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## Electronic Properties of Ga<sub>2</sub>O<sub>3</sub> Polymorphs

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Ga<sub>2</sub>O<sub>3</sub> is an ultra-wide-band-gap semiconductor especially promising for power electronic applications. One advantage of this material is its ability to exist in different phases, which may add flexibility to device design, namely through polarization engineering of two-dimensional electron gases. Although much is known about monoclinic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, much less is known about many basic electronic properties of other phases. In this work, four of the most common phases of Ga<sub>2</sub>O<sub>3</sub> ( $\alpha$ ,  $\beta$ ,  $\kappa$ , and  $\epsilon$ ) are investigated with first-principles calculations based on hybrid density functional theory. The structural and electronic properties of each phase are compared, and band offsets between the phases and other common wide-band-gap semiconductors are determined. All four phases of Ga<sub>2</sub>O<sub>3</sub> are found to exhibit self-trapping holes, large Mg acceptor ionization energies, deep oxygen vacancy donor levels, and low-lying valence-band maxima. In addition, all phases have large valence-band offsets but small or modest conduction-band offsets with GaN, SiC, and Si. With AlN and diamond, all Ga<sub>2</sub>O<sub>3</sub> phases have large conduction-band and valence-band offsets.

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The wide bandgap, large breakdown voltage, high device figures of merit, and substrate availability make Ga<sub>2</sub>O<sub>3</sub> a tantalizing material for power electronics applications.<sup>1–4</sup> Steady progress has been made in Ga<sub>2</sub>O<sub>3</sub>-based devices, for instance the demonstration of vertical transistors with breakdown voltages over 1 kV,<sup>5</sup> transistors with power added efficiency of 12% at 1 GHz,<sup>6</sup> modulation-doped field effect transistors with sheet charge densities of  $1.2 \cdot 10^{13} \text{ cm}^{-2}$ ,<sup>7</sup> and enhancement-mode transistors with high breakdown voltages.<sup>4</sup>

The polymorphism exhibited by Ga<sub>2</sub>O<sub>3</sub> may spur further breakthroughs in gallium-oxide-based power electronics. Nearly all Ga<sub>2</sub>O<sub>3</sub>-containing devices utilize the monoclinic  $\beta$  phase, the most stable and best-characterized polymorph.<sup>8,9</sup> However, other phases of Ga<sub>2</sub>O<sub>3</sub> have also received attention due to potentially favorable growth characteristics, and to the possibility of polarization engineering made possible by the polar nature of their crystal structures.<sup>10–12</sup> In principle, this polarization could be utilized to produce two-dimensional electron gases (2DEGs) in analogy with GaN/AlN-based transistors.<sup>13</sup> Although  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is not usually discussed in the context of polarization, its structure is analogous to the ground-state corundum phase of Al<sub>2</sub>O<sub>3</sub>. Thus understanding  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is necessary for understanding the characteristics of Ga<sub>2</sub>O<sub>3</sub>-Al<sub>2</sub>O<sub>3</sub> alloys<sup>14</sup> that will be crucial for gallium oxide heterostructures.<sup>15</sup>

Theoretical studies have revealed some properties of different Ga<sub>2</sub>O<sub>3</sub> phases.<sup>9,13,14,16</sup> Early work examined structural properties, and reported that the  $\beta$  phase was most stable (although the  $\epsilon$  phase was found to be stable above 1600 K).<sup>9</sup> More recent work has established the polarization properties of the  $\epsilon$  and  $\kappa$  phases.<sup>13,16</sup> Although band offsets were recently reported for the  $\alpha$  and  $\beta$  phases,<sup>14</sup> they have apparently not been reported for the  $\epsilon$  or  $\kappa$  phases. Furthermore, although hole trapping and acceptor properties have been discussed in the theoretical literature for the  $\beta$  phase of Ga<sub>2</sub>O<sub>3</sub>,<sup>17–20</sup> much less is known about these properties in other phases.

In this work, the basic electrical properties of the  $\alpha$ ,  $\beta$ ,  $\kappa$ , and  $\epsilon$  phases of Ga<sub>2</sub>O<sub>3</sub> are calculated and compared. In addition to lattice parameters and band gaps, the stability of self-trapped holes and Mg acceptor ionization energies are also calculated. Band offsets are determined and compared with the wide-band-gap nitride semiconductors, diamond, Si, and with the 4H and 6H phases of SiC, all of which may be useful when incorporated into Ga<sub>2</sub>O<sub>3</sub>-containing devices.

The properties of Ga<sub>2</sub>O<sub>3</sub> phases are examined here using hybrid density functional theory, which has been shown to accurately

describe the band structures of wide-band-gap semiconductors as well as deep defect behavior within them.<sup>18,19,21–23</sup> These calculations are based on generalized Kohn-Sham theory<sup>24,25</sup> and use the screened hybrid functional of Heyd-Scuseria-Ernzerhof (HSE)<sup>26</sup> with the projector-augmented wave method,<sup>27</sup> as implemented in VASP.<sup>28</sup> The fraction of Hartree-Fock exchange for all phases is set to 0.32, yielding a bandgap (4.87 eV)<sup>14</sup> in close agreement with experiment (4.90 eV)<sup>29</sup> for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, the best characterized phase of gallium oxide. For consistency, these parameters are also applied to the other three polymorphs. Encouragingly, the calculated bandgap for  $\epsilon$ -Ga<sub>2</sub>O<sub>3</sub> of 4.48 eV falls within the 4.4–4.6 eV range of band gaps reported from experiments.<sup>30,31</sup> Also, the bandgap of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub> is consistent with previous calculations,<sup>14</sup> and is consistent with the larger 5.3 eV bandgap observed in optical studies.<sup>32</sup>

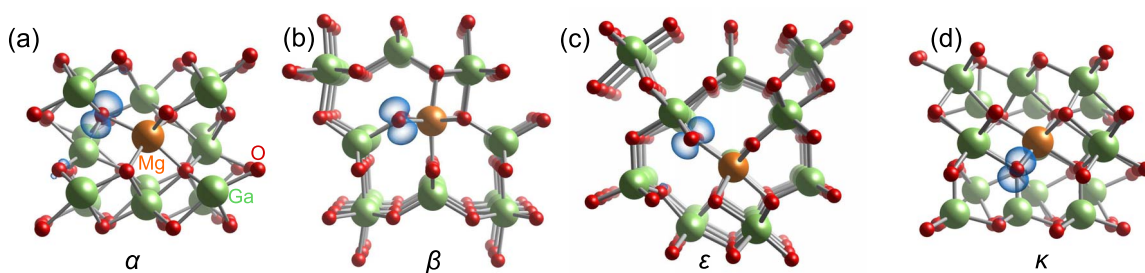
400 eV plane-wave cutoffs and two special k-points are used for all supercell calculations. For calculations of  $\alpha$ -Ga<sub>2</sub>O<sub>3</sub>,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub>, and  $\epsilon$ -Ga<sub>2</sub>O<sub>3</sub>, 120-atom supercells were employed, while for  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>, an 80-atom supercell was used. In each case, spin polarization is explicitly taken into account for the neutral charge states of impurities and for polarons. Ga 3d electrons were treated in the core, as their inclusion has a negligible effect on defect formation energies.<sup>33</sup> Calculations of defect thermodynamic transition levels are performed in accordance with the standard formalism,<sup>21</sup> and consistently with previous calculations in Ref. 19. The calculated lattice parameters for each phase of Ga<sub>2</sub>O<sub>3</sub> are shown in Table I. The lattice parameters are in good agreement with previous theoretical studies<sup>9,14,16,33</sup> and experimental observations.<sup>34</sup> Some controversy exists over the structural assignments for  $\epsilon$ - and  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub>.<sup>11,16,35</sup> In this study, the structure of  $\epsilon$ -Ga<sub>2</sub>O<sub>3</sub> is adopted from Ref. 34, and  $\kappa$ -Ga<sub>2</sub>O<sub>3</sub> is adopted from Ref. 16.

Different coordinations exist for both Ga and O atoms in Ga<sub>2</sub>O<sub>3</sub> polymorphs. For the sake of comparison, the same coordination is

**Table I. Calculated lattice parameters, direct band gaps, hole self-trapping energies ( $E_{\text{ST}}$ ), and Mg<sub>Ga</sub> acceptor ionization energies for the four phases of Ga<sub>2</sub>O<sub>3</sub> considered.**

phase	<i>a</i> (Å)	<i>b</i> (Å)	<i>c</i> (Å)	bandgap (eV)	$E_{\text{ST}}$ (eV)	Mg <sub>Ga</sub> (0/–) (eV)
$\alpha$	4.94	–	13.30	5.57	0.51	1.11
$\beta$	12.16	3.03	5.79	4.87	0.53	1.36
$\epsilon$	2.99	14.38	9.20	4.48	0.79	1.45
$\kappa$	5.04	8.61	9.19	4.84	0.87	1.46

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**Figure 1.** Spin densities of  $\text{Mg}_{\text{Ga}}^0$  acceptors in the four different  $\text{Ga}_2\text{O}_3$  polymorphs studied here: (a)  $\alpha$ , (b)  $\beta$ , (c)  $\epsilon$ , and (d)  $\kappa$ . Isosurfaces are set to 5% of the maximum in each case.

utilized for each defect center in each phase where possible.  $\text{Mg}_{\text{Ga}}$  is chosen to substitute for 6-fold coordinated Ga in each phase, as this configuration is the lowest in energy for  $\beta$ - $\text{Ga}_2\text{O}_3$ .<sup>19</sup> For the oxygen vacancy ( $V_{\text{O}}$ ), the 4-fold coordinated O sites are chosen, since it is the lowest-energy configuration for  $V_{\text{O}}^{2+}$  in  $\beta$ - $\text{Ga}_2\text{O}_3$ .<sup>33</sup> Hole self-trapping energies ( $E_{\text{ST}}$ ; as defined in Ref. 17) are reported for whichever site exhibits the largest  $E_{\text{ST}}$ . Band alignments between the different phases of  $\text{Ga}_2\text{O}_3$  and other wide-band semiconductors were calculated using surface-slab calculations with nonpolar orientations. This approach allows the electrostatic potential within each material to be referenced to vacuum, and provides natural band alignments that avoid strain or lattice mismatch effects.<sup>36,37</sup> For each surface calculation, supercells using at least 15 Å of material were used, together with at least 15 Å of vacuum to avoid spurious interactions. All atoms within 5 Å of the surface were allowed to relax.

Previous calculations indicated that holes spontaneously localize in the  $\beta$  phase of  $\text{Ga}_2\text{O}_3$ ;<sup>17</sup> in this work holes are found to self trap in all phases of  $\text{Ga}_2\text{O}_3$ . The  $E_{\text{ST}}$  of these hole polarons vary between 0.5–0.9 eV, and are listed in Table 1. This indicates that in every phase of  $\text{Ga}_2\text{O}_3$  holes prefer to spontaneously localize onto single O sites, as opposed to being delocalized free holes, even in the absence of any defect or impurity species. In all but one phase holes prefer to trap onto three-fold-coordinated oxygen sites. The exception is the  $\alpha$  phase, which does not contain such sites, thus holes can only localize onto a four-fold O site.  $\alpha$ - $\text{Ga}_2\text{O}_3$  has the lowest  $E_{\text{ST}}$  of all phases studied here, while  $\kappa$ - $\text{Ga}_2\text{O}_3$  has the highest. Should the Fermi level approach the  $E_{\text{ST}}$  level within  $\text{Ga}_2\text{O}_3$ , the polaron will act as a hole trap as has been previously predicted for AlN.<sup>38</sup>

Cation-site acceptors such as  $\text{Mg}_{\text{Ga}}$  have been shown to act as “polaronic” acceptors in  $\text{Ga}_2\text{O}_3$ ,<sup>18,19</sup> meaning that they trap holes onto nearest-neighbor O sites and act as deep acceptors. Because of the similarities in hole-trapping behavior found for all phases of  $\text{Ga}_2\text{O}_3$ , it is expected that the behavior of cation-site acceptors should be consistent in these compounds. This is confirmed by examining the ionization energies [equivalent to the acceptor’s (0/−) thermodynamic transition level] of the  $\text{Mg}_{\text{Ga}}$  acceptor in each  $\text{Ga}_2\text{O}_3$  phase, which are shown in Table 1. These energies range from 1.1–1.5 eV, indicating that Mg dopants will be deep in every phase of  $\text{Ga}_2\text{O}_3$ . Furthermore, these energies exhibit the same pattern as the self-trapped holes, with  $\alpha$ - $\text{Ga}_2\text{O}_3$  having the smallest  $\text{Mg}_{\text{Ga}}$  ionization energy and  $\kappa$ - $\text{Ga}_2\text{O}_3$  having the highest.

In Fig. 1, isosurfaces of the spin densities for  $\text{Mg}_{\text{Ga}}^0$  are shown for each phase of gallium oxide. In all cases, a highly localized hole is trapped at an O site adjacent to  $\text{Mg}_{\text{Ga}}$ . As was the case for hole polarons,  $\text{Mg}_{\text{Ga}}$  traps holes onto a three-fold-coordinated O site for all phases except for  $\alpha$ - $\text{Ga}_2\text{O}_3$ , where  $\text{Mg}_{\text{Ga}}$  traps a hole onto an adjacent four-fold site. Isosurface plots of the self-trapped holes (not shown) in each phase look nearly identical to those shown in Fig. 1, albeit with the absence of the  $\text{Mg}_{\text{Ga}}$  impurity.

The behavior of  $V_{\text{O}}$ , known to be a negative- $U$  donor in  $\beta$ - $\text{Ga}_2\text{O}_3$ ,<sup>33</sup> serves as another means for comparing the properties of gallium oxide phases. The negative- $U$  nature of  $V_{\text{O}}$  derives from the fact that the 1+ charge state is not stable, and only  $V_{\text{O}}^{2+}$  and  $V_{\text{O}}^0$  will be present in

thermal equilibrium. This same behavior is also observed for  $V_{\text{O}}$  in each phase of  $\text{Ga}_2\text{O}_3$ . As shown in Fig. 2,  $V_{\text{O}}$  is also a deep donor in each phase, and the (2+/0) transition level lies between 1.5–2.1 eV below the CBM of all phases of  $\text{Ga}_2\text{O}_3$ .

The results for the thermodynamic transition levels self-trapped holes,  $\text{Mg}_{\text{Ga}}$  acceptors and  $V_{\text{O}}$  donors are incorporated into Fig. 2, which also shows band offsets between  $\text{Ga}_2\text{O}_3$  phases as well as GaN, AlN, SiC, Si, and diamond. Note that due to the large Stokes shifts predicted for  $\text{Ga}_2\text{O}_3$  defects, optical transition levels will vary considerably from the thermodynamic transition levels.<sup>17,39</sup> Note also that all band edge positions are referenced to the valence band maximum (VBM) of GaN, which is set to 0 eV in Fig. 2. As would be expected due to the shared O 2p character of the valence band and of the polaronic states, the positions of  $\text{Mg}_{\text{Ga}}$  acceptor levels (and self-trapped holes) roughly follow the position of the gallium oxide VBM. The positions of the  $V_{\text{O}}$  donor levels roughly track each oxide CBM, as is expected since their defect states derive from the Ga s orbitals that comprise the conduction band. The overall variation in the absolute positions of the CBM and VBM of the gallium oxide phases is 0.9 eV or less, which is also roughly the level of variation in the band gaps, as shown in Table 1. In comparison with the other semiconductors shown in Fig. 2, the  $\text{Ga}_2\text{O}_3$  phases have much lower-lying VBM.

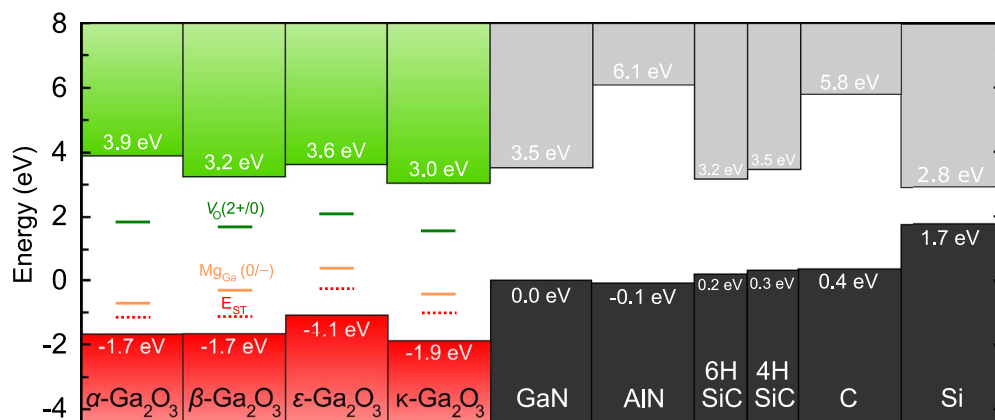
All  $\text{Ga}_2\text{O}_3$  phases have sizeable valence-band offsets (VBO) with GaN. The smallest VBO (1.1 eV) occurs for GaN/ $\epsilon$ - $\text{Ga}_2\text{O}_3$ ; all others are 1.7 eV or larger. The conduction band offsets (CBO) between GaN and  $\text{Ga}_2\text{O}_3$  phases are much smaller in comparison; none are larger than 0.5 eV. The calculated offsets for  $\beta$ - $\text{Ga}_2\text{O}_3$ /GaN (VBO = 1.7 eV, CBO = 0.3 eV) are in good agreement with an X-ray photoelectron spectroscopy (XPS) study<sup>40</sup> that reported a VBO of 1.4 eV and a CBO of 0.1 eV for the  $\beta$ - $\text{Ga}_2\text{O}_3$ /GaN heterostructure.

In contrast to GaN, the CBOs between  $\text{Ga}_2\text{O}_3$  and AlN are much larger, and exceed 2 eV for all phases. The calculated VBOs are all larger than 1 eV. This result conflicts with some recent XPS studies that report VBOs of −0.1 eV for plasma-assisted atomic-layer deposited (ALD) AlN and 0.7 eV for thermal ALD AlN in one study<sup>12</sup> and 0.6 eV in another<sup>41</sup> (giving resulting CBOs between 0.6 and 1.8 eV.<sup>12,41</sup>). However, it has been noted<sup>12</sup> that oxygen incorporation into AlN films may explain the large discrepancies between these values, which might also explain why they strongly deviate from the calculated offsets.

The calculated band offsets between diamond and the  $\text{Ga}_2\text{O}_3$  phases are similar to those of AlN, featuring large CBO and VBO. Both phases of SiC have small CBO with all  $\text{Ga}_2\text{O}_3$  phases, but have VBO that are greater than 1 eV in each case.

In the case of Si, all phases have substantial VBO that are all in excess of 2.5 eV. The calculated  $\beta$ - $\text{Ga}_2\text{O}_3$ /Si VBO of 3.4 eV is in good agreement with two XPS studies that report a 3.5 eV offset.<sup>42,43</sup> However, for CBOs with Si there is more variation between the  $\text{Ga}_2\text{O}_3$  phases. The CBOs are small for  $\kappa$ - $\text{Ga}_2\text{O}_3$ /Si (0.2 eV) and  $\beta$ - $\text{Ga}_2\text{O}_3$ /Si (0.4 eV). However, those for  $\epsilon$ - $\text{Ga}_2\text{O}_3$ /Si (0.8 eV) and  $\alpha$ - $\text{Ga}_2\text{O}_3$ /Si (1.1 eV) are more substantial.

In summary, hybrid functional calculations indicate that the  $\alpha$ ,  $\beta$ ,  $\kappa$ , and  $\epsilon$  phases of  $\text{Ga}_2\text{O}_3$  exhibit broadly similar electronic properties. In all phases, holes self trap as polarons, meaning that free holes will not



**Figure 2.** Band offsets between the  $\alpha$ ,  $\beta$ ,  $\epsilon$ , and  $\kappa$  phases of  $\text{Ga}_2\text{O}_3$  together with those of GaN, AlN, 4H-SiC, 6H-SiC, Si, and diamond. The dotted red lines in the oxide band gaps indicate the hole self-trapping energy relative to the VBM of the  $\text{Ga}_2\text{O}_3$  phases. The orange lines indicate the ionization energies of  $\text{Mg}_{\text{Ga}}$  acceptors in each phase, while the green lines indicate the  $\text{V}_{\text{O}}$  (2+/0) donor levels. Values in white indicate the energetic position of each band maximum or minimum.

be stable. This hole trapping also affects the properties of acceptors in these polymorphs, and  $\text{Mg}_{\text{Ga}}$  acceptors have large ionization energies in each phase. A comparison of band offsets indicates that all phases have low-lying VBM, 1 eV or more below that of GaN, AlN, SiC, Si and diamond. Furthermore, while all phases have CBM in the vicinity of those of GaN, SiC, and Si, they have substantial conduction band offsets with AlN and diamond.

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