## ELECTRONIC SUPPORTING INFORMATION

# Efficient Acceptor Groups for NLO Chromophores: Competing Inductive and Resonance Contributions in Heterocyclic Acceptors Derived from 2- 

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To ensure that the trends obtained by the SOS calculations are not methodological artifacts, we have also carried out Finite Field calculations. Here, we followed the local contribution formulation ${ }^{1,2}$ based on second derivatives of atomic charges. These were approximated by finite differences obtained from $\mathrm{INDO}^{3} / \mathrm{HF}$ Mulliken charges of molecules with applied electric fields between zero and $\pm 5.14 \times 10^{11}$ $\mathrm{V} / \mathrm{m}\left(10^{-3}\right.$ atomic units). A more detailed description of the applied methodology is given in Ref. [4].

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## TABLES

Table S1: Orientationally averaged second-order polarizabilities $\left(<\beta>=\left\{\beta_{\mathrm{x}}{ }^{2}+\beta_{\mathrm{y}}{ }^{2}+\beta_{\mathrm{z}}{ }^{2}\right\}^{1 / 2}\right)$, long-axis components of the second-order polarizability $\left(\beta_{\mathrm{xxx}}\right)$ for both SOS and FF values, and the second-order molecular polarizabilities $\left(\beta_{\mathrm{xxx}}\right)$ calculated with the two-state model for chromophores $\mathbf{1}$ to $\mathbf{1 0}$.

| chromophore | X | $<\beta>\left(10^{-30} \mathrm{esu}\right)$ <br>  | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | $\mathrm{SCI} / \mathrm{SOS}$ | $\mathrm{SCI} / \mathrm{SOS}$ | FF | two-state |  |
| 2 | $\mathrm{SiH}_{2}$ | 310 | 301 | 336 | model |
| 3 | $\mathrm{CH}_{2}$ | 315 | 314 | 355 | 502 |
| 4 | $\mathrm{C}=\mathrm{CH}_{2}$ | 352 | 350 | 387 | 520 |
| 5 | NH | 281 | 290 | 330 | 578 |
| 6 | CO | 467 | 443 | 487 | 471 |
| 7 | S | 323 | 321 | 362 | 755 |
| 8 | O | 316 | 320 | 368 | 532 |
| 9 | $\mathrm{C}=\mathrm{CHNO}_{2}$ | 506 | 489 | 543 | 499 |
| 10 | $\mathrm{SO}_{7}$ | $\mathrm{SO}_{2}$ | 393 | 500 | 379 |

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| chromophore | $E_{g e}(\mathrm{eV})$ | $\mu_{g, x}(D)$ | $\mu_{e, x}(D)$ | $\Delta \mu_{g e, x}(D)$ | $\mu_{g e, x}(D)$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 2.50 | 6.1 | 15.4 | 9.3 | 12.0 |
| 2 | 2.49 | 7.7 | 16.8 | 9.1 | 12.3 |
| 3 | 2.48 | 7.0 | 17.2 | 10.2 | 12.2 |
| 4 | 2.52 | 8.6 | 17.2 | 8.6 | 12.2 |
| 5 | 2.37 | 11.2 | 23.1 | 11.9 | 12.4 |
| 6 | 2.49 | 9.8 | 19.3 | 9.5 | 12.2 |
| 7 | 2.48 | 11.1 | 19.8 | 8.7 | 12.1 |
| 8 | 2.34 | 10.7 | 23.9 | 13.3 | 12.1 |
| 10 | 2.43 | 13.5 | 24.3 | 10.8 | 12.2 |

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Table S3: Orientationally averaged second-order polarizabilities $\left(<\beta>=\left\{\beta_{x}{ }^{2}+\beta_{y}{ }^{2}+\beta_{z}{ }^{2}\right\}^{1 / 2}\right)$, long-axis components of the second-order polarizability $\left(\beta_{\mathrm{xxx}}\right)$ for both SOS and FF values, and the second-order molecular polarizabilities ( $\beta_{\mathrm{xxx}}$ ) calculated with the two-state model for chromophores $\mathbf{1 1}$ to $\mathbf{2 0}$.

| chromophore | X | $\begin{gathered} \left\langle\beta>\left(10^{-30} \mathrm{esu}\right)\right. \\ \mathrm{SCI} / \mathrm{SOS} \end{gathered}$ | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ <br> SCI/SOS | $\begin{gathered} \beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right) \\ \mathrm{FF} \end{gathered}$ | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ <br> two-state <br> model |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 11 | $\mathrm{SiH}_{2}$ | 134 | 145 | 140 | 271 |
| 12 | $\mathrm{CH}_{2}$ | 130 | 140 | 140 | 257 |
| 13 | $\mathrm{C}=\mathrm{CH}_{2}$ | 156 | 170 | 160 | 307 |
| 14 | NH | 107 | 116 | 117 | 183 |
| 15 | CO | 210 | 225 | 203 | 427 |
| 16 | S | 134 | 144 | 140 | 267 |
| 17 | O | 122 | 131 | 133 | 233 |
| 18 | $\mathrm{C}=\mathrm{CHNO}_{2}$ | 191 | 205 | 187 | 339 |
| 19 | SO | 163 | 176 | 167 | 335 |
| 20 | $\mathrm{SO}_{2}$ | 212 | 226 | 209 | 404 |

## Supplementary material (ESI) for Journal of Materials Chemistry

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| chromophore | $E_{g e}(\mathrm{eV})$ | $\mu_{g, x}(D)$ | $\mu_{e, x}(D)$ | $\Delta \mu_{g e, x}(D)$ | $\mu_{g e, x}(D)$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 11 | 3.04 | 9.4 | 17.7 | 8.3 | 11.4 |
| 12 | 3.03 | 10.6 | 18.0 | 7.4 | 11.7 |
| 13 | 2.98 | 10.2 | 19.2 | 9.1 | 11.4 |
| 14 | 3.08 | 11.8 | 17.4 | 5.7 | 11.5 |
| 15 | 2.83 | 12.5 | 23.9 | 11.4 | 11.3 |
| 16 | 3.03 | 12.2 | 20.2 | 8.0 | 11.5 |
| 17 | 3.04 | 11.0 | 17.2 | 6.2 | 12.0 |
| 18 | 2.84 | 14.7 | 26.3 | 11.6 | 9.6 |
| 19 | 2.96 | 14.3 | 24.1 | 9.8 | 11.3 |
| 20 | 2.85 | 15.8 | 26.4 | 10.6 | 11.5 |

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Figure S1: Calculated values for $\beta_{\mathrm{xxx}}$ (a) derived from the two-state expression, considering only the first intense excited state, with the corresponding change in state dipole moment ( $\Delta \mu_{\mathrm{ge}, \mathrm{x},}$, open squares) and the reciprocal of the square of the transition energy ( $1 / E_{\mathrm{ge}}{ }^{2}$, solid squares) shown in (b), for chromophores 11 to 20.



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Table S5: Orientationally averaged second-order polarizabilities $\left(<\beta>=\left\{\beta_{x}{ }^{2}+\beta_{y}{ }^{2}+\beta_{z}{ }^{2}\right\}^{1 / 2}\right)$, long-axis components of the second-order polarizability $\left(\beta_{\mathrm{xxx}}\right)$ for both SOS and FF values, and the second-order molecular polarizabilities ( $\beta_{\mathrm{xxx}}$ ) calculated with the two-state model for chromophores $\mathbf{2 1}$ to $\mathbf{3 0}$.

| chromophore | X | $\begin{gathered} \hline<\beta>\left(10^{-30} \mathrm{esu}\right) \\ \mathrm{SCI} / \mathrm{SOS} \end{gathered}$ | $\begin{gathered} \beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right) \\ \mathrm{SCI} / \mathrm{SOS} \end{gathered}$ | $\begin{gathered} \beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right) \\ \mathrm{FF} \end{gathered}$ | $\beta_{\mathrm{xxx}}\left(10^{-30} \mathrm{esu}\right)$ <br> two-state <br> model |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 21 | $\mathrm{SiH}_{2}$ | 278 | 293 | 263 | 605 |
| 22 | $\mathrm{CH}_{2}$ | 277 | 290 | 263 | 520 |
| 23 | $\mathrm{C}=\mathrm{CH}_{2}$ | 279 | 288 | 264 | 487 |
| 24 | NH | 265 | 275 | 255 | 525 |
| 25 | CO | 364 | 375 | 327 | 443 |
| 26 | S | 308 | 323 | 286 | 558 |
| 27 | O | 296 | 307 | 277 | 569 |
| 28 | $\mathrm{C}=\mathrm{CHNO}_{2}$ | 326 | 334 | 301 | 423 |
| 29 | SO | 370 | 389 | 339 | 162 |
| 30 | $\mathrm{SO}_{2}$ | 454 | 476 | 405 | 650 |

Table S6: Transition energies ( $E_{\mathrm{ge}}$ ), x-component of the dipole moments for the ground state $\left(\mu_{\mathrm{g}, \mathrm{x}}\right)$ and the dominant electronic excited state $\left(\mu_{\mathrm{e}, \mathrm{x}}\right)$, the change in state dipole moments $\left(\Delta \mu_{\mathrm{ge}, \mathrm{x}}\right)$, and the transition dipole moment ( $\mu_{\mathrm{ge}, \mathrm{x}}$ ) values calculated with method for chromophores 21 to 30. In addition, the corresponding INDO/SCI/SOS calculated parameters $E_{\mathrm{ge}^{e},}, \mu_{\mathrm{e}^{\mathrm{e}}, \mathrm{x}}, \Delta \mu_{\mathrm{ge}^{e}, \mathrm{x}}$, and the transition dipole moments $\mu_{\mathrm{ge}^{\prime}, \mathrm{x}}$ and $\mu_{\mathrm{ee}, \mathrm{x}}$ (with their sign relative to $\mu_{\mathrm{ge}, \mathrm{x}}$ ) for the excited state e' are given for chromophores $\mathbf{2 5}$ and 28 to $\mathbf{3 0}$.

| chromophore | $E_{g e}$ | $\mu_{g, x}$ | $\mu_{e, x}$ | $\Delta \mu_{g e, x}$ | $\mu_{g e, x}$ | $E_{g e^{\prime}}$ | $\mu_{e e^{\prime}, x}$ | $\mu_{e^{\prime}, x}$ | $\Delta \mu_{g e^{\prime}, x}$ | $\mu_{g e^{\prime}, x}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
|  | $(\mathrm{eV})$ | $(D)$ | $(D)$ | $(D)$ | $(D)$ | $(e V)$ | $(D)$ | $(D)$ | $(D)$ | $(D)$ |
| 21 | 2.63 | 6.6 | 18.9 | 12.3 | 10.9 |  |  |  |  |  |
| 22 | 2.66 | 8.2 | 20.4 | 12.2 | 11.4 |  |  |  |  |  |
| 23 | 2.67 | 7.5 | 19.8 | 12.3 | 11.0 |  |  |  |  |  |
| 24 | 2.62 | 8.7 | 21.5 | 12.8 | 11.0 |  |  |  |  |  |
| 25 | 2.57 | 10.5 | 22.0 | 11.5 | 10.4 | 2.85 | -2.9 | 22.7 | 12.2 | -5.6 |
| 26 | 2.57 | 9.4 | 22.5 | 13.1 | 11.0 |  |  |  |  |  |
| 27 | 2.63 | 11.0 | 23.5 | 12.5 | 11.6 |  |  |  |  |  |
| 28 | 2.61 | 12.2 | 24.0 | 11.7 | 10.2 | 2.89 | 1.4 | 25.9 | 13.7 | -3.6 |
| 29 | 2.44 | 11.5 | 17.6 | 6.1 | 8.2 | 2.51 | -7.9 | 16.3 | 4.9 | -7.1 |
| 30 | 2.33 | 13.1 | 26.5 | 13.5 | 10.6 | 3.43 | -8.0 | 15.4 | 2.3 | -3.6 |

Figure S3: Evolution of $\beta_{\mathrm{xxx}}$ as a function of intermediate states normalized to the converged $\beta_{\mathrm{xxx}}$ value for chromophores $\mathbf{1}$ to 30. ${ }^{\text {a) }}$

a) The convergence plots for 21, 25, and 28-30 (lower right panel) have been separated from the rest of the series $\mathbf{2 2}$ 24, 26, 27 (lower left panel) to illustrate their distinctively different convergence behavior.

We note that the convergence pattern is similar in 1-10 (upper left panel); the first excited state $e$ gives rise to the largest contribution to $\beta$ which is partially compensated by a higher-lying state $e^{\prime}$. The ratio between the $e$ - and the $e^{\prime}$-related $\beta$ contributions essentially remains constant in $\mathbf{1 - 1 0}$. In 11-20 (upper right panel), again the first (or the second in $\mathbf{1 5}$ and 18) excited state $e$ gives rise to the largest contribution to $\beta_{\text {initial. }}$. This initial $\beta_{\text {initial }}$ value is partially compensated by either one or multiple higher-lying excited states; however, the combined contribution of the latter relative to $\beta_{\text {initial }}$ is constant. However, not all compounds in 21-30 (lower panels) share a convergence behavior similar to the one observed in 1-20. While in 21-24, 26 and 27 the initial $\beta_{\text {initial }}$ value is essentially built up by one excited state $e$, at least two low-lying excited states with large $\mu_{\mathrm{ge}}$ contribute to $\beta_{\text {initial }}$ in $\mathbf{2 5}$ and 28-30 (see also Figure S4). As a consequence, a two-state model cannot

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account for the evolution of the converged $\beta$ in 21-30, since it would not consider one of these two strongly participating states in 25 and 28-30.

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Figure S4: Evolution of $\beta_{\mathrm{xxx}}$ (normalized to the converged $\beta_{\mathrm{xxx}}$ ) as a function of intermediate states for chromophore 25.


Figure S5: Calculated values for $\beta_{\mathrm{xxx}}$ derived from the two-state (solid squares) and three-state (opensquare) expression, considering only the lowest (two) excited state(s) with significantly large oscillator strength for chromophores 21 to 30.


