## Paper

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# Efficient Measurement System to Investigate Micro-Doppler Signature of Ballistic Missile

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## Abstract

Micro-Doppler (MD) shift caused by the micro-motion of a ballistic missile (BM) can be very useful to identify it. In this paper, the MD signatures of three scale-model BMs are investigated using a portable measurement system. The measurement system consists of an X-band 2-by-2 phase comparison mono-pulse radar, and a mechanical device that can impart controlled spinning and coning motions simultaneously to a model to yield the MD signature that replicates the characteristic of each target and the corresponding micro-motion. The coning motion determined the overall period of MD, and the spinning motion increased its amplitude. MD was also dependent on aspect angle. The designed system is portable, and can implement many micro-motions; it will contribute to analysis of MD in various situations.

Key words: Micro-Motion, Mono-Pulse, Range-Doppler Algorithm, Ballistic Missile, Micro-Doppler

## 1. Introduction

Ballistic missiles (BM) are fast projectiles with small radar cross section (RCS) that follow a ballistic trajectory. Defending against BMs is a vital task in warfare. The range estimated using the two-way time-delay between the radar and BM, and the Doppler shift caused by the velocity are widely used to intercept BMs [1, 2]. However, a BM warhead can be accompanied by decoys, so efficient methods to distinguish the warhead from the decoys are required.

One efficient way to distinguish the warhead from the decoys is to exploit the differences in their micro-motions.

Representative micro-motion dynamics are spinning and coning, which form the micro-Doppler (MD) feature that facilitates warhead recognition [3, 4, 5]. A maneuvering warhead travels along the flight trajectory with inherent micro-motion dynamics, which can be regarded as a unique signature of the warhead.

Most previous research related to MD is based on simulation of a target composed of isotropic point scatterers; in this system, discrepancies between the real scattering mechanism and the point scatterer may lead to significant errors in recognition. Thus, MD analysis based on the measured data is required. Recent measurement systems

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use stepped-frequency waveform (SFW) radar to measure targets in an anechoic chamber [6, 7]. However, SFW requires compensation for inter-pulse phase errors [8], and the fixed motion equipment [4, 9], and the target must be portable so that its MD can be analyzed in various clutter environments. Furthermore, the rotation velocity of the target for micromotion dynamics must also be increased to resemble that of a real target.

In this paper, we investigated the dynamic MD signatures of scale models of three missiles. We used an efficient X-band measurement system that can measure MD in various cluttered environments. The measurement system is composed of a dual-motor system for micro-motion dynamics, and an X-band 2-by-2 phase-comparison monopulse (PCMP) radar to receive reflected radar signals from four antennas. Measurement results conducted outside the anechoic chamber demonstrates the efficiency of the devised system by clearly representing the MD signature of each target.

## 2. Micro-Doppler Principle and Measurement System

#### 2.1 Micro-Doppler principle of BM

Spinning and coning (Fig. 1) are rotations around an axis in a local coordinate (x, y, z). The axis of spinning is inside the target and the axis of coning is outside it.

We use the procedure introduced in [4, 10] to derive the principle of MD of BM because the procedure is very easy to understand and clearly explains MD of BM. The monostatic chirp radar signal used for the measurement system at fast time t is

$$s_0(t) = \exp\left(j2\pi\left(f_0 t + \frac{Bt^2}{2\tau}\right)\right) \times rect\left(\frac{t}{\tau}\right)$$
(1)

where  $f_0$  is the start frequency, *B* is the bandwidth,  $\tau$  is the pulse duration, and  $rect(\cdot) = 1$  over  $t-\tau/2 \le t \le t+\tau/2$  and 0 otherwise.

The reflected signal from a BM with a micro-motion is sampled at pulse repetition period  $T_s$  and is given by

$$g(t,t_s) = \sum_{k=1}^{K} A_k \exp\left(j2\pi \left(f_0(t-\tau_k(t_s)) + \frac{B(t-\tau_k(t_s))^2}{2\tau}\right)\right) \times rect\left(\frac{t-\tau_k(t_s)}{\tau}\right)$$
(2)

where  $t_s$  is the slow time which is a multiple of  $T_s$ , K is the number of the scatterers that comprise the BM,  $A_k$  is the amplitude of scatterer k, and  $\tau_k(t_s)$  is the two-way fast-time delay at  $t_s$  between the radar and k. Assuming that k located at  $(x_k, y_k)$  is observed at  $\theta(t_s)$  in the *x*-*y* plane,  $\tau_k(t_s)$  is given

by [10]

$$\tau_k(t_s) = \frac{2}{c} \left( R(t_s) + x_k(t_s) \cos \theta(t_s) + y_k(t_s) \sin \theta(t_s) \right)$$
(3)

where  $R(t_s)$  is the distance from the radar to the origin of the local coordinate, and *c* is the speed of the light.

Then high-resolution range profiles (RPs) can be obtained by matched-filtering  $g(t, t_s)$  by (1) in the *t* domain [11]. Because this operation is time-consuming, it is conducted in the fast-time frequency domain *f*; the resultant signal is

$$S_{M}(f,t_{s}) = \sum_{k=1}^{K} A_{k} \exp\left(-j2\pi f \tau_{k}(t_{s})\right)$$
(4)

If  $R(t_s)$  is removed by subtracting the known range to the target, and  $P_k(t_s)$  is defined as  $x_k(t_s)\cos\theta(t_s) + y_k(t_s)$  $\sin\theta(t_s)$ , then  $S_M(f, t_s)$  in the f domain can be inverse Fouriertransformed to

$$S_{M}(t,t_{s}) = \int_{t_{0}}^{t_{0}+B} S_{M}(f,t_{s}) \exp(j2\pi ft) df$$
  
=  $\sum_{k=1}^{K} A_{k}B \operatorname{sinc}\left(\pi B\left(t-\frac{2}{c}P_{k}(t_{s})\right)\right) \exp\left(j\pi B\left(t-\frac{2}{c}P_{k}(t_{s})\right)\right)$  (5)

Then by selecting the radar signal in range bin = ct/2 that contains the target, and differentiating the complex signal with respect to  $t_s$ , the time-varying Doppler frequency can be obtained as

$$f_D(t_s) = -\frac{B}{2} \frac{dP(t_s)}{dt_s}$$
(6)

Because  $P(t_s)$  depends on the time-varying micro-motion of BM, the number of

useful features that can be extracted increases after MD is represented in the time-frequency domain. In this paper, we represent MD by using the spectrogram (SP) given by

$$SP[x(t_s)](t_d, f_s) = \int_{-\infty}^{\infty} x(t_s) w(t_s - t_d) \exp(-j2\pi f t_s) dt_s$$
(7)

where  $w(\cdot)$  is a window function.

MD variation of a BM that is spinning and coning can be modeled by a rotation matrix formed by Rodrigues' rotation formula [12]. Assuming that a vector  $\overline{v}$  is rotated by  $\theta$  around an axis  $\overline{k}_e = (k_1, k_2, k_3)$ , the rotated vector  $\overline{v}_{rot}$  is expressed by

$$\overline{v}_{rot} = (I\cos\theta + [\overline{k}_e]_x\sin\theta + (1 - \cos\theta)\overline{k}_e\overline{k}_e^T) = \mathbf{M}\overline{v}$$
(8)

where

$$\begin{bmatrix} \bar{k} \end{bmatrix}_{x} = \begin{bmatrix} 0 & -k_{3} & k_{2} \\ k_{2} & 0 & -k_{1} \\ -k_{2} & k_{1} & 0 \end{bmatrix}$$
(9)

Similarly, for a scatter engaged in spinning and coning

motion starting from  $\overline{p}_0$ , the position at  $t_s$  is a product of two matrices

$$\vec{p}_t = \mathbf{M}_{\mathbf{c}}(t_s)\mathbf{M}_{\mathbf{s}}(t_s)\vec{p}_0 \tag{10}$$

where  $\mathbf{M}_{c}$  and  $\mathbf{M}_{s}$  are the rotation matrix constructed using the coning velocity vector and the spinning velocity vector, respectively (Fig. 1; Eq. 8). As a result, MD of the BM is given by

$$MD = \frac{2}{\lambda} \left[ \frac{d(\mathbf{M}_{c}\mathbf{M}_{s})}{dt_{s}} \overline{p}_{0} \right] \bullet \overline{a}_{R} = \frac{2}{\lambda} \left[ \left( \frac{d\mathbf{M}_{c}}{dt_{s}} \mathbf{M}_{s} + \mathbf{M}_{c} \frac{d\mathbf{M}_{s}}{dt_{s}} \right) \overline{r}_{0} \right] \bullet \overline{a}_{R}$$
(11)

where  $\overline{a}_{R}$  is the unit vector of the radar line of sight,  $\lambda$  is the wavelength, and • is the inner-product operator. MD is the combination of the spinning and the coning frequencies due to the multiplication of the elements in the coning and the spinning matrices.



Fig. 1. Micro-motion dynamics of BM.



(b) Hanges for coning ang

Fig. 2. Schematic of the micro-motion device.

#### 2.2 Measurement system

To simulate the micro-motion of BM, we designed a dualstructure motor system. composed of two inverter motors, one inside the missile and one outside it (Figs. 2, 3). The motor inside the missile is impart the spinning motion of the missile and the second motor imparts coning motion (Figs. 2a, 3a). The coning angle can be adjusted by changing the flange that connects the first and the second motors (Fig. 2b). To simulate the effect of various coning angles, the flange was designed to have tilt angles ranging from 0° to 15° in increments of 1°. Each motor can be operated up to 500 RPM by adjusting a knob in the control box (Fig. 3b) that controls the current flow in the inverter. Therefore, combined use of the two motors and the flanges can yield micro-motions with various spinning and coning velocities at various coning angles.

For the measured target, three scale models were designed. Each was made of aluminum with 3-mm thickness to represent the perfect inducting characteristic of the missile while minimizing the weight (Fig. 4). In addition, each target has a unique size of warhead and body, and number of wings. Model 1 (Fig. 4a) was manufactured to study the general MD signature of BM. Models 2 (Fig. 4b) and 3 (Fig. 4c) were to study the effect of spinning in general missiles such as cruise missiles. The whole system was supported by a steel frame that can be assembled. The second motor was covered with



(a) Micro-motion device



(b) Control box

Fig. 3. Manufactured micro-motion device.

radar-absorbing material to eliminate any effects of motor rotation (Fig. 5).

A portable X-band 2-by-2 radar system was designed to obtain the radar signal (Fig. 6, Table 1). The transmitter emits an up-chirp signal in which the frequency increases linearly. The reflected electromagnetic signal is sequentially received by the four-channel antenna receivers (Fig. 6a). The four signals are synchronized, then joined as a single complex signal by using a single-pole four-through (SP4T) switch. The trigger signal for synchronization between analogue to digital converter (ADC) and arbitrary waveform generator (AWG) is generalized by an Atmega 128 micro-processor. This processor also forms antenna selection signals for the













control of the SP4T switch. The micro-processor is controlled by a host computer and an RS-232 (Fig. 6b).

#### 2.3 Signal processing procedure

To demonstrate the operation of the designed measurement system, 4,096 pulses were received from a missile-shaped target with pulse repetition frequency = 1,000 Hz, and 4096 RPs were obtained by performing matched-



Fig. 5. Measurement setup. To eliminate the effect of micro-motion of the second motor, it is covered with radar-absorbing material.









filtering using the stored replica of the transmitted signal in (1). Then, the range bin containing the target, i.e., bt, was selected by finding the range bin with the maximum energy of the signal projected onto the fast time domain, and the signal st in this bin was obtained. Because the noise and clutter can distort the MD signal,  $s_t$  was digitally band-passed using a rectangular bandpass filter (BPF) in the frequency domain. Finally, (7) was used to transform the band-passed signal  $s_{bpf}$  to the timefrequency domain, and residual noise was removed by thresholding (Fig. 7).

#### Table 1. Specification of the radar

#### **3. Experimental Results**

#### 3.1 Experimental result of system verification

An experiment was performed to verify the overall operation of the designed system by using Model 2 engaged in micro-motion with arbitrary parameters in a wide-open hall of the LG building of Pohang University of Science Technology. After matched filtering, 4096 RPs were obtained with the target located at range bin  $b_t = 250$  (Fig. 8a). The radar signal  $s_t$  varied over time due to the micro-motion of



Fig. 8. Measurement result.

the target (Fig. 8b). The one-dimensional (1D) frequency spectrum obtained by Fourier transformation shows that most of the MD signal falls in the frequency range -330  $\leq f \leq$  330 Hz with the largest energy at 0 Hz (Fig. 8c). Thus, frequency components outside this range and the zero frequency were eliminated digitally (Fig. 8d).

In forming the spectrogram of  $s_{bpf}$ , for the sake of clarity the pixel with the largest value was set to 0 dB and those smaller than -10 dB were set to -10 dB. The spectrogram (Fig. 9) of the MD signal after time-frequency transform clearly represents the micro-motion of the target. Due to the coning motion of the main body composed of a cylinder and a cone, sinusoidal MD motion occurred at  $-50 \le f \le 50$  Hz. Because of the timevarying nature of the radar cross-section (RCS) of the tilted cone and the cylinder, the amplitude of coning motion in spectrogram was time-varying. Furthermore, micro-motion caused by a combination of coning and spinning of the three wings clearly appeared at  $-200 \le f \le 300$  Hz and  $200 \le f \le 300$ Hz. Close examination of these ranges shows that MD of the three wings are represented by two or three vertical lines that arise periodically due to the sharp variation in the RCS of a plate.

#### 3.2 Analysis of MD spectrogram

Effects of motion parameters were analyzed to demonstrate the accuracy of the designed system. The angular velocity  $\omega_s$  of spinning was set to 10.5 rad/s = 100 RPM or 15.7 rad/s = 150 RPM, and the angular velocity  $\omega_c$  of coning was set in the same manner. Two coning angles  $\theta_c = 7$  and 15° were used to study the effect of coning angle. Aspect angle  $\theta_a$  was varied between 0 and 90° to determine its effect on MD (Table 2).



Fig. 9. MD image after BPF

Table 2. Motion and observation parameters

Motion	Value
Spinning motion (rad/s)	10.5, 15.7
Coning motion (rad/s)	10.5, 15.7
Coning angle	7°, 15°
Aspect angle	0° - 90°

MD images (Fig. 10) of Model 1 rotating at  $\omega_c = 15.7$  rad/s with  $\theta_c = 15^\circ$  and observed at  $\theta_a = 90^\circ$  clearly represent the coning motion. MD of the body repeats every 0.4 s which corresponds to the coning frequency  $f_c = 2.5$  Hz. The range of MD is different from the ideal values due to the small RCS of the nose; the ideal maximum Doppler frequency  $f_{max}$  of the nose tip drawing a circle with a radius = 0.7tan(15°) = 0.19 m is 196.4 Hz and this is not seen in the MD image, whereas  $f_{max} = 100$  Hz that corresponds to a 0.36-cm cylinder is seen.

MD images (Fig. 11) of Model 2 rotating at  $\omega_c = \omega_s = 15.7$  rad/s with  $\theta_c = 15^\circ$  represents the addition of spinning and coning frequencies. As proved in (11), the spinning MD of the three wings at  $\theta_a = 90^\circ$  was added to and subtracted from the coning MD, yielding  $f_{\text{max}} = 426.3$  Hz, which is very close to



Fig. 10. MD image of model 1 ( $\omega_c = 15.7 \text{ rad/s}, \theta_c = 15^\circ, \theta_a = 90^\circ$ ).



Fig. 11. MD image of model 2 ( $\omega_s = \omega_c = 15.7 \text{ rad/s}, \theta_c = 15^\circ$ ).

the ideal value (Fig. 11a). In addition, four wings are clearly represented by four vertical lines that repeat at  $f_c = 2.5$  Hz. Comparison of MD images obtained at  $\theta_a = 90^\circ$  (Fig. 11a) and  $60^\circ$  (Fig. 11b) shows the angular dependency; ideal  $f_{max} = 366.6$  Hz because the velocity was scaled by  $\cos(30^\circ) = 0.86$ , but measured  $f_{max} = 230$  Hz due to the small RCS of the wing tip.

MD images (Fig. 12) of Model 3 obtained at  $\theta_a = 90^\circ$  and  $\theta_c = 15^\circ$  for four combinations of  $(\omega_s, \omega_c)$  demonstrate the effect of each component. Comparison of MD images for (10.5, 10.5) Hz and (10.5, 15.7) Hz proves that fc determines the repetition of the MD image of the target (Figs. 12a, b); MD in Fig. 12a repeats at  $f_c = 1.67$  Hz whereas that in Fig. 12b repeats at  $f_c = 2.5$  Hz. Comparison of MD images for (10.5, 10.5) Hz and (15.7, 10.5) Hz shows that the spinning MD adds the additional frequency (Figs. 12a, c) of approximately 50 Hz

due to the increased spinning frequency. Comparison of Fig. 12b and d shows the same result.

Finally, the effect of the coning angle was studied by comparing MD images of Model 3 with  $\theta_c = 7^{\circ}$  and 15°, observed at  $\omega_s = \omega_c = 15.7 \text{ rad/s}$ ,  $\theta_a = 90^{\circ}$  (Fig. 13). Although the MD repeated with the same period, MD image was very dependent on  $\theta_c$ ; the amplitude of the main body was much larger at  $\theta_c = 15^{\circ}$  than at  $\theta_c = 7^{\circ}$  (Fig. 13). In addition, the location of the MD of each blade changed due to RCS variation of the three wings, yielding totally different MD images.

## 4. Conclusion

In this paper, the dynamic MD signatures of BM were



Fig. 12. MD image of model 3 for  $(\omega_s, \omega_c)$   $(\theta_c = 15^\circ, \theta_a = 85^\circ)$ .



Fig. 13. MD image of model 3 for two  $\theta_c$  values ( $\omega_s = \omega_c = 15.7 \text{ rad/s}, \theta_a = 90^\circ$ ).

investigated using a portable micro-motion device, three scale models, and a portable X-band measurement system. The micro-motion device implemented spinning and coning motions of the scale models, and the designed radar system composed of the dual motor system and the X-band 2-by-2 PCMP radar successfully measured the MD signal. MD signals were represented as spectrograms after matched-filtering, selection of  $b_v$  and BPF. Analysis of MD signature showed that the coning motion determined the overall period of MD, and that the spinning motion increased its amplitude. MD was very dependent on the aspect angle because of the angular dependency of RCS of the models.

MD can occur in various types of targets engaged in micromotion and these targets may exist in various environments that can produce various clutter signals. Because the designed system can be moved to any environments and many of such micro-motions can be implemented by the designed motion-device, it will contribute to successful analysis of the MD signal in various situations

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