



DES14X3taz: A TYPE I SUPERLUMINOUS SUPERNOVA SHOWING A LUMINOUS, RAPIDLY COOLING INITIAL PRE-PEAK BUMP

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ABSTRACT

We present DES14X3taz, a new hydrogen-poor superluminous supernova (SLSN-I) discovered by the Dark Energy Survey (DES) supernova program, with additional photometric data provided by the Survey Using DECam for Superluminous Supernovae. Spectra obtained using Optical System for Imaging and low-Intermediate-Resolution Integrated Spectroscopy on the Gran Telescopio CANARIAS show DES14X3taz is an SLSN-I at $z = 0.608$. Multi-color photometry reveals a double-peaked light curve: a blue and relatively bright initial peak that fades rapidly prior to the slower rise of the main light curve. Our multi-color photometry allows us, for the first time, to show that the initial peak cools from 22,000 to 8000 K over 15 rest-frame days, and is faster and brighter than any published core-collapse supernova, reaching 30% of the bolometric luminosity of the main peak. No physical ^{56}Ni -powered model can fit this initial peak. We show that a shock-cooling model followed by a magnetar driving the second phase of the light curve can adequately explain the entire light curve of DES14X3taz. Models involving the shock-cooling of extended circumstellar material at a distance of $\simeq 400 R_{\odot}$ are preferred over the cooling of shock-heated surface layers of a stellar envelope. We compare DES14X3taz to the few double-peaked SLSN-I events in the literature. Although the rise times and characteristics of these initial peaks differ, there exists the tantalizing possibility that they can be explained by one physical interpretation.

Key words: supernovae: general

1. INTRODUCTION

Over the last 10 years, wide-field optical surveys have uncovered a new class of highly luminous transients: “superluminous” supernovae (SLSNe; see review of Gal-Yam 2012). At $M_{\text{peak}} < -21$ mag, SLSNe are 10–100 times brighter than classical core-collapse SN events, but are rarer, with $< 0.1\%$ of the rate (Quimby et al. 2013; McCrum et al. 2015; S. Prajs et al. 2016, in preparation) and only 30 well-studied examples. A physical understanding of these extreme events is still emerging. There are at least two distinct classes (Gal-Yam 2012): SLSNe-II show (narrow) hydrogen emission lines, believed to be generated by interaction with circumstellar material (CSM; Ofek et al. 2007; Smith et al. 2007; Gezari et al. 2009; Benetti et al. 2014), whereas SLSNe-I are hydrogen poor (Quimby et al. 2011), and would traditionally be classified as type Ic SNe. SLSNe-I almost invariably explode in galaxies that are low-mass, compact dwarfs (Neill et al. 2011; Lunnan et al. 2015) and that are metal-poor and strongly star-forming (Chen et al. 2013; Lunnan et al. 2013; Leloudas et al. 2015).

The power source of classical SNe Ic, the radioactive decay of ^{56}Ni , cannot easily explain SLSNe-I: several solar masses of ^{56}Ni are required to reach $M_{\text{peak}} < -21$, and the light curves are difficult to reproduce with a model that has an ejecta mass greater than the ^{56}Ni mass ($M_{\text{ej}} > M_{\text{Ni}}$; e.g., Chatzopoulos et al. 2009; Chomiuk et al. 2011; Inserra et al. 2013; Papadopoulos et al. 2015). Alternative power sources have been proposed, the most popular of which is energy input from a central engine, such as the spin-down of a newly formed and rapidly rotating magnetar (Kasen & Bildsten 2010; Woosley 2010) or accretion of fallback material onto a compact object (Dexter & Kasen 2013). Interactions between SN ejecta and a hydrogen-deficient CSM have also been studied (Chevalier & Irwin 2011; Chatzopoulos et al. 2013).

A new puzzle has now emerged: some SLSNe-I have double-peaked light curves, with an early initial peak a few days after the inferred explosion epoch (Nicholl & Smartt 2016). Two events in the literature have particularly prominent examples: SN2006oz (Leloudas et al. 2012) and LSQ14bdq (Nicholl et al. 2015). In both events, the initial peak lasts a few days in the rest-frame. Nicholl et al. (2015) modeled the initial peak of LSQ14bdq and found it inconsistent with being powered by ^{56}Ni , favoring instead luminosity from the shock-heated surface layers of the exploding star (e.g., Nakar & Sari

2010; Rabinak & Waxman 2011), followed by reheating by a central engine to drive the second peak. In this model, the brightness and duration of the LSQ14bdq initial peak would imply a progenitor radius of a few hundred solar radii. Other explanations that can also produce double-peaked light curves include shock heating of much more extended material (Piro 2015), or a shock breakout driven by a magnetar (Kasen et al. 2015).

Here, we present DES14X3taz, a double-peaked SLSN-I discovered by the Dark Energy Survey (DES; Flaugher et al. 2015) supernova (SN) program (DES-SN). Lying at a redshift $z = 0.608$, the DES-SN *griz* light curves probe rest-frame 3000–6000 Å and provide unique multi-color information of the initial peak. These multi-color data allow us, for the first time, to estimate the temperature and temperature evolution of the initial peak for comparison with various models. We assume $H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1}$ and a flat ΛCDM cosmology with $\Omega_{\text{matter}} = 0.3$ throughout.

2. OBSERVATIONS

DES14X3taz was discovered as a transient event at $r \simeq 23.0$ by DES-SN in DECam (Flaugher et al. 2015) *griz* images taken on 2014 December 21.1 (all dates UTC) at $\alpha = 02^{\text{h}}28^{\text{m}}04^{\text{s}}.46$, $\delta = -04^{\circ}05'12''.7$ (J2000), with previous non-detections on 2014 December 14.1. The field was imaged again, and DES14X3taz was designated as a transient object on 2014 December 27.1. As the light curve was monitored in *griz* by DES over the next 2–3 weeks, the existence of a pronounced initial peak in the light curve became clear, and the object was prioritized for spectroscopic observations. DES monitored DES14X3taz approximately weekly until 2015 February, and further light curve data were then obtained via the Search Using DECam for Superluminous Supernova (SUDSS) project in 2015 March and again in 2015 July. DES re-commenced 2015 August, with DES14X3taz still detected in the data. Further details of the DES-SN difference-imaging search pipeline can be found in Kessler et al. (2015) and Goldstein et al. (2015), with SLSN-specific details in Papadopoulos et al. (2015). Photometric measurements were made on the DES-SN and SUDSS data using the same pipeline as in Papadopoulos et al. (2015), which has also been used extensively in the literature (e.g., Firth et al. 2015 and references therein). The light curve data are presented in Table 1 and Figure 1; they are

Table 1
Light Curve Data for DES14X3taz

Calendar Date	MJD	Phase (days) ^a	f_g (counts) ^b	f_r (counts)	f_i (counts)	f_z (counts)
2014 Dec 10	57001.1	-49.1	-69.0 ± 81.1	-56.9 ± 55.8	0.3 ± 45.2	-7.9 ± 67.6
2014 Dec 14	57005.1	-46.6	-23.3 ± 68.1	4.3 ± 45.8	-5.6 ± 61.7	...
2014 Dec 14	57005.2	-46.5	28.6 ± 90.4
2014 Dec 21	57012.1	-42.2	1914.4 ± 61.4	1531.4 ± 41.8	1203.2 ± 58.5	...
2014 Dec 21	57012.2	-42.2	988.7 ± 120.9
2014 Dec 27	57018.1	-38.5	1973.4 ± 70.1	1672.0 ± 58.5
2014 Dec 27	57018.2	-38.5	1792.7 ± 87.0	...
2014 Dec 28	57019.1	-37.9	1391.8 ± 82.6
2015 Jan 03	57025.1	-34.1	...	1444.2 ± 118.3	1393.4 ± 104.0	...
2015 Jan 04	57026.1	-33.5	1152.0 ± 88.3
2015 Jan 07	57029.1	-31.7	438.4 ± 120.5
2015 Jan 10	57032.1	-29.8	...	1779.6 ± 83.4	1958.4 ± 69.7	...
2015 Jan 11	57033.1	-29.2	2084.6 ± 77.9
2015 Jan 14	57036.0	-27.3	1591.4 ± 59.6
2015 Jan 16	57038.1	-26.0	...	3308.9 ± 44.5
2015 Jan 17	57039.1	-25.5	3536.7 ± 63.8	...
2015 Jan 18	57040.1	-24.8	3297.5 ± 70.5
2015 Jan 20	57042.1	-23.6	3441.5 ± 71.8
2015 Jan 22	57044.1	-22.3	...	5290.8 ± 113.9	5195.1 ± 91.5	...
2015 Jan 23	57045.1	-21.7	4753.5 ± 110.3
2015 Jan 27	57049.1	-19.2	5757.2 ± 195.3
2015 Jan 29	57051.1	-18.0	...	7849.3 ± 85.4	7481.1 ± 85.2	...
2015 Jan 30	57052.1	-17.4	6690.9 ± 72.0
2015 Feb 07	57060.0	-12.4	9105.5 ± 100.7	11139.1 ± 97.2	10361.2 ± 89.6	...
2015 Feb 07	57060.1	-12.4	9206.2 ± 134.1
2015 Feb 25	57078.0	-1.2	11068.6 ± 342.2	14419.4 ± 376.7	14160.5 ± 263.2	12316.9 ± 321.3
2015 Mar 11	57092.0	7.5	10294.3 ± 574.2	15273.9 ± 389.2
2015 Jul 08	57211.4	81.7	775.2 ± 394.8	552.2 ± 318.1	1085.7 ± 247.7	2179.9 ± 201.8
2015 Jul 25	57228.3	92.3	1463.7 ± 133.5	1762.1 ± 201.5
2015 Jul 25	57228.4	92.3	165.7 ± 76.4	438.1 ± 77.0
2015 Aug 15	57249.3	105.3	224.0 ± 48.5	377.7 ± 39.1	714.3 ± 45.7	...
2015 Aug 15	57249.4	105.3	940.0 ± 52.7
2015 Aug 19	57253.4	107.8	66.8 ± 39.5	284.5 ± 33.6	633.2 ± 43.3	...
2015 Aug 21	57255.3	109.0	884.0 ± 75.8
2015 Aug 23	57257.3	110.3	96.5 ± 39.4	282.6 ± 32.2
2015 Aug 23	57257.4	110.3	646.9 ± 34.3	...
2015 Aug 31	57265.3	115.2	738.9 ± 102.5
2015 Aug 31	57265.4	115.3	702.4 ± 96.1	...
2015 Sep 03	57268.2	117.1	...	160.4 ± 186.7
2015 Sep 04	57269.2	117.7	245.0 ± 188.4
2015 Sep 08	57273.4	120.2	525.8 ± 58.8	879.2 ± 69.8
2015 Sep 13	57278.2	123.3	...	328.2 ± 47.3
2015 Sep 13	57278.3	123.3	89.6 ± 50.3
2015 Sep 18	57283.2	126.4	363.2 ± 58.9	782.6 ± 54.2
2015 Sep 22	57287.2	128.8	452.9 ± 210.8	178.4 ± 63.9
2015 Sep 22	57287.3	128.9	470.6 ± 76.0	678.8 ± 87.6

Notes.^a Relative to maximum light in the rest-frame.^b Fluxes f in each filter are given in counts. A zeropoint of 31.0 converts counts into AB magnitudes.

also available from the WISEREP archive⁴⁹ (Yaron & Gal-Yam 2012).

DES14X3taz was spectroscopically observed on 2015 January 26 (when the target was $r \simeq 21.5$) at the Gran Telescopio CANARIAS (GTC) using the Optical System for Imaging and low-Intermediate-Resolution Integrated Spectroscopy and the R500R grism. The spectrum was taken in relatively poor conditions (bright sky, 1'' seeing) and is of low signal-to-noise ratio, so it was re-observed on 2015 February 6

in dark conditions and 0'' seeing. The spectra were reduced using standard Image Reduction and Analysis Facility (IRAF) routines and have an effective wavelength coverage of 5200–9000Å (observer frame). The spectra are spectroscopically similar and available from WISEREP.

Figure 2 shows the spectrum obtained on 2015 February 6. Weak H β , [O II], and [O III] host-galaxy emission lines give a redshift of $z = 0.608$. A comparison with literature spectra using SUPERFIT (Howell et al. 2005) gives an excellent match to a hydrogen-poor SLSN. The broad absorption features

⁴⁹ <http://wiserep.weizmann.ac.il/>

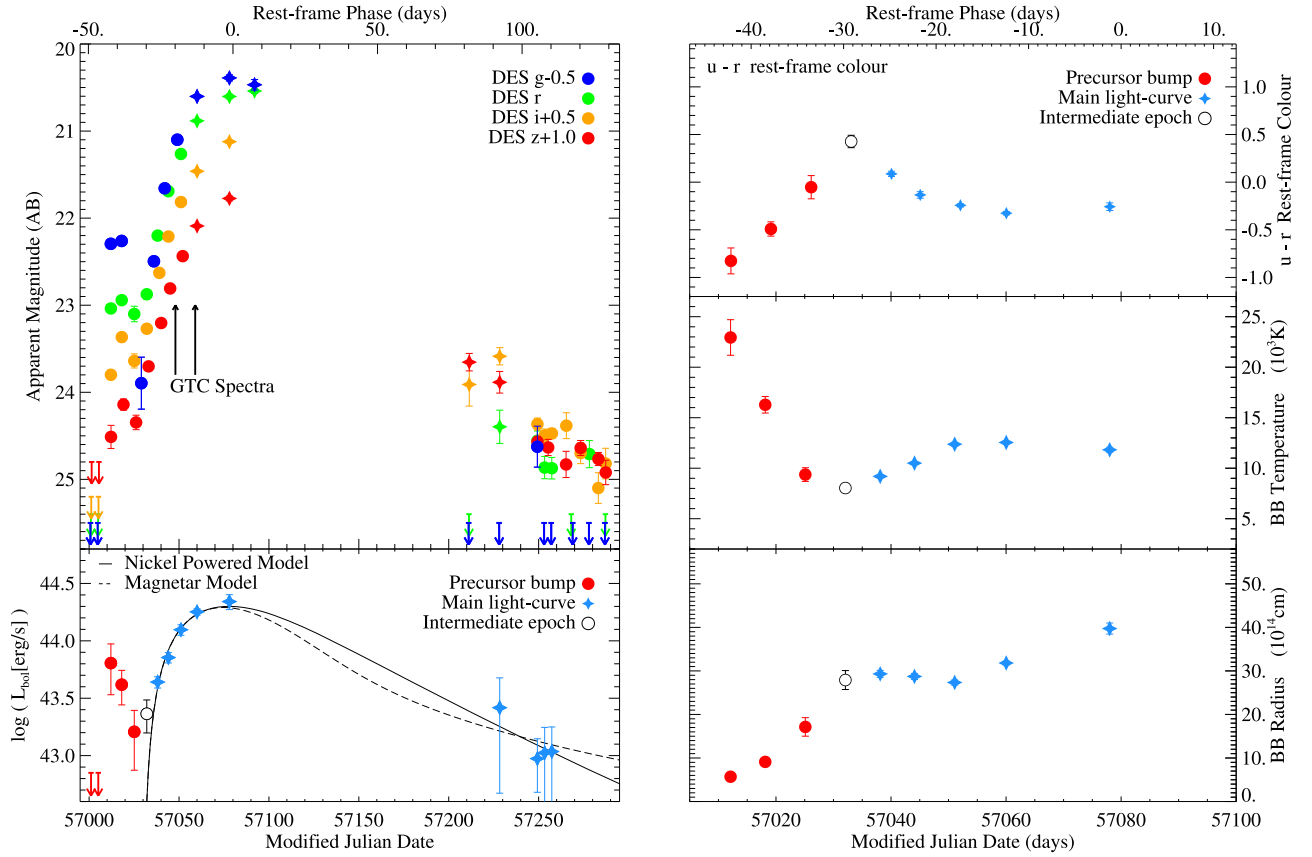


Figure 1. Photometric properties of DES14X3taz. Upper left: the multi-color observer-frame *griz* light curve from DES (circles) and SUDSS (crosses). The filters are offset for clarity. Upper limits are denoted with arrows, and the epochs of GTC spectroscopy are highlighted. Lower left: the bolometric light curve on epochs with three or more filter detections. Epochs associated with the initial peak are shown as red circles, with the main light curves plotted as blue crosses. Intermediate epochs are plotted as open circles. The main light curve has fits for both a ^{56}Ni -powered model (solid; $M_{\text{Ni}} = 26 M_{\odot}$) and a magnetar spin-down-powered model (dashed; $P_{\text{ms}} = 2.2 \text{ ms}$, $B_{14} = 1.3 \times 10^{14} \text{ G}$). Upper right: the rest-frame $u - r$ color evolution. Middle and lower right: the evolution of the temperature and radius inferred from the blackbody fits to individual epochs of photometry. Details of k -corrections and blackbody fits are in Section 3.

at rest-frame 4200 and 4500 Å are O II, seen in all SLSNe-I at this phase (e.g., Quimby et al. 2011; Inserra et al. 2013).

The host galaxy of DES14X3taz is detected in stacked images from DES that contain no SN light. Using SEXTRACTOR (Bertin & Arnouts 1996), we measure host-galaxy AB magnitudes (MAG_AUTO) of $(g, r, i, z) = (26.16 \pm 0.39, 25.07 \pm 0.13, 24.95 \pm 0.13, 25.00 \pm 0.18)$ after correcting for MW extinction. Using the Z-PEG photometric-redshift code (Le Borgne & Rocca-Volmerange 2002), a Kroupa (2001) initial mass function, and the redshift fixed at $z = 0.608$, we estimate a host-galaxy stellar mass of $\log(M/M_{\odot}) = 8.0^{+0.4}_{-0.2}$ and a star formation rate of $\log(\text{SFR}) = -0.8^{+0.6}_{-0.2} M_{\odot} \text{ yr}^{-1}$. These properties are consistent with the hosts of other SLSNe (e.g., Neill et al. 2011; Leloudas et al. 2015; Papadopoulos et al. 2015).

3. ANALYSIS

We next study the unusual light curve of DES14X3taz. The most striking feature (Figure 1, upper left panel) is the presence of a first peak in the light curve around MJD 57015 (2014 December 24), prior to the main peak at MJD 57080 (2015 February 27). The initial peak is most pronounced in the bluer filters, but is still distinct in the z -band.

We estimate the rest-frame *ugr* light curve from the observer-frame *riz* bands, which correspond to a similar

wavelength range at $z = 0.608$, using a standard k -correction procedure. With the limited wavelength coverage of our GTC spectra, pre- and at-max k -corrections are determined using the spectrally similar event LSQ14bdq, which has coverage down to 3600 Å (Nicholl et al. 2015) and late-time k -corrections using data from PTF12dam at +171 d (Nicholl et al. 2013). A polynomial fit to the rest-frame g -band light curve indicates DES14X3taz reached a peak brightness of $M_g = -21.39$ on MJD 57080 after correcting for foreground extinction and assuming zero internal reddening. We show the $u - r$ rest-frame color evolution in Figure 1 (upper right panel)—the initial peak is blue ($u - r \simeq -0.9$) at first, but quickly reddens by nearly 1.5 mag to $u - r \simeq 0.5$ (before becoming bluer again during the main rise to peak brightness). Such reddening behavior is typical of a rapid decrease in temperature due to cooling.

We analyze this rapid evolution in more detail by fitting a blackbody to each epoch of photometry with observations of at least 3 filters within 4 days, using Planck’s law and a luminosity distance of 3587 Mpc. (The epoch at MJD 57032 is excluded from these fits, as at this epoch there is potentially flux from both components of the light curve.) The evolution of the temperature and radius inferred from these fits is shown in Figure 1 (right panels) and demonstrates a rapid cooling from 22,000 to 8000 K over 15 rest-frame days. This is followed by a reheating phase during the main light curve, which then itself

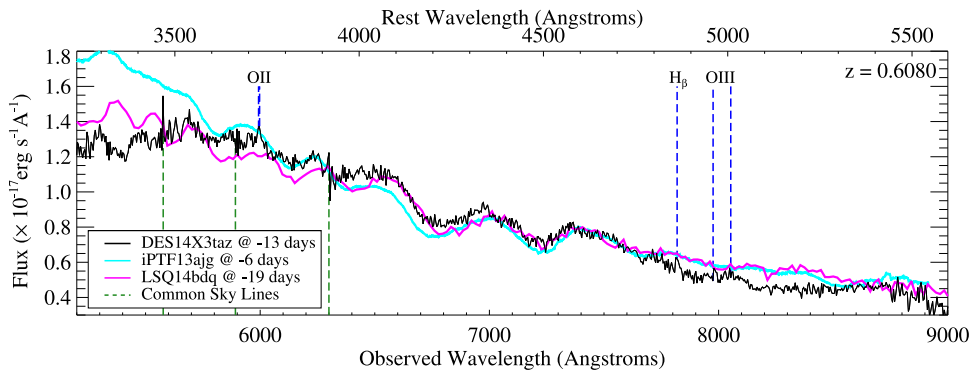


Figure 2. GTC optical spectrum of DES14X3taz from 2015 February 6 (black line). Galaxy emission features are highlighted in blue and prominent sky lines in green. Also shown are the spectra of the SLSNe-I iPTF13ajg at -6 d (cyan; Vreeswijk et al. 2014) and LSQ14bdq at -19 d (magenta; Nicholl et al. 2015).

cools as it approaches the main peak, consistent with other literature on SLSNe-I.

The bolometric light curve (Figure 1, lower left panel) is also constructed by integrating the best-fit blackbody function on each epoch. The two peaks are particularly clear in bolometric luminosity: one, the initial peak, prior to MJD 57030; and the other, the main peak.

3.1. The Main Peak

We fit the bolometric light curve of the main peak with two models: energy deposition from the radioactive decay of ^{56}Ni , and energy deposition from the spin-down of a rapidly rotating magnetar. In both cases, assuming a spherically symmetric and homologously expanding ejecta, the bolometric luminosity L as a function of time t since explosion is (Arnett 1982)

$$L(t) = e^{-(t/\tau_m)^2} \int_0^t \frac{2t'}{\tau_m^2} P(t') e^{(t'/\tau_m)^2} dt', \quad (1)$$

where $P(t)$ is the power function, the total absorbed power due to either ^{56}Ni or a central engine, and τ_m is the diffusion timescale parameter, a function of the ejecta mass (M_{ej}), opacity (κ), and kinetic energy (E_k). We assume a constant $\kappa = 0.1 \text{ cm}^2 \text{ g}^{-1}$ (see the discussion in Inserra et al. 2013).

The power function appropriate for ^{56}Ni depends on the ^{56}Ni mass (M_{Ni}) synthesized in the explosion, τ_m , and the time of explosion t_0 (e.g., Inserra et al. 2013; Papadopoulos et al. 2015). If we enforce $M_{\text{ej}} \geq M_{\text{Ni}}$, we obtain a ^{56}Ni mass of $M_{\text{Ni}} = 26 \pm 2.6 M_{\odot}$ with $M_{\text{ej}} = M_{\text{Ni}}$ (shown in Figure 1). Although the fit is statistically acceptable, with a χ^2 of 4.9 for 6 degrees of freedom (dof), such a model with $M_{\text{ej}} = M_{\text{Ni}}$ is difficult to reconcile with the observed spectra (see, e.g., the discussion in Inserra et al. 2013). Enforcing more physical constraints on the ratio of M_{Ni} to M_{ej} (e.g., Papadopoulos et al. 2015) then results in implied ejecta velocities significantly greater than observed in the spectra of SLSNe-I: for example, a limit of $M_{\text{Ni}} < 0.7 M_{\text{ej}}$ requires velocities of $> 22,000 \text{ km s}^{-1}$, compared to $10,000 \text{ km s}^{-1}$ measured from the DES14X3taz spectrum. Thus, like other SLSNe-I, the light curve of DES14X3taz cannot be fit with a physically plausible ^{56}Ni diffusion model.

The magnetar power function depends on two parameters. The initial spin period (P_{ms} , in milliseconds) and the initial magnetic field strength (B_{14} , in 10^{14} G). Assuming full-trapping of the magnetar radiation, we find $P_{\text{ms}} = 2.2 \pm 1.4 \text{ ms}$, $B_{14} = 1.25 \pm 0.30 \times 10^{14} \text{ G}$, and $\tau_m = 54.5 \text{ d}$. From this, we infer an ejecta mass of $M_{\text{ej}} = 9.5 M_{\odot}$. The model fit has a χ^2 of

5.4 with 5 dof. These values are similar to other SLSNe-I with initial peaks (e.g., LSQ14bdq: $P_{\text{ms}} = 1.7 \text{ ms}$ and $B_{14} = 0.6 \times 10^{14} \text{ G}$; Nicholl et al. 2015) and SLSNe-I in general (e.g., DES13S2cmm: $P_{\text{ms}} = 5.3 \text{ ms}$ and $B_{14} = 1.4 \times 10^{14} \text{ G}$; Papadopoulos et al. 2015).

3.2. Modeling the Initial Peak

We next examine the initial peak in the light curve. We consider both a ^{56}Ni power source (for example, from an underlying SN explosion) and various shock-cooling models.

3.2.1. ^{56}Ni -powered Models

As the main light curve is well modeled by a magnetar-driven power source, it is appropriate to consider whether the initial peak in the light curve is consistent with a normal ^{56}Ni -powered SN explosion, i.e., the SN that gave birth to the magnetar. We compare the observed multi-color light curve of the initial peak in DES14X3taz to various stripped envelope SNe (SESNe) from Bianco et al. (2014) and Modjaz et al. (2014). We use the templates of N. V. Karpenka et al. (2016, in preparation), where 28 events have been modeled using a simple parametric form allowing them to be placed at arbitrary redshifts (including k -corrections). As SLSNe-I (and DES14X3taz) are hydrogen-poor, we do not consider type II SNe.

None of the SESN templates are consistent with the initial peak in the DES14X3taz light curve: the initial peak rises faster, is significantly bluer at peak, and is 1.5 mag brighter than any of the template events. Similarly, fitting the ^{56}Ni -powered model from Section 3.1 to the initial peak requires $2.9 M_{\odot}$ of pure ^{56}Ni (i.e., $M_{\text{Ni}} = M_{\text{ej}}$) and a large inferred energy of $E_k = 4.1 \times 10^{52} \text{ erg}$. Typical core-collapse events produce $\sim 0.1 M_{\odot}$ of ^{56}Ni , and even a complete burning of $1 M_{\odot}$ of CO to ^{56}Ni produces only $1 \times 10^{51} \text{ erg}$ (Nicholl et al. 2015). Other power sources are clearly required.

3.2.2. Shock Cooling of a Stellar Envelope

We next consider the possibility that the initial peak of DES14X3taz is driven by shock cooling where, following an SN shock wave that heats the stellar envelope, the envelope expands and cools adiabatically, releasing energy. We use the analytical models of Rabinak & Waxman (2011) and consider various progenitor envelopes: a radiative H envelope (“blue supergiant,” BSG), a convective H envelope (“red supergiant,” RSG), He-dominated envelopes (“He envelopes”) and C/O-

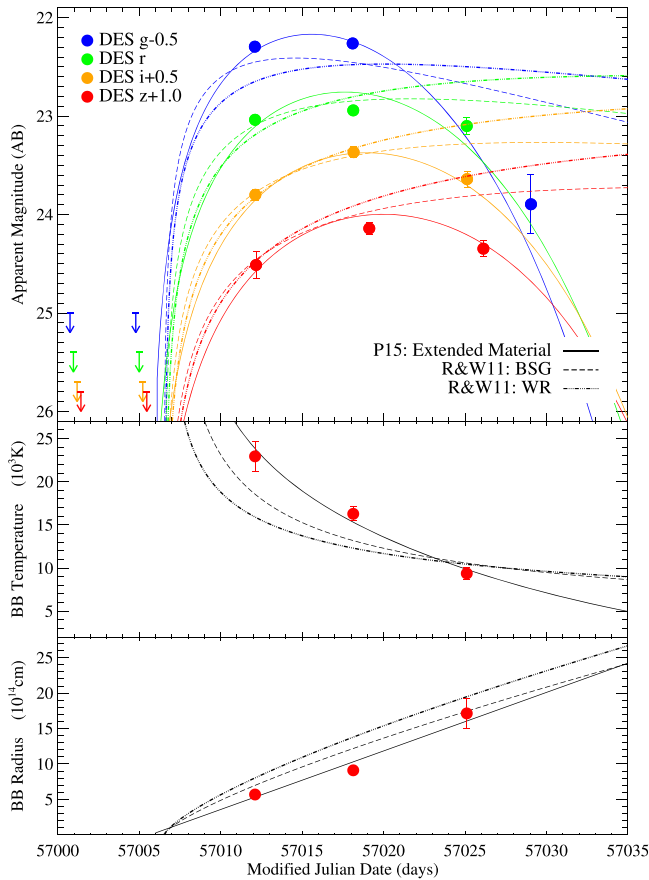


Figure 3. Top: various analytic models fit to the multi-color data of the initial peak of DES14X3taz. The best-fitting shock-cooling models from Rabinak & Waxman (2011) of an extended envelope (BSG; dashed; $M_{ej} = 39 M_{\odot}$, $R_* = 530 R_{\odot}$, $E_k = 1.5 \times 10^{52}$ erg) and a compact model (WR; dotted-dashed; $M_{ej} = 500 M_{\odot}$, $R_* = 50 R_{\odot}$, $E_k = 2.5 \times 10^{53}$ erg), together with an extended material model from Piro (2015; solid; $M_{core} = 9.5 M_{\odot}$, $M_{ext} = 2.7 M_{\odot}$, $R_{ext} = 400 R_{\odot}$, $E_{sn} = 6.0 \times 10^{51}$ erg) are shown. Middle and lower panels: the DES14X3taz temperature and radius evolution from blackbody fitting (red circles; from Figure 1) compared to the shock-cooling models discussed above.

dominated envelopes (“Wolf–Rayet”). The Rabinak & Waxman (2011) models are only valid at early times, so we cut off the models according to their Equation (17). We determine $L(t)$ from their relations and, assuming a blackbody, fit the models to the initial peak of DES14X3taz by determining fluxes in the DES bandpasses at $z = 0.608$.

We show the results in Figure 3. The RSG, BSG, and He envelopes are very similar, so we only show the BSG example for clarity. Although the rise and color of the initial peak can be reasonably well matched by these models, in detail the models do not drop fast enough to match the DES14X3taz light curve and do not well reproduce the inferred temperature evolution, with a χ^2 of 148 and 373 for the BSG and Wolf–Rayet models, respectively (for 8 dof). The size of the envelopes implied are also very large (e.g., $\simeq 520 R_{\odot}$ for the BSG envelope and $\simeq 30 R_{\odot}$ for the C/O-dominated envelope). These are a factor of ~ 10 larger than typical BSGs or Wolf–Rayet stars.

3.2.3. Shock Cooling of Extended Material

A similar model involves extended material around the progenitor star, but at much larger radii than the stellar

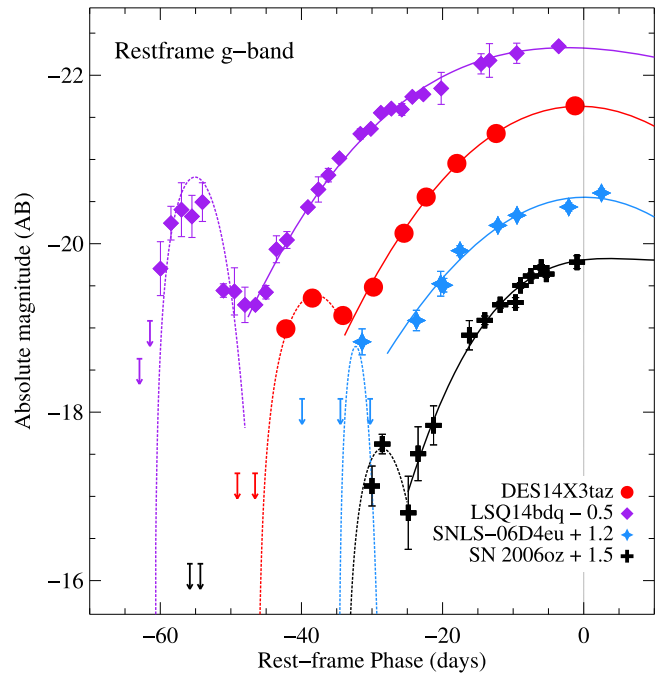


Figure 4. Rest-frame g -band light curve for DES14X3taz, compared to literature SLSNe-I with initial peaks. Offsets in absolute magnitude have been applied for clarity. Piro (2015) model fits are plotted over the initial peaks, and polynomial fits are overplotted on the main peak of each event. 3σ non-detections are shown.

envelope. If this material is sufficiently massive and extended—perhaps from stripping or earlier pre-explosion eruptions—the SN shock can propagate into it and produce a bright initial peak in the light curve (Ofek et al. 2010; Piro 2015) or, in the absence of further late-time energy input into the ejecta, an isolated bright, luminous transient (Drout et al. 2014).

We use the analytical relations of Piro (2015) to model this, showing the results in Figure 3. Due to degeneracies in the model between the core mass (M_{core}) and energy, we fix $M_{core} = 9.5 M_{\odot}$, the ejecta mass estimated from the magnetar fitting to the main peak (Section 3.1). The resultant fit has a χ^2 of 45 for 9 dof, with $\simeq 2.7 M_{\odot}$ of material at a radius of $\simeq 400 R_{\odot}$ and an energy $E_{sn} = 6.0 \times 10^{51}$ erg. (Relaxing this M_{core} constraint results in values of M_{core} between 5 and $100 M_{\odot}$, but does not significantly alter the mass of extended material or the χ^2 .) This model reproduces the color, photometric evolution, and temperature evolution of DES14X3taz.

4. DISCUSSION AND FUTURE WORK

The discovery of double-peaked SLSNe-I provides new insights into the progenitors of SLSNe. The unique multi-color information of the initial peak of DES14X3taz reveals a very hot event that cools rapidly. We propose that this is driven by a period of shock cooling, with the SLSN phase of the light curve driven by reheating from a central engine, consistent with a magnetar. Shock cooling from an extended stellar envelope can match the color of the early part of the initial peak, but does not reproduce the fast dropoff, providing a poor overall match. By contrast, models with extended material at $\simeq 500 R_{\odot}$ can replicate the entire initial peak.

The quantity of data on double-peaked light curves of SLSNe-I is increasing quickly. In Figure 4, we compare the

DES14X3taz rest-frame g -band light curve to other SLSNe-I from the literature that also exhibit initial peaks: SN2006oz ($z = 0.376$; Leloudas et al. 2012), SNLS-06D4eu ($z = 1.588$; Howell et al. 2013), and LSQ14bdq ($z = 0.345$; Nicholl et al. 2015). $griz$ data of the initial peaks are available for SN2006oz (covering $\approx 3400\text{--}6500\text{\AA}$ rest-frame) and SNLS-06D4eu ($\approx 1900\text{--}3400\text{\AA}$), but only data from a single broad “ $g + r$ ” filter are available for LSQ14bdq ($\approx 4200\text{\AA}$).

Figure 4 shows diversity in the luminosity and duration of the initial peaks, especially when compared to the rise-time and peak magnitude of the resulting SLSN. However, only ≈ 15 literature SLSNe have well-measured pre-explosion photometry, at least $\approx 30\%$ of which have double peaks. Other initial peaks may be present in the remainder, but lie below the photometric detection thresholds of the discovery surveys. The initial peaks in Figure 4 are all consistent with the Piro (2015) model, perhaps suggesting that one general physical interpretation can explain the evolution of all SLSN-I events.

Understanding the initial peaks of SLSNe-I in more detail will require spectral (and therefore velocity and composition) measurements. Such measurements can be used to search for any hydrogen or helium signatures that may be present, if the model of extended material is indeed correct. These observations are technically challenging, as SLSNe are rare and would need to be identified rapidly from other SN types based on only 1–2 epochs of photometry. However, DES is expected to find $\approx 5\text{--}6$ well-observed SLSNe-I per year (Scovaccicchi et al. 2016), and the distinct properties of these SLSNe at early times (very blue colors, very faint host galaxies) allow them to be cleanly selected, making spectral coverage a realistic possibility in the near future.

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