

# Current updates of the water-vapor line list in *HITRAN*: A new “Diet” for air-broadened half-widths

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## Abstract

The current edition of the *HITRAN* compilation employed a sophisticated algorithm for combining measurements available for the air-broadened half-widths of water-vapor absorption lines with theoretical values. Nevertheless, some of the values in the *HITRAN* database were found to be far from ideal, due to large dispersions that still exist in the experimental or theoretical methods. Therefore, new criteria were developed for introducing the best available air-broadened half-widths into *HITRAN*, based on physical principles and statistics. This update concerns the three most abundant isotopologues of water, with the values for H<sub>2</sub><sup>17</sup>O and H<sub>2</sub><sup>18</sup>O being the ones from analogous transitions of the principal isotopologues. The new parameters have been tested in different remote-sensing applications and improved constituent profiles were obtained. In total, air-broadened half-width values were updated for 11,787 transitions of water vapor in the *HITRAN* database (6789 for H<sub>2</sub><sup>16</sup>O, 2906 for H<sub>2</sub><sup>17</sup>O, and 2092 for H<sub>2</sub><sup>18</sup>O). Some additional updates to the water-vapor line list are also presented. The resultant file (01\_hit06.par) was uploaded to the *HITRAN* website (<http://www.cfa.harvard.edu/hitrان/>) in September 2006.

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## 1. Introduction

The line-by-line portion of the *HITRAN* database [1] consists of high-resolution spectroscopic parameters for 39 molecules of atmospheric interest, including many of their isotopologs. For every transition, *HITRAN* provides the vacuum line position, the line intensity, the Einstein *A*-coefficient and statistical weights, the air-broadened half-width ( $\gamma_{\text{air}}$ ), the self-broadened half-width, the lower state energy, the temperature-dependence

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exponent of  $\gamma_{\text{air}}$ , the air pressure-induced line shift, and lower and upper state vibrational and rotational quantum numbers.

The data for water vapor are very important for atmospheric sciences. Water vapor is the principal absorber of longwave radiation in the terrestrial atmosphere and it has a profound effect on the atmospheric energy budget in many spectral regions. The *HITRAN* database lists more than 64,000 significant transitions of water vapor ranging from the microwave region to the visible, with intensities that cover many orders of magnitude. These transitions are used, or have to be accounted for, in various remote-sensing applications.

Out of all water-vapor spectroscopic parameters in *HITRAN*, the intensity of weak lines and the overall pressure-broadening coefficients are the largest sources for uncertainty in remote-sensing retrievals [2]. For  $\text{H}_2\text{O}$ , the half-width parameters have the largest dynamic range of any molecule contained in *HITRAN* and they contribute on a par with the intensities to the radiance and transmission simulations in regimes of tropospheric pressure, i.e. where collision-broadening by air is significant. For accurate retrievals that are achievable with the high signal-to-noise and wide spectral coverage of current satellite-borne experiments, it is required to know the half-width and its temperature dependence better than a 3% uncertainty for strong lines and 10% for weak lines [3]. For the current compilation [1] (hereafter referred to as *HITRAN2004*), a great effort had been made to provide the most accurate value of the air-broadened half-width,  $\gamma_{\text{air}}$ , for every transition of the  $\text{H}_2^{16}\text{O}$ ,  $\text{H}_2^{18}\text{O}$ , and  $\text{H}_2^{17}\text{O}$  isotopologues of water. The algorithm used data from several theoretical, experimental, and semi-empirical datasets. The hierarchy of the sources from which  $\gamma_{\text{air}}$  was determined is described below.

### 1.1. The database of experimental measurements

The database of experimental measurements of collision-induced parameters was created by Gamache and Hartmann (hereafter referred to as GH database) [3]. It lists the vast majority of the experimental data from different sources, reported prior to the publication of Ref. [3]. The dataset spans the region  $0\text{--}22,640\text{ cm}^{-1}$ , listing values for  $\gamma_{\text{air}}$  (with reported experimental uncertainties) from over 40 sources. The  $\gamma_{\text{air}}$  values for over 3000 transitions have been measured more than once and those for over 6000 transitions have been measured only once. Overall there are more than 14,000 entries in the GH database.

### 1.2. Smoothed values

A database of smoothed values for collision-induced parameters has been created by Toth and is available on his website (<http://mark4sun.jpl.nasa.gov/data/spec/H2O>). The smoothing procedure is explained in detail in Refs. [4,5]. The transitions from  $600\text{ to }8000\text{ cm}^{-1}$  measured in Refs. [4,5] were grouped into subsets with  $v_2' = 0$  and  $v_2' > 0$  and then least-squares fitted to the empirical function (Eq. (4) in Ref. [5]) in terms of “families” of rotational transitions.

### 1.3. The database of values calculated using the complex Robert–Bonamy method

The complex Robert–Bonamy (CRB) method [6] was applied in Refs. [7,8] to calculate air-broadened half-widths of water vapor. A compilation of these calculations is available at [http://faculty.uml.edu/Robert\\_Gamache](http://faculty.uml.edu/Robert_Gamache). The calculations are obtained for 6040 transitions that involve states with  $J \leq 18$  in the  $0\text{--}3810\text{ cm}^{-1}$  region. The details of the calculations are given in the aforementioned references. The CRB values in general agree well with experiments [7,8], except for some high- $J$  transitions where the comparison is not always informative since such transitions are usually weak in experimental spectra and hence accurate experimental data are limited. In turn, the CRB calculations are not expected to be as accurate at high  $J$ 's due to the higher uncertainty in the wavefunctions and energies obtained from diagonalizing the Watson Hamiltonian [9]. While the experimental  $\gamma_{\text{air}}$  for transitions with  $|\Delta K| > 1$  are also limited or lack accuracy due to the relatively low intensity of such transitions, one should expect the accuracy of CRB values not to suffer since the information needed for the calculations is that of the states and is not dependent on the  $\Delta K$  of the transition.

#### 1.4. Database of semi-empirical values

In the semi-empirical treatment of Jacquemart et al. [10], experimental data from Refs. [11–15] and theoretical (CRB) data from Refs. [7,16,17] for each transition were taken and then separated into subsets of the data for transitions that would involve the same lower and upper state rotational quantum numbers, but where the vibrational quantum numbers were not necessarily the same. These data were then fit to the equation that describes the vibrational dependence of  $\gamma_{\text{air}}$  [15]:

$$\gamma[(v'_1, v'_2, v'_3)f \leftarrow (v''_1, v''_2, v''_3)i] = \gamma_{f \leftarrow i}^0 + A_{f \leftarrow i}(0.3\Delta v_1 + 0.7\Delta v_2 + 0.3\Delta v_3)^2, \quad (1)$$

where  $v_i$  represents the quantum numbers associated with the normal mode of vibration  $i$ . The prime and double prime are used, respectively, for the upper and lower levels of the transition, and  $\Delta v_i$  is equal to  $v'_i - v''_i$ . For water vapor, the notations  $i$  and  $f$  correspond, respectively, to the rotational quantum numbers ( $J'', K''_a, K''_c$ ) and ( $J', K'_a, K'_c$ ).  $\gamma_{f \leftarrow i}^0$  is equivalent to the half-width for a pure rotational transition which corresponds to  $(v'_1 = 0, v'_2 = 0, v'_3 = 0)f \leftarrow (v''_1 = 0, v''_2 = 0, v''_3 = 0)j$ . The coefficients  $\gamma_{f \leftarrow i}^0$  and  $A_{f \leftarrow i}$  deduced from the fit allow one to obtain any air-broadening coefficient of transitions having the same rotational quantum numbers but different vibrational quantum numbers. Obvious outliers were eliminated from the fit.

The algorithm used for compiling the air-broadened half-width data in *HITRAN2004* would first search the GH database: if  $\gamma_{\text{air}}$  for the transition was measured more than once the average of all experimentally determined values was taken, if it was measured just once the measurement was taken. It is worth noting that the data from Refs [13,15,18] were not included into the algorithm if they differed from the corresponding values in the semi-empirical database [10] by more than 20%. If the  $\gamma_{\text{air}}$  value for a given rotational–vibrational transition did not exist in the GH database, the algorithm then searched the smoothed-values database. If the transition was still not found, the search was extended to the database of CRB values. Finally, if a value for a particular transition was not found in any of the aforementioned databases, the  $\gamma_{\text{air}}$  value was derived using the semi-empirical approach of Jacquemart et al. [10]. For the transitions of  $\text{H}_2^{17}\text{O}$  and  $\text{H}_2^{18}\text{O}$  isotopologues, the values from corresponding transitions of  $\text{H}_2^{16}\text{O}$  were adapted if there were no direct measurements.

Overall, the above algorithm has provided a complete set of air-broadened half-widths for all assigned lines in *HITRAN2004*. The unassigned lines were given a default value of  $0.07 \text{ cm}^{-1} \text{ atm}^{-1}$ . Nevertheless, the values of the air-broadened half-width in *HITRAN2004* were found to be far from ideal when applied to some atmospheric transmission experiments [19,20] (see Section 3 for details). The reason that the sophisticated algorithm for adding air-broadened half-widths of water-vapor transitions into the *HITRAN2004* database does not always yield an optimum value is simply due to large dispersions in the experimental or theoretical methods that still exist. The experimental spectra are sometimes hard to interpret due to many reasons such as line overlaps, impurities in the cell, etc. The CRB calculations, despite being quite accurate in most of the cases, are still far from perfect due to the different approximations. Both theory and experiment are not accurate when dealing with weak lines with high- $J$  values. The semi-empirical calculations [10] are inheriting the problems of the experimental and theoretical sources even though obvious outliers were eliminated in the course of that work.

In this paper, we will describe the problems arising from using the *HITRAN2004* list in retrievals of the constituent profiles, the sources of the problems, and a new algorithm that yields a better dataset. The new dataset was validated in different remote-sensing missions. Other updates to the *HITRAN2004* water-vapor file will be discussed as well. In general, this manuscript explains the new parameters in the 01\_hit06.par file that were uploaded to the HITRAN website in September 2006.

## 2. An improved algorithm

In order to improve the  $\gamma_{\text{air}}$  values of water vapor, it was decided to create a hierarchical scheme that would favor one source over another based not on the general quality of the source but on its quality as applied to a particular transition or “family” of transitions. This approach is widely applied by pediatricians in order to control a patient’s weight. Every patient has a “normal weight” that can be achieved by losing or gaining

weight, but the recommended diet for every person is different as everyone has a different metabolism, daily activities, etc. We therefore called our new algorithm for adding air-broadened half-widths of water-vapor lines to the *HITRAN* database a “Diet”. In this section, the development of the Diet will be discussed and the steps of the algorithm will be summarized.

The GH database [3] aggregates the results of all experimental works that were performed before 2004, including ones dating back to the 1930s. It would be naive to assume that the quality of all experiments and their interpretation was ideal, and, as a first step, it was decided to find and omit the sources that provided consistently inaccurate values.

Fig. 1 shows the ratios of experimental  $\gamma_{\text{air}}$  values reported in different sources [4,12,21–30] in the 1000–2000  $\text{cm}^{-1}$  region (as an example) to corresponding values calculated by the CRB method plotted against the values of CRB  $\gamma_{\text{air}}$ . The values from a certain source are shown only if there were at least five measurements in the selected wavenumber range. While at low values of  $\gamma_{\text{air}}$  (effectively, transitions with high  $J$ ) the ratios are not expected to be necessarily very close to 1 since both experiments and theoretical calculations lack accuracy there, at higher values of  $\gamma_{\text{air}}$  the ratios are expected to be close to 1. One can see that data from larger datasets such as Zou and Varanasi [12] and Toth [4] agree well with CRB calculations (with rare exceptions), while those from some other references such as Refs. [21,22] often deviate. We have also examined the plots created by Gamache and Hartmann when their database was assembled [3]. Fig. 2 is an example of such an investigation. In Fig. 2 one can see air-broadened half-widths measured for four transitions by three or four independent laboratories, namely the works of Toth [4], Zou and Varanasi [12], Chang and Shaw [21], and Nicolaisen [24]. The error bars represent experimental uncertainties reported in these works. The dashed line is the average half-width that was put into the *HITRAN*2004 database for these transitions. The numbers AD and Max in the boxes are the average percent differences between any two measurements and the maximum percent difference between any two measurements. It is worth noting that difference is stated in percents of one of the values rather than of the average (as is more common in these types of comparisons). In all four plots (and through most of the dataset), the measurements performed by Chang and Shaw [21] disagree with those of Zou and Varanasi [12] and Toth [4], which are very consistent in general.

It is also clear from Figs. 2 and 3 that it is difficult to correctly estimate experimental uncertainties, as often error bars of different measurements do not overlap.

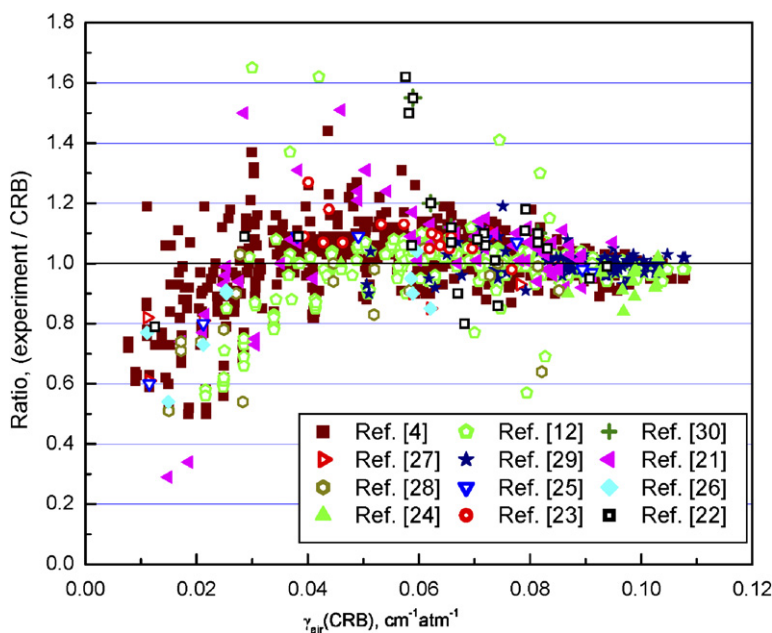


Fig. 1. The ratios of experimental  $\gamma_{\text{air}}$  to the corresponding CRB  $\gamma_{\text{air}}$  plotted against CRB  $\gamma_{\text{air}}$ .

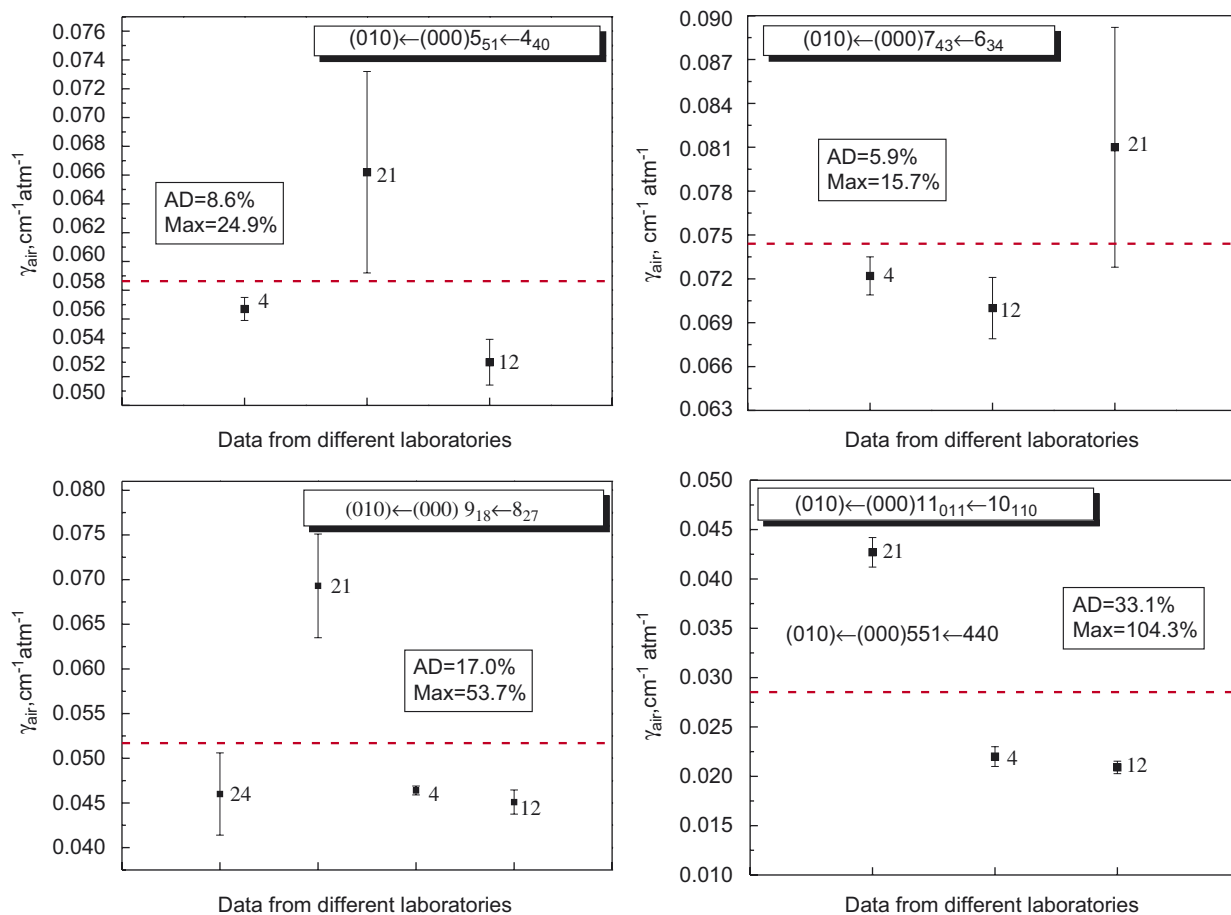


Fig. 2. An example of consistent disagreement of one source of experimental values with others. Upper and lower states involved in the transitions are given in the shadowed boxes as  $(v'_1 v'_2 v'_3) \leftarrow (v''_1 v''_2 v''_3)$ ,  $J'_{K_a' K_c'} \leftarrow J''_{K_a'' K_c''}$ . See text for other details.

After examination of the quality of data from different sources, the data from Refs. [21,22,28,31–35] were omitted from the GH database. The lower quality of some of the data in these sources is corroborated by Toth in Table 5 of Ref. [5]. After a final list was created, it was realized that Ref. [32] in general contains very good data, but the values from this work had not undergone a proper units conversion when incorporated into the GH database and were therefore making a bad contribution to the average. Nevertheless, there are plenty of other accurate measurements existing for the transitions studied in Ref. [32] and the quality of the final half-widths that were incorporated into the new *HITRAN* list had not suffered by omission of these measurements. The values from Ref. [32] are not included in Fig. 1. It should be stated that not all of the measurements in the above references are necessarily inaccurate.

Sometimes, however, it was hard to determine whether one or another experimental work was reliable, since it produced an equal amount of good- and bad-quality data. In Fig. 3a the measurement from Brown et al. [14] agrees better with that from Schermaul et al. [18] than with the one from Merienne et al. [13], whereas in Fig. 3b the measurement from Ref. [14] agrees better with that from Ref. [13] than with the one from Ref. [18]. This is a very common situation observed when scrutinizing the plots. To solve this problem, it was decided to compare the values in the GH database with those from semi-empirical calculations of Jacquemart et al. [10]. If the value in the GH database was more than 20% different from the semi-empirical value for the same transition, it was omitted from the Diet. The quality of the semi-empirical value is strongly dependent on the level of approximation used in the calculations in Ref. [10] (see Ref. [10] for details) and is also affected by the amount and quality of the input data. Nevertheless, it is considered to be within 20% of the real value of  $\gamma_{\text{air}}$

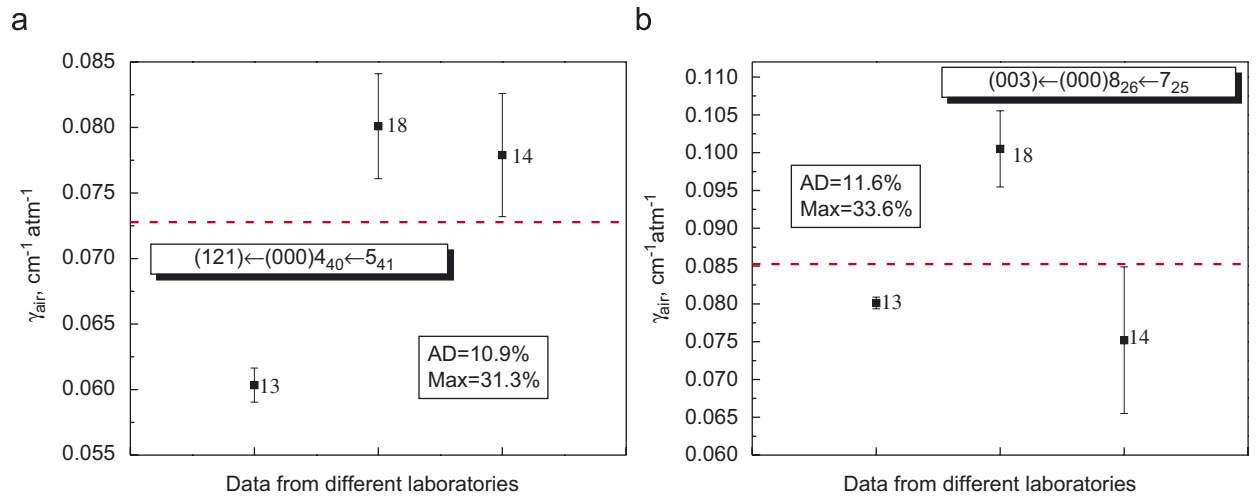


Fig. 3. Illustration of the difficulty of determining which source is more reliable. Upper and lower states involved in the transitions are given in the shadowed boxes as  $(v'_1 v'_2 v'_3) \leftarrow (v''_1 v''_2 v''_3)$ ,  $J'_{K_a' K_c'} \leftarrow J''_{K_a'' K_c''}$ . See text for details.

Table 1

Definition of  $U$  depending on the number of measurements in the GH database

$N$	Definition of $U$
$N > 2$	$U$ is an absolute percent difference between the average value and the measurement that deviates the most from it.
$N = 2$	$U$ is a percent difference between measured values with respect to the smaller value out of the two measurements.
$N = 1$	$U$ is a reported experimental uncertainty.

for most of the transitions. We also recommend the use of the semi-empirical dataset when one conducts new experiments, i.e. the new measurements should be expected to be close to the semi-empirical values and if they are not, a careful reanalysis of the particular measurement is in order.

However, 20% accuracy is not satisfactory and further treatment of data is required. The criterion of the data quality was an uncertainty ( $U$ ), which is defined differently depending on the number of measurements available for a particular transition. The definitions are given in Table 1.

The Diet starts with searching the GH database and it takes an experimental value (or average of experimental measurements if more than one exists) if the following requirements are satisfied:

- (i) for  $J_{\text{ave}} \leq 8$ ,  $U < 7\%$ ,
- (ii) for  $8 < J_{\text{ave}} \leq 13$ ,  $U < 10\%$ ,
- (iii) for  $J_{\text{ave}} > 13$ ,  $U < 15\%$ ,

where  $J_{\text{ave}} = (J' + J'')/2$ .

If the above conditions do not hold, then the experimental average is disregarded and the CRB value is used. If the theoretical value does not exist for this transition, or at  $J > 8$  either the upper or the lower state is near the prolate limit defined by  $(K_a/J) > 0.75$ , the semi-empirical value is used.

The  $J$ -dependence of the above constraints on uncertainty is necessary due to the fact that the air-broadened half-widths are getting narrower with increase of  $J$  while experimental limitations are the same. Therefore, the uncertainty  $U$  (in %) is increasing with  $J$  as well, assuming the same quality, although the experimental values may be lacking quality at higher  $J$  values, as corresponding transitions are generally weak. Likewise, the quality of the CRB calculations decreases with the increase of  $J$  as well. Several constraint limits (both on  $J$ 's and uncertainties) were tested and it seems that the limits above yield the best result.

For the case when transitions were measured three or more times, it may seem questionable why all measurements are disregarded if the requirements above are not satisfied, instead of eliminating only the measurements that are farthest from the average. The reason is that unfortunately in some cases the outlying measurement may actually be more accurate than the other measurements of the same transition despite their

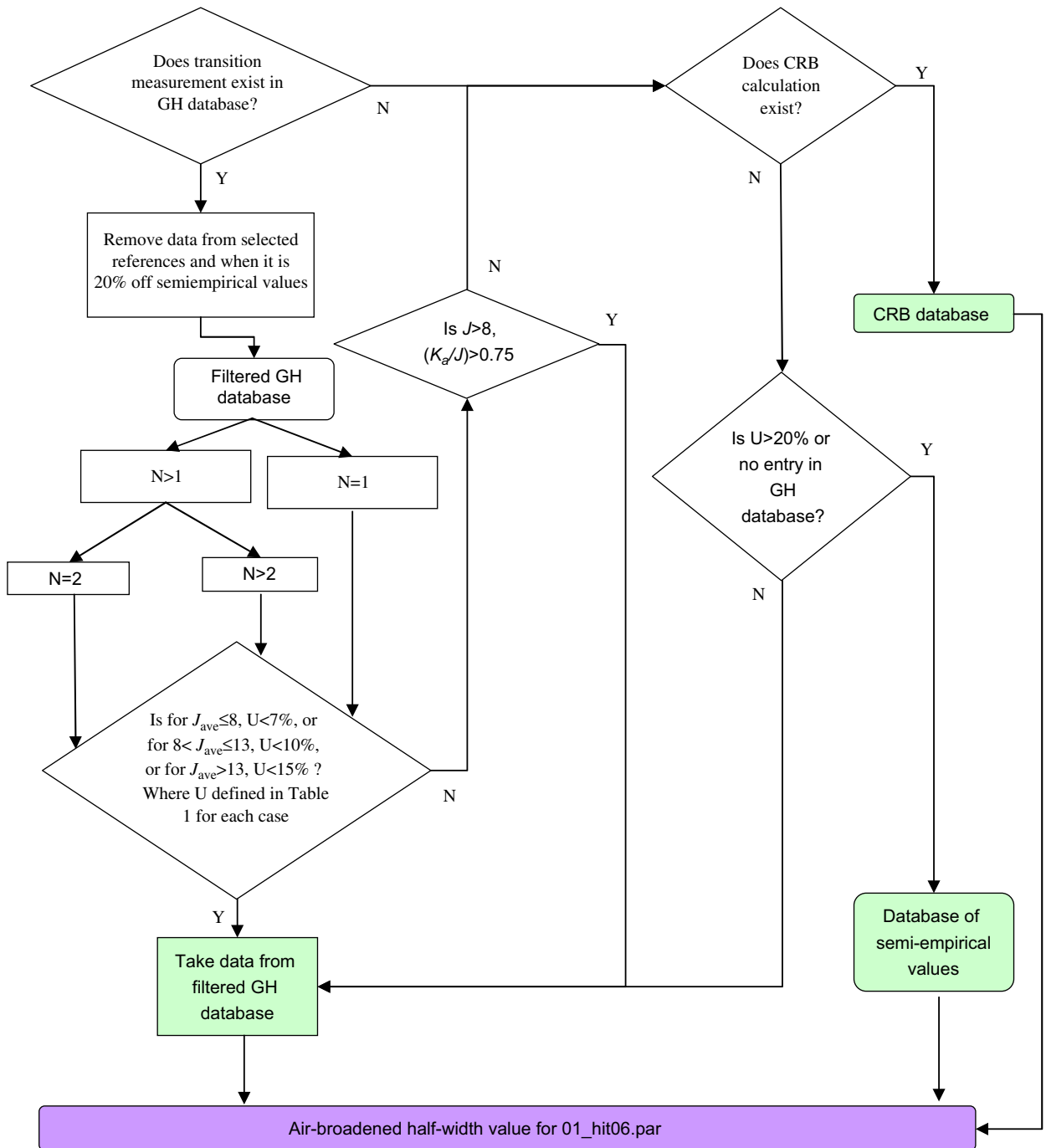


Fig. 4. A flow diagram of the new algorithm for choosing air-broadened half-widths from different sources.  $N$  is the number of measurements of air-broadened half-width for a transition.

agreement between each other. Besides, in most of the cases, three or more measurements are only available for moderate  $J$  values for which CRB calculations are expected to be accurate within 5%.

The Diet is summarized in the flow diagram given in Fig. 4. Fig. 5 shows the percent differences in the principal isotopologue  $\gamma_{\text{air}}$  values of the *HITRAN2004* database and result of the current work (01\_hit06.par file available on the *HITRAN* website). The transitions for which  $\gamma_{\text{air}}$  was unchanged as well as the lines that have changed by more than 120% are not shown, and different symbols are used for transitions with different  $\Delta K$  to give the reader an idea of the relative intensities of the plotted transitions. Below  $8000\text{ cm}^{-1}$  one can see a large amount of significant differences, while above that wavenumber the difference does not exceed 20% (with the exception of three transitions around  $12,000\text{ cm}^{-1}$ ), which is not surprising since in *HITRAN2004* the data from Refs. [13,15,18] were already tested against the semi-empirical calculations, for 20% agreement. For the three outlying transitions,  $\gamma_{\text{air}}$  values in *HITRAN2004* were originating from the experimental work of Lucchesini et al. [36] and apparently these values were significantly different from semi-empirical predictions. The listing of all transitions for which widths have been changed can be found in the supplementary file.

Although the above algorithm proves to be very efficient, one should not expect the accuracy of the  $\gamma_{\text{air}}$  values to be better than 5% (although in many cases they are). Therefore, if better accuracy is needed and potentially achievable, one has to determine the best value manually for a particular case. For example, for the important microwave transition at 22 GHz the Diet chooses the value of  $0.0942\text{ cm}^{-1}\text{ atm}^{-1}$ , which is an average of three experimental measurements ( $0.0918\text{ cm}^{-1}\text{ atm}^{-1}$  [37],  $0.0965\text{ cm}^{-1}\text{ atm}^{-1}$  [38], and  $0.0942\text{ cm}^{-1}\text{ atm}^{-1}$  [39]) that have cleared the filtering criteria because they were relatively close to each other, whereas the value from CRB calculations is  $0.0920\text{ cm}^{-1}\text{ atm}^{-1}$  and the value from the corresponding rotational transition in the  $\nu_2$  band is  $0.0909\text{ cm}^{-1}\text{ atm}^{-1}$  (current update). This suggests that the current value in the update is overestimated. For the future *HITRAN* edition, we will adapt the value from the newest CRB calculations ( $0.0918\text{ cm}^{-1}\text{ atm}^{-1}$ ) that include the explicit determination of the velocity integral. Similarly, the new CRB value will be adapted for another important microwave line at 183 GHz. These values have been applied in the MonRTM radiative transfer model and yielded a good agreement with ground-based radiometric measurements from atmospheric radiation measurement (ARM) sites in Oklahoma and Alaska [40].

Another important point to make here is that, in contradiction to our original point of view, the older data (say, before 1960) are not necessarily worse than newer data. Ref. [37] (1945) for example seems to provide a

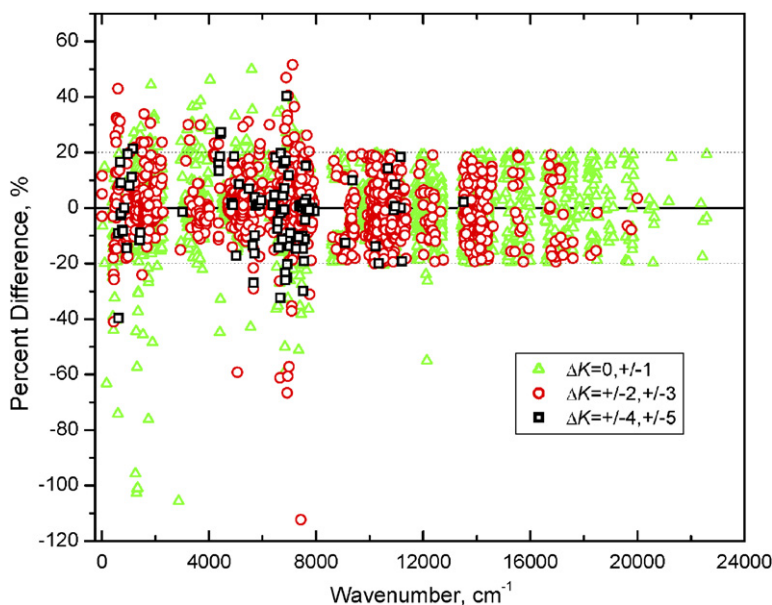


Fig. 5. Percent difference  $100 \times (\gamma_{\text{air}}(\text{HITRAN2004}) - \gamma_{\text{air}}(\text{Update})) / \gamma_{\text{air}}(\text{HITRAN2004})$  plotted against the wavenumber of the corresponding transition.



more accurate value than more recent (1969 and 1970) measurements in Refs. [38,39]. Therefore, it was decided not to remove the data from older references from the GH database as was planned originally.

### 3. Validation of the new algorithm

The new line list that was created in the course of this work was tested in application to several remote-sensing missions. Below some of these missions are described and results of their validation of our new line list are presented.

#### 3.1. Ground-based measurements: atmospherically emitted radiance interferometer (AERI-ER)

The water vapor  $\nu_2$  spectral region is commonly used for passive infrared remote sensing due to its large opacity range. The far-infrared pure rotational water-vapor band is an important spectral region for earth energy budget climate studies. The pure rotational region has a larger contribution to the longwave cooling rates [41], as it is closer to the peak of the blackbody radiance curve at terrestrial atmospheric temperatures. Therefore, it is important that the radiative transfer calculation, especially the spectroscopic parameters and the continuum, be as accurate as possible in the far-infrared region. Due to the limited utilization of the far infrared for remote sensing, there are not a lot of instruments monitoring the atmosphere in the infrared that can be used for validating the radiative transfer calculations in the far-infrared region. The Atmospheric Radiation and Measurement (ARM) program operates a ground-based extended range AERI-ER at the North Slopes of Alaska (NSA) that measures down to  $400\text{ cm}^{-1}$  in the infrared. Since the vertical distribution of water vapor in the atmosphere is significantly weighted towards the surface, the ground-based observations in the strong-absorbing pure-rotation band must be performed under conditions of very low atmospheric water-vapor loading in order for the lines not to become opaque close to the instrument. In addition, since the measurements are ground-based, the pressure-broadened widths are important in the line-by-line radiance calculations.

Presented in Fig. 6 is an example of a comparison of the Line-By-Line Radiative Transfer Model (LBLRTM) [42] calculations with AERI-ER observations in the far infrared that demonstrates the impact of recent versions of the water-vapor spectroscopic lines in the *HITRAN* database. The AERI-ER NSA downwelling radiance observations on March 11, 2001 were performed at the surface with a spectral resolution of  $0.48\text{ cm}^{-1}$ . The temperature and water-vapor profiles used in the LBLRTM calculation were obtained from a radiosonde launched from the NSA ARM site, which was coincident and co-located with the AERI-ER measurements. In order to account for inaccuracies in the radiosonde water-vapor measurements [43], the water-vapor profile was scaled with a retrieved total column water vapor. This case has a very low water-vapor loading of  $0.17$  perceptible cm, allowing for the observation of stronger water-vapor lines that are typically opaque. Fig. 6 shows that the radiance residuals (AERI-LBLRTM) in this region differ significantly when the calculations use *HITRAN2004* compared with *HITRAN2000* plus updates (*HITRAN2000+*) [44]. These residuals indicate that the air-broadened half-widths in the initial release of *HITRAN2004* were not as accurate as in the previous *HITRAN2000+*. This work was the basis of the re-evaluation and updates made to the initial *HITRAN2004* water-vapor line parameters outlined in this article. Radiance residuals in Fig. 6 show that the updated *HITRAN2004* (*HITRAN2004+*) water lines are an improvement over the initial *HITRAN2004* release and slightly better than the results obtained when using *HITRAN2000+* water-vapor lines in the LBLRTM calculation. In the first five rows of Table 2, a change of air-broadened half-widths in the *HITRAN* database for several transitions in the range of interest is shown. The frequencies are taken from the *HITRAN2004* edition. In the *HITRAN2004* database, some of the  $\gamma_{\text{air}}$  were determined by taking an average of the measurements available for the transition. As can be seen in that region, none of the measurements have cleared the filtering procedure and therefore were replaced with CRB calculations, which proved to be quite accurate.

The new line list was also tested against the ground-based Fourier transform spectrometer observations in the  $1000\text{--}2000\text{ cm}^{-1}$  region acquired during the Italian phase of the EAQUATE campaign [45]. The authors of Ref. [45] have tested the data from *HITRAN2000*, *HITRAN2004*, and the current update and concluded that the latter yields the highest consistency with the observations.

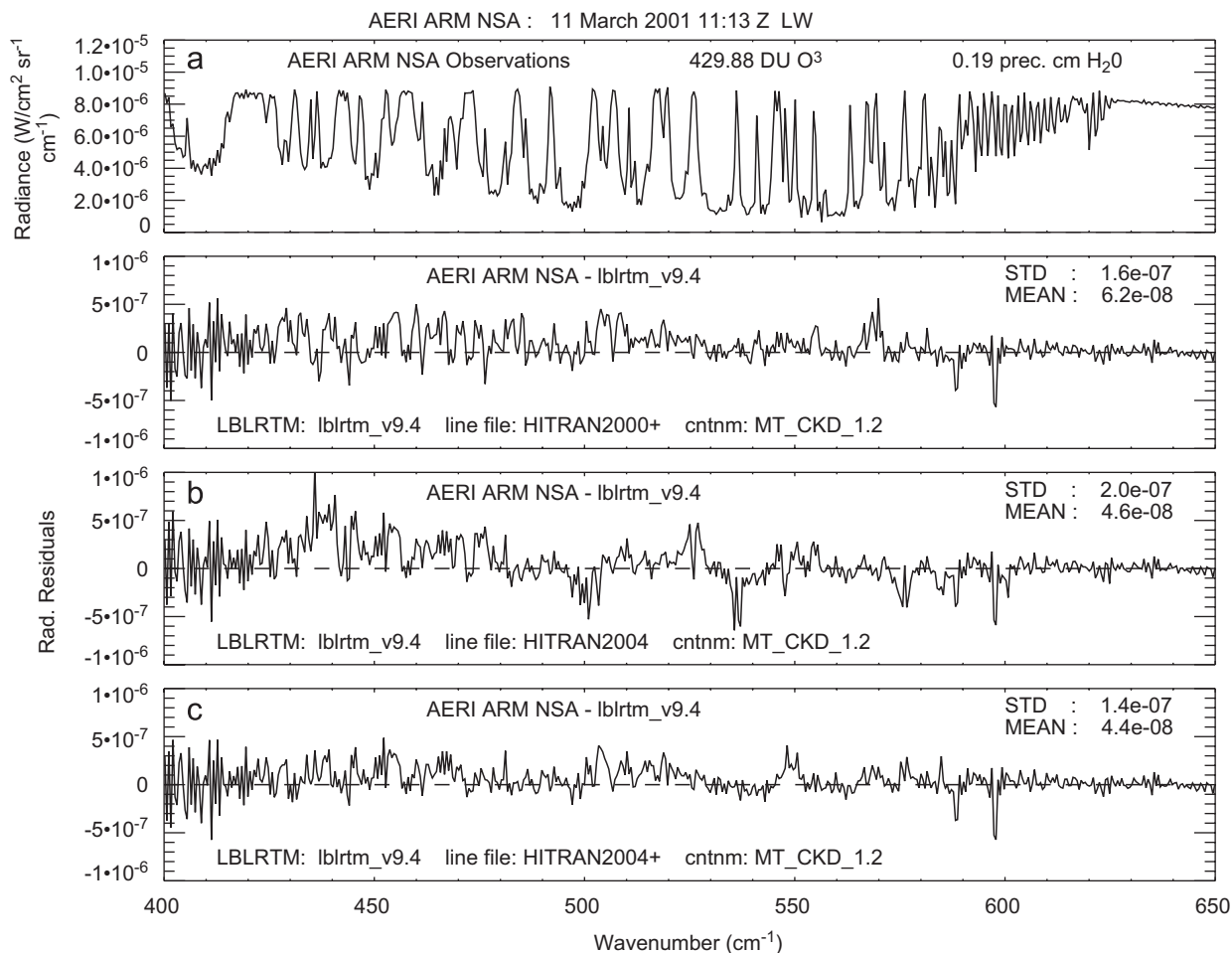


Fig. 6. A comparison example of LBLRTM [42] calculations with AERI-ER observations that demonstrates the impact of recent versions of the water-vapor spectroscopic lines in the *HITRAN* database.

Table 2

Examples of the changes of air-broadened half-widths in the *HITRAN* database

Wavenumber (cm <sup>-1</sup> )	Intensity (cm molecule <sup>-1</sup> )	$\gamma_{\text{air}}^{\text{—2004}}$ (cm <sup>-1</sup> atm <sup>-1</sup> )	Source in <i>HITRAN</i> 2004 <sup>a</sup>	$\gamma_{\text{air}}^{\text{—Update}}$ (cm <sup>-1</sup> atm <sup>-1</sup> )	Source in update
441.7140	3.645E–21	0.0605	[35]	0.0853	[8]
502.2564	1.733E–20	0.0769	[31,35]	0.0569	[8]
536.2459	1.355E–21	0.1051	[31,34]	0.0802	[8]
576.1144	6.235E–21	0.0631	[31,35]	0.0464	[8]
600.1034	2.382E–21	0.0702	[31]	0.0400	[8]
1987.3401	1.364E–23	0.1008	[4,33]	0.0927	[4]

<sup>a</sup>If more than one source is listed, it means that the average of the values reported in these sources is taken.

### 3.2. Limb measurements from satellite and balloons

The new air-broadened half-width values were also validated in different missions performing remote sensing above the lower troposphere.

The Atmospheric Chemistry Experiment (ACE), a satellite mission for remote sensing of the Earth's atmosphere developed by the Canadian Space Agency, features a high-resolution infrared Fourier transform spectrometer measuring in occultation mode [46]. In ACE retrievals, only small portions of the spectrum (typically  $0.3\text{--}1\text{ cm}^{-1}$ ) are being analyzed, rather than the entire spectrum. These microwindows are chosen in a way that most of their spectral features belong to the molecule of interest, assuming that known spectroscopic parameters for these features are accurate [47]. For  $\text{H}_2\text{O}$  retrievals from ACE, the new set of broadening parameters yields improved residuals at low altitudes relative to the residuals using *HITRAN2004* parameters, leading to a decrease in fitting chi-squared of the order of 10%. Considering that the altitudes in these retrievals are above 5 km where the pressure of air is less than that on the ground, and that parameters for the transitions in selected microwindows are considered to have sufficient accuracy already, the 10% improvement is a significant achievement.

In Fig. 7, an example of such improvement is presented, and the decrease in the residual is apparent when the update described in this paper is used. The evolution of the air-broadened half-widths of this line in the *HITRAN* database is presented in the last row of Table 2. In the *HITRAN2004* edition, the average of available experimental values was used; in the current update only the measurements from Ref. [4] have passed the filtering procedure, yielding a more accurate value for air-broadened half-width.

However, there remain significant w-shaped residuals in the fitting of  $\text{H}_2\text{O}$  lines for ACE spectra, features consistent with line shape effects from changes of velocity during collisions and/or the dependence of collisional parameters on absorber velocity [48]. This suggests that the Voigt function is not sufficiently accurate for  $\text{H}_2\text{O}$  in the ACE spectra, and future processing of the ACE measurements will therefore employ a more complex line-shape function for  $\text{H}_2\text{O}$ , such as the speed-dependent Voigt function [48].

The new update was also used in retrievals of temperature and water-vapor profiles from broad-band measurements of the atmospheric emission spectrum in the  $100\text{--}1000\text{ cm}^{-1}$  region acquired by a balloon-borne FT spectrometer (nadir sounding) [49]. The chi-square value of the fit was 0.960, whereas it was 0.988 when using the original *HITRAN2004* [50]. The improvement may seem to be very slight, but again the pressure of air is low and the effect of the accuracy of values of air-broadened half-widths on the retrievals is lower.

#### 4. Conclusions and potential improvements

When one has to pick the most accurate value out of several originating from different imperfect sources, there are numerous ways of accomplishing this feat. However, one should realize that none of these ways will be absolutely flawless. In this work, an algorithm for choosing values of air-broadened half-widths for inclusion into the *HITRAN* database was developed based on the physical principles and statistics. We do not claim that this algorithm is the best possible solution of the problem. Nevertheless, the new algorithm

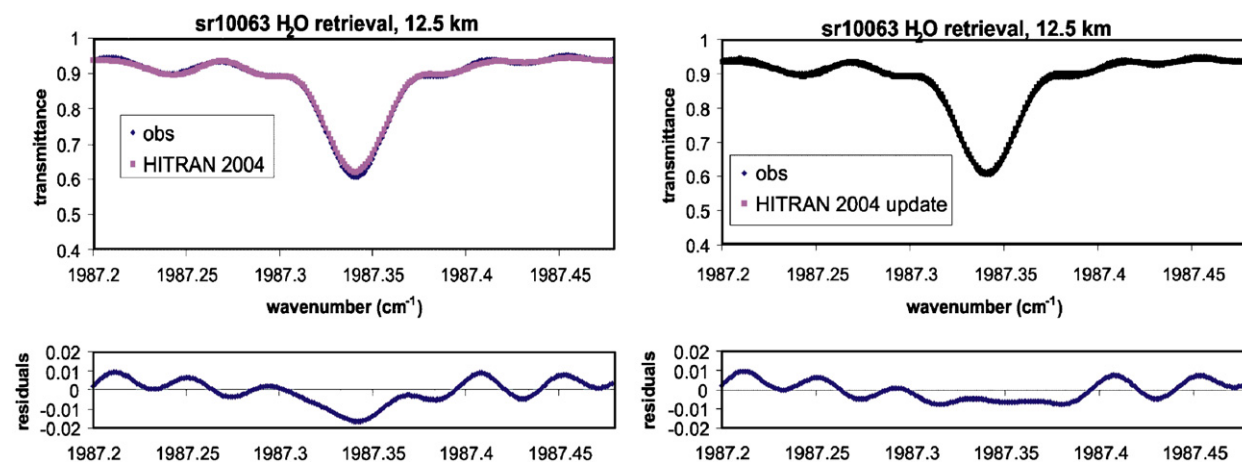


Fig. 7. Comparison of molecular profiles measured with ACE-FTS and the ones calculated using *HITRAN2004* and the update described in this paper.

provided significant improvement to the current *HITRAN2004* edition, as was demonstrated in the validation in application to different remote-sensing missions. This work is an excellent example of fruitful collaboration between the *HITRAN* project and end users, which is extremely important for the continuous improvement of the *HITRAN* database. In total, air-broadened half-width values were updated for 11,787 transitions of water vapor in the *HITRAN2004* database (6789 for H<sub>2</sub><sup>16</sup>O, 2906 for H<sub>2</sub><sup>17</sup>O, and 2092 for H<sub>2</sub><sup>18</sup>O). A supplementary file listing transitions with changed air-broadened half-widths is provided.

In order to improve the current list further, more experimental works are desirable, especially in the near-IR region. There have already been a number of publications extending data existing in the GH database. The CRB calculations have proven to be a good alternative to the experimental values, especially considering that it is a little easier to anticipate where CRB values will be inaccurate. The CRB values for shorter wavelength regions [17] will be included into the algorithm. Nevertheless, further developments of the CRB method are needed, to be more confident in the theoretical values throughout the entire frequency range. In particular, the intermolecular potential constants will need to be better determined and the use of wavefunctions from *ab initio* calculations will replace the Watson Hamiltonian approach. When a significant bulk of new theoretical and experimental values becomes available, it would also be useful to update existing semi-empirical calculations.

Other pressure-induced parameters of water vapor are needed for some improvements in *HITRAN*. For example, the temperature dependences of air-broadened half-widths will be updated based on the work of Toth et al. [51].

## 5. Other updates to the water-vapor parameters in *HITRAN2004*

The H<sub>2</sub><sup>18</sup>O parameters have been updated and H<sub>2</sub><sup>17</sup>O parameters have been added in the  $3\nu + \delta$  and  $4\nu$  polyad region using data from Tanaka et al. [52].

Twenty-five lines in the 14,468–14,558 cm<sup>-1</sup> region were removed from the database according to the recommendation by Tolchenov et al. [53] as these lines were in fact due to the oxygen molecule. In the near future, line assignments, intensities and positions for the principal isotopologue of water will be updated in the 9250–26,000 cm<sup>-1</sup> region using data from Ref. [53].

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## Appendix A. Supplementary Materials

Supplementary data associated with this article can be found in the online version at [doi:10.1016/j.jqsrt.2007.06.009](https://doi.org/10.1016/j.jqsrt.2007.06.009).

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