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Lin, Shengmao; Lampi, Marsha C.; Reinhart-King, Cynthia A.; Tsui, Gary C.P.; Wang, Jian; Nelson, Carl A.; and Gu, Linxia, "Eigenstrain as a mechanical set-point of cells" (2018). *Mechanical & Materials Engineering Faculty Publications*. 291. https://digitalcommons.unl.edu/mechengfacpub/291

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Published in *Biomechanics and Modeling in Mechanobiology* 17 (2018), 951–959. doi 10.1007/s10237-018-1004-0 Copyright © 2018 Springer-Verlag GmbH Germany, part of Springer Nature 2018. Used by permission. Submitted 3 August 2017; accepted 23 January 2018; published online 5 February 2018.

Eigenstrain as a mechanical set-point of cells

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Abstract

Cell contraction regulates how cells sense their mechanical environment. We sought to identify the set-point of cell contraction, also referred to as tensional homeostasis. In this work, bovine aortic endothelial cells (BAECs), cultured on substrates with different stiffness, were characterized using traction force microscopy (TFM). Numerical models were developed to provide insights into the mechanics of cell–substrate interactions. Cell contraction was modeled as eigenstrain which could induce isometric cell contraction without external forces. The predicted traction stresses matched well with TFM measurements. Furthermore, our numerical model provided cell stress and displacement maps for inspecting the fundamental regulating mechanism of cell mechanosensing. We showed that cell spread area, traction force on a substrate, as well as the average stress of a cell were increased in response to a stiffer substrate. However, the cell average strain, which is cell type-specific, was kept at the same level regardless of the substrate stiffness. This indicated that the cell average strain is the tensional homeostasis that each type of cell tries to maintain. Furthermore,

cell contraction in terms of eigenstrain was found to be the same for both BAECs and fibroblast cells in different mechanical environments. This implied a potential mechanical set-point across different cell types. Our results suggest that additional measurements of contractility might be useful for monitoring cell mechanosensing as well as dynamic remodeling of the extracellular matrix (ECM). This work could help to advance the understanding of the cell-ECM relationship, leading to better regenerative strategies.

Keywords: Eigenstrain, Tensional homoeostasis, Steady state, Cell contraction, TFM, Mechanosensing

1 Introduction

The extracellular matrix (ECM) is essential for regulating cell behavior and tissue function (Discher et al. 2005; Oria et al. 2017). Local ECM structure and mechanics are increasingly recognized as important mechanical effectors of cell responses and tissue regeneration (Yu et al. 2011; Lin 2017). This is illustrated by the fact that both ECM rigidity (Hadjipanayi et al. 2009) and local tension (Beloussov et al. 2000) regulate cellular mechanotransduction pathways, and their dysregulation results in many different types of disease (Rittweger et al. 2009). Following mechanical disturbance of the ECM, cells can alter their contractile force and morphology including cell spread area, focal adhesion quantity and strength, and cytoskeleton structure (Califano and Reinhart-King 2010; Prager-Khoutorsky et al. 2011; Solon et al. 2007; Kumar et al. 2016). It has been speculated that cell contraction, generated by the cross-bridging interaction of actin and myosin II motors, maintains tensional homeostasis, i.e., a stable intracellular state, also referred as the set-point of cells (Paszek et al. 2005).

Various physical parameters, including tensional force on cells and cell stiffness, were proposed as the index factors of cellular tensional homeostasis (Brown et al. 1998; Mizutani et al. 2004). The early experiments illustrated that external loads led to an instantaneous alternation of cell behaviors including tensional force and stiffness. After unloading, cells returned to their original mechanical state. This indicated that cell stiffness and tensional force correlated with the magnitude of mechanical stimuli, instead of its loading history. However, the connection between different levels of mechanical stimuli, including different ECM stiffnesses, were not identified, which is the basis for identifying the set-point of cells. Recent studies with traction force microscopy (TFM) suggest that cell geometry or strain energy could be a better index of tensional homeostasis (Califano and Reinhart-King 2010; Munevar et al. 2001; Oakes et al. 2014; Tan et al. 2003; Wang et al. 2002; Rape et al. 2011). However, integrating a variety of discrete experimental data to inspect the fundamental regulating factors of cells remains a challenge.

Virtual testing of cell tractions using numerical or theoretical models could be a potential tool for data integration (Lemmon and Romer 2010; Oers et al. 2014; Vermolen and Gefen 2015) and has demonstrated efficiency to further interpolate TFM measurements. Zielinski et al. (2013) simulated two-dimensional (2D) and three-dimensional (3D) cell-ECM interactions to quantify the sensitivity of traction force to parameters including cell stiffness, focal adhesion density, cell aspect ratio, and contractility. The cell contraction was modeled as the inward pressure. The cell stiffness and focal adhesion density were identified as two major parameters affecting the traction forces. Oakes et al. (2014) have proposed a physical model for adherent cells, considering contraction as a uniform inward pressure combined with a uniform boundary line tension. They observed that cell shape influenced the traction force distribution, and that cellular strain energy density exhibited minimal alternations in response to different substrates. Edwards and Schwarz (2011) have modeled cell contraction as thermal elasticity, in which temperature alternations induced cell contraction. It was then shown that traction forces concentrated around the cell periphery, especially along the stiffer direction of ECM. However, there is little data to quantify cellular tensional homeostasis.

To identify key index factors for maintaining cellular tensional homeostasis, a numerical model was developed in this work to mimic TFM measurements and capture the detailed cell–ECM interactions. Contraction was modeled as eigenstrain, which could induce isometric cell contraction without external forces and better match with cell physiology. We hypothesized that a cell strives to maintain its pre-set eigenstrain in response to mechanical disturbance. Numerical models were used to integrate the published experimental observations in 3T3 fibroblasts (Oakes et al. 2014) as well as our TFM results for bovine aortic endothelial cells (BAECs). The predicted cell mechanics were used to test our hypothesis. This work could help to advance the understanding of cell–ECM interactions, specifically cell mechanosensing, and may lead to better regenerative strategies.

2 Materials and methods

2.1 TFM measurements of BAECs

BAECs were maintained at 37 °C and 5% CO₂ in Medium 199 (Invitrogen) supplemented with 10% FetalClone III (HyClone), and 1% each of penicillin-streptomycin (Invitrogen), MEM amino acids (Invitrogen) and MEM vitamins (Invitrogen). Variably compliant polyacrylamide (PA) gels coated with 0.1 mg/mL rat tail type I collagen (BD Biosciences) were prepared by altering the ratio of acrylamide to bis-acrylamide (BioRad) in the polymerization solution as described previously (Califano and Reinhart-King 2008; Wang and Pelham 1998). Gels were synthesized with Young's moduli of 1 and 10 kPa to mimic physiologically relevant tissue stiffness (Engler et al. 2004). Cells were seeded on PA gels embedded with 0.5µm fluorescent beads (Invitrogen) and allowed to attach and spread overnight. Bead fields were imaged using a Zeiss Axio Observer.Z1m microscope with a Hamamatsu ORCA-ER camera using a 20× objective before (stressed configuration) and after removal (relaxed configuration) of the cells with trypsin (Invitrogen). These images allowed substrate displacements to be tracked with correlation-based optical flow (Marganski et al. 2003). Substrate displacements were then translated into a strain field that was used to compute the traction strain using Bayesian statistics that maximized the most likely traction field describing a given strain field. The substrate strains were converted to traction stresses using the LIBTRC analysis library developed by Professor Micah Dembo of Boston University (Dembo and Wang 1999).

2.2 Numerical modeling

Numerical models were developed to mimic the TFM measurements of both BAECs and 3T3 fibroblasts cultured on substrates of various stiffnesses (**Fig. 1a**). For shape-controlled 3T3 fibroblasts (Oakes et al. 2014), the round cell had a diameter of 30 μ m and a height of 5 μ m. The substrate was assumed as a cuboid (75 × 75 × 5 μ m) with a fixed bottom surface. A total of 205 focal adhesions were generated to connect the cellular bottom with the top surface of the substrate, excluding the nucleus area which is considered 10 μ m in diameter (Milo and Phillips 2015). The focal adhesions were modeled as truss



Fig. 1. a) Three-dimensional model of cell–substrate interactions; **b)** adherent cell stiffness in response to varied substrate stiffness.

elements with a length of 1 µm (Franz and Müller 2005) and a stiffness of 20 nN/µm. The cells were modeled as an isotropic elastic material with variable stiffness in response to substrate stiffness, which was adopted from measurements by Solon et al. (2007), as depicted in Fig. 1b. Specifically, cell stiffness increased with a stiffer substrate and reached a plateau as the substrate stiffness exceeded 16 kPa. The constitutive model for linear isotropic materials was represented by the generalized Hooke's law that relates the Cauchy stress tensor (σ_{ij}) and the infinitesimal strain tensor (ε_{ij}) as

$$\varepsilon_{ij} = \frac{1+v}{E} \sigma_{ij} - \frac{v}{E} \sigma_{kk} \delta_{ij}$$
(1)

where δ_{ij} is the Kronecker delta, *E* is Young's modulus and *v* is Poisson's ratio.

For BAECs, the cell geometry was reconstructed from corresponding TFM images in ImageJ software (NIH, Bethesda, MD). The cell stiffness was reverse-fitted as 1.5 and 3 kPa in response to substrate stiffness of 1 and 10kPa, respectively, in order to match the traction stress field measured by TFM.

Cell contractility was modeled as eigenstrain, which was imposed to induce isometric contraction as the initial condition (Edwards and Schwarz 2011; Pan and Zhong 2016). The eigenstrain was defined as any nonelastic strains in the material without any external forces, such as strains caused by thermal expansion mismatch or phase transformation, plastic strains, etc. (Jun and Korsunsky 2010). In this work, the eigenstrain was the contraction strain of a cell. If a cell is free to contract, i.e., an isolated cell without constraints from the environment, this leads to zero stress in the cell. In reality, each cell is constrained by its local mechanical environment, and this leads to tensional stress within the cell as well as traction stress exerted against the ECM. In this work, a -20% eigenstrain was imposed on the cell, which was reverse-fitted from the published experimental data (Oakes et al. 2014) to match the measured traction stresses. The strain energy on the substrate was calculated (Butler et al. 2002) as

$$W = \frac{1}{2} \int \mathbf{T}(\mathbf{r}) \cdot \mathbf{u}_{s}(\mathbf{r}) dA$$
 (2)

where $\mathbf{T}(\mathbf{r})$ is the cell contraction-induced traction stress on the substrate, $\mathbf{u}_{c}(\mathbf{r})$ is displacement, and A is the targeted area.

2.3 BAEC cytoskeletal visualization

BAECs were allowed to adhere and spread overnight on 1 and 10 kPa polyacrylamide gels. Cells were then fixed with 3.7% formaldehyde (VWR) and permeabilized using 1% triton (VWR), followed by blocking with 1% bovine serum albumin (Sigma) and incubation with Alexa Fluor 488 Phalloidin (Life Technologies) to visualize F-actin. Cells were imaged using a Zeiss Axio Observer.Z1m microscope with a Hamamatsu ORCA-ER camera using a 40× oil immersion objective.

3 Results

Our model was used to capture the published experimental results for 3T3 fibroblasts (Oakes et al. 2014). Briefly, the round cell shape was controlled to decouple it from the effect of substrate stiffness on cell geometry. The substrate mechanics were then characterized in terms of traction stress (kPa) and displacement (magnitude, μ m), as well as the average strain energy (pJ). The corresponding modeling results for substrate stiffnesses ranging from 2.8 to 30 kPa are shown in **Fig. 2.** Using an eigenstrain-induced cell contraction, a good match was observed between our model predictions and the published experimental and numerical data (Oakes et al. 2014). Moreover, our model prediction demonstrated a better fit than the existing numerical model (Oakes et al. 2014). Increased displacement and traction stress were found around cell edges. The substrate traction



Fig. 2. Cell-induced substrate mechanics depend on substrate stiffness. **a)** Substrate displacement and traction stressmaps obtained from numerical models; **b–d)** average substrate displacement, traction stress and strain energy are predicted by numerical models and compared to the published experimental data. Reproduced with permission from Oakes et al. (2014).

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stress and displacement were inversely related to an increase in substrate stiffness. Specifically, the same eigenstrain contraction of cells led to a larger displacement with lower stresses on soft substrates, compared to stiff ones (Fig. 2a). Furthermore, we quantified the average displacement, average traction stress, and strain energy exerted on the substrate (Fig. 2b–d). A substrate stiffness below 10 kPa led to rapid changes in both substrate traction stress and displacement. This trend plateaued on stiffer substrates. Interestingly, the effect of substrate stiffness on strain energy showed an oscillating pattern. This could be explained by the rapid shift between the average displacements and traction stresses.

Cell mechanics under varied substrate stiffness were further investigated (**Fig. 3**). In response to increased substrate stiffness, the average cell stress followed the same trend as the substrate traction stress depicted in Fig. 2c. Notably, the average maximum principal strain of the cell had minimal sensitivity to the substrate stiffness. As the substrate stiffness varied between 2.8 and 30 kPa, the average maximum principal strain of the cell was reduced by 7.8% (from 14.1 to 13.0%), compared to a 12.5% reduction assuming a constant cell stiffness independent of substrate rigidity. This indicated that dynamic tuning of cell stiffness in response to the mechanical environment is used to maintain a cell average strain.

Our hypothesis based on eigenstrain as the cellular set-point was also tested using BAECs with uncontrolled morphologies. Representative substrate traction stress maps (**Fig. 4**) demonstrated that peak



Fig. 3. Substrate stiffness altered the average stress in a single cell, while its average maximal principal strain remained constant.



Fig. 4. Representative traction stress map (kPa) for BAECs cultured on the substrate with a stiffness of 1 kPa (upper) and 10 kPa (lower), respectively.

stress is located at cell protrusions and is distant from the central nucleus area. Similar stress patterns were predicted from our computational models, which were reconstructed from TFM images of individual cells assuming a -20% eigenstrain (**Fig. 5a**). However, the peak substrate stress obtained from our simulations was generally less than what was measured using TFM. As the substrate stiffness changed



Fig. 5. Representative model predictions for BAECs cultured on substrates with a stiffness of 1 kPa (left) and 10 kPa (right): **a)** substrate traction stress map (kPa), **b)** substrate displacement map (μ m), **c)** cell displacement map (μ m), **d)** cell stress map (kPa).

from 1 to 10 kPa, peak substrate stresses of 0.31 and 0.57 kPa, respectively, were obtained from our simulations, compared to 0.52 and 2 kPa measured using TFM. Our computational model allowed us to probe the displacement or stress virtually anywhere within the computational domain, a feature that was not affected by experimental techniques. Specifically, Fig. 5b illustrates the cell contraction-induced substrate displacement map using our model. Figure 5c, d depicts the cellular response in terms of cell stress and displacement. It is clear that a stiffer substrate leads to higher stresses as well as decreased cell and substrate displacements (Fig. 5). The simulated integrin attachments between cells and their substrate resulted in larger cellular stresses and displacements than the substrate stresses and displacements. Table 1 summarizes both experimental results and the computational predictions. The predicted average substrate stress is within the range of the experimental measurements. It is worth noting that the change in the average cell strain was minimal compared to other parameters, such as stress, displacement, stiffness, or spread area. By comparing the cell average strain magnitudes between BAECs and fibroblast cells (Fig. 3), it is interesting to observe that the cell average strain is different for both cell types, but the cell contraction in terms of eigenstrain is the same.

4 Discussion

In this work, cell–ECM interactions were examined through combined TFM measurements and numerical models to quantify the tensional homeostasis of cells. The working hypothesis is that cellular behaviors on various substrates are regulated by the contraction strain of cells.

Substrate stiffness (kPa)	1	10	
Cell stiffness (kPa)	1.5	3	
Cell spread area (TFM) (µm ²)	1600 ± 148	1693.7 ± 283	
Average stress on substrate (TFM) (Pa)	81.3 ± 20.0	189.6 ± 23.7	
Average stress on substrate (Pa)	92.0 ± 5.5	208.9 ± 21.8	
Cell average stress (Pa)	228 ± 9	559 ± 23	
Cell average strain	5% ± 1.1%	3.73 ± 0.2%	

Table 1. Characteristics of bovine aortic endothelial cells from both TFM measurements (labeled) and numerical models.

A –20% eigenstrain was implemented to mimic the cell contraction on various substrates. This technique was validated in previous literature (Edwards and Schwarz 2011). Our TFM measurements of BAECs, as well as published TFM observations in 3T3 fibroblasts (Oakes et al. 2014), were used to validate the numerical model. Our results indicate that eigenstrain, an efficient set-point for cell contraction, regulates cell behaviors in response to various mechanical stimuli.

The motivation to determine the set-point of a cell went back to the work by Brown et al. (1998). They quantified how fibroblasts, cultured within three-dimensional collagen lattices, responded to predefined static and dynamic tensional loadings on the gel. The data indicated that increased loading led to an acute reduction in cell contractile force, and vice versa. However, after unloading, the gel embedded with living fibroblast cells maintained an equilibrium tensional force of approximately 40-60 dynes/million cells. This indicated that tensional force correlated with the magnitude of the external stimuli, instead of its loading history. At each loading level, the tensional force was the same. Reflected in our eigenstrain model, different loading magnitudes correspond to the different substrate stiffnesses. The tensional force at discrete substrate stiffness remained at the same level. Mizutani et al. (2004) investigated the same problem from the perspective of altered cell stiffness. They observed that external substrate stretch caused instantaneous cell stiffness alterations and fibroblasts returned to their initial cell stiffness after unloading. These data demonstrated that mechanical stimuli induced acute changes in cell stiffness that correlated with the magnitude of the applied stimuli. In addition, this indicated that our numerical model should consider the stiffness adaptation of cells to different substrates. Overall these observations illustrated that the mechanical behaviors of cells are correlated with the magnitude of mechanical stimuli they experience, instead of the loading history. However, the relationship between different levels of mechanical stimuli, including different substrate or gel stiffnesses, was not identified, which is the basis we used to identify the set-point of cells, i.e., tensional homeostasis.

Our eigenstrain model captured TFM results of 3T3 fibroblasts in terms of substrate displacement and traction stress maps. It is interesting to note that our model prediction of average substrate displacement, traction stress, and strain energy has demonstrated a better match than the published numerical model. In the literature (Oakes

et al. 2014), cell contraction was simulated as a contractile pressure of 2.4 kPa along with a uniform membrane tension of 0.7 nN/ μ m, compared to our model implementation using an eigenstrain of -20%. In addition, cell stiffness in the published model was assumed to be constant, whereas we adopted the measured acute cell stiffness, which varied in response to different substrate stiffnesses (Solon et al. 2007). The stiffness adaptation of cells correlated well with the observations that cells reorganize their actin cytoskeleton into bundles or stress fibers on stiff substrates (Solon et al. 2007; Yeung et al. 2005; Parker et al. 2002; Théry et al. 2006). The above-mentioned differences in numerical models led to different conclusions about the regulatory mechanisms of cells. For example, Oakes et al. (2014) found that the strain energy of a cell was correlated with the cell spread area, and the local curvature regulated the traction stress distributions on the substrate. Our main conclusion is that cells exhibit a constant strain regardless of altered substrate deformation or traction stress distributions in response to varied substrate stiffness. The eigenstrain and stiffness adaptations were two key factors that improved the performance of our model.

TFM measurements of BAECs demonstrated that traction stresses and cell area vary in response to substrate stiffness. The predicted traction stress distributions with higher stress located at cell protrusions (Fig. 5) agree with our TFM data as well as prior published TFM results (Edwards and Schwarz 2011). Our validated model allowed us to further investigate the cellular responses in different mechanical environments. Our model indicated that cells exhibit larger stresses and smaller displacements when cultured on a stiffer matrix, and followed the same trend as the TEM data when substrate mechanics were varied. But the predicted cellular displacements and stresses were found to be larger than the substrate displacements and stresses. This could be explained by the existence of focal adhesions. The peak substrate stress obtained from our simulations was relatively less than what was measured using TFM. This could be attributed to subcellular heterogeneity, as evidenced by the actin cytoskeleton from our BAEC fluorescent images (Fig. 6). BAECs reorganize their cytoskeleton in response to substrate stiffness, indicating alterations in cell stiffness. In our numerical model, we considered bulk cell stiffening but not the inhomogeneity within the cell. The homogeneous cell simplification will alter the traction stress localization, but will not impact our understanding



Fig. 6. Representative cytoskeletal maps of BAECs cultured on the substrate with a stiffness of 1 and 10 kPa.

of the tensional homeostasis of a cell. For unconstrained BAECs, we demonstrated that the cell average strain was constant regardless of cell morphology (Table 1). Moreover, the cell average strain in BAECs was less than half of the strain in fibroblast cells. This indicated that the cell average strain is cell type-specific. However, it is worth noting that the eigenstrains on both cells are the same. This suggests that different cell types have the same set-point level in terms of cell contraction. Further experimental evidence for additional cell types is needed to validate this observation.

The practical implication of this work is the identification of cellular strain as the tensional homeostasis state that each cell type maintains following mechanical disturbances, including alteration of substrate stiffness. Understanding the tensional homeostasis of cells is an important consideration in cellular biophysics and could provide a starting point for investigating many biological processes such as cell growth, differentiation, and migration. We found that cellular strain was independent of a cell's mechanical environment, and speculate that eigenstrain may be a cellular set-point that is conserved across cell types. This needs to be further tested through experimental measurements that include the cell morphology and the mechanics of focal adhesions. Measurements of contractility may be useful for monitoring cell mechanosensing as well as dynamic ECM remodeling. This work has the potential to elucidate why cells change their area and morphologies in response to substrate stiffness. Our mechanical model, based on the cell shape with the assumption of constant focal adhesion density, is a significant first step in quantifying the

tensional homeostasis of cells. The dynamical behaviors of focal adhesions might play a role in the cellular tensional hemostasis, which could be incorporated in our future work. The relevant cellular remodeling processes can also be addressed in future studies by incorporating the dynamic reorganization of cytoskeletal architecture and focal adhesion dynamics that occur when cells interact with their ECM (Discher et al. 2005; Kumar et al. 2016; Humphrey et al. 2014). In summary, our finding that eigenstrain is consistent could help to advance the understanding of the cell-ECM relationship, specifically with respect to cell mechanosensing, and may lead to better regenerative strategies (Chawla 2012; Zhong et al. 2007).

Acknowledgments — This work was supported in part by the National Science Foundation CAREER award (CBET-1254095) and the Foundation of Fujian Educational Committee (JAT170422). The authors also thank Mr. Hozhabr Mozafari for constructive comments.

Conflict of interest — The authors have no conflicts of interest.

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