</sup>). The density of  $3x10^{12}$ /cm<sup>2</sup> corresponds to 0.5 % of number of Ti atoms in 1 monolayer of STO, which makes low intensity of Ti<sup>3+</sup> respect to Ti<sup>4+</sup>.

#### Role of TMA precursor and pre-grown LAO and Al<sub>2</sub>O<sub>3</sub> layers (Figure S6)

The tri-methyl aluminum [Al(CH<sub>3</sub>)<sub>3</sub>, (TMA)], TMA molecules were introduced on the surface of Y<sub>2</sub>O<sub>3</sub>/STO heterostructure at 300 °C in order to confirm that TMA molecules can induce the redox reaction of STO. It should be noted that the interface of Y<sub>2</sub>O<sub>3</sub>/STO heterostructure is insulating as shown in Fig. 3 (b). This experiment was progressed with an identical ALD process of LAO or Al<sub>2</sub>O<sub>3</sub> except that the valve for water injection was closed during the experiment, so that the total process time was preserved. However, the STO substrate keep insulating property ( $R_s > 10^9 \Omega/sq.$ ) after the introducing of TMA molecules (55 cycles, which is much longer than the critical cycle  $\sim$ 15 cycles) on bare STO substrate, and also on Y<sub>2</sub>O<sub>3</sub>/STO heterostructure as shown Fig. S6 (a). The repetitive pulse (55 cycles) of TMA precursor cannot induce the reduction of STO substrate by itself, and also on Y<sub>2</sub>O<sub>3</sub> layers grown on STO substrate irrespective of Y<sub>2</sub>O<sub>3</sub> film thickness, which means that the redox reaction of STO cannot be progressed until when LAO, YAO and Al<sub>2</sub>O<sub>3</sub> films are thicker than those critical thicknesses. There still exists required critical thickness of LAO and Al<sub>2</sub>O<sub>3</sub> layers when those layers were grown on top of Y<sub>2</sub>O<sub>3</sub> layer (4.7 nm-thick) to create 2-DEG, which emphasizes the important role of pre-grown LAO and Al<sub>2</sub>O<sub>3</sub> layers (Fig. S6 (b)). LAO and Al<sub>2</sub>O<sub>3</sub> layers (thicker than the critical thickness) are necessary for the redox reaction even on Y<sub>2</sub>O<sub>3</sub>/STO heterostructure, thereby it constitutes the circumstance to overcome the energy barrier, so that the reduction of STO can occur during the TMA pulse. We claimed that  $Al_2O_3$  (>1.2 nm) acts as a catalyst for the reaction, thus the redox reaction of STO substrate can occur when the TMA molecules are introduced.

#### **Influence of the growth temperature (Figure S7)**

The influence of ALD growth temperature on the conductivity of LAO/STO heterostructure was investigated in order to confirm that the reaction is governed by kinetics. The density of carriers decreased drastically with decreasing growth temperature, which suggests that the reaction is limited by kinetics.

#### Mobility of LAO/STO at the low temperature (down to 30 K). (Figure S8)

The mobility of LAO/STO heterostructue was measured at the low temperature down to 30 K. The measuring temperature was controlled by means of liquid helium, which allowed the lowest temperature of the stage down to 30 K. The mobility of LAO/STO heterostructure was further increased with decreasing temperature down to 30 K.

**Confirm of ohmic contact (Figure S9)** 

In order to confirm the contact properties, I-V measurement was performed with the two Au contacts at the corners using KEITHLEY 2400 sourcemeter. The linearity of I-V curve demonstrates that ohmic contacts were made between two Au contacts as shown in Fig. S8.

#### **Stability of 2-DEG in the heterostructure**

The sheet resistance of the LAO/STO heterostructure was increased by  $>10^3$  in 8 months (in atmosphere at room temperature), from the fabrication of the heterostructure. However, the LAO/Y<sub>2</sub>O<sub>3</sub>/STO heterostructure kept conducting channel for much longer period than that of the LAO/STO, i. e., the sheet resistance of the LAO/Y<sub>2</sub>O<sub>3</sub>/STO heterostructure was increased by just 10 % in 8 months (in atmosphere at room temperature), from the fabrication of the heterostructure.

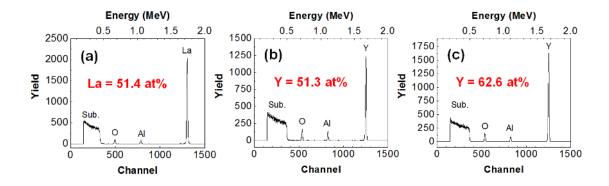


Figure S1 The composition of (a) LAO film (17 nm-thick) and YAO (18 nm) films with cycle ratios (Y:Al) of (b) 2:1 and (c) 4:1, evaluated by Rutherford back-scattering spectroscopy (RBS) using a 2 MeV He<sup>2+</sup> beam. The estimated atomic ratio of La/[La+Al] in the LAO film was 51.4 %, which indicates the grown film is very close to stoichiometry The estimated atomic ratios of Y/[Y+Al] in YAO films were 51.3 and 62.6% with cycle ratios of 2:1 and 4:1, respectively.

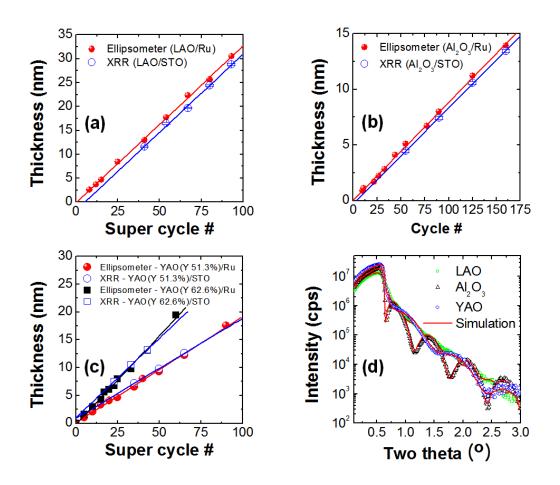


Figure S2 The thickness of (a) LAO, (b)  $Al_2O_3$  and (c) YAO films grown on an STO substrate estimated by X-ray reflectivity measurement. Films were grown on Ru substrates simultaneously and the thickness was measured by ellipsometry to confirm the linear growth behavior depends on the deposition cycle number. (d) Representative simulation results of LAO,  $Al_2O_3$  and YAO films grown on the STO substrate.

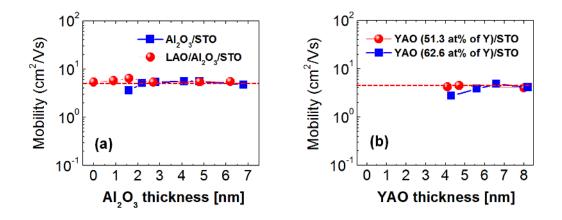


Figure S3 (a) The mobility of Al<sub>2</sub>O<sub>3</sub>/STO and LAO/Al<sub>2</sub>O<sub>3</sub>/STO heterostructures depend on the Al<sub>2</sub>O<sub>3</sub> film thickness. 6 nm-thick LAO films were grown on top of Al<sub>2</sub>O<sub>3</sub> films to fabricate LAO/Al<sub>2</sub>O<sub>3</sub>/STO heterostructures. (b) The mobility of stoichiometric (51.3% of Y) YAO/STO and Y-rich (62.6% of Y) YAO/STO heterostructures as a function of YAO film thickness. YAO films were grown by ALD using cycle ratios of 2:1 (Y:Al) and 4:1, respectively.

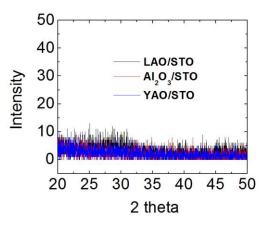


Figure S4 Glancing angle X-ray diffraction spectra of LAO/STO, Al $_2O_3$ /STO and YAO (51.3%

of Y)/STO heterostructures showing that amorphous films were grown on the STO substrates.

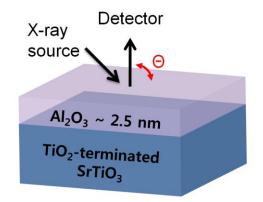


Figure S5 The illustration of the angle-resolved XPS analysis with a 2.5 nm-thick  $Al_2O_3/STO$  heterostructure. The analysis sampling depth changes as a function of the photoelectron take-off angle. The photoelectron take-off angle ( $\Theta$ ) is defined as the angle between the sample surface and detector.

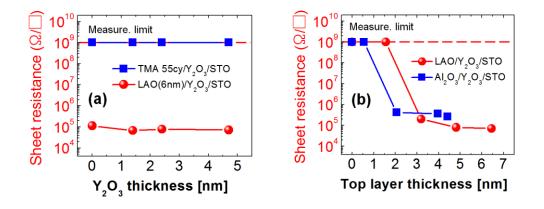


Figure S6 (a) Repetitive pulses of TMA (55 cycles) on  $Y_2O_3/STO$  heterostructure depend on the  $Y_2O_3$  film thickness without introducing an oxygen source (H<sub>2</sub>O), which did not generate the carriers. However, the carriers were generated by means of following the deposition of a 6 nm-thick LAO layer on top of the  $Y_2O_3/STO$  heterostructure. (b) A required critical thickness of LAO and  $Al_2O_3$  layers grown on top of the  $Y_2O_3$  layer (4.7 nm-thick) to create 2-DEG, which underlines the important role of LAO and  $Al_2O_3$  layers.

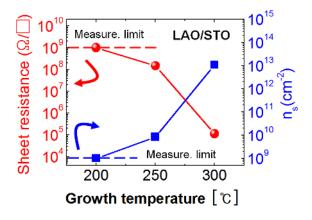


Figure S7 The influence of ALD growth temperature on the conductivity of LAO/STO heterostructure. The density of carriers decreased drastically with decreasing growth temperature.

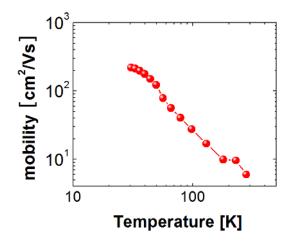


Figure S8 The mobility of LAO/STO heterostructure at the low temperature (down to 30 K).

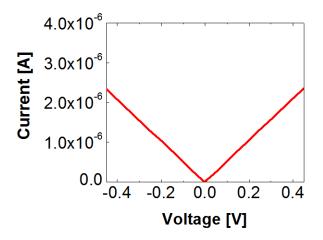


Figure S9 Ohmic I-V behavior between two Au electrodes using KEITHLEY 2400 sourcemeter.