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Differential reflection phase shift under conditions of attenuated internal reflection

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The angle-of-incidence dependence of the differential reflection phase shift Δ between p and s polarizations is considered a function of the real and imaginary parts of the relative complex dielectric function ε of an interface in the domain of fractional optical constants, i.e., under conditions of internal reflection. The constraint on complex ε such that oscillatory and monotonic angular responses are obtained is determined. A sensitive and stable technique, which is based on attenuated internal reflection ellipsometry between the Brewster angle and the critical angle, is proposed for measuring small induced absorption ($\varepsilon_i \sim 10^{-5}$) in the medium of refraction. © 1999 Optical Society of America [S0740-3232(99)01407-6]

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1. INTRODUCTION

When a plane wave of monochromatic light is incident at the planar interface between two homogeneous, optically isotropic, nonmagnetic media from the side of low index of refraction (i.e., in external reflection), the differential reflection phase shift (or ellipsometric angle) Δ decreases monotonically from π to 0 as the angle of incidence ϕ is increased from 0 (normal incidence) to 90° (grazing incidence), respectively.^{1,2} This assumes the presence of absorption in the medium of refraction, the $\exp(j\omega t)$ dependence for time-harmonic wave fields, and the rest of the Nebraska–Muller conventions.³

Under conditions of attenuated internal reflection (AIR), the Δ -versus- ϕ curve can be oscillatory, exhibiting one minimum and one maximum between normal and grazing incidence.^{4,5} In this paper we determine the constraint on the relative complex dielectric function

$$\varepsilon = \varepsilon_1/\varepsilon_0 = \varepsilon_r - j\varepsilon_i, \quad (1)$$

such that either monotonic or oscillatory phase-versus-angle response is obtained in AIR (i.e., in the domain of fractional optical constants). In Eq. (1) ε_0 and ε_1 are the dielectric constants of the media of incidence and refraction, respectively, at a given wavelength.

We also propose the use of AIR ellipsometry between the Brewster angle and the critical angle as a sensitive technique for measuring small absorption in the medium of refraction, as represented by small values of ε_i . Measurement of the refractive index on the basis of total internal reflection has received renewed interest recently.^{6,7}

2. MONOTONIC AND OSCILLATORY PHASE-VERSUS-ANGLE RESPONSES

The ratio $\rho = r_p/r_s$ of the complex-amplitude reflection coefficients for p and s linear polarizations, parallel and perpendicular to the plane of incidence, respectively, is given by¹

$$\rho = \frac{[\sin \phi \tan \phi - (\varepsilon - \sin^2 \phi)^{1/2}]}{[\sin \phi \tan \phi + (\varepsilon - \sin^2 \phi)^{1/2}]}, \quad (2)$$

and Δ is obtained from ρ by

$$\Delta = \arg(\rho). \quad (3)$$

For a given value of ε_r in the range $0 \leq \varepsilon_r \leq 1$, the Δ -versus- ϕ curve depends on ε_i . As an example, Fig. 1 shows a family of such curves for $\varepsilon_r = 0.6$. The response is oscillatory for $\varepsilon_i = 0, 0.005, 0.015, 0.045, 0.075$, and 0.10 (curves marked o, a, b, c, d, and e, respectively) and is monotonic for $\varepsilon_i = 0.1366, 0.2$, and 2.0 (curves f, g, and h, respectively). The limiting value $\varepsilon_i = 0.1366$ (curve f) is determined by numerical iteration as the upper bound on ε_i , below which an oscillatory response would occur. Figure 2 gives a clear three-dimensional view of the tran-

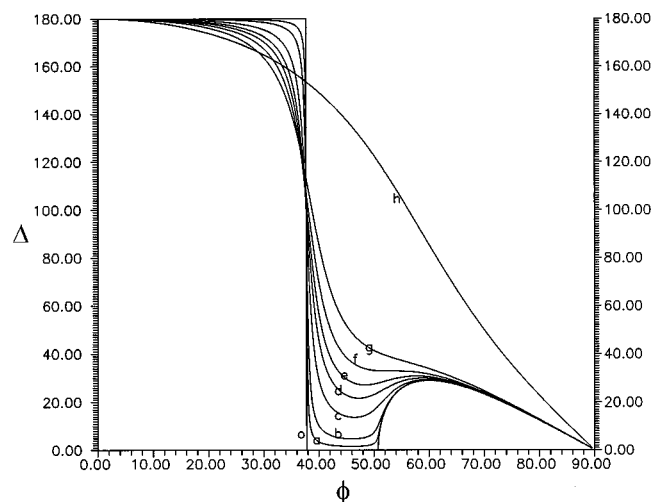


Fig. 1. Differential reflection phase shift Δ versus angle of incidence ϕ for different values of the imaginary part of the relative dielectric constant ε_i and the constant real part $\varepsilon_r = 0.6$. The curves marked o, a, b, c, d, e, f, g, and h correspond to $\varepsilon_i = 0, 0.005, 0.015, 0.045, 0.075, 0.1, 0.1366, 0.2$, and 2.0 , respectively.

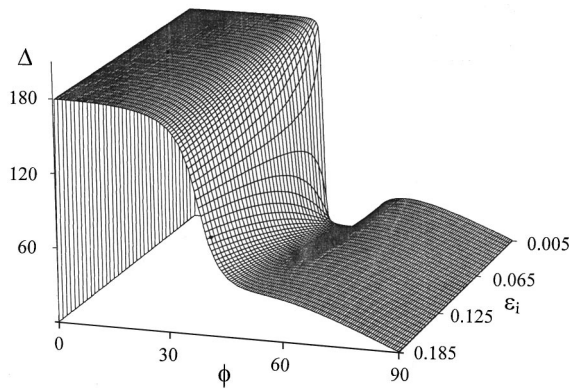


Fig. 2. Differential reflection phase shift Δ as a function of the angle of incidence ϕ and the imaginary part of the relative dielectric constant ϵ_i for a given value of the real part $\epsilon_r = 0.6$.

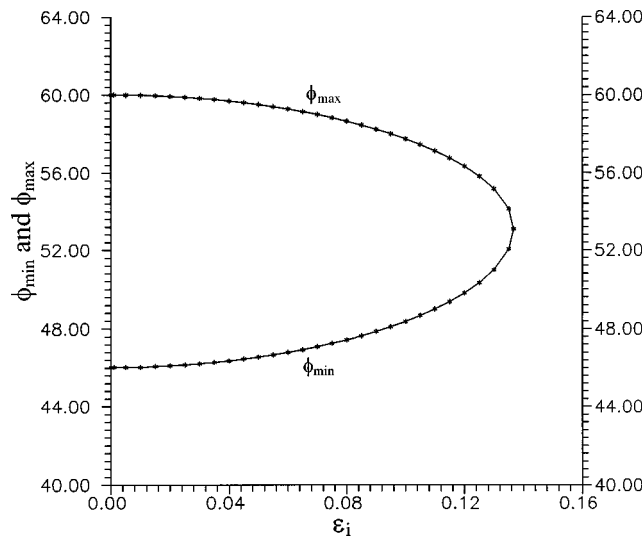


Fig. 3. Angles of incidence ϕ_{\min} and ϕ_{\max} for minimum and maximum differential reflection phase shift Δ as functions of the imaginary part of the relative dielectric constant ϵ_i when $\epsilon_r = 0.6$, corresponding to the data shown in Fig. 1.

3. ATTENUATED INTERNAL REFLECTION ELLIPSOMETRY BETWEEN THE BREWSTER ANGLE AND THE CRITICAL ANGLE

For internal reflection at a sharp interface between two transparent media, $\Delta = 0$ between the Brewster angle and the critical angle, as shown in Fig. 1 for $\epsilon_r = 0.6$ and $\epsilon_i = 0$. If absorption is introduced into the medium of refraction (e.g., by dissolving an absorbing species into an aqueous solution), ϵ_i is no longer 0 and $\Delta \neq 0$.

From Figs. 1 and 2 it is apparent that measurement of Δ by internal reflection ellipsometry between the Brewster angle and the critical angle provides a sensitive and stable technique for measuring small values of ϵ_i . Although the sensitivity of Δ with respect to ϵ_i is very high at and in the immediate neighborhood of the Brewster and critical angles, the associated measurement is highly susceptible to small angle-of-incidence errors. Therefore ellipsometry at these special angles is not recommended.

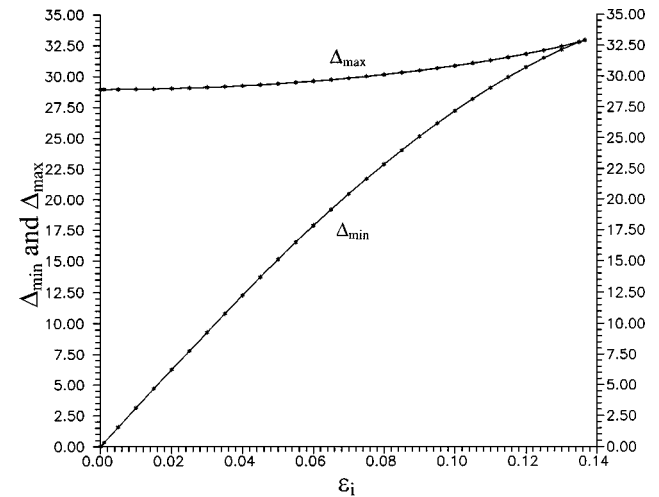


Fig. 4. Minimum and maximum differential reflection phase shift Δ_{\min} and Δ_{\max} as functions of the imaginary part of the relative dielectric constant ϵ_i when $\epsilon_r = 0.6$, corresponding to the data shown in Fig. 1.

sition from oscillatory to monotonic Δ -versus- ϕ response as ϵ_i is increased, for a constant value of $\epsilon_r = 0.6$.

In Fig. 1 the range of ϕ for which $\Delta=0$, when $\epsilon_r = 0.6$ and $\epsilon_i = 0$ (curve o for a transparent interface), extends from the Brewster angle $\phi_B = \arctan(0.6^{1/2}) = 37.76^\circ$ to the critical angle $\phi_c = \arcsin(0.6^{1/2}) = 50.77^\circ$. In Fig. 3, for each value of ϵ_i in the range $0 < \epsilon_i < 0.1366$, the angles of incidence ϕ_{\min} and ϕ_{\max} of minimum and maximum differential reflection phase shift Δ are calculated and are plotted as functions of ϵ_i . The two angles merge, $\phi_{\min} = \phi_{\max}$, at $\epsilon_i = 0.1366$. In Fig. 4 the corresponding Δ_{\min} and Δ_{\max} are shown as functions of ϵ_i .

By varying ϵ_r and by repeating the calculation for the upper bound on ϵ_i for each value of ϵ_r , we obtain the domain of fractional optical constants (ϵ_r, ϵ_i) for which the Δ -versus- ϕ response is oscillatory. The results are represented in Fig. 5 by the region of the complex ϵ plane between the curve and the ϵ_r axis. A monotonic response is obtained for all (ϵ_r, ϵ_i) pairs that correspond to points outside this region.

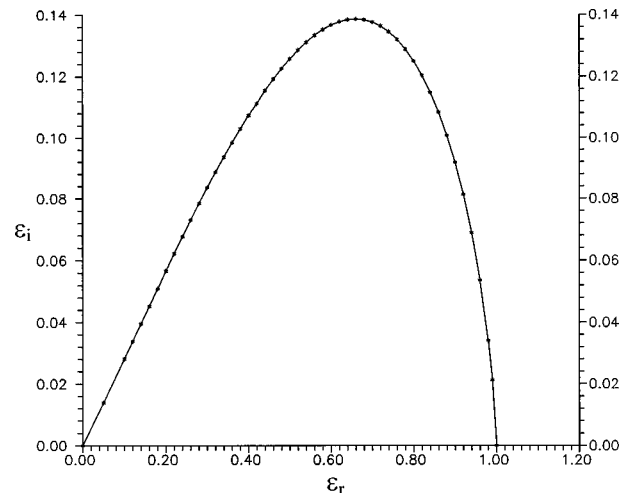


Fig. 5. Upper bound on the imaginary part of the relative dielectric constant ϵ_i for an oscillatory Δ -versus- ϕ response as a function of the real part ϵ_r .

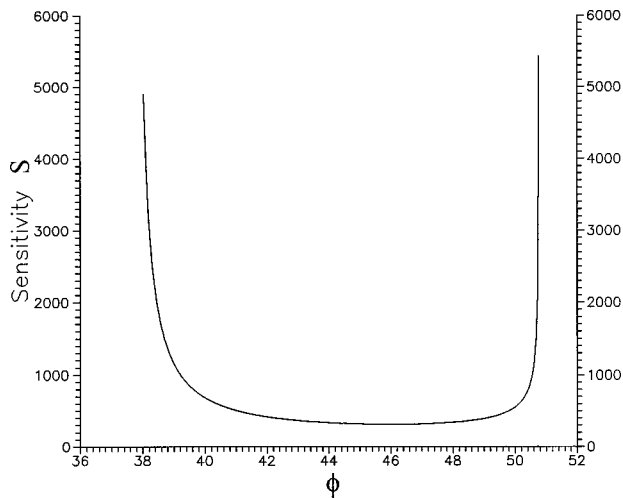


Fig. 6. Ellipsometric sensitivity S , Eq. (5), as a function of the angle of incidence ϕ for attenuated internal reflection between the Brewster angle and the critical angle of a transparent interface with $\epsilon_r = 0.6$.

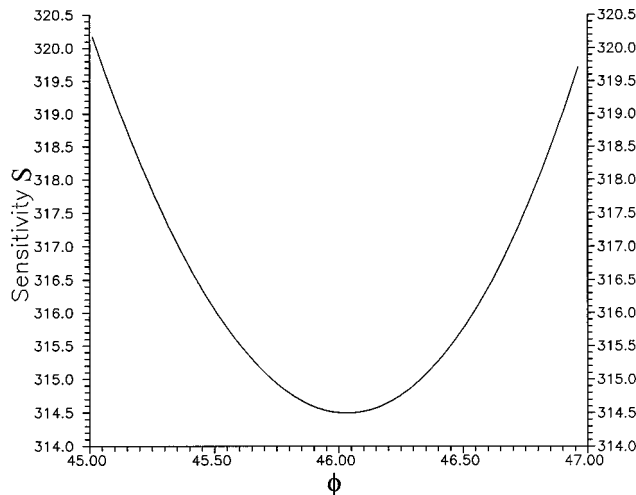


Fig. 7. Same as in Fig. 6, but for a range of $\pm 1^\circ$ of the incidence angle ϕ around the minimum.

It is instructive to define a sensitivity function as

$$S = [\partial\Delta/\partial\epsilon_i]_{\epsilon_i=0}. \quad (4)$$

From Eqs. (2) and (3), one obtains an explicit expression for S :

$$S = (180^\circ/\pi) \sin \phi \tan \phi / [(\epsilon_r - \sin^2 \phi)^{1/2} \times (\tan^2 \phi - \epsilon_r)]. \quad (5)$$

Equation (5) predicts singularities of the sensitivity function at the Brewster and critical angles. Figure 6 shows the variation of S with ϕ between these two angles for $\epsilon_r = 0.6$. The sensitivity is minimum at $\phi = 46.036^\circ$, and Fig. 7 shows $S(\phi)$ within $\pm 1^\circ$ of this angle. The

minimum sensitivity of 314.49 (degrees per unit change of ϵ_i) is quite high. Operation at or near the minimum of S is insensitive to small angle-of-incidence (or beam convergence/divergence) errors. Because changes of Δ of 0.03° are readily detectable in ellipsometry,¹ changes of ϵ_i from 0 to 10^{-5} should be measurable. Notice, however, that the proposed model assumes that the change of Δ is due to bulk absorption in the medium of refraction and not to the formation of a third thin-film phase at the boundary between the two media. Determination of complex refractive indices by attenuated total reflectance (not phase) measurements were previously reported.^{8,9}

4. CONCLUSION

In this paper we have found the conditions under which the differential phase shift on reflection at an interface between two media can become an oscillatory function of the angle of incidence. We have also proposed the use of attenuated internal reflection ellipsometry between the Brewster angle and the critical angle as a sensitive and stable technique for measuring the small imaginary part of the complex dielectric function of an absorbing medium of refraction.

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