

1 **ELASTOPLASTIC CONSOLIDATION SOLUTIONS FOR SCALING FROM**
2 **SHALLOW PENETROMETERS TO PIPELINES**

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8 **Abstract**

9 The build-up of friction on seabed pipelines is an important design consideration, affecting their
10 stability and the resulting in-service strain and fatigue. The consolidation beneath a partially
11 embedded pipeline has been investigated in the past and linked to the build-up of axial pipe-
12 soil resistance. This paper extends previous work by providing solutions for consolidation
13 around a new class of shallow penetrometer, to provide a basis to scale from site investigation
14 results directly to the build-up of pipeline friction. Small strain finite element analyses, using
15 Modified Cam clay soil model, are presented for the novel toroid and ball penetrometers. The
16 effects of initial penetrometer embedment, device roughness, strength gradient and overload
17 ratio have been explored in a comprehensive manner, and are compared with pipe results. The
18 toroid penetrometer shows excellent agreement with an element of an infinitely long pipe,
19 simplifying the scaling process. The ball penetrometer shows a faster consolidation response,
20 typically by a factor of 3, reflecting the more effective drainage mechanisms of a three
21 dimensional device compared to a plane strain device. The dissipation responses are fitted by
22 simple equations to aid application in design.

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24 INTRODUCTION

25 Subsea pipelines serve as a significant component of offshore oil and gas developments, to
26 connect wells with other facilities and for export of processed hydrocarbons, and are usually
27 laid directly on the seabed. After the pipeline-laying process, on soft clay excess pore pressure
28 is present in the surrounding soil. As it dissipates, a significant rise in pipe-soil resistance occurs
29 reflecting the increase in effective stress. This has an influence on the global stability of the
30 pipeline, including the lateral buckling and axial pipeline walking in response to thermal cycles
31 during operation. The same mechanism affects the capacity of shallow subsea foundations,
32 which rises after installation due to consolidation (Gourvenec et al. 2014).

33 Early work into the post-laying consolidation around pipelines by Gourvenec & White (2010)
34 and Krost et al. (2011) considered elastic, uniform soil and a smooth pipe-soil interface.
35 Coupled consolidation finite-element analyses were later presented by Chatterjee et al (2012)
36 using Modified Cam Clay model in ABAQUS by means of both large-deformation finite-
37 element (LDFE) and small-strain finite-element (SSFE) methods. The effects of embedment,
38 rough and smooth interface conditions and large deformations associated with the penetration
39 were investigated. These solutions have practical value in allowing prediction of the
40 consolidation-induced changes in bearing capacity of a pipeline, and the increase in pipe-soil
41 interface friction-build-up to be predicted. These allow the pipeline design to be optimised.

42 However, practical application of these solutions requires an estimate of the coefficient of
43 consolidation of the shallow near-surface soils, typically at a depth of < 0.5 m. Conventional
44 site investigation tools such as the cone penetrometer are not suited to these near-surface
45 conditions, as ~~it~~the dissipation process consolidation is governed by the dissipation drainage
46 towards ~~in~~ the far field, without the influence of permeable top soil surface (Chatterjee et al.,
47 2014). Yan et al. (2010, 2011) proposed a new class of shallow ball- and toroid-shaped
48 penetrometers specifically for investigating shallow seabed properties and determined bearing
49 factors for undrained penetration, allowing strength profiles to be back-calculated from
50 penetration resistance.

51 This paper ~~follows from~~ extends the previous studies into to present coupled analysis of
52 undrained penetration ~~then and~~ consolidation ~~for around~~ a pipe, by exploring the behaviour of
53 and the toroid and ball penetrometers, using the Modified Cam Clay model in ABAQUS. The
54 main aim is to explore and compare-quantify the consolidation characteristics of shallowly
55 embedded objects, in terms of the time-scale for consolidation, with the aim of allowing simple

56 scaling from the penetrometer results to ~~the~~ pipeline and foundation behaviour. The derived
 57 solutions can be used to interpret dissipation results from this new class of penetrometers, to
 58 provide estimates of the consolidation parameters. These interpretations therefore unlock a new
 59 method to accurately determine near-surface consolidation parameters to support pipeline and
 60 may also be applicable to the design of shallow foundations such as the steel mudmats used to
 61 support subsea infrastructure design.

62 A range of variables are allowed for, including for embedment depth (expressed as the depth of
 63 the invert of the pipe or penetrometer, w , normalised by the diameter D), over-load ratios
 64 (OLRs) relevant for field situations (Jewell and Ballard 2011; White et al. 2011), pipe interface
 65 roughness (extreme cases for fully smooth and rough) and consolidation coefficient profile (c_v
 66 is either uniform or increasing proportionally with depth according to the effective stress level).
 67 The OLR is the ratio between the vertical load applied to the seabed during consolidation (i.e.
 68 the submerged self-weight for the case of the pipeline), W , divided by the initial undrained
 69 bearing capacity (i.e. maximum penetration resistance) at that depth, V_{\max} .

70 **Kinematic mechanisms during undrained penetration and subsequent consolidation**

71 The study assumed a wished-in-place pipe, toroid or ball geometry with embedment ratio w/D
 72 ranging from 0.1 to 0.5. For each embedment depth, the pipe was displaced vertically by $0.1D$
 73 in an undrained manner in order to mobilise the bearing capacity at the pre-embedded depth.
 74 The specified overloading ratio, $OLR = V_{\max}/W$ (considering values of 1, 4 and 12) was then
 75 achieved by reducing the vertical load, which established the initial excess pore pressure
 76 distribution. The subsequent consolidation response was then examined, quantifying the time-
 77 related excess pore pressure dissipation.

78 During the whole consolidation responses, vertical equilibration on the pipe must be satisfied
 79 (Figure 1), so that:

$$80 \quad \int \sigma'_{ni} \cos \theta \delta A + \int \Delta u_{N,i} \cos \theta \delta A + \int \tau_{fi} \sin \theta \delta A = W \quad (1)$$

81 where θ is the inclination from the vertical, δA is a local element of surface area and σ'_{ni} , τ_{fi}
 82 and $\Delta u_{N,i}$ are the local effective contact stress, shear stress and local effective vertical stress and
 83 excess pore pressure (hydrostatic pressure being ignored) respectively. The three components
 84 are integrated over the surface area of the embedded objects, balancing the resultant vertical
 85 loading W .

86 **FINITE ELEMENT ANALYSES**87 **Soil Model and Parameters**

88 The soil was modelled using Modified Cam clay (Roscoe and Burland 1968), as implemented
89 in ABAQUS. The soil response was taken as linear elastic before yielding. All parameters used
90 for the numerical analyses are listed in Table 1. The selected soil parameters were chosen to be
91 similar to those measured for kaolin clay used for centrifuge model tests by Stewart (1992) and
92 House et al. (2001). For more detailed discussion refer to Lu (2004).

93 A difficulty when using the MCC model is to define a unique c_v , to normalise dissipation
94 processes and quantify the average consolidation characteristics. During the consolidation
95 response, the soil volumetric stiffness ($1/m_v$) changes with mean effective stress (and whether
96 the soil is loading or unloading) and hence the consolidation coefficient varies with the mean
97 effective stress and load path.

98 For convenience, an initial invert value of c_v is adopted for normalisation, where the c_v value
99 is expressed using the initial soil state and (plastic) isotropic compressibility, m_v as

$$100 \quad c_v = \frac{k}{m_v \gamma_w} = \frac{k(1+e_0)p'_0}{\lambda \gamma_w} \quad (2)$$

101 where k is permeability, and e_0 (initial void ratio) and p'_0 (initial effective stress) are taken as
102 the virgin (undisturbed) values at the depth of the object invert, prior to penetration.

103 In order to investigate how the timescale for consolidation varies with the distribution with
104 depth of c_v , two separate series of analyses were undertaken:

105 (a) homogenous case: with an artificial surcharge of 200 kPa applied at the soil surface
106 (including on top of the embedded part of the pipe, toroid or ball) (Figure 2), giving an
107 approximately uniform value of c_v within the soil domain.

108 (b) linear case: with a very small surcharge of 0.001 kPa, giving essentially a linear increase of
109 c_v with depth.

110 Comparison of these two series allowed assessment of the effective c_v for the latter case in
111 order to obtain similar consolidation timescale as for the homogeneous case.

112 In all analyses the soil was initially K_0 -consolidated (Wroth 1984), with K_0 given by

$$113 \quad K_{0nc} = 1 - \sin \phi'_{tc} = 0.6 \quad (\phi'_{tc} = 23.5^\circ) \quad (3)$$

114

115 where φ'_{tc} is the friction angle for triaxial compression conditions.

116 In situ effective stresses and pore pressures vary with depth according to the respective self-
117 weights. The initial size of yield locus is determined by p'_c , expressed as

$$118 \quad p'_c = \frac{q_0^2}{M^2 p'_0} + p'_0 \quad (4)$$

119 where p'_0 , and q_0 are the initial effective mean stress and deviatoric stress, respectively, at a
120 given depth. The initial void ratio e_0 , can be calculated from

$$121 \quad e_0 = e_N - \kappa \ln p'_0 - (\lambda - \kappa) \ln p'_c \quad (5)$$

$$122 \quad e_N = e_{cs} + (\lambda - \kappa) \ln(2)$$

123 where κ and λ are the usual swelling and compression indices in MCC.

124 For these initial conditions, the starting point of the analyses for a given depth is denoted by
125 'O' in $p' - q$ and $e - \ln p'$ spaces, as shown in Figure 3 (a). The stress path to reach critical state
126 for an element that is sheared during undrained penetration is denoted by OB.

127 Figure 3 (b) shows the regular Tresca hexagon plotted in the deviatoric plane and von Mises
128 circle (MCC failure criterion on the deviatoric plane). The comparison of the modelling using
129 two yield criteria (Tresca and MCC) provides an indication of the mode of deformation during
130 failure.

131 When triaxial conditions dominate the failure mechanism (Lode angle $\theta = \pm 30^\circ$),
132 $s_{utc_MCC} = s_{u_Tresca}$. When plane strain conditions dominate the failure mechanism (Lode angle
133 $\theta = 0^\circ$), $s_{ups} = 1.15 s_{u_Tresca} = 1.15 s_{utc_MCC}$.

134 The critical state corresponds to the following internal friction angles:

$$135 \quad \sin \varphi'_{tc} = \frac{3M}{6 + M} \quad (6)$$

$$136 \quad \sin \varphi'_{te} = \frac{3M}{6 - M}$$

137 When plane strain conditions prevail, with $\sigma'_2 = 0.5(\sigma'_1 + \sigma'_3)$, the critical state corresponds to
138 an internal friction angle of:

$$139 \quad \sin\phi'_{ps} = \frac{\sigma'_1 - \sigma'_3}{\sigma'_1 + \sigma'_3} = \frac{q}{\sqrt{3}p'} = \frac{M}{\sqrt{3}} \quad (7)$$

140 where M is the slope of the critical state line, and ϕ'_{tc} , ϕ'_{te} , ϕ'_{ps} , are the friction angles under
141 triaxial compression, triaxial extension and plane strain conditions.

142 For soil with $M = 0.92$, as assumed in this study, these equations lead to interface friction
143 coefficients for a rough interface of 0.53, which corresponds to shearing in plane strain
144 conditions.

145 For triaxial compression conditions, the undrained shear strength ratio s_u/σ'_v for K_0
146 consolidated soil can be calculated from the MCC parameters using (Wroth 1984)

$$147 \quad \frac{s_{ut}}{\sigma'_v} = \frac{\sin\phi'_{tc}}{2a} \left(\frac{a^2 + 1}{2} \right)^\Lambda \quad (8)$$

148 For plane strain conditions the undrained shear strength ratio can be expressed as

$$149 \quad \frac{s_{ups}}{\sigma'_v} = \frac{2}{\sqrt{3}} \frac{\sin\phi'_{tc}}{2a} \left(\frac{a^2 + 1}{2} \right)^\Lambda \quad (9)$$

150

151 Where

$$152 \quad a = \frac{3 - \sin\phi'_{tc}}{2(3 - 2\sin\phi'_{tc})} \quad (10)$$

$$153 \quad \Lambda = \frac{\lambda - \kappa}{\lambda}$$

154 This leads to $(s_{ups}/\sigma'_{v0})_{nc}$ of 0.29 and mudline (plane strain) strengths of 0 (actually
155 0.00029) kPa and 57 kPa for the 0.001 and 200 kPa surcharge cases, respectively.

156 **Pipe, toroid and ball penetrometer properties**

157 The pipe, toroid and ball were modelled as rigid bodies with unit weight equal to the saturated
158 unit weight of the soil, which facilitated reaching equilibrium under the geostatic stresses. The
159 penetration resistance V in the subsequent step, which was applied as an external force to the
160 rigid body, therefore did not include any component of soil buoyancy. The interface conditions
161 considered were fully rough (soil bonded to pipe, toroid and ball) and fully smooth (zero shear
162 stress at pipe, toroid and ball surface), with pore water flow normal to the pipe surface always

163 set to zero.

164 The ratio between the outer and inner diameters of the toroid was 2. This ratio was identified
165 by Yan et al. (2011) as sufficient to practically eliminate interference between opposite sides
166 during undrained penetration.

167 **Finite element mesh**

168 Although a plane strain model would have been sufficient for the pipe model, the analyses were
169 undertaken using a slice (normal to the pipeline axis) of three-dimensional eight-noded
170 hexahedral elements, with multiple constraints forcing an identical response of the
171 corresponding nodes on each lateral face of the slice, thus imposing longitudinally-uniform
172 conditions (Figure 4). The reason for using a three-dimensional model was that this model was
173 also used to explore axial motion of the pipe segment (Yan et al. 2014). One slice of the soil
174 domain for the pipe included 3602 elements.

175 Similarly, for the toroid and ball, the analysis was undertaken using a ten-degree-wedge of
176 eight-noded hexahedral elements, with multiple constraints forcing identical response of the
177 corresponding nodes on each circumferential face of the slice, thus imposing axisymmetric and
178 circumferentially-uniform conditions.

179 For all models, the soil domain extended $8D$ horizontally and $10D$ vertically from the centreline
180 of the embedded objects, with zero horizontal displacements on the lateral boundaries, zero
181 vertical displacement at the base and drainage allowed only at the upper surface. The ten-degree
182 soil models for the toroid and ball penetrometers comprised 2998 and 3310 elements
183 respectively. This method allowed the three dimensional problems to be modelled at
184 considerably reduced computational expense, by analysing only a small radial slice of the
185 model.

186 **Model and mesh validation**

187 The numerical FE model was validated in a step-by-step fashion to confirm the correct use of
188 the MCC soil model (for both surcharges of 200 kPa and 0.001 kPa). The mesh sensitivity using
189 the Tresca model is first validated against the published results, which shows sufficient
190 robustness (more details are provided in the next section). The same meshing strategy was
191 therefore adopted for the MCC soil model.

192 Figure 5 presents the undrained vertical capacity factors ($V/As_{u, \text{invert}}$, with A the projected area
193 in plan view, and $s_{u, \text{invert}}$ the plane strain invert shear strength calculated from the in situ profile

194 based on equations (8) and (9) for a deeply embedded object ($w/D = 0.5$ for pipe, toroid, and
195 ball) from the MCC models. These results are compared with the FE modelling for the
196 (inscribed) Tresca soil.

197 For all pipe and ball cases with the 0.001 kPa surcharge, the MCC model generally gives
198 consistent agreement with the Tresca model (though 4 to 7% higher). This suggests that the soil
199 in the plastic zone is mostly sheared under plane strain conditions.

200 All responses using the 200 kPa surcharge show a softer build-up of resistance, and do not quite
201 reach a plateau within the applied displacement of $0.1D$ for the rough cases. This is due to
202 differences in rigidity index (G/s_{u0}) for the MCC and Tresca model, which were ~ 150 and 333
203 respectively. A set of analyses ($w/D = 0.5$) undertaken using an identical G/s_u for the MCC and
204 Tresca model, gave a discrepancy of less than 4% at the plateau (reached within a displacement
205 of $0.1D$) for the pipe, and 7% for the ball. This also suggests that most soil in the plastic zone
206 is sheared under plane strain conditions for the different objects.

207 UNDRAINED PENETRATION RESPONSE

208 The limiting undrained penetration resistances, for embedment ratios of 0.1 to 0.5, are
209 compared with published values for pipe, toroid and ball in Table 2 (Randolph et al. 2000;
210 Randolph and White 2008; [Merifield et al., 2009](#); Yan et al. 2011). The penetration resistances
211 using the MCC soil model have been normalised by the relevant projected contact area in plan
212 view and the plane strain shear strength at invert level. ~~The results using a simple Tresca soil
213 model are also tabulated and show close agreement. Comparative results of the WIP (wished-
214 in-place) and PIP (pushed-in-place) analyses using a simple Tresca soil model are also shown.
215 The PIP results are up to 7% higher than the corresponding WIP results at shallower
216 embedment, and generally give close agreement at deep embedment.~~

217 The FE results for the nominal surcharge of 200 kPa for the three objects generally show
218 excellent agreement with rigid plastic limit analyses. For the nominal surcharge of 0.001 kPa,
219 the normalised resistances calculated from MCC are consistently higher than those calculated
220 from the Tresca soil model, around 6-8% for toroid and pipe, and around 3% for ball. ~~This
221 might be consistent with observations of centrifuge modelling on a toroid penetrometer (Yan
222 et al., 2011), that the Tresca model gives a conservative slight under-prediction (comparing the
223 smooth case with the centrifuge results).~~

224 These results provide further confirmation that the MCC model is performing correctly in
225 undrained conditions. The following consolidation analyses were undertaken following

226 undrained penetration displacement by $0.1D$.

227 **CONSOLIDATION RESPONSE**

228 **Pore pressure dissipation at object invert**

229 After penetration, the specified vertical load was applied, factoring the maximum penetration
230 load to reflect overloading, and then maintained constant while consolidation was permitted.

231 Contours of initial excess pore pressure normalised by the invert value for the pipe and ball
232 geometry are shown for $w/D = 0.5$ in Figure 6 and corresponding variations around the
233 periphery are shown in Figure 7. For the 0.001 kPa surcharge case, the excess pore pressure is
234 more concentrated towards the object invert, reflecting the increasing soil strength with depth,
235 and therefore the concentration of load at the object invert. In uniform soil conditions, the
236 excess pore pressure is almost uniform ($\pm 15\%$) over most of the surface of the embedded objects
237 ($0.4 < x/D < 0.4$). Figure 6 shows that the excess pore pressure field for the toroid follows very
238 closely that for the pipe, indicating that the adopted ratio of the internal and external toroid
239 diameters is adequate to eliminate interference between opposite sides. The ball penetrometer
240 has a more compact excess pore pressure field due to the three dimensional geometry. This
241 gives a shorter drainage path for a given embedment.

242 Figure 8 shows variations of excess pore water pressure at the invert of the objects, normalised
243 by the initial value, with non-dimensional time T . The invert pore pressure is relevant to the
244 interpretation of penetrometer tests, because the ball and toroid penetrometers are equipped
245 with pore pressure transducers at this position.

246 The majority of the results are well fitted by simple hyperbolic equations (shown as solid lines)
247 in the form of

$$248 \quad \frac{\Delta u}{\Delta u_i} = \frac{1}{1 + (T / T_{50})^m} \quad (11)$$

249 where T_{50} is the value of T for 50% dissipation and m is a constant.

250 It can be seen that, for a rough embedded object, there is an initial increase in invert excess pore
251 pressure for all embedment ratios. This is due to the Mandel-Cryer effect (Cryer 1963; Mandel
252 1963) as discussed for consolidation beneath a pipe by Gourvenec and White (2010) and a
253 skirted foundation by Gourvenec and Randolph (2010). The effect is essentially determined by
254 comparison of the early rate of development of effective stress and excess pore pressures at the

255 invert and the soil at the edges of the object. For the rough interface, the excess pore pressure
256 is distributed more evenly on the interface, and the soil near to the edge (and hence the free
257 surface) consolidates more quickly than the invert soil, which leads to the Mandel-Cryer effect
258 at the invert. The effect is evident for each embedment, and for all rough objects. Due to the
259 Mandel-Cryer effect, the hyperbolic fit does not capture the initial portion of the dissipation
260 responses for rough objects. The smooth pipe (and other objects) are unaffected by the Mandel-
261 Cryer effect, as the excess pore pressure is more concentrated at the invert.

262 Table 3 summarises the fitting values of T_{50} and m for different embedment levels for all three
263 objects. The dissipation for the toroid follows closely the behaviour for the pipe. The normalised
264 time factors for dissipation for the ball range between 2 and 4 times lower than for the pipe or
265 toroid, with an average ratio of about three for a given degree of dissipation. The implied higher
266 rate of excess pore pressure dissipation around the ball is consistent with the smaller volume of
267 soil involved during penetration, and more effective drainage due to the three dimensional
268 geometry.

269 For the extreme embedment ratios ($w/D = 0.1$ and 0.5), results are also shown in Figure 8 for
270 0.001 kPa surcharge (with a high depth gradient of c_v). For those cases, the c_v values that give
271 a good match to the uniform cases at T_{50} and during the latter part of the consolidation curve
272 are higher than the invert value. The ratios χ of the operative value $c_{v,operative}$ to the invert value
273 are summarised in Table 4. The ratios lie in the range $2.5 - 4.0$, indicating more rapid dissipation
274 than if the entire soil domain had the same c_v value as at the invert. This is linked to the higher
275 c_v within the consolidating soil beneath the pipe invert and the different initial pore pressure
276 field. An alternative interpretation of this effect is to consider the depths at which the operative
277 c_v is found, which are also shown in Table 4.

278 **Average pore pressure dissipation around object surface**

279 The decay in the average excess pore pressure around the pipe periphery U_{av} and the
280 corresponding rise in normalised average normal effective stress Σ are useful quantities. They
281 are related to the ~~average~~-volumetric change of the soil adjacent to the objects, which indicates
282 the increase in shear strength due to reconsolidation after installation. This ~~potentially~~ reflects
283 the build-up of axial ~~or sliding~~ resistance between ~~the~~ embedded objects and the seabed (Yan,
284 2013). The factors U_{av} and Σ can be defined as~~The decay in the average excess pore pressure~~
285 ~~around the pipe periphery U_{av} and the corresponding rise in normalised average normal~~
286 ~~effective stress Σ are useful quantities since they indicate the build up of potential axial~~

287 ~~resistance between the embedded objects and the seabed. The factor U_{av} , and Σ can be defined~~
288 ~~as~~

$$289 \quad U_{av} = \frac{U}{U_i} = \frac{\int \Delta u \delta A}{\int \Delta u_i \delta A} \quad (12)$$

$$290 \quad \Sigma = \frac{\int \sigma'_n \delta A - \int \sigma'_{n,init} \delta A}{\int \sigma'_{n,f} \delta A - \int \sigma'_{n,init} \delta A} \quad (13)$$

291 where U is the integrated excess pore pressure Δu around periphery, σ'_n denotes the normal
292 effective stress, and $\sigma'_{n,av,init}$ and $\sigma'_{n,av,f}$ are the values before and after dissipation.

293 These trends are shown in Figure 9 for the embedded objects. The averaged pore pressure and
294 inverted effective stress responses agree to within 5% throughout the decay process, indicating
295 that the changes in total normal stress on the object surface are small. The pipe results calculated
296 from the MCC model are similar to the elastic results, but show more rapid dissipation as
297 consolidation progresses compared with the elastic solution, reflecting increasing stiffness as
298 the effective stress rises.

299 The majority of the pore pressure dissipation results shown in Figure 9 are well fitted by simple
300 exponent equations in the form of

$$301 \quad U_{av} = 1 - \Sigma = 0.5 \left(\frac{T}{T_{50}} \right)^n \quad (14)$$

302

303 where U_{av} and Σ are the average excess pore pressure, and average normal effective stresses
304 around the pipe periphery. T_{50} is the value of T for 50% dissipation, n is a constant (summarised
305 in Table 5).

306 The rough objects exhibit more consistent consolidation responses during the initial dissipation,
307 up to 20% dissipation around the pipe or toroid periphery and up to 30% dissipation at the ball
308 periphery. This consistent trend of decay of pore pressure is due to the evenly distributed excess
309 pore pressure at the rough interface. As consolidation progresses, the time for dissipation is
310 prolonged for increasing embedment, reflecting the variation in the drainage distance with
311 increasing embedment. The fitted curves show decreasing fitting parameter n with increasing
312 embedment ratio, reflecting this feature.

313 Figure 10 summarises values of T_{50} observed (a) for the invert pore pressure dissipation, and

314 (b) for the averaged perimeter dissipation, with increasing w/D for the three objects. This
315 provides a simple comparison of the relative rates of consolidation. The T_{50} values for the
316 penetrometer invert can be compared to those for a piezocone (Teh and Houlsby 1991), of
317 around 0.5 to 1, depending on the soil rigidity index. The surface penetrometers therefore
318 exhibit much shorter consolidation times than for a deeply embedded cone.

319 It can also be observed that the excess pore pressure at the invert of rough objects generally
320 dissipates more slowly than for smooth objects, due to the Mandel-Cryer effect. By contrast,
321 for the averaged excess pore pressure around the periphery, rough objects show faster
322 dissipation, reflecting the effect of the initial excess pore pressure field distribution.

323 ~~It is notable that t~~The parameters summarised in Table 3 and Table 5 provides a method to
324 transfer the invert response to the average response around the perimeter of the objects
325 perimeter response, allowing for assessment of build-up ofof the increase in potential axial
326 resistance between the objects and soil. The more preferable method may be amounting more
327 pore pressure transducers along the device periphery for a more direct measurement of the
328 average response.

329 **Effect of overloading ratio and object roughness**

330 Excess pore pressure dissipation responses for the toroid and ball penetrometers, with
331 overloading ratios of 1, 4, and 12, are illustrated in Figure 11 for the extreme embedment ratios
332 of 0.1 and 0.5. The dissipation responses for high overloading show dilatory behaviour, with
333 pore pressure increasing from the initial value to a maximum followed by a decrease to the
334 hydrostatic value.

335 The initial excess pore pressure field generated during undrained penetration has a comparable
336 extent and magnitude for a given object with a given embedment and interface condition,
337 irrespective of the overloading ratio applied. The rough objects result in a pore pressure field
338 more evenly distributed on the periphery, while the smooth objects result in a pore pressure
339 field more concentrated at the invert.

340 The overloading event led to the generation of negative excess pore pressure around the
341 embedded objects, but positive excess pore pressure remains in the far field. For most cases
342 with $OLR > 1$, an increase of excess pore pressure (swelling) was observed during the initial
343 period of time, as flow from the far field towards the periphery exceeds the rate of dissipation
344 from the periphery to the free drainage surface. The dissipation time decreases for increasing
345 overloading ratio reflecting this neutralisation of excess pore pressure. Although dissipation is

346 initially faster for the rough objects than the smooth, the time histories of consolidation soon
347 become closely banded and the consolidation responses for rough and smooth objects
348 eventually converge to similar time factors for full consolidation.

349 Additional illustration of these phenomena is provided by the stress paths depicted in $e - \ln(p')$
350 space as shown in Figure 12 (for overloading ratios of 1, 4, and 12 under surcharge of 200 kPa).
351 To aid interpretation of the stress paths, the states are denoted with superscript 1, 2, and 3 for
352 overloading ratios of 1, 4, and 12 respectively. The initial state is denoted by O at the in situ
353 effective stress ($p' = \sigma'_{v0}(1+2K_0)/3$), from which state the soil is loaded along an undrained
354 stress path during penetration from O to B₁. In the unloading step to establish the overload ratio,
355 the excess pore pressure at the interface falls significantly to balance the residual applied force,
356 while the effective stress remains virtually constant (remaining at B') in $e - \ln(p')$ space. The soil
357 in the far field is largely unaffected by this unloading event, and the effective stress and the
358 excess pore pressures remain at a similar magnitude as for the OLR = 1 case. This forms a
359 drainage front advancing towards the surface of the object and the soil mass. This in turn
360 increases the excess pore pressure at the object (B₁ to B₂ for case of OLR = 4, and B₂ to B₃ for
361 case of OLR = 12). With time, the process begins to reverse and the dissipation at the invert
362 begins (B₂ to C₂ for case of OLR = 4; B₃ to C₃ for case of OLR = 12).

363 COMPARISON WITH FIELD ~~AND NUMERICAL DATA AND AVAILABLE~~ 364 ~~PUBLISHED RESULTS~~

365 The dissipation curves from numerical analyses in this study and large-deformation finite
366 element (LDFE) analyses (Chatterjee et al., 2012) for a half embedded ($w/D = 0.5$), smooth
367 pipe in homogeneous case ~~is~~are illustrated in Fig. 13. Only minimal discrepancy ~~of~~between
368 them is ~~found~~evident, which provides ~~the~~validation of the numerical solution in this study.
369 The field data (the average excess pore pressure from four ~~invert-mounted~~ transducers²
370 ~~recording~~ plotted against elapsed time) is extrapolated from a published field test (Hill and
371 Jacob, 2008). This test was implemented offshore via the Fugro SMARTPIPE device. During
372 the test, ~~such~~a model pipe with a length of 1.1 m and a diameter of 0.225 m ~~was~~ first penetrated
373 to a depth of $0.6D$ in a soft clay seabed. It was then held under a constant vertical load with the
374 decay of excess pore pressure being recorded by four transducers spaced along the pipe invert.
375 To obtain good agreement between the field data and the numerical result, an operative $c_{v,operative}$
376 of $48 \text{ m}^2/\text{year}$ is adopted to normalise the field data ~~based on the uniform soil case. The invert~~
377 ~~$c_{v,invert}$ of this model pipe can also be estimated from suggested scaling factor γ presented in~~

378 ~~Table 4. The value of χ is ranged within 1–2.5 for the case of $w/D = 0.5$, herein the value of χ~~
379 ~~is chosen as 2.5, which yields~~Using the linear soil case, an invert value of the $c_{v, \text{invert}}$ of
380 ~~20 m²/year applies. It can be seen the dissipation response of field data normalised by such~~
381 ~~$c_{v, \text{invert}}$ follows close agreement with numerical result of normally consolidated case~~This is a
382 ~~relatively narrow range of uncertainty, which could be reduced if the penetration resistance data~~
383 ~~was available, allowing the appropriate soil profile to be selected. Using this back calculated~~
384 ~~$c_{v, \text{invert}}$, the consolidation degree is around 50% after 2700 s, which corresponds to~~based on a
385 ~~T_{50} of 0.022; and 99% after 122035s, which corresponds to completion of primary~~
386 ~~consolidation.~~

387 The comparison is also extended to the available published solutions of measuring for
388 consolidation coefficient by based on different types of devices. Apart from In addition to the
389 results of toroid and ball penetrometers from this present study, the dissipation curves obtained
390 by the strain path method for the conventional cone penetration test (CPT) (Teh and Houlsby,
391 1991), and simulated by the coupled large-deformation finite element (LDFE) analyses for the
392 parkable piezoprobe test (PPP) (Chatterjee et al., 2014) in homogeneous clay are presented Fig.
393 14. The aforementioned dissipation responses for a pipe are also included as a reference for
394 comparison. The decay process of the toroid penetrometer presented in this paper here shows a
395 good agreement with the pipe results of pipe, due to their similar geometry. For the same
396 diameter, the toroid and pipe shows, but a faster response than the result of CPT, and most
397 importantly a different shape of response. This shows that the presence of the permeable top
398 surface alters the shape of the dissipation response as well as the overall rate, emphasising the
399 importance of using device-specific dissipation solutions to interpret the different types of
400 test implies the CPT data from deep penetrated position cannot reflect the speed-up drainage
401 rate due to the permeable top surface of soil.

402 The ball penetrometer performs shows a similar dissipation response as to the PPP. Both show
403 a faster decay compared to toroid and pipe for the same D, which indicates a more rapid
404 determination of $c_{v, \text{operative}}$ compared to the toroid penetrometer or pipe of the same diameter.
405 The value of $c_{v, \text{invert}}$ can also be predicted following a similar process of field data described
406 above.

407 These two values can be directly used for pipe design, as the scaling factor χ of ball
408 penetrometer is almost same compared to the pipe for each case (Table 4). Note that when
409 these results show measurement using that the hemiball penetrometer provides a rapid
410 method of is used to estimating the dissipation response of around a pipe, given the differences

411 ~~in the values of T_{50} (and m) need to be changed from that of ball penetrometer to pipe using~~
412 ~~suggestion of (Table 3). For instance, the dimensionless time scaling factors $T_{50} = 0.033$ for~~
413 ~~and $m = 1.3$ of the smooth ball penetrometer ($w/D = 0.5$) are required to be changed to almost~~
414 ~~three times ~~quicker~~ shorter than $T_{50} = 0.082$ and $m = 1.05$ of the smooth pipe ($w/D = 0.5$).~~

415 Throughout the comparison, the numerical solutions of pipe, toroid and ball penetrometers ~~infor~~
416 ~~the case of homogeneous ~~ease~~ soil, as presented here, in this paper~~ can be used to determine
417 $c_{v,operative}$ in their corresponding field tests. The estimation of $c_{v,invert}$ can also be achieved with
418 the aid of suggested scaling factors χ , which ~~is~~ span a relatively narrowed ~~into a small~~ range by
419 the numerical solutions ~~of them infor the~~ normally consolidated case. The efficiency in
420 estimating consolidation coefficient via ball penetrometer might be highlighted since the cost
421 of conducting such a field test ~~is~~ depends mainly ~~dependent~~ on the vessel time.

422 LIMITATIONS

423 Although the numerical solutions reported in this paper have been capable of determination of
424 the consolidation degree through a back-calculation of c_v , the effect of higher hydraulic
425 conductivity around the interface arising from the roughness and asperities of the pipe coating
426 ~~should be considered around~~ may have an influence at the pipe-soil interface (Jewell and
427 Ballard, 2011). A special drainage or consolidation condition along the interface could be
428 ~~introduced to avoid the unexpectedly~~ the cause of a higher coefficient of consolidation being
429 deducing from ~~fitting~~ field data (i.e. $c_{v,operative}$ of $36 \text{ m}^2/\text{year}$ in this study) relative to laboratory
430 tests, using devices such as the Rowe cell. However, such an effect may also exist at the surface
431 of a pipe, in which case the observed dissipation rate on the penetrometer is realistic for design.
432 ~~Also~~ Finally, the process of pipe installation ~~is~~ has been regarded as a monotonic penetration
433 followed by a consolidation, without the consideration of cyclic behaviour. ~~A more~~
434 ~~sophisticated hyperplasticity with a non-linear kinematic hardening based on the modified Cam~~
435 ~~clay soil model (Likitlersuang and Houlsby, 2006; Apriadi et al., 2009) may be adopted to avoid~~
436 ~~non-conservative elastic response (Houlsby et al., 2005)~~ It is possible that the consolidation rate
437 around a pipeline may be altered by a dynamic component of the installation process which
438 remoulds the surrounding soil and alters the initial pore pressure field. However, centrifuge
439 model testing shows that this effect is minimal (Cocjin et al. 2017).

440 CONCLUSIONS

441 This paper presented numerical results based on the Modified Cam clay model to investigate

442 the consolidation process after partial embedment of a pipeline and of shallow toroid and ball
443 penetrometers. This is an important consideration for design as pore pressure dissipation
444 governs the rate at which pipeline axial friction develops. These novel shallow penetrometers
445 offer an efficient basis to determine the relevant consolidation rates directly in situ. The effects
446 on consolidation rate of embedment, object-soil interface conditions and different overloading
447 ratios have been investigated.

448 For both smooth and rough pipes, toroids and balls, consolidation time increased with
449 increasing initial embedment, and was greater for the rough interface condition. An initial
450 increase in excess pore pressure was observed at the invert for rough embedded objects due to
451 the Mandel-Cryer effect. Simple hyperbolic or exponential equations were fitted to the
452 dissipation curves both at the object invert and averaged over the surface.

453 These results now provide an interpretation method for shallow ball and toroid penetrometers
454 to determine the consolidation properties of soft soils, giving these new tools practical value.
455 Also, the resulting values of c_v can be converted into average rates of pore pressure dissipation,
456 to assess the rate of effective stress recovery – for example to predict the build-up of friction
457 on seabed pipelines.

458 The consolidation responses for a toroid penetrometer generally show excellent agreement with
459 those for an infinitely long pipe, confirming that the adopted toroid shape – specifically the
460 ratio of internal and external diameters – is devoid of interaction effects. The shallow ball
461 penetrometer shows a faster consolidation response, typically by a factor of 3, reflecting the
462 more effective drainage mechanisms of a three dimensional device compared to a plane strain
463 device. The toroid dissipation response is therefore more directly applicable in pipeline
464 analysis, once the relative diameter of the two objects is accounted for. On the other hand, the
465 ball provides a more rapid determination of c_v , which offers improved time efficiency if
466 required during the survey operations.

467 The dissipation responses were also compared with those from elastic solutions, highlighting
468 the effects of different initial excess pore pressure distribution and some stiffness increase
469 during consolidation arising from the MCC model.

470 It is anticipated that these solutions will allow the hemiball and toroid penetrometer to gain
471 practical acceptance as improved tools for characterising the near-surface properties of soft
472 soils.

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548 **Table 1 Input parameters of numerical study**

Soil property Parameters	Values
Slope of critical state line (CSL) in p' - q space, M (friction angle in triaxial compression, ϕ'_{tc})	0.92 (23.5°)
Void ratio at $p' = 1$ kPa on (CSL), e_{cs}	2.14
Slope of the virgin compression line in e - $\ln(p')$ space, λ	0.205
Slope of the swelling and recompression line in e - $\ln(p')$ space, κ	0.044
Elastic shear modulus, G	$50p_0'$
Saturated bulk unit weight, γ_{sat} : kN/m ³	15.0
Unit weight of water, γ_w : kN/m ³	10
Permeability of soil, k : m/s	1.0×10^{-9}
Pipe/Toroid/Ball diameter, D : m (although all results are presented in non-dimensional form)	0.5

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Table 2 Summary of (fully mobilised) undrained bearing capacities for pipe, toroid and ball

Objects	<i>w/D</i>	This study								References								
		Tresca				MCC				Tresca								
		FE				FE				Upper bound ^{c,d}				FE ^e		FE ^f		
		Homo		Non-homo		Homo ^a		Non-homo ^b		Homo		Non-Homo		Homo		Homo		
		Rough	Smooth	Rough	Smooth	Rough	Smooth	Rough	Smooth	Rough	Smooth	Rough	Smooth	Rough	Smooth	Rough	Smooth	
Pipe	0.10	3.17	2.77	3.78	3.00	3.24	2.84	4.02	3.24	3.35	2.79	3.88	3.02	3.27	3.32	2.98		
	0.20	4.18	3.48	4.34	3.38	4.28	3.57	4.64	3.61	4.20	3.51	4.43	3.38	4.24	4.17	3.54		
	0.30	4.80	3.85	4.70	3.62	4.91	3.94	5.03	3.88	4.81	3.82	4.84	3.63	4.83	4.77	3.92		
	0.40	5.36	4.13	5.20	3.83	5.34	4.12	5.50	4.14	5.32	4.28	5.14	3.84	5.30	5.25	4.21		
	0.50	5.79	4.36	5.44	4.00	6.00	4.53	5.74	4.31	5.54	4.57	5.35	4.00	5.73	5.65	4.46		
Toroid	0.10	3.18	2.78	3.80	3.01	3.26	2.85	4.04	3.26					3.30				
	0.20	4.21	3.51	4.36	3.40	4.32	3.60	4.66	3.64					4.35				
	0.30	4.86	3.90	4.72	3.64	4.97	3.98	5.06	3.91					5.01				
	0.40	5.43	4.19	5.22	3.85	5.41	4.17	5.53	4.17					5.52				
	0.50	5.87	4.43	5.47	4.01	6.10	4.60	5.78	4.34					6.07				
Ball	0.10	2.17	1.76	1.70	1.31	2.14	1.77	1.75	1.37					2.16				
	0.20	3.85	2.92	2.69	2.05	3.81	2.92	2.80	2.18					3.91				
	0.30	5.19	3.84	3.50	2.63	5.10	3.77	3.63	2.73					5.37				
	0.40	6.25	5.60	4.18	3.10	6.08	4.46	4.34	3.29					7.02				
	0.50	7.16	5.12	4.80	3.51	6.94	4.96	4.96	3.71	7.65(UB) ^e				7.50				

^{a,b} Fully mobilised vertical capacity under 200 kPa and 0.001 surcharge;

^{c,d,e,f} Sources of solution: Upper bound solutions of pipe from Randolph and White (2008); ball from Randolph et al. (2000); FE solutions of pipe, toroid and

ball from Yan et al. (2011); FE solution for PIP (push-in-place) pipe results from Merifield et al. (2009).

Table 3 Values of T_{50} and constant m of hyperbolic fits to invert pore pressure dissipation

Initial embedment, w/D	Rough objects				Smooth objects			
	Pipe (toroid)		Ball		Pipe (Toroid)		Ball	
	T_{50}	m	T_{50}	m	T_{50}	m	T_{50}	m
0.1	0.028 (0.032)	1.05	0.012	1.3	0.022 (0.022)	1.05	0.005	1.3
0.2	0.055 (0.058)	1.05	0.018	1.3	0.040 (0.040)	1.05	0.012	1.3
0.3	0.072 (0.070)	1.05	0.026	1.3	0.056 (0.056)	1.05	0.018	1.3
0.4	0.095 (0.095)	1.05	0.032	1.3	0.072 (0.072)	1.05	0.025	1.3
0.5	0.110 (0.110)	1.05	0.042	1.3	0.082 (0.084)	1.05	0.033	1.3

Table 4 Operative c_v for different initial embedment values

Initial embedment, w/D	$\chi = c_{v,operative}/c_{v,invert}$		Depth of operative c_v (normalised by object diameter)		$\chi = c_{v,operative}/c_{v,invert}$		Depth of operative c_v (normalised by object diameter)	
	Smooth pipe (toroid)	Rough pipe (toroid)	Smooth pipe (toroid)	Rough pipe (toroid)	Smooth ball	Rough ball	Smooth ball	Rough ball
	0.1	2.60	4.20	0.28	0.46	2.60	4.20	0.28
0.2	2.70	3.00	0.58	0.64	2.70	3.00	0.58	0.64
0.3	2.60	3.00	0.84	0.98	2.60	3.00	0.84	0.98
0.4	2.40	2.80	1.02	1.20	2.40	2.40	1.02	1.02
0.5	2.50	2.80	1.33	1.50	2.80	2.80	1.50	1.50

Table 5 Values of T_{50} and constant n of exponents fits to periphery pore pressure dissipation

Initial embedment, w/D	Rough objects				Smooth objects			
	Pipe (Toroid)		Ball		Pipe (Toroid)		Ball	
	T_{50}	n	T_{50}	n	T_{50}	n	T_{50}	n
0.1	0.012 (0.014)	0.52	0.0038	0.58	0.015 (0.015)	0.52	0.004	0.60
0.2	0.026 (0.027)	0.50	0.0055	0.50	0.030 (0.030)	0.55	0.008	0.65
0.3	0.036 (0.035)	0.47	0.0070	0.49	0.044 (0.044)	0.6	0.012	0.66
0.4	0.040 (0.040)	0.44	0.0078	0.48	0.055 (0.055)	0.6	0.016	0.66
0.5	0.050 (0.050)	0.44	0.012	0.48	0.079 (0.080)	0.6	0.020	0.66

Figure captions

Figure 1 Schematic diagram of problem

Figure 2 Geometry and definitions (pipe case)

Figure 3 Schematic illustration of critical state model and undrained failure criteria

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Figure 8 Excess pore pressure dissipation time histories at invert for embedded objects
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Figure 13 Comparison of calculated and observed dissipation curves at pipe invert

Figure 14 Excess pore pressure dissipation time histories at invert for smooth pipe and
penetrometers in this study compared with CPT, pipe and PPP (Chatterjee et al., 2014) on
homogeneous soil

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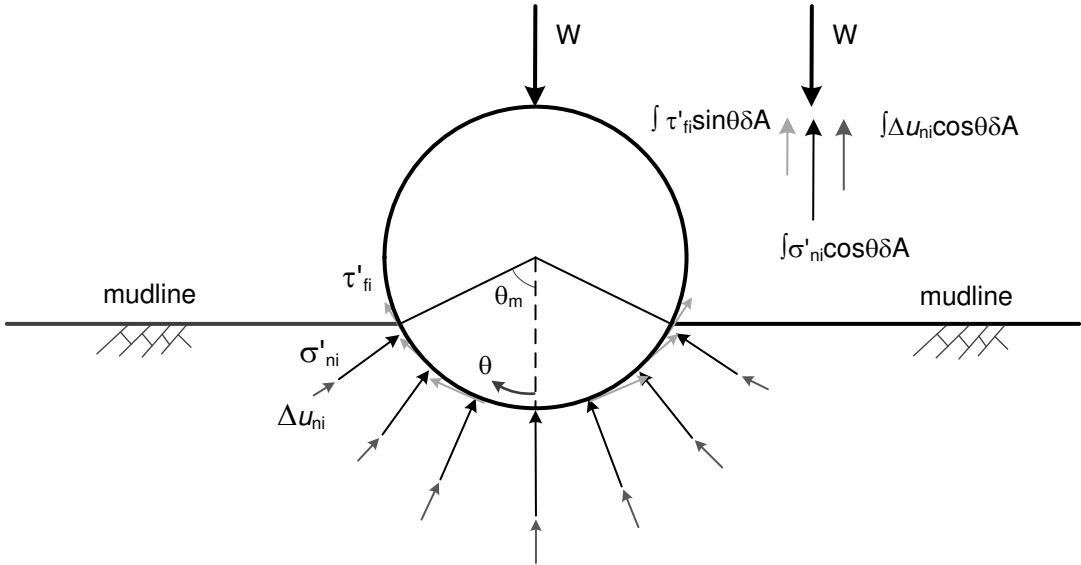
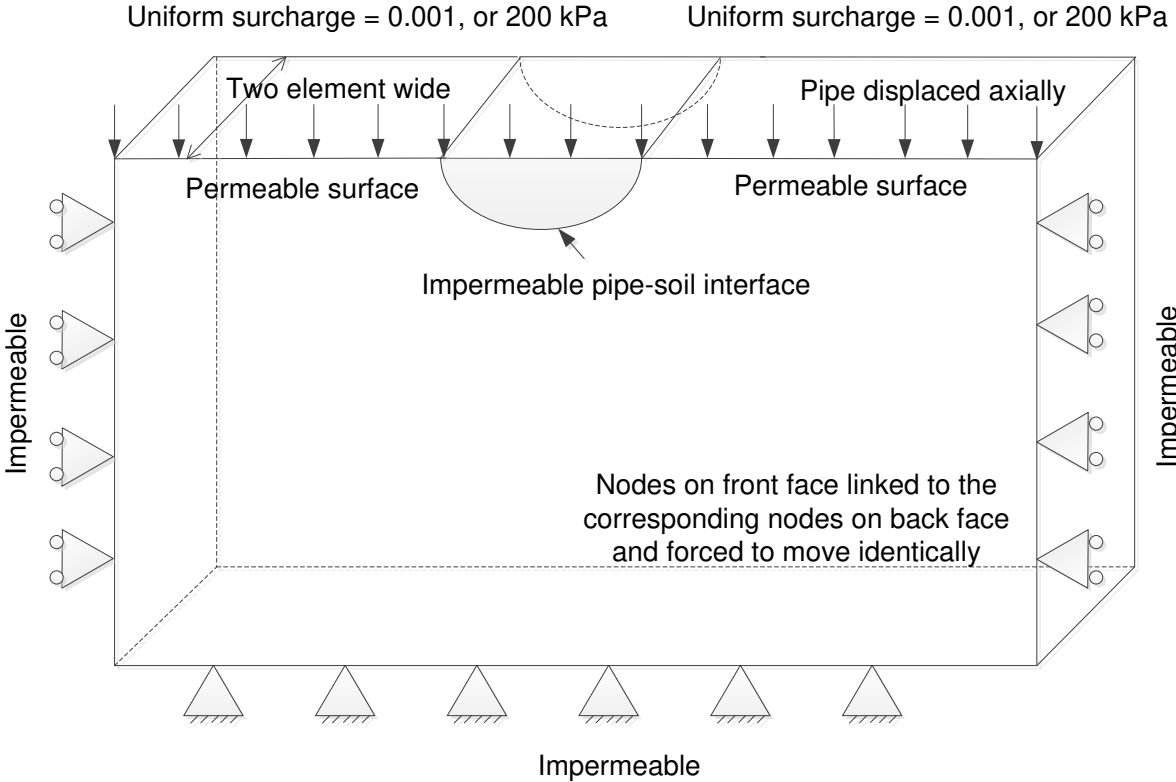
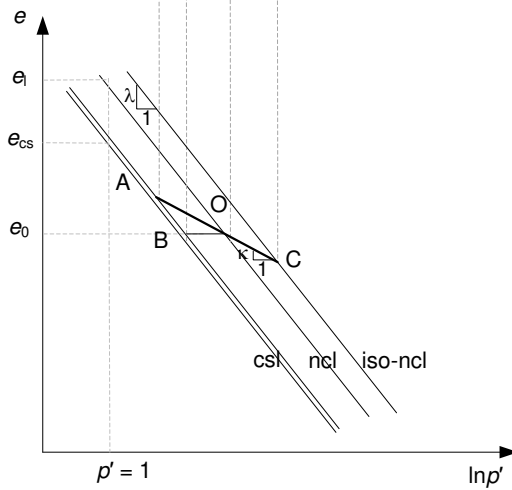
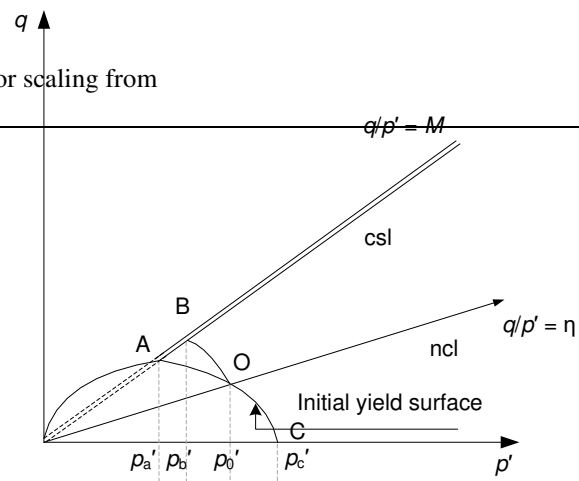


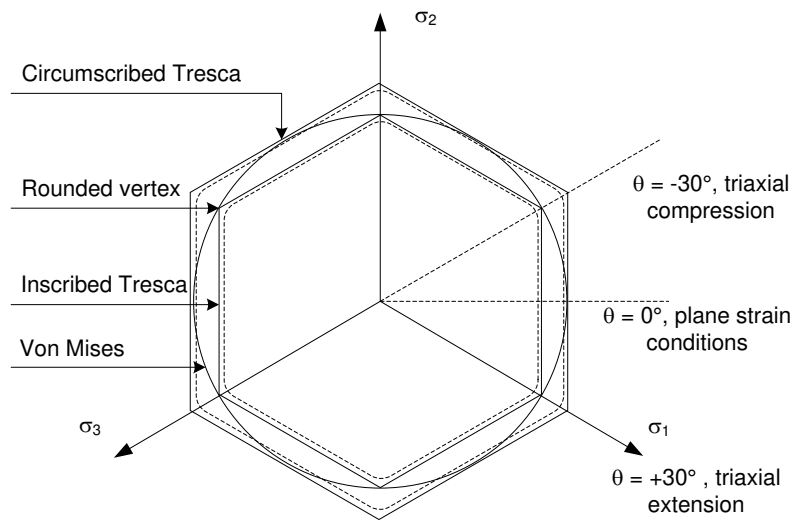
Figure 1 Schematic diagram of problem



[Figure 2 Geometry and definitions \(pipe case\)](#)~~Figure 2 Geometry and definitions (pipe case)~~



(a)



(b)

Figure 3 Schematic illustration of critical state model and undrained failure criteria

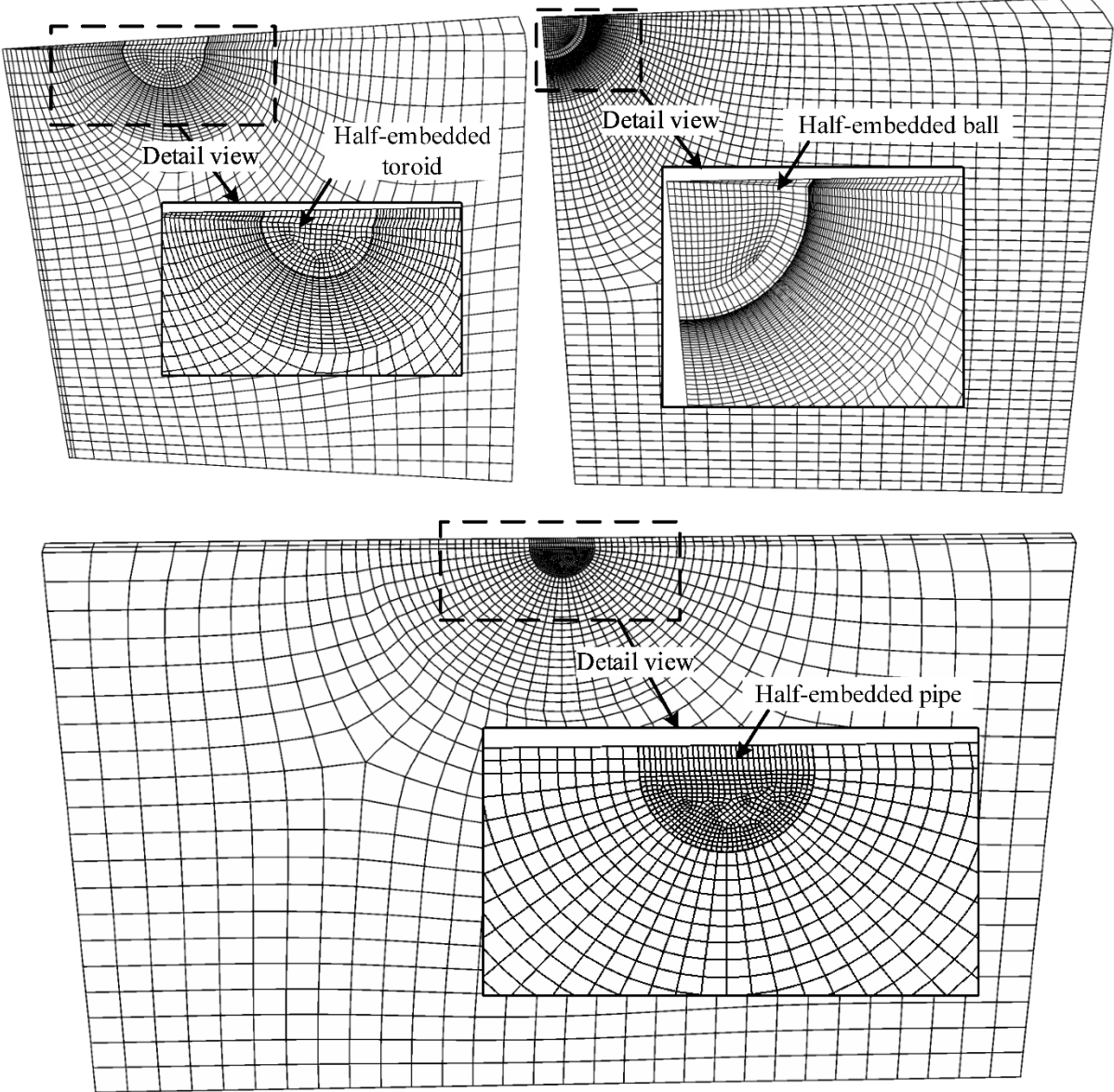
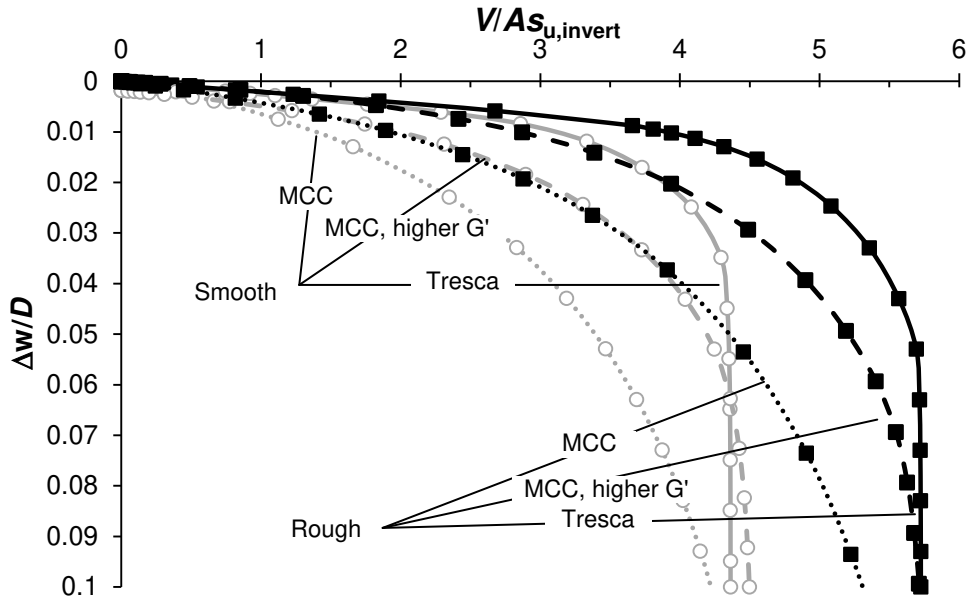
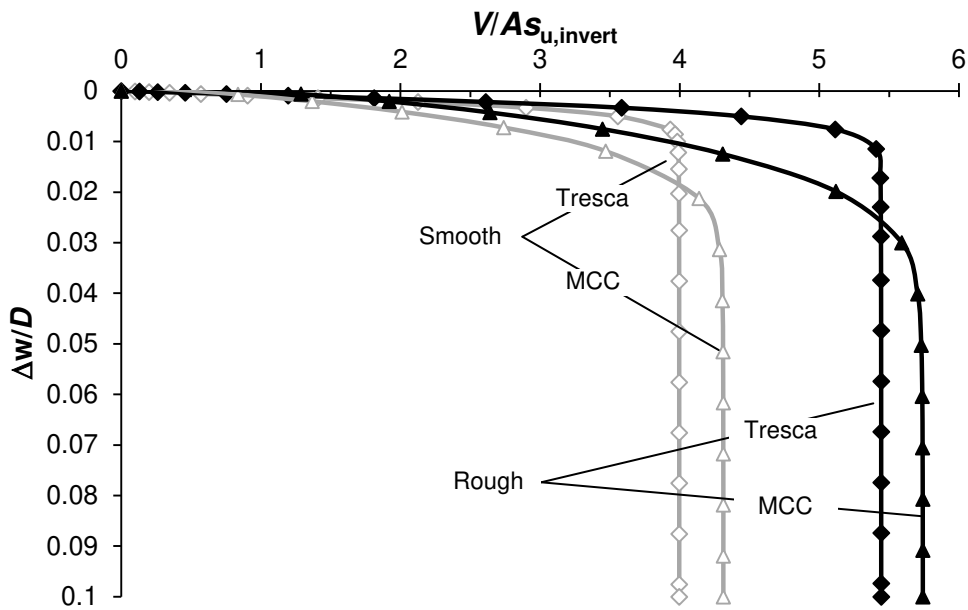


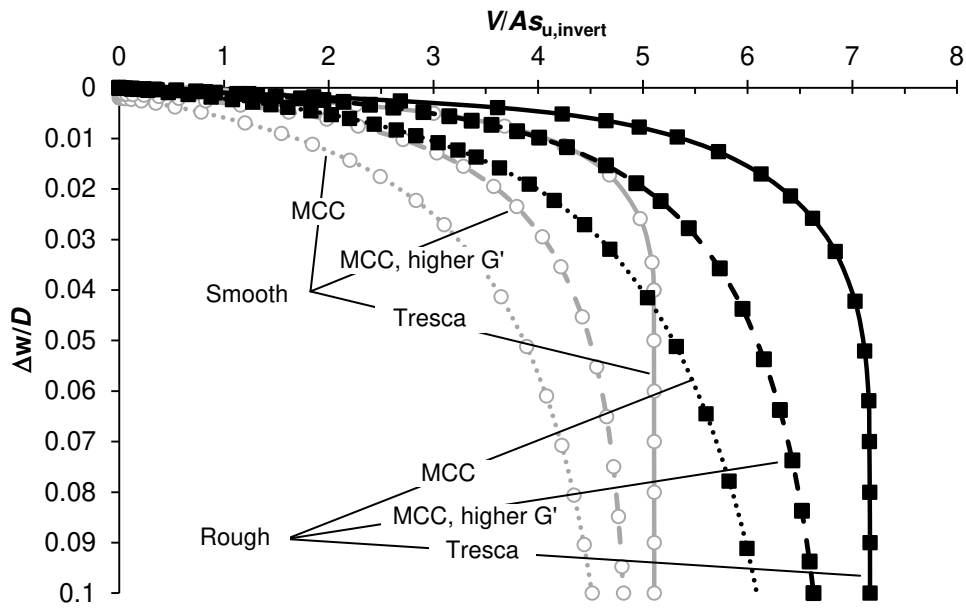
Figure 4 Finite element meshes (Oblique view of models) Figure 4 Finite element meshes (Oblique view of models)



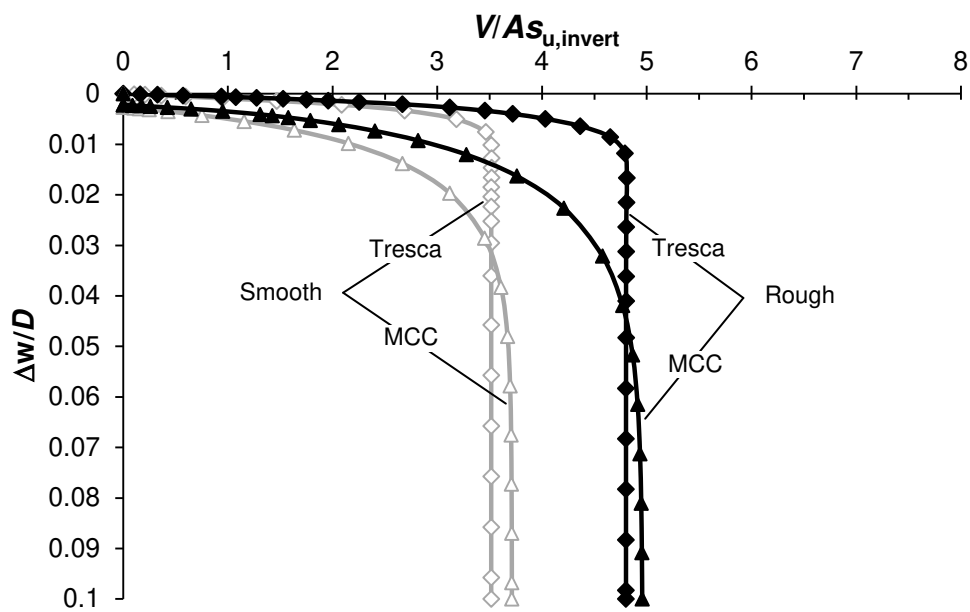
(a) Pipe, homogeneous strength



(b) Pipe, proportional strength



(c) Ball, homogeneous strength



(d) Ball, proportional strength

Figure 5 Mesh validation: Undrained penetration resistance ($w/D = 0.5$)

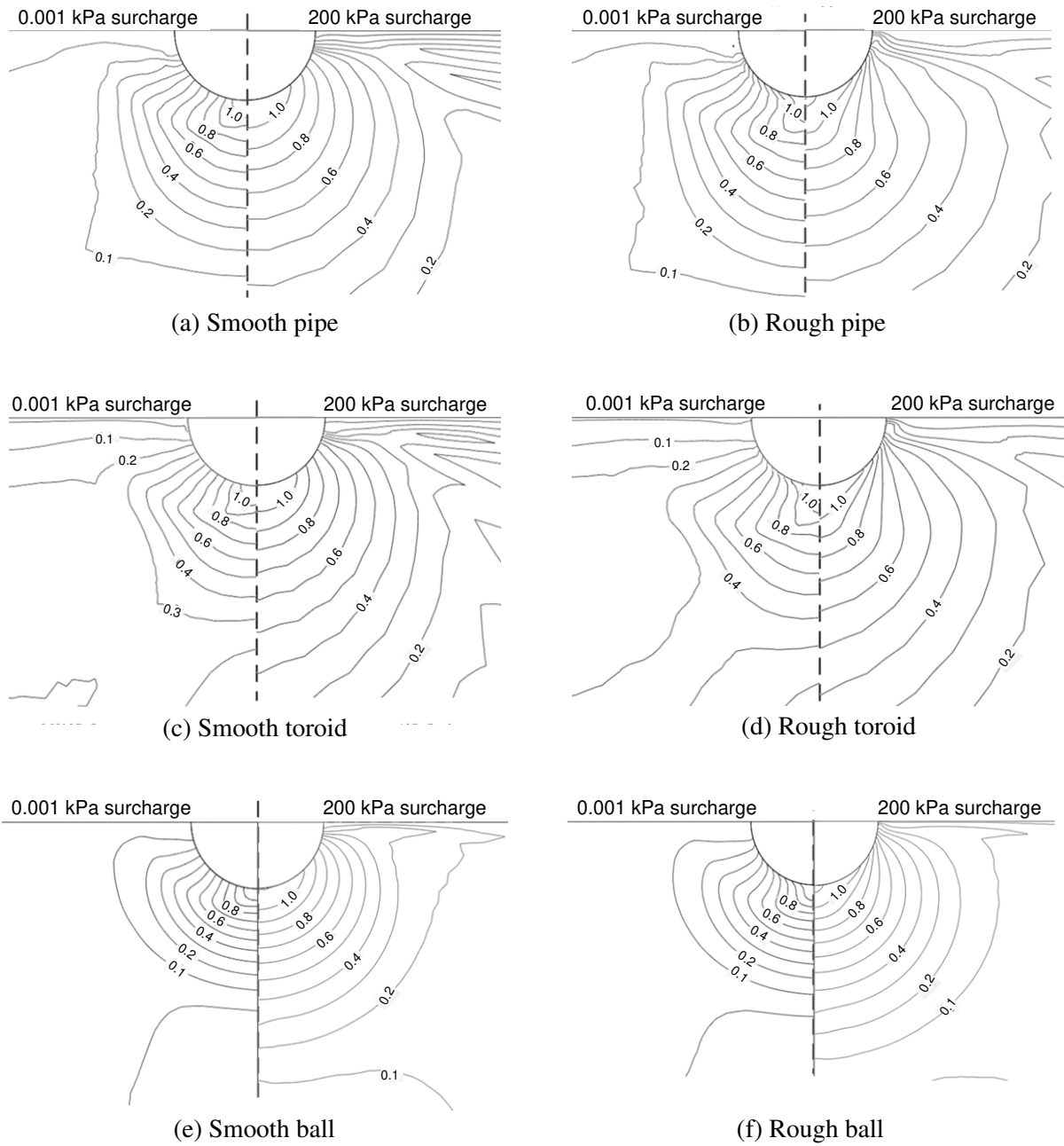


Figure 6 Excess pore pressure distributions after penetration

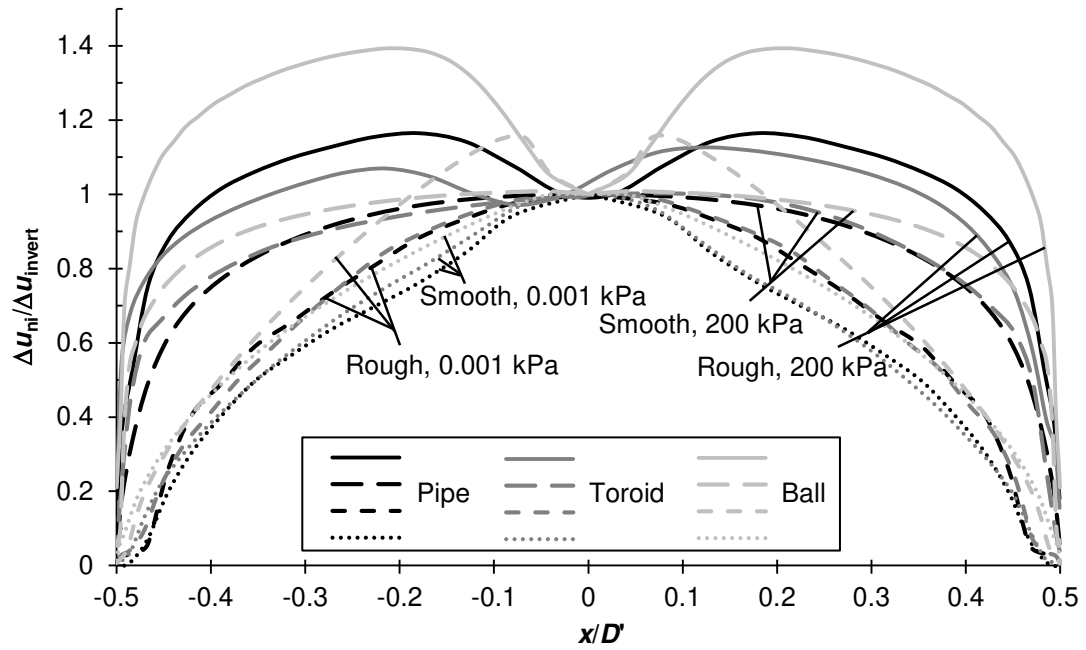


Figure 7 Excess pore pressure distributions around object periphery after penetration
Excess pore pressure distributions around object periphery after penetration

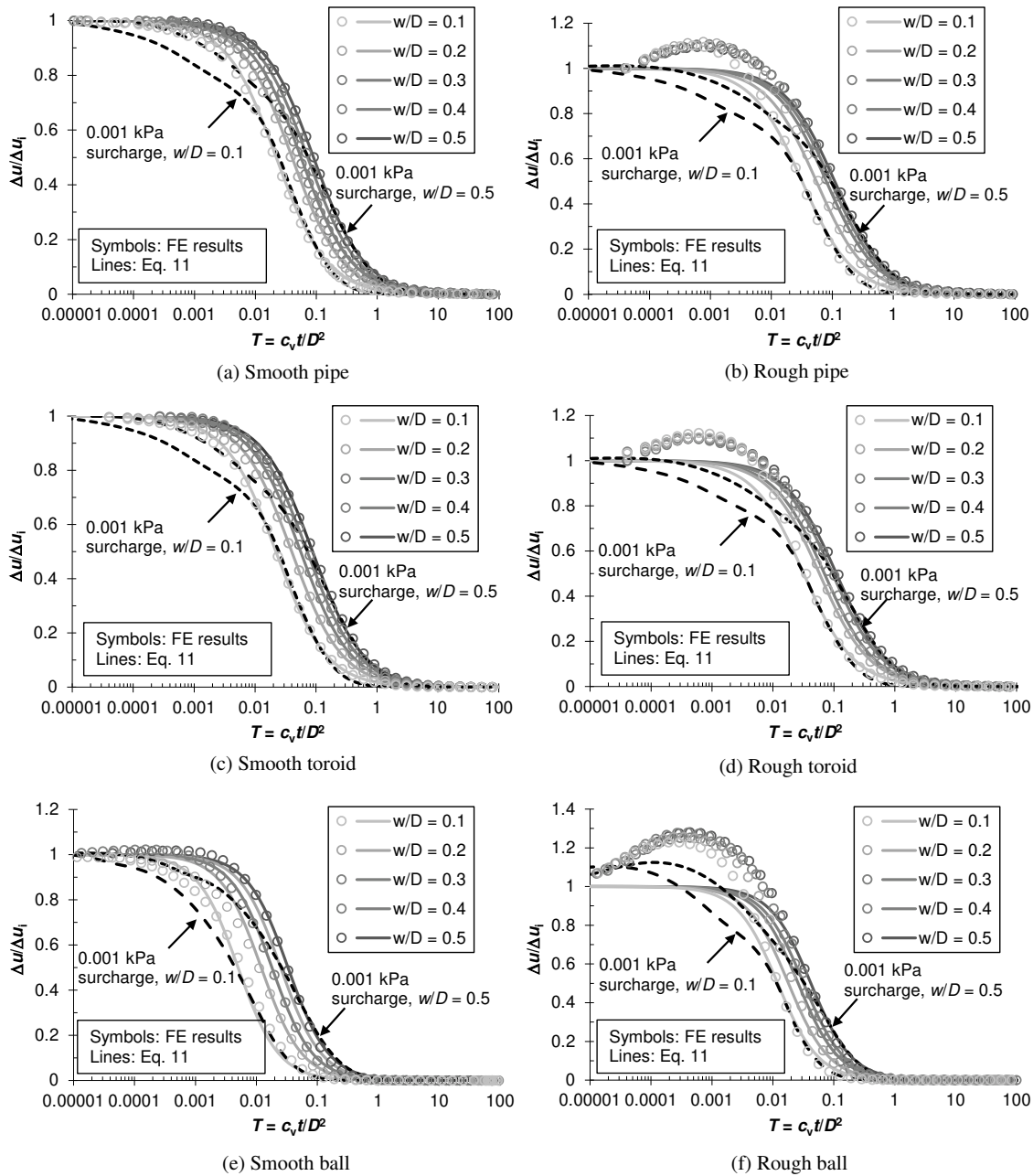


Figure 8 Excess pore pressure dissipation time histories at invert for embedded objects (200 kPa surcharge cases unless otherwise noted)

Figure 8 Excess pore pressure dissipation time histories at invert for embedded objects (200 kPa surcharge cases unless otherwise noted)

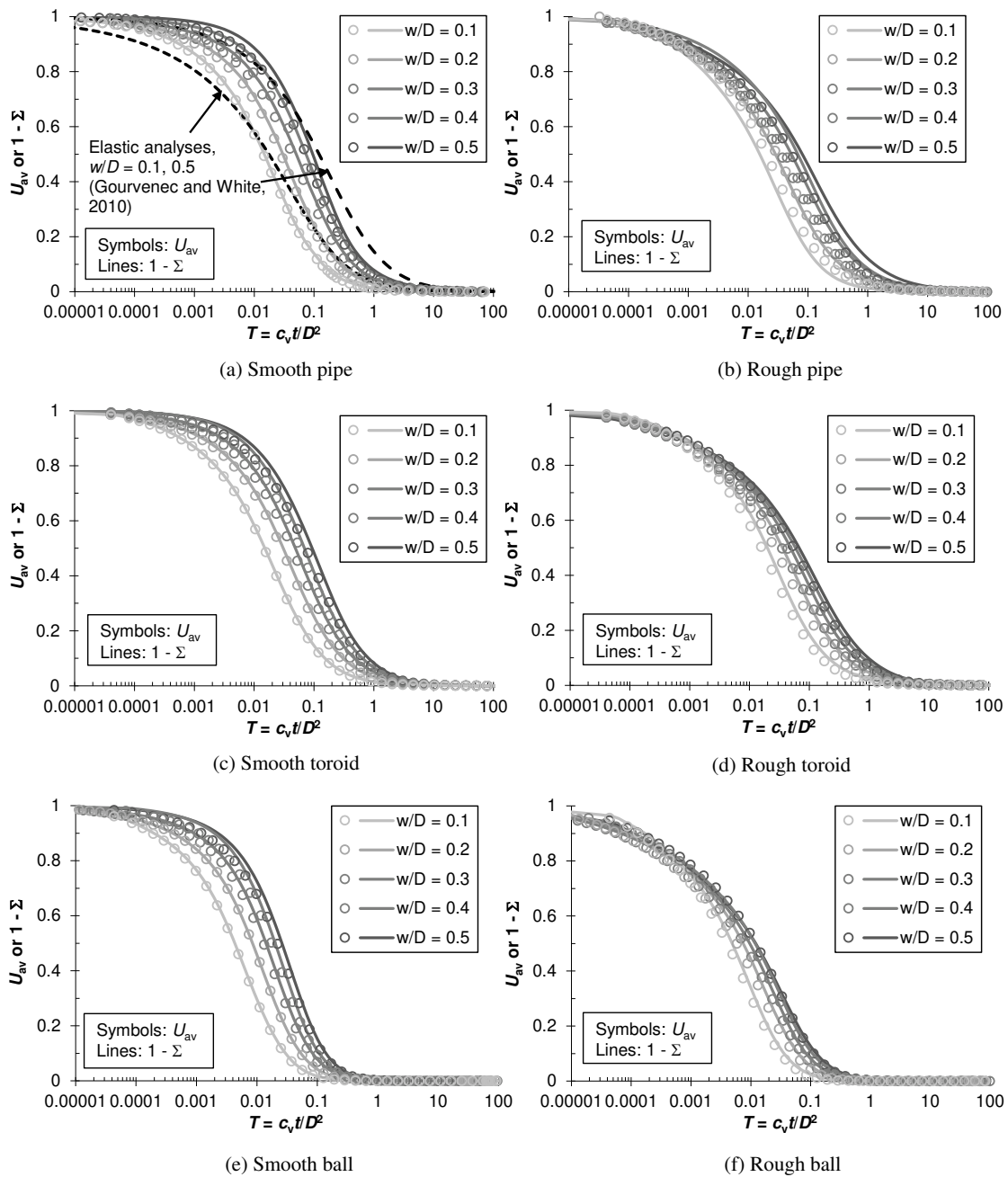
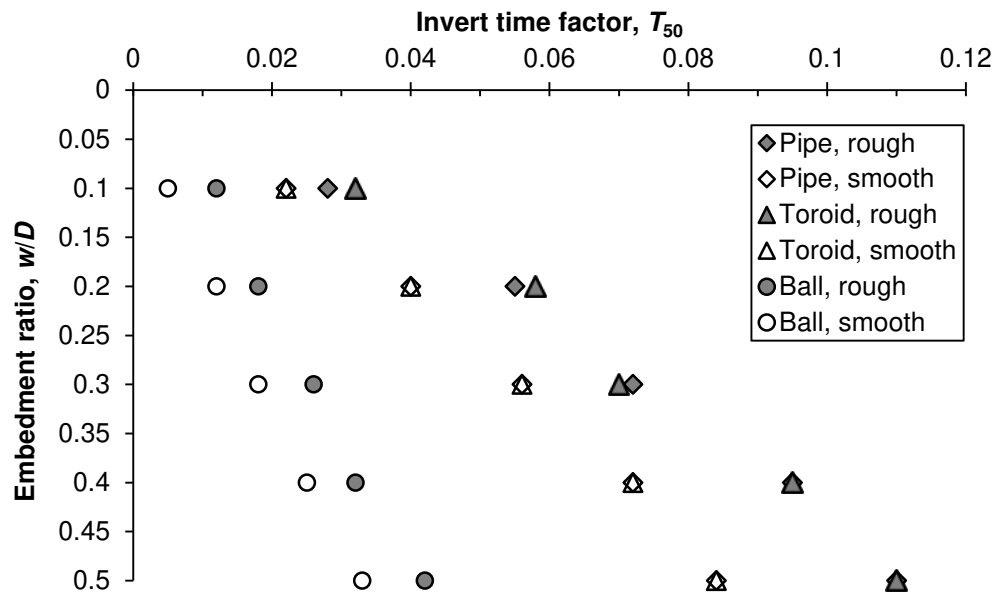
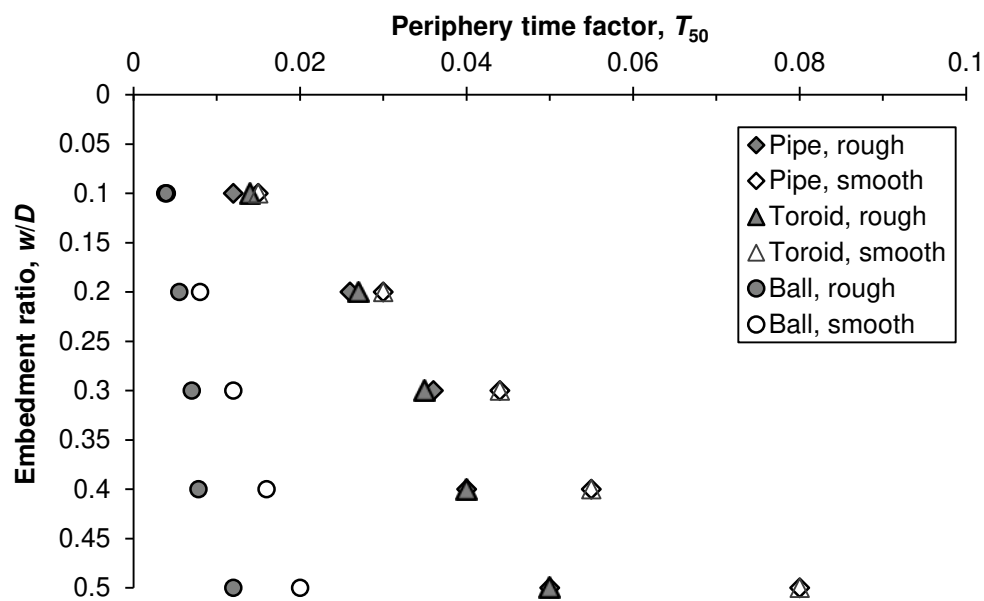


Figure 9 Average pore pressure dissipation and rise in effective stress around the object periphery

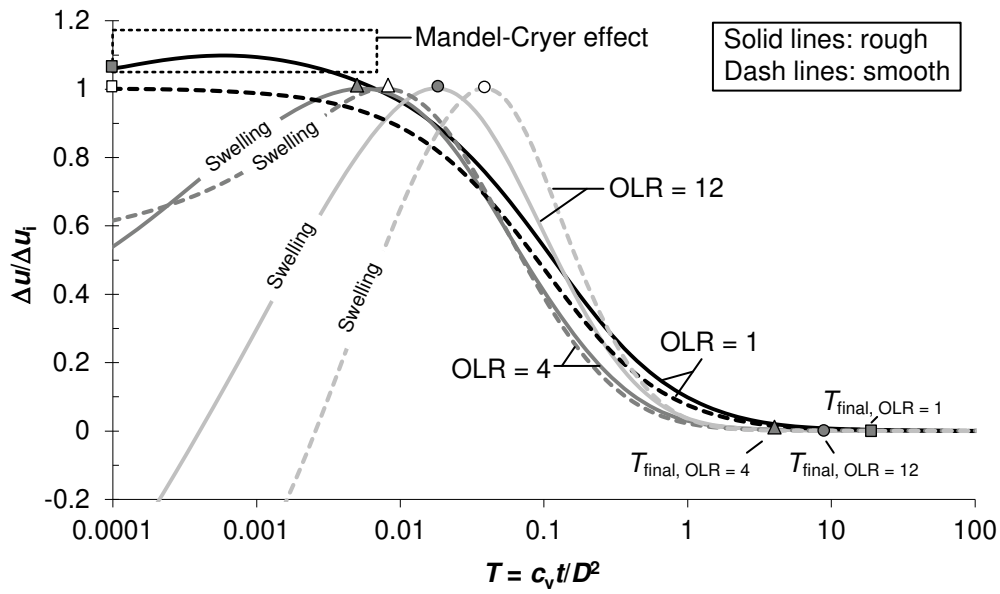


(a) Invert consolidation

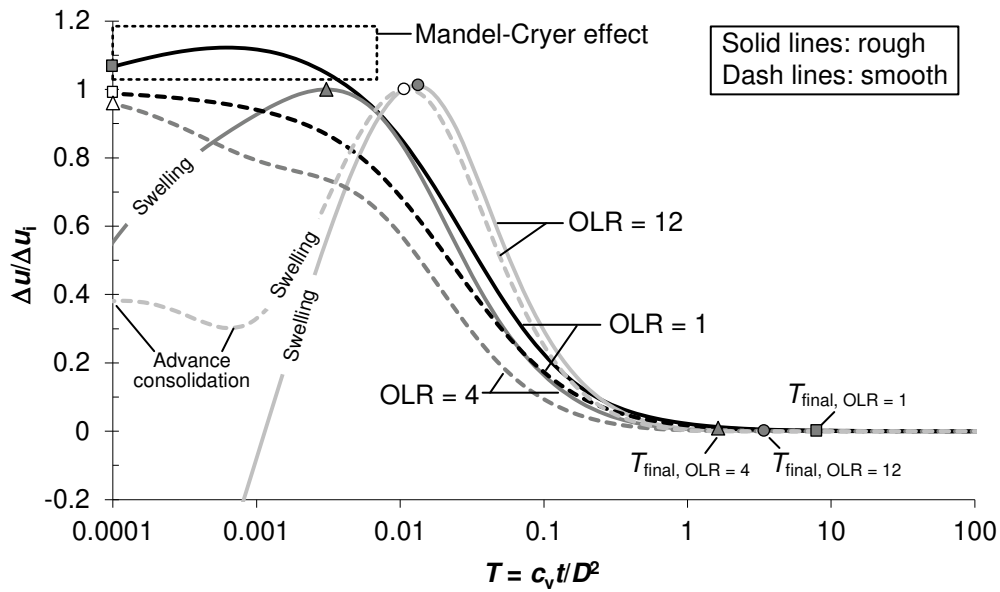


(b) Periphery consolidation

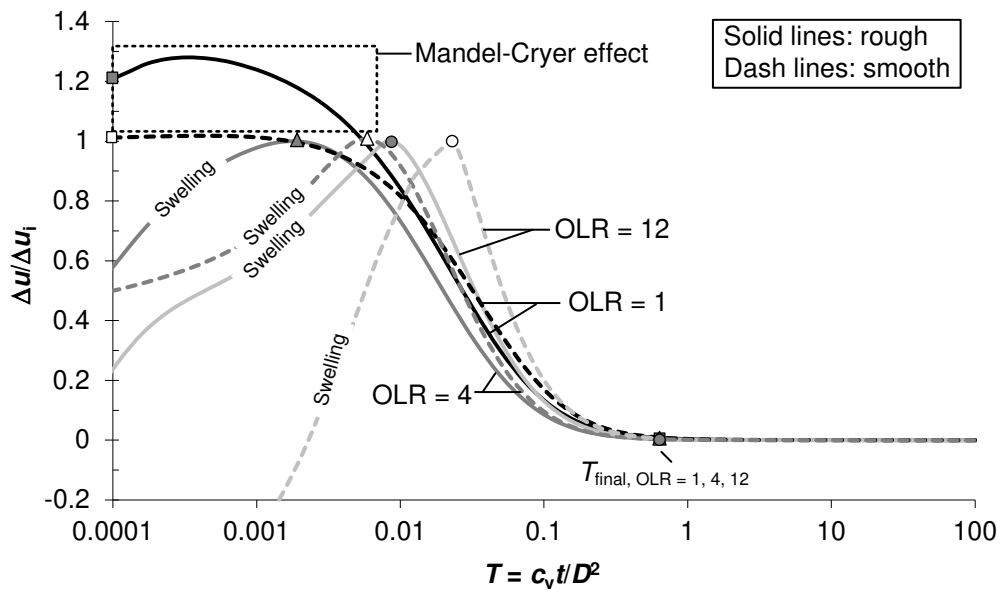
Figure 10 Summary of T_{50} for invert consolidation and average periphery consolidation



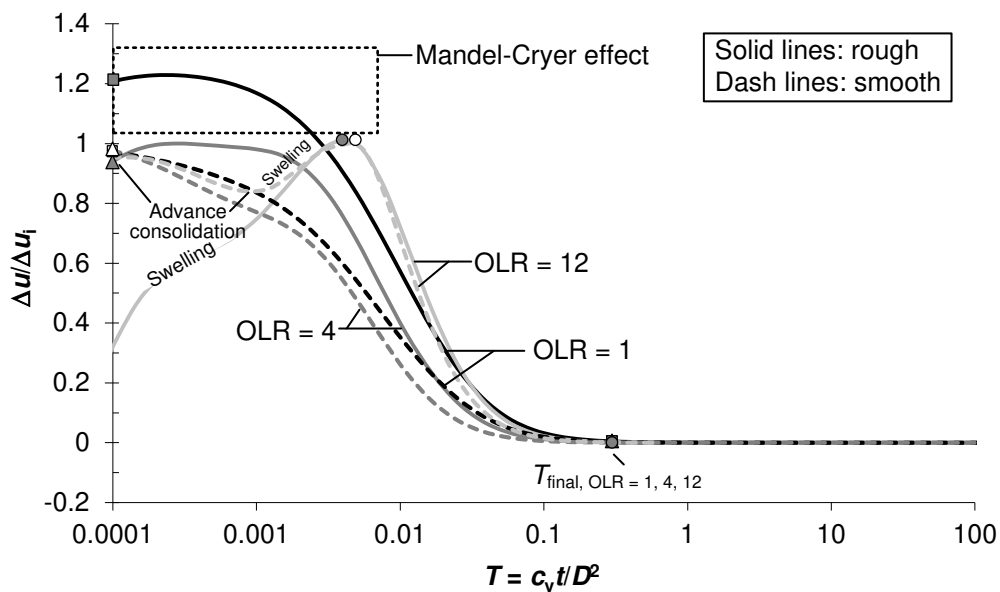
(a) Toroid with different OLR, $w/D = 0.5$



(b) Toroid with different OLR, $w/D = 0.1$



(c) Ball with different OLR, $w/D = 0.5$



(d) Ball with different OLR, $w/D = 0.1$

[Figure 11 Excess pore pressure dissipation time histories at invert for embedded toroid and pipe under varying OLR and \$w/D\$](#)

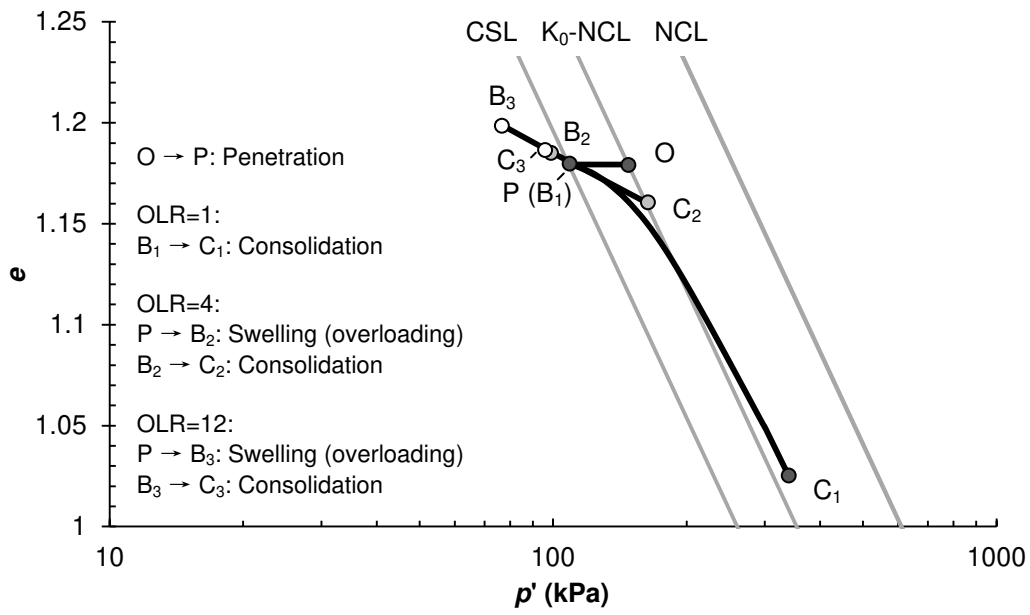


Figure 12 Stress path for toroid invert during penetration and consolidation

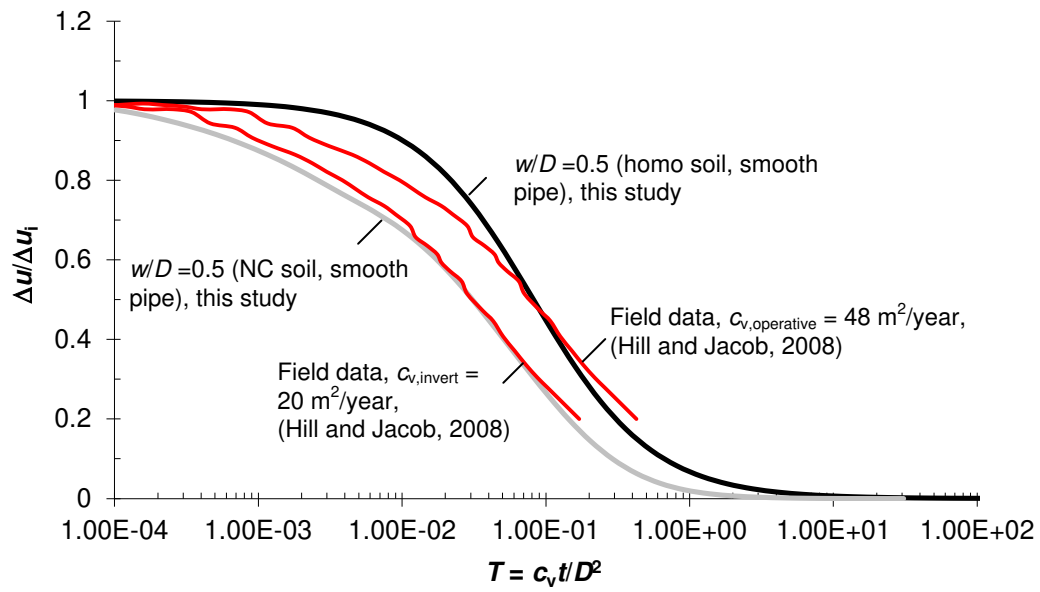
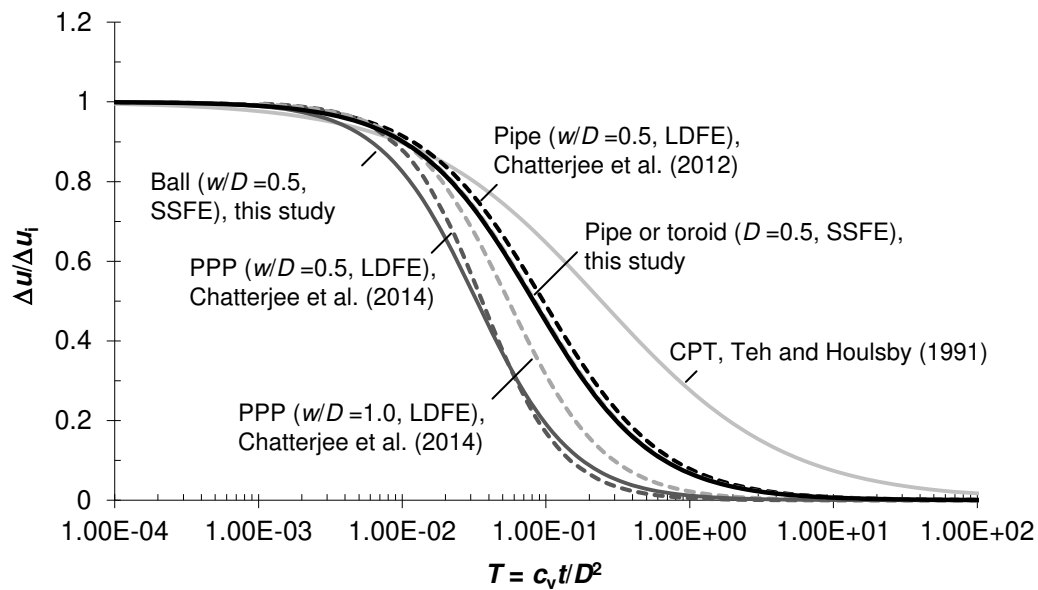


Figure 13 Comparison of calculated and observed dissipation curves at pipe invert



[Figure 14 Excess pore pressure dissipation time histories at invert for smooth pipe and penetrometers in this study compared with CPT, pipe and PPP \(Chatterjee et al., 2014\) on homogeneous soil](#)
~~Figure 14 Excess pore pressure dissipation time histories at invert for smooth pipe and penetrometers in this study compared with CPT, pipe and PPP (Chatterjee et al., 2014) on homogeneous soil~~