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# Effect of the Fabrication Technique on the Thermoelectric Performance of Mg-Based Compounds-A Case Study of n-Type Mg2Ge

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# Effect of the Fabrication Technique on the Thermoelectric Performance of Mg-Based Compounds-A Case Study of n-Type Mg2Ge

# Abstract

High performance, low cost, and low toxicity have been the main characteristics associated with magnesium-based thermoelectric materials. Nevertheless, the high volatility of magnesium creates challenges in the synthesis of these materials. In this work, n-type Mg2Ge is synthesized using a solid-state technique, fully characterized, and compared with Mg2Ge fabricated through different processes. We have found that Bi is an ineffective dopant in Mg2Ge and precipitates into Mg2Bi3. Regardless of the technique used, the loss of Mg by evaporation and formation of precipitates in Bi-doped samples resulted in a low charge carrier concentration and, consequently, a low power factor. The precipitates significantly reduced the lattice thermal conductivity, however, leading to a figure-of-merit, zT, of 0.4 at 725 K, improving the previously reported figure-of-merit, zT, of 0.2 for Sb-doped Mg2Ge. This work highlights the impact of the fabrication technique on the thermoelectric performance of Mg-based compounds.

# Disciplines

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# Effect of the Fabrication Technique on the Thermoelectric Performance of Mg-Based Compounds—A Case Study of n-Type Mg<sub>2</sub>Ge

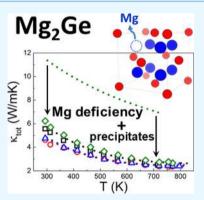
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**Supporting Information** 

**ABSTRACT:** High performance, low cost, and low toxicity have been the main characteristics associated with magnesium-based thermoelectric materials. Nevertheless, the high volatility of magnesium creates challenges in the synthesis of these materials. In this work, n-type Mg<sub>2</sub>Ge is synthesized using a solid-state technique, fully characterized, and compared with Mg<sub>2</sub>Ge fabricated through different processes. We have found that Bi is an ineffective dopant in Mg<sub>2</sub>Ge and precipitates into Mg<sub>2</sub>Bi<sub>3</sub>. Regardless of the technique used, the loss of Mg by evaporation and formation of precipitates in Bi-doped samples resulted in a low charge carrier concentration and, consequently, a low power factor. The precipitates significantly reduced the lattice thermal conductivity, however, leading to a figure-of-merit, zT, of 0.4 at 725 K, improving the previously reported figure-of-merit, zT, of 0.2 for Sb-doped Mg<sub>2</sub>Ge. This work highlights the impact of the fabrication technique on the thermoelectric performance of Mg-based compounds.



#### INTRODUCTION

Magnesium-based thermoelectric materials have garnered considerable attention because of their promising performance, low cost, and environmental friendliness.<sup>1,2</sup>  $Mg_2X$  (X = Si, Ge, and Sn) compounds crystallize in the Fm3m space group, and they exhibit similar electronic band structures and a split conduction band. Recent studies have focused on the thermoelectric performance of n-type ternary<sup>3-5</sup> and quaternary<sup>4,6,7</sup> alloys of Mg<sub>2</sub>Si, Mg<sub>2</sub>Ge, and Mg<sub>2</sub>Sn in various ratios to improve their conversion efficiency, defined by the figure-ofmerit,  $zT = S^2 T / \rho \kappa_{tot}$ , where S is the Seebeck coefficient, T is the absolute temperature,  $\rho$  is the electrical resistivity, and  $\kappa_{tot}$  is the total thermal conductivity, through the convergence of conduction bands<sup>8,9</sup> and increased mass-difference phonon scattering.<sup>10-12</sup> The binary compounds have received limited attention, however, because of their relatively low performances, especially Mg<sub>2</sub>Ge, with three reports on n-type bulk samples<sup>10,13,14</sup> and only one report on the thin film p-type Mg<sub>2</sub>Ge.<sup>15</sup> It is of great interest to understand the transport properties of these binary compounds and their interactions with the most common dopants to correctly design alloys with higher thermoelectric performance.

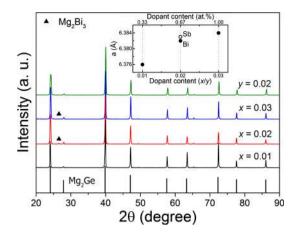
The fabrication of Mg-based semiconductors encounters intrinsic challenges because of the low vapor pressure and high reactivity of magnesium with potential crucibles. A large number of fabrication methods have been employed to synthesize Mg-based thermoelectric materials, including low-temperature ball milling,<sup>8,16</sup> solid-state reaction,<sup>12,17</sup> one-step

sintering by spark plasma sintering (SPS) using MgH<sub>2</sub> as a precursor for Mg,<sup>14,18</sup> melting and casting using crucibles made of tantalum,<sup>19</sup> alumina sealed with B<sub>2</sub>O<sub>3</sub>,<sup>20</sup> boron nitride,<sup>7</sup> graphite,<sup>12</sup> and molybdenum foil<sup>21</sup> to avoid the reaction of magnesium with quartz, where these are sealed under vacuum. Here, we have synthesized n-type Mg<sub>2</sub>Ge bulk samples by doping with Bi and Sb through the solid-state reaction. The transport properties of these samples are compared with those of previous studies,<sup>10,13</sup> including our recent report on n-type Mg<sub>2</sub>Ge samples, fabricated via one-step SPS of elemental Ge and Bi/Sb with MgH<sub>2</sub>.<sup>14</sup> A maximum *zT* of ~0.4 at 725 K is achieved for Mg<sub>2</sub>Ge<sub>0.98</sub>Bi<sub>0.02</sub>, an improvement over the *zT* of 0.2 for the Sb-doped Mg<sub>2</sub>Ge<sup>13</sup> because of a significant decrease in the thermal conductivity, while obtaining a similar power factor (PF). Our work underscores the impact that different synthesis techniques have on the properties of Mg-based thermoelectrics and the difficulty in comparing these materials when they are synthesized by different methods.

#### RESULTS AND DISCUSSION

X-ray diffraction (XRD) patterns of the sintered Mg<sub>2</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> (x = 0.01, 0.02, and 0.03) and Mg<sub>2</sub>Ge<sub>1-y</sub>Sb<sub>y</sub> (y = 0.02) samples (Figure 1) can be indexed to the cubic CaF<sub>2</sub> structure (*Fm3m*)

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**Figure 1.** Room-temperature XRD patterns of  $Mg_2Ge_{1-x}Bi_x$  (x = 0.01, 0.02, and 0.03) and  $Mg_2Ge_{1-y}Sb_y$  (y = 0.02) samples. The inset shows the lattice parameter as a function of the dopant type and content, as obtained by Rietveld refinement. Trace amounts of  $Mg_2Bi_3$  were detected in the Bi-doped samples with x = 0.02 and 0.03.

space group) of Mg<sub>2</sub>Ge. Small amounts of the secondary phase  $Mg_2Bi_3$  were detected in the Bi-doped samples with x = 0.02and 0.03. The inset of Figure 1 shows the increase in the lattice parameter of Mg<sub>2</sub>Ge with the increasing dopant concentration, as larger atoms of Bi and Sb substitute Ge atoms.<sup>22</sup> The difference between the lattice parameters of the samples with a Bi content of x = 0.02 and 0.03 is insignificant, which is consistent with the detection of Mg2Bi3 in these samples, suggesting that Bi added beyond 0.67 at. % has formed precipitates. The observed solubility of Bi in Mg<sub>2</sub>Ge obtained in the current study using a solid-state synthesis technique is significantly higher than 0.17 at. % in our previous report, based on the one-step SPS synthesis of Mg<sub>2</sub>Ge.<sup>14</sup> The latter method comprises a single step at a lower temperature that relies on the rapid reaction of Ge and MgH<sub>2</sub> upon hydrogen liberation, with an extremely short time for homogenization. The reaction of Mg with Bi also appears to have been preferred in the one-step SPS synthesis, reducing the amount of Bi available for doping. As the Mg<sub>2</sub>Bi<sub>3</sub> precipitates were already formed, the subsequent annealing step had no effect beyond stabilization of the structure. In this work, the mixture was reacted in two stages at

a higher temperature for a prolonged time, with grinding and homogenization of the powder in the interval between the stages, allowing for better dissolution of Bi in  $Mg_2Ge$ .  $Mg_2Bi_3$  precipitates still formed, however, albeit only at higher concentrations of Bi.

The backscattered electron microscopy (BSE) images of both Bi- and Sb-doped samples at the same dopant concentration of x = 0.02 and y = 0.02 (0.67 at. %) confirmed the presence of precipitates with high atomic number at grain boundaries (particle size < 1  $\mu$ m), which are indicated by arrows in Figure 2a, for the Bi-doped sample. The energydispersive spectroscopy (EDS) analysis (Figure 2c,d) reveals that these precipitates are rich in Bi, confirming the presence of the Mg<sub>2</sub>Bi<sub>3</sub> phase detected by XRD. We have also detected that these precipitates in the Bi-doped samples are synthesized by a one-step SPS process.<sup>14</sup> The Sb-doped sample is free of precipitates (Figure 2b).

Temperature-dependent Seebeck coefficient (S), electrical resistivity ( $\rho$ ), and PF, where PF =  $S^2/\rho$ , for Mg<sub>2</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> (x = 0.01, 0.02, and 0.03) and  $Mg_2Ge_{1-y}Sb_y$  (y = 0.02) samples are shown in Figure 3. All samples show a linear dependence of the Seebeck coefficient with a temperature of up to ~650 K, where it reaches a plateau (Figure 3a). Bi doping reduces the absolute S from ~210  $\mu$ V/K for x = 0.01 to -190  $\mu$ V/K for x = 0.02, after which it remains roughly constant, showing that further addition of Bi leads to an insignificant increase in the charge carrier concentration (Table 1), n, on which the position of the Fermi level and consequently the Seebeck coefficient are dependent. This result is consistent with the observation of Birich precipitates formed at grain boundaries in the x = 0.02sample. The Sb-doped sample exhibits a lower absolute S compared to the similarly Bi-doped samples, x = 0.02 and 0.03. The measured charge carrier concentrations of these samples are very similar (Table 1), indicating a comparable position of the Fermi level in relation to the bottom of the conduction band. According to eq 1

$$S = \frac{2k_{\rm B}^2 T}{3e\hbar^2} \left(\frac{\pi}{3n}\right)^{2/3} m^* (1+r)$$
(1)

where  $k_{\rm B}$  is the Boltzmann constant, *e* is the electron charge,  $\hbar$  is the reduced Planck constant,  $m^*$  is the density of states

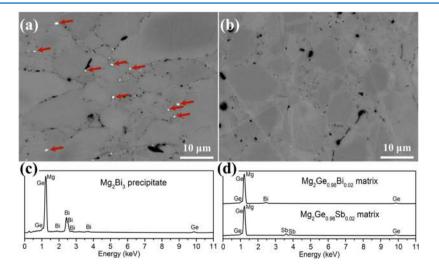
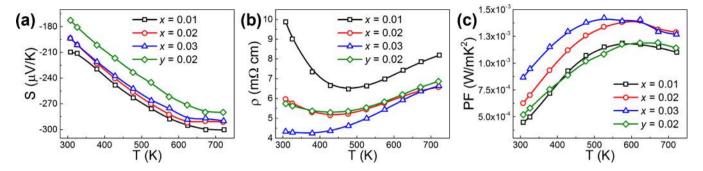


Figure 2. BSE images of  $Mg_2Ge_{0.98}X_{0.02}$ : (a) X = Bi and (b) X = Sb. Representative EDS spectra of (c)  $Mg_2Bi_3$  precipitates (indicated by the arrows in (a)) and (d)  $Mg_2Ge_{0.98}X_{0.02}$  (X = Bi and Sb).



**Figure 3.** Temperature dependence of the (a) Seebeck coefficient, (b) electrical resistivity, and (c) PF of  $Mg_2Ge_{1-x}Bi_x$  (x = 0.01, 0.02, and 0.03) and  $Mg_2Ge_{1-y}Sb_y$  (y = 0.02).

Table 1. Room-Temperature Seebeck Coefficient, Electrical Resistivity, Hall Coefficient ( $R_{\rm H}$ ), Carrier Concentration (n), and Hall Mobility ( $\mu_{\rm H}$ )<sup>*a*</sup>

	$S (\mu V/K)$	$\rho \ (m\Omega \cdot cm)$	$R_{\rm H}~({\rm cm}^3/{\rm C})$	$n (10^{18} \text{ cm}^3)$	$\mu_{\rm H}~({\rm cm}^2/{\rm V~s})$	z
Bi = 0.01	209.6	9.9	1.54	4.1	154	0.005
Bi = 0.02	193.6	5.9	0.92	6.8	154	0.010
Bi = 0.03	194.0	4.3	0.76	8.4	172	0.015
Sb = 0.02	148.7	5.7	0.78	8.0	136	0.010

<sup>*a*</sup>The final Mg deficiency, *z*, in Mg<sub>2-z</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> (x = 0.01, 0.02, and 0.03) and Mg<sub>2-z</sub>Ge<sub>1-y</sub>Sb<sub>y</sub> (y = 0.02) was estimated based on the measured carrier concentration.

effective mass, and *r* is the scattering parameter.<sup>23</sup> Note that the Seebeck coefficient is proportional to  $\frac{m*(1+r)}{n^{2/3}}$ . The carrier concentration values of samples doped with Bi and Sb are similar. Therefore, the difference in the Seebeck coefficient must arise from the differences in the effective mass, suggesting that Sb and Bi have different effects on the electronic band structure, as the scattering mechanism is unlikely to change between the dopants.

The electrical resistivity of the samples initially decreases with temperature, showing the typical behavior of intrinsic semiconductors, and exhibits metallic behavior at higher temperatures (Figure 3b). The magnitude and temperature of this behavioral change decrease with the dopant concentration, as the thermal energy needed to promote carriers to the conduction band decreases with the charge carrier concentration.<sup>24</sup> This effect is only visible because of the relatively low charge carrier concentrations in the current study, even at x = 0.03, as displayed in Figure 4, where a rapid decrease in the Hall

**Figure 4.** Temperature dependence of the Hall coefficient for  $Mg_2Ge_{1-x}Bi_x$  (x = 0.01, 0.02, and 0.03) and  $Mg_2Ge_{1-y}Sb_y$  (y = 0.02) samples.

coefficient ( $R_{\rm H}$ ) above ~200 K indicates the thermal excitation of electrons to the conduction band. The room-temperature Hall carrier mobility (Table 1),  $\mu_{\rm H}$ , has increased slightly by the addition of Bi. This is due to very low charge carrier concentrations of lightly doped samples, which show the behavior of transport properties of the nondegenerate semiconductor.

The stoichiometry of Mg is known to have a significant effect on the electrical characteristics of Mg-based thermoelectric materials.<sup>25</sup> Each excess Mg atom occupies interstitial sites, contributing two electrons to the conduction band,<sup>25</sup> whereas each Mg atom vacancy acts either as a double hole donor<sup>26,27</sup> or as an electron trap localized within the band gap.<sup>28</sup> Regardless of the principle regulating Mg vacancies, Mg deficiency severely decreases the carrier concentration of n-type samples. This explains the much lower electron carrier concentration obtained in the current study,  $n \approx 10^{18}$  cm<sup>-3</sup>, compared to the similarly Sb-doped Mg<sub>2</sub>Ge,<sup>13</sup> which resulted in a roomtemperature electrical resistivity of 0.195 m $\Omega$ -cm and a carrier concentration of 5.8 × 10<sup>20</sup> cm<sup>-3</sup> for Mg<sub>2.2</sub>Ge<sub>0.99</sub>Sb<sub>0.01</sub> with excess Mg. The deficiency of Mg (z) in Mg<sub>2-z</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> and  $Mg_{2-z}Ge_{1-y}Sb_y$  was estimated using  $z = \frac{1}{2}\left(\frac{V \times n}{4} - x\right)$ , where V is the unit cell volume, n is the carrier concentration, and x is the dopant stoichiometry, and the results are displayed in Table 1. This equation estimates the Mg deficiency by considering Bi atoms and Mg vacancies as the only sources of charge carriers/ vacancies in the system, that is, assuming that there are no Mg<sub>2</sub>Bi<sub>3</sub> precipitates.

The bipolar effect is observed at temperatures above ~650 K in the Seebeck coefficient plot (Figure 3a), similar to our previous study.<sup>14</sup> This effect is not visible on the thermal conductivity plot as a function of temperature (Figure 5a) because of the negligible bipolar component of the thermal conductivity<sup>29</sup> ( $\kappa_b$ ) (Figure S1 of the Supporting Information). The total ( $\kappa_{tot}$ ) and lattice ( $\kappa_{lat}$ ) thermal conductivities of the

samples are shown in Figure 5a. The Wiedemann–Franz relation,  $\kappa_e = L_0 T/\rho$ , is used to calculate the electronic thermal

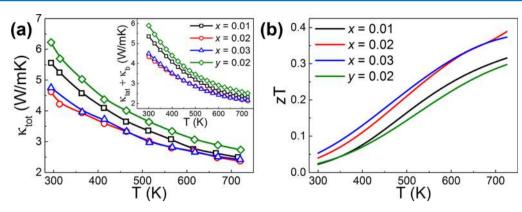


Figure 5. Temperature dependence of the (a) total thermal conductivity (with the sum of the lattice and bipolar components as an inset) and (b) figure-of-merit, zT, of Mg<sub>2</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> and Mg<sub>2</sub>Ge<sub>1-x</sub>Sb<sub>y</sub>.

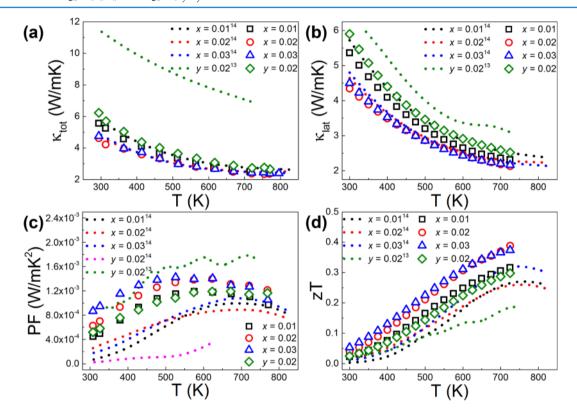


Figure 6. Temperature dependence of the (a) total thermal conductivity, (b) lattice component of the thermal conductivity, (c) PF, and (d) figureof-merit, *zT*, of Mg<sub>2</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> (x = 0.01, 0.02, and 0.03) and Mg<sub>2</sub>Ge<sub>1-y</sub>Sb<sub>y</sub> (y = 0.02) samples from this work compared with other reports on n-Mg<sub>2</sub>Ge.<sup>13,14</sup>

conductivity, where  $L_0$  is the Lorenz number and  $\rho$  is the electrical resistivity. The Lorenz number used in the calculation of  $\kappa_e$  is the theoretical limit for degenerate semiconductors 2.45  $\times 10^{-8} \text{ V}^2 \text{ K}^{-2}$ .<sup>30</sup> The lattice and bipolar thermal conductivities of all samples are calculated by deducting the electronic contribution to the total thermal conductivity using  $\kappa_{\text{tot}} = \kappa_{\text{lat}} + \kappa_e + \kappa_b$  (Figure 5a inset). The lattice thermal conductivity dominates  $\kappa_{\text{tot}}$  in all samples, and it decreases with the temperature and Bi dopant concentration. Even though the charge carrier density increased slightly with the added dopant, the electronic component of the thermal conductivity is too small to significantly contribute to the total thermal conductivity.

The total thermal conductivity of the samples produced in the current study is much lower than that in the previous report on heavily doped  $n-Mg_2Ge^{13}$  (Figure 6a). Figure 6b shows that

the lattice thermal conductivity in the previous report is similar to that of the Sb-doped sample in the current study, which is evidence of the significantly lower contribution of the electronic thermal conductivity, originating from the low carrier concentration of our samples. The lattice thermal conductivity of the Bi-doped samples is lower than that of the Sb-doped one owing to the increased phonon scattering by Mg<sub>2</sub>Bi<sub>3</sub> precipitates and point defects. Moreover, the effect of the precipitates is visible in the significant decrease in the lattice thermal conductivity with the increased Bi concentration. The formation of Mg<sub>2</sub>Bi<sub>3</sub> leads to further depression of the electron concentration by reducing the amount of available Bi, which acts as an electron donor in Mg2Ge, and by leading to the additional loss of Mg from the matrix. Furthermore, the low charge carrier concentration of the samples produced in this work makes  $\kappa_e$  practically negligible, representing only ~4% of

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 $\kappa_{totv}$  whereas in highly doped Mg<sub>2</sub>Ge,<sup>13</sup> it accounts for ~50%. Despite the significantly higher electrical resistivity obtained in the previous report<sup>14</sup> and in this work, when compared to the heavily Sb-doped Mg<sub>2</sub>Ge,<sup>13</sup> our samples exhibit a similar PF (Figure 6c). The final *zT* of our samples is therefore higher, owing to the reduced thermal conductivity (Figure 6d). These results suggest that studying the interaction of Bi with high *zT* Mg-based thermoelectrics alloyed with Ge might be of interest to improve our understanding of their high performance and perhaps guide further improvements. To take advantage of the impact of precipitates on the thermal conductivity while maintaining low electronic resistivity, we suggest a double-doping approach with Sb and Bi.

#### CONCLUSIONS

n-type Mg<sub>2</sub>Ge bulk samples doped with Bi or Sb were synthesized via a solid-state reaction and compared with previous reports on n-type Mg<sub>2</sub>Ge fabricated by different processes. Despite the attempt to compensate for the Mg loss in the current study, all samples exhibited deficiency of magnesium, heavily depressing the charge carrier concentration. Furthermore, Bi reacted with Mg to form Mg<sub>2</sub>Bi<sub>3</sub> precipitates, leading to further Mg loss and reduced doping of Mg<sub>2</sub>Ge by Bi. This led to a high electrical resistivity, a high Seebeck coefficient, and a low lattice thermal conductivity, whereas the electronic thermal conductivity was negligible. This resulted in a maximum zT of ~0.4 at 725 K for Mg<sub>2</sub>Ge<sub>0.98</sub>Bi<sub>0.02</sub>. These results highlight the impact of the fabrication technique on the magnesium stoichiometry in Mg-based thermoelectric compounds and indicate that Bi is an unsuitable dopant for Mg<sub>2</sub>Ge.

#### MATERIALS AND METHODS

Magnesium turnings (Mg, Alfa Aesar, 99.98%), germanium powder (Ge, Alfa Aesar, 99.999%), bismuth powder (Bi, Alfa Aesar, 99.999%), and antimony shots (Sb, Alfa Aesar, 99.999%) were used to synthesize  $Mg_2Ge_{1-x}Bi_x$  (x = 0.01, 0.02, and 0.03) and  $Mg_2Ge_{1-y}Sb_y$  (y = 0.02). Stoichiometric mixtures of elements were sealed in a boron nitride-coated vacuum-sealed quartz ampoule, with 50 wt % extra Mg to compensate for losses during the synthesis processes. The required amount of extra magnesium was optimized to obtain a single-phase compound, that is, free of unreacted Ge, as determined by the XRD analysis. This approach was necessary to directly compare the impact of different synthesis techniques on the characteristics of Mg<sub>2</sub>Ge samples with added Bi and Sb. The mixtures were reacted at 1123 K for 30 min and annealed at 673 K for 10 h. The product was hand-ground to a fine powder in an agate mortar and pestle, and the resulting powders were sintered in a 12 mm graphite die using SPS at 1073 K and an axial pressure of 50 MPa for 1 h under vacuum. All material and powder handling and preparation were performed in a glovebox under a positive-pressure argon atmosphere.

The crystal structure and phase characterization were obtained with a MAC Science X-ray diffractometer with Cu K $\alpha$  radiation ( $\lambda$  = 1.544 Å, 40 kV, 25 mA). The obtained powder XRD patterns were fitted using Rietveld refinement to estimate the lattice parameters of the synthesized samples. The microstructure and phase composition of the Mg<sub>2</sub>Ge<sub>1-x</sub>Bi<sub>x</sub> and Mg<sub>2</sub>Ge<sub>1-y</sub>Bi<sub>y</sub> samples were analyzed by a scanning electron microscopy (SEM; JEOL 7001F SEM) instrument, with the SEM equipped with EDS. The thermal diffusivity (*D*) was obtained using the laser flash method with a Linseis LFA 1000

instrument and used to calculate the thermal conductivity ( $\kappa$ ), from  $\kappa = d \cdot D \cdot C_p$ . The density (d) was measured using dimensions and weight of the sample, whereas the heat capacity ( $C_p$ ) was obtained from the literature.<sup>8,31,32</sup> The Seebeck coefficient (S) and the electrical conductivity ( $\sigma$ ) were obtained using the slope method in a quasi-steady-state mode of a Linseis LSR-3 instrument, with temperature differences between the probes of 2–10 K in the parallelepiped-shaped samples. Hall effect measurements were performed using a Physical Property Measurement System from Quantum Design from 5 to 400 K.

#### ASSOCIATED CONTENT

#### **Supporting Information**

The Supporting Information is available free of charge on the ACS Publications website at DOI: 10.1021/acsomega.7b01389.

Lattice and bipolar conductivity as a function of inverse temperature (PDF)

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#### Notes

The authors declare no competing financial interest.

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