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ELAS - A General-Purpose Computer Program for the Equilibrium Problems of Linear Structures

Volume II. Documentation of the Program

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## Preface

The work described in this report was performed by the Engineering Mechanics Division of the Jet Propulsion Laboratory.

The ELAS program was developed by Dr. Senol Utku and Dr. Fevzican A. Akyuz, and is dedicated to the memory of Professor M. Inan of the Technical University of Istanbul.

## Acknowledgment

The author is indebted to Vivia Crew for her help in editing all documents related with the ELAS program.

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#### Abstract

A general-purpose digital computer program (named ELAS) for the in-core solution of linear equilibrium problems of structural mechanics is described for potential and actual users in Volume I of this report, and is documented in Volume II. The program requires minimum input for the description of the problem. The solution is obtained by means of the displacement method and the finite element technique. Almost any geometry and structure may be handled because of the availability of lineal, triangular, quadrilateral, tetrahedral, hexahedral, conical, triangular torus, and quadrilateral torus elements. The assumption of piecewise linear deflection distribution insures monotonic convergence of the deflections from the stiffer side with decreasing mesh size. The stresses are provided by the best-fit strain tensors in the least-squares sense at the mesh points where the deflections are given. The selection of local coordinate systems whenever necessary is automatic. The core memory is efficiently used by means of dynamic memory allocation, an optional mesh-point relabelling scheme, imposition of the boundary conditions during the assembly time, and the straight-line storage of the rows of the stiffness matrix within variable bandwidth and the main diagonal. The number of unsuppressed degrees of freedom that can be handled in a given problem is 500 to 600 for a typical structure, but might far exceed these average values for special types of problems; the execution time of such problems is about four minutes in 32 K IBM 7094 Model I machines. The program is written in FORTRAN II language. The source deck consists of about 8000 cards and the object deck contains about 1400 binary cards. The physical program (standard ELAS) is available from COSMIC, the agency for the distribution of NASA computer programs.


## I. Introduction

Volume I, User's Manual, of this report gives a general description of ELAS,* a general-purpose digital computer program for the in-core solution of linear equilibrium problems of structural mechanics, and contains the information necessary for input preparation, arrangement of the physical program, and interpretation of output and error messages.

Volume II, Documentation of the Program, is published in two parts: the present volume-the basic Volume IIwhich gives the theoretical background of the program and contains tables and figures describing the COMMON variables, their meanings, and their arrangement in COMMON; and an Addendum to Volume II, which
*First two syllables of the word Elasticity.
contains program descriptions, flowcharts, and source program listings for all program elements of ELAS/Level 3. (The original version of the ELAS program made available from COSMIC ${ }^{* *}$ in April 1968 is designated ELAS/Level 0. Subsequent program corrections made in January 1969, March 1969, and May 1969 updated the program to ELAS/Level 1, ELAS/Level 2, and ELAS/ Level 3, respectively.)

In addition to the list of references cited in the text, a list of documented works related with the development of the ELAS program is given in the bibliography. A corrigenda for Volume I is given in the Appendix.

[^0]
## II. Theoretical Background

This section summarizes the mathematical formulation, the numerical method of solution, and the design features of the program.

## A. Mathematical Formulation

Let $V$ denote the material volume of the structure within the closed boundary S. Let $x_{\alpha}, \alpha=1,2,3$, denote a fixed right-handed Cartesian coordinate system. The Greek subscripts always refer to these axes; therefore, $\sigma_{\alpha \beta}$ is the stress tensor described in such a coordinate system. Let $u_{\alpha}$ denote the displacement vector; $\tilde{p}_{\alpha}$ the body force; $m$ the unit mass; double dots above, the second time derivative; and comma in the subscript the partial differentiation with respect to the space variable represented or implied by the subscript following the comma. If it is assumed that repeated subscripts imply summation over the range, the equilibrium of any particle within $V$ may be expressed as

$$
\begin{equation*}
\boldsymbol{\sigma}_{\beta \alpha_{\beta}}+\tilde{p}_{\alpha}=m \ddot{u}_{\alpha} \tag{1}
\end{equation*}
$$

In the equilibrium problems, the loading is such that

$$
\begin{equation*}
\ddot{\boldsymbol{u}}_{\alpha}=0 \tag{2}
\end{equation*}
$$

Therefore, substitution of $\ddot{u}_{\alpha}$ from Eq. (2) into Eq. (1) yields

$$
\begin{equation*}
\sigma_{\beta \alpha, \beta}+\tilde{p}_{\alpha}=0 \quad \text { in } V \tag{3}
\end{equation*}
$$

Let $S^{\prime}$ denote the portion of $S$ where the tractions are prescribed. The equilibrium condition on $S^{\prime}$ is

$$
\begin{equation*}
\sigma_{\alpha \beta} n_{\beta}+p_{\alpha}=0 \tag{4}
\end{equation*}
$$

where $n_{\beta}$ is the unit normal vector and $p_{\alpha}$ is the prescribed traction. The stress-strain relationship of the material is

$$
\begin{equation*}
\sigma_{\alpha \beta}=D_{\alpha \beta \gamma \delta}\left(\epsilon_{\gamma \delta}-\epsilon_{\gamma \delta}^{0}\right) \tag{5}
\end{equation*}
$$

where $\epsilon_{\gamma \delta}^{0}$ is the prescribed strain tensor, $\epsilon_{\gamma \delta}$ is the strain tensor, and $D_{\alpha \beta \gamma \delta}$ is the material matrix, which is positivedefinite and symmetrical, so that

$$
D_{\alpha \beta \gamma \delta}=D_{\gamma \delta \alpha \beta}=D_{\beta \alpha \gamma \delta}=D_{\alpha \beta \delta \gamma}
$$

The strain displacement relationships are

$$
\begin{equation*}
\epsilon_{\alpha \beta}=\frac{1}{2}\left(u_{\alpha, \beta}+u_{\beta, \alpha}\right) \tag{6}
\end{equation*}
$$

Finally, the displacement boundary conditions may be stated as

$$
\begin{equation*}
u_{\alpha}=u_{\alpha}^{0} \quad \text { on } S^{\prime \prime} \tag{7}
\end{equation*}
$$

where $u_{\alpha}^{0}$ denotes the prescribed displacements on $S^{\prime \prime}$. It should be noted that

$$
\begin{equation*}
S^{\prime}+S^{\prime \prime}=S \tag{8}
\end{equation*}
$$

In an equilibrium problem, usually $V, S^{\prime}, S^{\prime \prime}, p_{\alpha}, u_{\alpha^{\prime}}^{0} \tilde{p}_{\alpha}$, $D_{\alpha \beta \gamma \delta}, n_{\beta}$, and $\epsilon_{\gamma \delta}^{0}$ are given and $u_{\alpha}, \epsilon_{\alpha \beta}$, and $\sigma_{\alpha \beta}$ are requested.

Equations (3) through (8) constitute the differential equation formulation of the equilibrium problem in threedimensional continuum. A finite difference solution based on this formulation may be set up as follows. A regular mesh is placed in $V$ such that $S^{\prime}$ and $S^{\prime \prime}$ are determined by the mesh points. If $S$ is not defined by coordinate surfaces, such representation of $S^{\prime}$ and $S^{\prime \prime}$ is only approximate. The displacements $u_{\alpha}$ at the mesh points in $V$ and on $S$ are taken as the primary unknowns. The prescribed $u_{\alpha}^{0}$ displacements are assigned to the mesh points of $S^{\prime \prime}$. With the use of Eq. (6), $\epsilon_{\alpha \beta}$ at the mesh points in $V$ and on $S$ are approximated by the first differences of $u_{\alpha}$ and $u_{\alpha}^{0}$. Then, by the use of Eq. (5), the values of $\sigma_{\alpha \beta}$ are expressed at the mesh points. Finally, depending upon the mesh point in $V$ or on $S^{\prime}$, Eq. (3) or Eq. (4), respectively, is used to write the difference equations for the unknown displacements. After the unknown displacements from these equations have been computed, the strains and the stresses may be computed from the finite difference approximations of Eq. (6) and Eq. (5). Such a solution method has the following drawbacks:
(1) To minimize the truncation errors, a regular mesh in $V$ is used; however, this causes approximate representation of boundary $S$ and, therefore, increases the truncation errors in the finite difference approximations of Eqs. (4) and (7). Since the errors in the finite difference approximations of Eqs. (4) and (7) dominate in the solution more (Ref. 1) than the errors in the finite difference approximation of Eq. (3), either an irregular mesh in $V$ may be considered to represent $S$ more accurately, or higher-order formulas for the boundary conditions are used, although neither scheme is desirable in a general-purpose digital computer program.
(2) Because of the symmetry and the positive-definiteness of $D_{\alpha \beta \gamma \delta}$, the formulation given by Eqs. (3) through (8) is self-adjoint and positive-definite. However, the coefficient matrix of the unknown
displacements in the finite difference equilibrium equations is, in general, neither symmetric nor positive-definite. The loss of the two desirable qualities of the problem in the numerical formulation increases the storage requirements and solution time. Because of these setbacks, the mathematical formulation given by Eqs. (3) through (8) is modified slightly as explained in the following paragraph.

Consider the quantity $\pi$ defined as

$$
\begin{equation*}
\pi=\frac{1}{2} \int_{V} \epsilon_{\alpha \beta} \sigma_{\alpha \beta} d V-\int_{V} u_{\alpha} \tilde{p}_{\alpha} d V-\int_{S^{\prime}} u_{\alpha} p_{\alpha} d S \tag{9}
\end{equation*}
$$

where $d V$ is the volume element, $d S$ the area element, and the other symbols are as previously defined. Consider the displacement fields satisfying Eq. (7). For each such displacement field, by means of Eqs. (9), (6), and (5), a scalar $\pi$ may be computed. It can be shown that, for sufficiently smooth displacement fields satisfying Eq. (7), the stationary point of $\pi$, i.e., the point for which

$$
\begin{equation*}
\delta \pi=0 \tag{10}
\end{equation*}
$$

also satisfies Eqs. (3) and (4). In fact, by the methods of calculus of variations, Eq. (10) yields Eq. (3) as the Euler differential equation, and Eq. (4) as the additional boundary condition. Therefore, Eq. (10) is an equivalent statement of Eqs. (3) and (4). The quantity $\pi$ is known as the "total potential energy" of the system. Thus, the formulation given by Eq. (10) reduces the problem to that of locating the stationary point of the total potential energy functional. How the numerical solution is set up from this formulation (which is sometimes referred to as the extremum formulation of the problem), and its advantages, are discussed in the next subsection.

## B. Numerical Solution Method Based on the Extremum Formulation

A random mesh is placed in $V$ such that the mesh elements are line segments, triangles, quadrilaterals, tetrahedrons, hexahedrons, conical segments, or triangular or quadrilateral tori. Some of the mesh elements are shown in Fig. III-I (Vol. I). The types of mesh elements that may be used in different structures are given in Table III-2 (Vol. I). The randomness of the mesh enables the selection of mesh points that are exactly on the boundary S. For clarity, the mesh points are labelled sequentially, with integer numbers starting from 1 . If there are $s$ number of mesh points, there are $s$ ! different types of possible labelling. In the discussion that follows it will become obvious
that some of these labelling systems are more desirable than others.

It is assumed that one of the possible $s$ ! systems is selected. Next, the mesh elements are labelled sequentially with integers. If there are $\underline{m}$ number of mesh elements, there are $\underline{m}$ ! number of different labelling systems. It is supposed that one of the possible $\underline{m}!$ systems is selected. In what follows, superscript $m$ indicates the element label, and subscripts $t$ or $s$ indicate the mesh point label. To solve the equilibrium problem formulated in Section II-A numerically, instead of computing $u_{\alpha}$ at every point of $V$ and $S$, an attempt is made to find, at the mesh points, certain related quantities that define the distorted configuration of the structure in the same way as $u_{\alpha}$. These quantities are called deflections, which are displacements/rotations at the mesh points. Given a mesh point, the total number of independent deflection components is the number of degrees of freedom of that mesh point. In Table III-1 (Vol. I), the deflection components at a mesh point of different structures are shown as referred to an overall coordinate system. Let $\underline{k}$ denote the number of degrees of freedom at a mesh point of a structure. The value of $\underline{k}$ for different structures is given in the last column of Table III-1 (Vol. I). It will be assumed that the deflection components at a mesh point are ordered as shown in the table. The subscripts $k$ and $l$ will be used to indicate the sequence number implied by this ordering. If a prime appears on $k$ or $l$, this implies that a local coordinate system is used in defining the degree-of-freedom directions. Let $q_{k s}$ denote the $k$ th deflection component at mesh point $s$. A mesh element may be defined by the mesh points that are coincident with its vertices. For clarity in referencing, the convention of Table III-5 (Vol. I) is adopted in ordering the vertices of mesh elements. The type numbers shown in this table refer to the numbers shown on the shaded squares of Table III-2 (Vol. I). Subscripts $g$ and $h$ will be exclusively used to denote the sequence number of a vertex in the $g$ number of vertices of an element.

With the preceding definitions, the method used to obtain the stationary point of the total potential energy functional may now be explained. This is the classical Ritz procedure (Ref. 2), where the undetermined parameters of the problem are the unknown components of $q_{k s}$ deflections. Equation (9) is first written as

$$
\begin{equation*}
\pi=\frac{1}{2} \int_{V m} \epsilon_{\alpha \beta} \sigma_{\alpha \beta} d V^{m}-\int_{V^{m}} u_{\alpha} \tilde{p}_{\alpha} d V^{m}-\int_{S^{\prime} m} u_{\alpha} p_{\alpha} d S^{\prime m} \tag{11}
\end{equation*}
$$

where the repeated index implies the summation over the range. Next, an attempt is made to select a family of displacement fields that are sufficiently smooth, but otherwise arbitrary, ignoring for the time being the essential boundary conditions of Eq. (7). A piecewise linear displacement field is acceptable in this sense (Ref.3). Of course there are other piecewise continuous fields that are acceptably smooth. However, to simplify the understanding of the procedure, it is assumed that the displacement fields are piecewise linear. Such a field may be described mathematically for the $m$ th element in terms of the deflections of its vertices as

$$
\begin{equation*}
u_{\alpha^{\prime}} \approx B_{\alpha^{\prime} \beta^{\prime} l^{\prime} h} \widetilde{Q}_{l^{\prime} l} \mu_{h t}^{m} q_{l t} x_{\beta^{\prime}}+\text { rigid body movement } \tag{12}
\end{equation*}
$$

where the primes indicate the local coordinate system of the element. The coefficients $\mu_{h t}^{m}$ constitute a binary array such that, for a given $m$ and $h$, it is zero throughout the range of $t$, but 1 at the value of $t$ corresponding to the $h$ th vertex of the $m$ th element. In fact, $\mu_{n t}^{m} q_{l t}$ is the list of deflection components pertaining to the vertices of the $m$ th element. The matrix $\widetilde{Q}_{l, l}$ in Eq. (12) is the coordinate transformation matrix, where, for fixed $l^{\prime}$, it represents the direction cosines of the local axis related with $l^{\prime}$ degree-of-freedom direction in the coordinate system associated with $l$. The space variable $x_{\beta}$, is the distance measured along the $\beta^{\prime}$ th local axis. The coefficients $B_{\alpha^{\prime} \beta^{\prime} l^{\prime}{ }^{\prime} h}$ may be computed from the local coordinates of the vertices of the $m$ th element. In Table III-3 (Vol. I), for different mesh elements, the orientation of the local coordinate system relative to the overall coordinate system is given. It should be noted that Eq. (12) is an approximation of the true displacements in the $m$ th element, even if the exact values of deflection components $q_{l t}$ are known. However, it may be shown that the error decreases with decreasing mesh size. With the use of $u_{\alpha^{\prime}}$ from Eq. (12) in Eq. (6), the strains in the $m$ th element, as referred to the local coordinate system of the $m$ th element, may be expressed as

$$
\begin{equation*}
\boldsymbol{\epsilon}_{\alpha^{\prime} \beta}=B_{\alpha^{\prime}, l^{\prime} l^{\prime} h} \tilde{Q}_{l^{\prime} l} \mu_{h t}^{m} q_{l t} \tag{13}
\end{equation*}
$$

Let $D_{\alpha^{\prime} \beta^{\prime} \gamma^{\prime} \delta^{\prime}}$ denote the material constants of the $m$ th element as referred to the local coordinate system of the $m$ th element. If

$$
\begin{equation*}
K_{k g l h}^{m}=\tilde{Q}_{k^{\prime} k}\left(\int_{V^{m}} B_{\delta^{\prime} \gamma^{\prime} k^{\prime} g} D_{\delta^{\prime} \gamma^{\prime} \alpha^{\prime} \beta^{\prime}} B_{\alpha^{\prime} \beta^{\prime} l^{\prime} h} d V\right) \tilde{Q}_{l^{\prime} l} \tag{14}
\end{equation*}
$$

and

$$
\begin{equation*}
P_{k y}^{m}=\widetilde{Q}_{k^{\prime} k}\left(\int_{V^{\prime} m} B_{\alpha^{\prime} \beta^{\prime} k^{\prime} g} x_{\beta^{\prime}}, \tilde{p}_{\alpha^{\prime}} d V+\int_{S^{\prime} m} B_{\alpha^{\prime} \beta^{\prime} k^{\prime} g} x_{\beta^{\prime}} p_{\alpha^{\prime}} d \mathrm{~S}\right) \tag{15}
\end{equation*}
$$

are defined, $\pi$ of Eq. (11) may be expressed as

$$
\begin{equation*}
\pi=\frac{1}{2} q_{k s} \mu_{l g}^{m} K_{k g l h}^{m} \mu_{h t}^{m} q_{l t}-q_{k s}\left(\mu_{\mu_{g s}}^{m} p_{k g}^{m}+Q_{k s}\right) \tag{16}
\end{equation*}
$$

where $Q_{k s}$ denotes the prescribed concentrated loads at the mesh points. The deflection components $q_{k s}$ (or $q_{t t}$ ) should be such that, on $S^{\prime}$, they satisfy the essential boundary conditions of Eq. (7). Let $d_{j}$ denote the portion of $q_{l t}$, which is unknown. The essential boundary conditions may be expressed as

$$
\begin{equation*}
q_{l t}=e_{l t j} d_{j}+e_{l t}^{0} \tag{17}
\end{equation*}
$$

where the coefficients $e_{l t j}$ and $e_{i t}^{0}$ are quantities that may easily be determined from Eq. (7). For example, if there are no prescribed deflections in the problem, $e_{t t j}$ is a binary array containing only one 1 in the whole range of $j$ for a given $l t$, and $e_{l t}^{0}$ is zero throughout. Actually, the $d_{j}$ are the true undetermined parameters of the problem. If $q_{l t}$ is substituted from Eq. (17) into Eq. (16), the values of $d_{j}$ corresponding to the stationary point of $\pi$ may be obtained from the set of linear equations

$$
\begin{equation*}
\pi, d_{j}=0 \tag{18}
\end{equation*}
$$

since

$$
\begin{equation*}
\delta \pi=\pi_{, d_{j}} \delta d_{j} \tag{19}
\end{equation*}
$$

The equations given by Eq. (18) may be rewritten as

$$
\begin{equation*}
\ell \leftarrow_{i j} d_{j}=\mathscr{P}_{i} \tag{20}
\end{equation*}
$$

where

$$
\begin{equation*}
\partial \mathcal{X}_{i j}=e_{k s i} \mu_{g s}^{m} K_{k g l h}^{m} \mu_{h t}^{m} e_{l t j} \tag{21}
\end{equation*}
$$

and

$$
\begin{equation*}
\mathscr{P}_{i}=e_{k s i}\left(\mu_{g s}^{m} P_{k g}^{m}+Q_{k s}\right)-e_{k s i} \mu_{g s}^{m} K_{k g l h}^{m} \mu_{h t}^{m} e_{l t}^{0} \tag{22}
\end{equation*}
$$

The coefficients $K_{k g l h}^{m}$ constitute the element stiffness matrix of the $m$ th element and $P_{k g}^{m}$ is called the $m$ th element load vector. In Eq. (20), the coefficient matrix $\delta \mathcal{X}_{i j}$ is the stiffness matrix associated with the directions of $d_{j}$ deflections, and the right-hand-side vector $\mathscr{P}_{i}$ lists the loads in these directions. Equations (21) and (22) indicate how the coefficient matrix and the right-hand-side vector of the governing equations can be systematically generated from the element stiffness matrices and load vectors.

The operation implied by Eqs. (21) and (22) is referred to as the assembling of the elemental matrices. Because of the positive-definiteness and the symmetry of $D_{\alpha \beta \gamma \delta}$, Eq. (21) shows that

$$
\begin{equation*}
\mathscr{X}_{i j}=\mathscr{X}_{j i} \tag{23}
\end{equation*}
$$

and $\mathscr{K}_{i j}$ is also positive-definite. Once the unknown deflections $d_{j}$ are solved from Eq. (20), the complete deflections $q_{l t}$ are obtained by substituting $d_{j}$ into Eq. (17). After deflections $q_{k s}$ have been computed, the strains and the stresses at the mesh points may be computed as described in Ref. 4.

The method of solution described in the preceding has the following advantages:
(1) Since the mesh is random, the boundaries $S^{\prime}$ and $S^{\prime \prime}$ may be closely approximated, and thus minimize truncation errors.
(2) Any a priori knowledge about the variation of $\boldsymbol{u}_{\alpha}$ in $V$ or on $S$ may be used advantageously by varying the mesh size accordingly to minimize the truncation errors.
(3) The self-adjoint character, as well as the positivedefiniteness of the problem, is preserved, since $\varnothing \mathcal{L}_{i j}$ is always symmetric and positive definite.

## C. The Program

I. Criteria for Storage Allocation. All the input data previously mentioned (see Fig. IV-1, Vol. I) are stored permanently in COMMON after their validity is checked. No fixed-length block is assigned to these diversified data. The data are compactly stored as a string of variable length. This enables the program to compete in obtaining priorities efficiently in a multiprogramming environment. After reading the control card, the program determines the pointers of each of the data blocks relative to the beginning of COMMON. When the other input data become available, they are placed in the proper place in COMMON by the pointers. Although the locations of the data blocks vary from one job to another, the locations of the pointers and the control information provided by the control card have fixed locations at the beginning of COMMON. The remainder of the core is assigned for
the program instructions and temporary storage for the coefficient matrix and the right-hand-side vector of the governing equations. The program consists of four links, and since all the program instructions are not required simultaneously, only the instructions for each link in turn as needed are retained in the core.

A sketch of the governing equations in Eq. (20) is given in Fig. II-1 (Vol. I). Since the coefficient matrix is symmetric, the program allows storage only for the shaded area shown in the figure. From Eq. (20) it may be observed that, for a fixed $j, \partial \tau_{i j}$ represents a vector listing the forces that may develop when a unit deflection is applied in the fixed $j$ direction, while keeping all the other degree-of-freedom directions with zero deflections. The nonzero entries of this vector coincide with the deflection directions of only those mesh points that are connected with the disturbed mesh point by means of mesh elements and deflection boundary conditions. This shows that the $\ell_{i j}$ matrix is sparse and usually has a large zero area in the upper right-hand corner.

Before generating the coefficient matrix and the righthand side of Eq. (20), the program computes a pointer for each of the rows of the coefficient matrix, and a pointer for the right-hand-side vector so that the coefficients shown in the shaded area in Fig. II-1 (Vol. I) can be stored compactly in COMMON as a string. Actually, the pointers of the rows are the addresses of the words immediately preceding the diagonal elements. As discussed in Ref. 5, by the proper ordering of $d_{j}$ unknowns in Eq. (20), the zero area may be increased in the upper right-hand corner of $\partial \mathcal{X}_{i j}$. If the user chooses to assign zero into the ISHUF field of the control card, the unknowns are ordered as implied by the mesh-point labels; e.g., the unknown deflection components of the first mesh point are placed first, those of the second mesh point are placed second, etc. If the user assigns ISHUF $=1$ or 2 , the program first tries to find a better labelling system with the method given in Ref. 5, and uses these new mesh-point labels in ordering the unknowns $d_{j}$. For example, if mesh point with label 25 is the first mesh point in the new labelling system, the unknowns of this mesh point are listed first in $d_{j}$. If the user assigns ISHUF $=3$, the better labelling system is required by the program from input data cards (number 17 in Fig. IV-1, Vol. I). The method for relabelling described in Ref. 5 requires the generation of the mesh-point connectivity matrix $N_{s t}$, which is a binary matrix. If mesh point $s$ is connected to mesh point $t$ by a mesh element or by a deflection boundary condition, $N_{s t}=N_{t s}=1$; otherwise, $N_{s t}=N_{t s}=0$. It is always assumed that a mesh point is connected to itself.


Fig. II-I. Definition of $E_{l t j}$, augmented matrix
If point $s$ is completely constrained by the deflection boundary conditions, $N_{s t}=N_{t s}=0$ for all $t$, except $t=s$. The program generates $N_{s t}$ from the information provided by the element description data and deflection boundary conditions. The connectivity matrix $N_{s t}$ is always generated, since it is also used in determining the pointers of the rows of $\Theta \mathscr{C}_{i j}$.
2. Method of Assembly. To obtain the coefficient matrix and the right-hand-side vector of the governing equations, the mesh elements are processed, one at a time, first to obtain the element stiffness matrix and the element load vector for each, and then to assemble these according to Eqs. (21) and (22) and the allocated storage. Let $A^{r}$ denote the vector in COMMON and $R_{i}$ denote the pointer of the $i$ th equation in Eq. (20). Let us assume that the right-hand-side vector is stored after the coefficient matrix. Let $E_{k s i^{\prime}}$ (or $E_{l t j^{\prime}}$ ) denote the augmented matrix composed of $e_{k s i}$ (or $e_{t t}$ ) and $e_{k s}^{0}$ (or $e_{t t}^{0}$ ), as shown in Fig. II-1. Let a prime on the index indicate that the range of unprimed index is increased by 1 , and let an underlined index indicate the largest value within the range. With this notation, the assembly procedure may be summarized as

$$
\begin{align*}
A^{r}= & \gamma_{i}^{r}, E_{k s i}, Q_{k s}+\gamma_{i}^{\tau}, E_{k s i} \cdot \mu_{g s}^{m} P_{k g}^{m} \\
& +\delta_{i^{\prime}, j}^{r}, E_{k s s^{\prime}}, \mu_{g s}^{m} K_{k g l h}^{m} \mu_{h t}^{m} E_{l t j^{\prime}} \tag{25}
\end{align*}
$$

where
$\left.\begin{array}{lr}\delta_{i^{\prime} j^{\prime}}^{r}=1 & \text { if } i^{\prime} \leq j^{\prime}, j^{\prime} \leq \underline{i}, \text { and } r=R_{i^{\prime}}+j^{\prime}-i^{\prime} \\ \delta_{i^{\prime} j^{\prime}}^{r}=-1 & \text { if } i^{\prime}>\underline{i}, i^{\prime} \leq \underline{i}, \text { and } r=R_{\underline{i}}+i^{\prime} \\ \delta_{i^{\prime} j^{\prime}}^{r}=0 & \text { for all other possibilities }\end{array}\right\}$
and

$$
\left.\begin{array}{cc}
\gamma_{i^{\prime}}^{r}=1 & \text { if } i^{\prime} \leq \underline{i} \text { and } r=R_{\underline{i}}+i^{\prime}  \tag{27}\\
\gamma_{i^{\prime}}^{r}=0 & \text { for all other possibilities }
\end{array}\right\}
$$

In the program, only the nonzero $E_{k s i}$, constants are computed by the deflection boundary condition input units and the connectivity matrix. For each $k s$ the nonzero entries of $E_{k s i}$, are stored with their values and $i^{\prime}$ indices. The values and the indices of nonzero $Q_{k s}$ entries are directly provided by the concentrated load input cards. The binary coefficients $\delta_{i^{\prime} j^{\prime}}^{r}$ and $\gamma_{i}^{r}$, are not stored, but determined from Eqs. (26) and (27). If $m$ and $g$ are given, the $s$ value of the nonzero entry of $\mu_{g_{s}}^{m}$ is obtained from the element description data. Let $\widetilde{e}_{a}$ and $i_{a}^{\prime}$ denote the nonzero values and corresponding indices in $E_{k s i}$, for a given $k s$, and let $\underline{a}$ denote the maximum value of $a$, so that $1 \leq \underline{a} \leq a$ ( $b$ and $\underline{b}$ are alternate symbols). Let $\underline{d}$ denote the number of concentrated load input units. This notation is used in the flow diagram corresponding to Eq. (25) given in Fig. II-2. In the ELAS program, the summations implied by the first term in the right-hand side of Eq. (25) are implemented in Link 1 and the remainder in Link 2.

## 3. Method of Solution of the Governing Equations.

 Since $\ell \mathscr{C}_{i j}$ is a symmetric and positive-definite and bandlimited matrix for the solution of Eq. (20), the Cholesky algorithm may be applied. In this method, the decomposed stiffness matrix $B_{i j}$ from $\mathcal{L}_{i j}$ is first computed as$$
\begin{equation*}
B_{i j}, B_{j^{\prime} j}=Q\left\{_{i j}\right. \tag{28}
\end{equation*}
$$

where the range of $j^{\prime}$ equals that of $i$. Then, from

$$
\begin{equation*}
B_{i j} d_{j}^{\prime}=\mathscr{P}_{i} \tag{29}
\end{equation*}
$$

with a forward sweep, the auxiliary unknowns $d_{j}^{\prime}$ can be solved. Finally,

$$
\begin{equation*}
B_{i j} d_{i}=d_{j}^{\prime} \tag{30}
\end{equation*}
$$

yields the unknowns with a backward sweep. In Fig. II-1 (Vol. I) the border of the zero area in the upper righthand corner is not always defined by the last nonzero coefficient in each equation. This is because the shaded areas of $Q \mathcal{L}_{i j}$ and $B_{i j}$ are identical only when the border is selected as shown in Fig. II-1 (Vol. I). In the ELAS program, $\mathscr{C _ { i j } \text { coefficients in the shaded area of the figure are } { } ^ { \text { a } } \text { . }}$ first modified to those of the coefficients of $B_{i j}$, then $\mathscr{P}_{i}$ constants are changed to $d_{i}^{\prime}$, and finally, $d_{i}^{\prime}$ values are converted to the numerical values of the $d_{j}$ unknowns. Then, from Eq. (17), $q_{l t}$ is computed on the same area as $d_{j}$. These operations are carried out in Link 3 of ELAS.
4. Computation of Stresses. The computation of stresses in displacement methods poses a harder problem in structures of two- or three-dimensional continuum than that in truss and frame structures, which truly have a finite num-


Fig. II-2. Flow diagram corresponding to the summations implied by Eq. (25)
ber of deflection components for the determination of their distorted configuration. The difficulty arises from the fact that the structures of two- or three-dimensional continuum actually have infinitely many deflection components, and the relations of the type of Eq. (12) are only approximate.

After computing the deflections as mesh functions, the problem of stress computation with acceptable accuracy in reasonable machine times still remains. Experience has shown that the use of Eq. (13) and then Eq. (6) presents the following drawbacks: (I) the exact location of the point for which the stresses are computed is not known, and (2) the computed stresses may be largely different from the actual stresses. Despite these setbacks, stress computation of this type is being widely used because it is modular in elements, just as is the generation of the governing equations in Eq. (20), a feature that facilitates automation. In the ELAS program, the best-fit stress computation method of Ref. 4 is used for structures of continuum. This method is just as easy to automate and has the following advantages: (1) stresses are computed at the points where the deflections are obtained, (2) the accuracy in stresses is comparable with that of deflections, and (3) the stress boundary conditions of Eq. (4) may be satisfied during the computation of the stresses of boundary points. This scheme was initially devised for triangular finite elements (Ref. 6).

In the following paragraphs the stress computation in structures of two- and three-dimensional continua is explained. The computation of stresses in structures composed of elements of one-dimensional continuum is performed by multiplying element stiffness matrices with computed deflections.

Mesh Line Set. Suppose that the deflections at the mesh points of a structure of three-dimensional continuum are available and that the stresses at mesh point $s$ are requested. The question of how much deflection data should be included in the computation is of practical importance because the computation time rapidly increases with this quantity. Experience with the method of computing stresses in the element indicates that deflections of the set of elements meeting at mesh point $s$ are sufficient for the computation of its stresses with acceptable accuracy. The mesh points of the element set are called "mesh-point set" and the mesh lines meeting at the common mesh point $s$ are called "mesh-line set." The scheme adopted in ELAS is modular in the mesh-line set-the next-best unit after elements. The stress computation at a mesh point starts with the determination of the element set, and
consequently, the mesh-line set associated with this mesh point. Then, if this mesh point is on the boundary, the average boundary surface area associated with this node and the direction cosines of the outer normal are computed.

Selection of Local Coordinate Systems at the Mesh Points. In a given problem, it is desirable to have one fixed, right-handed coordinate system to express the stresses. However, this is not practical for structures composed of anisotropic material, at the boundary points where the outer normal is not coincident with the coordinate lines, and for shell structures. The following method is adopted in the ELAS program for the selection of local coordinate systems at the mesh points.

At an internal node, the local axes may be taken as the material axes unless the material is isotropic, in which case they should be taken (1) parallel to the overall coordinate system in plates and three-dimensional solids, and (2) as the principal curvature directions and the normal of the middle surface, or any other suitable system that the user inputs in shells. At a boundary node, the first local axis may be coincident with the outer normal, and the directions of the remaining local axes may be determined so that (1) the local third axis becomes the middle surface normal in plates and shells, and (2) the direction defined by the cross-product of outer normal with the overall axis, which makes the largest angle with the outer normal, then becomes the second local axis in three-dimensional solid structures.

Stress Computation at an Internal Mesh Point. Let $j$ be the label of a mesh line in the mesh-line set, with $\Delta \rho_{\gamma^{\prime}}$ the position vector and $\Delta u_{\gamma}$, the displacement vector of the far end of the mesh line relative to the mesh point where the local coordinates are defined. As the first approximation of the strain along the $j$ th mesh line, the following may be written:

$$
\begin{equation*}
\epsilon=\frac{\Delta \rho_{\gamma^{\prime}} \Delta u_{\gamma^{\prime}}}{\Delta \rho_{\mu^{\prime}} \Delta \rho_{\mu^{\prime}}} \tag{31}
\end{equation*}
$$

The same strain may be obtained from the strain tensor of mesh point $s$ as

$$
\begin{equation*}
\epsilon=\frac{\Delta \rho_{\alpha^{\prime}} \Delta \rho_{\beta^{\prime}} \epsilon_{\alpha^{\prime} \beta^{\prime}}}{\Delta \rho_{\mu^{\prime}} \cdot \Delta \rho_{\mu^{\prime}}} \tag{32}
\end{equation*}
$$

Equating Eq. (31) to Eq. (32) and cancelling the denominators results in the following expression:

$$
\begin{equation*}
\left(\Delta \rho_{\alpha^{\prime}} \Delta \rho_{\beta^{\prime}}\right)_{j} \epsilon_{\alpha^{\prime} \beta^{\prime}}=\left(\Delta \rho_{\gamma^{\prime}} \Delta u_{\gamma^{\prime}}\right)_{j} \tag{33}
\end{equation*}
$$

The number of equations in Eq. (33) is equal to the range $j$. Usually the range of $j$ is greater than the number of independent components of the strain tensor. (If not, the mesh may be readjusted by repeating the deflection computation.) Therefore, in Eq. (33), there are more equations than the unknown strain components. Such a set may be solved by least squares, first by multiplying both sides with the transpose of the coefficient matrix, then by inverting the new coefficient matrix. This leads to

$$
\begin{equation*}
\epsilon_{\alpha^{\prime} \beta^{\prime}}=\left[\left(\Delta \rho_{\alpha^{\prime}} \Delta \rho_{\beta^{\prime}}\right)_{i}\left(\Delta \rho_{\delta^{\prime}}, \Delta \rho_{v^{\prime}}\right)_{i}\right]^{-1}\left(\Delta \rho_{\delta^{\prime}} \Delta \rho_{v^{\prime}}\right)_{j}\left(\Delta \rho_{\gamma^{\prime}} \Delta u_{\gamma^{\prime}}\right)_{j} \tag{34}
\end{equation*}
$$

where the range of $i$ is equal to that of $j$. The stresses at the mesh point may be obtained by substituting $\epsilon_{\alpha^{\prime} \beta}$, from Eq. (34) into Eq. (5). If the problem is a plane-strain problem, one should first impose

$$
\begin{equation*}
\epsilon_{3^{\prime} 3^{\prime}}=\epsilon_{1^{\prime} 3^{\prime}}=\epsilon_{3^{\prime} 1^{\prime}}=\epsilon_{2^{\prime} 3^{\prime}}=\epsilon_{3^{\prime} 2^{\prime}}=0 \tag{35}
\end{equation*}
$$

on Eq. (33). For plane-stress problems, $D_{\alpha^{\prime} \beta^{\prime} \gamma^{\prime} \delta^{\prime}}$ in Eq. (5) should be modified to guarantee

$$
\begin{equation*}
\epsilon_{1^{\prime} 3^{\prime}}=\epsilon_{3^{\prime} 1^{\prime}}=\epsilon_{2^{\prime} 3^{\prime}}=\epsilon_{3^{\prime} 2^{\prime}}=\sigma_{3^{\prime} 3^{\prime}}=0 \tag{36}
\end{equation*}
$$

For the bending of plates and shells, $\epsilon_{\alpha^{\prime}, \beta}$, should be interpreted as curvature changes and, in Eq. (33), ( $\left.\Delta \rho_{\gamma}, \Delta u_{\gamma^{\prime}}\right)_{j}$ should be taken as the projection of the rotations vector of the far-end mesh point relative to the current mesh point in the $j$ th mesh line on $\vec{i}_{\xi} \times \Delta \vec{\rho}$ direction, where $\vec{i}_{\xi}$ is the unit vector of the third local axis. Also, the conditions in Eq. (36) should be imposed on $D_{\alpha^{\prime} \beta^{\prime} \gamma^{\prime} \delta^{\prime}}$ and Eq. (5) should be replaced by

$$
\begin{equation*}
M_{\gamma^{\prime} \delta}=-\frac{t^{3}}{12} D_{\gamma^{\prime} \delta^{\prime} \alpha^{\prime} \beta}, \epsilon_{\alpha^{\prime}, \beta^{\prime}} \tag{37}
\end{equation*}
$$

where $M_{\gamma^{\prime} \delta}$ denotes the bending moments and $t$ is the thickness. The membrane case of shells is identical with the plane-stress case, provided that Eq. (5) is replaced with

$$
\begin{equation*}
N_{\gamma^{\prime} \delta^{\prime}}=t D_{\gamma^{\prime} \delta^{\prime} \alpha^{\prime} \beta^{\prime}} \epsilon_{\alpha^{\prime} \beta^{\prime}} \tag{38}
\end{equation*}
$$

where $N_{\gamma^{\prime} \delta}$, denotes the membrane forces.
Stress Computation at a Boundary Mesh Point. The procedure for stress computation at a boundary mesh point is basically the same as the computation at an internal mesh point. Here, the stress boundary conditions, expressed in terms of the strains, are included in Eq. (33) before the application of the least-squares scheme for their solution. The stress boundary conditions may be written as

$$
\begin{equation*}
\frac{n_{\beta^{\prime}}}{D_{1^{\prime} 1^{\prime} 1^{\prime} 1^{\prime}}} D_{\alpha^{\prime} \beta^{\prime} \gamma^{\prime} \delta^{\prime}} \epsilon_{\gamma^{\prime} \delta^{\prime}}=\frac{-p_{\alpha^{\prime}}}{D_{1^{\prime} 1^{\prime} 1^{\prime} 1^{\prime}}} \tag{39}
\end{equation*}
$$

where $p_{\alpha^{\prime}}$ represents the prescribed boundary stresses. If at the boundary mesh point the deflections, in place of the stresses, are prescribed, $R_{\alpha^{\prime}}$ reaction forces may be found from the equilibrium equations of the boundary node, and the following may be written:

$$
\begin{equation*}
p_{\alpha^{\prime}}=-\frac{R_{\alpha^{\prime}}}{A} \tag{40}
\end{equation*}
$$

where $A$ is the average boundary surface area associated with the mesh point. In Eq. (39), the reason for division by $D_{1^{\prime} 1^{\prime} 1^{\prime} 1^{\prime}}$ is to reduce the coefficients of strains in the stress boundary equations to the same order of magnitude as those of Eq. (33). The procedure described here for a three-dimensional solid may be readily extended to other types of structures with the help of previous paragraphs.

# III. COMMON Variables and COMMON Blocks of the Program 

The memory organization in each of the four links of ELAS is illustrated in Fig. III-1. Table III-1 lists the blocks of COMMON sequentially, and gives a short description of each block. In the table, the variableaddress blocks are listed with increasing COMMON addresses. It should be noted that the variable-address blocks in COMMON are packed in a string, one after the other, without any waste of core locations. Such blocks may be properly located by means of pointers, which are also in COMMON. A pointer is a word whose content is one less than the COMMON address of the first word
of the associated COMMON block. The constituents of Block Group 1 are listed in Table VI-3 (Vol. I), in the order in which they appear in COMMON. These constituents are alphabetically ordered with their symbolic names in Table III-2. In Table III-3, the meanings of entries of important vectors, especially those defined by the pointers, are given. The additional COMMON variables of Link 4 are listed alphabetically in Table III-4, and with increasing COMMON addresses in Table III-5. Table III-5 also contains a short description of these variables.
Table III-1. Sequence and descriptions of COMMON blocks**


Table III-2. Alphabetical listing of the constituents of COMMON block group $\mathbf{1}^{\text {a }}$

| Symbol | Location in COMMON | Brief description | Symbol | Location in COMMON | Brief description |
| :---: | :---: | :---: | :---: | :---: | :---: |
| AA | 1-. $\cdot$ | Name of whole COMMON block for floating- | IELT | 28 | Element type number |
|  |  |  | IERR | 79 | Error indicator |
| ACEL | 39 | Body force per unit volume | IGEM | 78 | Indicator for structures inscribed in |
| ALI | 83 | Thermal expansion coefficient of an element |  |  | $(z=0)$-plane |
|  |  | in first material axis direction | IH | 10 | Maximum number of vertices |
| Al2 | 84 | Thermal expansion coefficient of an element in second material axis direction | IIA | 62 | Pointer for thermal expansion coefficient array |
| Al3 | 85 | Thermal expansion coefficient of an element | IIC | 74 | Pointer for dbe unit constants array |
|  |  | in third material axis direction | IID | 61 | Pointer for material constants array |
| CONS | 45 | Constant for element load vector | IIS | 77 | Pointer for subelement stiffness matrix $\text { (IIS }=350)$ |
| DG | 82 | Temperature gradient for an element in direction $y$ (or $z$ ) | IMAT | 7 | Number of material types |
| DGY | 332 | Temperature gradient along local $y$-axis for | IMES | 326 | Indicator for mesh topology input |
|  |  | an element | IMET | 31 | Material type number |
| DGZ | 331 | Temperature gradient along local z-axis for an element | IMFI | 15 | Number of angle types |
|  |  |  | IMMX | 12 | Number of torsion consfants types |
| DT | 81 | Value of temperature change for an element | IMMY | 13 | Number of $\boldsymbol{\gamma}$-moment of inertia types |
| D21 | 86-106 | Material constants for an element | IMMZ | 14 | Number of z-moment of inertia types |
| G1 | 47 | First direction cosine of acceleration vector | IMS | 34 | Number of vertices of current element |
| G2 | 48 | Second direction cosine of acceleration vector | IN | 1 | Total number of nodal points |
| G3 | 49 | Third direction cosine of acceleration vector | IND | 33 | IND $=$ IDEG * IN |
| IA | 1-* | Name of whole COMMON block for fixedpoint references | INP | 42 | Indicator for output level |
| IARE | 16 | Number of cross-sectional area types | INX | 9 | Number of last link to be executed |
| 1BB | 59 | Pointer for IBB array | IORD | 37 | Number of words allocated for the reduced stiffness mafrix |
| IBN | 2 | Total number of dbc input units | 1001 | 38 | IORD1 $=10 \mathrm{RD}+1$ |
| 1BO | 60 | Pointer for IBO array | IP | 4 | Total number of nonzero concentrated load |
| IBUN | 327 | Indicator for boundary conditions input |  |  | components |
| ICAR | 66 | Pointer for cross-sectional areas array | IPBG | 43 | Integer constant for element load vector |
| ICFI | 70 | Pointer for angles array | IPEN | 44 | Integer constant for element load vector |
| ICIX | 67 | Pointer for torsional constants array | IPIR | 329 | Indicator for local coordinate axes selection |
| ICIY | 68 | Pointer for $\boldsymbol{y}$-moments of inertia array | IPR | 333 | Pointer for pressure array |
| 1CIZ | 69 | Pointer for z-moments of inertia array | IPRS | 5 | Number of pressure fypes |
| ICOR | 328 | Indicator for coordinates input | ISDT | 348 | Number of temperature change types |
| IDEF | 75 | Pointer for unknown deflections array (initially loads array) | ISDY | 347 | Number of temperature gradients along local $y$-axis |
| IDEG | 8 | Degrees of freedom at a node | ISDZ | 346 | Number of temperature gradients along local $z$-axis |
| IDS | 36 | Order of the subelement stiffness malrix | ISHUF | 35 | Relabeling indicator |
| 1DT | 63 | Pointer for temperature changes array | IST | 76 | Pointer for reduced stiffness matrix of the whole structure |
| IDY | 64 | Pointer for temperature gradients array ( $y$-direction) | ISTR | 27 | Indicator for plane-strain case |
| IDZ | 334 | Pointer for temperature gradients array (z-direction) | ISUM <br> IT | $32$ $3$ | Number of equations in the reduced system <br> Total number of elements |

[^1] tion of the block is partly excluded.

Table III-2 (contd)

| Symbol | Location in COMMON | Brief description | Symbol | Location in COMMON | Brief description |
| :---: | :---: | :---: | :---: | :---: | :---: |
| ITAP | 41 | Chain program tape number | $J 4$ | 53 | Pointer for J4W array |
| ITAS | 335 | Scratch tape number | J5 | 54 | Pointer for J5W array |
| ITE | 65 | Pointer for thicknesses array | J6 | 55 | Pointer for J6W array |
| ITEM | 29 | Temperature change type number | J7 | 56 | Pointer for JTW array |
| ITIC | 30 | Thickness type number | 38 | 57 | Pointer for J8W array |
| ITYPE | 6 | Indicator for material type | J9 | 345 | Pointer for J9W array |
| IU | 46 | Pointer for diagonal-element count vector | 110 | 344 | Pointer for Jlow array |
| IXX | 71 | Pointer for $X$-coordinates array | M | 25 | Label of current element |
| IYY | 72 | Pointer for Y-coordinates array | $\mathrm{N}_{\mathrm{i}}$ | 17-24 | Labels of vertices of an element |
| IZZ | 73 | Pointer for Z-coordinates array | NTIC | 349 | Number of thickness types |
| 18 | 11 | Maximum number of words to describe an element | P PRES | $107-130$ 330 | Load vector of a subelement |
|  |  |  | PRES | 330 | Pressure value for an element |
| JARE | 340 | Type number of cross-sectional area | $S$ | 351-** | Subelement stiffness matrix |
| JMFI | 15 | Number of angle types | TE | 80 | Value of thickness for an element |
| JMMX | 339 | Type number of the torsional constant about local $x$-axis | UV | 131-154 | Deflections due to temperature changes for an element |
| JMMY | 338 | Type number of the sectional moment of inertia about local $y$-axis | X XD | $155-162$ $179-185$ | Overall X-coordinates for vertices of an element $X$-coordinates of vertices, other than the first |
| JMMZ | 337 | Type number of the sectional moment of inertia about local $z$-axis |  |  | vertex, of an element relative to the first vertex |
| JPRS | 343 | Type number of pressure | Y | 163-170 | Overall $Y$-coordinates of vertices of an element |
| JSDY | 342 | Type number of temperature gradient along local $y$-axis | YD | 186-192 | $Y$-coordinates of vertices, other than the first verlex, of an element relative to the first vertex |
| JSDZ | 341 | Type number of temperafure gradient along local $z$-axis |  | 171-178 | Overall Z-coordinates of vertices of an element |
| $J 1$ | 50 | Pointer for JIW array | ZD | 193-199 | $Z$-coordinates of vertices, other than the first vertex, of an element relative to the first |
| J2 | 51 | Pointer for J2W array |  |  |  |
| $J 3$ | 52 | Pointer for J3W array | ZGEM | 40 | Floating-point equivalent of IGEM |

Table III-3. Meanings of the entries of important vectors

| Vector in COMMON | Meaning of rth entry of the vector (all divisions are in infeger arithmetic sense) |
| :---: | :---: |
| AA vector | The rth component of the total COMMON vector in floating-point mode |
| D21 vector | The rth component of a row-listed upper $6 \times 6$ material matrix (see Fig. Ill-2b, Vol. I), if it exists |
| IA vector | The rth component of the total COMMON vector in fixed-point mode |
| IBB-pointer-related vector | IBB value of Jth degree of freedom direction at ith node (user's label); $i=1+(r-1) / I D E G$, $J=r-(i-1)^{*}$ IDEG (see Table VI-2, Vol. I) |
| IBO-pointer-related vector | IBO value of $J$ th degree of freedom direction at ith node (user's label); $i=1+(r-1) /$ IDEG, $J=r-(i-1)^{*}$ IDEG (see Table VI-2, Vol. I) |
| ICAR-poinfer-related vector | Value of $r$ th-fype cross-sectional area, if it exists |
| ICFI-pointer-related vector | Value of rth-type angle defining principal axes of cross section, if it exists |
| ICIX-pointer-related vector | Value of rth-type torsional constant, if it exists |
| ICIY-pointer-related vector | Value of rth-type $\boldsymbol{\gamma}$-moment of inertia, if it exists |
| ICIZ-pointer-related vector | Value of rih-type z-moment of inertia, if it exists |
| IDEF-pointer-related vectors | (1) Value of prescribed concentrated load in Jth degree of freedom direction at node $i$ (user's label); $i=1+(r-1) /$ IDEG, $J=r-(i-1)^{*}$ IDEG |
|  | (2) Value of rth component of reduced load vector (the right-hand-side vector in Fig. II-1, Vol. I) |
|  | (3) Value of rih component of reduced deflection vector (\{d\} vector in Fig. II-1, Vol. I) <br> (4) Value of deflection at the $J$ th degree of freedom direction at node $;$ (user's label); $i=1+$ $(r-1) /$ IDEG, $J=r-(i-1)^{*}$ IDEG |
|  | For (2) and (3) the node $;$ (user's label) and direction J associated with the rth entry may be obtained as follows: Let $r^{\prime \prime}$ be the entry number of the word, in IBB-pointer-related vector, where the absolute value is $r$ and $r^{\prime \prime}$ th entry of IBO-pointer-related vector is -1 . Then $i=1+\left(r^{\prime \prime}-1\right)$ /IDEG, and $J=r^{\prime \prime}-(i-1)^{*}$ IDEG |
| IDT-pointer-relafed vector | Value of rth-type temperature increase, if it exists |
| IDY-pointer-related vector | Value of rth-type temperature gradient in $\boldsymbol{\gamma}$-direction, if it exists |
| IDZ-pointer-related vector | Value of rth-type femperature gradient in z-direction, if it exists |
| IIA-pointer-related vector | Value of thermal expansion coefficient in the $J$ th material axes direction in element $; \boldsymbol{i}=1+$ $(r-1) / k_{2}, J=r-k_{2}^{*}(i-1)$, where $k_{2}$ is 1,2 , or 3 , depending upon whether ITYPE is 0,1 , or 2 , respectively |
| IIC-pointer-related vector | Value $C$ of $J$ th degree of freedom direction at $i$ th node (user's label); $i=1+(r-1) /$ IDEG, $J=r-(i-1)^{*}$ IDEG (see Table V1-2, Vol. I) |
| UD-pointer-related vector | Value of th material constant of material type $; i=1+(r-1) / k_{1}, j=r-k_{1}{ }^{*}(i-1)$, where $k_{1}$ is 2,9 , or 21 , depending upon whether ITYPE is 0,1 , or 2 , respectively |
| IIS-pointer-related vector | Element $k_{m n}$ of the free-free subelement stiffness matrix, $m=1+(r-1) / I D S^{\prime}, n=r-$ IDS $^{\prime} *$ ( $m-1$ ); $m$ corresponds to $m^{\prime}$ th degree of freedom direction ( $m^{\prime}=1+(m-1) / 1 \mathrm{MS}^{\prime}$ ) af vertex $m^{\prime \prime}\left(m^{\prime \prime}=m-1 M 5^{*}(m-1) ; n\right.$ corresponds to $n^{\prime}$ th degree of freedom direction ( $\left.n^{\prime}=1+(n-1) / I M S^{\prime}\right)$ at vertex $n^{\prime \prime}\left(n^{\prime \prime}=n-I D S^{\prime *}\left(n^{\prime}-1\right)\right.$; $1 M S^{\prime}$ is the number of vertices of subelement, and IDS ${ }^{\prime}=1$ MS ${ }^{*}$ IDEG |
| IPR-pointer-related vector | Value of rth-type pressure, if it exists |
| IST-pointer-related vector | (1) Element $K_{m n}$ of the stiffness matrix of the supported structure. To find mth direction, enter IBB-pointer-related vector with the entry number $r^{\prime}$ of the word, in IU-poinfer-related vector, which is closesf to, but not greater than $r$. Let $r^{\prime \prime}$ be the entry number of the word, in IBB-pointerrelated vector, whose absolute value is $r^{\prime}$ and the $r^{\prime \prime}$ th entry in IBO-poinfer-related vector is $-1 ; m$ th direction corresponds to $m^{\prime \prime}$ th degree of freedom direction at node $m^{\prime}$ (user's label); $m^{\prime}=1+\left(r^{\prime \prime}-1\right) /$ IDEG, $m^{\prime \prime}=r^{\prime \prime}-\left(m^{\prime}-1\right)^{*}$ IDEG. To find $n$th direction, determine $s^{\prime}$ by adding to $r^{\prime}$ the difference between $r$ and the $r^{\prime}$ th entry of IU-pointer-related vector. Let $s^{\prime \prime}$ be the entry number of the word in IBB-pointer-related vector, whose absolute value is $s^{\prime}$ and the $s^{\prime \prime}$ th entry in IBO-pointer-related vector is - 1 ; $n$th direction corresponds to $n^{\prime \prime}$ th degree of freedom direction at node $n^{\prime}$ (user's label); $n^{\prime}=1+\left(s^{\prime \prime}-1\right) / 1$ DEG, $n^{\prime \prime}=s^{\prime \prime}-\left(n^{\prime}-1\right)^{*}$ IDEG <br> (2) Value of residual force acting at node $i$ in direction $J$. where $i=1+(r-1$ I/IDEG, $J=r-$ ( $\mathbf{i}-1$ )* IDEG |

## Table III-3 (contd)

| Vector in COMMON | Meaning of rth entry of the vector (all divisions are in integer arithmetic sense) |
| :---: | :---: |
| ITE-poinfer-related vector | Value of rth-type thickness, if it exists |
| IU-pointer-related vector | Entry number in IST-pointer-related vector of rth diagonal element of the reduced stiffness matrix |
| IXX-pointer-related vector | $X$-coordinate of noder $r$ (user's label) |
| IYY-pointer-related vector | $\boldsymbol{Y}$-coordinate of node $\boldsymbol{r}$ (user's label) |
| IZZ-pointer-related vector | Z-coordinate of node $\boldsymbol{r}$ (user's label), if it exists |
| J1-pointer-related vector | JIW value of rth element (see Table IV-3, Vol. I) |
| J2-pointer-related vector | J2W value of r th element (see Table IV-3, Vol. I) |
| J3-pointer-related vector | J3W value of rth element (see Table IV-3, Vol. I) |
| j4-poinfer-related vector | J4W value of rth element (see Table IV-3, Vol. I) |
| 55-pointer-related vector | J5W value of rth element (see Table IV-3, Vol. I) |
| J6-pointer-related vector | J6W value of $\mathrm{r}^{\text {th }}$ element ( (see Table IV-3, Vol. I), if it exists |
| J7-pointer-related vector | J7W value of $r$ th element (see Table IV-3, Vol. 1), if it exists |
| J8-pointer-related vector | J8W value of r (h element (see Table IV-3, Vol. I), if it exists |
| J9-pointer-related vector | J9W value of rth element (see Table IV-3, Vol. I), if it exists |
| J10-pointer-related vector | J10W value of $r$ th element (see Table IV-3, Vol. 1), if it exists |
| N vector | The label (user's) of the rth vertex of an element |
| MAX-pointer-related vector | Number of nonzero entries above rth diagonal element of the decomposed reduced stiffness matrix (see Fig. II-1, Vol. I) |
| P vector | Element load acting in direction $J$ of $i$ th vertex of a subelement; $J=1+(r-1) / 1 \mathrm{MS}^{\prime}, i=r-$ ( $J-1)^{*} I \mathrm{MS}^{\prime}$ (IMS' $=$ number of vertices of the subelement) |
| S matrix | See IIS-pointer-related vector |
| UV vector | Deflection in direction $J$ of ith vertex of a subelement subjected to temperature change in local coordinates; $J=1+(r-1) / I M S^{\prime}, i=r-(J-1)^{*} I M S^{\prime}\left(I M S^{\prime}=\right.$ number of vertices of subelement) |
| $X$ vector | $X$-coordinate of r th vertex of an element |
| XD vector | X -coordinate, relative to the first vertex, of ( $r+1$ ) st vertex of an element |
| Y vector | $Y$-coordinate of $r$ th vertex of an element |
| YD vector | $\boldsymbol{\gamma}$-coordinate, relative to first vertex, of (r+1) st vertex of an element |
| $Z$ vector | $Z$-coordinate of the $r$ th vertex of an element |
| ZD vector | Z-coordinate, relative to the first vertex, of $(r+1)$ st vertex of an element |

Table III-4. Alphabetical list of additional COMMON variables for Link $4^{a}$

| Symbol | COMMON location | Symbol | COMMON location | Symbol | COMMON location | Symbol | COMMON location |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| A | 14786-15415 | ICAS | 212 | IWG | 14660-14749 | NU | 292-294 |
| ANGIE | 211 | ICLA | 206 | JM1 | 293 | QF | 253-258 |
| ARE | 205 | ICLAS | 274-277 | JPI | 292 | QN | 247-252 |
| AST | 203 | ICN | 201 | JS1 | 294 | RED | 265-270 |
| B | 15416-15479 | ICOL | 295 | LM | 202 | RES | 259-264 |
| BAS | 271-273 | ICON | 210 | MAC | 14400-14659 | SIR | 223-225 |
| BIR | 220-222 | IDR | 297 | MSET | 15596-15695 | SR | 235-240 |
| BST | 217 | IE | 213 | MB | 215 | W | 15696-15704 |
| C | 15480-15495 | IM | 208 |  | 215 | W | 15696-15704 |
| DD | 14750-14785 | IMEL | 207 | NB | 214 | XF | 244-246 |
| DIN | 226-234 | INBON | 204 | NBAN | 278-287 | XN | 241-243 |
| ETA | 229-231 | IONE | 200 | NEL | 14000-14399 | XII | 226-228 |
| FF | 14000-15704 | IRIG | 296 | NES | 295-297 | ZTA | 232-234 |
| IC | 209 | IROT | 216 | NSET | 15496-15595 |  |  |

Table III-5. List of additional COMMON variables for Link $4^{\text {a }}$

| Location in COMMON | Symbol | Brief description |
| :---: | :---: | :---: |
| 1-199 |  | This portion of COMMON is as in Table VI-3 of Vol. 1 |
| 200 | IONE | Total number of one-dimensional elements in the structure |
| 201 | ICN | Label of current mesh point (ICN varies from 1 to (N) |
| 202 | LM | Total number of non-one-dimensional elements meeting at mesh point ICN |
| 203 | AST | Indicator containing * or BCD blank, depending upon whether mesh point ICN is on boundary or not, respectively |
| 204 | INBON | Indicator containing 1 or 0 , depending upon whether mesh point ICN is on boundary or not, respectively |
| 205 | ARE | Average boundary surface area for mesh point ICN, if it is on boundary |
| 206 | ICLA | Total number of class ${ }^{\text {b }}$ types for elements of material type group IM at mesh point ICN |
| 207 | IMEL | Total number of material types at mesh point ICN |
| 208 | IM | Current material type group number (IM varies from 1 to (MEL) |
| 209 | IC | Current class ${ }^{\text {b }}$ type group number (IC varies from 1 to ICLA) |
| 210 | ICON | Sequence number of current strain-deflection equation at mesh point ICN for material group IM and for class ${ }^{\text {b }}$ group IC |
| 211 | ANGLE | Angle between XII local axis and the 1-2 line of the lowest labeled shell element attached to mesh point ICN |
| 212 | ICAS | Class ${ }^{\text {b }}$ type number of $1 \mathrm{Cth}^{\text {class }}{ }^{\text {b }}$ group of 1 Mth material group at mesh point ICN |
| 213 | IE | Number of mesh elements (of class ${ }^{\text {b }}$ group IC of material group IM) plus 1 at mesh point ICN |
| 214 | NB | Total number of mesh points in node set at mesh point ICN |
| 215 | MB | Number of boundary points attached to mesh point ICN |
| 216 | IROT | Indicator containing 0 or 1, depending upon whether local axes at mesh point ICN are parallel to overall axes or not, respectively |
| 217 | BST | Indicator containing BCD blank or **, depending upon whether local axes at mesh point ICN are parallel to overall or not, respectively |
| 220-222 | BIR | Direction cosines of outer unit normal yector at mesh point ICN, if it is on boundary |
| aThis fable is not applicable to subrautine DIMI of Link 4. ${ }^{\mathrm{b}}$ Class types are those of Table VI-6, Vol. I. |  |  |

Table III-5 (contd)

| Location in COMMON | Symbol | Brief destription |
| :---: | :---: | :---: |
| 223-225 | SIR | Vector heading towards structure of mesh point ICN, if if is on boundary |
| 226-234 | DIN | Direction cosines of local axes in overall coordinate system at mesh poinf ICN (the columns of DIN are named as XII, ETA, and ZTA) |
| 235-240 | SR | Independent components of stress tensor for ICth class ${ }^{\text {b }}$ group of IMth material group at mesh point ICN |
| 241-243 | XN | Overall coordinates of mesh point ICN |
| 244-246 | XF | Overall coordinates of IIth vertex of ILth element of ICth class ${ }^{\text {b }}$ group of $\operatorname{IMth}$ material group at mesh poinf ICN |
| 247-252 | QN | Deflection components in overall coordinates of mesh point ICN |
| 253-258 | QF | Deflection components in overall coordinates of mesh point whose overall coordinates are in XF |
| 259-264 | RES | Residual forces ${ }^{\text {c }}$ in overall coordinates at mesh point ICN, if on boundary |
| 265-270 | RED | Relative deflections (in overall coordinates) of mesh point related with XF vector with respect to mesh point ICN |
| 271-273 | BAS | Direction cosines of 1-2 line of the lowest labeled element of class ${ }^{\text {b }}$ group IC of material group IM at mesh point ICN |
| 274-277 | ICLAS | Number of class ${ }^{\text {b }}$ groups in each material group (maximum 4) of mesh point ICN |
| 278-287 | NBAN | List of labels of boundary mesh points attached to mesh point ICN |
| 292-294 | NU | Vector containing the sequence numbers of the vertices after (JP1), before (JMI), and above (JS1) mesh point ICN in the lth mesh element of the node set (with Table III-5, vol. I) |
| 292 | JP1 | See NU (1) |
| 293 | JM1 | See NU (2) |
| 294 | JSI | See NU (3) |
| 295-297 | NES | Vector containing number of independent strain components (ICOL), number of right-hand sides (IRIG) and indicator of right-hand-side arrangement (IDR) (IDR $=0$ means lineal strains first, IDR $=1$ means rotational strains first) for current ICN/IM/IC |
| 295 | ICOL | See NES (1) |
| 296 | IRIG | See NES (2) |
| 297 | IDR | See NES (3) |
| 329-349 |  | This portion of COMMON is as shown in Table VI-3, Vol. I. See also Fig. III-3, Link 4 |
| 349-13999 |  | See Fig. III-1, Link 4 |
| 14000-15704 | FF | Vector containing information for stress conputation at mesh point ICN |
| 14000-14399 | NEL | Element set information of mesh point ICN (see Table VI-7, Vol. I) |
| 14400-14659 | MAC | Table for classes and material of element set at mesh point ICN (see Table VI-7, Vol. I) |
| 14660-14749 | IWG | Vector of weights of strain-deflection equations for current ICN/IM/IC |
| 14750-14785 | DD | Material matrix for current ICN/IM/IC |
| 14786-15415 | A | Augmented matrix of strain-deflection equations for current ICN/IM/IC |
| 15416-15479 | B | Coefficient matrix (or its inverse) of the least-squares equations for strain for current ICN/IM/IC |
| 15480-15495 | C | Right-hand-side vector(s) of the least-squares equations for strains for current ICN/IM/IC |
| 15496-15595 | NSET | List of labels of mesh points on the boundary and attached to mesh point ICN |
| 15596-15695 | MSET | Auxiliary array for NSET |
| 15696-15704 | W | Direction cosines of new material axes in the old for current ICN/IM/IC |
| ${ }^{\text {c Residual forces are those }}$ | Uutput Ite | Sect. VI-D and VI-E, VoI. II. |

## Appendix

## Corrigenda for Volume I

Corrigenda for Volume I

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[^0]:    *) Computer Software Management and Information Center, Computer Center, University of Georgia, Athens, Georgia, 30601, telephone 404-542-3265.

[^1]:    a See Table III-1 for sequence and descriptions of COMMON blocks. Table III-2 is a reordering of Table VI-3 (Vol. I), in both of which the "General Descriptors IV" sec-

