

## Dynamical Supersymmetry Breaking in Vector-Like Gauge Theories

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We provide vector-like gauge theories which break supersymmetry dynamically.

There is a piece of folklore which holds that vector-like gauge theories cannot break supersymmetry dynamically. In this paper, we point out remarkable exceptions to this piece of folklore.

Let us consider a supersymmetric  $SU(2)$  gauge theory with four doublet chiral superfields  $Q_i$ . We also introduce six singlet chiral superfields  $Z^{ij} = -Z^{ji}$ . Here  $i$  and  $j$  denote the flavor indices ( $i, j=1, \dots, 4$ ).

The tree-level superpotential of our model is given by<sup>\*\*)</sup>

$$W_{\text{tree}} = \lambda_{ij}^{kl} Z^{ij} Q_k Q_l, \quad (1)$$

where the  $\lambda_{ij}^{kl}$  denote generic coupling constants with  $\lambda_{ij}^{kl} = -\lambda_{ji}^{kl} = -\lambda_{ij}^{lk}$ . The peculiarity of this superpotential resides in the fact that it raises all the  $D$ -flat directions in the doublets  $Q_i$ , which is a necessary condition for supersymmetry to break down.<sup>1)</sup> Of course, supersymmetry remains unbroken perturbatively in this model.

The exact effective superpotential of the model, which takes into account the full nonperturbative effects, may be written in terms of gauge-invariant low-energy degrees of freedom<sup>2)</sup>

$$V_{ij} = -V_{ji} \sim Q_i Q_j \quad (2)$$

as follows:

$$W_{\text{eff}} = X(\text{Pf } V_{ij} - \Lambda^4) + \lambda_{ij}^{kl} Z^{ij} V_{kl}, \quad (3)$$

where  $X$  is an additional chiral superfield,  $\text{Pf } V_{ij}$  denotes the Pfaffian of the antisymmetric matrix  $V_{ij}$ , and  $\Lambda$  is a dynamical scale of the  $SU(2)$  gauge interaction.<sup>2),3)</sup> This is none other than a superpotential of the O'Raifeartaigh type.<sup>4)</sup> Namely, this effective superpotential yields conditions for supersymmetric vacua

$$\text{Pf } V_{ij} = \Lambda^4, \quad \lambda_{ij}^{kl} V_{kl} = 0, \quad (4)$$

which cannot be satisfied simultaneously as long as  $\Lambda \neq 0$ . Therefore we conclude that supersymmetry is dynamically broken in our model.

We note that this conclusion is not in contradiction with the index argument.<sup>5)</sup>

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<sup>\*\*)</sup> This tree-level superpotential is natural since it possesses two global symmetries. One is an axial  $U(1)$  symmetry associated with a  $Q_i$  phase transformation and the other is an anomaly-free  $R$  symmetry.

The doublets  $Q_i$  cannot be decoupled by means of mass terms  $m^{ij}Q_iQ_j$ , since the apparent masses may be absorbed in the shifts of the singlets  $Z^{ij}$ .

It is straightforward to generalize the above model to an  $Sp(N)$  gauge theory<sup>6)</sup> with  $2(N+1)$  chiral superfields in the  $2N$  representation. Here we adopt the notation  $Sp(1)=SU(2)$ .

These vector-like models might serve as a supersymmetry-breaking mechanism in the hidden<sup>7)</sup> or visible sector.

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