

Distance-Engineered Plasmon-Enhanced Light Harvesting in CdSe Quantum Dots

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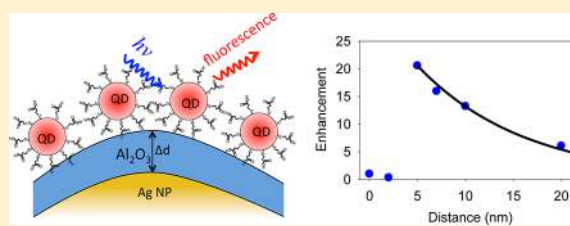
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S Supporting Information

ABSTRACT: Improvement of light harvesting in semiconductor quantum dots (QDs) is essential for the development of efficient QD-based solar energy conversion systems. In this study, plasmon-enhanced light absorption in CdSe QDs sensitized on silver (Ag) nanoparticle (NP) films was examined as a function of interparticle (QD to Ag NP) distance. Up to 24-fold plasmonic enhancement of fluorescence from QDs was observed when the particle separation distance was ≥ 5 nm. The enhancement effect was observed to largely sustain the exciton lifetimes in QDs and to strongly depend on the incident photon wavelength following the plasmon resonant strength of Ag NPs, confirming that the enhanced photoluminescence was mainly due to the enhancement in photoabsorption in CdSe QDs by the plasmon of Ag NPs. This study suggests applications of Ag NPs in QD-based solar energy conversion for significantly improving light harvesting in QDs.

SECTION: Plasmonics, Optical Materials, and Hard Matter



Semiconductor quantum dots (QDs) have received extensive attention because of their attractive chemical and photophysical properties, such as ease of synthesis, tunable band gap, broad absorption spectra, large extinction coefficient, and long photogenerated exciton lifetimes.^{1–3} These properties make QDs a promising class of materials, potentially superior to inorganic or organic molecules, to be utilized as light-harvesting and charge-separation components in solar energy conversion schemes.^{4,5} For example, effective and robust photogeneration of H₂ from water facilitated by semiconductor QDs (or hybrid nanostructures) with catalysts (e.g., platinum, nickel) has been previously reported in the literature.^{6–13} The fundamental processes in these photocatalytic reactions normally involve light harvesting by a QD, followed by charge transfer to the catalyst, where the eventual turnover of reactants to fuels occurs. Because most photocatalytic reactions require multiple charges, effective charge accumulation is essential to realize high efficiency. This requires not only schemes with effective charge separation and suppressed charge recombination but also the development of effective approaches for the enhancement of light harvesting by QDs. Utilizing the localized surface plasmon of metal nanoparticles (NPs) or structures to increase the light harvesting in QDs can be a promising strategy to improve the efficiency in a QD-based solar energy conversion system. The localized surface plasmon resonance (LSPR) of metal particles can strongly intensify the optical field in the vicinity of the metal surfaces and thus can improve the absorption of incident photons in the light-harvesting materials

nearby.^{14–17} This plasmonic effect has been exploited to improve efficiencies in solar cells.^{18–23}

Although plasmon-enhanced photoluminescence in QDs by adjacent gold (Au) or silver (Ag) NPs or nanostructures has been broadly reported,^{24–41} studies focusing on plasmon-enhanced light harvesting in QDs, particularly considering their applications in solar energy conversion, have rarely been described. For most previously reported examples, plasmon enhancements of QD photoluminescence were accompanied by striking decreases in exciton lifetime.^{28–38} Thus, metal NP/QD coupling may reduce the exciton lifetime by increasing the rate of radiative relaxation and/or by opening up energy transfer or other nonradiative relaxation pathways.^{26,31,33,34,39,42–45} Considering the application of QDs in solar energy conversion, a reduction of exciton lifetime may diminish the benefits of plasmon-enhanced absorption. Briefly, the shorter lifetimes could decrease yields for desired charge transfer or charge separation. They also could reduce the exciton diffusion length. Ideally, a plasmonic enhancement that improves photon absorption by QDs without significantly attenuating exciton lifetimes is preferred for applications in solar energy conversion.

Herein, we report the distance-engineered plasmon-enhanced photon absorption in CdSe QDs by using Ag NPs

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but without significantly shortening the QDs' exciton lifetimes. The sample configuration is shown in Figure 1. The distance

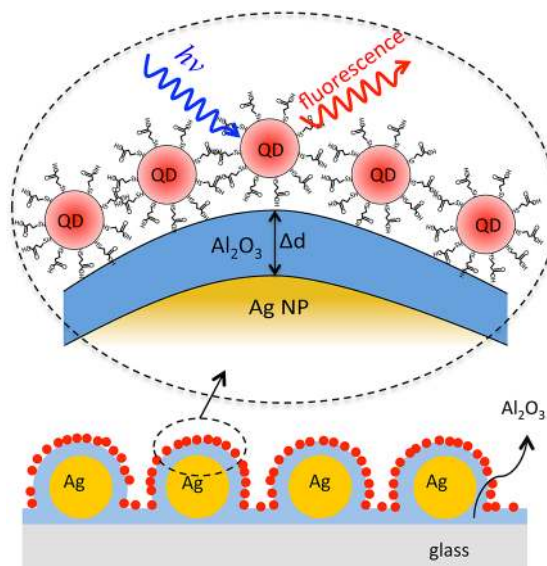


Figure 1. A schematic diagram of an Al_2O_3 -coated Ag NP film with CdSe QDs sensitized on the surface. The distance between QDs and the Ag NP is tuned by changing the thickness (Δd) of the Al_2O_3 layer. Upon illumination, the fluorescence and fluorescent lifetime of QDs as a function of Al_2O_3 thickness are recorded to examine the distance-dependent plasmonic effect of Ag NPs.

from Ag NPs to CdSe QDs is precisely tuned by changing the thickness of the Al_2O_3 layer grown on top of Ag NPs by atomic layer deposition (ALD). Importantly, the ALD approach produces a pinhole-free spacer layer with subnanometer control over the spacer thickness,⁴⁶ which enables a precise study of plasmon-enhanced absorption by QDs as a function of charge-transfer distance and dielectric. By avoiding inhomogeneity in the spacer thickness, it has been reported that it may be possible to further increase the excitation of QDs through plasmon enhancement.⁴⁰ The plasmonic effect of Ag NPs was examined by detecting the fluorescence from QDs. Both fluorescence quenching and enhancement were observed as a function of distance (Al_2O_3 thickness). At a short distance of 2 nm, the fluorescence of the QDs was quenched by $\sim 70\%$ with a significantly shortened exciton lifetime, effects attributable to ultrafast energy transfer from QDs to Ag NPs. In contrast, a 5 nm Ag NP-to-QD distance produced dramatically plasmon-enhanced fluorescence with only slightly shortened exciton lifetimes. Beyond 5 nm, the enhancement decays exponentially with increasing separation distance. The plasmon-enhanced fluorescence of QDs can be mainly attributed to the enhanced absorption of incident photons rather than an enhanced radiative versus nonradiative time constant (photoluminescent quantum yield). This study suggests exploiting Ag NPs in QD-based solar energy conversion for significantly improving light harvesting in QDs without changing their exciton lifetimes.

The synthesis of Ag NPs with an average diameter of 40 ± 6 nm is based on a previously reported method.⁴⁷ Details regarding the synthesis and the fabrication of Al_2O_3 -coated Ag NP films are in the Supporting Information (SI). The Al_2O_3 layer was grown on the top of Ag NPs by ALD. The thickness of Al_2O_3 was controlled by the number of ALD cycles. The deposited Al_2O_3 layers with 2, 5, 7, 10, and 20 nm thicknesses

were grown using 18, 45, 64, 91, and 182 cycles, respectively. A representative AFM image of the Ag NP film with 5 nm Al_2O_3 is shown in the inset in Figure 2, indicating a uniform

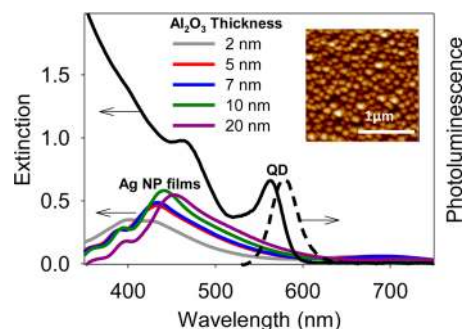


Figure 2. (a) UV-vis extinction spectra of a series of Ag NP films coated with various thicknesses (2, 5, 7, 10, and 20 nm) of Al_2O_3 . Also plotted are the extinction (black solid line) and emission (black dashed line) spectra of the CdSe QDs in water. Shown in the inset is the AFM image of a Ag NP film coated with 5 nm Al_2O_3 , indicating a uniform distribution of Ag NPs on the surface.

distribution of Ag NPs on the substrate (~ 50 NPs per μm^2). On the basis of a simple calculation, these Ag NPs (with 5 nm Al_2O_3 coating) are estimated to cause a $\sim 24\%$ increase in the surface area relative to a bare glass coverslip. An Al_2O_3 -coated (by ALD) glass coverslip (without Ag NPs) was used as a control substrate (designated $\text{Al}_2\text{O}_3/\text{glass}$). Water-soluble CdSe QDs, with an average diameter of 5.5 nm, were purchased from Ocean Nanotech. The QDs were capped by 3-mercaptopropionic acid (MPA). Similar QDs have been used as sensitizers for QD-based solar cells and energy conversion processes.^{5,7,8} To sensitize the Ag NP films with QDs, the Ag/ Al_2O_3 films were soaked for 2 h in an aqueous solution of CdSe QDs having a concentration of 1.1 mg/mL. The QD-sensitized Ag NP films were washed and then dried for time-resolved fluorescence detections. We note that the coverage of the QDs on the film is not directly resolved by AFM images due to the high surface heterogeneity of Ag NP films. However, the coverage can be estimated to be $\sim 50\%$ on the surface of the substrate based on the coverage of QDs on $\text{Al}_2\text{O}_3/\text{glass}$ films under the same experimental condition (Figure S1, SI). Detection of the fluorescence from QD-covered Ag NP films was accomplished with a home-built confocal microscope system coupled with a time-correlated single photon counting (TCSPC) module. The details of the setup are described in the SI. The excitation laser (tuned from 400 to 515 nm) was focused on the samples through an air objective (Olympus, 100 \times , 0.9 N.A.), and the fluorescence from QDs was collected by the same objective and captured by the detector. The detected photons were used to construct the fluorescence decay curves of QDs. The collection of photons was conducted for the same length of time (2 min) for QDs on different substrates. During detection, the samples were rastered at a speed of 0.5 $\mu\text{m}/\text{s}$ to avoid photodegradation of the QDs under laser excitation. Final decay curves were obtained by averaging over three to four data scans for each sample.

UV-vis extinction spectra of the Ag NP films with five Al_2O_3 thicknesses (2, 5, 7, 10, and 20 nm) are shown in Figure 2, indicating a broad spectral response region of the Ag NP plasmon. The plasmonic peak shows a red shift as the Al_2O_3 thickness increases, which is due to changes of average

dielectric constant in the medium surrounding the Ag NP, that is, combinations of Al_2O_3 and air. The UV-vis absorption and emission spectra of the CdSe QDs in water are shown in Figure 2. The absorption spectrum of the QDs overlaps the plasmonic region of the Ag NPs, suggesting the probability of plasmon-enhanced absorption of the CdSe QDs on Ag NP films. In contrast, the emission of the CdSe QDs centered at 590 nm is weakly coupled with the plasmon resonance of Ag NPs and thus can be expected to undergo little extinction.

In order to examine the plasmonic effect of Ag NPs on CdSe QDs, time-resolved fluorescence decays were acquired. The samples consisted of QDs on Al_2O_3 -coated Ag NP films with different Al_2O_3 thicknesses and a control substrate (Al_2O_3 /glass) without Ag NPs. The decay curves are shown in Figure 3a. Because these curves were constructed by the photons collected in the same duration of times, the areas below the curves indicate the total number of emitting photons collected from the QDs during the same time (the steady-state fluorescence intensity). The steady-state fluorescence intensities (I) and the corresponding fluorescence enhancement factors (EFs) of the QDs on different substrates are calculated according to the following equations

$$I = \int_0^t i(t) dt \quad (1)$$

$$EF = \frac{I_{\text{Ag}}}{I_0} \quad (2)$$

where $i(t)$ is the intensity at the delay time t in the decay curves and I_{Ag} and I_0 are the fluorescence intensities of QDs on Ag NP films and the Al_2O_3 /glass coverslip, respectively. The fluorescence intensity of QDs as a function of Al_2O_3 thickness is plotted in Figure 3b. The data point at zero distance/thickness is from QDs on Al_2O_3 /glass (without Ag). The Al_2O_3 layer coated on Ag NPs works as a spacer, and its thickness determines the distance from QDs to Ag NPs. Compared with the QDs on Al_2O_3 /glass, QDs on the Ag NP films exhibit dramatic changes in fluorescence intensity as a function of the Al_2O_3 thickness. When the Al_2O_3 thickness is 2 nm, the fluorescence of QDs is largely quenched, and the lifetime is significantly shortened compared with that for the QDs on the control substrate. However, when the thickness of the Al_2O_3 layer is increased to 5 nm, the fluorescence of QDs is considerably enhanced compared with that on Al_2O_3 /glass. This enhancement becomes less significant when the thickness of Al_2O_3 is further increased. This enhancement in fluorescence should not be due to the increase (by only ~24%) of surface area (hence, more sensitized QDs on the surface) of Ag NP films relative to the control substrate.

The correlation between the fluorescence intensity of QDs and the thickness of the Al_2O_3 layer suggests distance-dependent plasmonic effects of the Ag NPs including both fluorescence quenching and enhancement. First, when the interparticle distance is 2 nm, the fluorescence of QDs is significantly quenched, and their lifetime is considerably shortened. Similar quenching effects for dye molecules or QDs on Au particles have been previously reported when the separation distance was less than 5 nm.^{39,42–44,48} Such quenching has been ascribed to nonradiative energy transfer from chromophores to metal NPs. Given that this quenching mechanism results in shorter decay times, we assume that a related mechanism is operative here. Another possible quenching mechanism is charge transfer from QDs to Ag

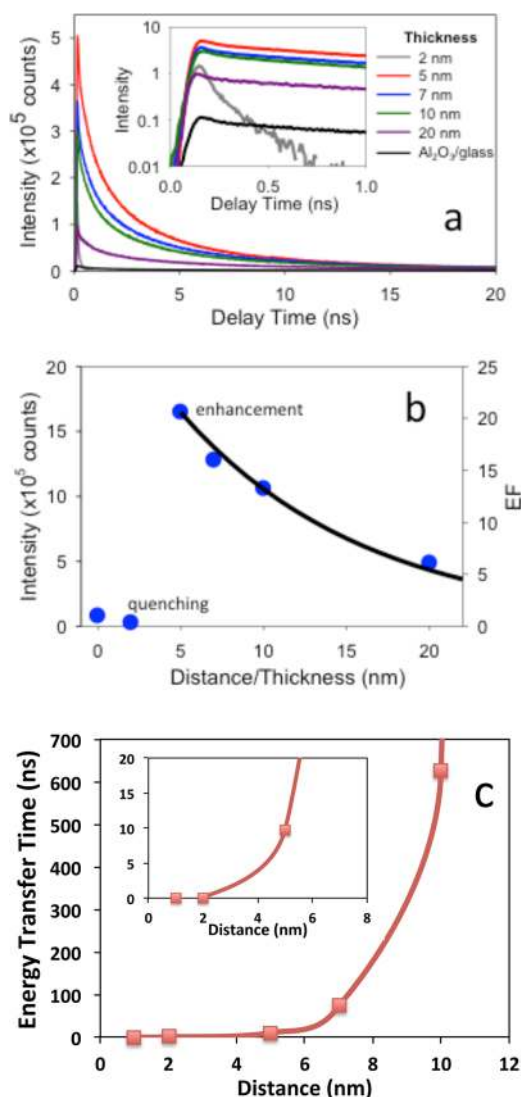


Figure 3. (a) Fluorescence decay curves of CdSe QDs sensitized on Al_2O_3 /glass (control substrate without Ag NPs) and Al_2O_3 -coated Ag NP films with different Al_2O_3 thicknesses (2, 5, 7, 10, and 20 nm). These curves are constructed by photons collected within the same duration of times under identical experimental conditions (excitation at 400 nm). (b) The EF in fluorescence intensity of QDs as a function of distance (Al_2O_3 thickness) with an excitation at 400 nm. The data point at 0 nm is from the QDs on the control substrate (Al_2O_3 /glass). The solid line is a single-exponential fit described in the text. (c) Calculated energy transfer time as a function of distance from QDs to Ag NPs based on the FRET mechanism. The details of the calculation are in the SI. The inset shows the extended view of the energy-transfer time at the short distances.

NPs. The electron-transfer time (without the spacer) from similar CdSe nanorods to Au NPs has been reported to be a few to tens of picoseconds.⁴⁹ However, Al_2O_3 is an insulating material. According to the distance-dependent CdSe electron-transfer dynamics reported in previous literature,⁵⁰ the electron-transfer time from the CdSe QD to Ag NPs at a distance of 2 nm can be estimated to be a few to tens of nanoseconds. Therefore, a <80 ps quenching time is too fast for the electron transfer with a 2 nm Al_2O_3 insulator layer. We therefore conclude that this mechanism is negligible. The energy transfer from QDs to Ag NPs can follow either the surface energy transfer (SET)^{51–54} or the Förster resonance

energy transfer (FRET) mechanisms.⁵⁵ A recent study of energy transfer from CdSe/ZnS QDs to Au NPs has indicated that the energy transfer followed the FRET mechanism when the Au NPs showed a strong LSPR.⁵⁵ The energy transfer times from QDs to the Ag NP films as a function of separation distance based on the SET and FRET mechanisms are calculated (see the SI for the details) and compared in Figure S2 (SI). At a separation distance of 2 nm, the energy-transfer times for both SET and FRET are calculated to be <80 ps, agreeing with the observed shortened lifetime (<80 ps, limited by the instrument response function) of QDs on the Ag NP film coated with Al₂O₃ of thickness 2 nm. As the distance increases to 5 nm, the energy-transfer times are calculated to be 2.5 ns for SET and 9.8 ns for FRET. However, the energy-transfer time of 2.5 ns is not observed in the fluorescence decay of QDs on Ag NPs films with 5 nm Al₂O₃, whose nonradiative decay time is calculated to be ~8 ns, more consistent with the energy-transfer time calculated by FRET. We therefore believe that the energy transfer from QDs to Ag NPs very likely follows the FRET mechanism. The FRET times as a function of distance are shown in Figure 3c.

As the thickness of Al₂O₃ increases, the nonradiative energy-transfer rate (the FRET time from QDs to the Ag NP films) is significantly attenuated to the nanosecond time scale. As well, the QD fluorescence intensity dramatically increases to values well above that for Ag-free samples. We surmise that under these conditions, plasmonic enhancement of QD absorption now dominates nonradiative transfer. At a thickness of 5 nm, the fluorescence enhancement reaches ~24. As the Al₂O₃ thickness increases beyond 5 nm, the fluorescence enhancement decays slowly from this maximum value. This distance dependence is broadly consistent with those reported in the literature.^{31,56,57} The correlation between fluorescence enhancement and distance can be described as an exponential decay function, which is logically close to the local density of optical states produced by the dipole in the Purcell enhancement of the plasmon's radiation rate.⁵⁸ As shown in Figure 3b, the plasmonic fluorescence enhancement (starting from the maximum point, F_{\max} at 5 nm) as a function of distance d (Al₂O₃ thickness) is well fit by a single-exponential function, yielding a decay parameter $D = 0.89 \text{ nm}^{-1}$

$$EF(d) = EF_{\max} \exp(-Dd) \quad (d \geq 5 \text{ nm}) \quad (3)$$

The enhancement in fluorescence of QDs on Ag NP films can, in principle, originate from plasmonic effects on both absorption and emission. The spectral position of the plasmon resonance of Ag NPs relative to the absorption and emission spectra of QDs is determinant.^{25,27,31,36,38,59,60} The steady-state fluorescence intensity (I) of QDs is controlled by the excitation rate (k_{exc}) and the quantum yield (Φ)

$$I = k_{\text{exc}} \Phi \quad (4)$$

$$\Phi = \frac{k_r}{k_r + k_{\text{nr}}} = k_r \tau \quad (5)$$

where k_r and k_{nr} are the radiative and nonradiative decay constants, respectively, and τ is the fluorescence lifetime. Because the absorption of the CdSe QDs overlaps well with the spectrum of the Ag NP plasmon resonance (see Figure 2), a dramatic enhancement in k_{exc} in QDs is reasonably expected with the excitation at 400 nm. Such an enhancement is due to the intensification of the local laser excitation field by Ag NPs, leading to a higher frequency of photon absorption in QDs.

Additionally, an enhancement in k_r can also occur once the emission spectrum of the luminescent species is coupled with the plasmon resonance of Ag NPs, leading to an increase in the fluorescence quantum yield and shortened fluorescence lifetime (if k_{nr} is not $\gg k_r$). However, because the emission spectrum of the studied CdSe QDs is red-shifted by ~150 nm relative to the extinction peak of the Ag NP plasmon film (see Figure 2), the plasmon enhancement in k_r is negligible in the currently studied system due to such a weak coupling. This conclusion is consistent with the comparison of fluorescence decay curves of QDs on different substrates, as shown in Figure S3 (SI). The decay curves show multiple exponential components. Compared with the lifetime on the control substrate (Al₂O₃/glass), the QDs on Ag NP films (with Al₂O₃ thickness ≥ 5 nm) exhibit a slightly shortened lifetime in the slow component, while the fast component (corresponding to ~70% in the amplitude) remains unchanged (see Figure S3 and Table S1 (SI) for the lifetimes of QDs on different substrates). Decay curve differences decrease as the thickness of the Al₂O₃ layer on Ag NPs increases. The slightly shortened slow-component lifetimes of QDs on Ag NPs compared to QDs on Al₂O₃/glass are attributed to the presence of the energy transfer from QDs to Ag NPs (on a tens–hundreds of nanoseconds time scale when the distance is ≥ 5 nm, as shown in Figure 3c) and/or the change in the dielectric constant of the substrate caused by the variation of Al₂O₃ thickness. We therefore conclude that the observed dramatic increase in fluorescence for QDs on Ag NP films primarily originates from the plasmon enhancement effect in the photoabsorption (k_{exc}).

The excitation-wavelength-dependent plasmonic effect further confirms that the plasmon-enhanced fluorescence is primarily due to the enhanced photon absorption for QDs on Ag NPs. Besides the excitation at 400 nm, the plasmon-enhanced fluorescence at three additional excitation wavelengths (450, 490, and 515 nm) for QDs on a Al₂O₃-coated Ag NP film with a 5 nm Al₂O₃ thickness was examined. (See the fluorescence decay curves of QDs at different excitation wavelengths in Figure S4, SI.) The fluorescence EF as a function of excitation wavelength is shown in Figure 4b. With

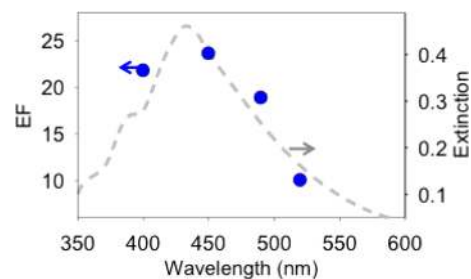


Figure 4. Fluorescence EF at different excitation wavelengths (blue dots) for QDs on an Al₂O₃-coated Ag NP film with a 5 nm Al₂O₃ thickness. Also shown in the panel is the extinction spectrum (dashed line) of the Ag NP film.

the excitation at 450 nm, a slightly higher EF (~24) than that at 400 nm is achieved. However, the EF decreases to 18 and 10, respectively, as the excitation wavelength is extended to 490 and 515 nm. The distribution of EFs at different excitation wavelengths is consistent with the shape of the extinction spectrum associated with the Ag NP plasmon resonance. An excitation wavelength with stronger coupling to the plasmon resonance leads to a larger enhancement effect in fluorescence.

This is consistent with the primary source of plasmon enhancement originating from increased photon absorption for the studied QDs on Ag NPs.

Plasmon enhancement of photon absorption by QDs on Ag NP films suggests that Ag NPs could be used to increase light harvesting by QDs for solar energy conversion. The broad and size-tunable optical spectra of QDs allow for strong (weak) coupling between the photon absorption (emission) in QDs and the plasmon resonance in Ag NPs. We have demonstrated that this coupling can lead to a dramatic enhancement in light harvesting in QDs and meanwhile, unlike in previously reported studies, maintain an unchanged photon-generated exciton lifetime with a properly tuned Ag NP-to-QD distance. We note that at a short interparticle distance (2 nm), the coupling between QDs and Ag NPs significantly reduced the exciton lifetime to <80 ps, presumably by activating ultrafast energy transfer from QDs to Ag NPs. In the context of QDs as sensitizers in solar energy conversion systems, such a shortened exciton lifetime could significantly decrease overall efficiencies by decreasing quantum yields for charge transfer from QDs to catalysts or external charge acceptors. The time scales for these processes typically range from picoseconds to nanoseconds^{7,9,61–63} and thus diminish and even eliminate the enhancement effect in light absorption. An ideal interparticle distance from QDs to Ag NPs is found to be ~5 nm, at which the light absorption in QDs is enhanced by ~20-fold with only slight attenuation of the exciton lifetime.

In conclusion, we have studied the distance-engineered plasmonic effect of Ag NPs on the photophysical properties of semiconductor QDs. Both plasmon-induced fluorescence quenching and enhancement were observed as the interparticle distance from QDs to Ag NPs was tuned from 2 to ≥5 nm. At 2 nm, ultrafast energy transfer from QDs to Ag NPs was activated to quench the fluorescence of QDs. When the distance was increased to 5 nm or greater, the plasmonic enhancement effect became dominant, leading to an intensified fluorescence by more than 20-fold with only slightly shortened exciton lifetimes. We have confirmed that the enhanced fluorescence of QDs primarily originated from the enhancement of the local excitation field by Ag NPs. Our study indicates that with a properly tuned interparticle distance and spectral coupling between QD and Ag NP plasmon resonance (strong coupling in absorption but weak in emission), dramatic plasmon enhancement of photon absorption by QDs without shortening exciton lifetimes is achievable. These findings point toward the architectures necessary for efficient use of Ag NP plasmons to boost light harvesting in QD-based solar energy conversion schemes.

■ ASSOCIATED CONTENT

Supporting Information

The details of Ag NP synthesis and ALD fabrications; experimental setup; AFM image of QDs sensitized on a glass coverslip; detailed calculation of the energy transfer rate from QDs on Ag NP films at different interparticle distances; fitting of fluorescence decays of QDs on different substrates; and fluorescence enhancement for QDs on a Ag NP film at different excitation wavelengths. This material is available free of charge via the Internet at <http://pubs.acs.org>.

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Notes

The authors declare no competing financial interest.

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