# Description of limits of ranges of iterations of stochastic integral mappings of infinitely divisible distributions 

Ken-iti Sato

Hachiman-yama 1101-5-103, Tenpaku-ku, Nagoya, 468-0074 Japan
E-mail address: ken-iti.sato@nifty.ne.jp

Abstract. For infinitely divisible distributions $\rho$ on $\mathbb{R}^{d}$ the stochastic integral mapping $\Phi_{f} \rho$ is defined as the distribution of improper stochastic integral $\int_{0}^{\infty-} f(s) d X_{s}^{(\rho)}$, where $f(s)$ is a non-random function and $\left\{X_{s}^{(\rho)}\right\}$ is a Lévy process on $\mathbb{R}^{d}$ with distribution $\rho$ at time 1 . For three families of functions $f$ with parameters, the limits of the nested sequences of the ranges of the iterations $\Phi_{f}^{n}$ are shown to be some subclasses, with explicit description, of the class $L_{\infty}$ of completely selfdecomposable distributions. In the critical case of parameter 1 , the notion of weak mean 0 plays an important role. Examples of $f$ with different limits of the ranges of $\Phi_{f}^{n}$ are also given.

## 1. Introduction

Let $I D=I D\left(\mathbb{R}^{d}\right)$ be the class of infinitely divisible distributions on $\mathbb{R}^{d}$, where $d$ is a fixed finite dimension. For a real-valued locally square-integrable function $f(s)$ on $\mathbb{R}_{+}=[0, \infty)$, let

$$
\Phi_{f} \rho=\mathcal{L}\left(\int_{0}^{\infty-} f(s) d X_{s}^{(\rho)}\right),
$$

the law of the improper stochastic integral $\int_{0}^{\infty-} f(s) d X_{s}^{(\rho)}$ with respect to the Lévy process $\left\{X_{s}^{(\rho)}: s \geqslant 0\right\}$ on $\mathbb{R}^{d}$ with $\mathcal{L}\left(X_{1}^{(\rho)}\right)=\rho$. This integral is the limit in probability of $\int_{0}^{t} f(s) d X_{s}^{(\rho)}$ as $t \rightarrow \infty$. The domain of $\Phi_{f}$, denoted by $\mathfrak{D}\left(\Phi_{f}\right)$, is the class of $\rho \in I D$ such that this limit exists. The range of $\Phi_{f}$ is denoted by $\mathfrak{R}\left(\Phi_{f}\right)$. If $f(s)=0$ for $s \in\left(s_{0}, \infty\right)$, then $\Phi_{f} \rho=\mathcal{L}\left(\int_{0}^{s_{0}} f(s) d X_{s}^{(\rho)}\right)$ and $\mathfrak{D}\left(\Phi_{f}\right)=I D$. For many choices of $f$, the description of $\mathfrak{R}\left(\Phi_{f}\right)$ is known; they are quite diverse. A seminal example is $\mathfrak{R}\left(\Phi_{f}\right)=L=L\left(\mathbb{R}^{d}\right)$, the class of selfdecomposable distributions on $\mathbb{R}^{d}$, for $f(s)=e^{-s}$ (Wolfe (1982), Sato (1999), Rocha-Arteaga and Sato (2003)). The iteration $\Phi_{f}^{n}$ is defined by $\Phi_{f}^{1}=\Phi_{f}$ and $\Phi_{f}^{n+1} \rho=\Phi_{f}\left(\Phi_{f}^{n} \rho\right)$ with $\mathfrak{D}\left(\Phi_{f}^{n+1}\right)=\left\{\rho \in \mathfrak{D}\left(\Phi_{f}^{n}\right): \Phi_{f}^{n} \rho \in\right.$ $\left.\mathfrak{D}\left(\Phi_{f}\right)\right\}$. Then

$$
I D \supset \mathfrak{R}\left(\Phi_{f}\right) \supset \mathfrak{R}\left(\Phi_{f}^{2}\right) \supset \cdots .
$$

We define the limit class

$$
\Re_{\infty}\left(\Phi_{f}\right)=\bigcap_{n=1}^{\infty} \mathfrak{R}\left(\Phi_{f}^{n}\right) .
$$

If $f(s)=e^{-s}$, then $\mathfrak{R}\left(\Phi_{f}^{n}\right)$ is the class of $n$ times selfdecomposable distributions and $\Re_{\infty}\left(\Phi_{f}\right)$ is the class $L_{\infty}$ of completely selfdecomposable distributions, which is the smallest class that is closed under convolution and weak convergence and contains all stable distributions on $\mathbb{R}^{d}$. This sequence and the class $L_{\infty}$ were introduced by Urbanik (1973) and studied by Sato (1980) and others. If $f(s)=(1-s) 1_{[0,1]}(s)$, then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty}$, which was established by Jurek (2004) and Maejima and Sato (2009); in this case $\mathfrak{R}\left(\Phi_{f}\right)$ is the class of $s$-selfdecomposable distributions in the terminology of Jurek (1985). The paper of Maejima and Sato (2009) showed $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty}$ in many cases including (1) $f(s)=(-\log s) 1_{[0,1]}(s),(2) s=\int_{f(s)}^{\infty} u^{-1} e^{-u} d u(0<s<\infty)$, (3) $s=\int_{f(s)}^{\infty} e^{-u^{2}} d u\left(0<s<s_{0}=\sqrt{\pi} / 2\right)$. The classes $\mathfrak{R}\left(\Phi_{f}\right)$ corresponding to (1)(3) are Goldie-Steutel-Bondesson class $B$, Thorin class $T$ (see Barndorff-Nielsen et al. (2006)), and the class $G$ of generalized type $G$ distributions, respectively. These results pose the problem what classes other than $L_{\infty}$ can appear as $\Re_{\infty}\left(\Phi_{f}\right)$ in general. For $-\infty<\alpha<2, p>0$, and $q>0$, we consider the three families of functions $\bar{f}_{p, \alpha}(s), l_{q, \alpha}(s)$, and $f_{\alpha}(s)$ as in $[\mathrm{S}]$ (we refer to Sato (2010) as $[\mathrm{S}]$ ). We define $\bar{\Phi}_{p, \alpha}$, $\Lambda_{q, \alpha}$, and $\Psi_{\alpha}$ to be the mappings $\Phi_{f}$ with $f(s)$ equal to these functions, respectively. In this paper we will prove the following theorem on the classes $\Re_{\infty}\left(\Phi_{f}\right)$ of those mappings. The case $\alpha=1$ is delicate. There the notion of weak mean 0 plays an important role.

Theorem 1.1. (i) If $\alpha \leqslant 0, p \geqslant 1$, and $q>0$, then

$$
\mathfrak{R}_{\infty}\left(\bar{\Phi}_{p, \alpha}\right)=\Re_{\infty}\left(\Lambda_{q, \alpha}\right)=\Re_{\infty}\left(\Psi_{\alpha}\right)=L_{\infty} .
$$

(ii) If $0<\alpha<1, p \geqslant 1$, and $q>0$, then

$$
\mathfrak{R}_{\infty}\left(\bar{\Phi}_{p, \alpha}\right)=\mathfrak{R}_{\infty}\left(\Lambda_{q, \alpha}\right)=\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)=L_{\infty}^{(\alpha, 2)} .
$$

(iii) If $\alpha=1, p \geqslant 1$, and $q=1$, then

$$
\Re_{\infty}\left(\bar{\Phi}_{p, 1}\right)=\Re_{\infty}\left(\Lambda_{1,1}\right)=\Re_{\infty}\left(\Psi_{1}\right)=L_{\infty}^{(1,2)} \cap\{\mu \in I D: \mu \text { has weak mean } 0\} .
$$

(iv) If $1<\alpha<2, p \geqslant 1$, and $q>0$, then

$$
\mathfrak{R}_{\infty}\left(\bar{\Phi}_{p, \alpha}\right)=\mathfrak{R}_{\infty}\left(\Lambda_{q, \alpha}\right)=\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)=L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D: \mu \text { has mean } 0\}
$$

Let us explain the concepts used in the statement of Theorem 1.1. A distribution $\mu \in I D$ belongs to $L_{\infty}$ if and only if its Lévy measure $\nu_{\mu}$ is represented as

$$
\nu_{\mu}(B)=\int_{(0,2)} \Gamma_{\mu}(d \beta) \int_{S} \lambda_{\beta}^{\mu}(d \xi) \int_{0}^{\infty} 1_{B}(r \xi) r^{-\beta-1} d r
$$

for Borel sets $B$ in $\mathbb{R}^{d}$, where $\Gamma_{\mu}$ is a measure on the open interval $(0,2)$ satisfying $\int_{(0,2)}\left(\beta^{-1}+(2-\beta)^{-1}\right) \Gamma_{\mu}(d \beta)<\infty$ and $\left\{\lambda_{\beta}^{\mu}: \beta \in(0,2)\right\}$ is a measurable family of probability measures on $S=\left\{\xi \in \mathbb{R}^{d}:|\xi|=1\right\}$. This $\Gamma_{\mu}$ is uniquely determined by $\nu_{\mu}$ and $\left\{\lambda_{\beta}^{\mu}\right\}$ is determined by $\nu_{\mu}$ up to $\beta$ of $\Gamma_{\mu}$-measure 0 (see $[\mathrm{S}]$ and Sato (1980)). For a Borel subset $E$ of the interval ( 0,2 ), the class $L_{\infty}^{E}$ denotes, as in $[\mathrm{S}]$, the totality of $\mu \in L_{\infty}$ such that $\Gamma_{\mu}$ is concentrated on $E$. The classes $L_{\infty}^{(\alpha, 2)}$ and $L_{\infty}^{(1,2)}$ appearing in Theorem 1.1 are for $E=(\alpha, 2)$ and (1,2), respectively. Let $C_{\mu}(z)$ $\left(z \in \mathbb{R}^{d}\right), A_{\mu}$, and $\nu_{\mu}$ be the cumulant function, the Gaussian covariance matrix, and the Lévy measure of $\mu \in I D$. A distribution $\mu \in I D$ is said to have weak mean $m_{\mu}$ if $\lim _{a \rightarrow \infty} \int_{1<|x| \leqslant a} x \nu_{\mu}(d x)$ exists in $\mathbb{R}^{d}$ and if

$$
C_{\mu}(z)=-\frac{1}{2}\left\langle z, A_{\mu} z\right\rangle+\lim _{a \rightarrow \infty} \int_{|x| \leqslant a}\left(e^{i\langle z, x\rangle}-1-i\langle z, x\rangle\right) \nu_{\mu}(d x)+i\left\langle m_{\mu}, z\right\rangle .
$$

This concept was introduced by $[\mathrm{S}]$ recently. If $\mu \in I D$ has mean $m_{\mu}$ (that is, $\int_{\mathbb{R}^{d}}|x| \mu(d x)<\infty$ and $\int_{\mathbb{R}^{d}} x \mu(d x)=m_{\mu}$ ), then $\mu$ has weak mean $m_{\mu}$ (Remark 3.8 of [S]).

Section 2 begins with exact definitions of $f_{\alpha}, \bar{f}_{p, \alpha}$, and $l_{q, \alpha}$ and expounds existing results concerning $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$. Then, in Section 3, we will prove Theorem 1.1. In Section 4 we will give examples of $\Phi_{f}$ for which $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$ is different from those appearing in Theorem 1.1. Section 5 gives some concluding remarks.

## 2. Known Results

Let $-\infty<\alpha<2, p>0$, and $q>0$ and let

$$
\begin{aligned}
\bar{g}_{p, \alpha}(t) & =\frac{1}{\Gamma(p)} \int_{t}^{1}(1-u)^{p-1} u^{-\alpha-1} d u, \quad 0<t \leqslant 1 \\
j_{q, \alpha}(t) & =\frac{1}{\Gamma(q)} \int_{t}^{1}(-\log u)^{q-1} u^{-\alpha-1} d u, \quad 0<t \leqslant 1 \\
g_{\alpha}(t) & =\int_{t}^{\infty} u^{-\alpha-1} e^{-u} d u, \quad 0<t \leqslant \infty
\end{aligned}
$$

Let $t=\bar{f}_{p, \alpha}(s)$ for $0 \leqslant s<\bar{g}_{p, \alpha}(0+), t=l_{q, \alpha}(s)$ for $0 \leqslant s<j_{q, \alpha}(0+)$, and $t=f_{\alpha}(s)$ for $0 \leqslant s<g_{\alpha}(0+)$ be the inverse functions of $s=\bar{g}_{p, \alpha}(t), s=j_{q, \alpha}(t)$, and $s=g_{\alpha}(t)$, respectively. They are continuous, strictly decreasing functions. If $\alpha<0$, then
$\bar{g}_{p, \alpha}(0+), j_{q, \alpha}(0+)$, and $g_{\alpha}(0+)$ are finite and we define $\bar{f}_{p, \alpha}(s), l_{q, \alpha}(s)$, and $f_{\alpha}(s)$ to be zero for $s \geqslant \bar{g}_{p, \alpha}(0+), j_{q, \alpha}(0+)$, and $g_{\alpha}(0+)$, respectively. Let $\bar{\Phi}_{p, \alpha}, \Lambda_{q, \alpha}$, and $\Psi_{\alpha}$ denote $\Phi_{f}$ with $f=\bar{f}_{p, \alpha}, l_{q, \alpha}$, and $f_{\alpha}$, respectively. Let $K_{p, \alpha}, L_{q, \alpha}$, and $K_{\infty, \alpha}$ be the ranges of $\bar{\Phi}_{p, \alpha}, \Lambda_{q, \alpha}$, and $\Psi_{\alpha}$, respectively. These mappings and classes were systematically studied in Sato (2006) and [S]. In the following cases we have explicit expressions:

$$
\begin{aligned}
\bar{f}_{1, \alpha}(s) & =l_{1, \alpha}(s)= \begin{cases}(1-|\alpha| s)^{1 /|\alpha|} 1_{[0,1 /|\alpha|]}(s) & \text { for } \alpha<0, \\
e^{-s} & \text { for } \alpha=0, \\
(1+\alpha s)^{-1 / \alpha} & \text { for } \alpha>0,\end{cases} \\
\bar{f}_{p,-1}(s)=\left\{1-(\Gamma(p+1) s)^{1 / p}\right\} 1_{[0,1 / \Gamma(p+1)]}(s), & p>0, \\
l_{q, 0}(s) & =\exp \left(-(\Gamma(q+1) s)^{1 / q}\right), \quad q>0, \\
f_{-1}(s) & =(-\log s) 1_{[0,1]}(s) .
\end{aligned}
$$

In the case $p=q=1$ we have $\bar{\Phi}_{1, \alpha}=\Lambda_{1, \alpha}$ and $K_{1, \alpha}=L_{1 . \alpha}$, which are in essence treated earlier by Jurek (1988, 1989); $\bar{\Phi}_{1, \alpha}=\Lambda_{1, \alpha}$ were studied by Maejima et al. (2010a), and Maejima and Ueda (2010b) with the notation $\Phi_{\alpha}$. The mapping $\Lambda_{q, 0}$ and the class $L_{q, 0}$ with $q=1,2, \ldots$ coincide with those introduced by Jurek (1983) in a different form. A variant of $\Psi_{\alpha}$ is found in Grigelionis (2007).

A related family is

$$
G_{\alpha, \beta}(t)=\int_{t}^{\infty} u^{-\alpha-1} e^{-u^{\beta}} d u, \quad 0<t \leq \infty,
$$

for $-\infty<\alpha<2$ and $\beta>0$. Let $t=G_{\alpha, \beta}^{*}(s)$ for $0 \leq s<G_{\alpha, \beta}(0+)$ be the inverse function of $s=G_{\alpha, \beta}(t)$. If $\alpha<0$, then $G_{\alpha, \beta}(0+)$ is finite and we define $G_{\alpha, \beta}^{*}(s)=0$ for $s \geq G_{\alpha, \beta}(0+)$. Let $\Psi_{\alpha, \beta}$ denote $\Phi_{f}$ with $f=G_{\alpha, \beta}^{*}$. This was introduced by Maejima and Nakahara (2009) and studied by Maejima and Ueda (2010b) and, in the level of Lévy measures, by Maejima et al. (2010c). Clearly, $\Psi_{\alpha, 1}=\Psi_{\alpha}$. We have

$$
G_{-\beta, \beta}^{*}(s)=(-\log \beta s)^{1 / \beta} 1_{[0,1 / \beta]}(s), \quad \beta>0 .
$$

Earlier the mappings $\Psi_{0,2}$ and $\Psi_{-\beta, \beta}$ were treated in Aoyama et al. (2008) and Aoyama et al. (2010), respectively; $\Psi_{-2,2}$ appeared also in Arizmendi et al. (2010).

Maejima and Sato (2009) proved the following two results.
Proposition 2.1. Let $0<t_{0} \leqslant \infty$. Let $h(u)$ be a positive decreasing function on $\left(0, t_{0}\right)$ such that $\int_{0}^{t_{0}}\left(1+u^{2}\right) h(u) d u<\infty$. Let $g(t)=\int_{t}^{t_{0}} h(u) d u$ for $0<t \leqslant t_{0}$. Let $t=f(s), 0 \leqslant s<g(0+)$, be the inverse function of $s=g(t)$ and let $f(s)=0$ for $s \geqslant g(0+)$. Then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty}$.

Proposition 2.2. $\mathfrak{R}_{\infty}\left(\Psi_{0}\right)=L_{\infty}$.

It follows from Proposition 2.1 that $\Re_{\infty}\left(\Phi_{f}\right)=L_{\infty}$ for $f=\bar{f}_{p, \alpha}$ with $p \geqslant 1$ and $-1 \leqslant \alpha<0, f=l_{q, \alpha}$ with $q \geqslant 1$ and $-1 \leqslant \alpha<0, f=f_{\alpha}$ with $-1 \leqslant \alpha<0$, and $f=G_{\alpha, \beta}^{*}$ with $-1 \leqslant \alpha<0$ and $\beta>0$. The function $f_{0}$ for $\Psi_{0}=\Phi_{f_{0}}$ does not satisfy the condition in Proposition 2.1 but Proposition 2.2 is proved using the identity $\Psi_{0}=\Lambda_{1,0} \Psi_{-1}=\Psi_{-1} \Lambda_{1,0}$.

In November 2007-January 2008, Sato wrote four memos, showing the part related to $\Psi_{\alpha}$ in (ii), (iii), and (iv) of Theorem 1.1. But assertion (iii) for $\Psi_{1}$ was shown with the set $\{\mu \in I D: \mu$ has weak mean 0$\}$ replaced by the set of $\mu \in L_{\infty}$ satisfying some condition related to (4.6) of Sato (2006). At that time the concept of weak mean was not yet introduced. Those memos showed that some proper subclasses of $L_{\infty}$ appear as limit classes $\Re_{\infty}\left(\Phi_{f}\right)$.

Sato's memos were referred to by a series of papers Maejima and Ueda (2009a, b, 2010a, b) and Ichifuji et al. (2010). In Maejima and Ueda (2010a, c) they characterized $\mathfrak{R}\left(\Lambda_{1, \alpha}^{n}\right),-\infty<\alpha<2$, for $n=1,2, \ldots$, in relation to a decomposability which they called $\alpha$-selfdecomposability, and found $\mathfrak{R}_{\infty}\left(\Lambda_{1, \alpha}\right)$ for $-\infty<\alpha<2$. But the description of $\mathfrak{R}_{\infty}\left(\Lambda_{1,1}\right)$ was similar to Sato's memos. In Maejima and Ueda (2010b) they showed that $\Psi_{\alpha, \beta}$ with $-\infty<\alpha<2$ and $\beta>0$ satisfies $\mathfrak{R}_{\infty}\left(\Psi_{\alpha, \beta}\right)=\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)$, under the condition that $\alpha \neq 1+n \beta$ for $n=0,1,2, \ldots$. For $\Psi_{0,2}$ and $\Psi_{-\beta, \beta}$ with $\beta>0$, this result was earlier obtained by Aoyama et al. (2010). Further it was shown in Maejima and Ueda (2009b) that $\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)=\mathfrak{R}_{\infty}\left(\Lambda_{1, \alpha}\right)$ for $-\infty<\alpha<2$. An application of the result in Maejima and Ueda (2010a) was given in Ichifuji et al. (2010).

If $f(s)=b 1_{[0, a]}(s)$ for some $a>0$ and $b \neq 0$, then it is clear that $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=$ $\mathfrak{R}\left(\Phi_{f}\right)=I D$. A first example of $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$ satisfying $L_{\infty} \varsubsetneqq \mathfrak{R}_{\infty}\left(\Phi_{f}\right) \varsubsetneqq I D$ was given by Maejima and Ueda (2009a); they showed that if $f(s)=b^{-[s]}$ for a given $b>1$ with $[s]$ being the largest integer not exceeding $s$, then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty}(b)$, the smallest class that is closed under convolution and weak convergence and contains all semi-stable distributions on $\mathbb{R}^{d}$ with $b$ as a span; in this case $\mathfrak{R}\left(\Phi_{f}\right)$ is the class $L(b)$ of semi-selfdecomposable distributions on $\mathbb{R}^{d}$ with $b$ as a span. See Sato (1999) for the definitions of semi-stability, semi-selfdecomposability, and span. See Maejima et al. (2000) for characterization of $L_{\infty}(b)$ as the limit of the class $L_{n}(b)$ of $n$ times $b$-semi-selfdecomposable distributions and for description of the Lévy measures of distributions in $L_{\infty}(b)$. Recall that $L_{\infty} \varsubsetneqq L_{\infty}(b)$.

We have the following result in $[\mathrm{S}]$.
Proposition 2.3. The assertions related to $\Lambda_{q, \alpha}$ in (i), (ii), and (iv) of Theorem 1.1 are true.

Indeed, in [S], Theorem 7.3 says that $\Lambda_{q+q^{\prime}, \alpha}=\Lambda_{q^{\prime}, \alpha} \Lambda_{q, \alpha}$ for $\alpha \in(-\infty, 1) \cup(1,2)$, $q>0$, and $q^{\prime}>0$, and hence $\Lambda_{q, \alpha}^{n}=\Lambda_{n q, \alpha}$, and further, Theorem 7.11 combined with Proposition 6.8 describes $\bigcap_{q>0} L_{q, \alpha}$ for $\alpha \in(-\infty, 1) \cup(1,2)$.

## 3. Proof of Theorem 1.1

We prepare some lemmas. We use the terminology in $[\mathrm{S}]$ such as radial decomposition, monotonicity of order $p$, and complete monotonicity. In particular, our complete monotonicity implies vanishing at infinity. The location parameter $\gamma_{\mu}$ of $\mu \in I D$ is defined by

$$
C_{\mu}(z)=-\frac{1}{2}\left\langle z, A_{\mu} z\right\rangle+\int_{\mathbb{R}^{d}}\left(e^{i\langle z, x\rangle}-1-i\langle z, x\rangle 1_{\{|x| \leqslant 1\}}(x)\right) \nu_{\mu}(d x)+i\left\langle\gamma_{\mu}, z\right\rangle .
$$

Let $K_{p, \alpha}^{\mathrm{e}}\left[\right.$ resp. $\left.K_{\infty, \alpha}^{\mathrm{e}}\right]$ denote the class of distributions $\mu \in I D$ for which there exist $\rho \in I D$ and a function $q_{t}$ from $[0, \infty)$ into $\mathbb{R}^{d}$ such that $\int_{0}^{t} f_{p, \alpha}(s) d X_{s}^{(\rho)}-q_{t}$ [resp. $\left.\int_{0}^{t} f_{\alpha}(s) d X_{s}^{(\rho)}-q_{t}\right]$ converges in probability as $t \rightarrow \infty$ and the limit has distribution $\mu$.

Lemma 3.1. Let $-\infty<\alpha<2$ and $p>0$. The domains of $\bar{\Phi}_{p, \alpha}$ and $\Psi_{\alpha}$ are as follows:

$$
\begin{aligned}
& \mathfrak{D}\left(\bar{\Phi}_{p, \alpha}\right)=\mathfrak{D}\left(\Psi_{\alpha}\right) \\
& \quad= \begin{cases}I D & \text { for } \alpha<0, \\
\left\{\rho \in I D: \int_{|x|>1} \log |x| \nu_{\rho}(d x)<\infty\right\} & \text { for } \alpha=0, \\
\left\{\rho \in I D: \int_{|x|>1}|x|^{\alpha} \nu_{\rho}(d x)<\infty\right\} & \text { for } 0<\alpha<1, \\
\left\{\rho \in I D: \int_{|x|>1}|x| \nu_{\rho}(d x)<\infty, \int_{\mathbb{R}^{d}} x \rho(d x)=0,\right. & \\
\left.\quad \lim _{a \rightarrow \infty} \int_{1}^{a} s^{-1} d s \int_{|x|>s} x \nu_{\rho}(d x) \text { exists in } \mathbb{R}^{d}\right\} & \text { for } \alpha=1, \\
\left\{\rho \in I D: \int_{|x|>1}|x|^{\alpha} \nu_{\rho}(d x)<\infty, \int_{\mathbb{R}^{d}} x \rho(d x)=0\right\} & \text { for } 1<\alpha<2 .\end{cases}
\end{aligned}
$$

This is found in Sato (2006) or Theorems 4.2, 4.4 and Propositions 4.6, 5.1 of [S].
Lemma 3.2. Let $-\infty<\alpha<2$ and $p>0$. The class $K_{p, \alpha}^{\mathrm{e}}\left[\right.$ resp. $\left.K_{\infty, \alpha}^{\mathrm{e}}\right]$ is the totality of $\mu \in I D$ for which $\nu_{\mu}$ has a radial decomposition $\left(\lambda_{\mu}(d \xi), u^{-\alpha-1} k_{\xi}^{\mu}(u) d u\right)$ such that $k_{\xi}^{\mu}(u)$ is measurable in $(\xi, u)$ and, for $\lambda_{\mu}$-a.e. $\xi$, monotone of order $p$ [resp. completely
monotone] on $\mathbb{R}_{+}^{\circ}=(0, \infty)$ in $u$. The classes $K_{p, \alpha}$ and $K_{\infty, \alpha}$, that is, the ranges of $\bar{\Phi}_{p, \alpha}$ and $\Psi_{\alpha}$, are as follows:

$$
\begin{gathered}
K_{p, \alpha}= \begin{cases}K_{p, \alpha}^{e} & \text { for }-\infty<\alpha<1, \\
\left\{\mu \in K_{p, 1}^{e}: \mu \text { has weak mean } 0\right\} & \text { for } \alpha=1, \\
\left\{\mu \in K_{p, \alpha}^{e}: \mu \text { has mean } 0\right\} & \text { for } 1<\alpha<2,\end{cases} \\
K_{\infty, \alpha}= \begin{cases}K_{\infty, \alpha}^{\mathrm{e}} & \text { for }-\infty<\alpha<1, \\
\left\{\mu \in K_{\infty, 1}^{e}: \mu \text { has weak mean } 0\right\} & \text { for } \alpha=1, \\
\left\{\mu \in K_{\infty, \alpha}^{e}: \mu \text { has mean } 0\right\} & \text { for } 1<\alpha<2 .\end{cases}
\end{gathered}
$$

See Theorems 4.18, 5.8, and 5.10 of [S]. Note that if $\mu$ is in $K_{\infty, \alpha}^{\mathrm{e}}$ or $K_{p, \alpha}^{\mathrm{e}}$ with $0<\alpha<2$, then $\int_{\mathbb{R}^{d}}|x|^{\beta} \mu(d x)<\infty$ for $\beta \in(0, \alpha)$ (Propositions 4.16 and 5.13 of [S]). It follows from the lemma above that $K_{p, \alpha}^{\mathrm{e}} \supset K_{p^{\prime}, \alpha}^{\mathrm{e}}$ and $K_{p, \alpha} \supset K_{p^{\prime}, \alpha}$ for $p<p^{\prime}$ and that $K_{\infty, \alpha}^{\mathrm{e}}=\bigcap_{p>0} K_{p, \alpha}^{\mathrm{e}}$ and $K_{\infty, \alpha}=\bigcap_{p>0} K_{p, \alpha}$. In fact, this is the reason why we use the notation $K_{\infty, \alpha}^{\mathrm{e}}$ and $K_{\infty, \alpha}$.

Lemma 3.3. Let $\rho \in L_{\infty}$.
(i) Let $0<\alpha<2$. Then $\int_{\mathbb{R}^{d}}|x|^{\alpha} \rho(d x)<\infty$ if and only if $\Gamma_{\rho}((0, \alpha])=0$ and $\int_{(\alpha, 2)}(\beta-\alpha)^{-1} \Gamma_{\rho}(d \beta)<\infty$.
(ii) $\int_{|x|>1} \log |x| \rho(d x)<\infty$ if and only if $\int_{(0,2)} \beta^{-2} \Gamma_{\rho}(d \beta)<\infty$.

Proof. Assertion (i) is shown in Proposition 7.15 of [S]. Since

$$
\begin{gathered}
\int_{|x|>1} \log |x| \nu_{\rho}(d x)=\int_{(0,2)} \Gamma_{\rho}(d \beta) \int_{S} \lambda_{\beta}^{\rho}(d \xi) \int_{1}^{\infty}(\log |r \xi|) r^{-\beta-1} d r \\
=\int_{(0,2)} \Gamma_{\rho}(d \beta) \int_{1}^{\infty}(\log r) r^{-\beta-1} d r=\int_{(0,2)} \beta^{-2} \Gamma_{\rho}(d \beta)
\end{gathered}
$$

assertion (ii) follows.
Lemma 3.4. Let $\mu$ and $\rho$ be in $L_{\infty}^{(1,2)}$. Suppose that $\Gamma_{\rho}(d \beta)=(\beta-1) b(\beta) \Gamma_{\mu}(d \beta)$ and $\lambda_{\beta}^{\rho}=\lambda_{\beta}^{\mu}$ with a nonnegative measurable function $b(\beta)$ such that $(\beta-1)^{-1}(b(\beta)-1)$ is bounded on $(1,2)$. Then, $\int_{1}^{a} s^{-1} d s \int_{|x|>s} x \nu_{\rho}(d x)$ is convergent in $\mathbb{R}^{d}$ as $a \rightarrow \infty$ if and only if $\mu$ has weak mean $m_{\mu}$ for some $m_{\mu}$.

Proof. Notice that $b(\beta)$ is bounded on $(1,2)$ and that $\int_{|x|>1}|x| \nu_{\rho}(d x)<\infty$ by Lemma 3.3. We have

$$
\begin{gathered}
\int_{1}^{a} s^{-1} d s \int_{|x|>s} x \nu_{\rho}(d x)=\int_{1}^{a} s^{-1} d s \int_{(1,2)} \Gamma_{\rho}(d \beta) \int_{S} \xi \lambda_{\beta}^{\rho}(d \xi) \int_{s}^{\infty} r^{-\beta} d r \\
=\int_{(1,2)} b(\beta) \Gamma_{\mu}(d \beta) \int_{S} \xi \lambda_{\beta}^{\mu}(d \xi) \int_{1}^{a} s^{-\beta} d s=I_{1} \quad \text { (say) }
\end{gathered}
$$

and

$$
\int_{1<|x| \leqslant a} x \nu_{\mu}(d x)=\int_{(1,2)} \Gamma_{\mu}(d \beta) \int_{S} \xi \lambda_{\beta}^{\mu}(d \xi) \int_{1}^{a} r^{-\beta} d r=I_{2} \quad \text { (say). }
$$

Hence

$$
I_{1}-I_{2}=\int_{(1,2)}(b(\beta)-1) \Gamma_{\mu}(d \beta) \int_{S} \xi \lambda_{\beta}^{\mu}(d \xi) \int_{1}^{a} r^{-\beta} d r
$$

Since

$$
\left|(b(\beta)-1) \int_{1}^{a} r^{-\beta} d r\right| \leqslant(\beta-1)^{-1}|b(\beta)-1|
$$

and $\int_{1}^{a} r^{-\beta} d r$ tends to $(\beta-1)^{-1}, I_{1}-I_{2}$ is convergent in $\mathbb{R}^{d}$ as $a \rightarrow \infty$. Hence $I_{1}$ is convergent if and only if $I_{2}$ is convergent.

Lemma 3.5. Let $f$ and $h$ be locally square-integrable functions on $\mathbb{R}_{+}$. Assume that there is $s_{0} \in(0, \infty)$ such that $h(s)=0$ for $s \geqslant s_{0}$ and that $\Phi_{h}$ is one-to-one. Then $\Phi_{f} \Phi_{h}=\Phi_{h} \Phi_{f}$.

Proof. Let $f_{t}(s)=f(s) 1_{[0, t]}(s)$. Then $\Phi_{f_{t}} \Phi_{h}=\Phi_{h} \Phi_{f_{t}}$ by Lemma 3.6 of Maejima and Sato (2009). Let $\rho \in \mathfrak{D}\left(\Phi_{f}\right)$. Then $\Phi_{f_{t}} \rho \rightarrow \Phi_{f} \rho$ as $t \rightarrow \infty$ by the definition of $\Phi_{f}$. Hence $\Phi_{h} \Phi_{f_{t}} \rho \rightarrow \Phi_{h} \Phi_{f} \rho$ by (3.1) of Maejima and Sato (2009). It follows that $\Phi_{f_{t}} \Phi_{h} \rho \rightarrow \Phi_{h} \Phi_{f} \rho$. Since the convergence of $\int_{0}^{t} f(s) d X_{s}^{\left(\Phi_{h} \rho\right)}$ in law implies its convergence in probability, $\Phi_{h} \rho$ is in $\mathfrak{D}\left(\Phi_{f}\right)$ and $\Phi_{f} \Phi_{h} \rho=\Phi_{h} \Phi_{f} \rho$. Conversely, suppose that $\rho \in I D$ satisfies $\Phi_{h} \rho \in \mathfrak{D}\left(\Phi_{f}\right)$. Then $\Phi_{h} \Phi_{f_{t}} \rho=\Phi_{f_{t}} \Phi_{h} \rho \rightarrow \Phi_{f} \Phi_{h} \rho$ as $t \rightarrow \infty$. Looking at (3.8) of Maejima and Sato (2009), we see that $\int_{0}^{s_{0}} h(s) \neq 0$ from the one-to-one property of $\Phi_{h}$. Hence $\left\{\Phi_{f_{t}} \rho: t>0\right\}$ is precompact by the argument in pp. 138-139 of Maejima and Sato (2009). Hence, again from the one-to-one property of $\Phi_{h}, \Phi_{f_{t}} \rho$ is convergent as $t \rightarrow \infty$, that is, $\rho \in \mathfrak{D}\left(\Phi_{f}\right)$.

Lemma 3.6. Let $f$ be locally square-integrable on $\mathbb{R}_{+}$. Suppose that there is $\beta \geqslant 0$ such that any $\mu \in \mathfrak{R}\left(\Phi_{f}\right)$ has Lévy measure $\nu_{\mu}$ with a radial decomposition $\left(\lambda_{\mu}(d \xi)\right.$, $\left.u^{\beta} l_{\xi}^{\mu}(u) d u\right)$ where $l_{\xi}^{\mu}(u)$ is measurable in $(\xi, u)$ and decreasing on $\mathbb{R}_{+}^{\circ}$ in $u$. Then

$$
\mathfrak{R}_{\infty}\left(\Phi_{f}\right) \subset \mathfrak{R}_{\infty}\left(\Lambda_{1,-\beta-1}\right)=L_{\infty}
$$

Proof. Clearly $l_{\xi}^{\mu} \geqslant 0$ for $\lambda_{\mu}$-a. e. $\xi$. Since $\int_{|x|>1} \nu_{\mu}(d x)<\infty$, we have $\lim _{u \rightarrow \infty} l_{\xi}^{\mu}(u)=0$ for $\lambda_{\mu}$-a. e. $\xi$. Hence we can modify $l_{\xi}^{\mu}(u)$ in such a way that $l_{\xi}^{\mu}(u)$ is monotone of order 1 in $u \in \mathbb{R}_{+}^{\circ}$. Recall that a function is monotone of order 1 on $\mathbb{R}_{+}^{\circ}$ if and only if it is decreasing, right-continuous, and vanishing at infinity (Proposition 2.11 of $[\mathrm{S}]$ ). Then it follows from Theorem 4.18 or 6.12 of $[\mathrm{S}]$ that

$$
\begin{equation*}
\mathfrak{R}\left(\Phi_{f}\right) \subset \mathfrak{R}\left(\Lambda_{1,-\beta-1}\right) . \tag{3.1}
\end{equation*}
$$

Let us write $\Lambda=\Lambda_{1,-\beta-1}$ for simplicity. We have $\Phi_{f} \Lambda=\Lambda \Phi_{f}$ by virtue of Lemma 3.5 , since $\Lambda$ is one-to-one (Theorem 6.14 of $[\mathrm{S}]$ ). If $\Phi_{f} \Lambda^{n}=\Lambda^{n} \Phi_{f}$ for some integer $n \geqslant 1$, then

$$
\Phi_{f} \Lambda^{n+1}=\Phi_{f} \Lambda \Lambda^{n}=\Lambda \Phi_{f} \Lambda^{n}=\Lambda \Lambda^{n} \Phi_{f}=\Lambda^{n+1} \Phi_{f}
$$

Hence $\Phi_{f} \Lambda^{n}=\Lambda^{n} \Phi_{f}$ for $n=1,2, \ldots$. Now we claim that

$$
\begin{equation*}
\mathfrak{R}\left(\Phi_{f}^{n}\right) \subset \mathfrak{R}\left(\Lambda^{n}\right) \tag{3.2}
\end{equation*}
$$

for $n=1,2, \ldots$. Indeed, this is true for $n=1$ by (3.1); if (3.2) is true for $n$, then any $\mu \in \mathfrak{R}\left(\Phi_{f}^{n+1}\right)$ has expression

$$
\mu=\Phi_{f}^{n+1} \rho=\Phi_{f} \Phi_{f}^{n} \rho=\Phi_{f} \Lambda^{n} \rho^{\prime}=\Lambda^{n} \Phi_{f} \rho^{\prime}=\Lambda^{n} \Lambda \rho^{\prime \prime}=\Lambda^{n+1} \rho^{\prime \prime}
$$

for some $\rho \in \mathfrak{D}\left(\Phi_{f}^{n+1}\right)$, $\rho^{\prime} \in \mathfrak{D}\left(\Lambda^{n}\right)$ with $\Phi_{f}^{n} \rho=\Lambda^{n} \rho^{\prime}$, and $\rho^{\prime \prime} \in \mathfrak{D}(\Lambda)$ with $\Phi_{f} \rho^{\prime}=\Lambda \rho^{\prime \prime}$, which means (3.2) for $n+1$. It follows from (3.2) that $\mathfrak{R}_{\infty}\left(\Phi_{f}\right) \subset \mathfrak{R}_{\infty}(\Lambda)$. The equality $\mathfrak{R}_{\infty}(\Lambda)=L_{\infty}$ is from Proposition 2.3.

Proof of the part related to $\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)$ in Theorem 1.1. The result for $-1 \leqslant \alpha \leqslant 0$ is already known (see Propositions 2.1 and 2.2). But the proof below also includes this case. First, using Lemma 3.2, notice that Lemma 3.6 is applicable to $\Phi_{f}=\Psi_{\alpha}$ and $\beta=(-\alpha-1) \vee 0$.

Case $1(-\infty<\alpha<0)$. We have $\mathfrak{D}\left(\Psi_{\alpha}\right)=I D$ in Lemma 3.1. Let us show that

$$
\begin{equation*}
\Psi_{\alpha}\left(L_{\infty}\right)=L_{\infty} \tag{3.3}
\end{equation*}
$$

Let $\rho \in L_{\infty}$ and $\mu=\Psi_{\alpha} \rho$. Then for $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$, where $\mathcal{B}\left(\mathbb{R}^{d}\right)$ is the class of Borel sets in $\mathbb{R}^{d}$,

$$
\begin{aligned}
\nu_{\mu}(B) & =\int_{0}^{\infty} d s \int_{\mathbb{R}^{d}} 1_{B}\left(f_{\alpha}(s) x\right) \nu_{\rho}(d x)=\int_{0}^{\infty} t^{-\alpha-1} e^{-t} d t \int_{\mathbb{R}^{d}} 1_{B}(t x) \nu_{\rho}(d x) \\
& =\int_{0}^{\infty} t^{-\alpha-1} e^{-t} d t \int_{(0,2)} \Gamma_{\rho}(d \beta) \int_{S} \lambda_{\beta}^{\rho}(d \xi) \int_{0}^{\infty} 1_{B}(\operatorname{tr} \xi) r^{-\beta-1} d r \\
& =\int_{(0,2)} \Gamma(\beta-\alpha) \Gamma_{\rho}(d \beta) \int_{S} \lambda_{\beta}^{\rho}(d \xi) \int_{0}^{\infty} 1_{B}(u \xi) u^{-\beta-1} d u
\end{aligned}
$$

Hence $\mu \in L_{\infty}$ with

$$
\begin{equation*}
\Gamma_{\mu}(d \beta)=\Gamma(\beta-\alpha) \Gamma_{\rho}(d \beta) \quad \text { and } \quad \lambda_{\beta}^{\mu}=\lambda_{\beta}^{\rho} . \tag{3.4}
\end{equation*}
$$

Let us show the converse. Let $\mu \in L_{\infty}$. In order to find $\rho \in L_{\infty}$ satisfying $\Psi_{\alpha} \rho=\mu$, it suffices to choose $\Gamma_{\rho}, \lambda_{\beta}^{\rho}, A_{\rho}$, and $\gamma_{\rho}$ such that (3.4) holds and

$$
\begin{equation*}
A_{\mu}=\int_{0}^{\infty} f_{\alpha}(s)^{2} d s A_{\rho} \tag{3.5}
\end{equation*}
$$

$$
\begin{equation*}
\gamma_{\mu}=\int_{0}^{\infty-} f_{\alpha}(s) d s\left(\gamma_{\rho}+\int_{\mathbb{R}^{d}} x\left(1_{\left\{\left|f_{\alpha}(s) x\right| \leqslant 1\right\}}-1_{\{|x| \leqslant 1\}}\right) \nu_{\rho}(d x)\right) \tag{3.6}
\end{equation*}
$$

(see Proposition 3.18 of $[\mathrm{S}]$ ). This choice is possible, because $\inf _{\beta \in(0,2)} \Gamma(\beta-\alpha)>0$, $\int_{0}^{\infty} f_{\alpha}(s) d s=\int_{0}^{\infty} t^{-\alpha} e^{-t} d t=\Gamma(1-\alpha), \int_{0}^{\infty} f_{\alpha}(s)^{2} d s=\int_{0}^{\infty} t^{1-\alpha} e^{-t} d t=\Gamma(2-\alpha)$, and

$$
\begin{aligned}
\int_{0}^{\infty} & f_{\alpha}(s) d s \int_{\mathbb{R}^{d}}|x|\left|1_{\left\{\left|f_{\alpha}(s) x\right| \leqslant 1\right\}}-1_{\{|x| \leqslant 1\}}\right| \nu_{\rho}(d x) \\
& =\int_{0}^{\infty} t^{-\alpha} e^{-t} d t \int_{\mathbb{R}^{d}}|x|\left|1_{\{|t x| \leqslant 1\}}-1_{\{|x| \leqslant 1\}}\right| \nu_{\rho}(d x) \\
& =\int_{0}^{1} t^{-\alpha} e^{-t} d t \int_{1<|x| \leqslant 1 / t}|x| \nu_{\rho}(d x)+\int_{1}^{\infty} t^{-\alpha} e^{-t} d t \int_{1 / t<|x| \leqslant 1}|x| \nu_{\rho}(d x) \\
& =\int_{|x|>1}|x| \nu_{\rho}(d x) \int_{0}^{1 /|x|} t^{-\alpha} e^{-t} d t+\int_{|x| \leqslant 1}|x| \nu_{\rho}(d x) \int_{1 /|x|}^{\infty} t^{-\alpha} e^{-t} d t<\infty
\end{aligned}
$$

since $\int_{0}^{1 /|x|} t^{-\alpha} e^{-t} d t \sim(1-\alpha)^{-1}|x|^{\alpha-1}$ as $|x| \rightarrow \infty$ and $\int_{1 /|x|}^{\infty} t^{-\alpha} e^{-t} d t \sim|x|^{\alpha} e^{-1 /|x|}$ as $|x| \downarrow 0$. Therefore (3.3) is true. It follows that $\Psi_{\alpha}^{n}\left(L_{\infty}\right)=L_{\infty}$ for $n=1,2 \ldots$. Hence $\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right) \supset L_{\infty}$. On the other hand, $\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right) \subset L_{\infty}$ by virtue of Lemma 3.6.

Case $2(0 \leqslant \alpha<1)$. Since $\mathfrak{D}\left(\Psi_{\alpha}\right)$ is as in Lemma 3.1, it follows from Lemma 3.3 that

$$
L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)= \begin{cases}\left\{\rho \in L_{\infty}: \int_{(0,2)} \beta^{-2} \Gamma_{\rho}(d \beta)<\infty\right\}, & \alpha=0 \\ \left\{\rho \in L_{\infty}^{(\alpha, 2)}: \int_{(\alpha, 2)}(\beta-\alpha)^{-1} \Gamma_{\rho}(d \beta)<\infty\right\}, & 0<\alpha<1\end{cases}
$$

We have

$$
\begin{equation*}
\Psi_{\alpha}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)\right)=L_{\infty}^{(\alpha, 2)} \tag{3.7}
\end{equation*}
$$

where $L_{\infty}^{(0,2)}=L_{\infty}$. Indeed, if $\rho \in L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)$ and $\mu=\Psi_{\alpha} \rho$, then we have $\mu \in L_{\infty}^{(\alpha, 2)}$ and (3.4), using $\Gamma(\beta-\alpha)=(\beta-\alpha)^{-1} \Gamma(\beta-\alpha+1)$ for $0 \leqslant \alpha<1$. Conversely, if $\mu \in$ $L_{\infty}^{(\alpha, 2)}$, then we can find $\rho \in L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)$ satisfying $\mu=\Psi_{\alpha} \rho$ in the same way as in Case 1 ; when $\alpha=0$, we have $\int_{(0,2)} \beta^{-2} \Gamma_{\rho}(d \beta)<\infty$ since $\Gamma_{\rho}(d \beta)=\beta(\Gamma(\beta+1))^{-1} \Gamma_{\mu}(d \beta)$ and $\int_{(0,2)} \beta^{-1} \Gamma_{\mu}(d \beta)<\infty$. Hence (3.7) holds. Now we have

$$
\begin{equation*}
\Psi_{\alpha}^{n}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n}\right)\right)=L_{\infty}^{(\alpha, 2)} \tag{3.8}
\end{equation*}
$$

for $n=1,2, \ldots$. Indeed, it is true for $n=1$ by (3.7) and, if (3.8) is true for $n$, then

$$
\begin{aligned}
L_{\infty}^{(\alpha, 2)} & =\Psi_{\alpha}^{n}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n}\right)\right)=\Psi_{\alpha}^{n}\left(L_{\infty}^{(\alpha, 2)} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n}\right)\right) \\
& =\Psi_{\alpha}^{n}\left(\Psi_{\alpha}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)\right) \cap \mathfrak{D}\left(\Psi_{\alpha}^{n}\right)\right) \\
& =\Psi_{\alpha}^{n}\left(\Psi_{\alpha}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n+1}\right)\right)\right)=\Psi_{\alpha}^{n+1}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n+1}\right)\right) .
\end{aligned}
$$

It follows from (3.8) that $L_{\infty}^{(\alpha, 2)} \subset \mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right)$. Next we claim that

$$
\begin{equation*}
\mathfrak{R}\left(\Psi_{\alpha}\right) \cap L_{\infty} \subset L_{\infty}^{(\alpha, 2)} . \tag{3.9}
\end{equation*}
$$

Let $\mu \in \mathfrak{R}\left(\Psi_{\alpha}\right) \cap L_{\infty}$. Then $\mu$ has a radial decomposition $\left(\lambda_{\mu}(d \xi), r^{-\alpha-1} k_{\xi}^{\mu}(r) d r\right)$ with the property stated in Lemma 3.2. On the other hand,

$$
\begin{aligned}
\nu_{\mu}(B) & =\int_{(0,2)} \Gamma_{\mu}(d \beta) \int_{S} \lambda_{\beta}^{\mu}(d \xi) \int_{0}^{\infty} 1_{B}(r \xi) r^{-\beta-1} d r \\
& =\int_{S} \bar{\lambda}_{\mu}(d \xi) \int_{(0,2)} \Gamma_{\xi}^{\mu}(d \beta) \int_{0}^{\infty} 1_{B}(r \xi) r^{-\beta-1} d r
\end{aligned}
$$

for $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$, as there are a probability measure $\bar{\lambda}_{\mu}$ on $S$ and a measurable family $\left\{\Gamma_{\xi}^{\mu}\right\}$ of measures on $(0,2)$ satisfying $\int_{(0,2)}\left(\beta^{-1}+(2-\beta)^{-1}\right) \Gamma_{\xi}^{\mu}(d \beta)=$ const such that $\Gamma_{\mu}(d \beta) \lambda_{\beta}^{\mu}(d \xi)=\bar{\lambda}_{\mu}(d \xi) \Gamma_{\xi}^{\mu}(d \beta)$. Hence, by the uniqueness in Proposition 3.1 of [S], there is a positive, finite, measurable function $c(\xi)$ such that $\lambda_{\mu}(d \xi)=c(\xi) \bar{\lambda}_{\mu}(d \xi)$ and, for $\lambda_{\mu}$-a.e. $\xi, r^{-\alpha-1} k_{\xi}^{\mu}(r) d r=c(\xi)^{-1}\left(\int_{(0,2)} r^{-\beta-1} \Gamma_{\xi}^{\mu}(d \beta)\right) d r$. Hence $k_{\xi}^{\mu}(r)=$ $c(\xi)^{-1} \int_{(0,2)} r^{\alpha-\beta} \Gamma_{\xi}^{\mu}(d \beta)$, a. e. $r$. Since $k_{\xi}^{\mu}(r)$ is completely monotone, it vanishes as $r$ goes to infinity. Hence $\Gamma_{\xi}^{\mu}((0, \alpha])=0$ for $\lambda_{\mu}$-a. e. $\xi$. Hence $\Gamma_{\mu}((0, \alpha])=0$, that is, $\mu \in L_{\infty}^{(\alpha, 2)}$, proving (3.9). Now, using Lemma 3.6, we obtain $\mathfrak{R}_{\infty}\left(\Psi_{\alpha}\right) \subset \mathfrak{R}\left(\Psi_{\alpha}\right) \cap L_{\infty} \subset$ $L_{\infty}^{(\alpha, 2)}$.

Case $3(\alpha=1)$. Let us show that

$$
\begin{equation*}
\Psi_{1}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{1}\right)\right)=L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\} \tag{3.10}
\end{equation*}
$$

Let $\rho \in L_{\infty} \cap \mathfrak{D}\left(\Psi_{1}\right)$, that is, $\rho \in L_{\infty}^{(1,2)}, \int_{(1,2)}(\beta-1)^{-1} \Gamma_{\rho}(d \beta)<\infty, \int_{\mathbb{R}^{d}} x \rho(d x)=0$, and $\lim _{a \rightarrow \infty} \int_{1}^{a} s^{-1} d s \int_{|x|>s} x \nu_{\rho}(d x)$ exists in $\mathbb{R}^{d}$. Let $\mu=\Psi_{1} \rho$. Then, as in Case $1, \mu \in L_{\infty}^{(1,2)}$ and (3.4) holds with $\alpha=1$. By Lemma 3.2, $\mu$ has weak mean 0 . Conversely, let $\mu \in L_{\infty}^{(1,2)} \cap\{\mu \in I D$ : weak mean 0$\}$. Choose $\rho \in L_{\infty}^{(1,2)}$ such that $\Gamma_{\rho}(d \beta)=(\Gamma(\beta-1))^{-1} \Gamma_{\mu}(d \beta), \lambda_{\beta}^{\rho}=\lambda_{\beta}^{\mu}, A_{\rho}=A_{\mu}$, and $\gamma_{\rho}=-\int_{|x|>1} x \nu_{\rho}(d x)$ (note that $\int_{(1,2)}(\beta-1)^{-1} \Gamma_{\rho}(d \beta)<\infty$ and hence $\int_{|x|>1}|x| \nu_{\rho}(d x)<\infty$ by Lemma 3.3). Then $\int_{\mathbb{R}^{d}} x \rho(d x)=0$ (see Lemma 4.3 of $[\mathrm{S}]$ ). Since $\mu$ has weak mean, $\int_{1}^{a} s^{-1} d s \int_{|x|>s} x \nu_{\rho}(d x)$ is convergent as $a \rightarrow \infty$ by application of Lemma 3.4 with $b(\beta)=1 / \Gamma(\beta)$. Hence $\rho \in \mathfrak{D}\left(\Psi_{1}\right)$. We have $\nu_{\Psi_{1} \rho}=\nu_{\mu}, A_{\Psi_{1} \rho}=A_{\mu}$, and $\Psi_{1} \rho$ has weak mean 0 . Among distributions $\mu^{\prime} \in I D$ having $\nu_{\mu^{\prime}}=\nu_{\mu}$ and $A_{\mu^{\prime}}=A_{\mu}$, only one distribution has weak mean 0 . Hence $\Psi_{1} \rho=\mu$. This proves (3.10). We have

$$
\begin{equation*}
\Psi_{1}^{n}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{1}^{n}\right)\right)=L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\}, \quad n=1,2, \ldots \tag{3.11}
\end{equation*}
$$

from (3.10) by the same argument as in Case 2. Hence

$$
\begin{equation*}
L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\} \subset \mathfrak{R}_{\infty}\left(\Psi_{1}\right) . \tag{3.12}
\end{equation*}
$$

Next

$$
\begin{equation*}
\mathfrak{R}\left(\Psi_{1}\right) \cap L_{\infty} \subset L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\} . \tag{3.13}
\end{equation*}
$$

Indeed, $\mathfrak{R}\left(\Psi_{1}\right) \cap L_{\infty} \subset L_{\infty}^{(1,2)}$ by the same argument as in Case 2. Any $\mu \in \mathfrak{R}\left(\Psi_{1}\right)$ has weak mean 0 by Lemma 3.2. Now it follows from Lemma 3.6 that

$$
\begin{equation*}
\Re_{\infty}\left(\Psi_{1}\right) \subset L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\} \tag{3.14}
\end{equation*}
$$

Case $4(1<\alpha<2)$. We show that

$$
\begin{equation*}
\Psi_{\alpha}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)\right)=L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D: \text { mean } 0\} \tag{3.15}
\end{equation*}
$$

Let $\rho \in L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)$, that is, $\rho \in L_{\infty}^{(\alpha, 2)}, \int_{(\alpha, 2)}(\beta-\alpha)^{-1} \Gamma_{\rho}(d \beta)<\infty$, and $\int_{\mathbb{R}^{d}} x \rho(d x)=0$ (Lemmas 3.1 and 3.3). Let $\mu=\Psi_{\alpha} \rho$. Then $\mu \in L_{\infty}^{(\alpha, 2)}$ and (3.4) holds. Hence $\int_{\mathbb{R}^{d}}|x| \mu(d x)<\infty$ by Lemma 3.3 and $\mu$ has mean 0 by Lemma 3.2. Conversely, if $\mu \in L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D:$ mean 0$\}$, then we can find $\rho \in L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}\right)$ satisfying $\Psi_{\alpha} \rho=\mu$, similarly to Case 3 . Hence (3.15) is true. It follows that

$$
\Psi_{\alpha}^{n}\left(L_{\infty} \cap \mathfrak{D}\left(\Psi_{\alpha}^{n}\right)\right)=L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D: \text { mean } 0\}, \quad n=1,2, \ldots
$$

similarly to Cases 2 and 3 . Hence

$$
\begin{equation*}
L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D: \text { mean } 0\} \subset \Re_{\infty}\left(\Psi_{\alpha}\right) . \tag{3.16}
\end{equation*}
$$

We can also prove

$$
\mathfrak{R}\left(\Psi_{\alpha}\right) \cap L_{\infty} \subset L_{\infty}^{(\alpha, 2)} \cap\{\mu \in I D: \text { mean } 0\}
$$

similarly to Cases 2 and 3. Hence the reverse inclusion of (3.16) follows from Lemma 3.6.

Proof of the part related to $\mathfrak{R}_{\infty}\left(\bar{\Phi}_{p, \alpha}\right)$ in Theorem 1.1. We assume $p \geqslant 1$. Since monotonicity of order $p \in[1 . \infty$ ) implies monotonicity of order 1 (Corollary 2.6 of $[\mathrm{S}])$, it follows from Lemma 3.2 that Lemma 3.6 is applicable with $\beta=(-\alpha-1) \vee 0$. Hence $\Re_{\infty}\left(\bar{\Phi}_{p, \alpha}\right) \subset L_{\infty}$. If $\rho \in L_{\infty} \cap \mathfrak{D}\left(\bar{\Phi}_{p, \alpha}\right)$ and $\bar{\Phi}_{p, \alpha} \rho=\mu$, then $\rho \in L_{\infty}^{(\alpha, 2)}$ (understand that $L_{\infty}^{(\alpha, 2)}=L_{\infty}$ for $\alpha \leqslant 0$ ) and, noting that

$$
\begin{aligned}
\nu_{\mu}(B) & =\int_{0}^{\infty} d s \int_{\mathbb{R}^{d}} 1_{B}\left(\bar{f}_{p, \alpha}(s) x\right) \nu_{\rho}(d x)=\frac{1}{\Gamma(p)} \int_{0}^{1} t^{-\alpha-1}(1-t)^{p-1} d t \int_{\mathbb{R}^{d}} 1_{B}(t x) \nu_{\rho}(d x) \\
& =\frac{1}{\Gamma(p)} \int_{0}^{1} t^{-\alpha-1}(1-t)^{p-1} d t \int_{(0,2)} \Gamma_{\rho}(d \beta) \int_{S} \lambda_{\beta}^{\rho}(d \xi) \int_{0}^{\infty} 1_{B}(t r \xi) r^{-\beta-1} d r \\
& =\int_{(0,2)} \frac{\Gamma(\beta-\alpha)}{\Gamma(\beta-\alpha+p)} \Gamma_{\rho}(d \beta) \int_{S} \lambda_{\beta}^{\rho}(d \xi) \int_{0}^{\infty} 1_{B}(u \xi) u^{-\beta-1} d u
\end{aligned}
$$

and recalling Lemmas 3.1 and 3.3, we obtain $\mu \in L_{\infty}^{(\alpha, 2)}$ with

$$
\begin{equation*}
\Gamma_{\mu}(d \beta)=\frac{\Gamma(\beta-\alpha)}{\Gamma(\beta-\alpha+p)} \Gamma_{\rho}(d \beta) \quad \text { and } \quad \lambda_{\beta}^{\mu}=\lambda_{\beta}^{\rho} . \tag{3.17}
\end{equation*}
$$

Now the proof of assertions (i), (ii), and (iv) can be given in parallel to the corresponding assertions for $\Psi_{\alpha}$. Note that, if $-\infty<\alpha<1$, then

$$
\int_{0}^{\infty} \bar{f}_{p, \alpha}(s) d s \int_{\mathbb{R}^{d}}|x|\left|1_{\left\{\left|\bar{f}_{p, \alpha}(s) x\right| \leqslant 1\right\}}-1_{\{|x| \leqslant 1\}}\right| \nu_{\rho}(d x)<\infty
$$

similarly. We also use the fact that $k_{\xi}^{\mu}(r)$ vanishes at infinity if it is monotone of order $p \in[1, \infty)$.

For assertion (iii) in the case $\alpha=1$, we have to find another way, as Lemma 3.4 is not applicable if $\beta>1$. Let us show

$$
\begin{equation*}
\bar{\Phi}_{p, 1}\left(L_{\infty} \cap \mathfrak{D}\left(\bar{\Phi}_{p, 1}\right)\right)=L_{\infty}^{(1,2)} \cap\{\mu \in I D: \text { weak mean } 0\} . \tag{3.18}
\end{equation*}
$$

Suppose that $\rho \in L_{\infty} \cap \mathfrak{D}\left(\bar{\Phi}_{p, 1}\right)$ and $\bar{\Phi}_{p, 1} \rho=\mu$. Then $\rho \in L_{\infty}^{(1,2)}, \int_{(1,2)}(\beta-1)^{-1} \Gamma_{\rho}(d \beta)<$ $\infty, \mu \in L_{\infty}^{(1,2)}$ with (3.17), and $\mu$ has weak mean 0 by Lemma 3.2. Conversely, suppose that $\mu \in L_{\infty}^{(1,2)}$ with weak mean 0 . As in $[\mathrm{S}]$, let $\mathfrak{M}^{L}$ be the class of Lévy measures of infinitely divisible distributions on $\mathbb{R}^{d}$ and let $\bar{\Phi}_{p, 1}^{L}$ be the transformation of Lévy measures associated with the mapping $\bar{\Phi}_{p, 1}$. Define $\Gamma_{0}(d \beta)=\frac{\Gamma(\beta-1+p)}{\Gamma(\beta-1)} \Gamma_{\mu}(d \beta)$. Then $\int_{(1,2)}(2-\beta)^{-1} \Gamma_{0}(d \beta)<\infty$. Define

$$
\nu_{0}(B)=\int_{(1,2)} \Gamma_{0}(d \beta) \int_{S} \lambda_{\beta}^{\mu}(d \xi) \int_{0}^{\infty} 1_{B}(r \xi) r^{-\beta-1} d r
$$

for $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$. We have $\nu_{0} \in \mathfrak{M}^{L}$. We see

$$
\begin{aligned}
\nu_{\mu}(B) & =\int_{(1,2)} \frac{\Gamma(\beta-1)}{\Gamma(\beta-1+p)} \Gamma_{0}(d \beta) \int_{S} \lambda_{\beta}^{\mu}(d \xi) \int_{0}^{\infty} 1_{B}(u \xi) u^{-\beta-1} d u \\
& =\int_{0}^{\infty} d s \int_{\mathbb{R}^{d}} 1_{B}\left(\bar{f}_{p, 1}(s) x\right) \nu_{0}(d x)
\end{aligned}
$$

from the calculation above. Since $\nu_{\mu} \in \mathfrak{M}^{L}$, we have $\nu_{0} \in \mathfrak{D}\left(\bar{\Phi}_{p, 1}^{L}\right)$ and $\bar{\Phi}_{p, 1}^{L} \nu_{0}=$ $\nu_{\mu}$. Then it follows from Theorem 4.10 of $[\mathrm{S}]$ that $\nu_{\mu}$ has a radial decomposition $\left(\lambda_{\mu}(d \xi), u^{-2} k_{\xi}^{\mu}(u) d u\right)$ such that $k_{\xi}^{\mu}(u)$ is measurable in $(\xi, u)$ and, for $\lambda_{\mu}$-a. e. $\xi$, monotone of order $p$ in $u \in \mathbb{R}_{+}^{\circ}$. Hence $\mu \in \mathfrak{R}\left(\bar{\Phi}_{p, 1}\right)$ from Lemma 3.2. Since $\bar{\Phi}_{p, 1}^{L} \nu_{0}=\nu_{\mu}$ and $\bar{\Phi}_{p, 1}^{L}$ is one-to-one (Theorem 4.9 of $[\mathrm{S}]$ ), we have $\mu=\bar{\Phi}_{p, 1} \rho$ for some $\rho \in \mathfrak{D}\left(\bar{\Phi}_{p, 1}\right)$ with $\nu_{\rho}=\nu_{0}$. It follows that $\rho \in L_{\infty}$. This finishes the proof of (3.18). Now we can show (3.11)-(3.14) with $\bar{\Phi}_{p, 1}$ in place of $\Psi_{1}$ similarly to Case 3 in the preceding proof.

Proof of the part related to $\mathfrak{R}_{\infty}\left(\Lambda_{q, \alpha}\right)$ in Theorem 1.1. Since we have Proposition 2.3, it remains only to consider $\Lambda_{1,1}$. But the assertion for $\mathfrak{R}_{\infty}\left(\Lambda_{1,1}\right)$ is obviously true, since $\Lambda_{1,1}=\bar{\Phi}_{1,1}$.

## 4. Some examples of $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$

We present some examples of $\Phi_{f}$ for which the class $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$ is different from those appearing in Theorem 1.1.

Define $T_{a}$, the dilation by $a \in \mathbb{R} \backslash\{0\}$, as $\left(T_{a} \mu\right)(B)=\int 1_{B}(a x) \mu(d x)=\mu((1 / a) B)$, $B \in \mathcal{B}\left(\mathbb{R}^{d}\right)$, for measures on $\mathbb{R}^{d}$. Define $P_{t}$, the raising to the convolution power $t>0$, in such a way that, for $\mu \in I D, P_{t} \mu$ is an infinitely divisible distribution with characteristic function $\widehat{P_{t} \mu}(z)=\widehat{\mu}(z)^{t}$. The mappings $T_{a}$ (restricted to $I D$ ), $P_{t}$, and $\Phi_{f}$ are commutative with each other. A measure $\mu$ on $\mathbb{R}^{d}$ is called symmetric if $T_{-1} \mu=\mu$. A distribution $\mu$ on $\mathbb{R}^{d}$ is called shifted symmetric if $\mu=\rho * \delta_{\gamma}$ with some symmetric distribution $\rho$ and some $\delta$-distribution $\delta_{\gamma}$. Let $I D_{\text {sym }}=I D_{\text {sym }}\left(\mathbb{R}^{d}\right)$ $\left[\right.$ resp. $I D_{\text {sym }}^{\text {shift }}=I D_{\text {sym }}^{\text {shift }}\left(\mathbb{R}^{d}\right)$ ] denote the class of symmetric [resp. shifted symmetric] infinitely divisible distributions on $\mathbb{R}^{d}$.

Example 4.1. Let $f(s)=b 1_{[0, a]}(s)-b 1_{(a, 2 a]}(s)$ with $a>0$ and $b \neq 0$. Then $\Re_{\infty}\left(\Phi_{f}\right)=I D_{\text {sym }}$.

Indeed, for $\rho \in I D$,

$$
C_{\Phi_{f} \rho}(z)=\int_{0}^{a} C_{\rho}(b z) d s+\int_{a}^{2 a} C_{\rho}(-b z) d s=a C_{\rho}(b z)+a C_{\rho}(-b z)=C_{P_{a} T_{b}\left(\rho * T_{-1} \rho\right)}(z)
$$

for $z \in \mathbb{R}^{d}$, and hence $\Phi_{f} \rho=P_{a} T_{b}\left(\rho * T_{-1} \rho\right)$. Define $U \rho=P_{1 / 2} \rho * T_{-1} P_{1 / 2} \rho$. Then $U \rho \in I D_{\text {sym }}$ for any $\rho \in I D$. If $\rho \in I D_{\text {sym }}$, then $U \rho=\rho$. Hence $U^{n} \rho=U \rho$ for $n=1,2, \ldots$. Since $\Phi_{f}=P_{a} T_{b} P_{2} U=P_{2 a} T_{b} U$, we have $\Phi_{f}^{n}=P_{2 a}^{n} T_{b}^{n} U=U P_{2 a}^{n} T_{b}^{n}$ and $U=\Phi_{f}^{n} P_{1 /(2 a)}^{n} T_{1 / b}^{n}$. Hence $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=\mathfrak{R}(U)=I D_{\text {sym }}$.

Example 4.2. Let $f(s)=b 1_{[0, a]}(s)-b 1_{(a, a+c]}(s)$ with $a>0, c>0, a \neq c$, and $b \neq 0$. Then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=I D_{\text {sym }}^{\text {shift }}$.

To see this, notice that

$$
C_{\Phi_{f} \rho}(z)=a C_{\rho}(b z)+c C_{\rho}(-b z)=(a+c)\left(a_{1} C_{T_{b} \rho}(z)+\left(1-a_{1}\right) C_{T_{b} \rho}(-z)\right)
$$

for $\rho \in I D$, where $a_{1}=a /(a+c)$. That is, $\Phi_{f} \rho=P_{a+c} T_{b}\left(P_{a_{1}} \rho * P_{1-a_{1}} T_{-1} \rho\right)$. Let us define $V \rho=P_{a_{1}} \rho * P_{1-a_{1}} T_{-1} \rho$. Note that $V$ is the stochastic integral mapping $\Phi_{f}$ in the case $a+c=1$ and $b=1$. We have

$$
\begin{equation*}
V^{n} \rho=P_{a_{n}} \rho * P_{1-a_{n}} T_{-1} \rho \tag{4.1}
\end{equation*}
$$

for $n=1,2, \ldots$, where $a_{n}$ is given by $a_{n}=1-a_{1}+a_{n-1}\left(2 a_{1}-1\right)$. Indeed, if (4.1) is true for $n$, then it is true for $n+1$ in place of $n$, since

$$
V^{n+1} \rho=P_{a_{n}} V \rho * P_{1-a_{n}} T_{-1} V \rho=P_{a_{n}} V \rho * P_{1-a_{n}} V T_{-1} \rho
$$

$$
\begin{aligned}
& =P_{a_{n}}\left(P_{a_{1}} \rho * P_{1-a_{1}} T_{-1} \rho\right) * P_{1-a_{n}}\left(P_{a_{1}} T_{-1} \rho * P_{1-a_{1}} \rho\right) \\
& =P_{a_{n} a_{1}+\left(1-a_{n}\right)\left(1-a_{1}\right)} \rho * P_{a_{n}\left(1-a_{1}\right)+\left(1-a_{n}\right) a_{1}} T_{-1} \rho \\
& =P_{a_{n+1}} \rho * P_{1-a_{n+1}} T_{-1} \rho .
\end{aligned}
$$

We see that $0<a_{n}<1$ for all $n$. We have $\Phi_{f}^{n}=P_{a+c}^{n} T_{b}^{n} V^{n}=V^{n} P_{a+c}^{n} T_{b}^{n}$ and $V^{n}=$ $P_{1 /(a+c)}^{n} T_{1 / b}^{n} \Phi_{f}^{n}=\Phi_{f}^{n} P_{1 /(a+c)}^{n} T_{1 / b}^{n}$. Therefore $\mathfrak{R}\left(\Phi_{f}^{n}\right)=\mathfrak{R}\left(V^{n}\right)$ and hence $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=$ $\mathfrak{R}_{\infty}(V)$. Next let us show that

$$
\begin{equation*}
\mathfrak{R}_{\infty}(V)=I D_{\text {sym }}^{\text {shift }} . \tag{4.2}
\end{equation*}
$$

If $\rho \in I D_{\text {sym }}$, then $V \rho=\rho$. Hence $I D_{\text {sym }} \subset \mathfrak{R}_{\infty}(V)$. If $\rho=\delta_{\gamma}$, then $V \rho=$ $\delta_{a_{1} \gamma} * \delta_{-\left(1-a_{1}\right) \gamma}=\delta_{\left(2 a_{1}-1\right) \gamma}$. Now $\delta_{\gamma}=V \delta_{\left(1 /\left(2 a_{1}-1\right)\right) \gamma}$, since $a_{1} \neq 1 / 2$. Hence all $\delta$-distributions are in $\mathfrak{R}\left(V^{n}\right)$ and hence in $\mathfrak{R}_{\infty}(V)$. Since $\mathfrak{R}_{\infty}(V)$ is closed under convolution, we obtain $I D_{\text {sym }}^{\text {shift }} \subset \Re_{\infty}(V)$. To show the converse, assume that $\mu \in$ $\mathfrak{R}_{\infty}(V)$. Then $\mu=V^{n} \rho_{n}$ for some $\rho_{n} \in I D$. It follows from (4.1) that $\nu_{\mu}=a_{n} \nu_{\rho_{n}}+$ $\left(1-a_{n}\right) T_{-1} \nu_{\rho_{n}}$. Let $\sigma_{n} \in I D$ be such that $\left(A_{\sigma_{n}}, \nu_{\sigma_{n}}, \gamma_{\sigma_{n}}\right)=\left(0, \nu_{\rho_{n}}, 0\right)$. It follows from $a_{n}=1-a_{1}+a_{n-1}\left(2 a_{1}-1\right)$ and from $0<a_{n}<1$ that $a_{n} \rightarrow 1 / 2$ as $n \rightarrow \infty$. Hence $a_{n}>1 / 3$ for all large $n$. We see that the set $\left\{\sigma_{n}: n=1,2, \ldots\right\}$ is precompact, since $\nu_{\sigma_{n}} \leqslant a_{n}^{-1} \nu_{\mu} \leqslant 3 \nu_{\mu}$ for all large $n$. Thus we can choose a subsequence $\left\{\sigma_{n_{k}}\right\}$ convergent to some $\mu^{\prime} \in I D$. Since $\int \varphi(x) \nu_{\sigma_{n_{k}}}(d x) \rightarrow \int \varphi(x) \nu_{\mu^{\prime}}(d x)$ for any bounded continuous function $\varphi$ which vanishes on a neighborhood of the origin and since $a_{n} \rightarrow 1 / 2$, we obtain $\nu_{\mu}=(1 / 2) \nu_{\mu^{\prime}}+(1 / 2) T_{-1} \nu_{\mu^{\prime}}$. Hence $\nu_{\mu}$ is symmetric. Hence $\mu * \delta_{-\gamma_{\mu}}$ is symmetric. It follows that $\mu \in I D_{\text {sym }}^{\text {shift }}$. This proves (4.2) and therefore $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=I D_{\text {sym }}^{\text {shift }}$.

Example 4.3. Let $\alpha<0$. Let $h(s)$ be one of $f_{\alpha}(s), \bar{f}_{p, \alpha}(s)$, and $l_{q, \alpha}(s)(p \geqslant 1, q>0)$. Let $s_{0}=\sup \{s: h(s)>0\}$. Then $0<s_{0}<\infty$. Define

$$
f(s)= \begin{cases}h(s), & 0 \leqslant s \leqslant s_{0} \\ -h\left(2 s_{0}-s\right), & s_{0}<s \leqslant 2 s_{0} \\ 0, & s>2 s_{0}\end{cases}
$$

Then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty} \cap I D_{\text {sym }}$.
Proof is as follows. First, recall that $\mathfrak{D}\left(\Phi_{f}\right)=\mathfrak{D}\left(\Phi_{h}\right)=I D$. We have, for $\rho \in I D$,

$$
\begin{aligned}
C_{\Phi_{f} \rho}(z) & =\int_{0}^{s_{0}} C_{\rho}(h(s) z) d s+\int_{s_{0}}^{2 s_{0}} C_{\rho}\left(-h\left(2 s_{0}-s\right) z\right) d s \\
& =\int_{0}^{s_{0}} C_{\rho}(h(s) z) d s+\int_{0}^{s_{0}} C_{\rho}(-h(s) z) d s \\
& =C_{\Phi_{h} \rho}(z)+C_{\Phi_{h} T_{-1} \rho}(z) .
\end{aligned}
$$

It follows that $\Phi_{f} \rho=\Phi_{h}\left(\rho * T_{-1} \rho\right)=\Phi_{h} P_{2} U \rho=U P_{2} \Phi_{h} \rho$, where $U$ is the mapping used in Example 4.1. It follows that $\Phi_{f}^{n}=\Phi_{h}^{n} P_{2}^{n} U=U P_{2}^{n} \Phi_{h}^{n}$ for $n=1,2, \ldots$. Hence $\mathfrak{R}\left(\Phi_{f}^{n}\right) \subset \mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}$. Conversely, assume that $\rho \in \mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}$. Then $\mu=\Phi_{h}^{n} \rho$ for some $\rho$ and $T_{-1} \mu=\Phi_{h}^{n} T_{-1} \rho$. Since $\Phi_{h}$ is one-to-one (see [S]), we have $\rho=T_{-1} \rho$. Hence $\Phi_{f}^{n} \rho=\Phi_{h}^{n} P_{2}^{n} U \rho=\Phi_{h}^{n} P_{2}^{n} \rho=P_{2}^{n} \mu$ and thus $\mu=\Phi_{f}^{n} P_{1 / 2}^{n} \rho \in \mathfrak{R}\left(\Phi_{f}^{n}\right)$. In conclusion, $\mathfrak{R}\left(\Phi_{f}^{n}\right)=\mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}$ and hence $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=\mathfrak{R}_{\infty}\left(\Phi_{h}\right) \cap I D_{\text {sym }}=$ $L_{\infty} \cap I D_{\text {sym }}$.

Example 4.4. Let $h(s)$ and $s_{0}$ be as in Example 4.3. Define

$$
f(s)= \begin{cases}h\left(s_{0}-s\right), & 0 \leqslant s \leqslant s_{0} \\ h\left(s-s_{0}\right), & s_{0}<s \leqslant 2 s_{0} \\ -h\left(3 s_{0}-s\right), & 2 s_{0}<s \leqslant 3 s_{0} \\ 0, & s>3 s_{0}\end{cases}
$$

Then $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty} \cap I D_{\text {sym }}^{\text {shift }}$.
To see this, notice that

$$
\begin{aligned}
C_{\Phi_{f} \rho}(z)= & \int_{0}^{s_{0}} C_{\rho}\left(h\left(s_{0}-s\right) z\right) d s+\int_{s_{0}}^{2 s_{0}} C_{\rho}\left(h\left(s-s_{0}\right) z\right) d s \\
& +\int_{2 s_{0}}^{3 s_{0}} C_{\rho}\left(-h\left(3 s_{0}-s\right) z\right) d s \\
= & \int_{0}^{s_{0}} C_{\rho}(h(s) z) d s+\int_{0}^{s_{0}} C_{\rho}(h(s) z) d s+\int_{0}^{s_{0}} C_{\rho}(-h(s) z) d s \\
= & 2 C_{\Phi_{h} \rho}(z)+C_{\Phi_{h} \rho}(-z) \\
= & 3\left(\frac{2}{3} C_{\Phi_{h} \rho}(z)+\frac{1}{3} C_{\Phi_{h} \rho}(-z)\right) .
\end{aligned}
$$

Hence $\Phi_{f} \rho=P_{3} V \Phi_{h} \rho$, where $V \rho=P_{2 / 3} \rho * P_{1 / 3} T_{-1} \rho$. This mapping $V$ is a special case of $V$ in Example 4.2 with $a_{1}=2 / 3$. Hence (4.1) holds with $a_{n}=2^{-1}\left(1+3^{-n}\right)$ and $1-a_{n}=2^{-1}\left(1-3^{-n}\right)$. Now $\Phi_{f}^{n}=P_{3}^{n} V^{n} \Phi_{h}^{n}=\Phi_{h}^{n} P_{3}^{n} V^{n}=V^{n} P_{3}^{n} \Phi_{h}^{n}$. Hence $\mathfrak{R}\left(\Phi_{f}^{n}\right) \subset \mathfrak{R}\left(\Phi_{h}^{n}\right) \cap \mathfrak{R}\left(V^{n}\right)$. It follows that $\mathfrak{R}_{\infty}\left(\Phi_{f}\right) \subset \mathfrak{R}_{\infty}\left(\Phi_{h}\right) \cap \mathfrak{R}_{\infty}(V)=L_{\infty} \cap I D_{\text {sym }}^{\text {shift }}$ from Theorem 1.1 and (4.2). Let us also show the converse inclusion $L_{\infty} \cap I D_{\text {sym }}^{\text {shift }} \subset$ $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$. It is enough to show

$$
\begin{equation*}
\mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}^{\text {shift }} \subset \mathfrak{R}\left(\Phi_{f}^{n}\right) \tag{4.3}
\end{equation*}
$$

For any $\gamma \in \mathbb{R}^{d}$ we have

$$
C_{\Phi_{h} \delta_{\gamma}}(z)=\int_{0}^{s_{0}} C_{\delta_{\gamma}}(h(s) z) d s=i \int_{0}^{s_{0}}\langle\gamma, h(s) z\rangle d s=i c\langle\gamma, z\rangle=C_{\delta_{c \gamma}}(z),
$$

where $c=\int_{0}^{s_{0}} h(s) d s>0$. That is, $\Phi_{h} \delta_{\gamma}=\delta_{c \gamma}$. Hence $\Phi_{f} \delta_{\gamma}=P_{3} \Phi_{h} V \delta_{\gamma}=$ $P_{3} \Phi_{h}\left(\delta_{(2 / 3) \gamma} * \delta_{-(1 / 3) \gamma}\right)=\Phi_{h} \delta_{\gamma}=\delta_{c \gamma}$. Hence $\Phi_{f}^{n} \delta_{\gamma}=\delta_{c^{n} \gamma}$ and $\delta_{\gamma}=\Phi_{f}^{n} \delta_{c^{-n} \gamma}$. Hence
all $\delta$-distributions are in $\mathfrak{R}\left(\Phi_{f}^{n}\right)$. Similarly all $\delta$-distributions are in $\mathfrak{R}\left(\Phi_{h}^{n}\right)$. Let $\mu \in \mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}^{\text {shift }}$. Then $\mu * \delta_{\gamma} \in \mathfrak{R}\left(\Phi_{h}^{n}\right) \cap I D_{\text {sym }}$ for some $\gamma$. Letting $\mu^{\prime}=\mu * \delta_{\gamma}$, we have $\mu^{\prime}=\Phi_{h}^{n} \rho^{\prime}$ for some $\rho^{\prime}$. Since $\mu^{\prime}=T_{-1} \mu^{\prime}=\Phi_{h}^{n} T_{-1} \rho^{\prime}$, we have $\rho^{\prime}=T_{-1} \rho^{\prime}$ from the one-to-one property of $\Phi_{h}$. Thus $V^{n} \rho^{\prime}=\rho^{\prime}$ and $\Phi_{f}^{n} \rho^{\prime}=\Phi_{h}^{n} P_{3}^{n} \rho^{\prime}=P_{s}^{n} \mu^{\prime}$. Hence $\mu^{\prime}=P_{1 / 3}^{n} \Phi_{f}^{n} \rho^{\prime}=\Phi_{f}^{n} P_{1 / 3}^{n} \rho^{\prime} \in \mathfrak{R}\left(\Phi_{f}^{n}\right)$. It follows that $\mu=\mu^{\prime} * \delta_{-\gamma} \in \mathfrak{R}\left(\Phi_{f}^{n}\right)$. This proves (4.3). Hence $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=L_{\infty} \cap I D_{\text {sym }}^{\text {shift }}$.

Example 4.5. Let $b>1$. Let $f(s)=b 1_{[0,1]}(s)+1_{(1,2]}(s)$. Then $L_{\infty}(b) \subset \mathfrak{R}_{\infty}\left(\Phi_{f}\right) \varsubsetneqq$ $I D$. We do not know whether $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$ equals $L_{\infty}(b)$. Here $L_{\infty}(b)$ is the $b$-semianalogue of the class $L_{\infty}$, mentioned in Section 2.

Let us show that $L_{\infty}(b) \subset \mathfrak{R}_{\infty}\left(\Phi_{f}\right)$. For $0<\alpha \leqslant 2$ define $\mathfrak{S}_{\alpha}(b)=\mathfrak{S}_{\alpha}\left(b, \mathbb{R}^{d}\right)$ as follows: $\rho \in \mathfrak{S}_{\alpha}(b)$ if and only if $\rho$ is a $\delta$-distribution or a non-trivial $\alpha$-semi-stable distribution with $b$ as a span, that is,

$$
\mathfrak{S}_{\alpha}(b)=\left\{\rho \in I D: P_{b^{\alpha}} \rho=T_{b} \rho * \delta_{\gamma} \text { for some } \gamma \in \mathbb{R}^{d}\right\} .
$$

We have $C_{\Phi_{f} \rho}(z)=C_{\rho}(b z)+C_{\rho}(z)$ for $\rho \in I D$, that is, $\Phi_{f} \rho=T_{b} \rho * \rho$. If $\rho \in \mathfrak{S}_{\alpha}(b)$ with $P_{b^{\alpha}} \rho=T_{b} \rho * \delta_{\gamma}$, then $\mu=\Phi_{f} \rho$ satisfies $\mu=T_{b} \rho * \rho=P_{b^{\alpha}} \rho * \delta_{-\gamma} * \rho=$ $P_{b^{\alpha}+1} \rho * \delta_{-\gamma}$ and $\mu \in \mathfrak{S}_{\alpha}(b)$. If $\mu \in \mathfrak{S}_{\alpha}(b)$ with $P_{b^{\alpha}} \mu=T_{b} \mu * \delta_{\gamma^{\prime}}$, then $\mu=\Phi_{f} \rho$ for $\rho=P_{1 /\left(b^{\alpha}+1\right)}\left(\mu * \delta_{(1 /(b+1)) \gamma^{\prime}}\right) \in \mathfrak{S}_{\alpha}(b)$. Therefore $\Phi_{f}\left(\mathfrak{S}_{\alpha}(b)\right)=\mathfrak{S}_{\alpha}(b)$. Hence $\mathfrak{S}_{\alpha}(b) \subset$ $\mathfrak{R}\left(\Phi_{f}^{n}\right)$ for $0<\alpha \leqslant 2$ and $n=1,2, \ldots$ It follows from Proposition 3.2 of Maejima and Sato (2009) that $\mathfrak{R}\left(\Phi_{f}^{n}\right)$ is closed under convolution and weak convergence. Hence $L_{\infty}(b) \subset \mathfrak{R}\left(\Phi_{f}^{n}\right)$ and thus $L_{\infty}(b) \subset \Re_{\infty}\left(\Phi_{f}\right)$. In order to show $\mathfrak{R}_{\infty}\left(\Phi_{f}\right) \varsubsetneqq I D$, let $\mu$ be such that $\nu_{\mu}=\delta_{a}$ with $a \neq 0$. Suppose that $\mu=\Phi_{f} \rho$ for some $\rho \in I D$. Then $\nu_{\mu}=T_{b} \nu_{\rho}+\nu_{\rho}$. If $\nu_{\rho} \neq 0$, then the support of $\nu_{\rho}$ contains at least one point $a^{\prime} \neq 0$ and hence the support of $\nu_{\mu}$ contains at least two points $\left\{a^{\prime}, b a^{\prime}\right\}$, which is absurd. If $\nu_{\rho}=0$, then $\nu_{\mu}=0$, which is also absurd. Therefore $\mu \notin \mathfrak{R}\left(\Phi_{f}\right)$ and hence $\mu \notin \Re_{\infty}\left(\Phi_{f}\right)$.

## 5. Concluding remarks

The limit class $\Re_{\infty}\left(\Phi_{f}\right)$ is not known in many cases. For instance it is not known for the following choices of $f(s): l_{q, 1}(s)$ with $q \in(0,1) \cup(1, \infty)$ in $[\mathrm{S}] ; \bar{f}_{p, \alpha}(s)$ with $p \in(0,1)$ and $\alpha \in(-\infty, 2)$ in $[\mathrm{S}] ; \cos \left(2^{-1} \pi s\right)$ in Maejima et al. (2010b); $e^{-s} 1_{[0, c]}(s)$ with $c \in(0, \infty)$ in Pedersen and Sato (2005); $G_{\alpha, \beta}^{*}(s)$ with $\alpha \in[1,2)$ and $\beta>0$ satisfying $\alpha=1+n \beta$ for some $n=0,1, \ldots$ in Maejima and Ueda (2010b). Another instance is $\Phi_{f}=\Upsilon^{\alpha}$ with $\alpha \in(0,1)$ related to the Mittag-Leffler function, introduced in Barndorff-Nielsen and Thorbjørnsen (2006).

Consider, as in Sato (2007), a stochastic integral mapping

$$
\Phi_{f} \rho=\mathcal{L}\left(\int_{0+}^{a} f(s) d X_{s}^{(\rho)}\right)
$$

with $0<a<\infty$ for a function $f(s)$ locally square-integrable on the interval $(0, a]$ and study $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)=\bigcap_{n=1}^{\infty} \mathfrak{R}\left(\Phi_{f}^{n}\right)$. Under appropriate choices of $f$ we obtain $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$ equal to $L_{\infty}^{(0, \alpha)} \cap I D_{0}$ with $\alpha \in(1,2), L_{\infty}^{(0, \alpha)} \cap I D_{0} \cap\{\mu \in I D: \mu$ has drift 0$\}$ with $\alpha \in(0,1)$, or a certain subclass of $L_{\infty}^{(0,1)} \cap I D_{0}$. This will be shown in a forthcoming paper.

It is an interesting problem what other classes can appear as $\mathfrak{R}_{\infty}\left(\Phi_{f}\right)$.

## References

[1] T. Aoyama, A. Lindner and M. Maejima. A new family of mappings of infinitely divisible distributions related to the Goldie-Steutel-Bondesson class. Elect. J. Probab. 15, 1119-1142 (2010).
[2] T. Aoyama, M. Maejima and J. Rosiński. A subclass of type $G$ selfdecomposable distributions. J. Theoret. Probab. 21, 14-34 (2008).
[3] O. Arizmendi, O. E. Barndorff-Nielsen and V. Pérez-Abreu. On free and classical type $G$ laws. Braz. J. Prob. Stat. (2010). To appear.
[4] O. E. Barndorff-Nielsen, M. Maejima and K. Sato. Some classes of infinitely divisible distributions admitting stochastic integral representations. Bernoulli 12, 1-33 (2006).
[5] O. E. Barndorff-Nielsen and S. Thorbjørnsen. Regularising mappings of Lévy measures. Stoch. Proc. Appl. 116, 423-446 (2006).
[6] B. Grigelionis. Extended Thorin classes and stochastic integrals. Liet. Mat. Rink. 47, 497-503 (2007).
[7] K. Ichifuji, M. Maejima and Y. Ueda. Fixed points of mappings of infinitely divisible distributions on $\mathbb{R}^{d}$. Statist. Probab. Lett. 80, 1320-1328 (2010).
[8] Z. J. Jurek. The class $L_{m}(Q)$ of probability measures on Banach spaces. Bull. Polish Acad. Sci. Math. 31, 51-62 (1983).
[9] Z. J. Jurek. Relations between the $s$-selfdecomposable and selfdecomposable measures. Ann. Probab. 13, 592-608 (1985).
[10] Z. J. Jurek. Random integral representations for classes of limit distributions similar to Lévy class $L_{0}$. Probab. Theory Relat. Fields 78, 473-490 (1988).
[11] Z. J. Jurek. Random integral representations for classes of limit distributions similar to Lévy class $L_{0}$, II. Nagoya Math. J. 114, 53-64 (1989).
[12] Z. J. Jurek. The random integral representation hypothesis revisited: new classes of $s$ selfdecomposable laws. Abstract and Applied Analysis, Proc. Intern. Conf., Hanoi, 2002, World Scientific, pp. 479-498 (2004).
[13] M. Maejima, M. Matsui and M. Suzuki. Classes of infinitely divisible distributions on $\mathbb{R}^{d}$ related to the class of selfdecomposable distributions. Tokyo J. Math. (2010a). To appear.
[14] M. Maejima and G. Nakahara. A note on new classes of infinitely divisible distributions on $\mathbb{R}^{d}$. Elect. Comm. Probab. 14, 358-371 (2009).
[15] M. Maejima, V. Pérez-Abreu and K. Sato. Type $A$ distributions: Infinitely divisible distributions related to arcsine density (2010b). Preprint.
[16] M. Maejima, V. Pérez-Abreu and K. Sato. Four-parameter fractional integral transformations producing Lévy measures of infinitely divisible distributions (2010c). Preprint.
[17] M. Maejima and K. Sato. The limits of nested subclasses of several classes of infinitely divisible distributions are identical with the closure of the class of stable distributions. Probab. Theory Relat. Fields 145, 119-142 (2009).
[18] M. Maejima, K. Sato and T. Watanabe. Completely operator semi-selfdecomposable distributions. Tokyo J. Math. 23, 235-253 (2000).
[19] M. Maejima and Y. Ueda. Stochastic integral characterizations of semi-selfdecomposable distributions and related Ornstein-Uhlenbeck type processes. Comm. Stoch. Anal. 3, 349-367 (2009a).
[20] M. Maejima and Y. Ueda. The relation between $\lim _{m \rightarrow \infty} \mathfrak{R}\left(\Psi_{\alpha}^{m+1}\right)$ and $\lim _{m \rightarrow \infty} \mathfrak{R}\left(\Phi_{\alpha}^{m+1}\right)$ (2009b). Private Communication.
[21] M. Maejima and Y. Ueda. Nested subclasses of the class of $\alpha$-selfdecomposable distributions. Tokyo J. Math. (2010a). To appear.
[22] M. Maejima and Y. Ueda. Compositions of mappings of infinitely divisible distributions with applications to finding the limits of some nested subclasses. Elect. Comm. Probab. 15, 227-239 (2010b).
[23] M. Maejima and Y. Ueda. $\alpha$-selfdecomposable distributions and related Ornstein-Uhlenbeck type processes (2010c). Preprint.
[24] J. Pedersen and K. Sato. The class of distributions of periodic Ornstein-Uhlenbeck proccesses driven by Lévy processes. J. Theor. Probab. 18, 209-235 (2005).
[25] A. Rocha-Arteaga and K. Sato. Topics in Infinitely Divisible Distributions and Lévy Processes. Aportaciones Matemáticas, Investigación 17, Sociedad Matemática Mexicana, México (2003).
[26] K. Sato. Class $L$ of multivariate distributions and its subclasses. J. Multivar. Anal. 10, 207-232 (1980).
[27] K. Sato. Lévy Processes and Infinitely Divisible Distributions, Cambridge University Press, Cambridge (1999).
[28] K. Sato. Two families of improper stochastic integrals with respect to Lévy processes. ALEA Lat. Am. J. Probab. Math. Statist. 1, 47-87 (2006).
[29] K. Sato. Transformations of infinitely divisible distributions via improper stochastic integrals. ALEA Lat. Am. J. Probab. Math. Statist. 3, 67-110 (2007).
[30] K. Sato. Fractional integrals and extensions of selfdecomposability. Lecture Notes in Mathematics (Springer), Lévy Matters I, 1-91 (2010). (http://ksato.jp/)
[31] K. Urbanik. Limit laws for sequences of normed sums satisfying some stability conditions. In: Multivariate Analysis-III (ed. P. R. Krishnaiah), Academic Press, pp. 225-237 (1973).
[32] S. J. Wolfe. On a continuous analogue of the stochastic difference equation $X_{n}=\rho X_{n-1}+B_{n}$, Stoch. Proc. Appl. 12, 301-312 (1982).

